

Axion Searches in light of string theory

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arxiv:2408.06652; 2410.12634

String Theory

- Still the most plausible theory of quantum gravity.
- Key Properties to phenomenology:
 - a. Extra Dimensions**

Resulting KK modes of particles .
 - b. Landscape**

Physical constants could be variable.

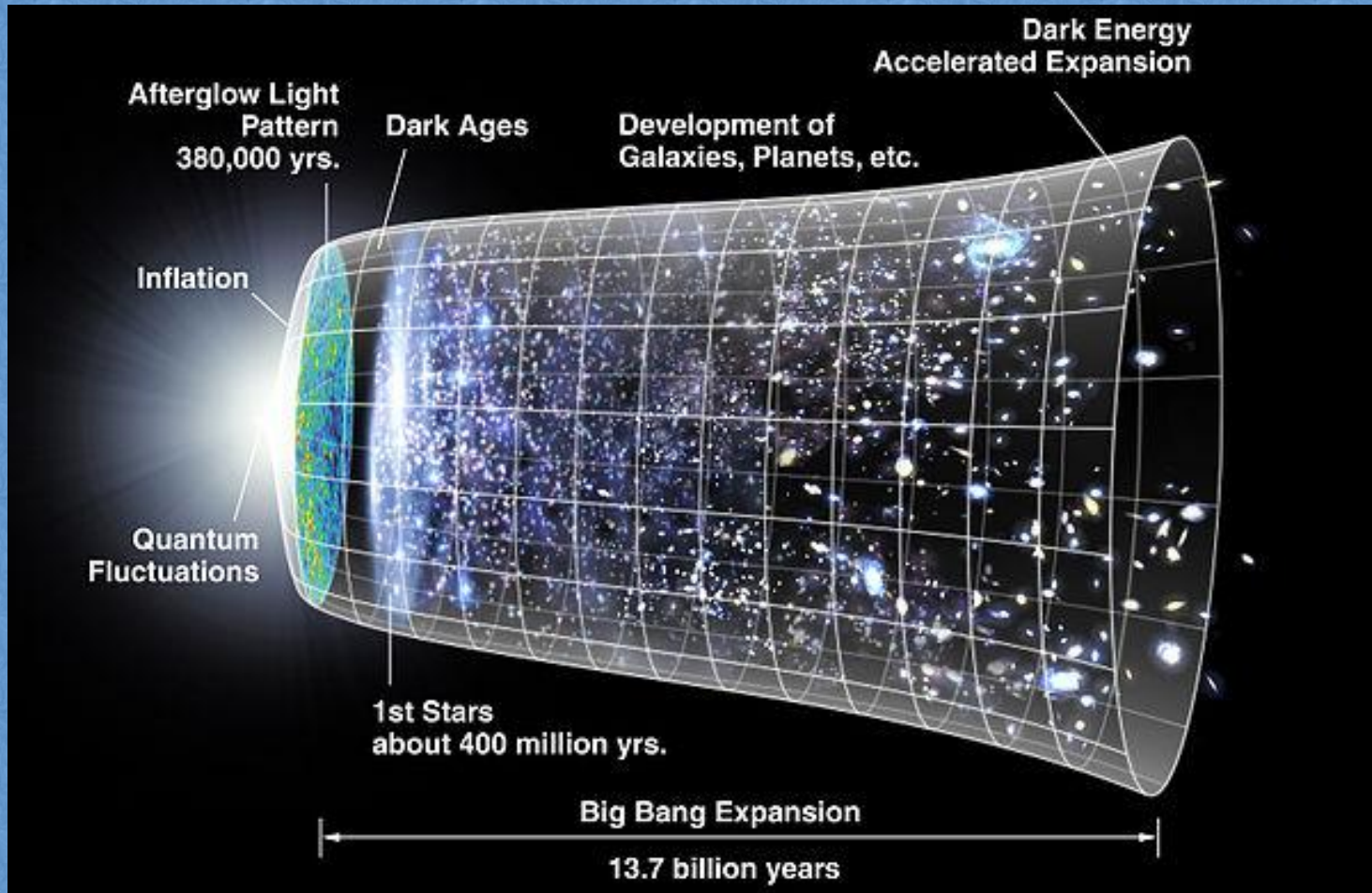
String Theory

- Key Properties to phenomenology:
 - a. Many axion-like particles**

Naively $\binom{6}{2} = 15$ actually many more:
Axiverse .
 - b. could solve the strong CP problem**

One of the alps potentiates is
dominated by QCD instanton.

History of the Universe



The General compactification

Low energy effective lagrangin

$$\mathcal{L} \supset \frac{f_a^2}{2} (\partial a)^2 - \Lambda^4 U(a),$$

$$\Lambda^4 = \mu^4 e^{-S} \quad f_a \sim \frac{M_{Pl}}{S} \sim 10^{16} \text{ GeV}$$

QCD Axion mass-symmetry breaking relations

$$m_0 \approx 6 \times 10^{-5} \text{eV} \left(\frac{10^{11} \text{GeV}}{f_a} \right)$$

General string theory: $f_a \sim 10^{16} \text{GeV}$

Dark Dimension: $f_a \sim 10^{10} \text{GeV}$

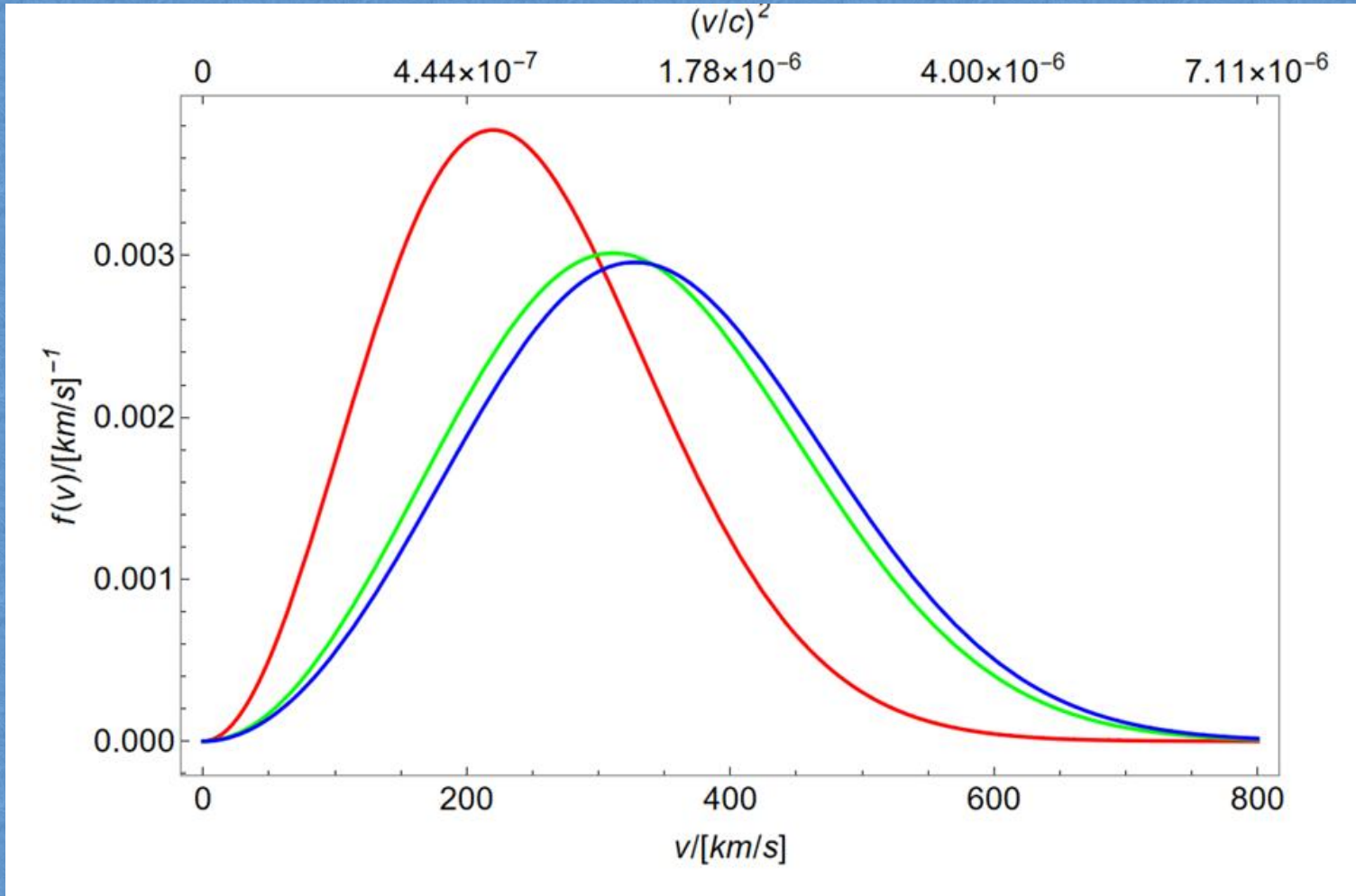
The Quantum picture of wave like dark matter

What we do know with certainties

density distribution

$$\rho_{DM} \approx 0.43 \text{GeV}/\text{cm}^3$$

The uncertainties: velocity distribution



$$f_1(v)dv = \frac{4v^2}{\sqrt{\pi}v_c^3} e^{-v^2/v_c^2} dv$$

The Quantum picture of wave like dark matter

wave function for a single particle:

$$D_i(v, t) = \frac{\sqrt{2\rho_{DM}/N}}{m} \cos \left[m \left(1 + \frac{v_i^2}{2} \right) t + \phi_i \right]$$

wave function of the total system:

$$D(t) = \frac{\sqrt{\rho_{DM}}}{m} \sum_j \alpha_j \sqrt{f(v_j)\Delta v} \\ \times \cos \left[m \left(1 + \frac{v_j^2}{2} \right) t + \phi_j \right]$$

The Quantum transitions in the Cavity

$$\begin{aligned} H_I &= - \int d^3x \mathcal{L}_{a\gamma\gamma} = -g_{a\gamma\gamma} B_0 a(t) \int d^3x \vec{\mathbf{E}} \cdot \hat{z} \\ &= -\frac{ig_{a\gamma\gamma} B_0 \sqrt{\rho_{DM}}}{\sqrt{2}m_a} \cdot \sum_{j,k} \alpha_j \sqrt{f(v_j) \Delta v} \\ &\times \cos(\omega_j t + \phi_j) \sqrt{\omega_k} (a_k e^{-i\omega_k t} U_k - c.c.) . \end{aligned}$$

The Quantum transitions in the Cavity

$$\begin{aligned}
 P &= \frac{\pi g_{a\gamma\gamma}^2 B_0^2 \rho_{DM} t V}{4m_a^2} \sum_{jk} C_k \delta(\omega_k - \omega_j) \omega_k \alpha_j^2 f(v_j) \Delta v \\
 &\approx \frac{\pi g_{a\gamma\gamma}^2 B_0^2 \rho_{DM} t V \langle \alpha^2 \rangle}{4m_a^2} \int dv f(v) \int d\omega_k C_k \frac{\omega_k}{d\omega_k} \delta(\omega_k - \omega_v) \\
 &\approx \frac{\pi g_{a\gamma\gamma}^2 B_0^2 t V}{2m_a^2} \rho_{DM} \int_0^\Delta dv_a f(v_a) C_{\omega_a} Q_c ,
 \end{aligned}$$

The axion spectrum density I_a is high

$$I_a = \frac{\rho_{CDM}}{(1/2)m_a \delta v^2}$$

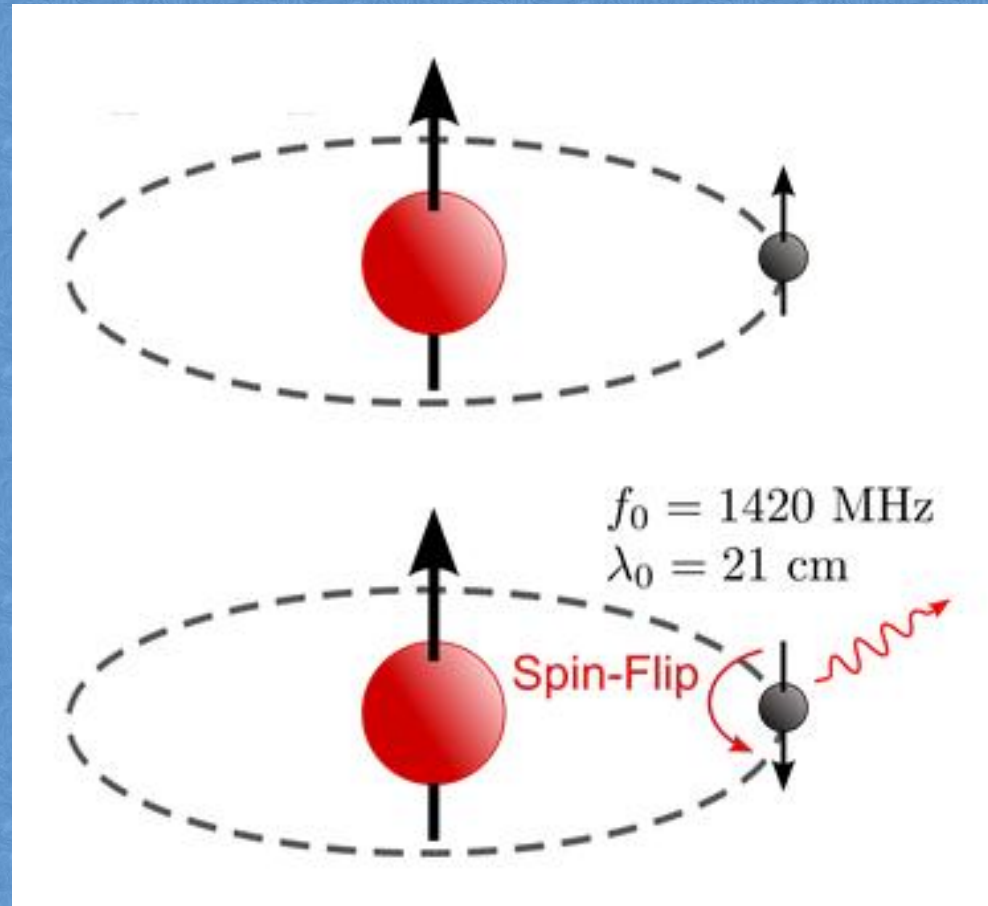
The interaction term in non-relativistic limit is:

$$H_{int} = \frac{1}{f_a} \sum g_f (\partial_t a \frac{\vec{p}_f \cdot \vec{\sigma}_f}{m_f} + \vec{\sigma}_f \cdot \vec{\nabla} a)$$

The transition rate is

$$R = \frac{\pi}{f_a^2} \left| \sum g_f \langle f | (\vec{v} \cdot \vec{\sigma}_f) | i \rangle \right|^2 I_a$$

Hydrogen 1S state transitions



Hydrogen atoms are ideal targets for the quantum window

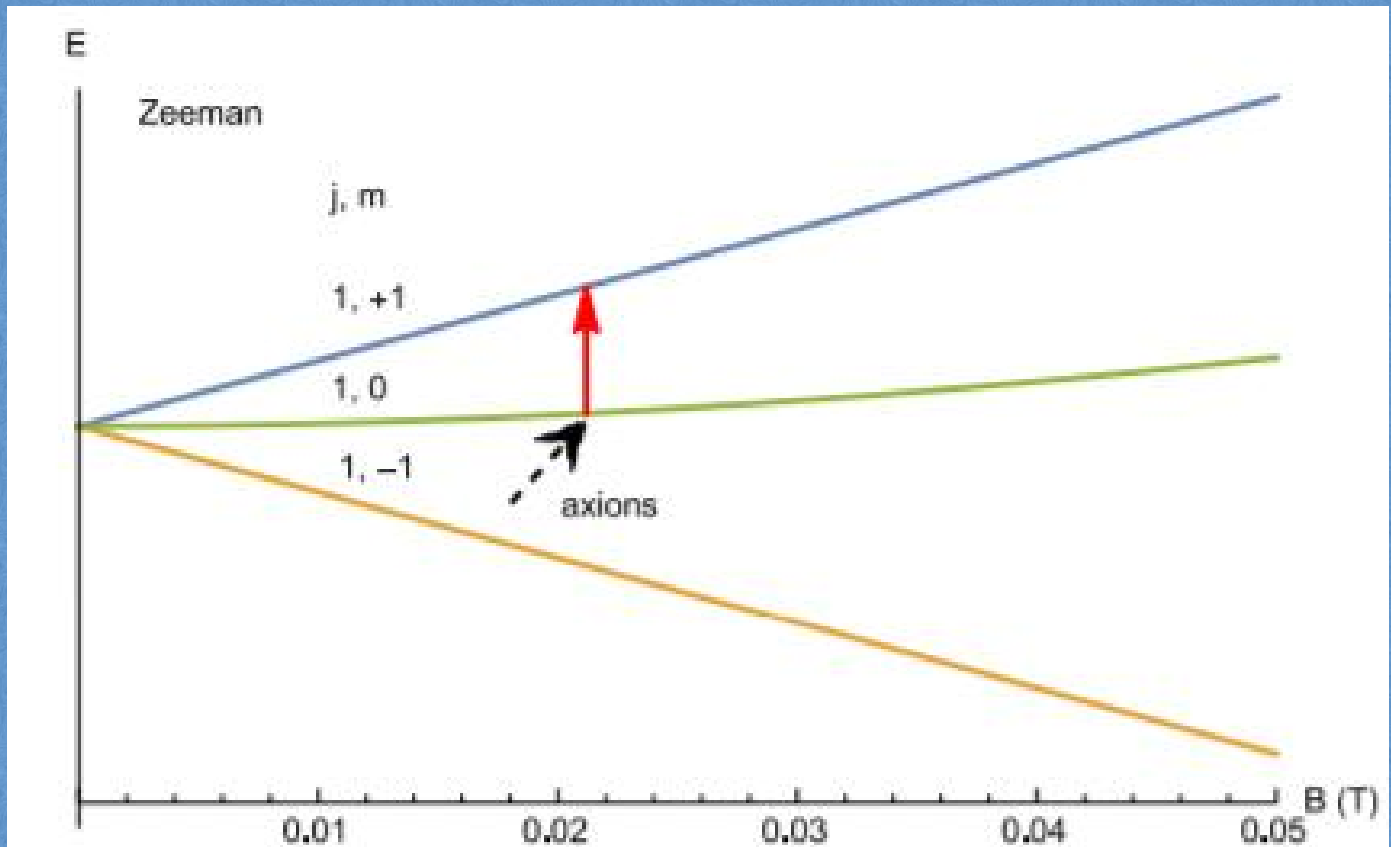


FIG. 3: The splitting of the hydrogen 1S triplet state. For the anthropic window $|1, 0\rangle \rightarrow |1, 1\rangle$ transition is suitable for the axion detection.

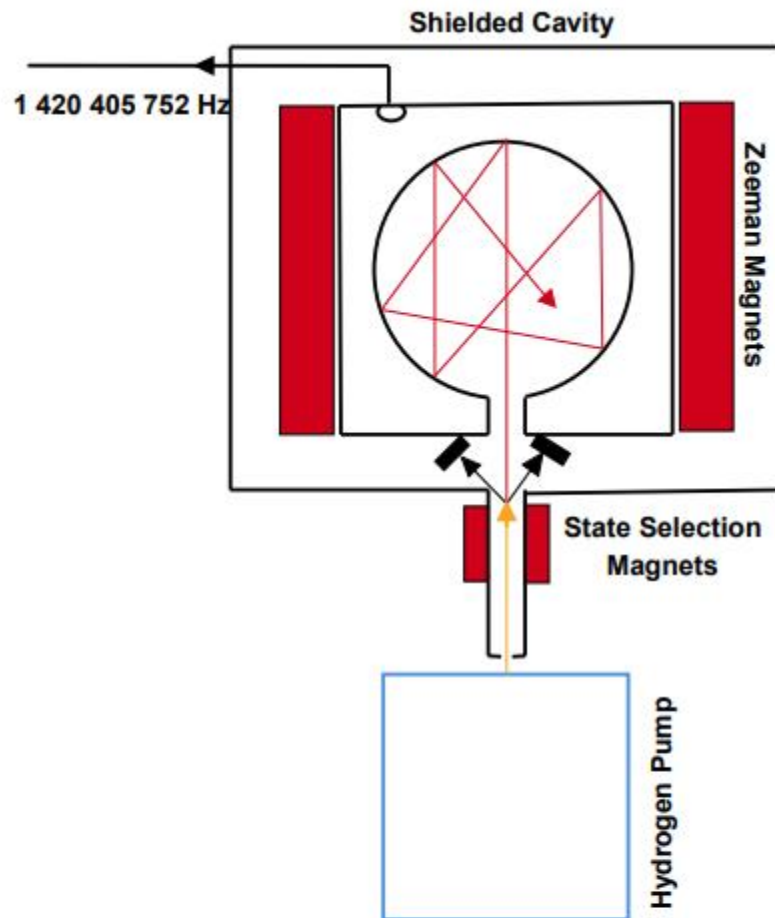


FIG. 5: An illustration of a typical hydrogen maser. With some modifications such as adding Zeeman magnets it may realize the scheme shown in Fig. 4.

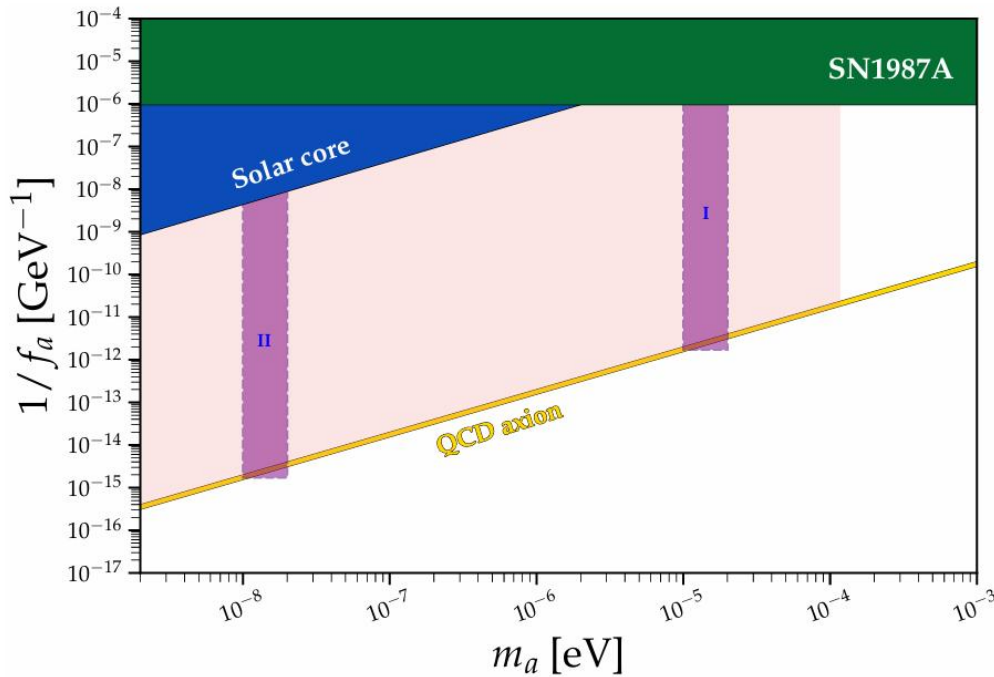
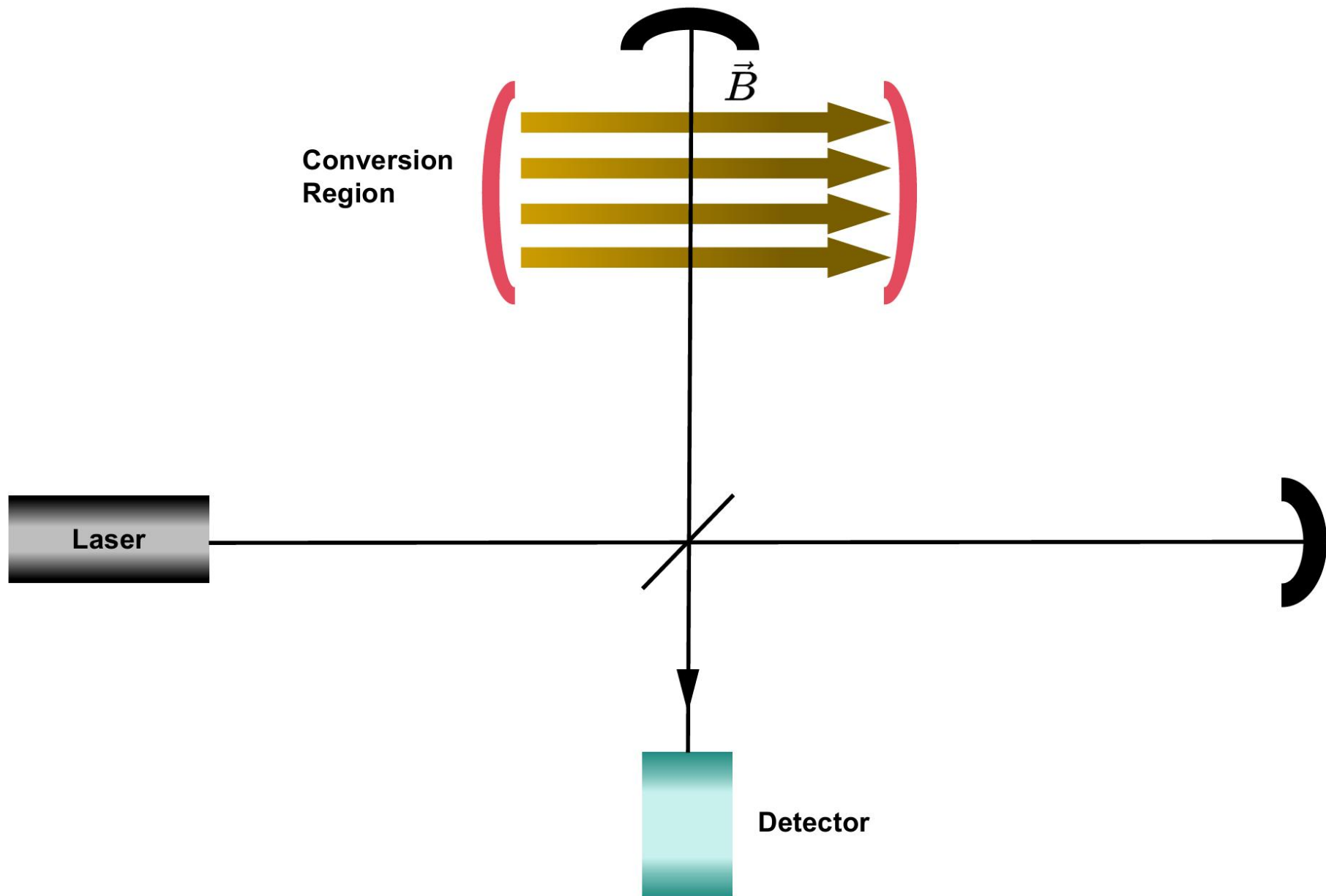


FIG. 3: The mass range that can be covered is determined by the available Zeeman field strength. To cover Region I ($m_a \in [10^{-5} \text{ eV}, 2 \times 10^{-5} \text{ eV}]$), a field strength of 0.15 Tesla is needed. Assuming 10% of this region is covered in a one-year data-taking period using 1 mole of H atoms, the sensitivity is $f_a = 5.8 \times 10^{11} \text{ GeV}$. For Region II ($m_a \in [10^{-8} \text{ eV}, 2 \times 10^{-8} \text{ eV}]$), a field strength of 0.002 Tesla is required. Assuming 10% of this region is covered using 10^3 moles of H over a one-year data-taking period, the sensitivity is $f_a = 5.8 \times 10^{14} \text{ GeV}$. The yellow band defines the QCD axion parameter space. The solid-colored regions, SN1987A and Solar Core, have been excluded by experiments and observations. The light coral region can be probed using this hydrogen splitting method if a 1 Tesla Zeeman field and an adequate amount of hydrogen atoms are deployed.

Dark Dimension

- Key Properties to phenomenology:
 - a. 1 Large Extra Dimension**
a lot of KK modes of dark matter.
 - b. Standard Model particles are confined on 4D.**
QCD axions can only account 1% dark matter.



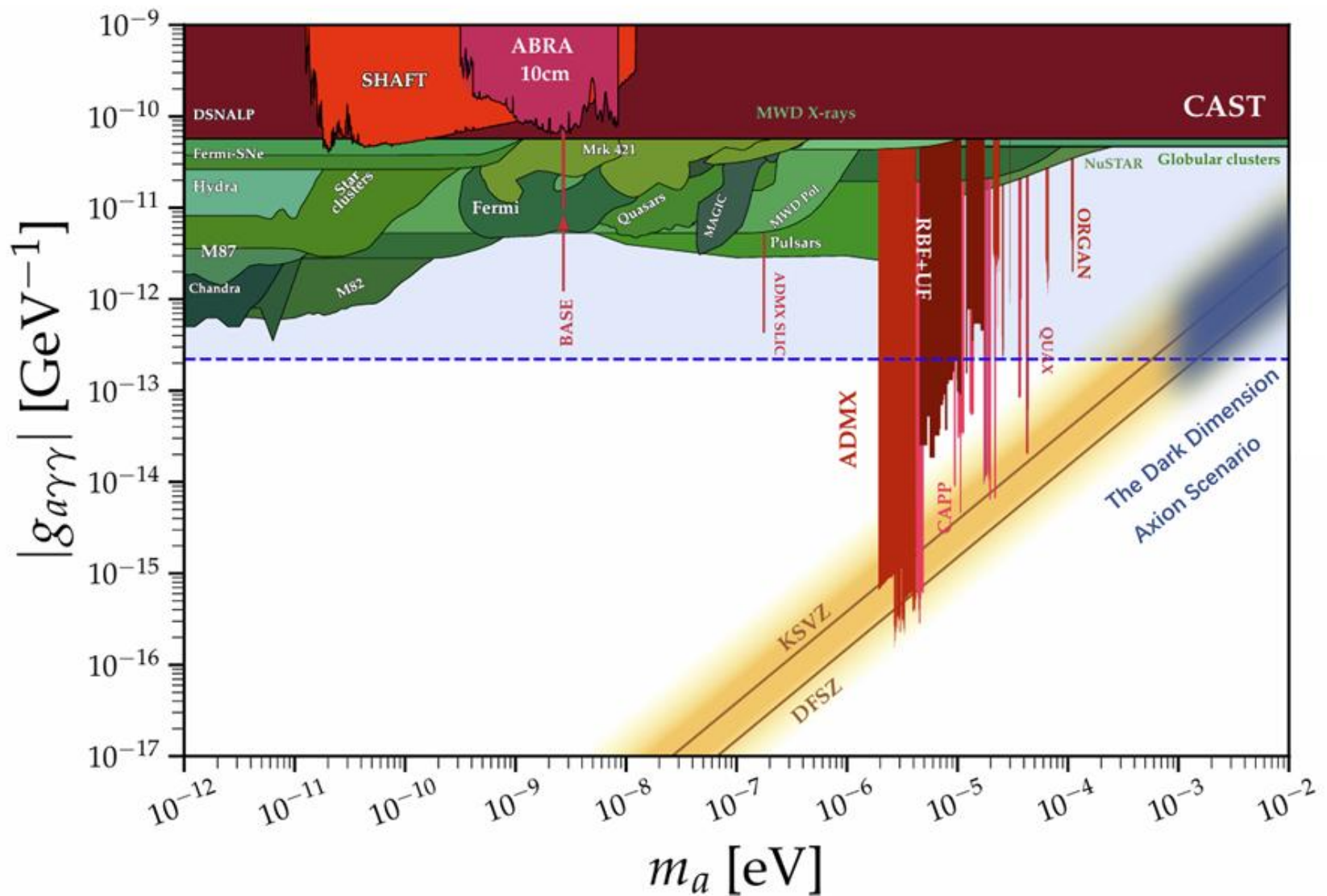
Axion in magnetic field

The conversion rate :

$$\eta_{\gamma \rightarrow a} \approx \frac{1}{4} (g_{a\gamma\gamma} BL)^2$$

Amplitude modulation:

$$\vec{E}_{in} = \vec{E}_0 (1 + \beta \sin \omega_m t) e^{i\omega t}$$



Assuming $n \sim 10^6$, $B \sim 40$ T, $L \sim 30$ m with a 10 W ($\lambda = 1 \mu\text{m}$) laser, and squeezed light achieving approximately O(10) dB shot noise reduction, after 10 days of operation.