

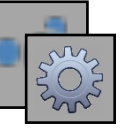
B&D介子三体衰变过程中的共振态粒子的虚贡献

王文飞 (山西大学)

第三届高能物理理论与实验融合发展研讨会
辽宁师范大学 2024-11-3



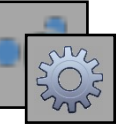
Outline



- **3-body B and D meson decays and the virtual contribution**
- **Some examples of the resonance virtual contributions**
- **Conclusions**

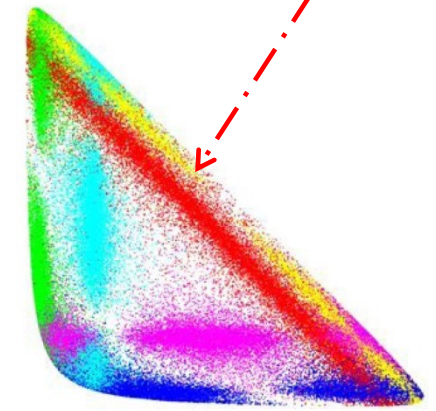
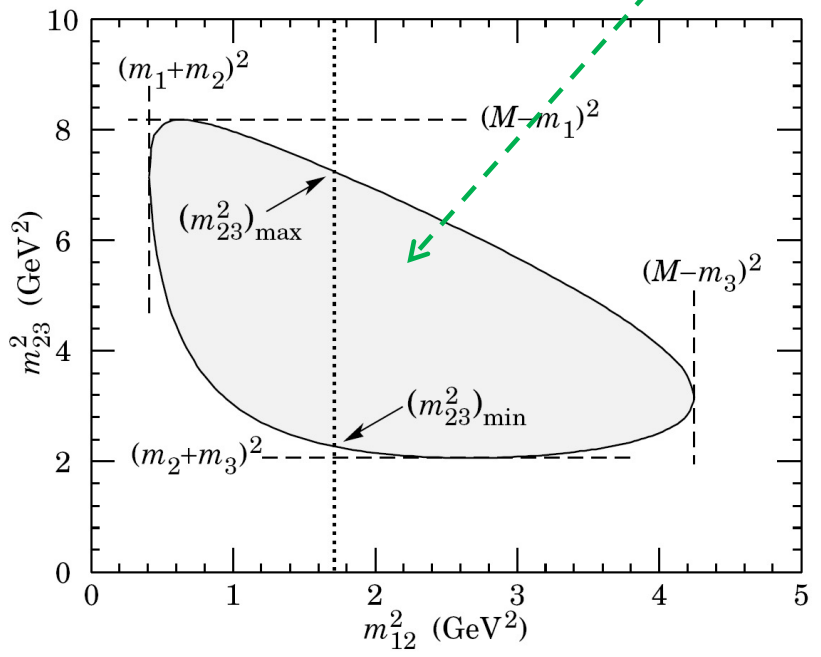
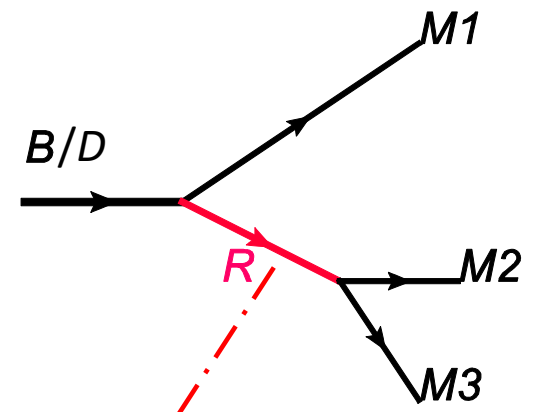
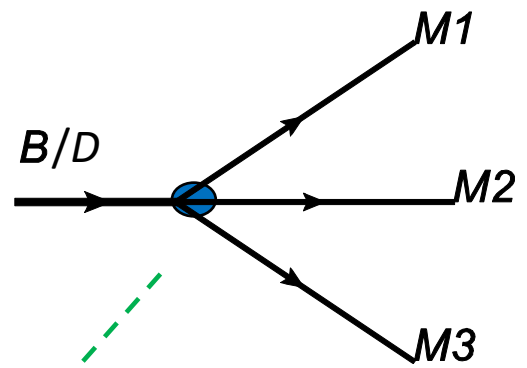


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- Some examples of the resonance virtual contributions
- Conclusions

B&D \rightarrow M1 M2 M3

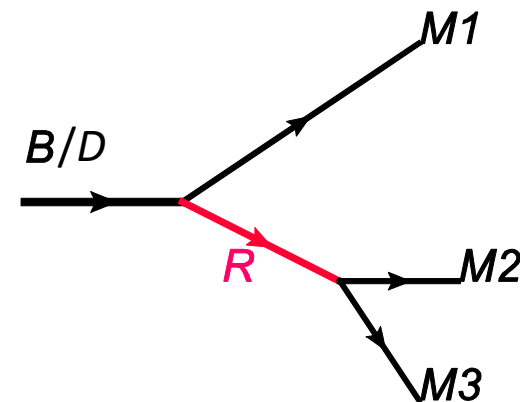
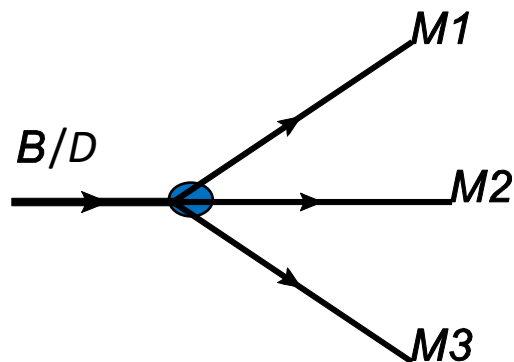


Dalitz Plot



3-body B and D meson decays and the virtual contribution

B&D → M1 M2 M3



The weak effective Hamiltonian:

1.

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \left\{ V_{ub} V_{uq}^* [C_1(\mu) Q_1^u(\mu) + C_2(\mu) Q_2^u(\mu)] - V_{tb} V_{tq}^* \left[\sum_{i=3}^{10} C_i(\mu) Q_i(\mu) \right] \right\} + \text{H.c.}, \quad (2)$$

where $q = d, s$. The functions Q_i ($i = 1, \dots, 10$) are the local four-quark operators:

The total amplitude within isobar approach:

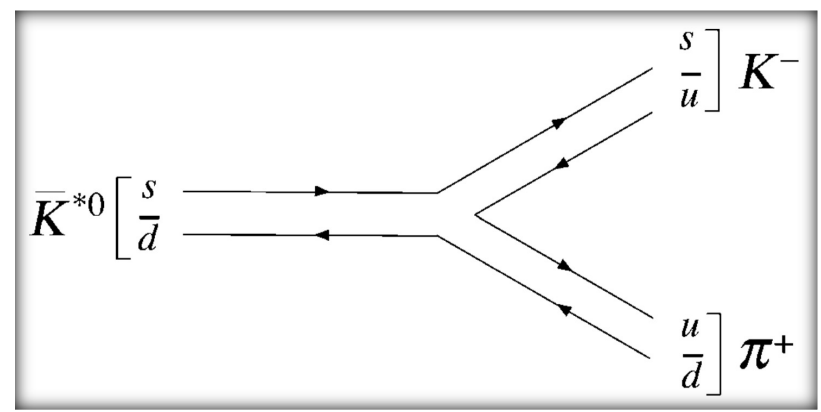
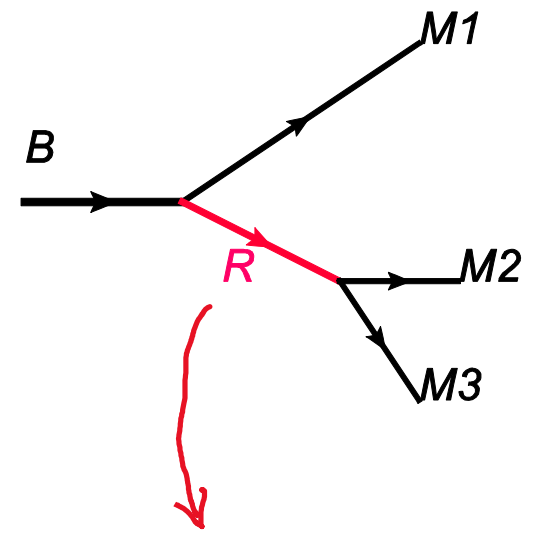
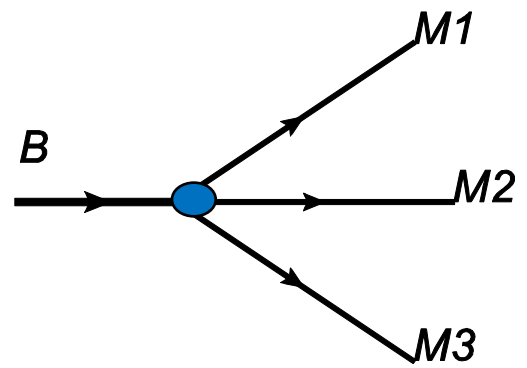
2.

$$\mathcal{A}_{\text{total}} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$



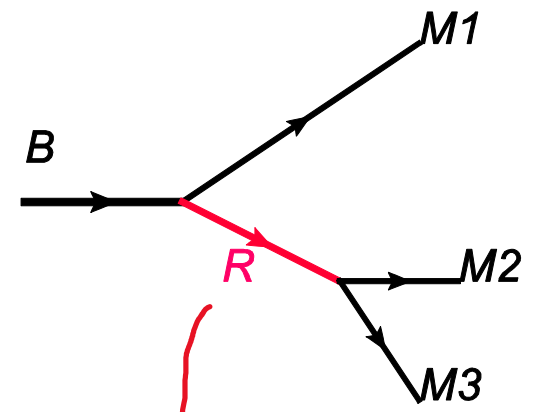
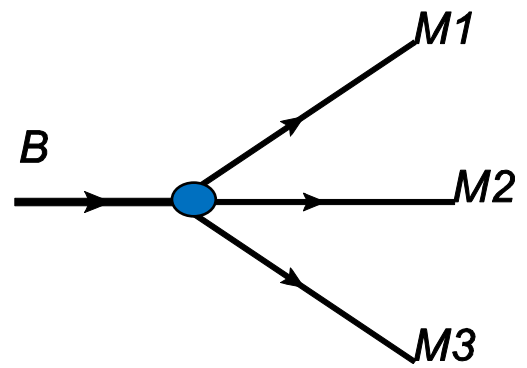
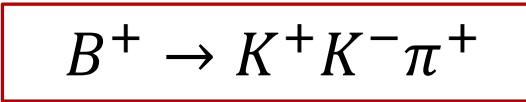
3-body B and D meson decays and the virtual contribution

$$B^+ \rightarrow K^+ K^- \pi^+$$

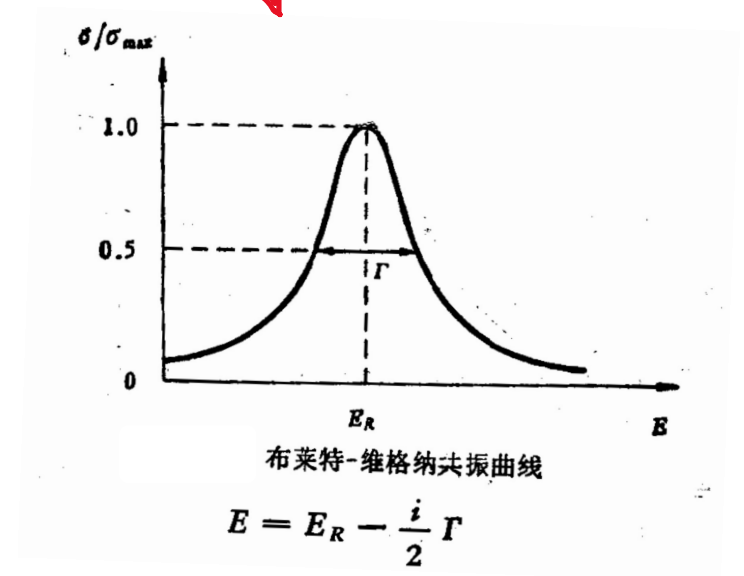




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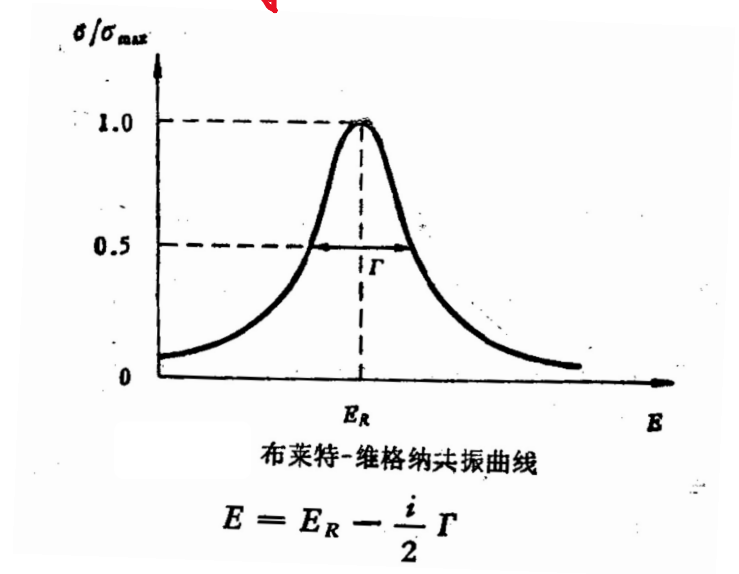
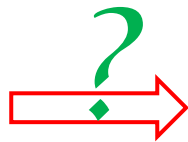
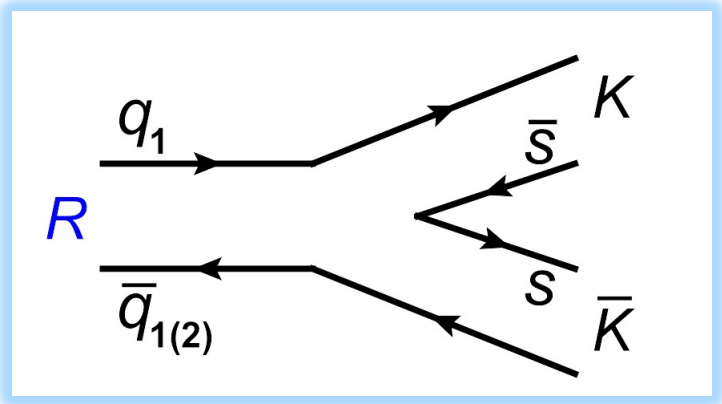
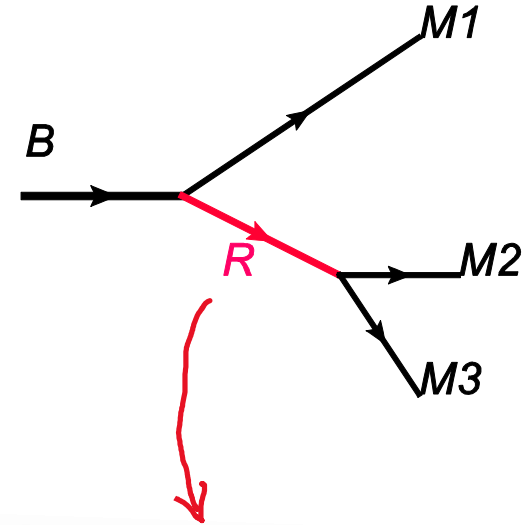
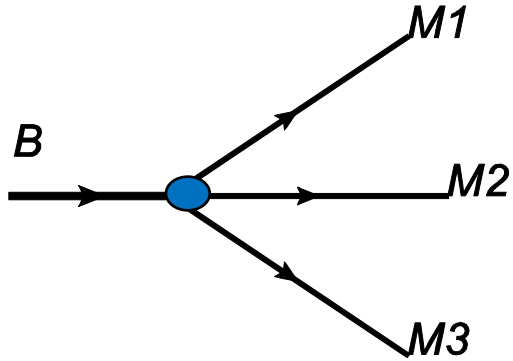


$$BW_R = \frac{m_R^2}{m_R^2 - s - im_R \Gamma_R(s)}$$



$$B^+ \rightarrow K^+ K^- \pi^+$$

$$\rho(770)^0 \rightarrow K^+ K^-$$





PHYSICAL REVIEW D

VOLUME 15, NUMBER 11

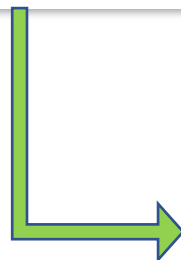
1 JUNE 1977

High-statistics study of the reactions $\pi^- p \rightarrow K^- K^+ n$ and $\pi^+ n \rightarrow K^- K^+ p$ at 6 GeV/c*

A. J. Pawlicki, D. S. Ayres, D. Cohen, R. Diebold, S. L. Kramer, and A. B. Wicklund

Argonne National Laboratory, Argonne, Illinois 60439

(Received 23 December 1976; revised manuscript received 15 March 1977)



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Physics Letters B

Volume 25, Issue 4, 4 September 1967, Pages 294-297



Further study of the $I = 1 K \bar{K}$ structure near threshold

A. Astier, J. Cohen-Ganouna, M. Della Negra, B. Maréchal, L. Montanet, M. Tomas[†], M. Baubillier, J. Duboc

3) The $I = 1 K \bar{K}$ channel is dominated, at threshold, by a virtual bound state resonance



PHYSICAL REVIEW D **94**, 072001 (2016)

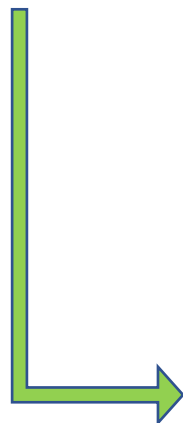
Amplitude analysis of $B^- \rightarrow D^+ \pi^- \pi^-$ decays

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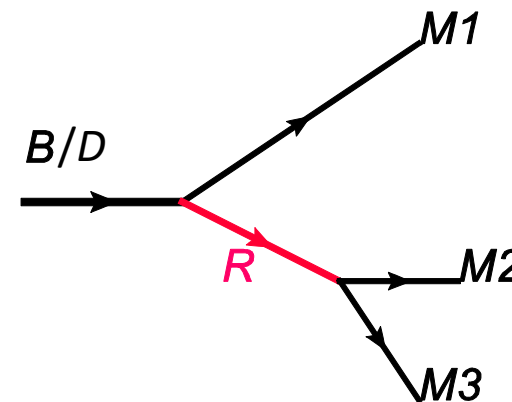
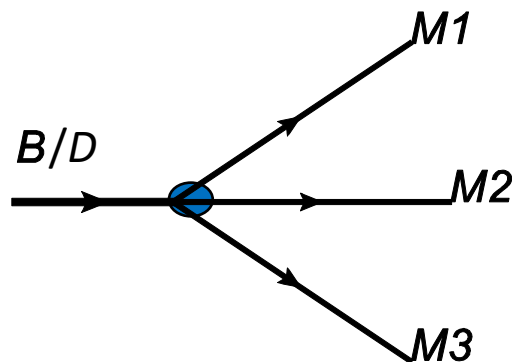


Contribution	Fit fraction (%)
$D_2^*(2460)^0$	35.7 ± 0.6
$D_1^*(2680)^0$	8.3 ± 0.6
$D_3^*(2760)^0$	1.0 ± 0.1
$D_2^*(3000)^0$	0.23 ± 0.07
$D_v^*(2007)^0$	10.8 ± 0.7
B_v^{*0}	2.7 ± 1.0
Total S wave	57.0 ± 0.8
Total fit fraction	115.7



3-body B and D meson decays and the virtual contribution

B&D → M1 M2 M3



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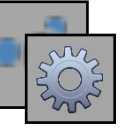
The total amplitude within isobar approach:

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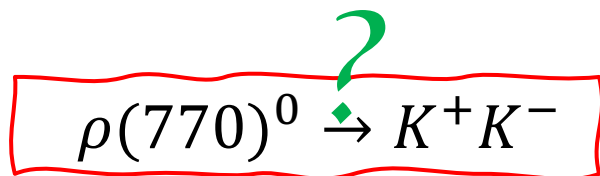
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Example - 1

PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

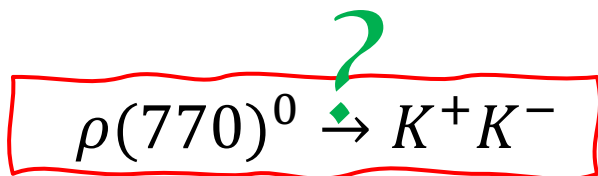
Amplitude Analysis of $B^\pm \rightarrow \pi^\pm K^+ K^-$ Decays

R. Aaij *et al.**
(LHCb Collaboration)

(Received 12 June 2019; revised manuscript received 15 October 2019; published 6 December 2019)

The first amplitude analysis of the $B^\pm \rightarrow \pi^\pm K^+ K^-$ decay is reported based on a data sample corresponding to an integrated luminosity of 3.0 fb^{-1} of pp collisions recorded in 2011 and 2012 with the LHCb detector. The data are found to be best described by a coherent sum of five resonant structures plus a nonresonant component and a contribution from $\pi\pi \leftrightarrow KK$ S -wave rescattering. The dominant contributions in the $\pi^\pm K^\mp$ and $K^+ K^-$ systems are the nonresonant and the $B^\pm \rightarrow \rho(1450)^0 \pi^\pm$ amplitudes, respectively, with fit fractions around 30%. For the rescattering contribution, a sizable fit fraction is observed. This component has the largest CP asymmetry reported to date for a single amplitude of $(-66 \pm 4 \pm 2)\%$, where the first uncertainty is statistical and the second systematic. No significant CP violation is observed in the other contributions.

DOI: [10.1103/PhysRevLett.123.231802](https://doi.org/10.1103/PhysRevLett.123.231802)



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$(30.7 \pm 1.2 \pm 0.9)\%$

kaon pair. This fit fraction implies a branching fraction $(1.60 \pm 0.14) \times 10^{-6}$ for the quasi-two-body decay $B^+ \rightarrow \pi^+ \rho(1450)^0 \rightarrow \pi^+ K^+ K^-$ [15]; this is in view of the

$$\rho(770)^0 \rightarrow K^+ K^-$$

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DOI: [10.1103/PhysRevLett.123.231802](https://doi.org/10.1103/PhysRevLett.123.231802)

PHYSICAL REVIEW D **101**, 111901(R) (2020)

Rapid Communications

Will the subprocesses $\rho(770, 1450)^0 \rightarrow K^+ K^-$ contribute large branching fractions for $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays?

Wen-Fei Wang ^{*}

Institute of Theoretical Physics, Shanxi University, Taiyuan, Shanxi 030006, China

(Received 21 April 2020; accepted 28 May 2020; published 9 June 2020)

We analyze the quasi-two-body decays $B^\pm \rightarrow \pi^\pm \rho(770, 1450)^0 \rightarrow \pi^\pm K^+ K^-$ in the perturbative QCD approach. The results in this work do not support that large branching fractions contributed by the resonances $\rho(770, 1450)^0$ in the $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays. The virtual contribution for $K^+ K^-$ from the tail of the resonance $\rho(770)^0$ which has been ignored in the experimental studies is about 1.5 times of the $\rho(1450)^0 \rightarrow K^+ K^-$ contribution, with the predicted branching fractions $\mathcal{B}_v = (1.31 \pm 0.27) \times 10^{-7}$ and $\mathcal{B} = (8.96 \pm 2.61) \times 10^{-8}$, respectively, for these two subprocesses in the $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays. The absence of $\rho(770)^0 \rightarrow K^+ K^-$ for the decay amplitude of a three-body hadronic B decay involving charged kaon pair could probably result in a larger proportion for the contribution from the resonance $\rho(1450)^0$ in experimental analysis.

DOI: [10.1103/PhysRevD.101.111901](https://doi.org/10.1103/PhysRevD.101.111901)

$$\rho(770)^0 \rightarrow K^+ K^-$$

Example - 1

PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

Amplitude Analysis of B^\pm

Rapid Communications

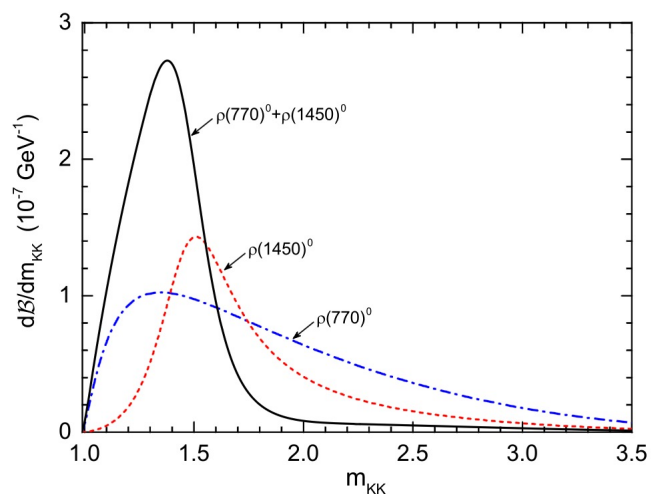


FIG. 2. The differential branching fractions for the quasi-two-body decays $B^\pm \rightarrow \pi^\pm \rho(770)^0 \rightarrow \pi^\pm K^+ K^-$, $B^\pm \rightarrow \pi^\pm \rho(1450)^0 \rightarrow \pi^\pm K^+ K^-$, and $B^\pm \rightarrow \pi^\pm (\rho(770)^0 + \rho(1450)^0) \rightarrow \pi^\pm K^+ K^-$.

PHYSICAL REVIEW D **101**, 111901(R) (2020)

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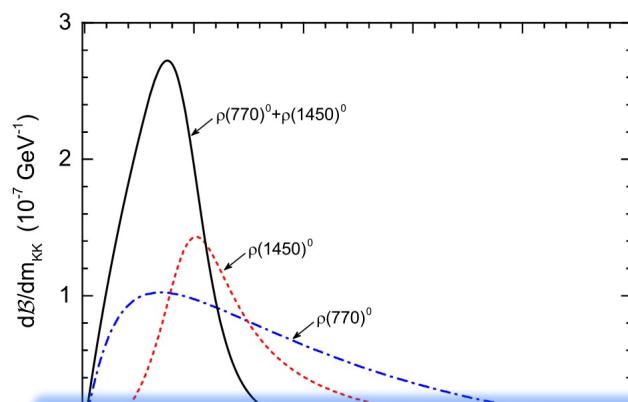
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PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

Amplitude Analysis of $B^{\pm} \rightarrow \pi^{\pm} K^+ K^-$

Rapid Communications



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PHYSICAL REVIEW D **103**, 056021 (2021)

Contributions for the kaon pair from $\rho(770)$, $\omega(782)$ and their excited states in the $B \rightarrow K\bar{K}h$ decays

Wen-Fei Wang^{1,2,*}

$= (1.31 \pm 0.27) \times 10^{-7}$ and $B^{\pm} \rightarrow \pi^{\pm} K^+ K^-$ decays. The virtual B decay involving charged resonances from the resonance $\rho(1450)^0$ in

FIG
qua
 B^{\pm}
 $\rho(1450)^0$

Example - 2

PHYSICAL REVIEW D **92**, 012013 (2015)

Study of D^{**} production and light hadronic states in the $\bar{B}^0 \rightarrow D^{*+} \omega \pi^-$ decay

D. Matvienko,^{4,50} A. Kuzmin,^{4,50} S. Eidelman,^{4,50} A. Abdesselam,⁶¹ I. Adachi,^{15,11} H. Aihara,⁶⁶ S. Al Said,^{61,30}
K. Arinstein,^{4,50} D. M. Asner,⁵² V. Aulchenko,^{4,50} T. Aushev,^{41,24} R. Ayad,⁶¹ V. Babu,⁶² I. Badhrees,^{61,29} S. Bahinipati,¹⁷
A. M. Bakich,⁶⁰ V. Bansal,⁵² V. Bhardwaj,⁵⁸ B. Bhuyan,¹⁸ J. Biswal,²⁵ A. Bobrov,^{4,50} A. Bondar,^{4,50} G. Bonvicini,⁷⁰
A. Bozek,⁴⁸ M. Bračko,^{37,25} T. E. Browder,¹⁴ D. Červenkov,⁵ A. Chen,⁴⁵ B. G. Cheon,¹³ K. Chilikin,²⁴ R. Chistov,²⁴
K. Cho,³¹ V. Chobanova,³⁸ S. -K. Choi,¹² Y. Choi,⁵⁹ D. Cinabro,⁷⁰ J. Dalseno,^{38,63} J. Dingfelder,³ Z. Doležal,⁵ Z. Drásal,⁵
A. Drutskoy,^{24,40} D. Dutta,⁶² D. Epifanov,⁶⁶ H. Farhat,⁷⁰ J. E. Fast,⁵² T. Ferber,⁸ B. G. Fulsom,⁵² V. Gaur,⁶² N. Gabyshev,^{4,50}

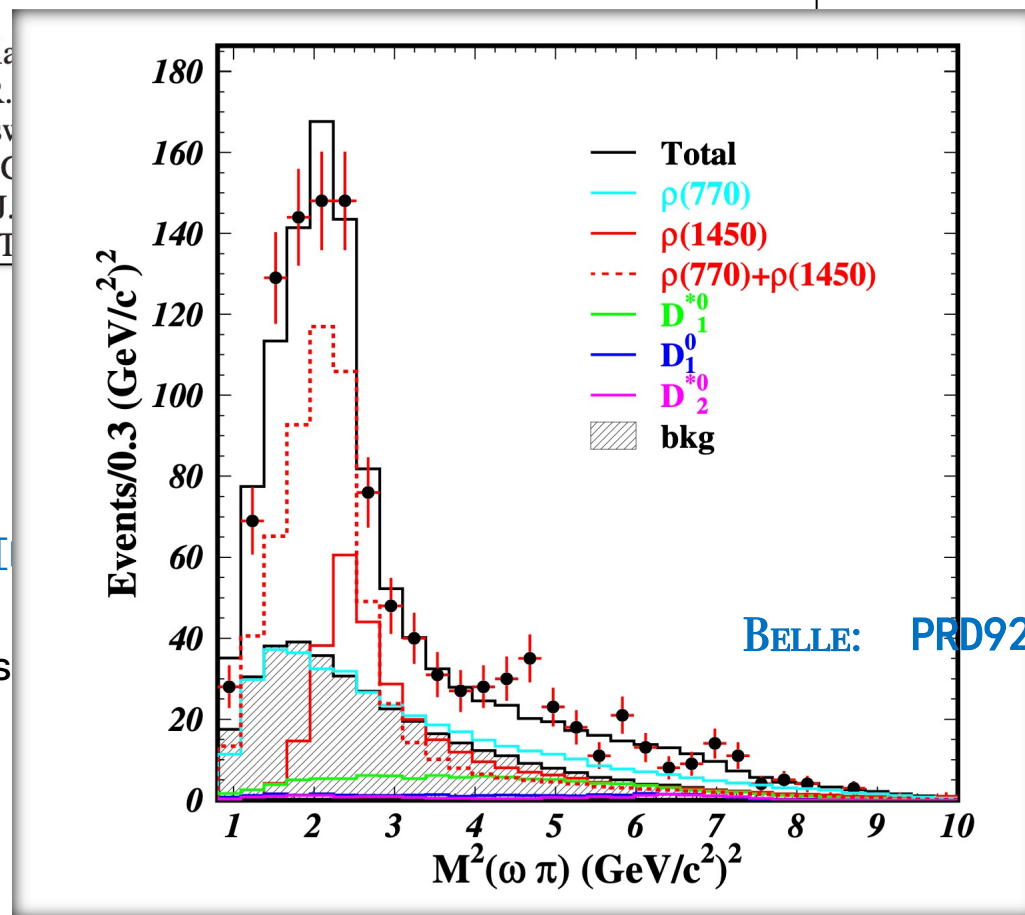
- ❖ $B^+ \rightarrow \bar{D}^{(*)0} \omega \pi^+$ and $B^0 \rightarrow D^{(*)-} \omega \pi^+$ have been measured but without any theoretical predictions [\[PRD64-092001, PRD74-012001, PRD92-012013\]](#)
- ❖ The $\omega \pi^+$ is related to the resonances $\rho(1450)$ and $\rho(770)$

Example - 2

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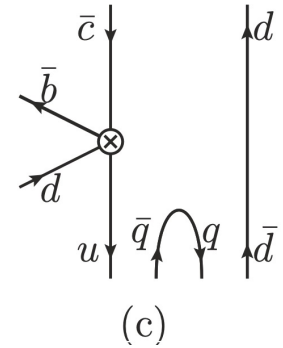
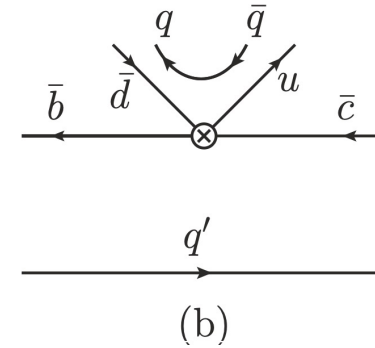
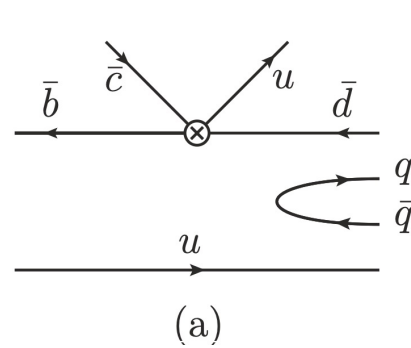
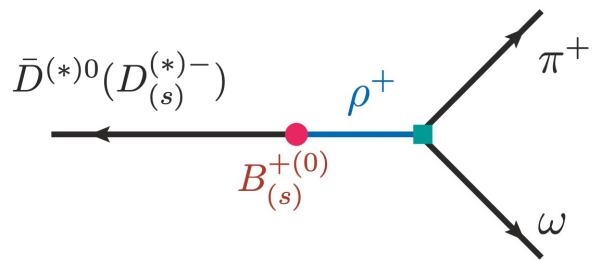


❖ $B^+ \rightarrow \bar{D}^{(*)0} \omega \pi^+$ and $B^0 \rightarrow D^{(*)-} \omega \pi^+$ have theoretical predictions

❖ The $\omega \pi^+$ is related to the resonances

Example - 2

$\rho \rightarrow \omega\pi$

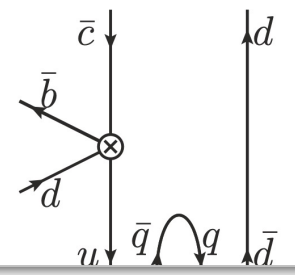
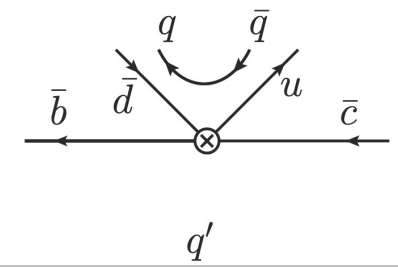
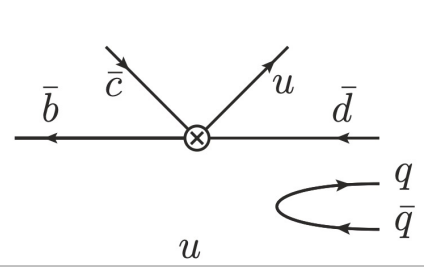
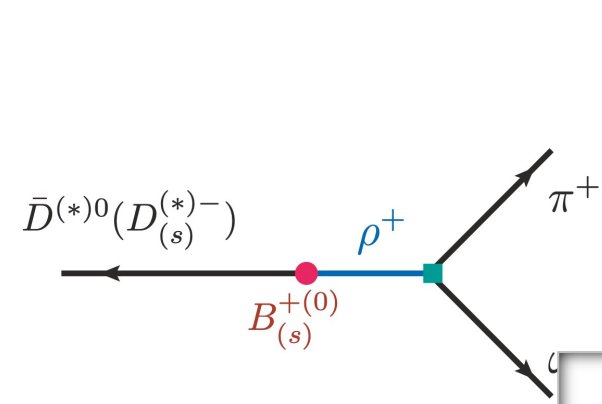


With PQCD approach

$$\begin{aligned} \mathcal{A} &= \langle (\omega\pi)_{P\text{-wave}} D^{(*)} | \mathcal{H}_{\text{eff}} | B \rangle \\ &= \phi_B \otimes \mathcal{H} \otimes \phi_{\omega\pi}^{P\text{-wave}} \otimes \phi_{D^{(*)}} \end{aligned}$$

$\rho \rightarrow \omega\pi$

Example - 2



Physics Letters B 763 (2016) 29–39

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With PQCD approach

Quasi-two-body decays $B \rightarrow K\rho \rightarrow K\pi\pi$ in perturbative QCD approach

Wen-Fei Wang^{a,b}, Hsiang-nan Li^{a,*}





PUBLISHED FOR SISSA BY SPRINGER

RECEIVED: November 3, 2023

REVISED: December 13, 2023

ACCEPTED: December 21, 2023

PUBLISHED: January 10, 2024

The $\rho(770, 1450) \rightarrow \omega\pi$ contributions for three-body decays $B \rightarrow \bar{D}^{(*)}\omega\pi$

Yu-Shan Ren ^{a,c}, Ai-Jun Ma ^b and Wen-Fei Wang ^{d,a,*}

^a*Institute of Theoretical Physics and State Key Laboratory of Quantum Optics and Quantum Optics Devices, Shanxi University, Taiyuan, Shanxi 030006, China*

^b*School of Mathematics and Physics, Nanjing Institute of Technology, Nanjing, Jiangsu 211167, China*

^c*School of Physics, University of Electronic Science and Technology of China, Chengdu 610054, China*

^d*Departament de Física Quàntica i Astrofísica and Institut de Ciències del Cosmos (ICCUB), Facultat de Física, Universitat de Barcelona, Martí i Franquès 1, 08028, Barcelona, Spain*

E-mail: ryskxky@yeah.net, theoma@163.com, wfwang@ub.edu

ABSTRACT: The decays $B \rightarrow \bar{D}^{(*)}\omega\pi$ are very important for the investigation of ρ excitations and the test of factorization hypothesis for B meson decays. The $B^+ \rightarrow \bar{D}^{(*)0}\omega\pi^+$ and $B^0 \rightarrow D^{(*)-}\omega\pi^+$ have been measured by different collaborations but without any predictions for their observables on theoretical side. In this work, we study the contributions of $\rho(770, 1450) \rightarrow \omega\pi$ for the cascade decays $B^+ \rightarrow \bar{D}^{(*)0}\rho^+ \rightarrow \bar{D}^{(*)0}\omega\pi^+$, $B^0 \rightarrow D^{(*)-}\rho^+ \rightarrow D^{(*)-}\omega\pi^+$ and $B_s^0 \rightarrow D_s^{(*)-}\rho^+ \rightarrow D^{(*)-}\omega\pi^+$. We introduce $\rho(770, 1450) \rightarrow \omega\pi$ subprocesses into the distribution amplitudes for $\omega\pi$ system via the vector form factor $F_{\omega\pi}(s)$ and then predict the branching

JHEP01(2024)047

Example - 2

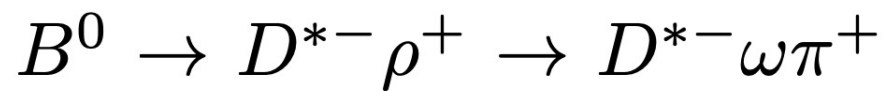
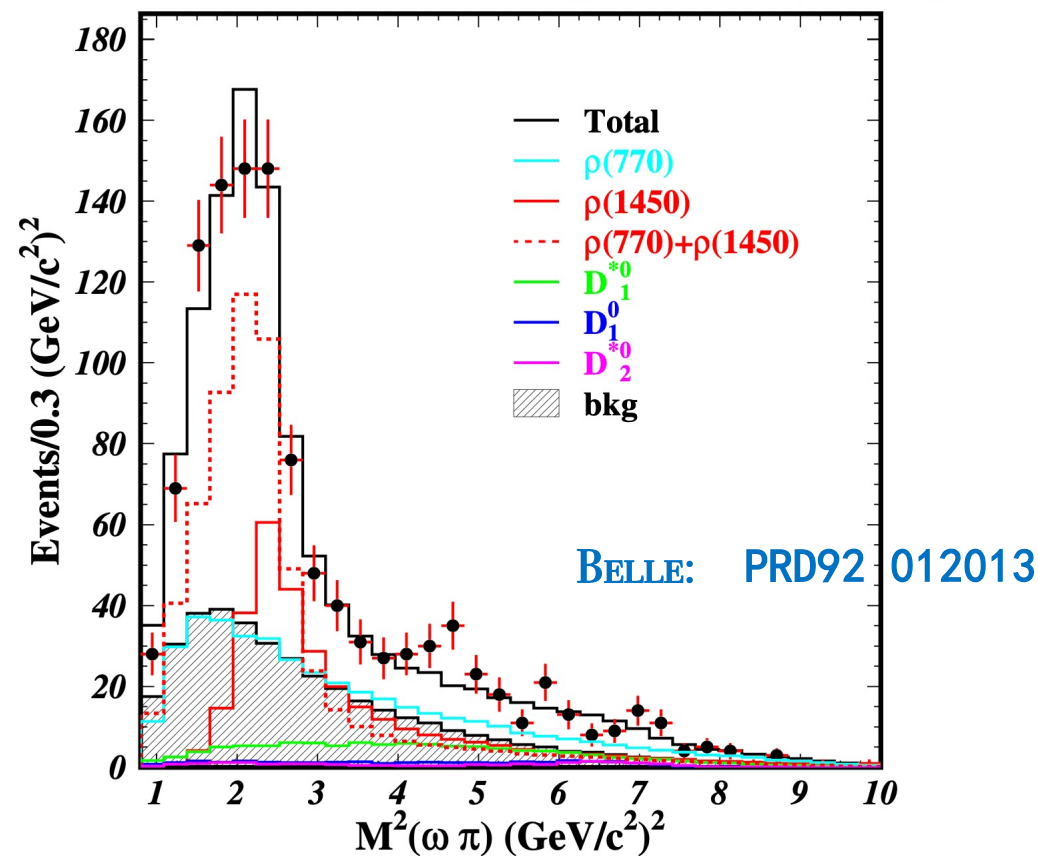
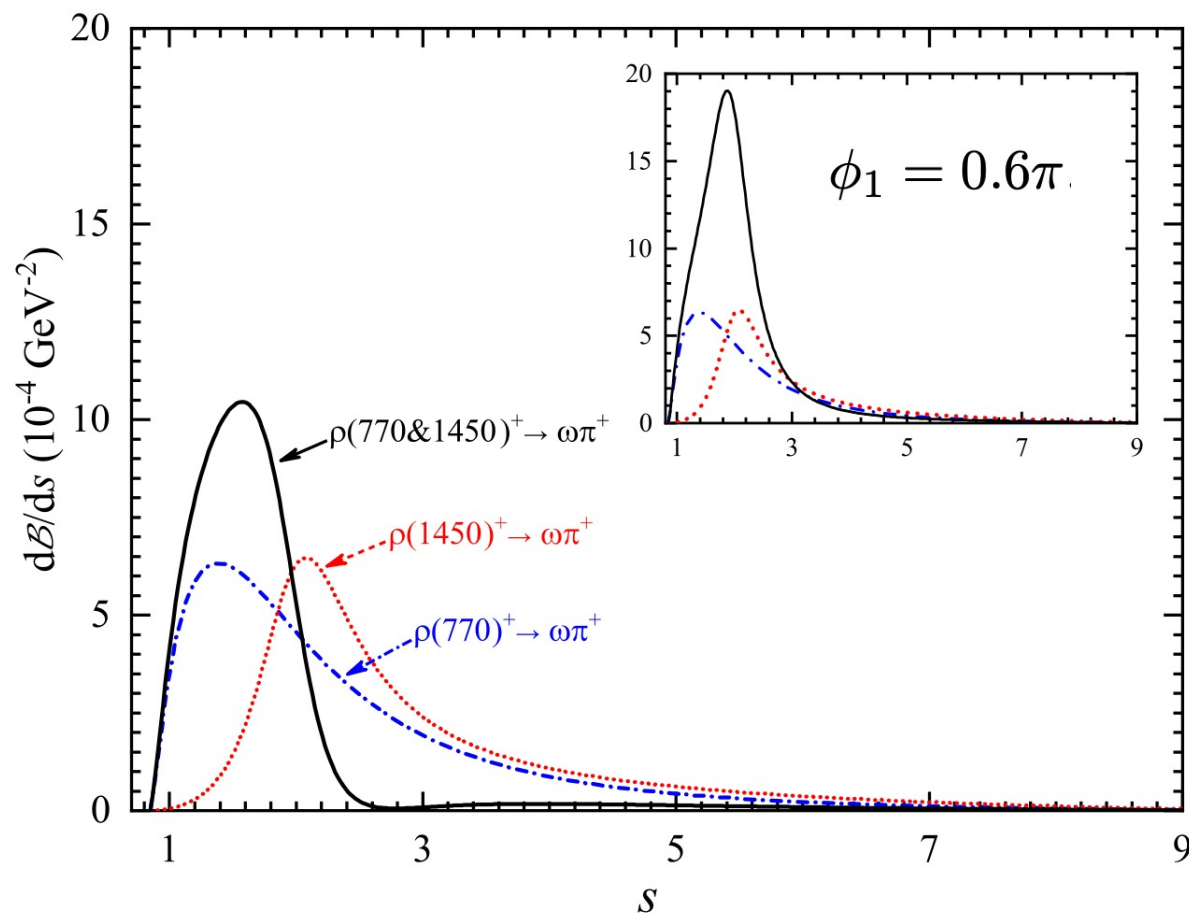
Example - 2

Decay modes	\mathcal{B} (in 10^{-3})	Γ_L/Γ
$B^+ \rightarrow \bar{D}^{*0}[\rho(770)^+ \rightarrow]\omega\pi^+$	$1.21^{+0.17+0.09+0.05+0.07}_{-0.17-0.09-0.03-0.07}$	$0.74^{+0.02}_{-0.02}$
$B^+ \rightarrow \bar{D}^{*0}[\rho(1450)^+ \rightarrow]\omega\pi^+$	$0.87^{+0.12+0.07+0.03+0.37}_{-0.12-0.07-0.02-0.37}$	$0.67^{+0.02}_{-0.02}$
$B^0 \rightarrow D^{*-}[\rho(770)^+ \rightarrow]\omega\pi^+$	$1.20^{+0.18+0.09+0.02+0.07}_{-0.18-0.08-0.01-0.07}$	$0.68^{+0.02}_{-0.02}$
$B^0 \rightarrow D^{*-}[\rho(1450)^+ \rightarrow]\omega\pi^+$	$0.89^{+0.13+0.06+0.02+0.38}_{-0.13-0.06-0.02-0.38}$	$0.63^{+0.01}_{-0.01}$
$B_s^0 \rightarrow D_s^{*-}[\rho(770)^+ \rightarrow]\pi\pi^+$	$1.03^{+0.11+0.08+0.00+0.05}_{-0.11-0.08-0.00-0.05}$	$0.65^{+0.01}_{-0.01}$
$B_s^0 \rightarrow D_s^{*-}[\rho(1450)^+ \rightarrow]\pi\pi^+$	$0.77^{+0.08+0.06+0.00+0.32}_{-0.08-0.06-0.00-0.32}$	$0.59^{+0.01}_{-0.01}$

BELLE: PRD92-012013

$$\left\{ \begin{array}{l} \mathcal{B} = (1.48 \pm 0.27^{+0.15+0.21}_{-0.09-0.56}) \times 10^{-3} \\ \mathcal{B} = (1.07^{+0.15+0.06+0.40}_{-0.31-0.13-0.02}) \times 10^{-3} \end{array} \right.$$

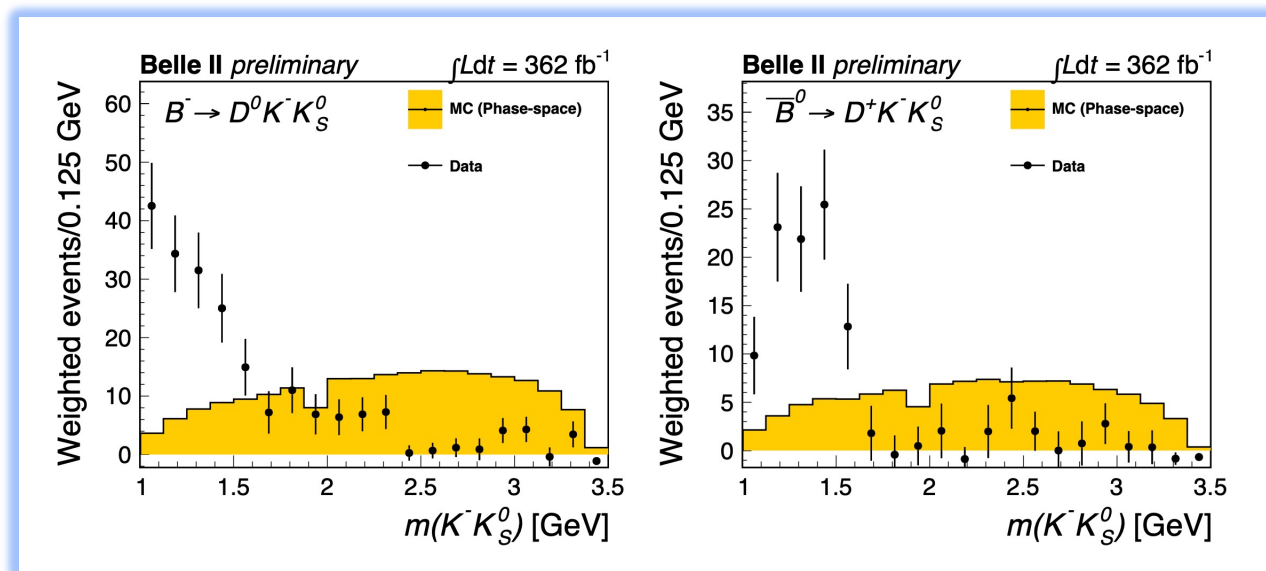
Example - 2



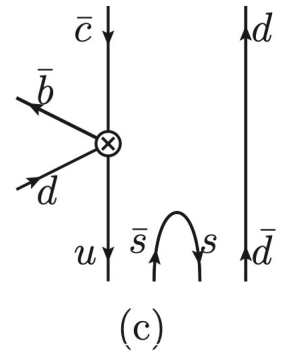
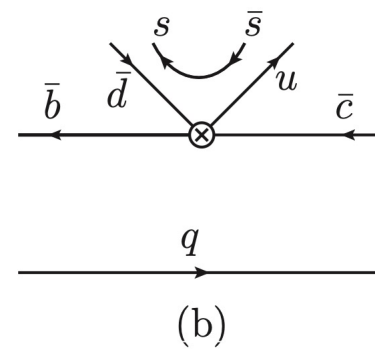
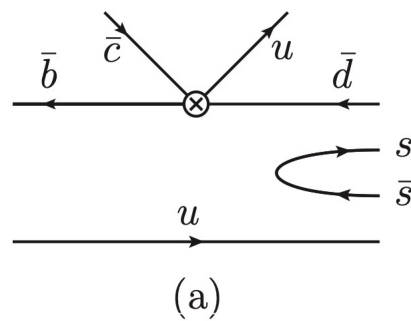
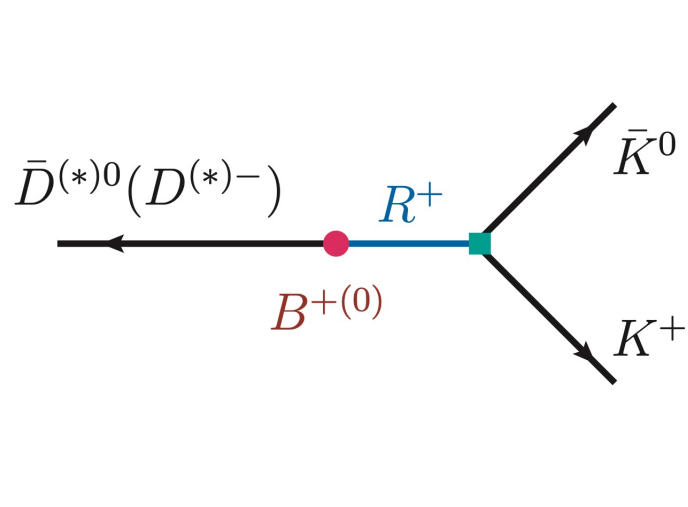
Belle-II arXiv:2305.01321

Example - 3

$$\begin{aligned}\mathcal{B}(B^- \rightarrow D^0 K^- K_S^0) &= (1.89 \pm 0.16 \pm 0.10) \times 10^{-4}, \\ \mathcal{B}(\bar{B}^0 \rightarrow D^+ K^- K_S^0) &= (0.85 \pm 0.11 \pm 0.05) \times 10^{-4}, \\ \mathcal{B}(B^- \rightarrow D^{*0} K^- K_S^0) &= (1.57 \pm 0.27 \pm 0.12) \times 10^{-4}, \\ \mathcal{B}(\bar{B}^0 \rightarrow D^{*+} K^- K_S^0) &= (0.96 \pm 0.18 \pm 0.06) \times 10^{-4},\end{aligned}$$



Example - 3



Isospin $I = 1$ $K^+ \bar{K}^0$

- $\rho(770)^+ & \rho(1450)^+ & \rho(1700)^+$
- $a_0(980)^+ & a_0(1450)^+$
- $a_2(1320)^+ & a_2(1700)^+$



**Low-mass enhancement of kaon pairs
in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays**

Wen-Fei Wang^{1,2,*} Li-Fei Yang,¹ Ai-Jun Ma,³ and Àngels Ramos^{2,†}

Example - 3

TABLE III: PQCD predictions for the branching fractions of the concerned quasi-two-body decays with the subprocess $\rho^+ \rightarrow K^+ \bar{K}^0$, here $\rho^+ = \rho(770)^+ + \rho(1450)^+$.

Mode	Unit	\mathcal{B}
$B^+ \rightarrow \bar{D}^0[\rho^+ \rightarrow] K^+ \bar{K}^0$	10^{-4}	$1.68^{+0.20+0.17+0.12}_{-0.20-0.15-0.12}$
$B^0 \rightarrow D^-[\rho^+ \rightarrow] K^+ \bar{K}^0$	10^{-4}	$0.98^{+0.06+0.13+0.06}_{-0.06-0.12-0.06}$
$B^+ \rightarrow \bar{D}^{*0}[\rho^+ \rightarrow] K^+ \bar{K}^0$	10^{-4}	$1.33^{+0.21+0.11+0.05}_{-0.21-0.11-0.07}$
$B^0 \rightarrow D^{*-}[\rho^+ \rightarrow] K^+ \bar{K}^0$	10^{-4}	$1.16^{+0.19+0.08+0.02}_{-0.19-0.09-0.02}$

Belle-II arXiv:2305.01321

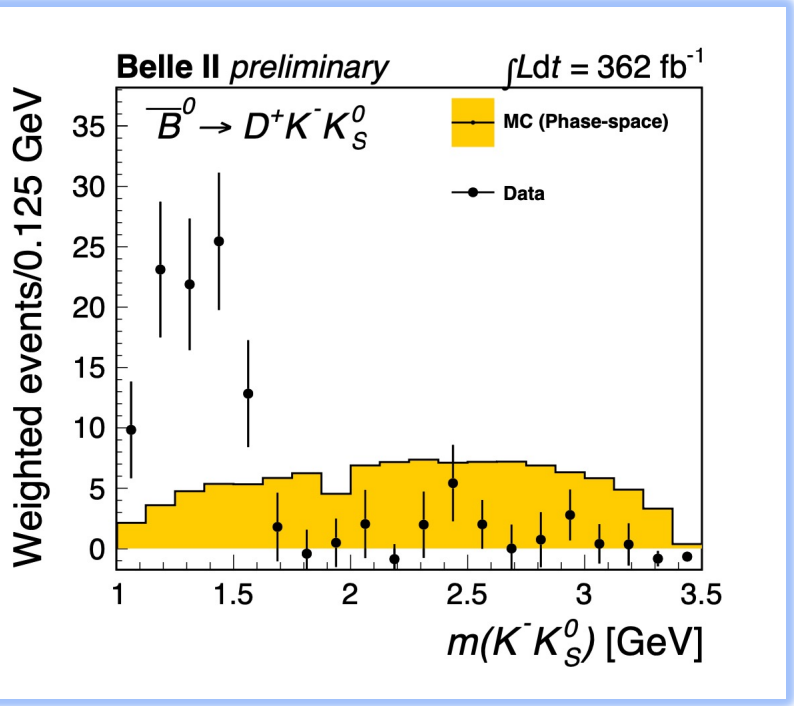
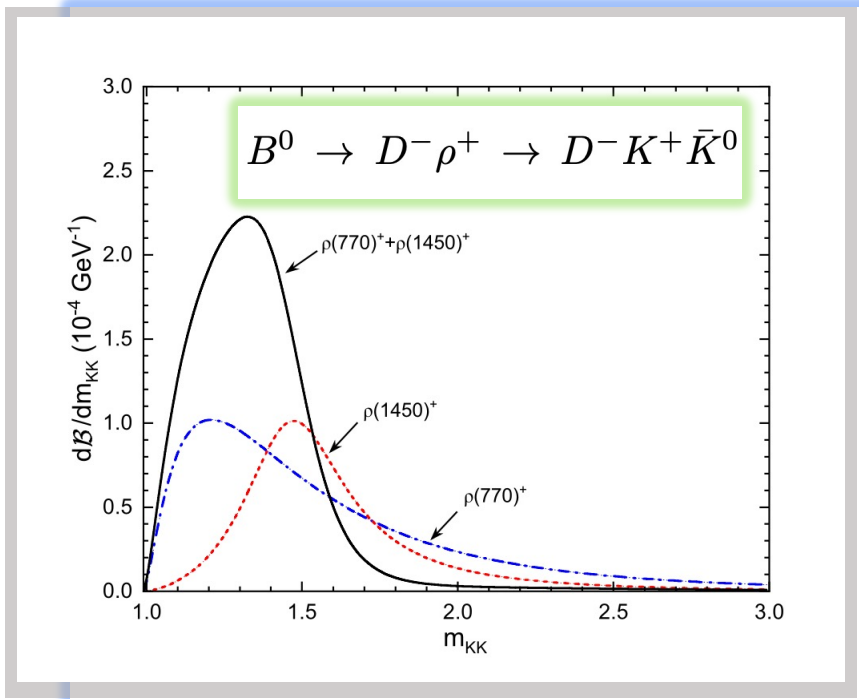
$$\mathcal{B}(B^- \rightarrow D^0 K^- K_S^0) = (1.89 \pm 0.16 \pm 0.10) \times 10^{-4},$$

$$\mathcal{B}(\bar{B}^0 \rightarrow D^+ K^- K_S^0) = (0.85 \pm 0.11 \pm 0.05) \times 10^{-4},$$

$$\mathcal{B}(B^- \rightarrow D^{*0} K^- K_S^0) = (1.57 \pm 0.27 \pm 0.12) \times 10^{-4},$$

$$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} K^- K_S^0) = (0.96 \pm 0.18 \pm 0.06) \times 10^{-4},$$

Example - 3



PHYSICAL REVIEW D **104**, 116019 (2021)

Subprocesses $\rho(770, 1450) \rightarrow K\bar{K}$ for the three-body hadronic D meson decays

Wen-Fei Wang 

*Institute of Theoretical Physics, Shanxi University, Taiyuan, Shanxi 030006, China
and State Key Laboratory of Quantum Optics and Quantum Optics Devices, Shanxi University,
Taiyuan, Shanxi 030006, China*

 (Received 15 October 2021; accepted 2 December 2021; published 21 December 2021)

We construct the theoretical framework for quasi-two-body D meson decays with the help of pion and kaon electromagnetic form factors and with which we study the contributions of the subprocesses $\rho(770, 1450) \rightarrow K\bar{K}$ for the three-body D decays within the flavor $SU(3)$ symmetry. Because of the limitations imposed by phase space and strong coupling, the contributions for the kaon pair from the virtual bound state $\rho(770)$ are channel dependent and generally small for the concerned three-body D decays, but some quasi-two-body processes could still be observed in the Dalitz plot analyses for related decays, such as $D^0 \rightarrow K^-\rho(770)^+ \rightarrow K^-K^+K_S^0$ and $D^+ \rightarrow K_S^0\rho(770)^+ \rightarrow K_S^0K^+K_S^0$, they are predicted to have the branching fractions $\mathcal{B} = (0.82 \pm 0.04) \times 10^{-4}$ and $\mathcal{B} = 0.47^{+0.05}_{-0.03} \times 10^{-4}$, which are $(1.86 \pm 0.16)\%$ and $(1.84^{+0.21}_{-0.16})\%$, respectively, of the total branching fractions for the corresponding three-body D decays. We find in this work that the normal subprocesses like $\rho(1450)^+ \rightarrow \pi^+\pi^0$ or $\rho(1450)^+ \rightarrow K^+\bar{K}^0$, which are bound by the masses of decaying initial states, will provide virtual contributions in some special decays.

DOI: [10.1103/PhysRevD.104.116019](https://doi.org/10.1103/PhysRevD.104.116019)

Example - 4

Example - 5

PHYSICAL REVIEW D **107**, 052010 (2023)

Observation of the decay $D_s^+ \rightarrow \omega\pi^+\eta$

M. Ablikim,¹ M. N. Achasov,^{12,b} P. Adlarson,⁷² M. Albrecht,⁴ R. Aliberti,³³ A. Amoroso,^{71a,71c} M. R. An,³⁷ Q. An,^{68,55}
Y. Bai,⁵⁴ O. Bakina,³⁴ R. Baldini Ferroli,^{27a} I. Balossino,^{28a} Y. Ban,^{44,g} V. Batozskaya,^{1,42} D. Becker,³³ K. Begzsuren,³⁰
N. Berger,³³ M. Bertani,^{27a} D. Bettoni,^{28a} F. Bianchi,^{71a,71c} E. Bianco,^{71a,71c} J. Bloms,⁶⁵ A. Bortone,^{71a,71c} I. Boyko,³⁴
R. A. Briere,⁵ A. Brueggemann,⁶⁵ H. Cai,⁷³ X. Cai,^{1,55} A. Calcaterra,^{27a} G. F. Cao,^{1,60} N. Cao,^{1,60} S. A. Cetin,^{59a}
J. F. Chang,^{1,55} W. L. Chang,^{1,60} G. R. Che,⁴¹ G. Chelkov,^{34,a} C. Chen,⁴¹ Chao Chen,⁵² G. Chen,¹ H. S. Chen,^{1,60}

M. L. ...
W. S. ...
R. E. d...

M. ABLIKIM *et al.*

PHYS. REV. D **107**, 052010 (2023)



(Received 9 February 2023; accepted 13 March 2023; published 27 March 2023)

Using 7.33 fb^{-1} of e^+e^- collision data collected by the BESIII detector at c.m. energies between 4.128 GeV and 4.226 GeV, we observe for the first time the decay $D_s^\pm \rightarrow \omega\pi^\pm\eta$ with a statistical significance of 7.6σ . The measured branching fraction of this decay is $(0.54 \pm 0.12 \pm 0.04)\%$, where the first uncertainty is statistical and the second is systematic.

DOI: [10.1103/PhysRevD.107.052010](https://doi.org/10.1103/PhysRevD.107.052010)

Example - 5

M. ABLIKIM *et al.*

PHYS. REV. D **107**, 052010 (2023)

 (Received 9 February 2023; accepted 13 March 2023; published 27 March 2023)

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DOI: 10.1103/PhysRevD.107.052010

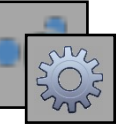
PDG-2024

74 Meson Summary Table

$2\pi^+\pi^-\pi^0$	—	935
$\eta\pi^+$	[ddd] (1.67 ± 0.09) %	S=1.1 902
$\omega\pi^+$	[ddd] (1.92 ± 0.30) $\times 10^{-3}$	822
$\omega\pi^+, \omega \rightarrow \pi^+\pi^-$	(3.9 ± 0.5) $\times 10^{-5}$	—
$3\pi^+2\pi^-$	(7.8 ± 0.8) $\times 10^{-3}$	899
$2\pi^+\pi^-2\pi^0$	—	902
$\eta\rho^+$	[ddd] (8.9 ± 0.8) %	724



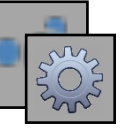
Outline



- 3-body B and D meson decays and the virtual contribution
- Some examples of the resonance virtual contributions
- **Conclusions**



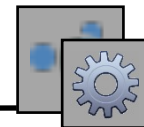
Conclusions



1.
$$\mathcal{A}_{total} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$

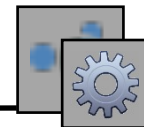


Conclusions



1.
$$\mathcal{A}_{total} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$

2. **Virtual contributions from various resonances exist in three-body B & D meson decays**

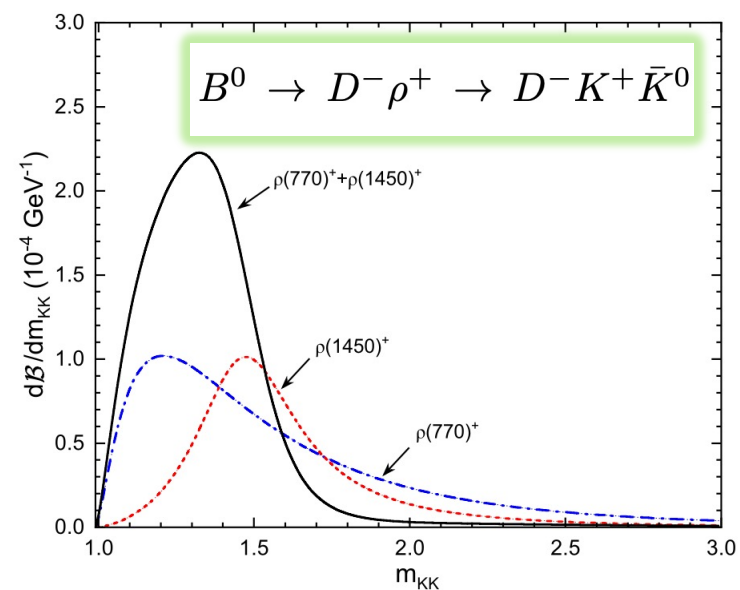


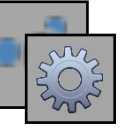
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$$\mathcal{A}_{total} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$

2. Virtual contributions from various resonances exist in three-body B & D meson decays

3.
$$D_{\rho_i}(s) = m_{\rho_i}^2 - s - i\sqrt{s}\Gamma_{\rho_i}(s)$$

4. A bump does not always mean a new state



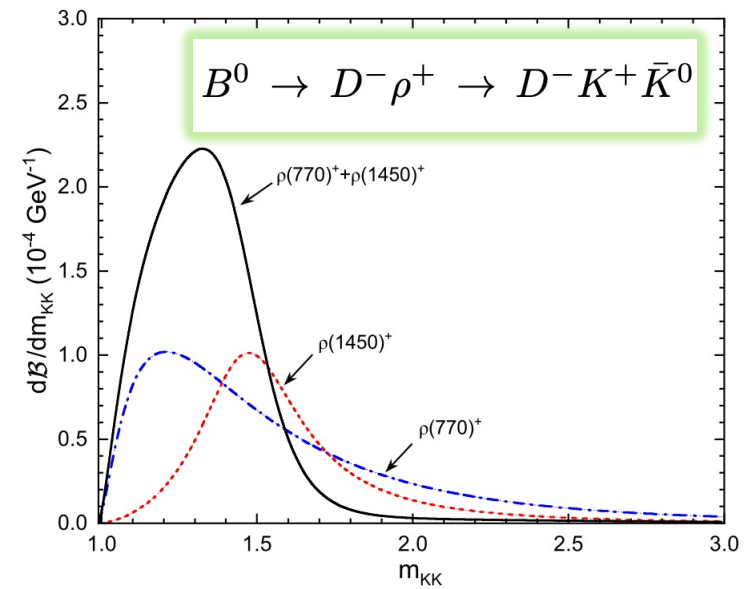


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2. Virtual contributions from various resonances exist in three-body B & D meson decays

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Thank You!