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### Motivation

### Method & Material

### Results

### Summary





## Motivation







### How are elements heavier than iron produced in the universe? ——One of the unsolved problems of physics in 21st century



**The Big Bang Theory** can explain the origin of light elements such as H, He, Li et al.













- How are elements heavier than iron produced in the universe?
  - ——One of the unsolved problems of physics in 21st century
    - The origin of the element from cabon to iron can be
    - explained by the nuclear fusion process in star.
      - For the superiron elements, two factors have caused
      - the nuclear fusion process in stars unable to
      - produce them.
      - **Coulomb barrier**
      - **Binding energy**







### How are elements heavier than iron produced in the universe? ——One of the unsolved problems of physics in 21st century



Neutron capture cross section of superiron elements is the key parameters for nuclear astrophysics

- Neutron is no charged;
- The neutron capture cross section of nuclieds tend to increase with decreasing energy;

The neutron capture reaction is the main pathway for the synthesis of superiron elements in the













## Motivation: Nuclear Data

### **Experimental data is needed to clarify the differences** between evaluation database.





- **Experimental data in resonance energy region** needs to be supplemented;
- There are obvious differences among different evaluation databases, and high-priecision experimental data clarification is needed;
- The experiments in Back-n can provide the experimental support for new version of **CENDL** database.



















## Method & Material

![](_page_7_Picture_2.jpeg)

![](_page_7_Picture_3.jpeg)

## **Experimental setup**

- by the neutron capture reaction.
- technique (PHWT).
- based on R Matrix.

![](_page_8_Figure_5.jpeg)

\* Figure from reference: Q.W. Zhang et al., Acta Phys. Sin. 70(22): 222801. (2021)

![](_page_8_Picture_7.jpeg)

![](_page_8_Picture_10.jpeg)

![](_page_8_Picture_11.jpeg)

![](_page_9_Picture_0.jpeg)

### **Oerview of Experimental Research on Neutron Capture Cross Section**

### Conducted by USC&SARI at Back-n Facility(As of May.2025)

Date	Target	Diameter (mm)	Thickness (mm)	BeamPower (kW)	Publish
2019.01	<sup>197</sup> Au	50.0	1.0		X.R. Hu et al., Nucl. Sci. Tech. (2021)
	natSe	50.0	2.0	~34	X.R. Hu et al., Chin. Phys.B (2022)
	<sup>89</sup> Y	50.0	1.0		Analyzing by CDUT
2019.05	<sup>197</sup> Au	30.0	1.0		X.X. Li et al., Nuclear Techniques (in Chinese)
	<sup>nat</sup> Er	50.0	1.0	~50	X.X. Li et al., Phys.Rev.C (2021, 2022)
	<sup>nat</sup> Sm	50.0	1.0		X.X. Li et al., Nucl. Sci. Tech. (2025) (in press)
2020.10	<sup>63</sup> Cu	30.0	0.1		Analyzing
	<sup>65</sup> Cu	30.0	0.1	~00	Anaryzing
2021.04	<sup>107</sup> Ag	30.0	0.1	100	VVIi at al Chin Dhua D (2022)
	<sup>109</sup> Ag	30.0	0.1	~100	<b>A.A. LI et al., Chin. Phys.d (2022)</b>
2022.07	<sup>nat</sup> Ag	30.0	0.1	~150	Analyzing
2022.11	<sup>nat</sup> Yb	30.0	0.1	~150	Y.J. Chen., Phys. Rev. C (2025)
2023.10	<sup>65</sup> Cu	30.0	0.1	150	
	<sup>209</sup> Bi	40.0	1.0	~130	Analyzing
2024.11	<sup>45</sup> Sc	50.0	0.2, 1.0	~ 170	

## **Key Step: Background Determination**

- The general experimental process and data analysis methods have  $\bullet$ been introduced in the previous reports.
- This report will introduce something different in experimental  $\bullet$ data analysis.

### **Background composition in the** neutron capture cross-section experiment.

BKG	Meaning	Measurement
$\boldsymbol{B}_{0}$	Activation background	no beam
$B_{\text{empty}}(t_n)$	Target-independent background	n+empty
$B_{sn}(t_n)$	Scattered neutron background	n+carbon/lead
$B_{s\gamma}(t_n)$	In-beam gamma background	n+lead

$$B_{\text{total}} = B_0 + B_{\text{empty}}(t_n) + B_{sn}(t_n) + B_{s\gamma}(t_n)$$

The most complex one among these backgrounds

For in-beam  $\gamma$  background, n+lead target can determine its shape (time structure)

![](_page_10_Figure_8.jpeg)

![](_page_10_Picture_9.jpeg)

![](_page_11_Picture_0.jpeg)

### The scale of in-beam $\gamma$ background is determined

by the "black resonance" method.

![](_page_11_Figure_3.jpeg)

- Tantalum and cobalt filters, each with a thickness of 1 mm, are placed in the beam line to absorb neutrons at specific energies.
- The neutron at 4.28 eV, 10.34 eV, 132 eV and 5.02 keV  $\bullet$ will be absorbed by filter.
- The experimental results is the results of the interaction between the gamma rays and the target.

![](_page_11_Picture_8.jpeg)

![](_page_11_Picture_9.jpeg)

![](_page_11_Picture_10.jpeg)

## **Key Step:** S<sub>n</sub> for natural target

### For the caculation of S<sub>n</sub> for natural target, different references provide different methods.

### • Phys. Rev. C 85, 044615 (2012)

### $^{nat}Sn = {}^{i}Sn, i = find(NA_i = max(NA))$

As shown in Eq. (3), the efficiency, and, hence, the calculated capture yield, is inversely proportional to  $E_{c}$ . However, only a single  $S_n$  (typically chosen to be that of the most abundant isotope in the sample) can be used in weighting the data. Therefore, in the analysis of the capture data, the abundances of the other Mg isotopes in the sample must be scaled according to their  $S_n$  value. In particular, the neutron separation energies used were 7.33, 11.09, and 6.44 MeV for  ${}^{25}$ Mg,  ${}^{26}$ Mg, and  ${}^{27}$ Mg, respectively.

<sup>nat</sup>Sn is the same as that of the isotope

with the greatest abundance.

• Chin. Phys. B 31, 080101 (2022)  $natSn = mean(NA_i*iSn)$ 

the measurement of the saturated 4.9 eV resonance in the thick <sup>197</sup>Au sample.  $\varepsilon_{E_n}$  is the detection efficiency, which for weighted counts is equal to the excitation energy  $(S_n + E_n)$  according to Eq. (3), and  $S_n = 7.054$  MeV for <sup>nat</sup>Se.  $C_W$  and  $B_W$  are the weighted pulse height spectra for the <sup>nat</sup>Se sample and the total background, respectively. The neutron fluence was

<sup>nat</sup>Sn is euqal to the weighted average of

isotope Sn according to their abundances.

![](_page_12_Picture_11.jpeg)

![](_page_12_Picture_12.jpeg)

![](_page_12_Picture_13.jpeg)

## **Key Step:** S<sub>n</sub> for natural target

- In the above calculations, natural Sn is a constant independent of the neutron energy.
- Actually, for natural target, differentt resonance peaks are contributed by different isotopes. So the natural Sn should be changed in different resnance energies.
- In this work, the neutron separation energy for natural target is determined by following equation:

 $^{nat}Sn(E) = {}^{i}Sn, i = find(A_i = max(A));$ 

 $\mathbf{A} = \mathbf{N}\mathbf{A}_{\mathbf{i}}^{*} \boldsymbol{\sigma}_{\mathbf{i}}$ 

![](_page_13_Figure_6.jpeg)

![](_page_13_Picture_7.jpeg)

![](_page_13_Picture_8.jpeg)

## **Key Step:** S<sub>n</sub> for natural target

### Selection of evaluation database

The resonance positions and values can vary across different evaluation databases, which may impact the calculation results of Sn. We compared the results of mai evaluation databases and selected the appropriate ones.

Eliminate inconsistent evaluation databases based on the

positions	of the	formant	peaks in	the ex	periment
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Database	The peak position is consistent with the experiment?	Results	
ENDF/B-VIII.0	Y	Sn1(E)	
JEFF-3.3	Ν	excluded	
JENDL-5	Y	Sn2(E)	
<b>TENDL-2023</b>	Ν	excluded	
BROND-3.1	Y	Sn3(E)	

- In this case, the database of JEFF-3.3 and TENDL-2023 was excluded, and Sn1(E) == Sn2(E) == Sn3(E).
- The selection of the three evaluation databases ENDF/B-VIII.0, JENDL-5 and BROND-3.1 will not affect the result of Sn.
- The result of ENDF/B-VIII.0 is shown in the figure.

![](_page_14_Figure_9.jpeg)

![](_page_14_Picture_10.jpeg)

![](_page_14_Picture_11.jpeg)

![](_page_15_Picture_0.jpeg)

### Results

![](_page_15_Picture_3.jpeg)

![](_page_15_Picture_4.jpeg)

![](_page_16_Picture_0.jpeg)

![](_page_16_Figure_1.jpeg)

TABLE III: Resonance parameters of <sup>168</sup>Yb extracted from the SAMMY code of experimental data.

	$E_R (eV)$	$\Gamma_{\gamma} ({\rm meV})$	$\Gamma_n \text{ (meV)}$	k
This work	$0.58 \pm 0.01$	$65.83 \pm 1.27$	$1.87 \pm 0.02$	$1.82 \pm 0.15$
$\mathrm{EX}\#1$	$0.600 \pm 0.008$	$90\pm 6$	$2.66 \pm 0.26$	$2.58 \pm 0.24$
$\mathrm{EX}\#2$		-	-	-
EX#3	$0.590 \pm 0.005$	$66\pm3$	$2.2 \pm 0.1$	$2.12 \pm 0.10$
DB#1-4	0.60	64.00	2.20	2.13
DB#5	0.60	98.01	1.00	1.85
This work	$22.53 \pm 0.45$	$63.48 \pm 4.69$	$32.76 \pm 0.66$	$21.61 \pm 1.81$
$\mathrm{EX}\#1$	$22.60 \pm 0.05$	-	$10.5 \pm 1.1$	-
$\mathrm{EX}\#2$	$22.5 \pm 0.04$	-	$6.1 \pm 0.4$	-
EX#3	$22.44 \pm 0.05$	$24.6 \pm 1.0$	$172 \pm 90$	$21.52 \pm 0.99$
DB#1-4	22.50	64.00	25.60	18.29
DB#5	22.50	104.34	40.84	29.34
This work	$27.27 \pm 0.55$	$65.88 \pm 6.12$	3.21+0.24	$3.06 \pm 0.22$
$\mathrm{EX}\#1$	$27.48 \pm 0.1$	-	$1 \pm 0.4$	-
EX#2	2 –	-	_	-
EX#3	$27.17 \pm 0.06$	-	$2.45 \pm 0.20$	-
DB#1-4	27.17	64.00	2.45	2.35
DB#5	27.17	104.34	4.76	4.55

- the R-Matrix code SAMMY was used to fit experimental yield and extract the nutron parameters;
- The existing experimental data and different evaluation data were used to compare with this work.

DB#1-5: ENDF/B-VIII.0, JENDL-4.0, BROND-3.1, JEFF-3.3, TENDL-2019

- EX#1: 1971年 V. P. Vertebnyj et al. (Neutron Phys. Conf. 1, 181 (1971))
- EX#2: 1973年 H. I. Liou et al. (Phys. Rev. C 7, 823 (1973))
- EX#3: 1981年 V. A. Anufriev et al. (Sov. At. Energy 49, 560 (1980))

### The data of DB#1-4 are consistent and have been merged.

- The results extracted from the experiment are inconsistent with all evaluation databases and existing experimental data.
- The existing experimental results are all inconsistent.
- None of the experiments is consistent with the results of any evaluation library.
- There are significant differences between DB#1-4 and DB#5.
- The first experimental extraction of the resonance parameters of <sup>168</sup>Yb at 27.3 eV.
- The experimental results are consistent with DB1-4 rather than DB5.

![](_page_16_Picture_17.jpeg)

Y.J. Chen, X.X. Li\*, W. Luo\* et al., Phys. Rev. C 111, 054607 (2025)

![](_page_16_Picture_19.jpeg)

![](_page_16_Picture_20.jpeg)

![](_page_16_Picture_21.jpeg)

![](_page_17_Picture_0.jpeg)

## Summary

![](_page_17_Picture_3.jpeg)

![](_page_17_Picture_4.jpeg)

![](_page_18_Picture_0.jpeg)

• The resonance parameter value of <sup>168</sup>Yb at 22.5 eV deserves further experimental verification. • The measurement and data analysis methods for the neutron capture cross-section of high cross-section targets have matured. • Experimental techniques and analytical methods of accurate measurement of low cross-section targets and low abundance targets are prerequisites for obtaining high-level physical results.

![](_page_18_Picture_3.jpeg)

![](_page_18_Picture_4.jpeg)

### Collaboration

![](_page_19_Picture_1.jpeg)

### W. Luo, S. Feng, X.X. Li, B. Jiang, Z. Y.J. Chen, W. Xie, Z.A. Lin, T. Wu, T

![](_page_19_Picture_3.jpeg)

![](_page_19_Picture_4.jpeg)

![](_page_19_Picture_6.jpeg)

![](_page_19_Picture_7.jpeg)

![](_page_19_Picture_9.jpeg)

![](_page_19_Picture_10.jpeg)

S. Zhang, D.X. Wang

![](_page_19_Picture_12.jpeg)

![](_page_19_Picture_13.jpeg)

J.J. He

![](_page_19_Picture_15.jpeg)

C.W. Ma

![](_page_19_Picture_17.jpeg)

	We appreciate support from all collaborators.
	We thank the staff members of the Back-n white neutron
	facility (https://cstr.cn/31113.02.CSNS.Back-n) at the Ch
CI	Spallation Neutron Source (CSNS)
. Jiang et al.	(https://cstr.cn/31113.02.CSNS), for providing technical
	support and assistance in data collection and analysis.

H.W. Wang, G.T. Fan, L.X. Liu, Y. Zhang, Z.R. Hao

J.G. Chen, X.H. Wang, J.F. Hu, B.S. Huang

![](_page_19_Picture_21.jpeg)

![](_page_19_Picture_22.jpeg)

X.R. Hu

![](_page_19_Picture_24.jpeg)

![](_page_19_Picture_25.jpeg)

![](_page_19_Picture_26.jpeg)

### University of South China

![](_page_20_Picture_1.jpeg)

## THANKS

![](_page_20_Picture_4.jpeg)

![](_page_20_Picture_5.jpeg)