

Towards High-Precision Lattice QCD Calculations of Heavy Meson LCDAs

Based on [arXiv: 25xx.xxxxx](#) & [arXiv: 2403.17492](#),
& [Phys.Rev.D 111 \(2025\) 3, 034503](#)

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Outline

- Motivation
- Theoretical framework
 - Sequential effective theory
 - Discussion of renormalization
- Lattice simulation
- Summary and prospect



Motivation

Factorization

[Beneke, Buchalla, Neubert, Sachrajda, 1999]

$$\langle \pi(p') \pi(q) | Q_i | \bar{B}(p) \rangle = f^{B \rightarrow \pi}(q^2) \int_0^1 dx T_i^{\text{I}}(x) \phi_\pi(x) + \int_0^1 d\xi dx dy T_i^{\text{II}}(\xi, x, y) \phi_B(\xi) \phi_\pi(x) \phi_\pi(y)$$

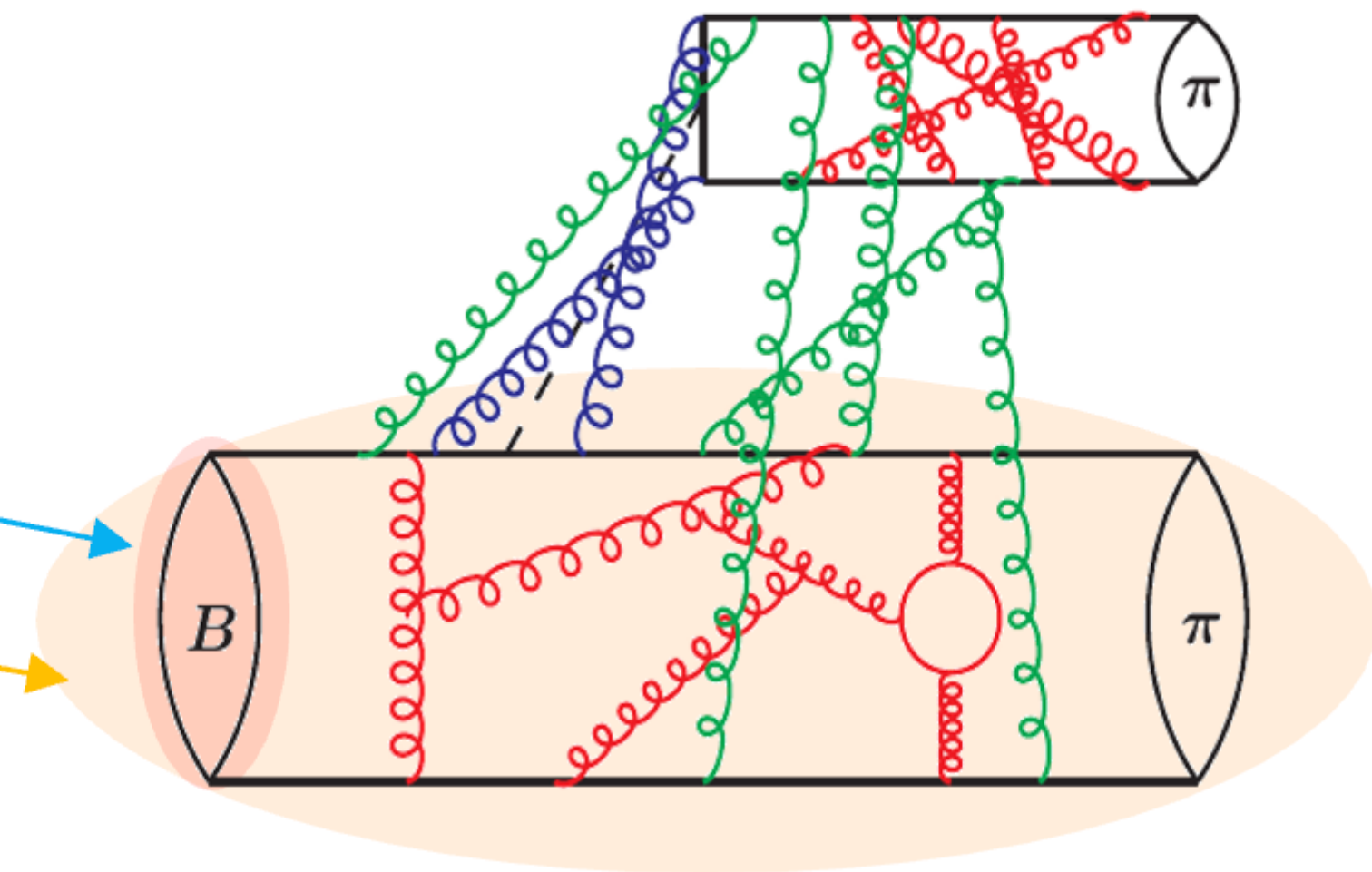
Form factor

Hard kernel $\otimes$ LCDA

Heavy meson LCDA

Hard kernel: Perturbative

Meson LCDAs: Nonperturbative



$$\mathcal{V}_{B \rightarrow K^*}(0) = 0.359^{+0.141}_{-0.085} \Big|_{\lambda_B} \begin{matrix} +0.019 \\ -0.019 \end{matrix} \Big|_{\sigma_1} \begin{matrix} +0.001 \\ -0.062 \end{matrix} \Big|_{\mu} \begin{matrix} +0.010 \\ -0.004 \end{matrix} \Big|_{M^2} \begin{matrix} +0.016 \\ -0.017 \end{matrix} \Big|_{s_0} \begin{matrix} +0.153 \\ -0.079 \end{matrix} \Big|_{\varphi_{\pm}(\omega)},$$

Model dependence

$$f_{B \rightarrow \pi}^0(0) = 0.122 \times \left[1 \pm 0.07 \Big|_{S_0^\pi} \pm 0.11 \Big|_{\Lambda_q} \pm 0.02 \Big|_{\lambda_E^2/\lambda_H^2} \begin{matrix} +0.05 \\ -0.06 \end{matrix} \Big|_{M^2} \pm 0.05 \Big|_{2\lambda_E^2+\lambda_H^2} \begin{matrix} +0.06 \\ -0.10 \end{matrix} \Big|_{\mu_h} \pm 0.04 \Big|_{\mu} \begin{matrix} +1.36 \\ -0.56 \end{matrix} \Big|_{\lambda_B} \begin{matrix} +0.25 \\ -0.43 \end{matrix} \Big|_{\sigma_1, \sigma_2} \right]$$

Inverse moment and log moment

[Gao, Lu, Shen, Wang, Wei, 2020; Cui, Huang, Shen, Wang, 2023]

🤔 Model dependence of heavy meson LCDAs

Models of leading twist heavy meson LCDAs

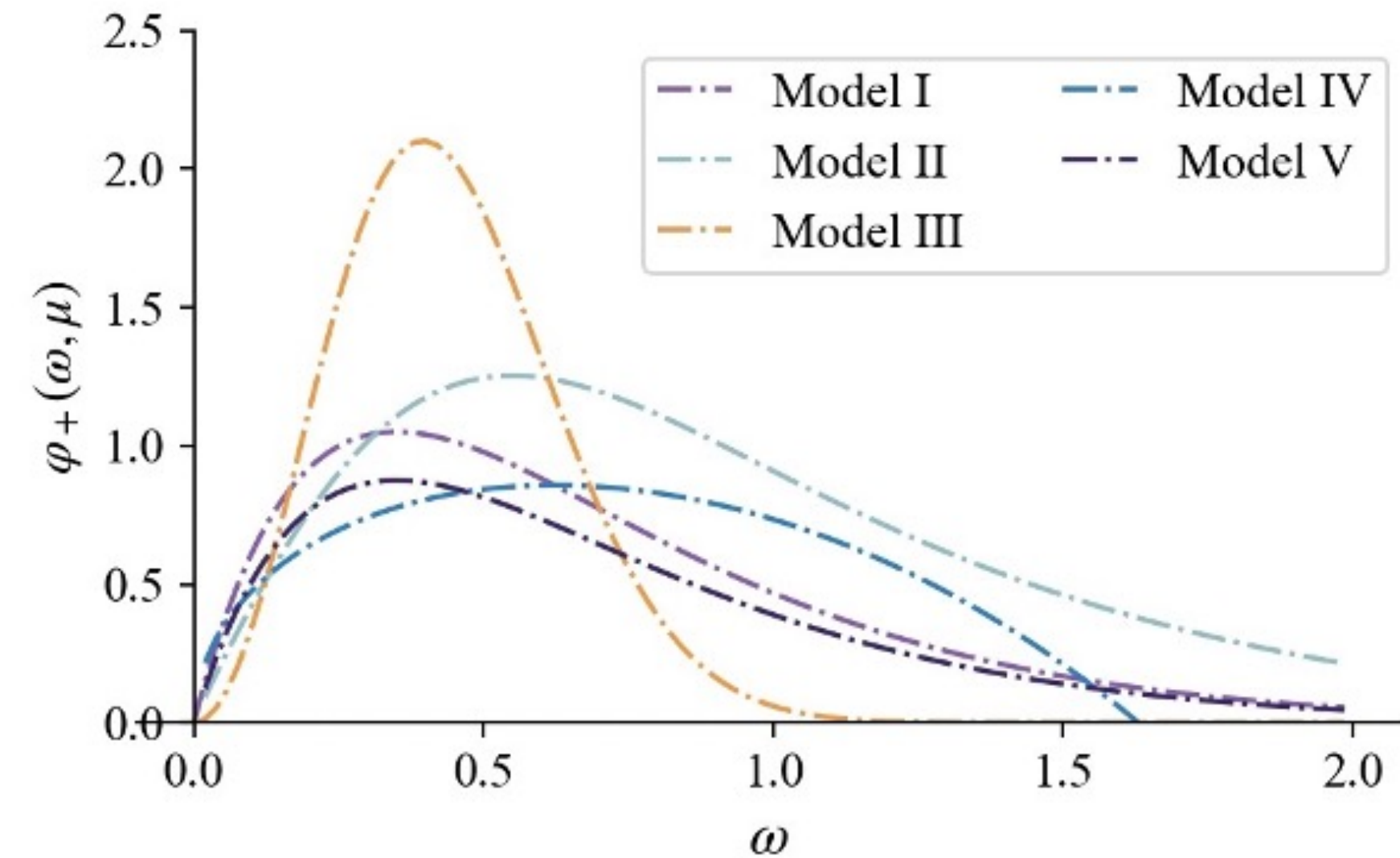
$$\varphi_I^+(\omega, \mu_0) = \frac{\omega}{\omega_0^2} e^{-\omega/\omega_0},$$

$$\varphi_{II}^+(\omega, \mu_0) = \frac{4}{\pi\omega_0} \frac{k}{k^2 + 1} \left[\frac{1}{k^2 + 1} - \frac{2 \left(\sigma_B^{(1)} - 1 \right)}{\pi^2} \ln k \right],$$

$$\varphi_{III}^+(\omega, \mu_0) = \frac{2\omega^2}{\omega_0\omega_1^2} e^{-(\omega/\omega_1)^2},$$

$$\varphi_{IV}^+(\omega, \mu_0) = \frac{\omega}{\omega_0\omega_2} \frac{\omega_2 - \omega}{\sqrt{\omega(2\omega_2 - \omega)}} \theta(\omega_2 - \omega),$$

$$\varphi_V^+(\omega, \mu_0) = \frac{\Gamma(\beta)}{\Gamma(\alpha)} \frac{\omega}{\omega_0^2} e^{-\omega/\omega_0} U(\beta - \alpha, 3 - \alpha, \omega/\omega_0),$$



$$\mathcal{V}_{B \rightarrow K^*}(0) = 0.359_{-0.085}^{+0.141} \Big|_{\lambda_B} \begin{matrix} +0.019 \\ -0.019 \end{matrix} \Big|_{\sigma_1} \begin{matrix} +0.001 \\ -0.062 \end{matrix} \Big|_{\mu} \begin{matrix} +0.010 \\ -0.004 \end{matrix} \Big|_{M^2} \begin{matrix} +0.016 \\ -0.017 \end{matrix} \Big|_{s_0} \begin{matrix} +0.153 \\ -0.079 \end{matrix} \Big|_{\varphi_{\pm}(\omega)}, \quad \text{Model dependence}$$

[Gao, Lu, Shen, Wang, Wei, 2020; Cui, Huang, Shen, Wang, 2023]

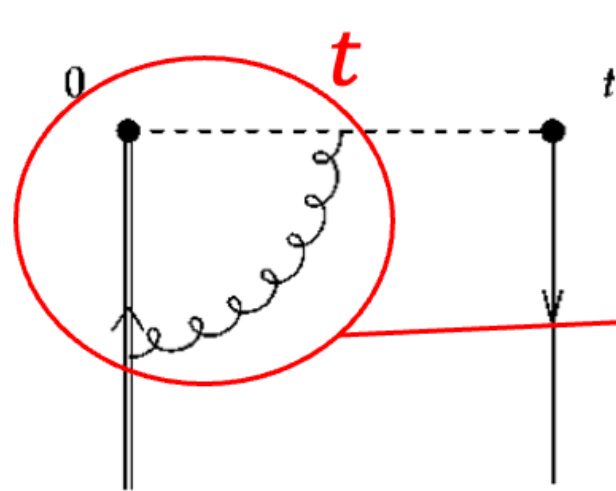


First-principle determination

Leading twist heavy meson LCDAs

$t \equiv z \cdot v$

$\langle 0 | O_v^P(t) | H(p_H) \rangle = i \tilde{f}_H m_H n_+ \cdot v \int_0^\infty d\omega e^{i\omega t n_+ \cdot v} \varphi^+(\omega; \mu), \quad [\text{Braun, Ivanov, Korchemsky, 2004}]$



(a)

$$O_v^{P,\text{ren}}(t, \mu) = O_v^{P,\text{bare}}(t) + \frac{\alpha_s C_F}{4\pi} \left\{ \left(\frac{4}{\hat{\epsilon}^2} + \frac{4}{\hat{\epsilon}} \ln(it\mu) \right) O_v^{P,\text{bare}}(t) - \frac{4}{\hat{\epsilon}} \int_0^1 du \frac{u}{1-u} [O_v^{P,\text{bare}}(ut) - O_v^{P,\text{bare}}(t)] \right\}$$

Diverge at $t \rightarrow 0$!

- Diverge at $t \rightarrow 0 \iff$ **No local limit**

- Non-negative moments $\int dk k^n \varphi_+(k)$ for $n=0,1,2,\dots$ are **not related to OPE**

- Cannot obtain φ_B^+ from lattice QCD through moments.

► How to solve the problem?

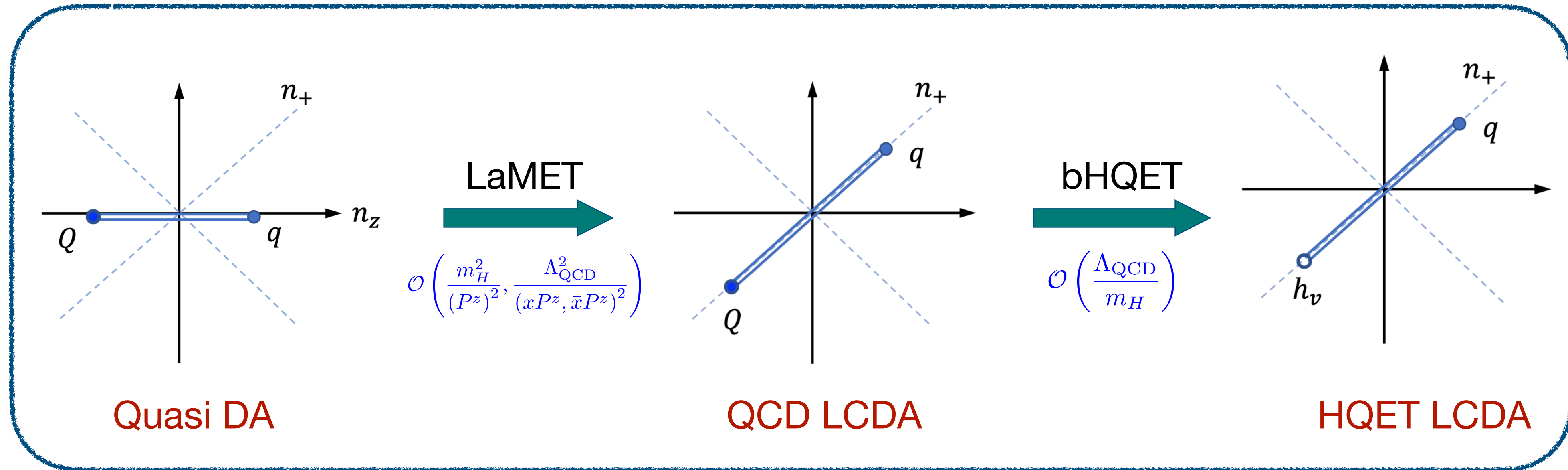
🤔 Off light-cone Wilson line, still heavy quark field h_v ;
[\[Wang², Xu, Zhao, 2020; Xu, Zhang, 2022; Hu, Wang, Xu, Zhao, 2024\]](#)

🤔 No $h_v \Rightarrow$ QCD heavy quark.

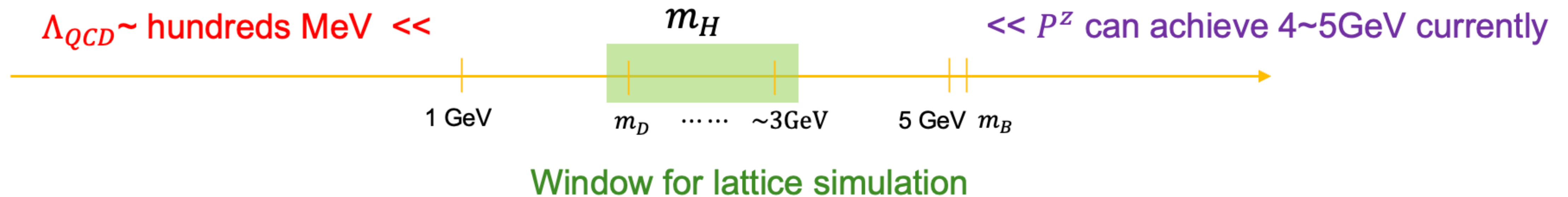
[\[Han, Wang, Zhang, et.al, 2403.17492; Han, Wang, Zhang, Zhang, 2408.13486; Deng, Wang, Wei, Zeng, 2409.00632\]](#)



Sequential effective theory



[Han, Hua, Ji, Lü, et al., 2024; Han, Hua, Ji, Lü, et al., 2025]





Lattice Setup

← Previous work

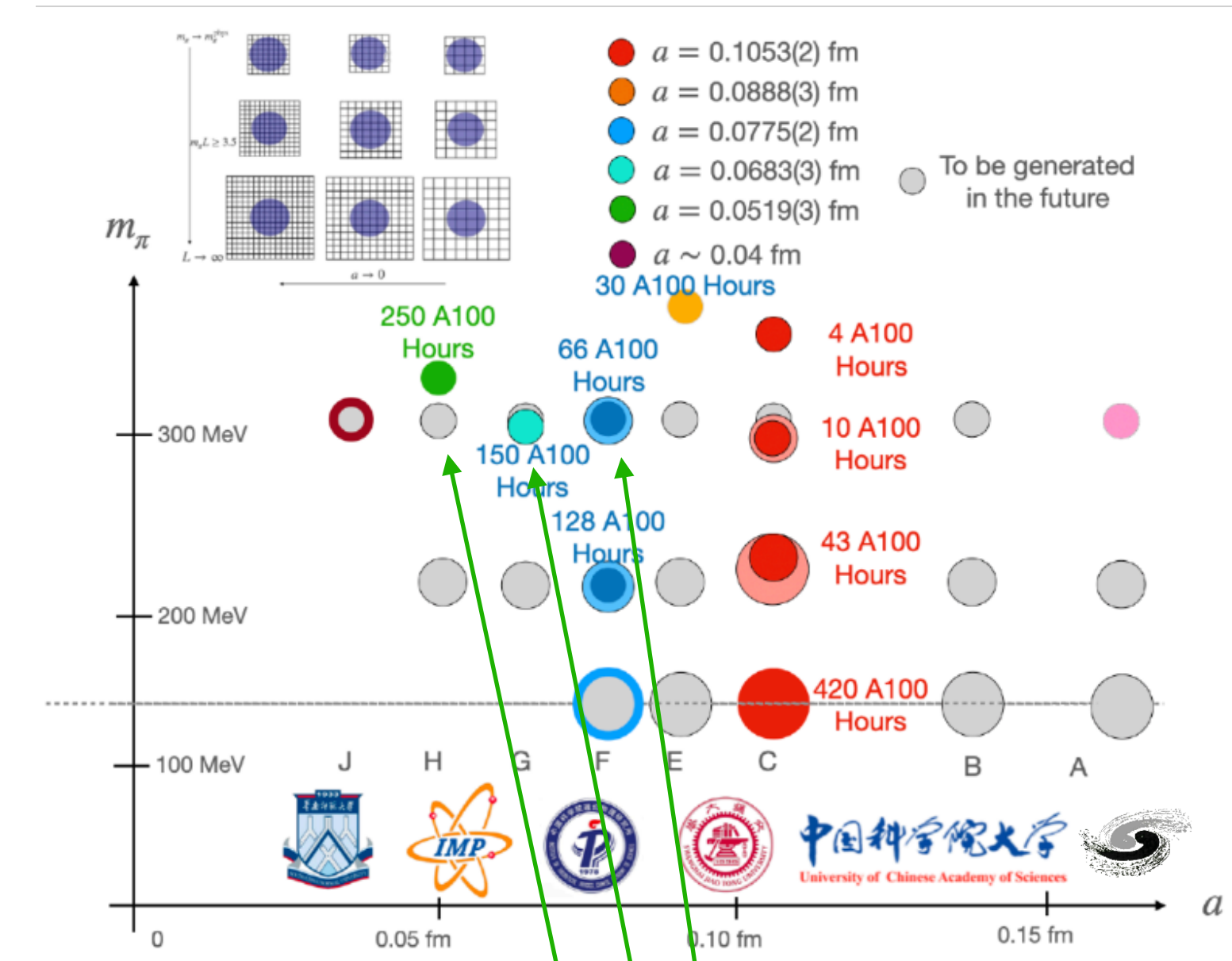
→ Towards precision calculation

Action	Tadpole improved Wilson clover fermion action	
Improvement	—	HYP smear for Wilson line; Multi-source enhancement
Lattice spacing	$a=0.05187\text{fm}$	$a=0.0775\text{fm}, 0.06826\text{fm}, 0.05187\text{fm}$
NPR	Simplified hybrid scheme	Hybrid scheme based on self-renormalization
Pz extrapolation	—	Infinite momentum extrapolation
Statistics	~5k	60k~100k

[arXiv: 2403.17492](https://arxiv.org/abs/2403.17492)

[Phys.Rev.D 111 \(2025\) 3, 034503](#)

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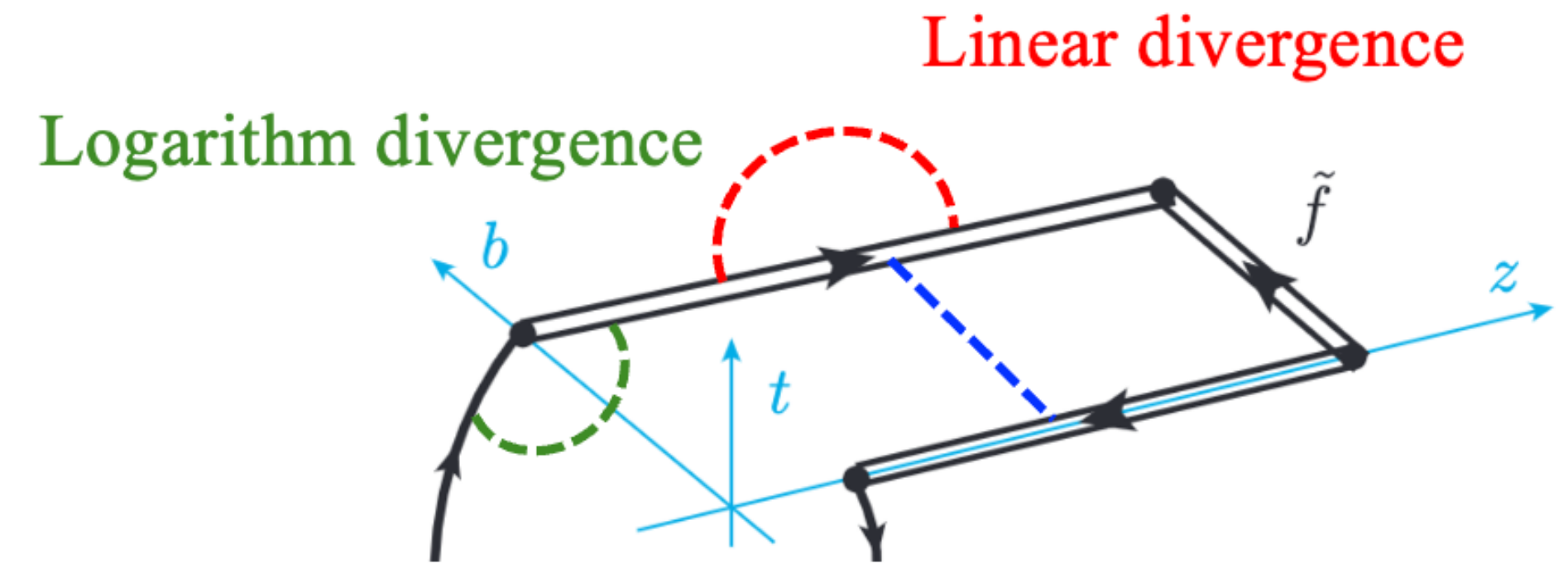


Ensembles we use



NPR of nonlocal operators

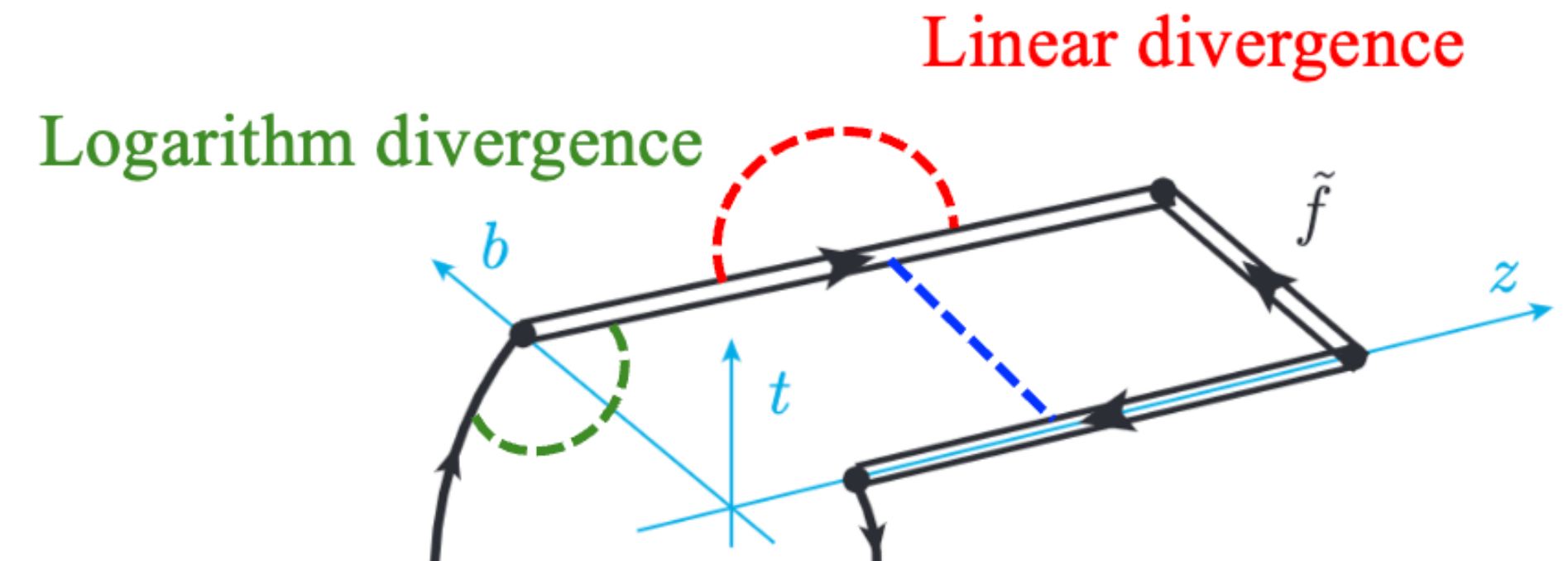
$$\left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_B = \underbrace{e^{\delta m|z|}}_{\text{Linear divergence + ...}} \underbrace{Z(a)}_{\text{Log divergence}} \left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_R$$





NPR of nonlocal operators

$$\left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_B = \underbrace{e^{\delta m|z|}}_{\text{Linear divergence + ...}} \underbrace{Z(a)}_{\text{Log divergence}} \left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_R$$

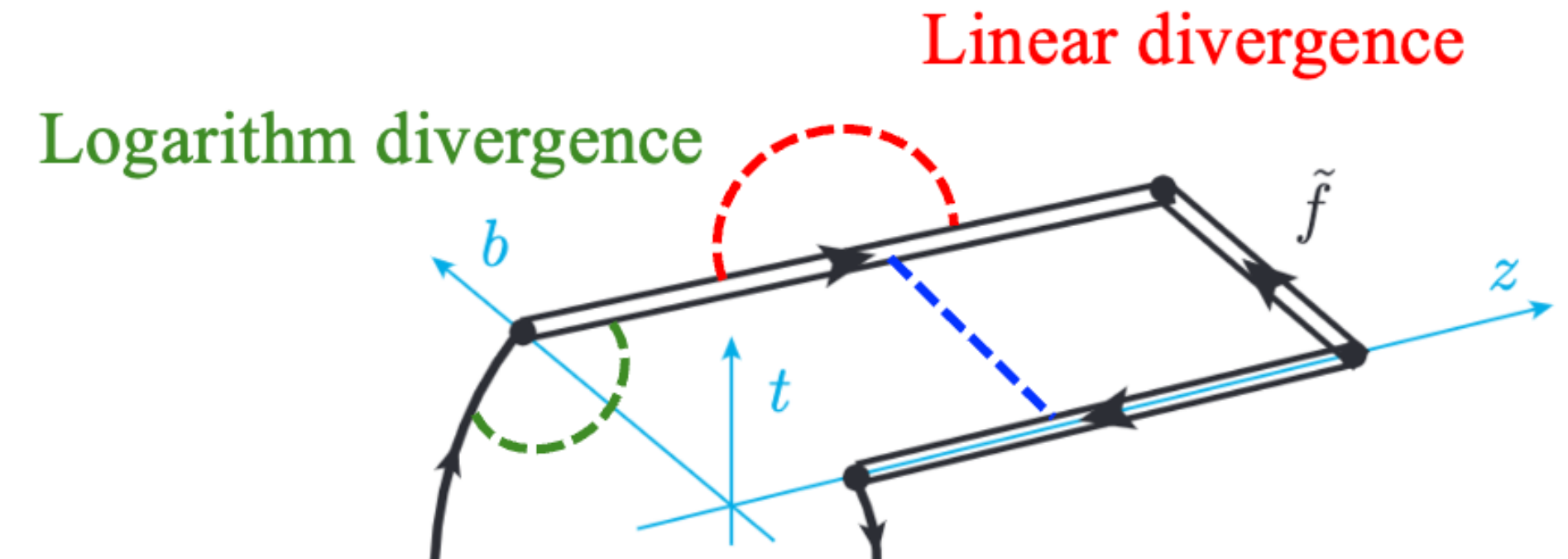


	RI/MOM scheme	Ratio scheme	Self-renormalization
Short distance	Local	Well	—
Large distance	Non-perturbative effects; IR logarithms of z	Invalid OPE	Cancel the linear divergence and discretization effect



NPR of nonlocal operators

$$\left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_B = \underbrace{e^{\delta m|z|}}_{\text{Linear divergence + ...}} \underbrace{Z(a)}_{\text{Log divergence}} \left[\bar{\psi}(z) \Gamma W(z,0) \psi(0) \right]_R$$



$$\langle O_\Gamma(z) \rangle = \Gamma \left(1 + \gamma g^2 \log(z^2/a^2) - m_{-1} \frac{z}{a} + \dots \right)$$

	← v1	← v2	→ Current work
Data	Very limited	Limited	Relatively sufficient
Short distance	Ratio scheme	Ratio scheme	Ratio scheme
Large distance	Ratio scheme	RGR+LRR	Self renormalization

[arXiv: 2403.17492](https://arxiv.org/abs/2403.17492)

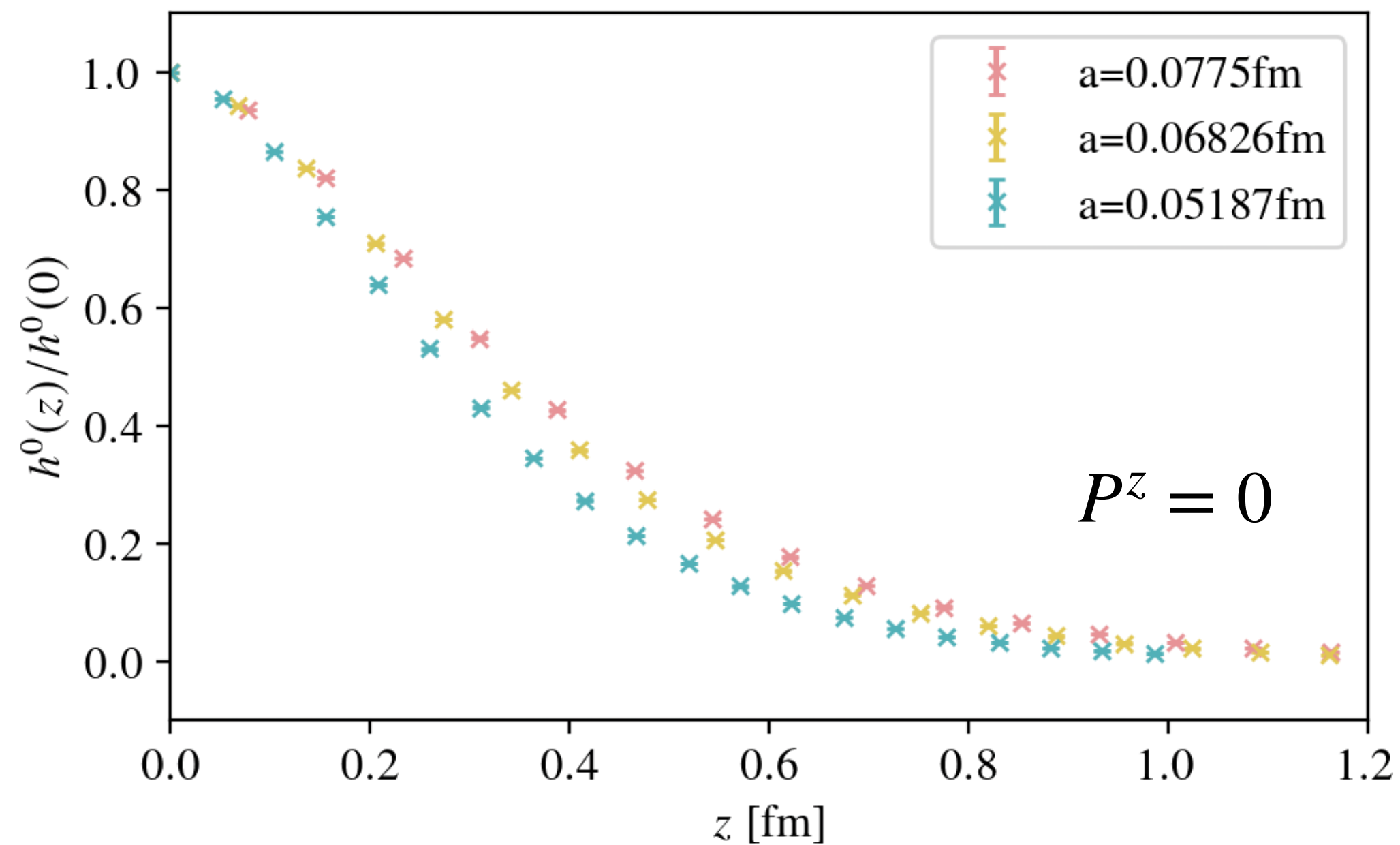
[Phys.Rev.D 111 \(2025\) 3, 034503](https://arxiv.org/abs/2403.17492)



Extraction of renormalization factor

[Hua, Su, Gui, et al. 2021]

Bare matrix element



$$\ln \tilde{h}_B^\pi(z, P_z = 0, 1/a) = \frac{kz}{a \ln[a \Lambda_{\text{QCD}}]} + f(z)a^2 + \frac{3C_F}{b_0} \ln \left[\frac{\ln(1/a \Lambda_{\text{QCD}})}{\ln(\mu/\Lambda_{\text{QCD}})} \right] + \frac{1}{2} \ln \left[1 + \frac{d}{\ln[a \Lambda_{\text{QCD}}]} \right]^2 + \begin{cases} \ln[C_{0,\text{NLO}}(z, \mu)] + m_0 z & \text{if } z_0 \leq z \leq z_1 \\ g(z) & \text{if } z_1 < z \end{cases}, \quad (21)$$

$$C_0^{\text{NLO}}(z, \mu) = 1 + \frac{\alpha_s C_F}{2\pi} \left[\frac{3}{2} \log \left(\frac{1}{4} e^{2\gamma_E} z^2 \mu^2 \right) + A \right]$$

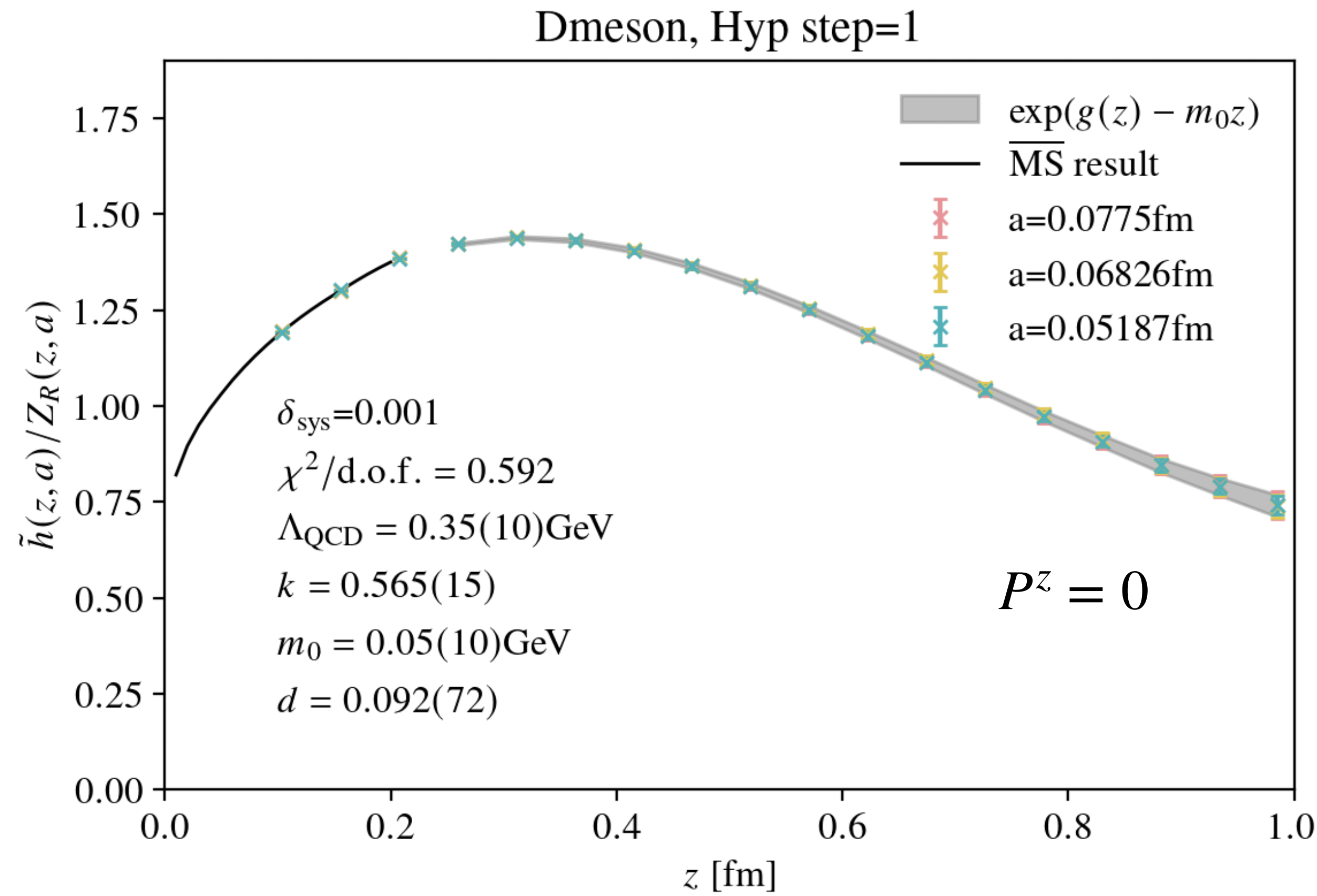
$$C_0^{\text{NLO+LRR}}(z, \mu) = C_0^{\text{NLO}}(z, \mu) + z\mu \left(C_{\text{PV}}(z, \mu) - \sum_{i=0}^{k-1} \alpha_s^{i+1}(\mu) r_i \right)$$

- Hypercubic smearing improved the SNR

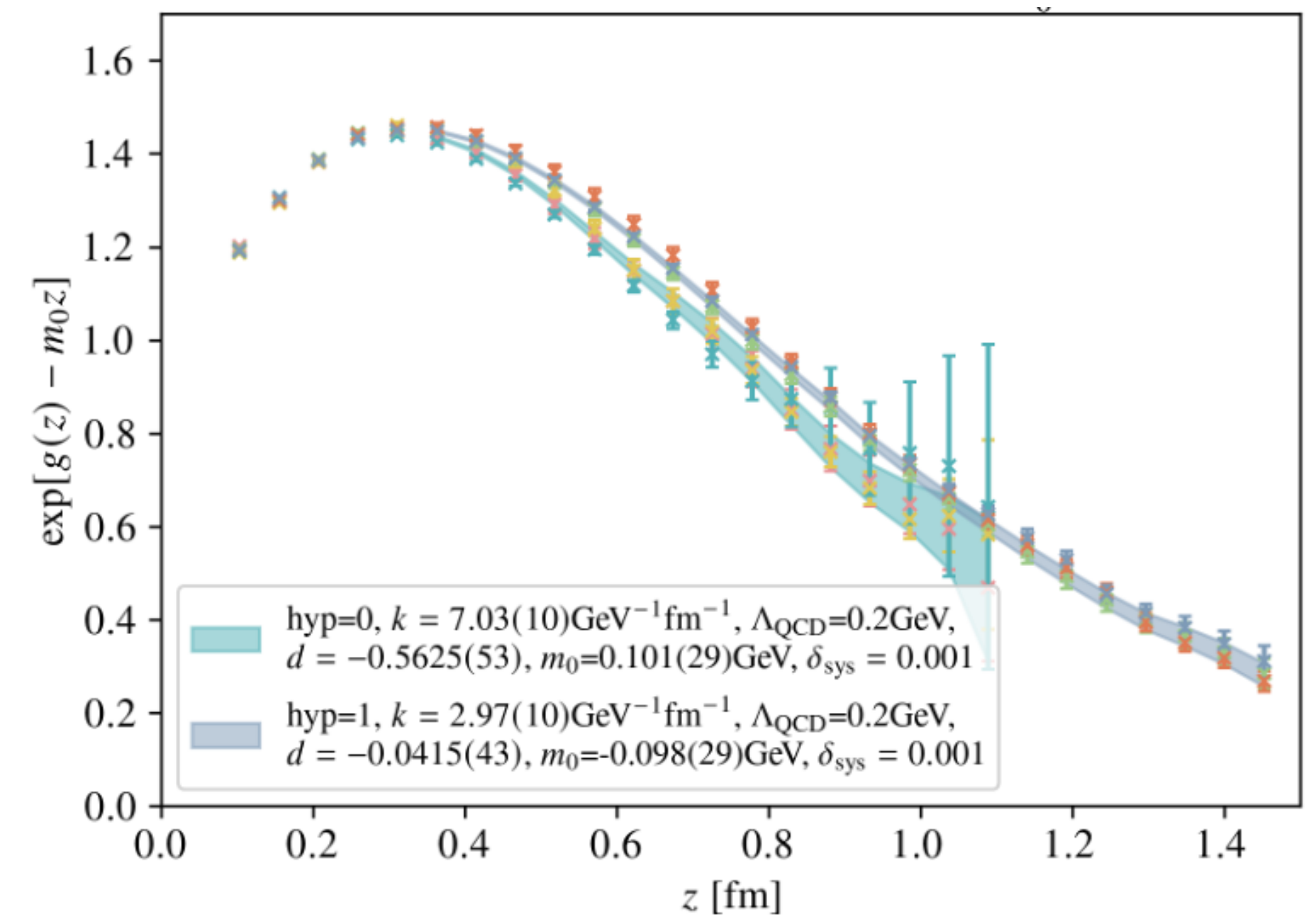


Extraction of renormalization factor

Renormalized matrix element



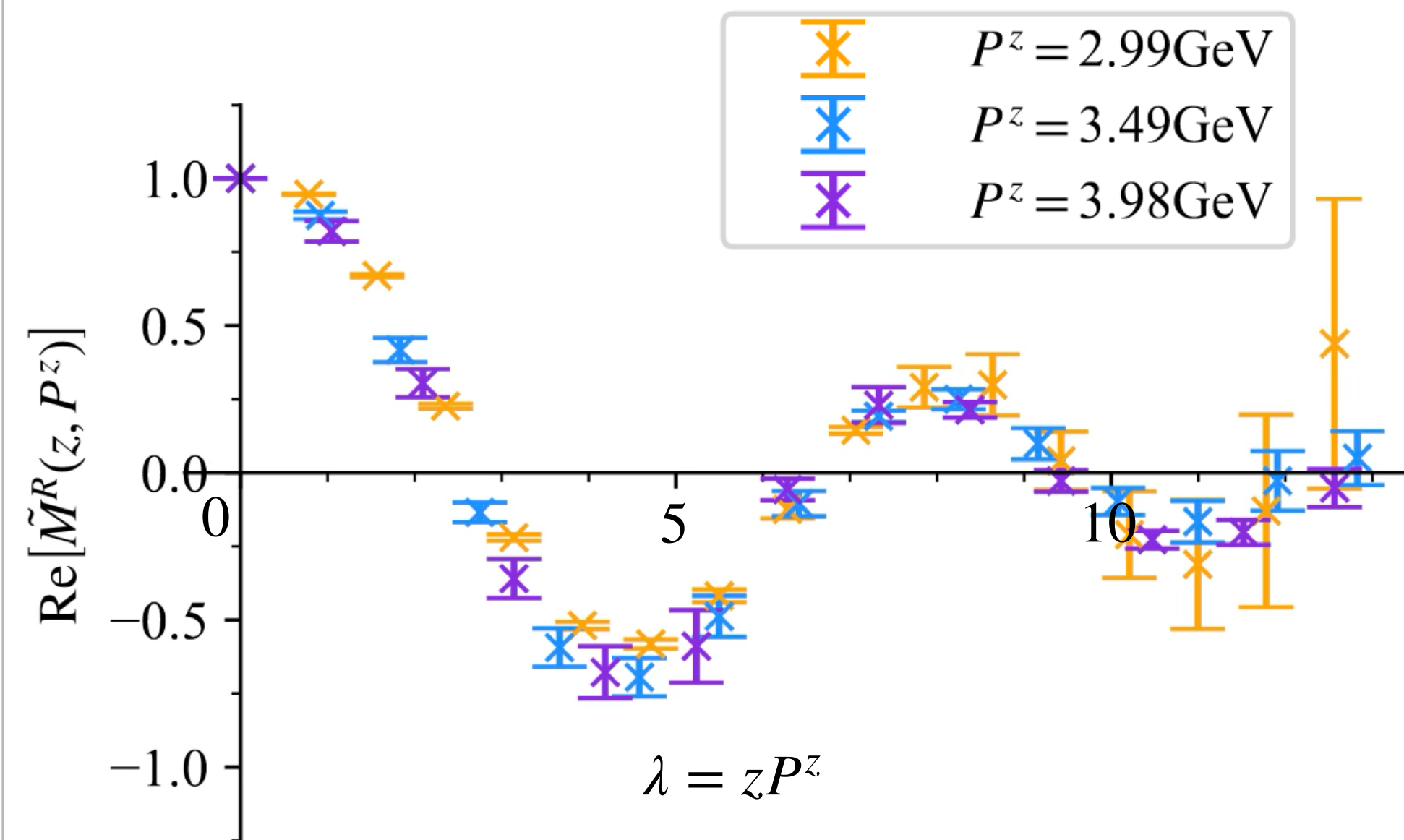
- Hypercubic smearing has only a minor impact on renormalization.



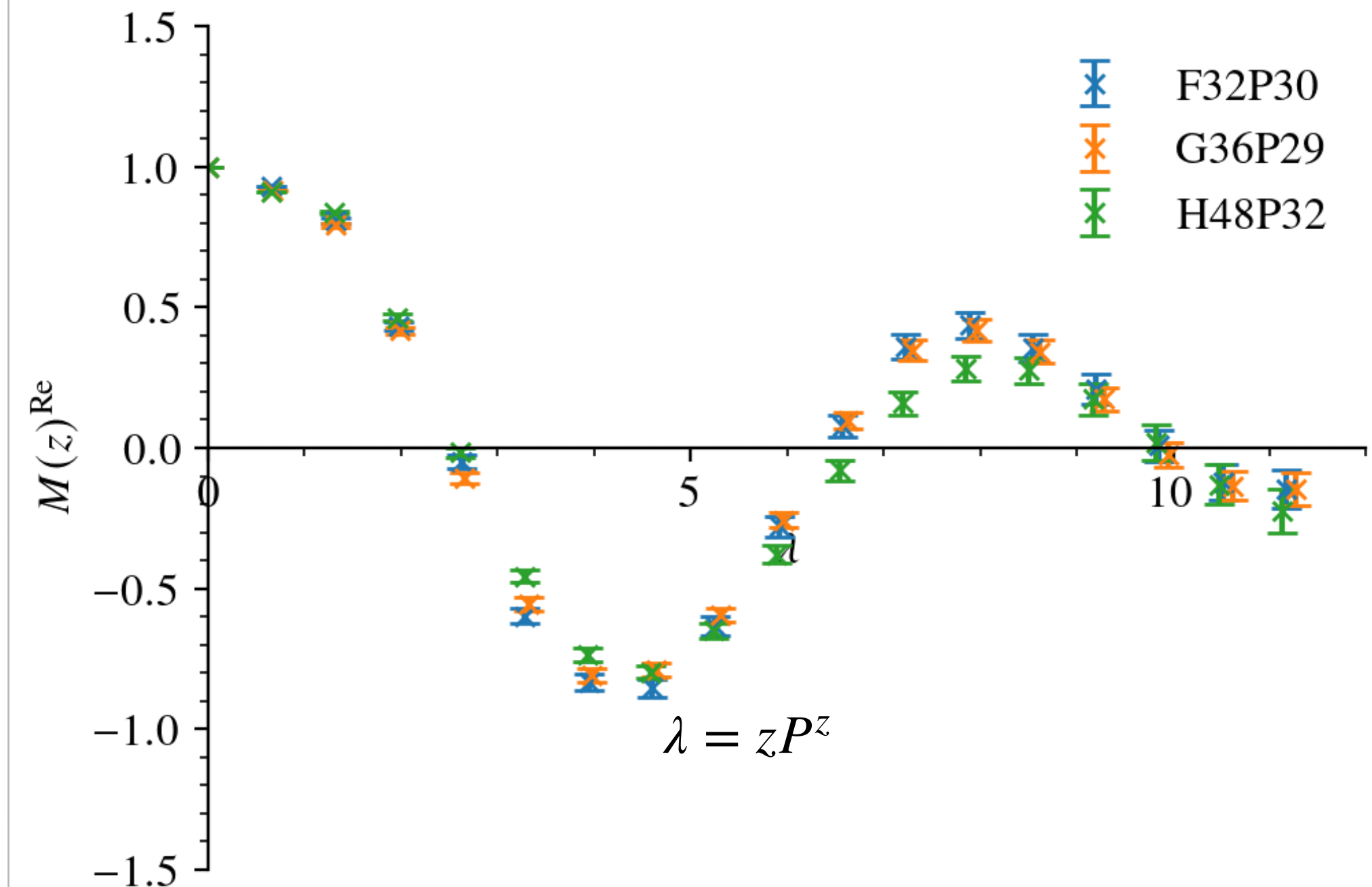


Renormalized matrix elements

← Previous work



→ Towards precision calculation



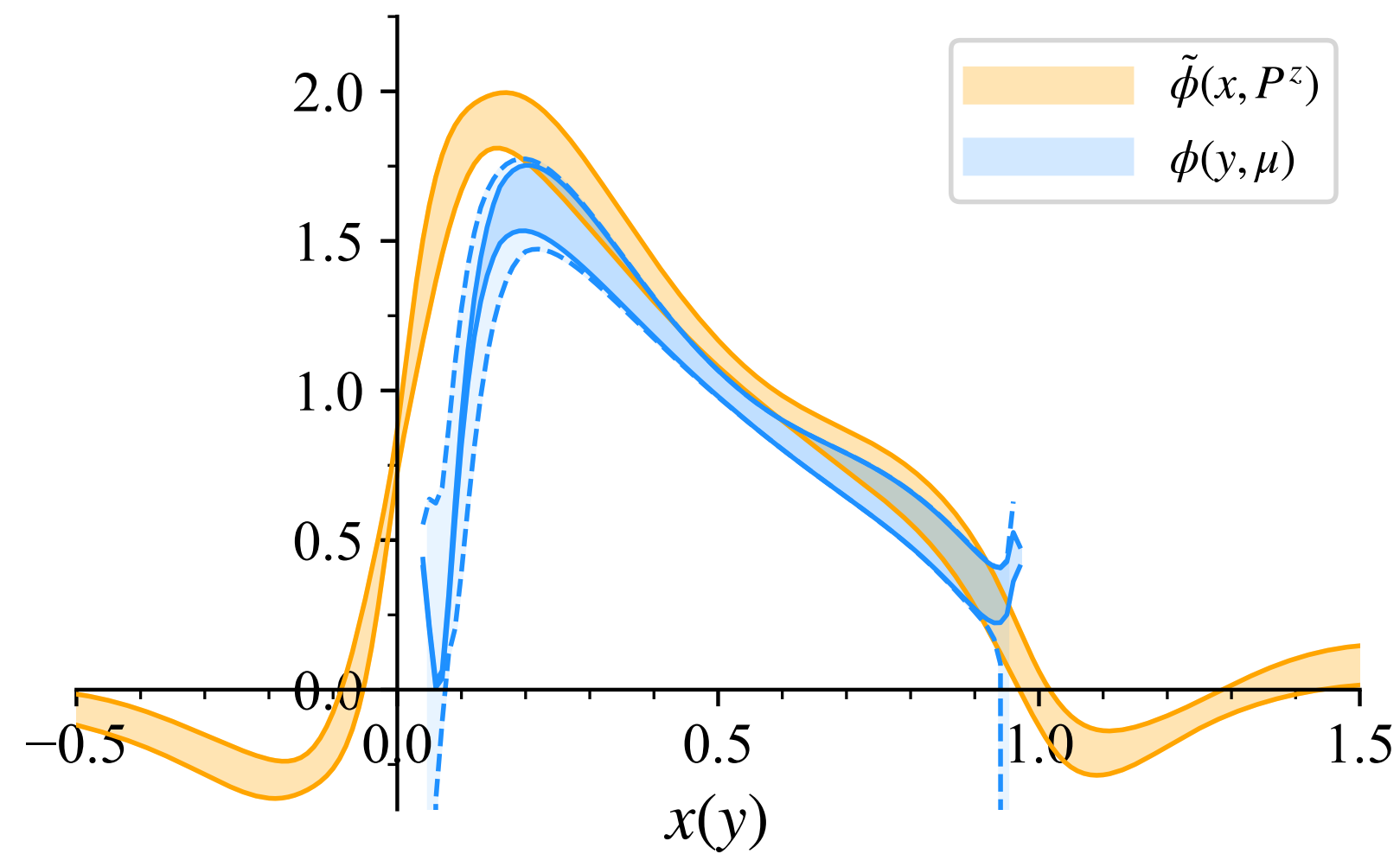
- Obtain a more precise results at large λ
- Multi lattice spacings to perform the continuum extrapolation

Matching for previous results

LaMET matching

$$\tilde{\phi}(x, P^z) = \int_0^1 C\left(x, y, \frac{\mu}{P^z}\right) \phi(y, \mu) + \mathcal{O}\left(\frac{m_H^2}{(P^z)^2}, \frac{\Lambda_{\text{QCD}}^2}{(xP^z, \bar{x}P^z)^2}\right)$$

[Liu, Wang, Xu, Zhang, Zhao, 2019]

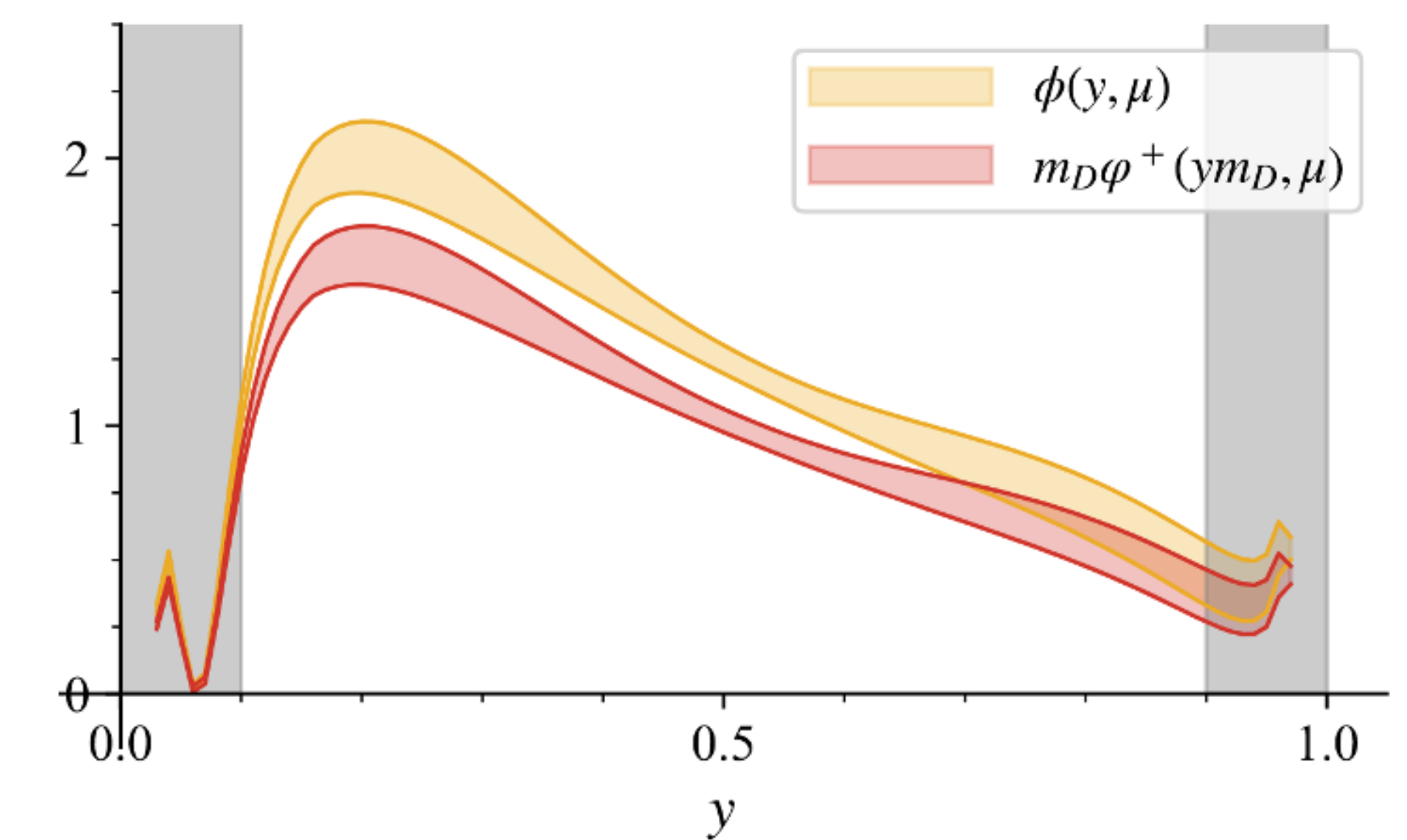


Quasi DA & QCD LCDA

bHQET matching

$$\varphi_{\text{peak}}^+(\omega, m_H) = \frac{f_H}{\tilde{f}_H} \frac{1}{\mathcal{J}_{\text{peak}}} \phi(y, \mu; m_H) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{m_H}\right)$$

[Beneke, Finauri, Vos, Wei, 2023; Ishaq, Jia, et al. 2019]



QCD LCDA & HQET LCDA ($\omega \sim \Lambda_{\text{QCD}}$)

[Han, Hua, Ji, Lü, et al., 2024; Han, Hua, Ji, Lü, et al., 2025]

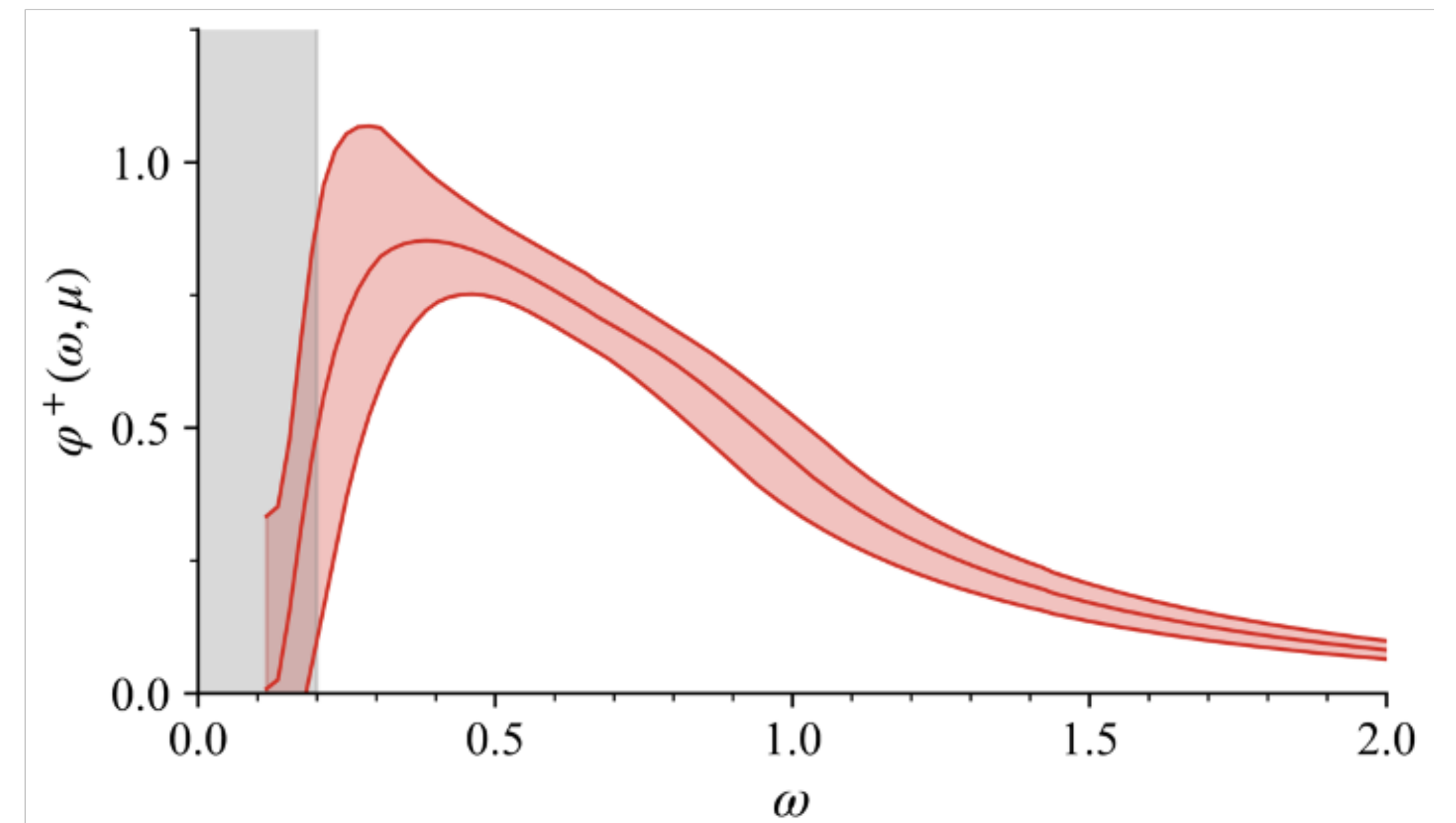
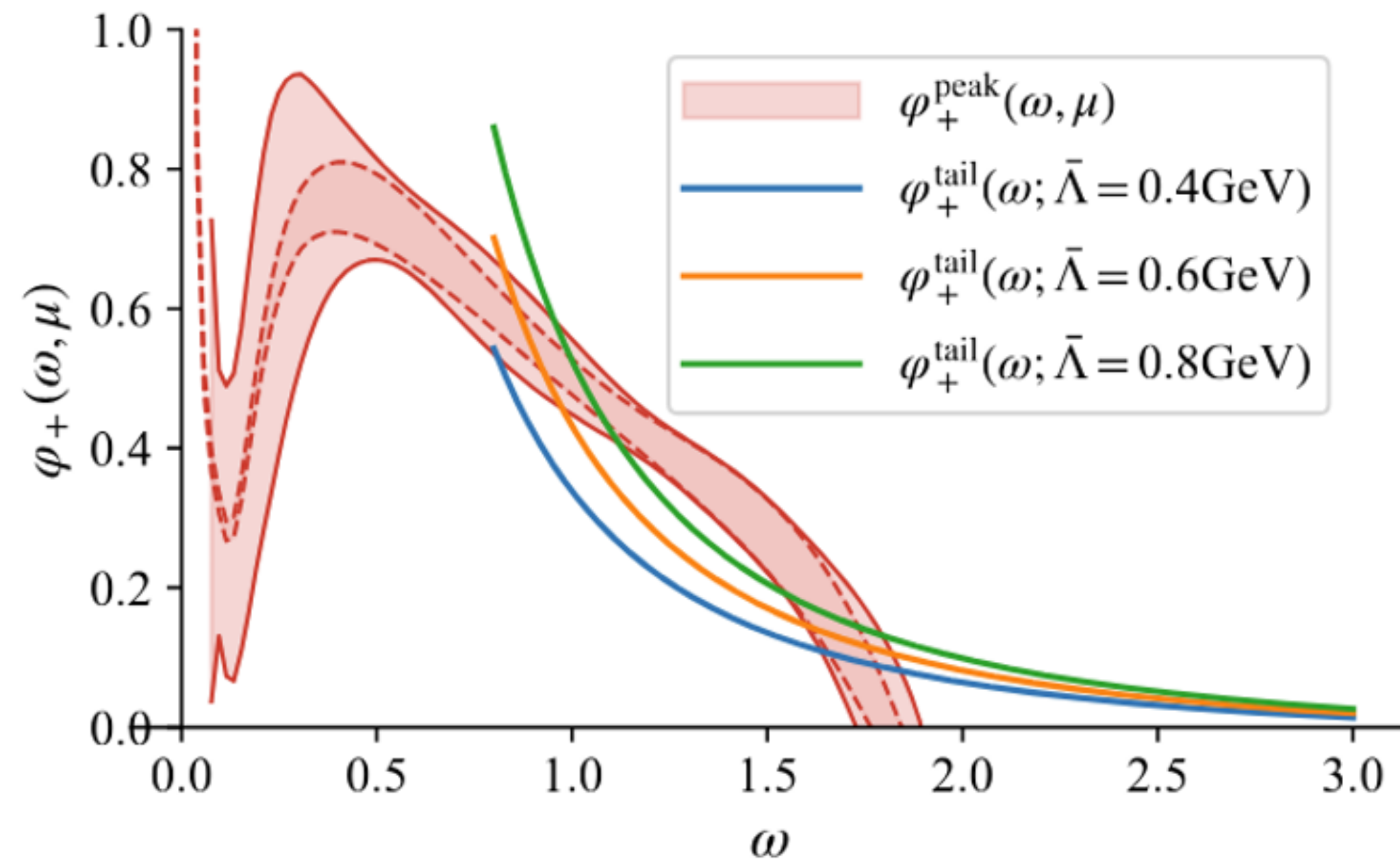
Matching for previous results

- Combine with $\omega \gg \Lambda_{\text{QCD}}$ region

$$\varphi_{\text{tail}}^+ = \frac{\alpha_s C_F}{\pi \omega} \left[\left(\frac{1}{2} - \ln \frac{\omega}{\mu} \right) + \frac{4\bar{\Lambda}}{3\omega} \left(2 - \ln \frac{\omega}{\mu} \right) \right]$$

[Lee, Neubert, 2005]

$$\varphi^+(\omega, \mu) = \varphi_{\text{peak}}^+(\omega, \mu) \theta(\omega_b - \omega) + \varphi_{\text{tail}}^+(\omega, \mu) \theta(\omega - \omega_b)$$



HQET LCDA

Inverse moments

- λ_B and $\sigma_B^{(1)}$ at $\mu = 1\text{GeV}$

		λ_B (GeV)	$\sigma_B^{(1)}$
Our results	$N=1$	0.389(35)	1.63(8)
	$N=2$	0.393(37)	1.62(7)
	$N=3$	0.381(59)	1.63(12)
Experiment	<i>Belle 2018</i>	> 0.24	
Other theoretical approach	<i>Khodjamirian, Mandal, Mannel, 2020</i>	0.383(153)	
	<i>Gao, Lu, Shen, Wang, Wei, 2020</i>	$0.343^{+0.064}_{-0.079}$	
	<i>Lee, Neubert, 2005</i>	0.48(11)	1.6(2)
	<i>Braun, Ivanov, Korchemsky, 2004</i>	0.46(11)	1.4(4)
	<i>Grozin, Neubert, 1997</i>	0.35(15)	
	<i>Mandal, Nandi, Ray, 2024</i>	0.338(68)	



Summary and prospect

✓ What we achieved:

- First-principle strategy for heavy meson LCDA
- Hybrid renormalization:
 - Short-distance: ratio scheme (robust and clean)
 - Long-distance: self-renormalization (non-perturbative, removes linear divergence & discretization effects)
- Better StN ratio at large z via smearing & statistics

🌱 What's next:

- Continuum extrapolation & large P^z extrapolation;
- A more reliable extraction of HQET LCDA
- Extract $\lambda_B, \sigma_B^{(1)}$
- Apply to higher-twist & B meson.

Thank you!