

用SU(3)味对称性/破缺方法研究

$D \rightarrow P/V/S/A \ell^+ \nu_\ell$ 半轻衰变

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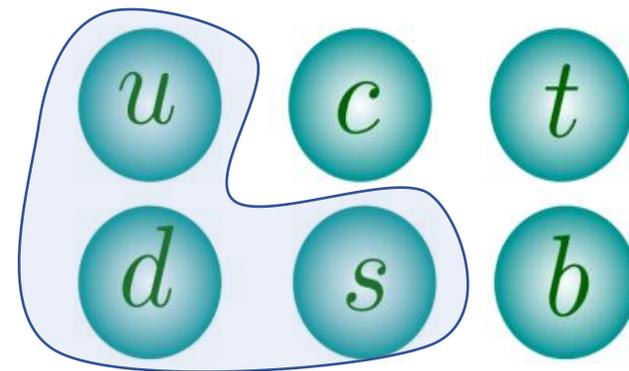
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研究背景

- 一些 $B(D \rightarrow P/V \ell \nu)$ 已经很好的测量，而 $B(D \rightarrow S/A \ell \nu)$ 还没有很好的测量到；
- 在BESIII, LHCb和BelleII上产生很多粲强子的实验数据；
- 计算 $D \rightarrow P/V \ell^+ \nu_\ell$ 与 $D \rightarrow S/A \ell^+ \nu_\ell$ 的形状因子；
- 对称性方法可以为粒子物理提供一些非常重要的信息。

SU(3) 味对称方法

- 不可约表示方法（哈密顿量，动力学细节）
- 拓扑图方法（夸克线，关联振幅）



优点

不依赖于动力学细节（不需要计算形状因子）

缺点

不能决定自身振幅的大小

用SU(3) 味对称方法研究介子衰变

decays \ baryons	Charmed mesons	Bottom mesons
$B/D \rightarrow M_1 M_2$	<ul style="list-style-type: none"> ① Hai -Yang Cheng et. al., PRD 86, 014014(2012); PRD 93, 114010(2016) ② Xiao-Gang He, Wei Wang et. al., EPJC 80, 359(2020) ; CPC 42, 103108(2018). ③ D. Pirtskhalava, PLB 742, 81 (2012). ④ 	<ul style="list-style-type: none"> ① Si -Hong Zhou et. al., PRD 92094016(2015); EPJC 77(2017) 2, 125; PRD 92, 094016(2015). ② Hai -Yang Cheng et. al., JHEP 09, 024(2011); PRD 91, 014011(2015). ③ Xiao-Gang He, Wei Wang et. al., EPJC 80, 359(2020) ; CPC 42, 103108(2018); PRD 93, 114002(2016); JHEP 08, 065(2013); PRD 69, 074002(2004); PRD 64, 034002(2001); EPJC 9, 443(1999); PRL 75, 1703(1995). ④
$B/D \rightarrow M l \nu$ $B/D \rightarrow M_1 M_2 l \nu$	$D \rightarrow P/V/S l \nu$ NPB995, 116349(2023) $D \rightarrow P P l \nu$ PRD107, 056022(2023) $D \rightarrow A(A \rightarrow VP) l \nu$ arXiv: 2404. 03857	$B \rightarrow V P l \nu$ working



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介子多重态

介子 $q\bar{q}$

$$D_i = (D^0(c\bar{u}), D^+(c\bar{d}), D_s^+(c\bar{s}))$$

$$P = \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} + \frac{\eta_1}{\sqrt{3}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} + \frac{\eta_1}{\sqrt{3}} & K^0 \\ K^- & \bar{K}^0 & -\frac{2\eta_8}{\sqrt{6}} + \frac{\eta_1}{\sqrt{3}} \end{pmatrix}$$

$$V = \begin{pmatrix} \frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} + \frac{\omega_1}{\sqrt{3}} & \rho^+ & K^{*+} \\ \rho^- & -\frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} + \frac{\omega_1}{\sqrt{3}} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & -\frac{2\omega_8}{\sqrt{6}} + \frac{\omega_1}{\sqrt{3}} \end{pmatrix}$$

$$S = \begin{pmatrix} \frac{a_0^0}{\sqrt{2}} + \frac{\sigma}{\sqrt{2}} & a_0^+ & K_0^+ \\ a_0^- & -\frac{a_0^0}{\sqrt{2}} + \frac{\sigma}{\sqrt{2}} & K_0^0 \\ K_0^- & \bar{K}_0^0 & f_0 \end{pmatrix}$$

$$\begin{aligned} \sigma &= u\bar{u}d\bar{d}, & f_0 &= (u\bar{u} + d\bar{d})s\bar{s}/\sqrt{2}, \\ a_0^0 &= (u\bar{u} - d\bar{d})s\bar{s}/\sqrt{2}, & a_0^+ &= u\bar{d}s\bar{s}, & a_0^- &= d\bar{u}s\bar{s}, \\ K_0^+ &= u\bar{s}d\bar{d}, & K_0^0 &= d\bar{s}u\bar{u}, & \bar{K}_0^0 &= s\bar{d}u\bar{u}, & K_0^- &= s\bar{u}d\bar{d}, \end{aligned}$$

S_{jm}^{im}

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos\theta_P & -\sin\theta_P \\ \sin\theta_P & \cos\theta_P \end{pmatrix} \begin{pmatrix} \eta_8 \\ \eta_1 \end{pmatrix} \quad \begin{pmatrix} f_0(980) \\ f_0(500) \end{pmatrix} = \begin{pmatrix} \cos\theta_S & \sin\theta_S \\ -\sin\theta_S & \cos\theta_S \end{pmatrix} \begin{pmatrix} f_0 \\ \sigma \end{pmatrix}$$

$M_j^i \quad i, j = 1, 2, 3 \text{ for } u, d, s.$

$$A = \begin{pmatrix} \frac{a_1^0}{\sqrt{2}} + \frac{f_1}{\sqrt{3}} + \frac{f_8}{\sqrt{6}} & a_1^+ & K_{1A}^+ \\ a_1^- & -\frac{a_1^0}{\sqrt{2}} + \frac{f_1}{\sqrt{3}} + \frac{f_8}{\sqrt{6}} & K_{1A}^0 \\ K_{1A}^- & \bar{K}_{1A}^0 & \frac{f_1}{\sqrt{3}} - \frac{2f_8}{\sqrt{6}} \end{pmatrix}$$

$J^{PC} = 1^{++}$

$$B = \begin{pmatrix} \frac{b_1^0}{\sqrt{2}} + \frac{h_1}{\sqrt{3}} + \frac{h_8}{\sqrt{6}} & b_1^+ & K_{1B}^+ \\ b_1^- & -\frac{b_1^0}{\sqrt{2}} + \frac{h_1}{\sqrt{3}} + \frac{h_8}{\sqrt{6}} & K_{1B}^0 \\ K_{1B}^- & \bar{K}_{1B}^0 & \frac{h_1}{\sqrt{3}} - \frac{2h_8}{\sqrt{6}} \end{pmatrix}$$

$J^{PC} = 1^{+-}$

$$\begin{pmatrix} K_1(1270) \\ K_1(1400) \end{pmatrix} = \begin{pmatrix} \sin\theta_{K_1} & \cos\theta_{K_1} \\ \cos\theta_{K_1} & -\sin\theta_{K_1} \end{pmatrix} \begin{pmatrix} K_{1A} \\ K_{1B} \end{pmatrix}$$

$$\begin{pmatrix} f_1(1285) \\ f_1(1420) \end{pmatrix} = \begin{pmatrix} \cos\theta_{3P1} & \sin\theta_{3P1} \\ -\sin\theta_{3P1} & \cos\theta_{3P1} \end{pmatrix} \begin{pmatrix} f_1 \\ f_8 \end{pmatrix}$$

$$\begin{pmatrix} h_1(1170) \\ h_1(1415) \end{pmatrix} = \begin{pmatrix} \cos\theta_{1P1} & \sin\theta_{1P1} \\ -\sin\theta_{1P1} & \cos\theta_{1P1} \end{pmatrix} \begin{pmatrix} h_1 \\ h_8 \end{pmatrix}$$

衰变分支比

$$\mathcal{H}_{eff}(c \rightarrow q\ell^+\nu_\ell) = \frac{G_F}{\sqrt{2}} V_{cq} \bar{q} \gamma^\mu (1 - \gamma_5) c \bar{\nu}_\ell \gamma_\mu (1 - \gamma_5) \ell,$$

$$\mathcal{M}(D \rightarrow M\ell^+\nu_\ell) = \frac{G_F}{\sqrt{2}} \sum_{mm'} g_{mm'} L_m^{\lambda_\ell \lambda_\nu} H_{m'}^{\lambda_M},$$

$$L_m^{\lambda_\ell \lambda_\nu} = \epsilon_\alpha(m) \bar{\nu}_\ell \gamma^\alpha (1 - \gamma_5) \ell,$$

$$H_{m'}^{\lambda_M} = \begin{cases} V_{cq}^* \epsilon_\beta^*(m') \langle P/S(p_{P/S}) | \bar{q} \gamma^\beta (1 - \gamma_5) c | D(p_D) \rangle \\ V_{cq}^* \epsilon_\beta^*(m') \langle V/A(p_{V/A}, \epsilon^*) | \bar{q} \gamma^\beta (1 - \gamma_5) c | D(p_D) \rangle \end{cases},$$

$$\frac{d\mathcal{B}(D_{(s)} \rightarrow M\ell^+\nu_\ell)}{dq^2} = \frac{\tau_D G_F^2 |V_{cq}|^2 \lambda^{1/2} (q^2 - m_\ell^2)^2}{24(2\pi)^3 M_{D_{(s)}}^3 q^2} \mathcal{H}_{total},$$

$$\mathcal{H}_{total} \equiv (\mathcal{H}_U + \mathcal{H}_L) \left(1 + \frac{m_\ell^2}{2q^2} \right) + \frac{3m_\ell^2}{2q^2} \mathcal{H}_S.$$

D \rightarrow M $l^+ \nu_l$ 衰变强子螺旋度振幅

$$\blacktriangleright H(D \rightarrow M l^+ \nu_l) = c_0^M D_i M^i_j H^j$$

$$\blacktriangleright \Delta H(D \rightarrow M l^+ \nu_l) = c_1^M D_a W^a_i M^i_j H^j + c_2^M D_i M^i_a W^a_j H^j$$

其中, c_i^M 是非微扰系数, H^j 是CKM 矩阵元。

$$\text{SU}(3) \text{ 味破缺系数矩阵 } W = (W_j^i) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

D \rightarrow P/V $\ell^+ \nu_\ell$ 形状因子

C₁: 所有的形状因子都被视为常数, 没有对强子动量转移 q^2 的依赖性

C₂: 具有SU(3)味对称, $F_i(q^2)$ 对的 q^2 依赖性的极点修正模型

$$F_i(q^2) = \frac{F_i(0)}{\left(1 - \frac{q^2}{m_{pole}^2}\right) \left(1 - \alpha_i \frac{q^4}{m_{pole}^4}\right)},$$

C₃: 具有SU(3)味对称 $F_i(q^2) = \frac{F_i(0)}{\left(1 - \frac{q^2}{m_{pole}^2}\right) \left(1 - \sigma_{1i} \frac{q^2}{m_{pole}^2} + \sigma_{2i} \frac{q^4}{m_{pole}^4}\right)}$ ($f_+^P(q^2)$ 和 $V(q^2)$),

$$F_i(q^2) = \frac{F_i(0)}{\left(1 - \sigma_{1i} \frac{q^2}{m_{pole}^2} + \sigma_{2i} \frac{q^4}{m_{pole}^4}\right)} \quad (f_0^P(q^2), A_1(q^2) \text{ 和 } A_2(q^2)),$$

C₄: 考虑SU(3)味破缺 c_{12}^M 的形状因子



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D \rightarrow P $\ell^+\nu_\ell$ 衰变振幅

TABLE I: The hadronic helicity amplitudes for the $D \rightarrow P\ell^+\nu$ decays including both the SU(3) flavor symmetry and the SU(3) flavor breaking contributions. $A_1 \equiv c_0^P + c_1^P - 2c_2^P$, $A_2 \equiv c_0^P - 2c_1^P - 2c_2^P$, $A_3 \equiv c_0^P + c_1^P + c_2^P$, $A_4 \equiv c_0^P - 2c_1^P + c_2^P$. $A_1 = A_2 = A_3 = A_4 = c_0^P$ if neglecting the SU(3) flavor breaking c_1^P and c_2^P terms.

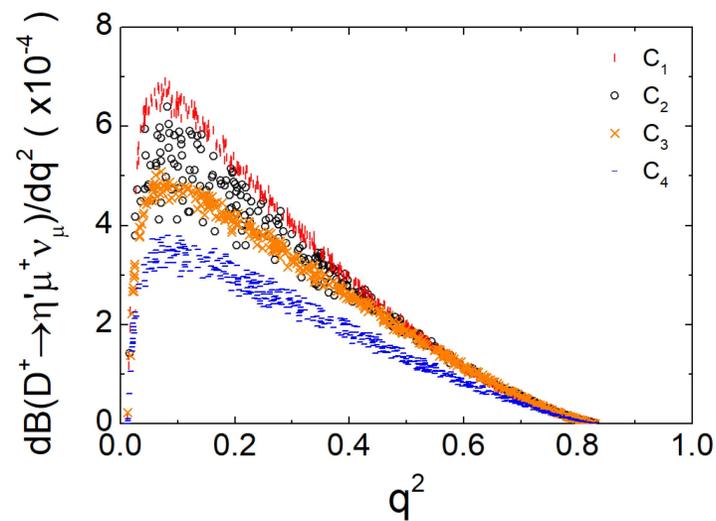
Hadronic helicity amplitudes	SU(3) flavor amplitudes
$H(D^0 \rightarrow K^-\ell^+\nu_\ell)$	$A_1V_{cs}^*$
$H(D^+ \rightarrow \bar{K}^0\ell^+\nu_\ell)$	$A_1V_{cs}^*$
$H(D_s^+ \rightarrow \eta\ell^+\nu_\ell)$	$(-\cos\theta_P\sqrt{2/3} - \sin\theta_P/\sqrt{3})A_2V_{cs}^*$
$H(D_s^+ \rightarrow \eta'\ell^+\nu_\ell)$	$(-\sin\theta_P\sqrt{2/3} + \cos\theta_P/\sqrt{3})A_2V_{cs}^*$
$H(D_s^+ \rightarrow \pi^0\ell^+\nu_\ell)$	$-\delta(-\cos\theta_P\sqrt{2/3} - \sin\theta_P/\sqrt{3})A_2V_{cs}^*$
$H(D^0 \rightarrow \pi^-\ell^+\nu_\ell)$	$A_3V_{cd}^*$
$H(D^+ \rightarrow \pi^0\ell^+\nu_\ell)$	$-\frac{1}{\sqrt{2}}A_3V_{cd}^*$
$H(D^+ \rightarrow \eta\ell^+\nu_\ell)$	$(\cos\theta_P/\sqrt{6} - \sin\theta_P/\sqrt{3})A_3V_{cd}^*$
$H(D^+ \rightarrow \eta'\ell^+\nu_\ell)$	$(\sin\theta_P/\sqrt{6} + \cos\theta_P/\sqrt{3})A_3V_{cd}^*$
$H(D_s^+ \rightarrow K^0\ell^+\nu_\ell)$	$A_4V_{cd}^*$

D \rightarrow P $\ell^+\nu_\ell$ 衰变分支比

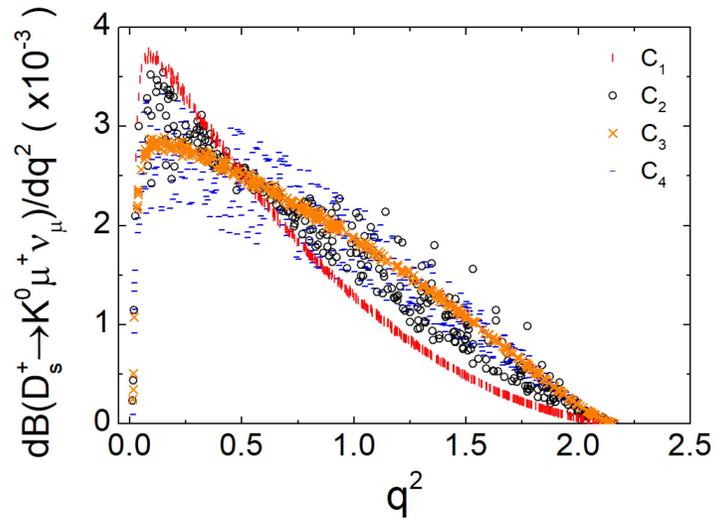
$10^{-3} \sim 10^{-5}$

Branching ratios	Exp. data	Ones in C_1	Ones in C_2	Ones in C_3	Ones in C_4	Previous ones
$\mathcal{B}(D^+ \rightarrow \bar{K}^0 e^+ \nu_e)(\times 10^{-2})$	8.72 ± 0.18	8.84 ± 0.06	8.83 ± 0.07	8.84 ± 0.06	8.83 ± 0.07	
$\mathcal{B}(D^+ \rightarrow \pi^0 e^+ \nu_e)(\times 10^{-3})$	3.72 ± 0.34	3.75 ± 0.05	$5.40 \pm 1.33^\dagger$	$5.04 \pm 0.12^\dagger$	3.70 ± 0.11	
$\mathcal{B}(D^+ \rightarrow \eta e^+ \nu_e)(\times 10^{-3})$	1.11 ± 0.14	1.15 ± 0.05	1.20 ± 0.05	1.20 ± 0.05	0.92 ± 0.08	
$\mathcal{B}(D^+ \rightarrow \eta' e^+ \nu_e)(\times 10^{-4})$	2.0 ± 0.8	2.59 ± 0.14	2.22 ± 0.34	2.09 ± 0.14	1.50 ± 0.20	
$\mathcal{B}(D^0 \rightarrow K^- e^+ \nu_e)(\times 10^{-2})$	3.549 ± 0.052	3.52 ± 0.02	3.52 ± 0.03	3.52 ± 0.03	3.52 ± 0.02	
$\mathcal{B}(D^0 \rightarrow \pi^- e^+ \nu_e)(\times 10^{-3})$	2.91 ± 0.08	2.95 ± 0.03	$4.23 \pm 1.03^\dagger$	$3.97 \pm 0.09^\dagger$	2.89 ± 0.06	
$\mathcal{B}(D_s^+ \rightarrow \eta e^+ \nu_e)(\times 10^{-2})$	2.32 ± 0.16	2.37 ± 0.11	2.34 ± 0.14	2.36 ± 0.12	2.32 ± 0.16	
$\mathcal{B}(D_s^+ \rightarrow \eta' e^+ \nu_e)(\times 10^{-3})$	8.0 ± 1.4	9.05 ± 0.04	8.25 ± 1.13	8.04 ± 0.43	8.02 ± 1.38	
$\mathcal{B}(D_s^+ \rightarrow K^0 e^+ \nu_e)(\times 10^{-3})$	3.4 ± 0.8	3.10 ± 0.08	3.56 ± 0.39	3.54 ± 0.12	3.40 ± 0.80	
$\mathcal{B}(D_s^+ \rightarrow \pi^0 e^+ \nu_e)(\times 10^{-5})$...	1.51 ± 0.07	2.10 ± 0.56	1.96 ± 0.10	1.92 ± 0.13	2.65 ± 0.38 [76]
$\mathcal{B}(D^+ \rightarrow \bar{K}^0 \mu^+ \nu_\mu)(\times 10^{-2})$	8.76 ± 0.38	8.56 ± 0.06	8.69 ± 0.15	8.61 ± 0.06	8.61 ± 0.06	
$\mathcal{B}(D^+ \rightarrow \pi^0 \mu^+ \nu_\mu)(\times 10^{-3})$	3.50 ± 0.30	3.67 ± 0.05	$5.32 \pm 1.31^\dagger$	$4.96 \pm 0.12^\dagger$	3.64 ± 0.10	
$\mathcal{B}(D^+ \rightarrow \eta \mu^+ \nu_\mu)(\times 10^{-3})$	1.04 ± 0.22	1.11 ± 0.05	1.18 ± 0.07	1.17 ± 0.05	0.90 ± 0.08	1.21 [7] 0.75 ± 0.15 [79]
$\mathcal{B}(D^+ \rightarrow \eta' \mu^+ \nu_\mu)(\times 10^{-4})$...	2.42 ± 0.13	2.10 ± 0.33	1.96 ± 0.13	1.41 ± 0.19	2.11 [7] 1.06 ± 0.20 [79]
$\mathcal{B}(D^0 \rightarrow K^- \mu^+ \nu_\mu)(\times 10^{-2})$	3.41 ± 0.08	3.41 ± 0.02	3.44 ± 0.05	3.43 ± 0.02	3.43 ± 0.02	
$\mathcal{B}(D^0 \rightarrow \pi^- \mu^+ \nu_\mu)(\times 10^{-3})$	2.67 ± 0.24	2.89 ± 0.02	$4.17 \pm 1.01^\dagger$	$3.90 \pm 0.09^\dagger$	2.85 ± 0.06	
$\mathcal{B}(D_s^+ \rightarrow \eta \mu^+ \nu_\mu)(\times 10^{-2})$	2.4 ± 1.0	2.30 ± 0.10	2.30 ± 0.17	2.31 ± 0.12	2.26 ± 0.16	
$\mathcal{B}(D_s^+ \rightarrow \eta' \mu^+ \nu_\mu)(\times 10^{-2})$	1.1 ± 1.0	0.86 ± 0.03	0.79 ± 0.11	0.77 ± 0.04	0.76 ± 0.13	
$\mathcal{B}(D_s^+ \rightarrow K^0 \mu^+ \nu_\mu)(\times 10^{-3})$...	3.01 ± 0.08	3.51 ± 0.38	3.46 ± 0.11	3.33 ± 0.78	3.9 [7] 3.85 ± 0.76 [79]
$\mathcal{B}(D_s^+ \rightarrow \pi^0 \mu^+ \nu_\mu)(\times 10^{-5})$...	1.48 ± 0.07	2.09 ± 0.53	1.93 ± 0.10	1.89 ± 0.13	
$\mathcal{B}(D_s^+ \rightarrow \pi^0 \tau^+ \nu_\tau)(\times 10^{-10})$...	3.45 ± 0.21	160.34 ± 149.53	4.20 ± 0.26	4.08 ± 0.34	$(27 \sim 36)$ [76]

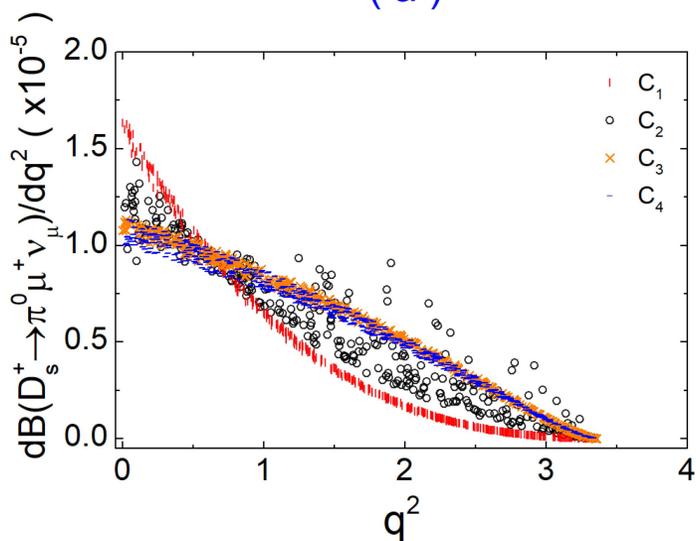
D \rightarrow P $\ell^+\nu_\ell$ 分支比对 q^2 的依赖关系



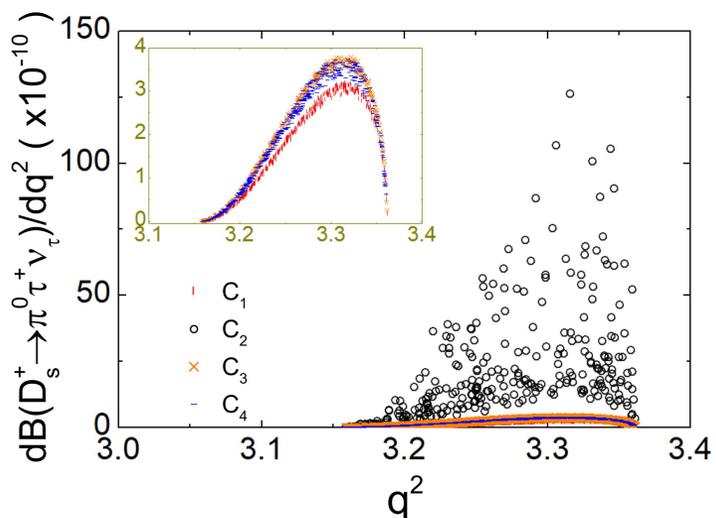
(a)



(b)



(c)



(d)

C₁ 红色
C₂ 黑色
C₃ 黄色
C₄ 蓝色

(a) $D^+ \rightarrow \eta' \mu^+ \nu_\mu$

(b) $D_s^+ \rightarrow K^0 \mu^+ \nu_\mu$

(c) $D_s^+ \rightarrow \pi^0 \mu^+ \nu_\mu$

(d) $D_s^+ \rightarrow \pi^0 \tau^+ \nu_\tau$

D \rightarrow $A\ell^+\nu_\ell$ 衰变振幅

$$A = c_0^A, B = c_0^B$$

TABLE II: The hadronic helicity amplitudes for the $D \rightarrow A\ell^+\nu$ decays including both the SU(3) flavor symmetry and the SU(3) flavor breaking contributions. $A_1 \equiv c_0^A + c_1^A - 2c_2^A$, $A_2 \equiv c_0^A - 2c_1^A - 2c_2^A$, $A_3 \equiv c_0^A + c_1^A + c_2^A$, $A_4 \equiv c_0^A - 2c_1^A + c_2^A$. $B_1 \equiv c_0^B + c_1^B - 2c_2^B$, $B_2 \equiv c_0^B - 2c_1^B - 2c_2^B$, $B_3 \equiv c_0^B + c_1^B + c_2^B$, $B_4 \equiv c_0^B - 2c_1^B + c_2^B$. $A_1 = A_2 = A_3 = A_4 = c_0^A$, $B_1 = B_2 = B_3 = B_4 = c_0^B$ if neglecting the SU(3) flavor breaking c_1^A, c_2^A and c_1^B, c_2^B terms.

Decay modes	SU(3) hadronic amplitudes	Decay modes	SU(3) hadronic amplitudes
$D^0 \rightarrow K_1(1270)^-\ell^+\nu_\ell$	$(\sin\theta_{K_1}A_1 + \cos\theta_{K_1}B_1) V_{cs}^*$	$D^0 \rightarrow K_1(1400)^-\ell^+\nu_\ell$	$(\cos\theta_{K_1}A_1 - \sin\theta_{K_1}B_1) V_{cs}^*$
$D^+ \rightarrow \bar{K}_1(1270)^0\ell^+\nu_\ell$	$(\sin\theta_{K_1}A_1 + \cos\theta_{K_1}B_1) V_{cs}^*$	$D^+ \rightarrow \bar{K}_1(1400)^0\ell^+\nu_\ell$	$(\cos\theta_{K_1}A_1 - \sin\theta_{K_1}B_1) V_{cs}^*$
$D_s^+ \rightarrow f_1(1285)\ell^+\nu_\ell$	$(\frac{1}{\sqrt{3}}\cos\theta_{3P_1} - \sqrt{\frac{2}{3}}\sin\theta_{3P_1}) A_2 V_{cs}^*$	$D_s^+ \rightarrow f_1(1420)\ell^+\nu_\ell$	$(-\frac{1}{\sqrt{3}}\sin\theta_{3P_1} - \sqrt{\frac{2}{3}}\cos\theta_{3P_1}) A_2 V_{cs}^*$
$D_s^+ \rightarrow h_1(1170)\ell^+\nu_\ell$	$(\frac{1}{\sqrt{3}}\cos\theta_{1P_1} - \sqrt{\frac{2}{3}}\sin\theta_{1P_1}) B_2 V_{cs}^*$	$D_s^+ \rightarrow h_1(1415)\ell^+\nu_\ell$	$(-\frac{1}{\sqrt{3}}\sin\theta_{1P_1} - \sqrt{\frac{2}{3}}\cos\theta_{1P_1}) B_2 V_{cs}^*$
$D^0 \rightarrow a_1(1260)^-\ell^+\nu_\ell$	$A_3 V_{cd}^*$	$D^0 \rightarrow b_1(1235)^-\ell^+\nu_\ell$	$B_3 V_{cd}^*$
$D^+ \rightarrow a_1(1260)^0\ell^+\nu_\ell$	$-\frac{1}{\sqrt{2}} A_3 V_{cd}^*$	$D^+ \rightarrow b_1(1235)^0\ell^+\nu_\ell$	$-\frac{1}{\sqrt{2}} B_3 V_{cd}^*$
$D^+ \rightarrow f_1(1285)\ell^+\nu_\ell$	$(\frac{1}{\sqrt{3}}\cos\theta_{3P_1} + \frac{1}{\sqrt{6}}\sin\theta_{3P_1}) A_3 V_{cd}^*$	$D^+ \rightarrow f_1(1420)\ell^+\nu_\ell$	$(-\frac{1}{\sqrt{3}}\sin\theta_{3P_1} + \frac{1}{\sqrt{6}}\cos\theta_{3P_1}) A_3 V_{cd}^*$
$D^+ \rightarrow h_1(1170)\ell^+\nu_\ell$	$(\frac{1}{\sqrt{3}}\cos\theta_{1P_1} + \frac{1}{\sqrt{6}}\sin\theta_{1P_1}) B_3 V_{cd}^*$	$D^+ \rightarrow h_1(1415)\ell^+\nu_\ell$	$(-\frac{1}{\sqrt{3}}\sin\theta_{1P_1} + \frac{1}{\sqrt{6}}\cos\theta_{1P_1}) B_3 V_{cd}^*$
$D_s^+ \rightarrow K_1(1270)^0\ell^+\nu_\ell$	$(\sin\theta_{K_1}A_4 + \cos\theta_{K_1}B_4) V_{cd}^*$	$D_s^+ \rightarrow K_1(1400)^0\ell^+\nu_\ell$	$(\cos\theta_{K_1}A_4 - \sin\theta_{K_1}B_4) V_{cd}^*$

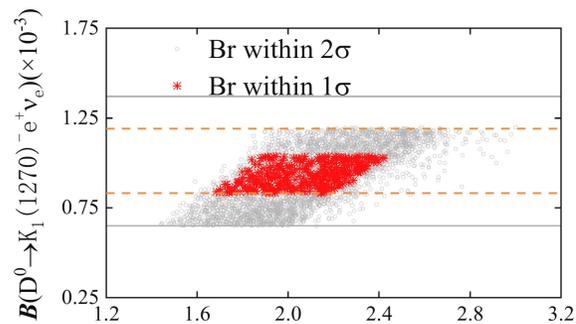
D \rightarrow Ae $^+$ ν_e 衰变分支比

$$\mathcal{B}(D^+ \rightarrow b_1(1235)^0 e^+ \nu_e, b_1(1235)^0 \rightarrow \omega \pi^0) \leq 1.75 \times 10^{-4},$$

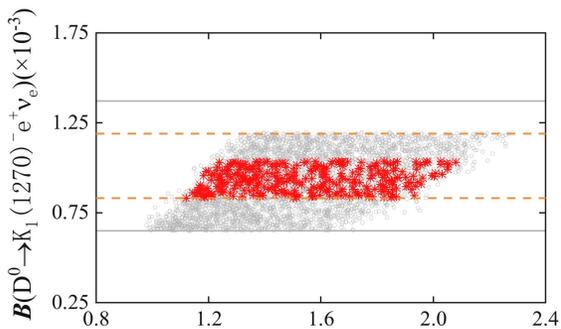
$$\mathcal{B}(D^0 \rightarrow b_1(1235)^- e^+ \nu_e, b_1(1235)^- \rightarrow \omega \pi^-) \leq 1.12 \times 10^{-4}.$$

Branching ratios	Ones within 1 σ error	Ones within 2 σ error	Previous results
$\mathcal{B}(D^0 \rightarrow K_1(1270)^- e^+ \nu_e)(\times 10^{-3})$	0.93 ± 0.10	0.92 ± 0.27	...
$\mathcal{B}(D^0 \rightarrow K_1(1400)^- e^+ \nu_e)(\times 10^{-3})$	0.02 ± 0.02	0.08 ± 0.08	...
$\mathcal{B}(D^+ \rightarrow \bar{K}_1(1270)^0 e^+ \nu_e)(\times 10^{-3})$	2.43 ± 0.27	2.39 ± 0.71	3.2 ± 0.40 [43]
$\mathcal{B}(D^+ \rightarrow \bar{K}_1(1400)^0 e^+ \nu_e)(\times 10^{-3})$	0.05 ± 0.05	0.20 ± 0.20	{0.005, 0.02} [43]
$\mathcal{B}(D_s^+ \rightarrow f_1(1285) e^+ \nu_e)(\times 10^{-3})$	0.86 ± 0.73	1.16 ± 1.16	{0.06, 0.36} [43]
$\mathcal{B}(D_s^+ \rightarrow f_1(1420) e^+ \nu_e)(\times 10^{-3})$	0.21 ± 0.21	0.39 ± 0.39	0.25 ± 0.05 [43]
$\mathcal{B}(D_s^+ \rightarrow h_1(1170) e^+ \nu_e)(\times 10^{-3})$	0.26 ± 0.26	0.43 ± 0.43	{0, 0.197} [43]
$\mathcal{B}(D_s^+ \rightarrow h_1(1415) e^+ \nu_e)(\times 10^{-3})$	0.28 ± 0.16	0.31 ± 0.24	0.64 ± 0.07 [43]
$\mathcal{B}(D^0 \rightarrow a_1(1260)^- e^+ \nu_e)(\times 10^{-5})$	4.46 ± 2.32	7.72 ± 6.54	6.90 [40], $5.261_{-0.639}^{+0.745}$ [39]
$\mathcal{B}(D^0 \rightarrow b_1(1235)^- e^+ \nu_e)(\times 10^{-5})$	2.61 ± 1.44	2.95 ± 2.06	4.85 [40]
$\mathcal{B}(D^+ \rightarrow a_1(1260)^0 e^+ \nu_e)(\times 10^{-5})$	5.79 ± 3.00	9.95 ± 8.41	9.38 [40], $6.673_{-0.811}^{+0.947}$ [39]
$\mathcal{B}(D^+ \rightarrow b_1(1235)^0 e^+ \nu_e)(\times 10^{-5})$	3.41 ± 1.88	3.84 ± 2.71	7.4 ± 0.70 [43], 6.58 [40]
$\mathcal{B}(D^+ \rightarrow f_1(1285) e^+ \nu_e)(\times 10^{-5})$	1.88 ± 1.88	3.33 ± 3.33	3.7 ± 0.80 [43]
$\mathcal{B}(D^+ \rightarrow f_1(1420) e^+ \nu_e)(\times 10^{-5})$	0.64 ± 0.54	0.84 ± 0.84	{0.02, 0.14} [43]
$\mathcal{B}(D^+ \rightarrow h_1(1170) e^+ \nu_e)(\times 10^{-5})$	5.28 ± 3.00	5.59 ± 4.50	14 ± 1.50 [43]
$\mathcal{B}(D^+ \rightarrow h_1(1415) e^+ \nu_e)(\times 10^{-5})$	0.11 ± 0.11	0.18 ± 0.18	{0, 0.02} [43]
$\mathcal{B}(D_s^+ \rightarrow K_1(1270)^0 e^+ \nu_e)(\times 10^{-5})$	11.77 ± 1.39	11.70 ± 3.73	17 ± 2.00 [43]
$\mathcal{B}(D_s^+ \rightarrow K_1(1400)^0 e^+ \nu_e)(\times 10^{-5})$	0.32 ± 0.32	1.17 ± 1.17	{0.05, 0.14} [43]

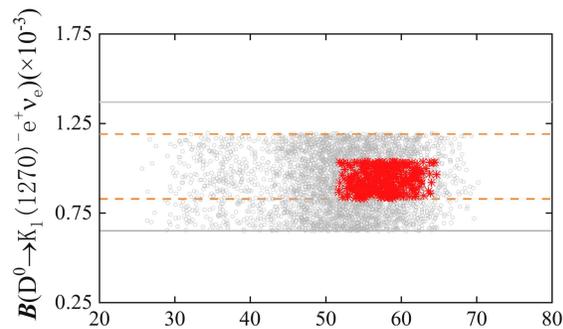
D \rightarrow A $l^+ \nu_l$ 对参数变化的敏感度



A
(a₁)

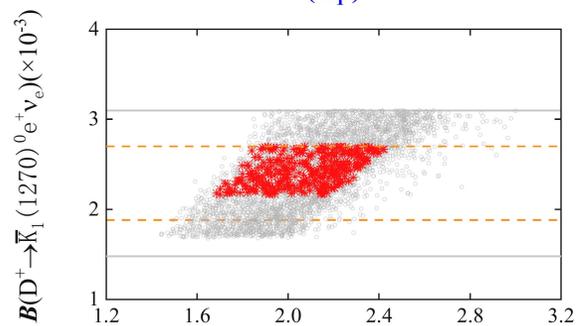


B
(a₂)

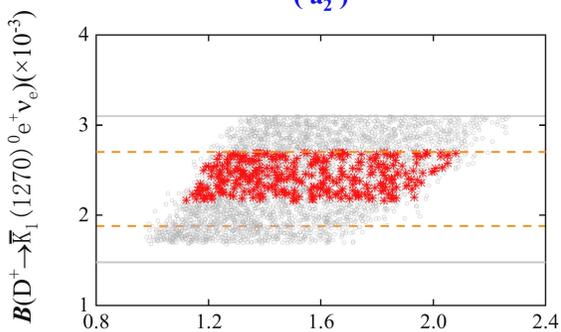


θ_{K1}
(a₃)

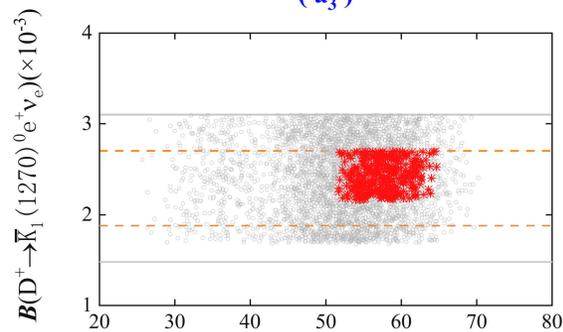
1.37 UL within 2σ
1.19 UL within 1σ
0.83 LL within 1σ
0.65 LL within 2σ



A
(b₁)

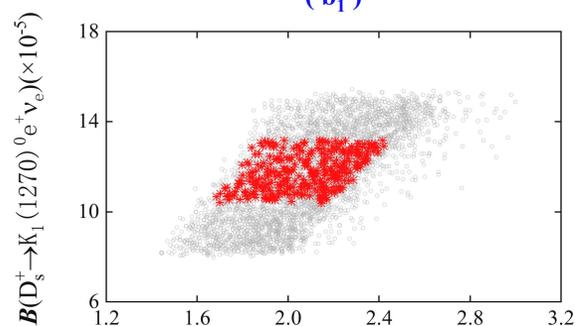


B
(b₂)

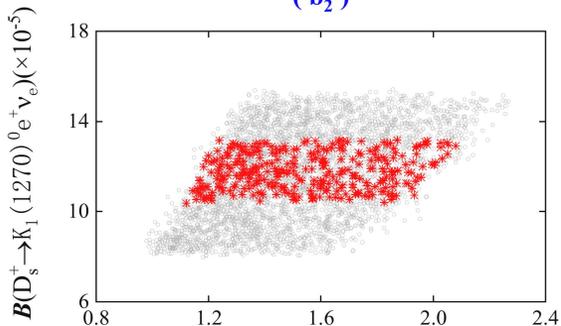


θ_{K1}
(b₃)

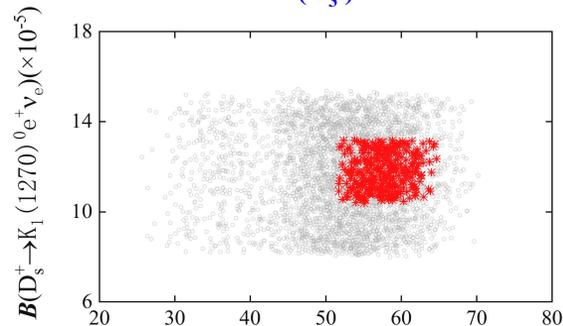
3.10 UL within 2σ
2.70 UL within 1σ
1.88 LL within 1σ
1.48 LL within 2σ



A
(c₁)

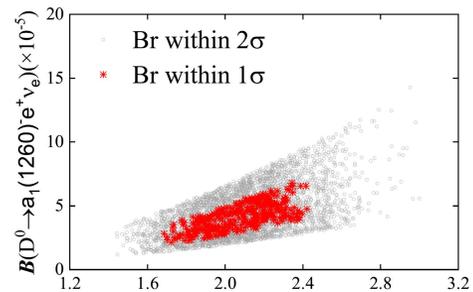


B
(c₂)

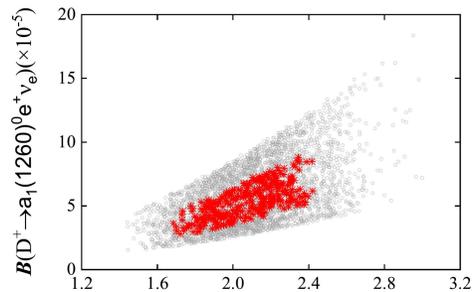


θ_{K1}
(c₃)

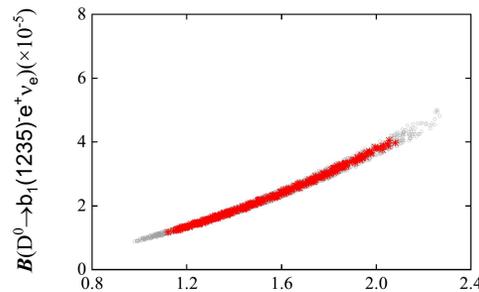
D → Ae⁺v_e 数值分析



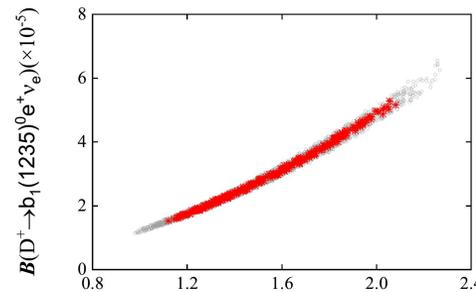
A
(a₁)



A
(a₂)



B
(b₃)



B
(b₄)

预言 θ_{3P1} :

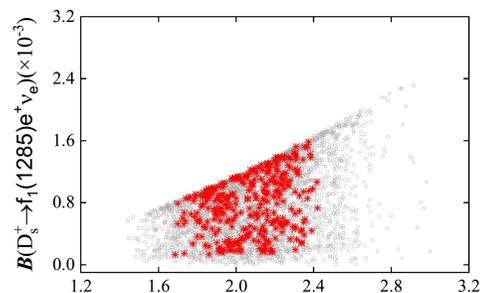
[55°, 125°] 1σ

[30°, 145°] 2σ

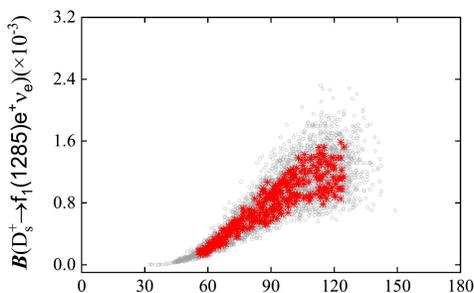
预言 θ_{1P1} :

[8°, 58°] 1σ

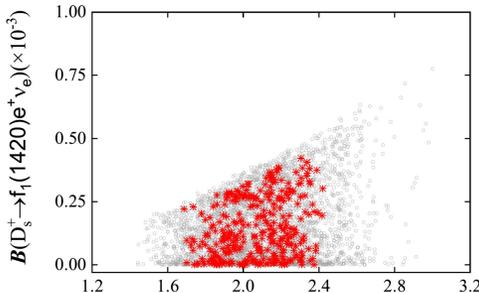
[-10°, 75°] 2σ



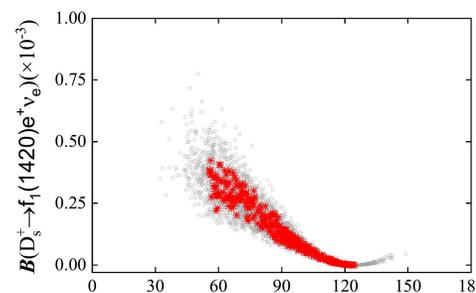
A
(b₁)



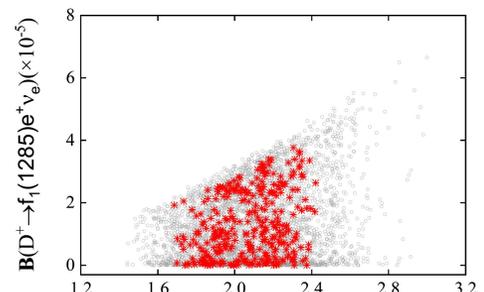
θ_{3P1}
(b₂)



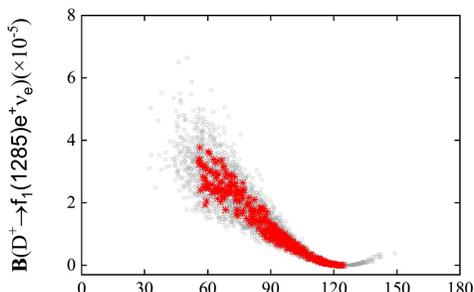
A
(b₃)



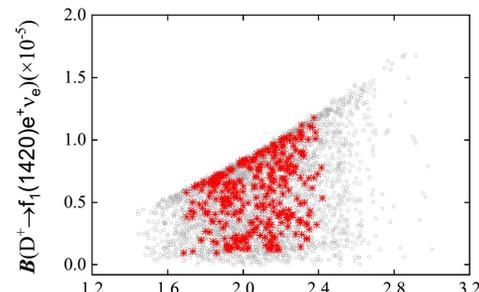
θ_{3P1}
(b₄)



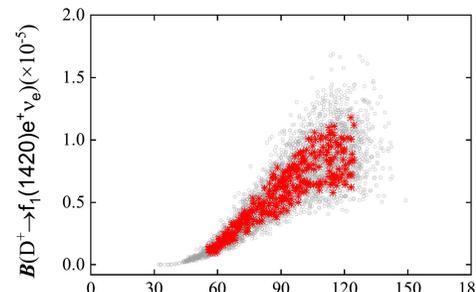
A
(c₁)



θ_{3P1}
(c₂)



A
(c₃)



θ_{3P1}
(c₄)



目录

01 研究背景与方法

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04 总结与展望

总结与展望

- 用SU(3)味对称方法研究D介子半轻三体衰变;
- 预言了一些实验上尚未观测到的衰变过程, 并且有一些衰变过程是首次预言的;
- 我们的预言结果未来有可能在BESIII, LHCb 和 BelleII上观测到;

谢谢!