Energy Correlator for Hadronic Structures

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味物理讲座, Oct 10, 2024

XL and Zhu, **PRL** 130 (2023), 9, 9 Liu, XL, Pan, Yuan and Zhu, **PRL** 130 (2023) 18, 18 Cao, XL and Zhu, **PRD** 107 (2023) 11, 114008 Chen, XL and Ma, **PRL** (2024) accepted + ...





Outline

- **O** Non-perturbative Structure studies
- O Energy Correlator
- **O** Nucleon energy Correlators (NECs)

 - **O** Phenomenology
- O Quarkonium Energy Correlator

O Definition, measurement, factorization and properties

O New insights into the non-perturbative structures ???



Quark and gluon internal motion

Major focus of the EicC, EIC ...

Collinear Parton Distribution Functions (PDFs)

hard probe, e.g., DIS

O inclusive over X, clean.O not differential enough, lose information

Transverse Moment Dependent-PDFs (TMDs)

$$f_{q/p}(x,k_t) = \int_{-\infty}^{\infty} \frac{dy^- dy_t}{(2\pi)^3} e^{ixp^+ y^-} e^{ik_t \cdot y_t} \frac{\gamma^+}{2} \langle P | \bar{\psi}(0,t) \rangle dt$$





SIDIS, Breit Frame





O Major tool for structure studiesO Enforce the b-to-b configuration

Transverse Moment Dependent-PDFs (TMDs)

$$f_{q/p}(x,k_t) = \int_{-\infty}^{\infty} \frac{dy^- dy_t}{(2\pi)^3} e^{ixp^+ y^-} e^{ik_t \cdot y_t} \frac{\gamma^+}{2} \langle P | \bar{\psi}(0,t) \rangle dt$$





- Major tool for structure studies
- Soft contamination
- Sudakov suppression $\sigma(k_T) \propto \frac{1}{q_T^2} e^{-\frac{Q^2}{q_T^2}}$
- O Distort azimuthal asymmetry

Hatta, Xiao, Yuan, Zhou, PRL 2021

Transverse Moment Dependent-PDFs (TMDs)

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$(0) \mathscr{L} \psi(y_t, y^-) | P \rangle$



• Still lose information

• How $c\bar{c} \rightarrow J/\psi$?

- NRQCD: encoded in $\langle \mathcal{O}_1 \rangle$, $\langle \mathcal{O}_8 \rangle$
 - O remains largely unknown: amount of energy released? Energy Distribution?



• • •

Andres, Basham, Belitsky, Brown, H Chen, Dixon, Dominguez, Elayavalli, Ellis, J Gao, Hofman, Hohenegger, Holguin, Jaarsma, ZB Kang, Kologlu, Korchemsky, Kravchuk, Komiske, Lee, HT Li, YB Li, Love, MX Luo, Maldacena, Meçaj, Marquet, Moult, Pathak, Procura, DY Shao, Simmons-Duffin, Sokatchev, Thaler, van Velzen, W. Wang, X-N Wang, Waalewijn, M Xiao, K Yan, TZ Yang, F Yuan, Zhang, Zhiboedov, HX Zhu +



Energy-Energy-Correlator (EEC)



Sterman, 1975 Bashman, et al. 1978



Energy-Energy-Correlator (EEC)



 $= \frac{1}{\Omega^2} \langle J^{\mu}(x) \mathscr{E}(n_1) \mathscr{E}(n_2) J^{\nu}(0) \rangle_{\Omega}$

 $\mathscr{E}(n) = \int_{0}^{\infty} dt \lim_{r \to \infty} T_{0\vec{n}}(t, \vec{n}r)r^{2}$ detector by the light-ray operator

• Casy to implement, "Jet w/o jet" **O** Perturbatively predictable e.g., Gao, Li, Moult, Zhu, 2023, Chen, et al. 2024 **O** Dual description

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Jet Structure in e+ e- Annihilation with Massless Hadrons

George F. Sterman (Illinois U., Urbana) Dec, 1975 10 pages Report number: ILL-TH-75-32 View in: KEK scanned document

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Our ensembles will thus be specified in terms of set states. To make this idea more quantitative we defi "angular energy current" in the e^+e^- CM frame:

$$j_{a}(\Omega) = \sum_{i=1}^{n} \eta_{i} \delta(\Omega - \omega_{i})$$



Ian Moult, MITP talk 2024

Conformal collider physics: Energy and charge correlations

Diego M. Hofman^a and Juan Maldacena^b

^a Joseph Henry Laboratories, Princeton University, Princeton, NJ 08544, USA ^bSchool of Natural Sciences, Institute for Advanced Study Princeton, NJ 08540, USA



$\mathscr{E}(n_1)\mathscr{E}(n_2)\sim \theta^{-2+\gamma(3)}\mathcal{O}, \theta\to 0$

Scaling rule by Hofman, Maldecena, 2008 conformal theory



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Energy Correlators As of 2024





Yang, He, Moult, Wang, 2024 PRL + ...

A new probe of QGP



12th International Conference on Hard and Electromagnetic Probes of High-Energy Nuclear Collisions







Full Spectrum with high precision



Yen-Jie Lee, MITP talk 2024

Full Spectrum with high precision



Yen-Jie Lee, MITP talk 2024

Nucleon Energy Correlator

Operator Definition

Nucleon EEC XL and Zhu, Phys. Rev. Lett. 130 (2023), 9, 9

$$f_{q,EEC}(x,\theta) = \int_{-\infty}^{\infty} \frac{dy^{-}}{2\pi} e^{ixp^{+}y^{-}} \frac{\gamma^{+}}{2} \langle P | \bar{\psi}(0) \mathcal{E} \rangle$$

$$\mathscr{E}(n) = \int_0^\infty dt \lim_{r \to \infty} T_{0\vec{n}}(t, \vec{n}r) r^2$$

- **O** Energy correlator in the forward region.
- **O** Probe directly the broken proton
- Purely collinear object, insensitive to soft radiations, e.g. no Sudakov suppression
- **O** Transverse dynamics through $\mathscr{E}(\theta)$
- Can be generalized to multiple-point correlation



 $S(\theta) \mathscr{L}\psi(y^{-}) | P \rangle$

Nucleon EEC XL and Zhu, Phys. Rev. Lett. 130 (2023), 9, 9

$$\circ \Sigma_N(Q^2, \theta) = \sum_i \int dx_B x_B^{N-1} \frac{E_i}{E_P} d\sigma(x_B, Q^2, p_i) \Theta(\theta - \theta_i)$$



hard probe for trigger



O Measurement in DIS

- Full inclusive measurement, no jet/hadrons, weighted by E_i
- Different θ 's probe different physics



Nucleon EEC XL and Zhu, Phys. Rev. Lett. 130 (2023), 9, 9

$$\circ \Sigma_N(Q^2,\theta) = \sum_i \int dx_B x_B^{N-1} \frac{E_i}{E_P} d\sigma(x_B,Q^2)$$



 $(p_i)\Theta(\theta - \theta_i)$

- When $\theta Q \ll Q$, DIS type factorization $\Sigma_N(Q^2, \theta) = \left[u^{N-1} \hat{\sigma} \left(u, Q^2, \mu \right) f_{\text{EEC}}(N, \ln \frac{\theta Q}{u\mu}) \right]$
 - O Derived by SCET Cao, XL, Zhu, 2303.01530
 - **O** rigorous QCD derivation by relating to the fracture function through sum rules

Chen, Ma, Tong, 2406.08559 **O** Free of soft contribution

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Breit Frame LO



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Breit Frame NLO



 $(\theta, p_i)\Theta(\theta - \theta_i)$

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•
$$\Sigma_N(Q^2, \theta) = \sum_i \int dx_B x_B^{N-1} \frac{E_i}{E_P} d\sigma(x_B, Q^2)$$



Similar factorization form

 $(p_i)\Theta(\theta - \theta_i)$

• When
$$\theta Q \ll Q$$
, DIS type factorization
 $\Sigma_N(Q^2, \theta) = \int u^{N-1} \hat{\sigma}(u, Q^2, \mu) f_{\text{EEC}}(N, \ln \theta)$

• Space like version of the EEC in e^+e^-

$$\Sigma = \int du u^2 \sigma(u, \mu) J(\mu, \ln \frac{\theta u Q}{\mu})$$
Dixon, Moult, Zh

















Moult talk







$$\Lambda_{\text{QCD}} \ll \theta Q \ll Q$$

• Perturbatively calculable

$$f_{\text{EEC}}^{(0)}(\theta) \propto \left[\frac{1}{\theta^2}(1-x)P(x)\right] \times \left[\xi f(\xi)\right]$$

• Dynamics dominated by coll. splittin
• Power law: $\theta^{-2+\gamma}$, γ by $P(N)$ + coll.



Moult talk







Deep NP region **O** Un-correlated distribution $d\Sigma/d\theta \sim \theta$



Moult talk





NP region

- O Enhanced NP region, vs. TMD
- **O** To be determined by future measurements

Encodes info. on proton intrinsic 0 structure and NP dynamics





Phenomenology



- **O** It will be nice if experiments can confirm the framework.
- **O** Not that forward to probe the scaling rule region and the onset of the transition region, $y \leq 2.5$, for $Q \leq 10 \,\text{GeV}$ @HERA, CEBAF@JLab maybe perfect for the deep NP region
- Possible precision measurement of TMD in the b-to-b region





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• The small θ region and TMD region can be related





 $\vec{k}_t = -\sum_{i \in X} \vec{p}_{i,t} = -\sum_{i \in X} E_i \sin \theta_i (\cos \phi_i, \sin \phi_i)$









 $\vec{k}_t = -\int d\theta d\phi \sum_i E \sin \theta (\cos \phi, \sin \phi) \delta(\Omega - \Omega_i)$ $= -\int d\theta d\phi \sin \theta (\cos \phi, \sin \phi) \mathscr{E}(\Omega)$ $\int^{\mu} dk_t k_t^n f(k_t) = (-)^n \int^R \prod d\Omega w(\Omega_1) \dots w(\Omega_n) \langle P | \dots \mathscr{E}(\Omega_1) \dots \mathscr{E}(\Omega_n) \dots | P \rangle$

For TMD TMM see e.g.: del Rio, Prokudin, Scimemi, Vladimirov, arXiv:2402.01836v1

XL, Zhu, arxiv: 2403.08874 XL, Shao, Zhu, in preparation



$\mathscr{E}(\Omega) = \sum E_i \delta(\Omega - \Omega_i)$ $i \in X$



$$\int dk_t k_t^n f(k_t) = (-)^n \int \prod_n d\Omega w(\Omega_1) \dots w(\Omega_n) \langle P | \dots \mathscr{E}(\Omega_1) \dots$$

- **O** TMD PDFs (moment) can be obtained by measuring N-pt Nucleon Energy Correlator, by suitably selecting $w(\Omega)$
- O Inclusive measurement! Do not force b-to-b limit, no jets/fragmentation function involved.
- Nucleon Energy Correlator can be regarded as a generating observable, contains more comprehensive information

XL, Zhu, arxiv: 2403.08874 XL, Shao, Zhu, in preparation

 $\mathscr{E}(\Omega_n) \dots |P\rangle$





Motivation

Quarkonium Physics

- O regarded as an excellent place to study nonpert phenomenon for a long time
- **O** How $c\bar{c} \rightarrow J/\psi$?
 - **ONRQCD:** encoded in $\langle \mathcal{O}_1 \rangle$, $\langle \mathcal{O}_8 \rangle$
 - O remains largely unknown: amount of energy released? Energy Distribution?



Motivation

Recent attempts using jet

Probing Quarkonium Production Mechanisms with Jet Substructure Matthew Baumgart^a,¹ Adam K. Leibovich^b,² Thomas Mehen^c,³ and Ira Z. Rothstein^{d1} ¹Department of Physics, Carnegie Mellon University, Pittsburgh, PA 15213 ²Pittsburgh Particle Physics Astrophysics and Cosmology Center (PITT PACC) $z = E_{J/\psi}/E_J$ Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA 15260 ³Department of Physics, Duke University, Durham, NC 27708 (Dated: June 27, 2018)

Unlike light hadron fragmentation, $D_{q \rightarrow J/\psi}(z)$ dominated by perturbative radiations: $E_I \rightarrow 2m_O$



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Chance to "see" hadronziation?



Jet-jet-correlation

$$\Sigma_{\text{EEC}} = \frac{1}{\sigma} \int d\sigma \sum_{ij} \frac{E_i E_j}{Q^2} \delta(\chi - \theta_{ij})$$

EEC

Quarkonium EC J/ψ "quarkonim-Jet correlation" In quarkonium rest frame!! $\Sigma_{QEC}(\chi) \propto \frac{1}{\sigma_{J/\psi}} \int d\sigma_{J/\psi} \frac{E_i}{M} \delta(\chi - \chi_i)$

~ average energy emitted at the angle χ





$$\circ \Sigma_{QEC} = \Sigma_{QEC,P.T.} + \Sigma_{QEC,had.}$$

- Hadronization enters as an additive correction, not in the form of convolution
- **O** Hadronization could be large

$$\Sigma_{QEC,P.T.} \sim \alpha_s(\mu) \frac{E(\chi)}{M} E^2(\chi) \langle \mathcal{O}_{1,8} \rangle,$$

$$\Sigma_{QEC,had.} \sim \frac{Mv}{M} M^2 v^2 \langle \mathcal{O}_{1,8} \rangle$$



$$\circ \Sigma_{QEC} = \Sigma_{QEC,P.T.} + \Sigma_{QEC,had.}$$

- Hadronization enters as an additive correction, not in the form of convolution
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$$\Sigma_{QEC,P.T.} \sim \alpha_s(\mu) \frac{E(\chi)}{M} E^2(\chi) \langle \mathcal{O}_{1,8} \rangle,$$

$$\Sigma_{QEC,had.} \sim \frac{Mv}{M} M^2 v^2 \langle \mathcal{O}_{1,8} \rangle$$

Quarkonium EC

for $J/\psi \quad \alpha_{\rm s}(M) \sim v^2, v \sim 0.5$ $\Sigma_{QEC,had}$ / $\Sigma_{QEC,P.T.} \sim \frac{Mv}{\alpha_s E} \frac{M^2 v^2}{E^2} \sim \frac{v^3}{\alpha_s} \frac{M^3}{E(\chi)^3}$ If $M/E(\chi) \sim 1$

 $\Sigma_{QEC,had} / \Sigma_{QEC,P,T} \sim 50\%!$

Excellent place to study the non-perturbative physics!



Generic J/ψ production configuration in pQCD



Chen, XL and Ma, PRL (2024) accepted

dead-cone effects Dokshitzer et al., J. Phys. G

$$d\sigma_{Q \to Qg} \sim \frac{\alpha_s C_F}{\pi} \frac{dE_g}{E_g} \frac{\theta^2 d\theta^2}{[\theta^2 + \theta_0^2]^2}$$

$$\theta_0 \sim \frac{M}{E_{\bar{J}/\psi}} \sim \frac{2M}{\sqrt{\hat{s}}} = 2\sqrt{r}$$

Generic J/ψ production configuration in pQCD



 $E_{J_{near}} \sim M$ And further suppressed by dead cone $E_{J_{away}}/E_{J_{near}} \sim \frac{1}{2}$ boost factor² $\sim \frac{2}{r}$ hard radiation depopulated, $E_{J_{near}} \sim M$ -1 cosχ 0







Sizable hadronization effect!!



Sizable hadronization effect!!



cosχ

Chen, XL and Ma, PRL (2024) accepted

Ignore interference, rotational covariant $E(\chi) = E$ [8] interference, boost covariant $E \sim p_T / \sin \chi$ 8 **|X**| [8] 000000 1.0 Wilson line



Relative size between non-inter vs interference Ignore interfer



Chen, XL and Ma, PRL (2024) accepted

Ignore interference, rotational covariant



interference, boost covariant



1.0





Conclusions This Talk



O PropertiesO A new tool for NP physics



Outlook

New perspective to the NP hadron structures



$$\langle \mathcal{E}(\vec{n}_1')\mathcal{E}(\vec{n}_2')\rangle = (\frac{q^0}{4\pi})^2 \left[1 + \frac{6\pi^2}{\lambda}(\cos^2\theta_{12} - \frac{1}{3}) - \frac{1}{3}\right]$$

Hofman, Maldacena, 2008



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Thanks

