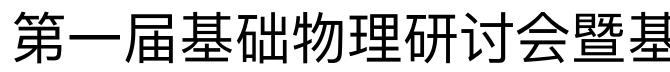


Light QCD exotics at BESIII

Institute of High Energy Physics, CAS

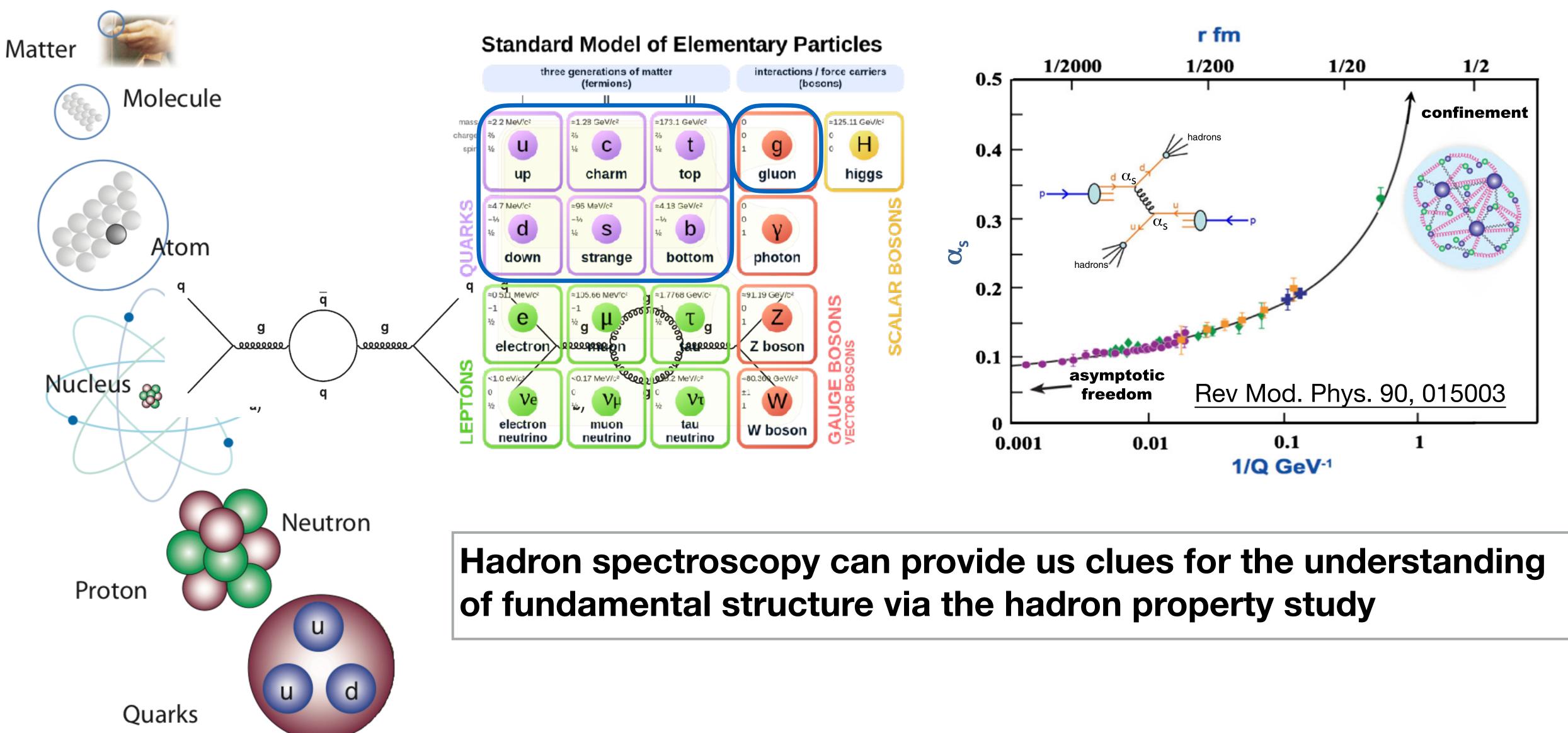


Yanping Huang

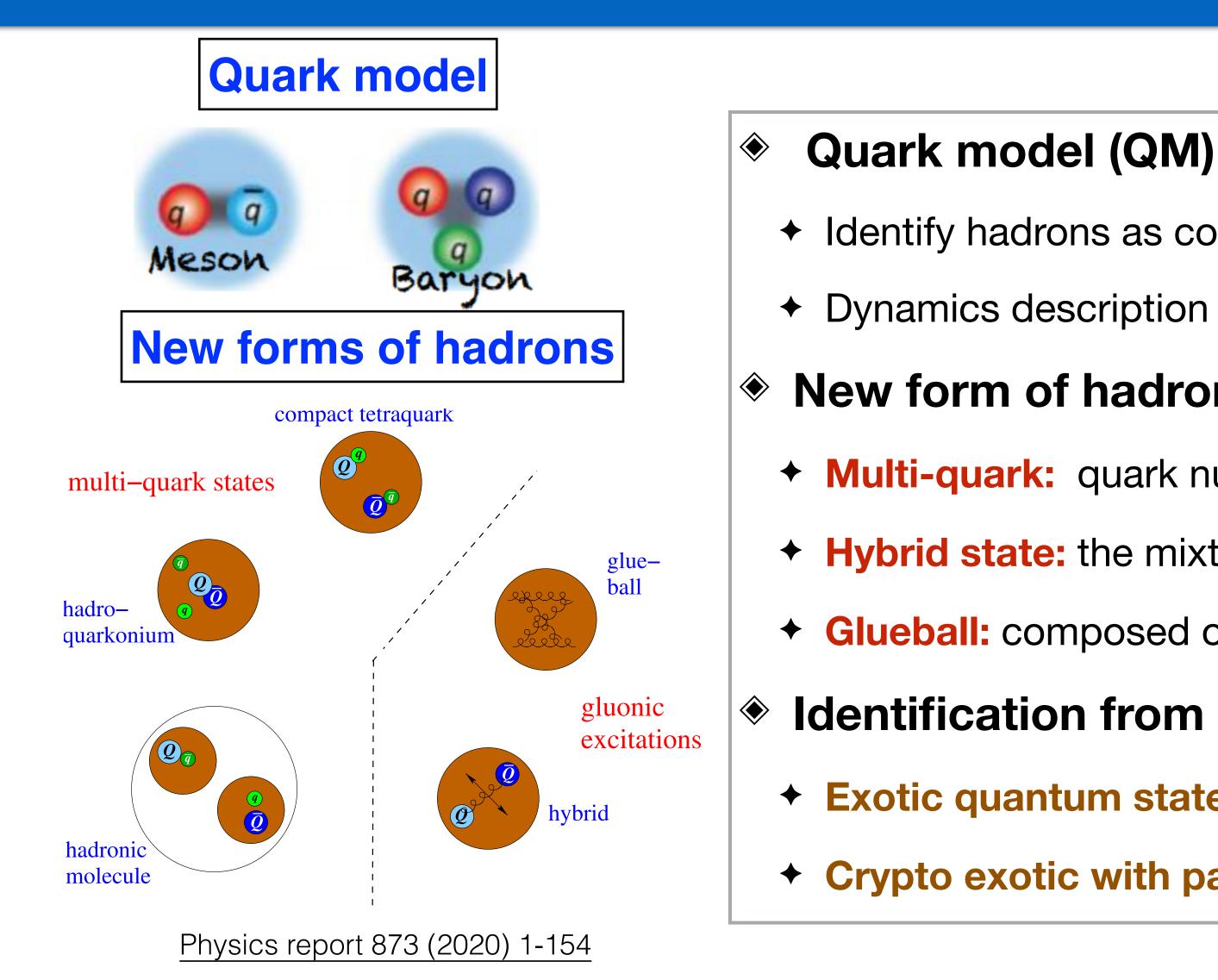
第一届基础物理研讨会暨基础物理平台年会,2024.11.1,河南科学院

Fundamental Structure of Matters









Many candidates, but no unambiguous hadrons with nonstandard structure have established

Forms of hadrons

- Identify hadrons as compound objects consisting of quarks and antiquarks
- Dynamics description inside hadrons

New form of hadrons:

- Multi-quark: quark number >= 4
- Hybrid state: the mixture of quark and gluon
 - **Glueball:** composed of gluons
- Identification from QM: challenging
 - **Exotic quantum states**
 - **Crypto exotic with particular properties**



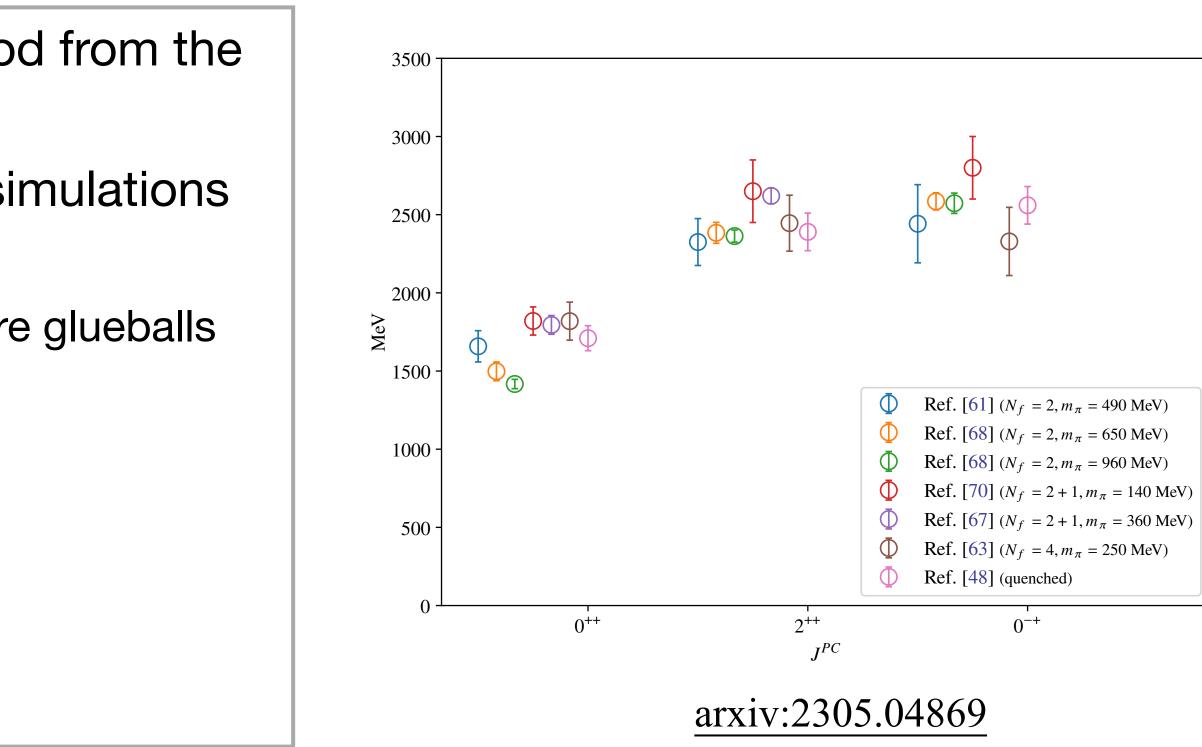




- - Glueballs to QCD is just as important as Higgs Boson to EW
- Lattice QCD (LQCD) is a non-perturbative method from the first principles in theory.
- **Different lattice QCD groups** (including lattice simulations with dynamical quarks)
 - Predictions on masses and production rates of pure glueballs
 - Consistent results and expected to be reliable.
- Lattice QCD predictions on glueball masses:
 - **0++ ground state:** 1.5 1.7 GeV/c²
 - ◆ 2++ ground state: 2.3 2.4GeV/c²
 - ◆ 0-+ ground state: 2.3 2.6GeV/c²

Glueballs

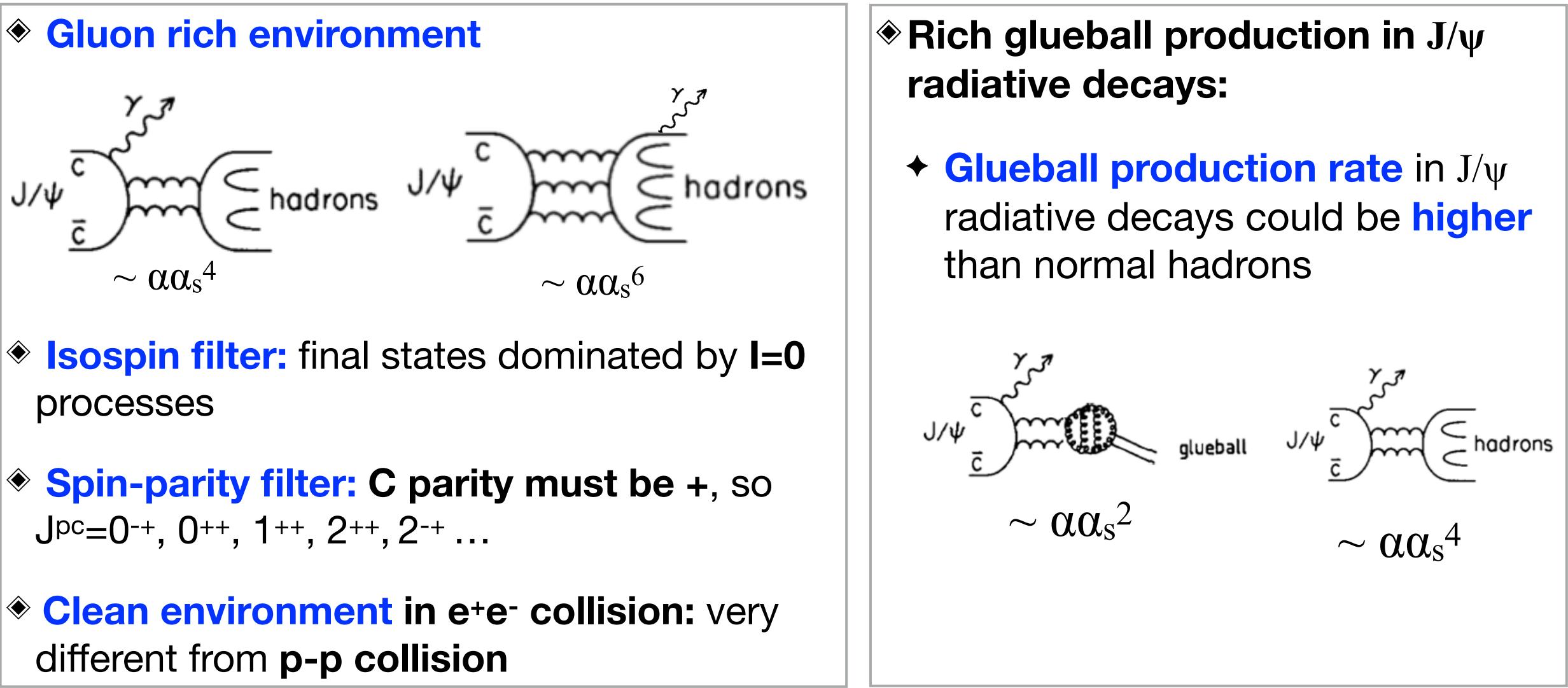
The basic theory for strong interactions is quantum chromodynamics (QCD) Gluon self-interaction: prediction of non-Abelian Gauge SU(3) QCD theory + Glueballs are unique particles formed with force carriers via self-interactions





4

Glueball production in J/\psi Radiative decay



\Rightarrow J/ ψ Radiative decay is an ideal place to search for glueballs

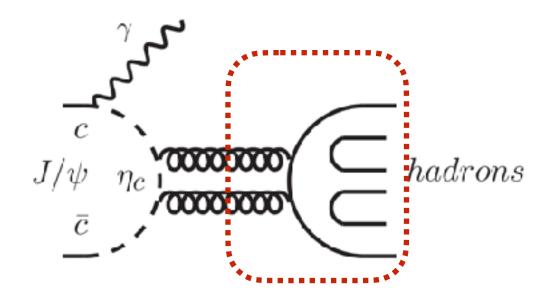




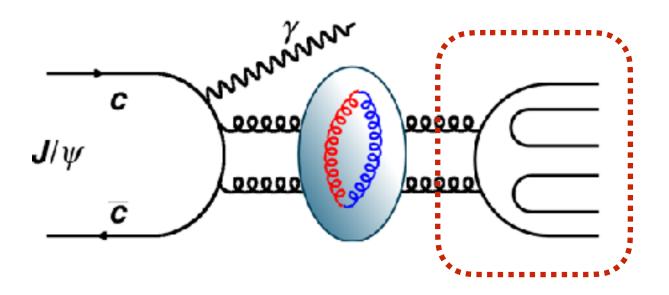
Flavor symmetric decays

No rigorous predictions on decay patterns and their branching ratios

 \bullet e.g. the 0⁻⁺ glueball could have similar decays of η_c



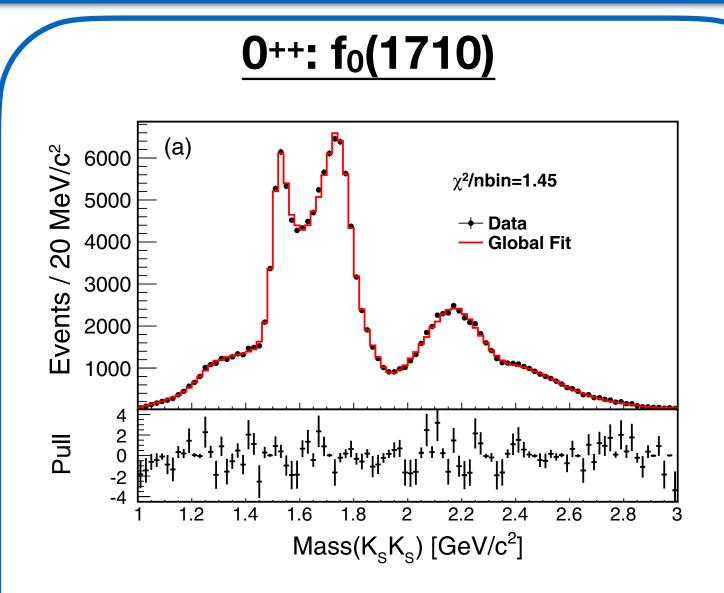
The glueball decays could be the analogy to Charmonium decays since they all decay via gluons (OZI suppression) [PLB 380 189(1996), Commu. Theor. Phys. 23.373 (1995)]







Historical Glueball Candidates



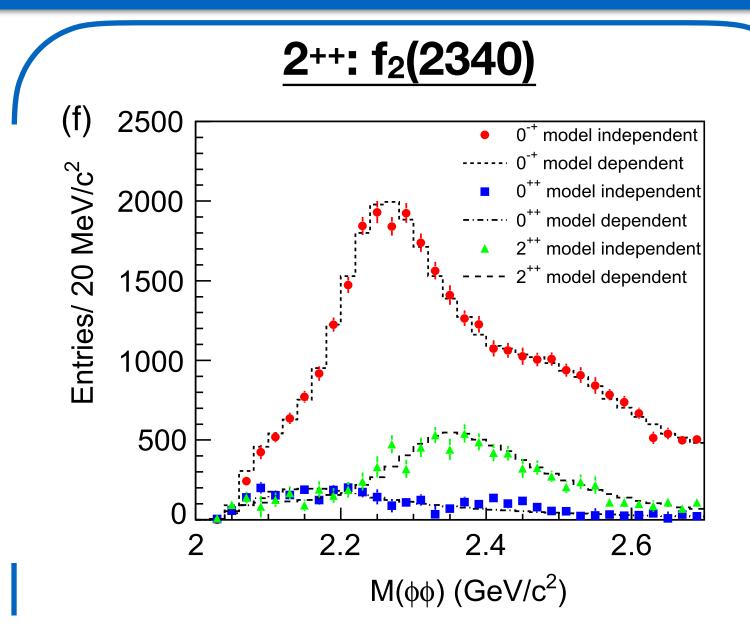
High production rate in $J/\psi \rightarrow \gamma f_0(1710)$:

 $B[J/\psi \to \gamma f_0(1710) \to \gamma \pi \pi] = (4.0 \pm 1.0) \times 10^{-4}$ $B[J/\psi \rightarrow \gamma f_0(1710) \rightarrow \gamma K_s^0 K_s^0] = (2.00^{+0.03}_{-0.02} + 0.31) \times 10^{-4}$

Decay suppression in f_0(1710) \rightarrow \eta \eta':

 $B[f_0(1710) \rightarrow \eta \eta' / f_0(1710) \rightarrow \pi \pi] < (2.9 \pm ^{+1.1}_{-0.9}) \times 10^{-3}$

$f_0(1710)$ favors to be a scalar glueball or large glueball content (mixing mechanism)

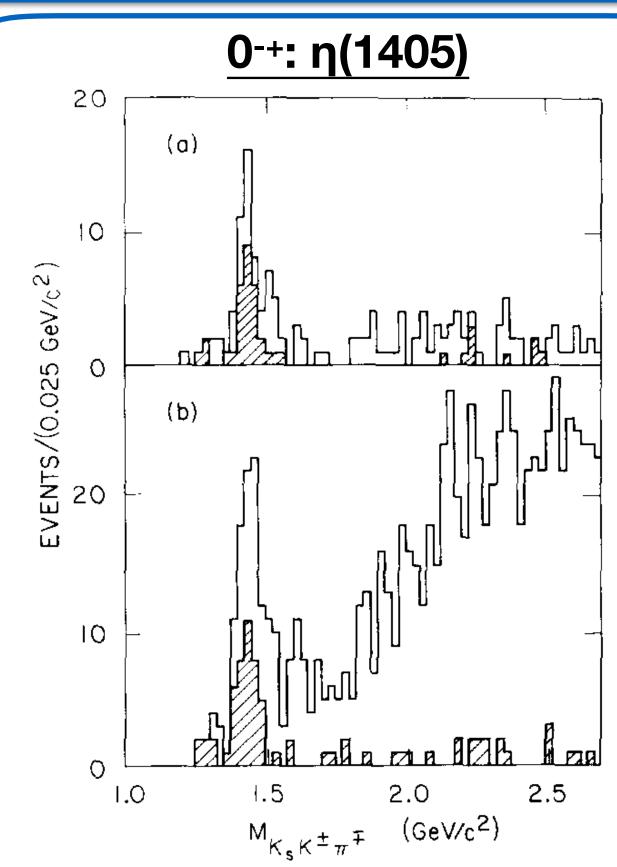


B(J/ψ→γf₂(2340)→γφφ) =(1.91 ± 0.14^{+0.72}_{-0.73} × 10⁻⁴ $B(J/\psi \rightarrow \gamma f_2(2340) \rightarrow \gamma K_s K_s) = (5.54^{+0.34}_{-0.40} + 3.82 \times 10^{-5})$ **B(J/ψ→**γ**f**₂(2340)→γη'η') =(8.67 ± 0.70^{+0.16}_{-1.67} × 10⁻⁶)

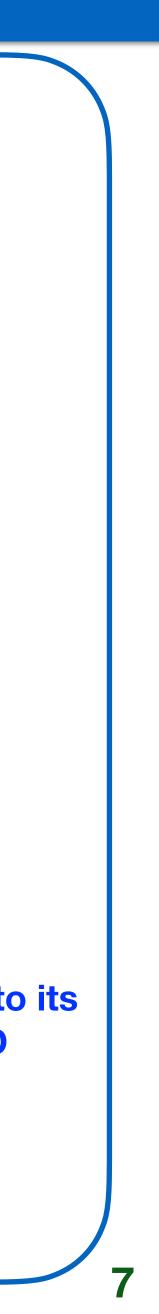
Production rate lower than LQCD prediction

Large overlaps between many broad tensor meson cause the its measurement difficulty

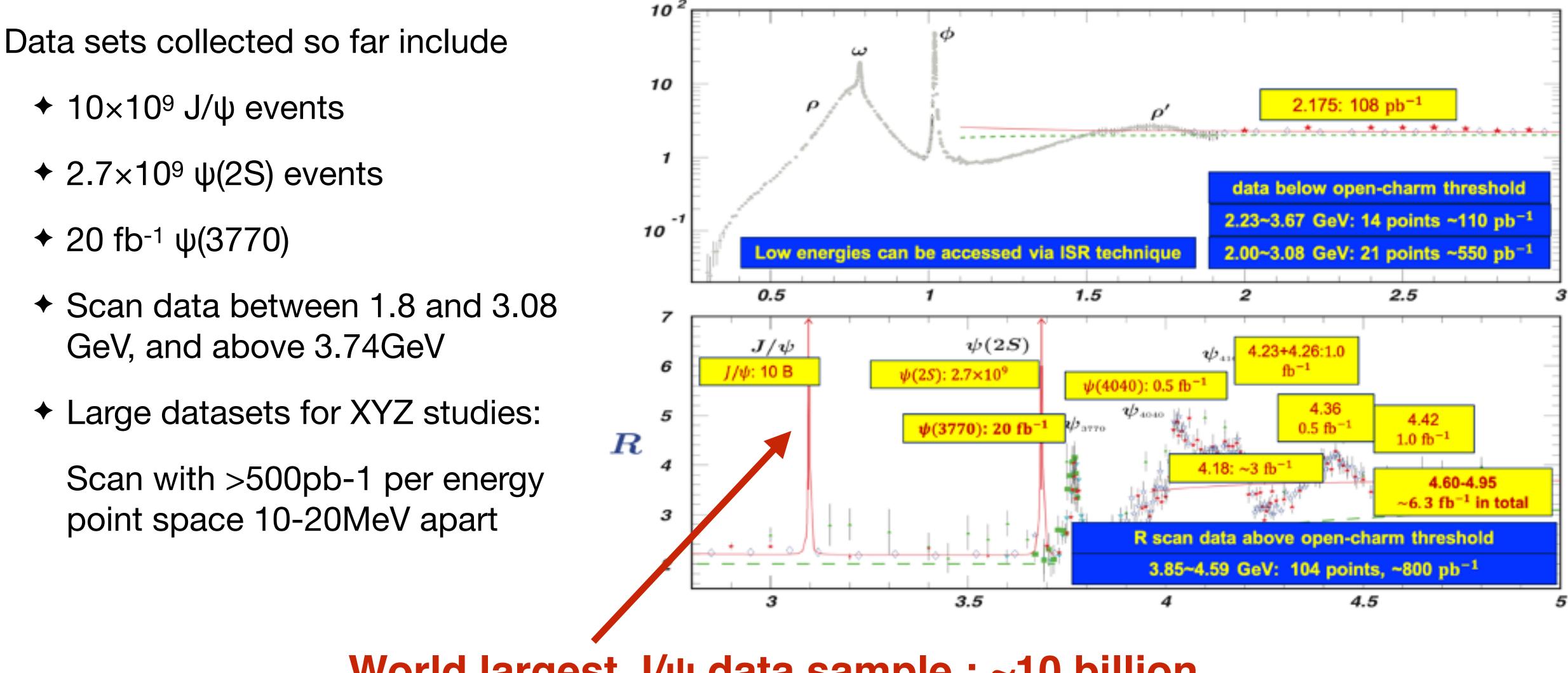
High production rate in $J/\psi \rightarrow \gamma f_2(2340)$: $B(J/\psi \rightarrow \gamma f_2(2340) \rightarrow \gamma \eta \eta) = (3.8^{+0.62}_{-0.66} + 2.37) \times 10^{-5}$



No longer a 0⁻⁺ glueball candidate due to its large mass difference from latest LQCD prediction



BESIII Data samples



Totally about 50fb⁻¹ integrated luminosity

World largest J/ψ data sample : ~10 billion

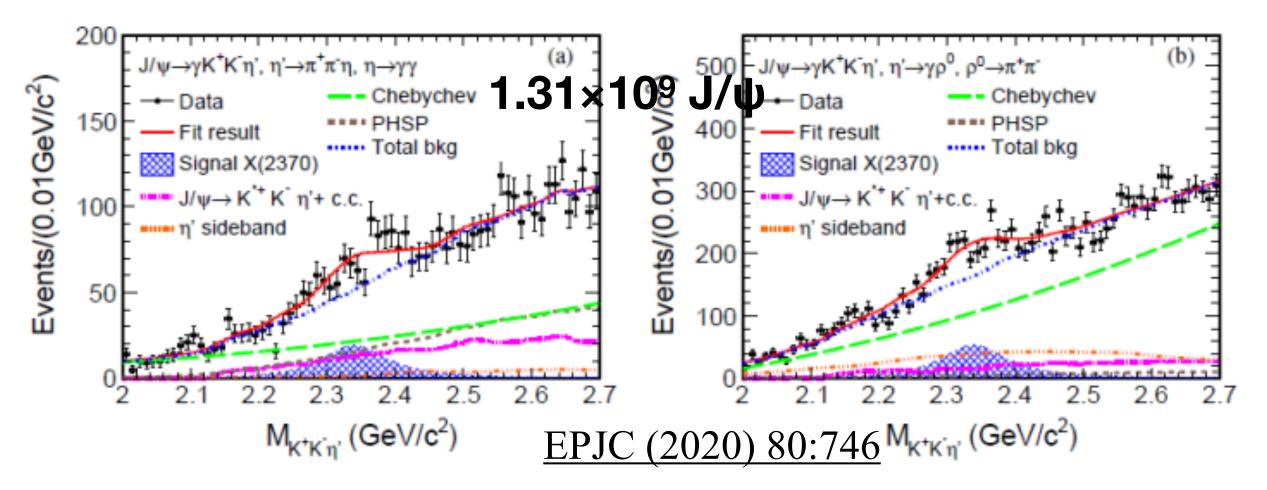




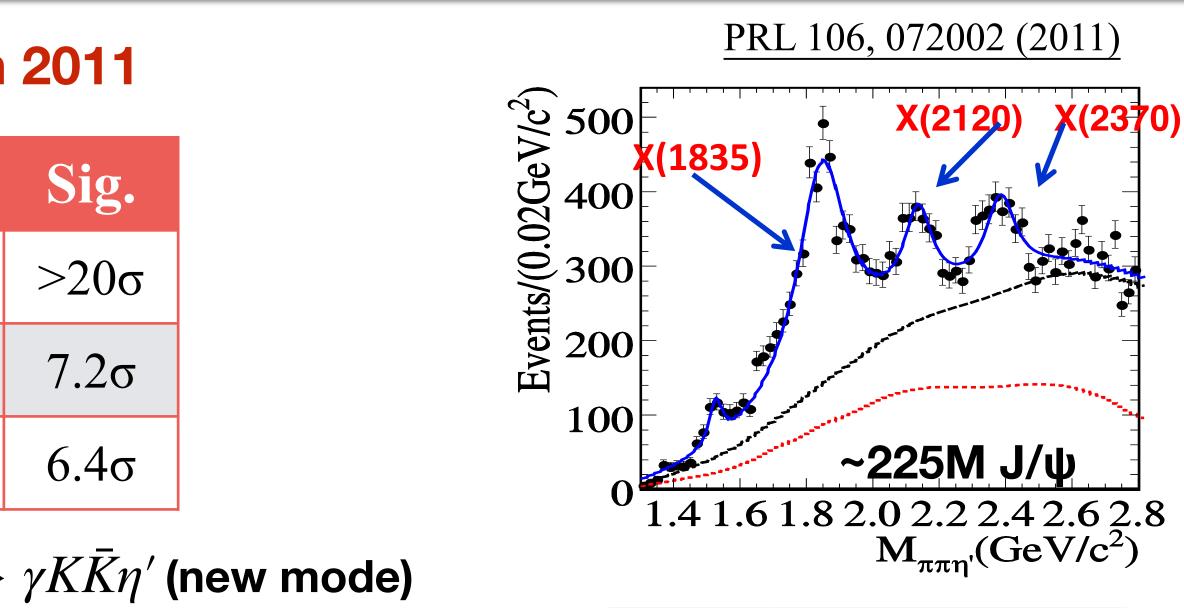
Discovered by BESIII in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ in 2011

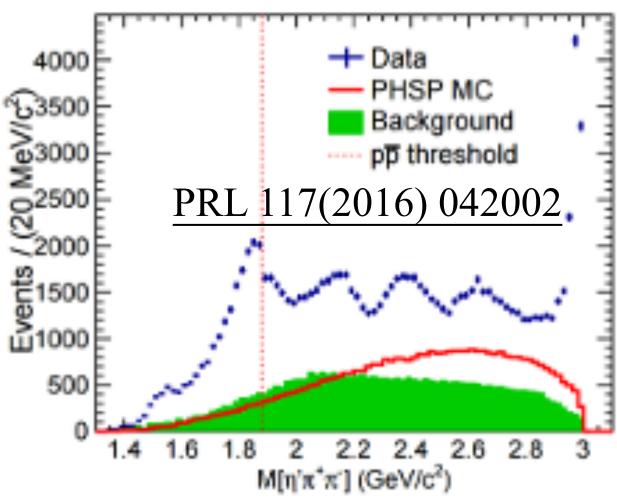
	M(MeV/c ²)	$\Gamma(MeV/c^2)$	
X(1835)	1836.5±3.0+5.6-2.1	190.1±9.0+38-36	-
X(2120)	2122.4±6.7 ^{+4.7} -2.7	$83 \pm 16^{+31}$ -11	
X(2370)	$2376.3 \pm 8.7^{+3.2}_{-4.3}$	83±17+44-6	

Confirmed by BESIII in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ and $J/\psi \rightarrow \gamma K \bar{K} \eta'$ (new mode) ۲

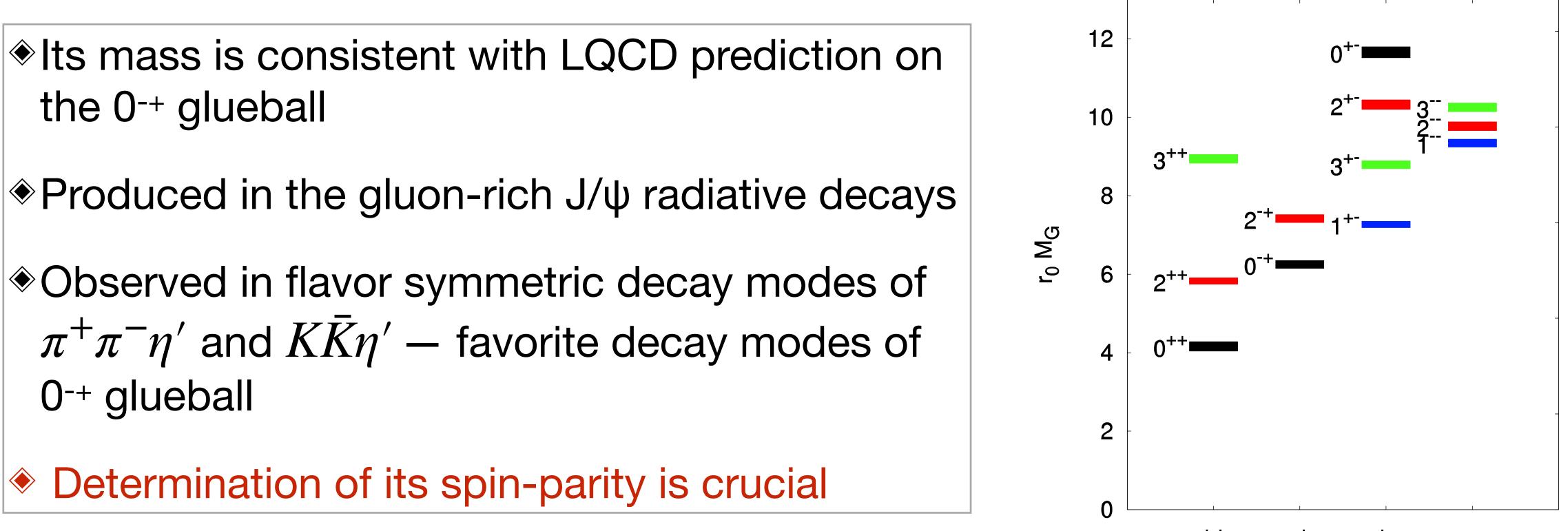


X(2370)







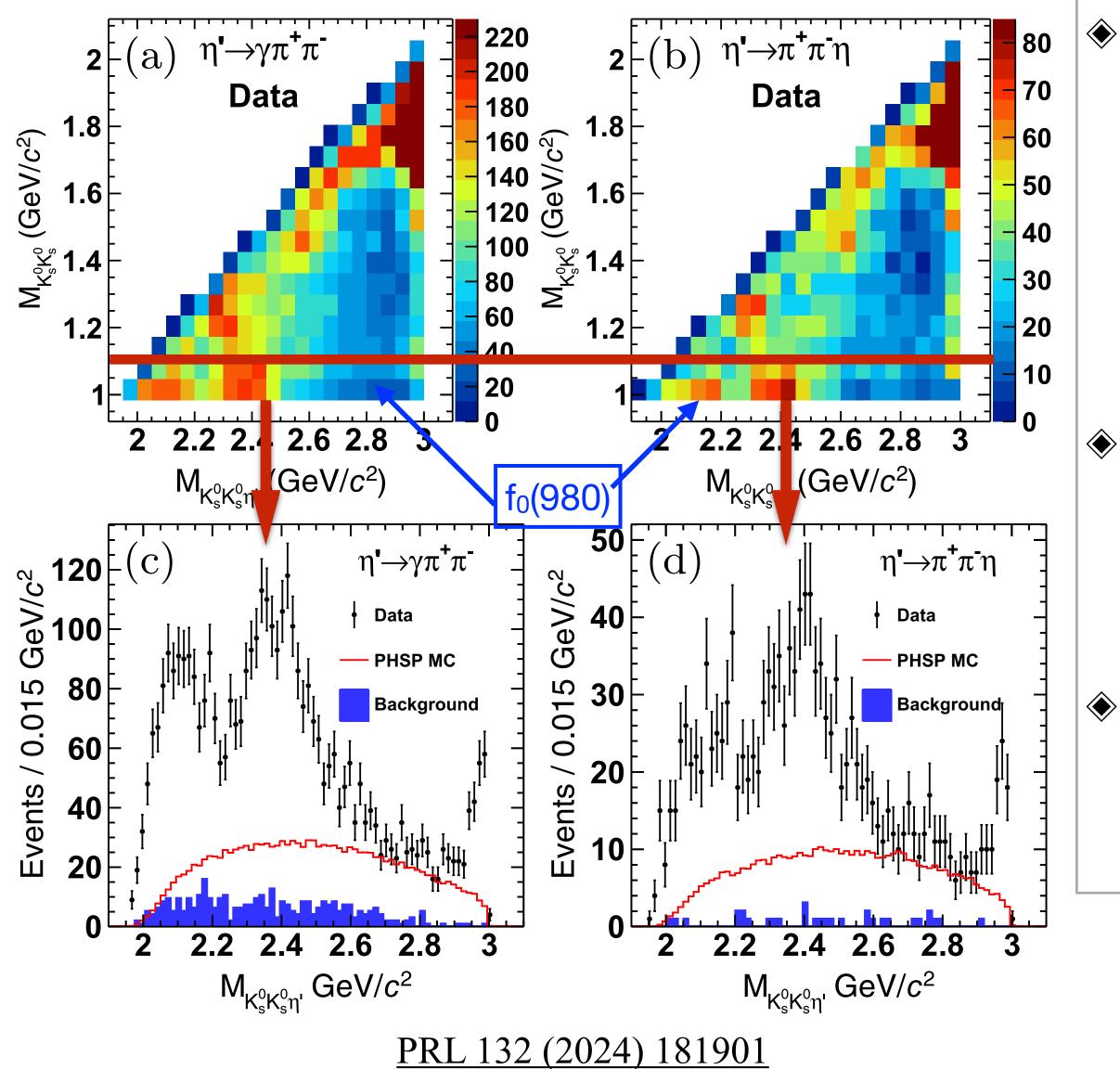


X(2370) - good candidate of 0⁻⁺ glueball



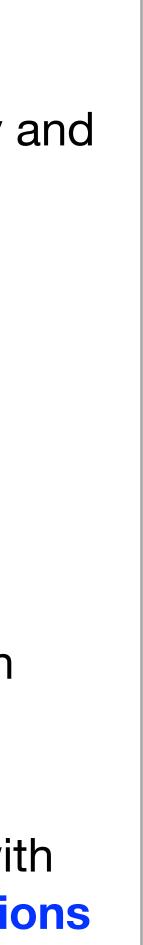


Spin-Parity determination of the X(2370) in $J/\psi \rightarrow \gamma K^0_s K^0_s \eta^2$



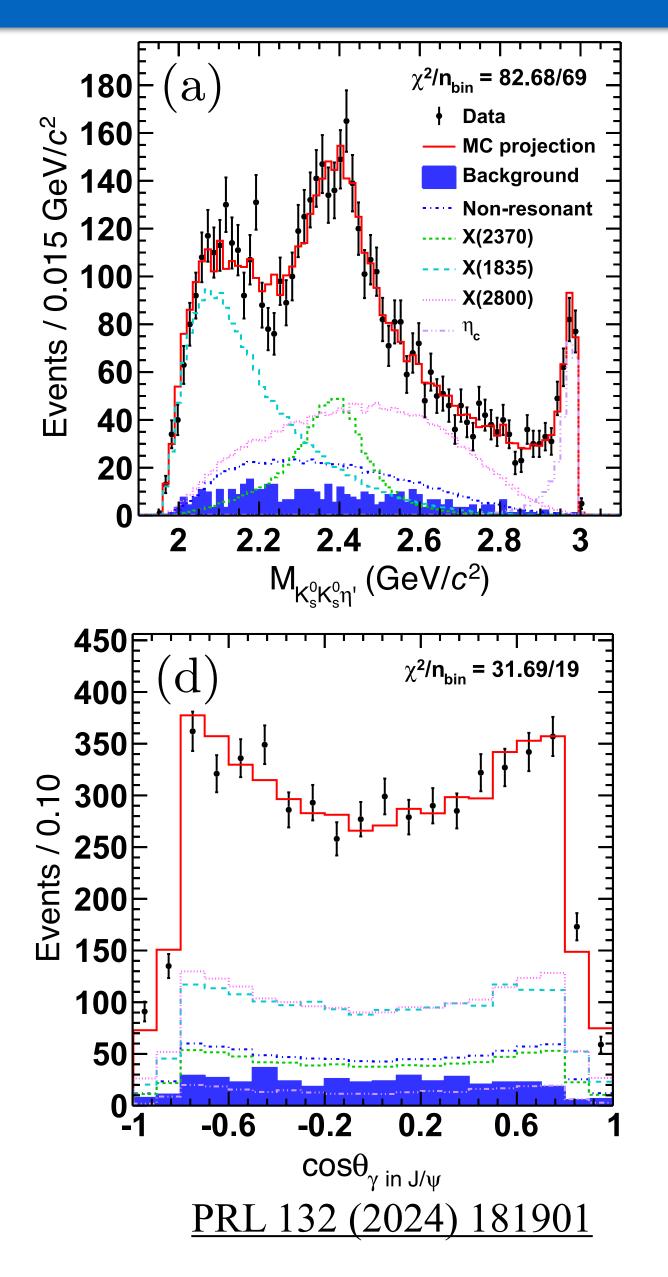
- Analysis advantage of $J/\psi \rightarrow \gamma K^{0}_{s}K^{0}_{s}\eta'$:
- Almost background free channel (exchange symmetry and C-parity conservation)
- + 10billion J/ψ data
- Very good BESIII detector performance
- Similar structures in $\eta' \rightarrow \pi^+\pi^-\eta / \gamma\pi^+\pi^-$ modes:
 - Evident f₀(980) in K⁰_sK⁰_s mass threshold
 - + Clear signal of X(1835), X(2370), η_c with f₀(980) selection
 - Best PWA fit can well describe the data:
 - + Spin-parity of the X(2370) is determined to be 0-+ with significance larger than 9.8σ w.r.t. other J^{pc} assumptions







Spin-Parity determination of the X(2370) in $J/\psi \rightarrow \gamma K^0_s K^0_s \eta^2$



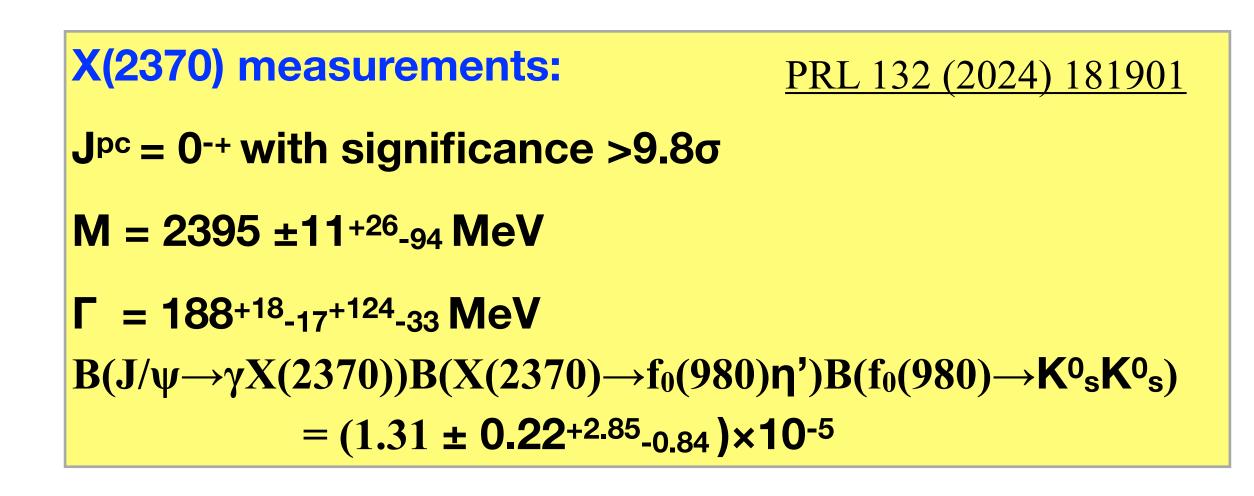
- Analysis advantage of $J/\psi \rightarrow \gamma K^{0}_{s}K^{0}_{s}\eta'$:
 - Almost background free channel
 - + 10billion J/ψ data
 - Very good BESIII detector performance
- Similar structures in $\eta' \rightarrow \pi^+\pi^-\eta / \gamma\pi^+\pi^-$ modes:
 - Evident f₀(980) in K⁰_sK⁰_s mass threshold
 - Clear signal of X(1835), X(2370), η_c with f₀(980) selection
- Best PWA fit can well describe the data:
 - Spin-parity of the X(2370) is determined to be 0⁻⁺ with significance larger than 9.8 w.r.t. other J^{pc} assumptions



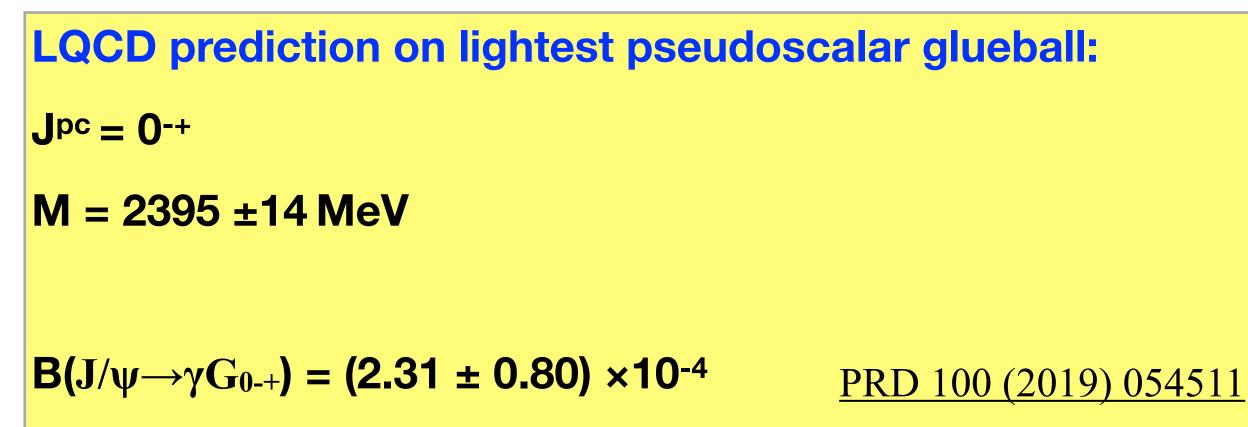




Compared with LQCD prediction on Lightest 0-+ Glueball



- - + The spin-parity of the X(2370) is determined to be 0⁻⁺ for the first time
 - Mass is in a good agreement with LQCD predictions
 - (assuming ~5% decay rate, $B(J/\psi \rightarrow \gamma X(2370)) = (10.7^{+22.8} 7) \times 10^{-4})$



The measurements are in a good agreement with the predictions on lightest pseudoscalar glueball

+ The estimation on B(J/ $\psi \rightarrow \gamma X(2370)$) and prediction on B(J/ $\psi \rightarrow \gamma G_{0-+}$) are consistent within errors







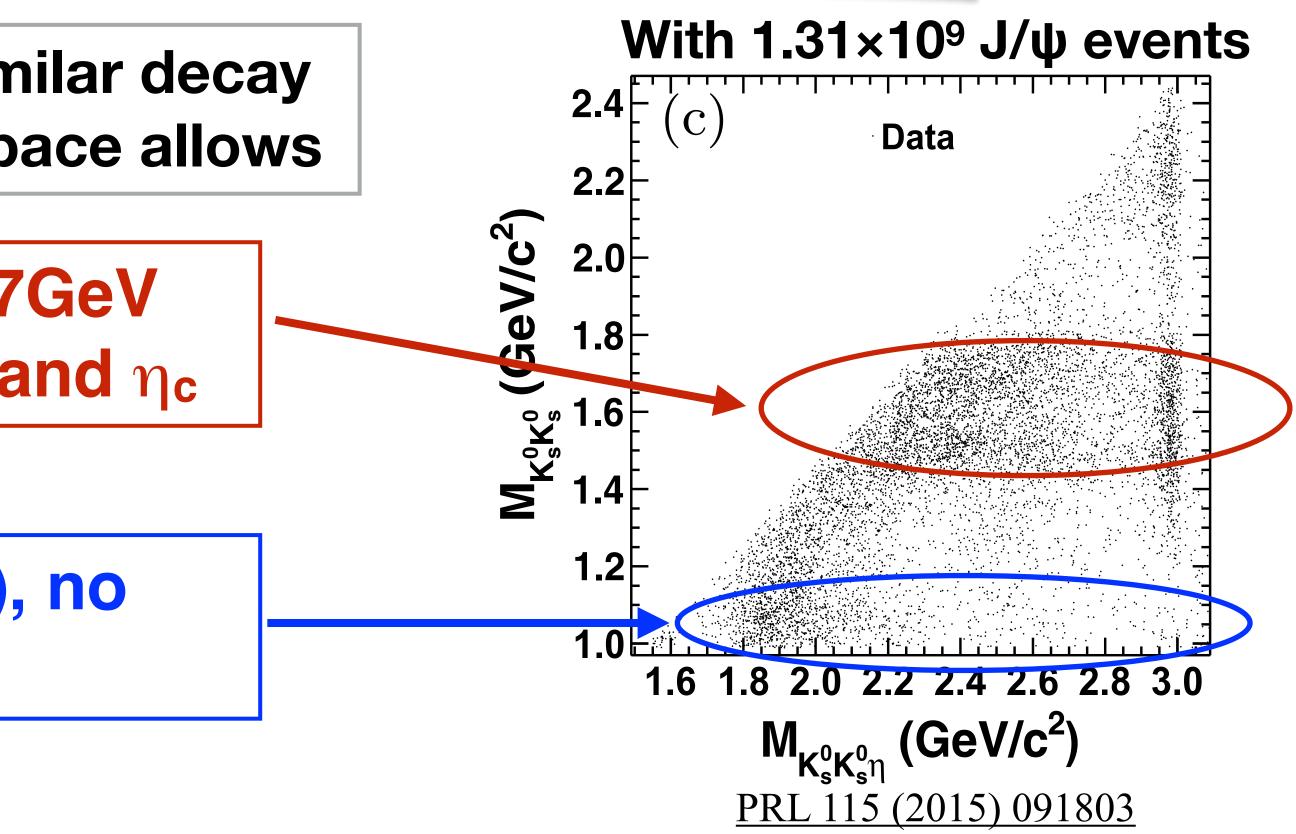
Observation and Spin-Parity Determination of the X(1835) in $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta$

Qualitatively, we can clearly observe: similar decay patterns of the X(2370) and η_c if phase space allows

In the upper KK mass band of 1.5-1.7GeV range, clear signals of both X(2370) and η_c

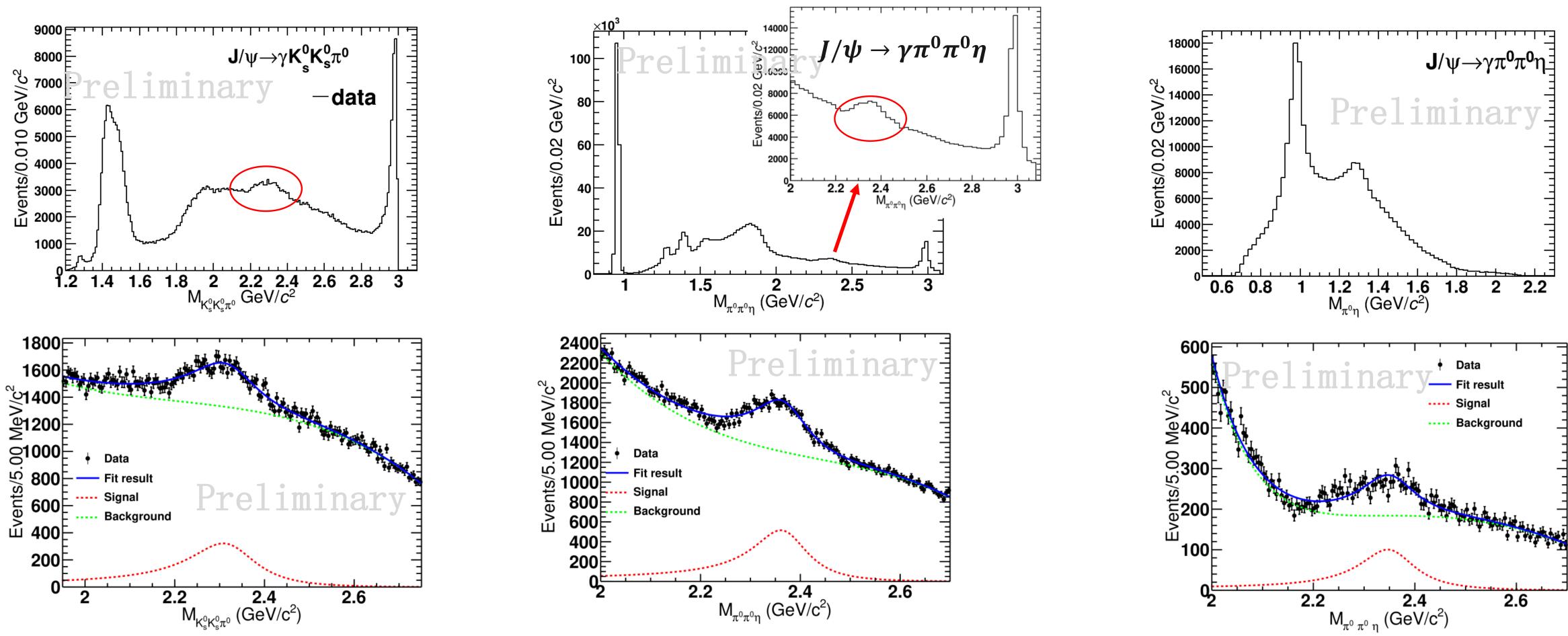
In the lower KK mass band of f₀(980), no **X(2370), nor** η_c

X(2370) in $J/\psi \rightarrow \gamma K^0_s K^0_s \eta$





Observation of new decay modes of the X(2370)



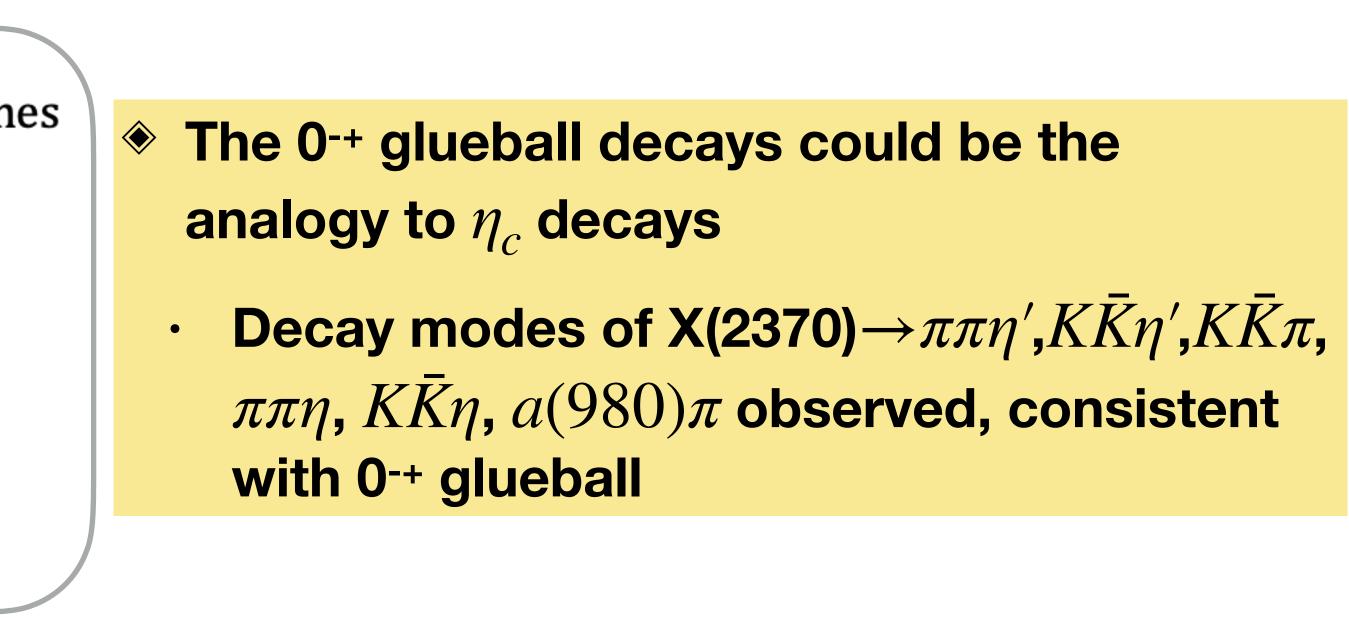
First observation of $X(2370) \rightarrow K_s^0 K_s^0 \pi^0$, $X(2370) \rightarrow \pi^0 \pi^0 \eta$ and $X(2370) \rightarrow a(980)\eta$ with significances >> 5σ and accompanied with η_c







		5 ma	jor η_c decay modes (from PDG)
(-	- 5		modes in 0 ⁻⁺ glueball traditional search
			Decays involving hadronic resonances
	Γ_1	$\eta'(958) \pi \pi$ $\eta'(958) K \overline{k}$	(1.87±0.26) %
	Γ2	$\eta'(958) K\overline{P}$	\overline{K} (1.61±0.25)%
			Decays into stable hadrons
		$K\overline{K}\pi$	$(7.0 \pm 0.4)\%$
		$K\overline{K}\eta$	$(1.32\pm0.15)\%$
	Г ₃₆	$\eta \pi^+ \pi^-$	(1.7 ± 0.5) %



Such high similarity between the X(2370) and η_c decay modes strongly supports the glueball interpretation of the X(2370)



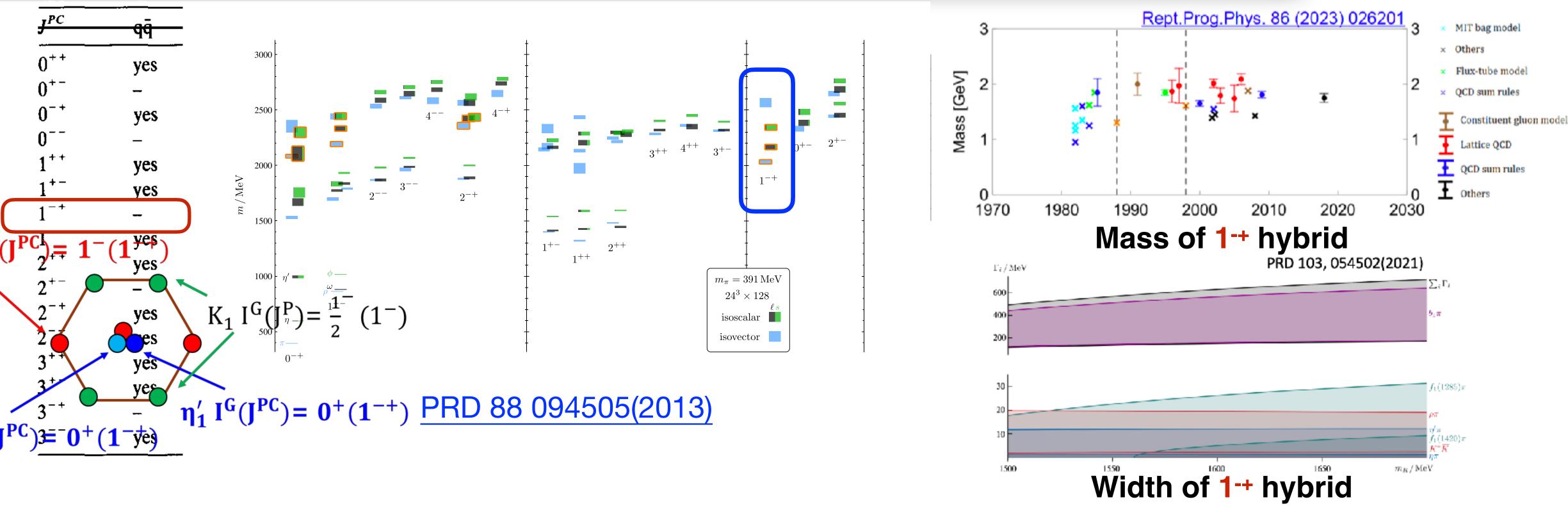
- \bullet With 10 billion J/ψ data, we should be able to answer this key question Either we find them or exclude them
 - In the above 2.3 GeV mass region as LQCD 0-+ glueball prediction
 - X(2370) is the unique 0-+ particle produced in these "5 golden modes" and in J/ψ radiative decays
- \Rightarrow 10 billion J/ψ data make us face a situation: Either we finally identify X(2370) as 0-+ glueball, or LQCD may face a big challenge in the glueball predictions

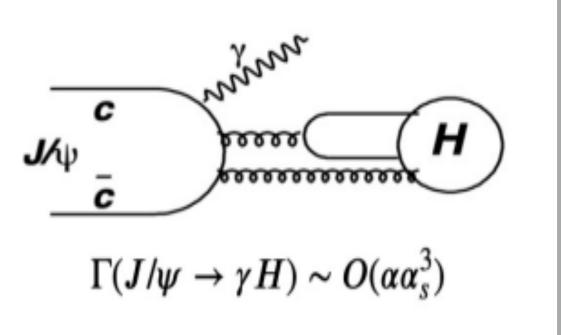






Exotic 1-+ state





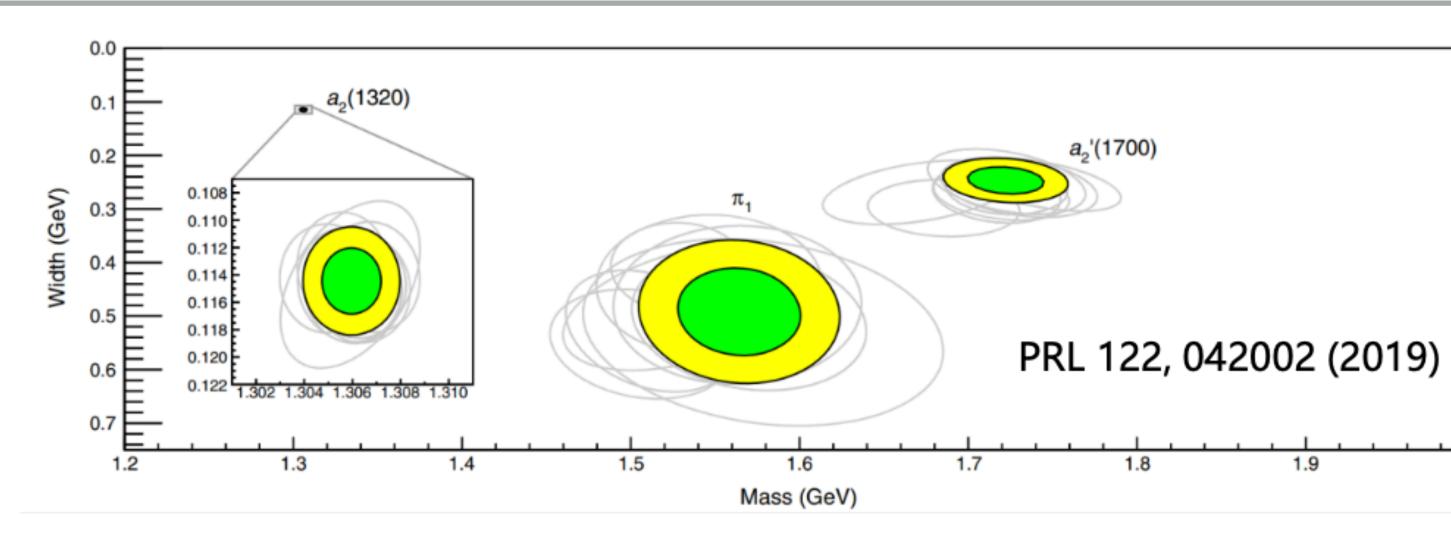
Spin-exotic state of 1⁻⁺: forbidden in conventional quark model Exotic state 1-+ provide an unique way for hybrid search: LQCD predicts the lightest nonet of 1-+ hybrids: 1.7 - 2.1GeV Can be produced in the gluon-rich charmonium decays



18



- Over 3 decades, experimental evidence for 3 candidates with 1-+ state:
 - ✦ All 1⁻⁺ iso-vectors
 - $\pi_1(1400)$: seen in $\eta\pi$
 - $\pi_1(1600)$: seen in $\rho \pi$, $\eta' \pi$, $b_1 \pi$, $f_1 \pi$
 - + $\pi_1(2015)$: seen in $b_1\pi$ and $f_1\pi$
- Some claims are controversial
- $\pi_1(1400)$ and $\pi_1(1600)$ can be one pole

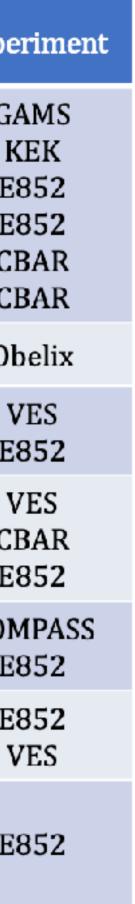


Spin-exotic mesons

1.9

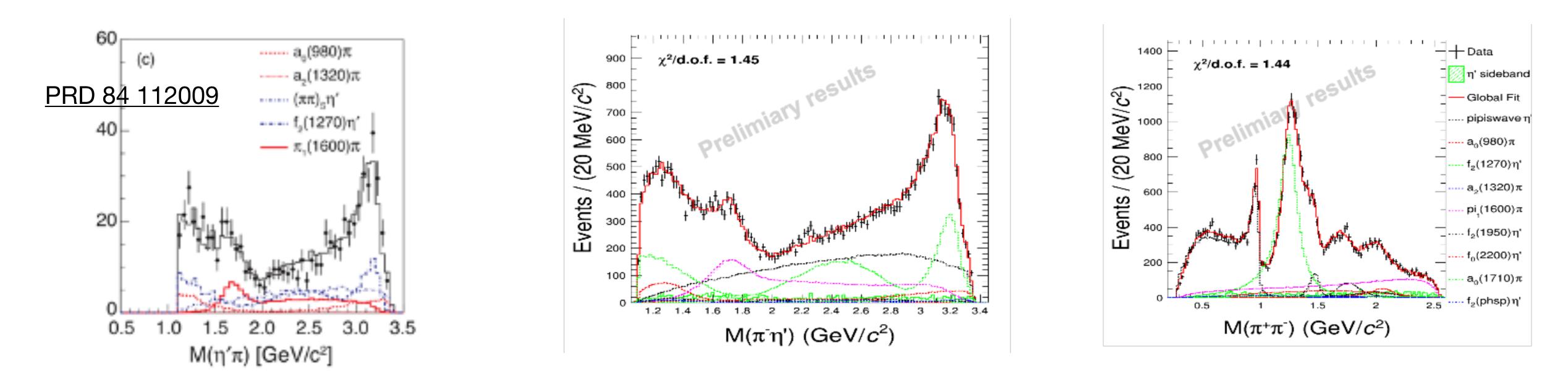
2.0

ates with 1-+ st	ate:				
			Decay mode	Reaction	Exper
		π ₁ (1400)	ηπ	$\pi^{-}p \rightarrow \pi^{-}\eta p$ $\pi^{-}p \rightarrow \pi^{0}\eta n$ $\pi^{-}p \rightarrow \pi^{-}\eta p$ $\pi^{-}p \rightarrow \pi^{0}\eta n$ $\bar{p}n \rightarrow \pi^{-}\pi^{0}\eta$ $\bar{p}p \rightarrow \pi^{0}\pi^{0}\eta$	GA KI E8 E8 CB CB
		ρ	$ ho\pi$	$ar{p}p ightarrow 2\pi^+ 2\pi^-$	Ob
			η'π	$\pi^{-}Be \rightarrow \eta' \pi^{-} \pi^{0}Be$ $\pi^{-}p \rightarrow \pi^{-} \eta' p$	V E8
	π ₁ (1600)	$b_1\pi$	$\pi^{-}Be ightarrow \omega\pi^{-}\pi^{0}Be$ $\bar{p}p ightarrow \omega\pi^{+}\pi^{-}\pi^{0}$ $\pi^{-}p ightarrow \omega\pi^{-}\pi^{0}p$	V) CB E8	
			ρπ	$\pi^{-}Pb \rightarrow \pi^{+}\pi^{-}\pi^{-}X$ $\pi^{-}p \rightarrow \pi^{+}\pi^{-}\pi^{-}p$	COM E8
			$f_1\pi$	$\pi^- p ightarrow p\eta \pi^+ \pi^- \pi^- \pi^- \pi^- A$ $\pi^- A ightarrow \eta \pi^+ \pi^- \pi^- A$	E8 V
	$f_1\pi$	$\pi^- p \rightarrow \omega \pi^- \pi^0 p$	E		
		π ₁ (2015)	$f_1\pi$ $b_1\pi$	$\pi^- p \to p \eta \pi^+ \pi^- \pi^-$	E8
042002 (2019)					





Observation of Exotic 1⁻⁺ Isovector state $\pi(1600)$



- PWA in $\psi' \to \gamma \chi_{c1}(\chi_{c1} \to \pi^+ \pi^- \eta')$ with higher ψ' data sample @ BESIII:

+ First observation of Exotic 1⁻⁺ Isovector state $\pi(1600)$ with a significance >10 σ better than other J^{PC} assumption

+ The significance of phase motion is also greater than 10σ

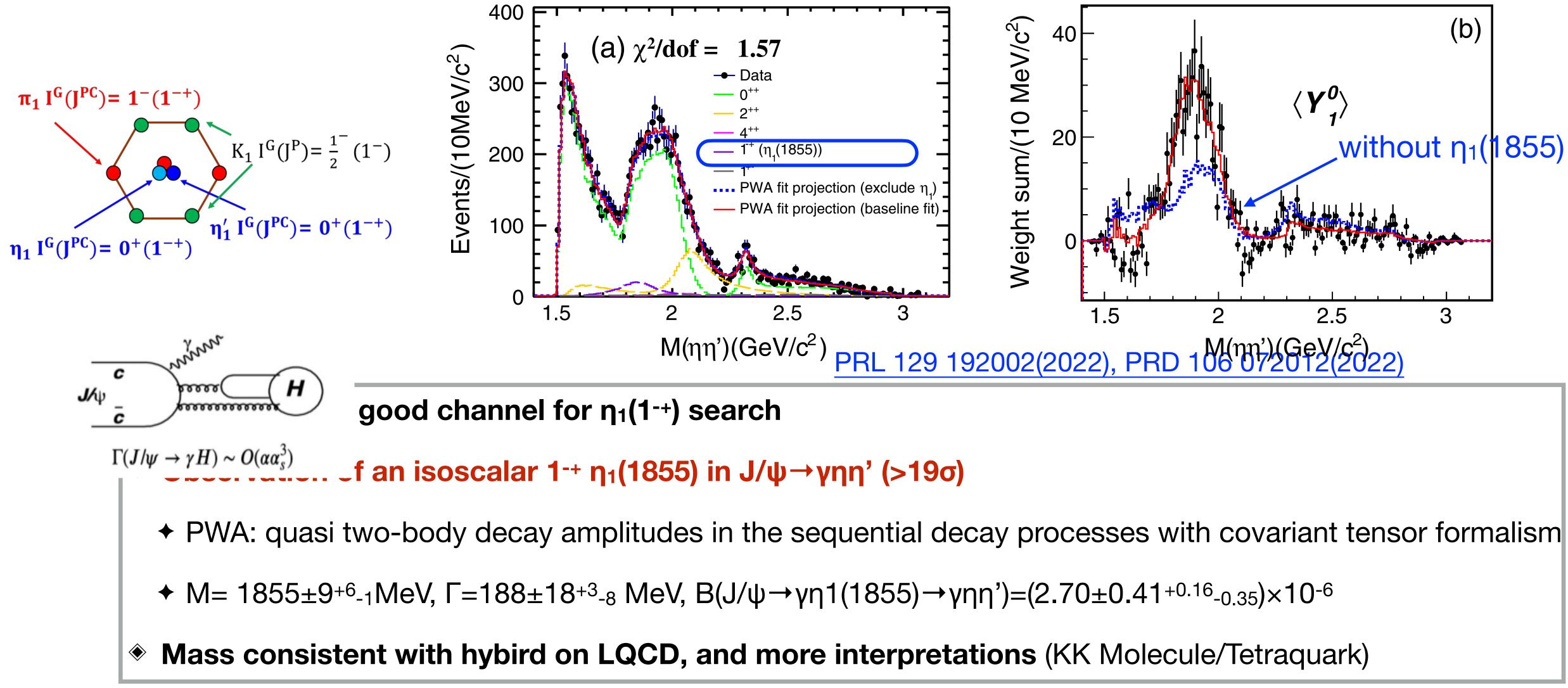
• CLEO-c results: evidence of an exotic P-wave $\eta'\pi$ amplitude with 4σ and but no significant phase motion





Observation of An Exotic 1-+ Isoscalar $\eta_1(1855)$

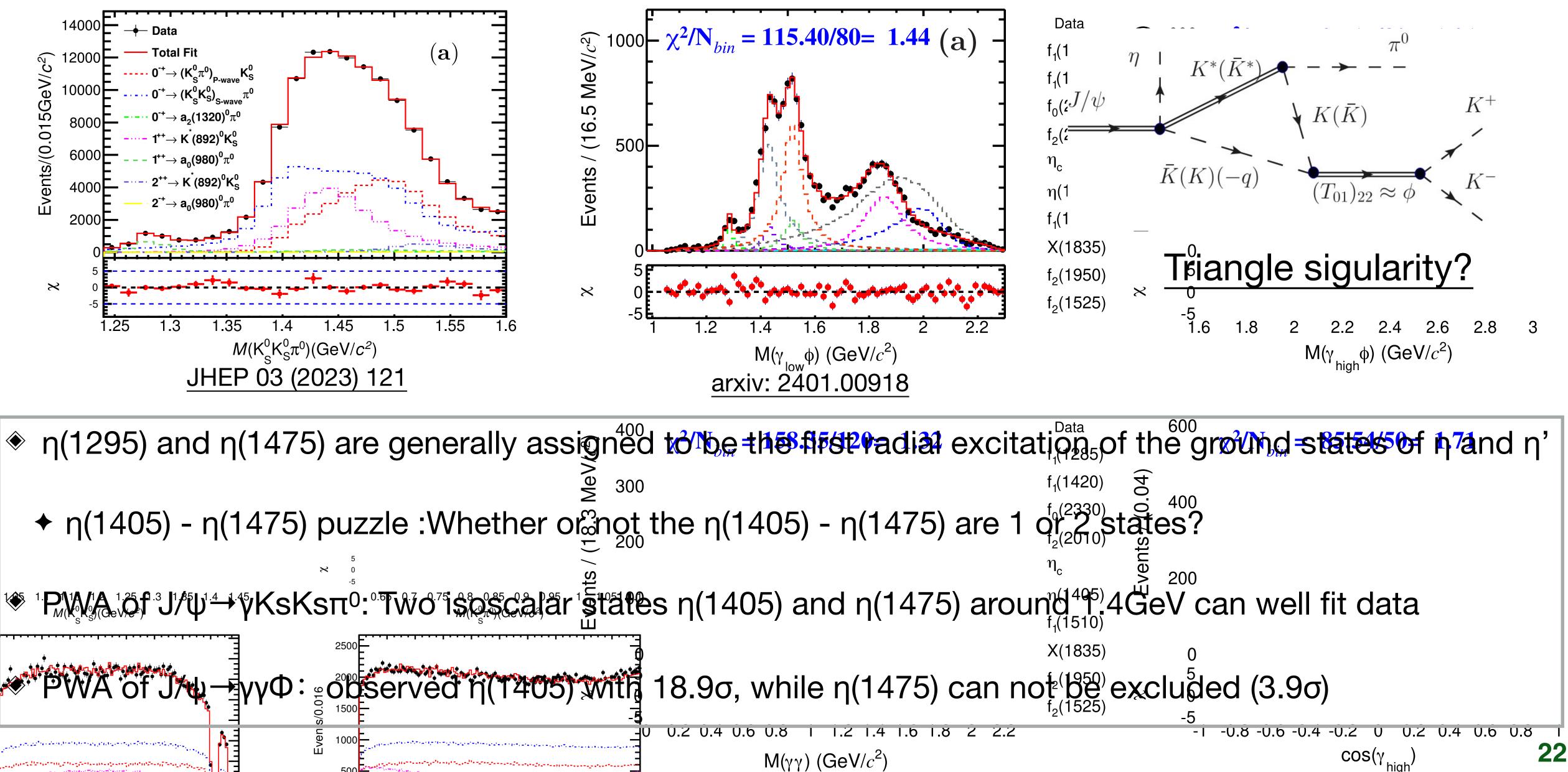
Isoscalar (1-+) is critical to establish the nonet hybrid multiplet: partners for the Isovector (1-+)



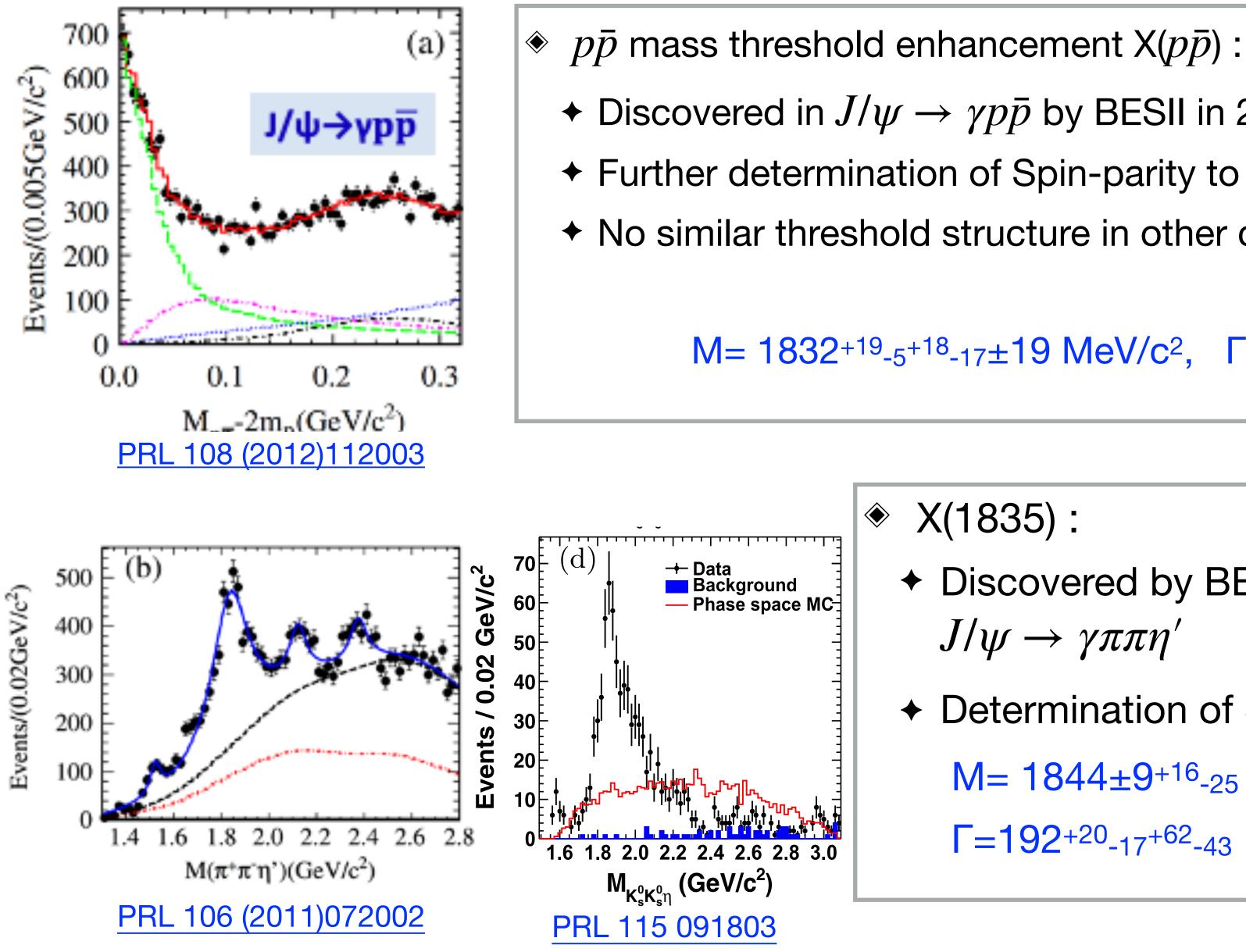




Shed new lights on $\eta(1405)/\eta(1475)$ puzzle



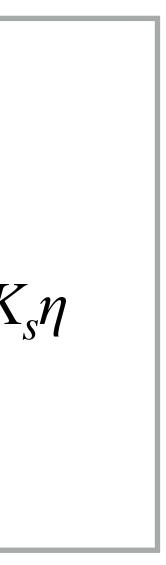
Observation of $X(p\bar{p})$ **and** X(1835)



- + Discovered in $J/\psi \rightarrow \gamma p \bar{p}$ by BESII in 2003 and confirmed by BESIII and CLEO-c Further determination of Spin-parity to be 0++
- + No similar threshold structure in other channels \rightarrow It can not be pure FSI effect
 - $M = 1832^{+19}_{-5}^{+18}_{-17} \pm 19 \text{ MeV/c}^2$, $\Gamma = 13 \pm 19 \text{ MeV/c}^2$ (<76 MeV/c²@90% C.L.)

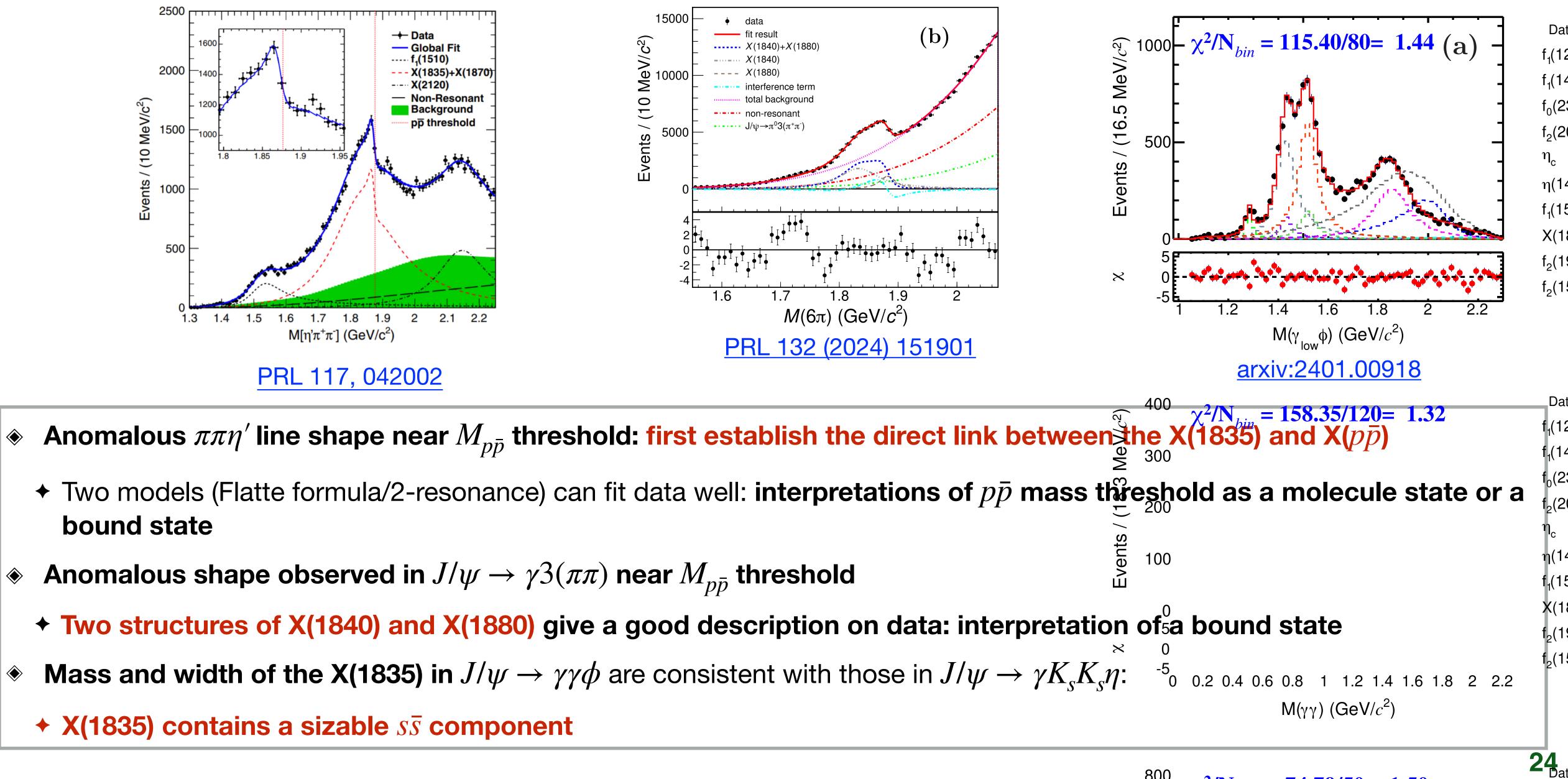
- ♦ X(1835) :
 - Discovered by BESII and confirmed by BESIII in $J/\psi \rightarrow \gamma \pi \pi \eta'$
 - + Determination of Spin-parity to be 0⁻⁺ in $J/\psi \rightarrow \gamma K_s K_s \eta$
 - $M = 1844 \pm 9^{+16} 25 MeV/c^{2}$
 - $\Gamma = 192^{+20}_{-17}^{+62}_{-43} \text{ MeV/c}^2$



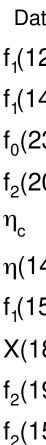




$M(6\pi)$ (GeV/ c^2) Direct link between the $X(p\bar{p})$ and X(1835)



- - bound state
- ۲
- ۲
 - + X(1835) contains a sizable $s\bar{s}$ component





- A set of interesting and important results from the light hadron spectroscopy achieved: •
- **Discovery of a glueball-like particle: X(2370)**
 - + Strong correlation between the X(1835) and $M_{p\bar{p}}$ threshold enhancement. A molecule state or a bound state?
 - + Observation of An Exotic 1⁻⁺ Isoscalar state $\eta_1(1855)$ and Isovector state $\pi(1600)$

+ ...

With the more data, the more extensive and intensive investigation are ongoing, looking forward to new results in the near future.



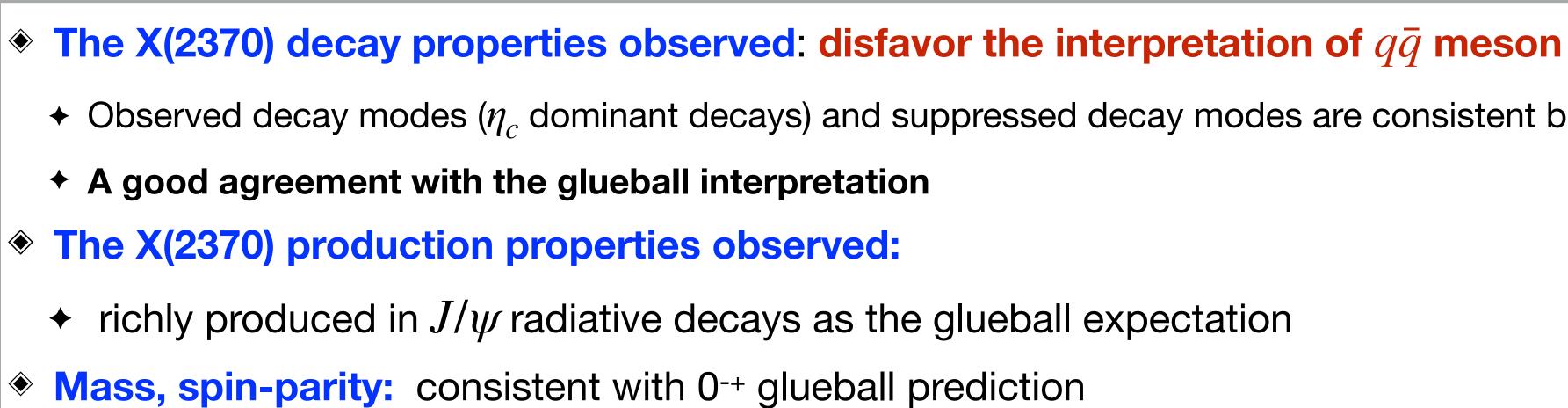






Interpretation

	X(2370)	η	
f ₀ (980)η'	\checkmark	\checkmark	Γ
f₀(980)η	Suppressed	Suppressed	
f ₀ (1500)η	\checkmark	\checkmark	



In the mass region larger than 2GeV, the only particle X(2370) for the 0⁻⁺ glueball candidate in J/ψ radiative decays and five golden decay modes $(\pi\pi\eta', K\bar{K}\eta', K\bar{K}\pi, \pi\pi\eta, K\bar{K}\eta)$

Interpertation on the X(2370)

Disfavors $q\bar{q}$ meson with pure $u\bar{u}/d\bar{d}$ component

Disfavors $q\bar{q}$ meson with pure $s\bar{s}$ component

Disfavors $q\bar{q}$ meson with pure $s\bar{s}$ component

+ Observed decay modes (η_c dominant decays) and suppressed decay modes are consistent between the X(2370) and η_c







