



河南师范大学
HENAN NORMAL UNIVERSITY

BESIII

BESIII上粲物理的研究

李惠静

河南师范大学

(On behalf of BESIII Collaboration)

2024年11月2日

第一届基础物理研讨会暨基础物理平台年会

10月31日-11月2日，郑州.



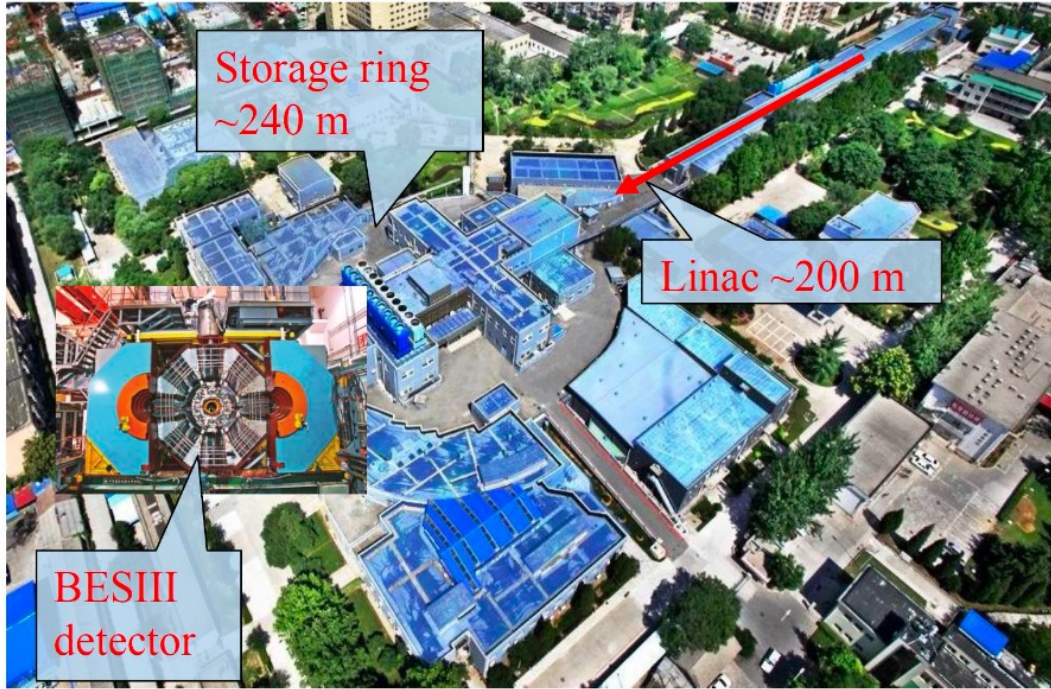
Outline

✍ **Hadronic decays**

✍ **(Semi-) leptonic decays**

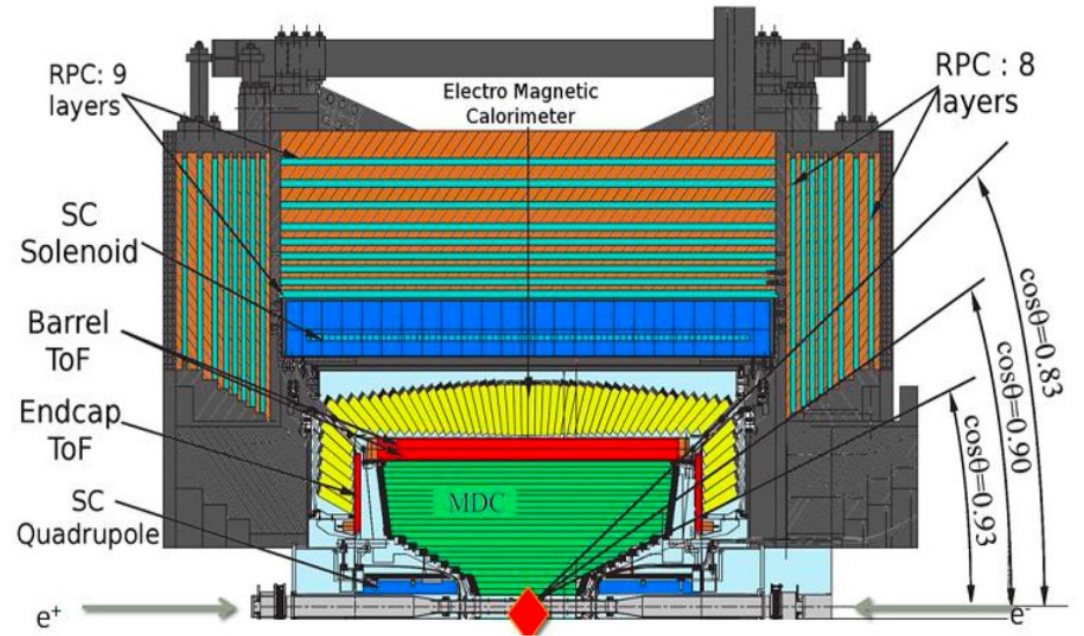
✍ **Rare decays**

Beijing Electron Positron Collider(BEPCII)



Double-ring, symmetry, multi-bunch $e^+ e^-$ collider
 $E_{cm} = 1.84$ to 4.95 GeV
 Energy spread: $\Delta E \approx 5 \times 10^{-4}$
 Peak luminosity in continuously operation @ $E_{cm} = 3.77$ GeV: $1.1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

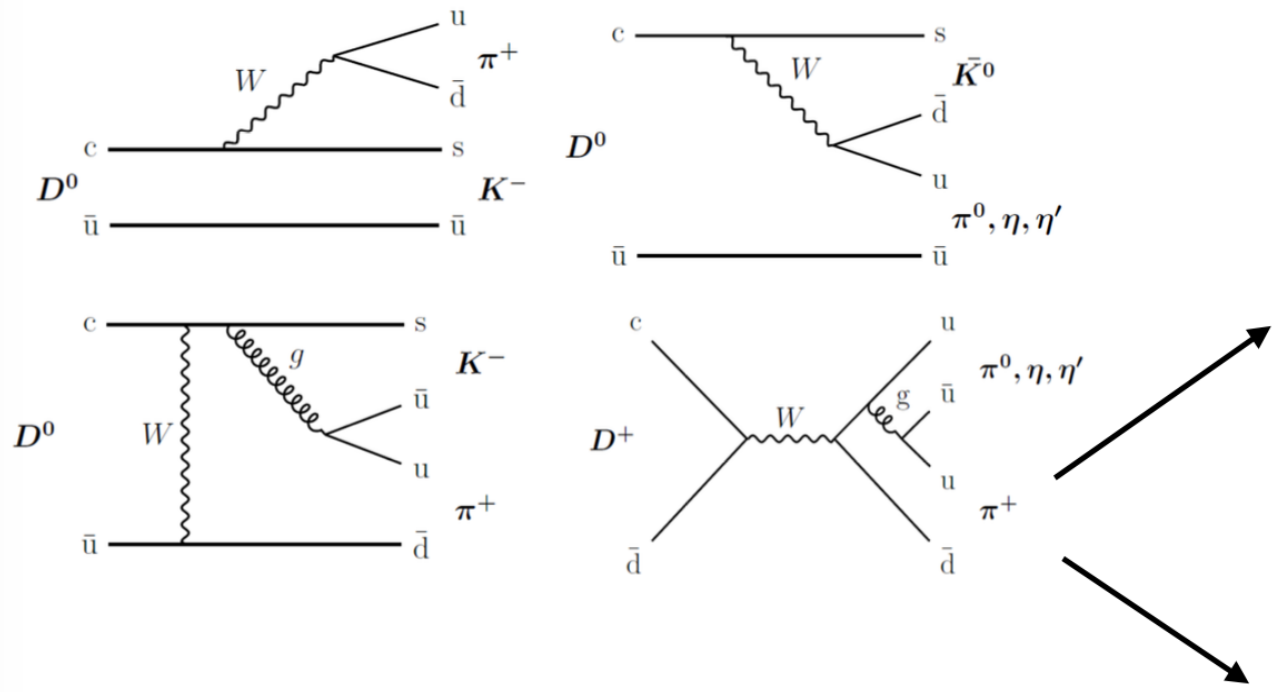
Beijing Spectrometer(BESIII)



Main Drift Chamber	Time Of Flight	Electromagnetic Calorimeter	Muon Counter
Small cell, 43 layer	Plastic scintillator	CsI(Tl): L=28 cm	RPC
$\sigma_{xy}=130 \mu\text{m}$	$\sigma_T(\text{barrel}): 65 \text{ ps}$	Barrel $\sigma_E=2.5\%$	Barrel: 9 layers
$dE/dx \sim 6\%$	$\sigma_T(\text{endcap}): 110 \text{ ps}$	Endcap $\sigma_E=5.0\%$	Endcap: 8 layers
$\sigma_p/p = 0.5\%$ at 1 GeV	(update to 60 ps with MRPC)		$\sigma_{\text{spatial}}: 1.48 \text{ cm}$

Why charm meson physics?

Hadronic decay



- ✓ Probe non-perturbative QCD
- ✓ Help to understand hadron spectroscopy
- ✓ Probe the weak decay mechanisms in DCS decays

Amplitude analysis

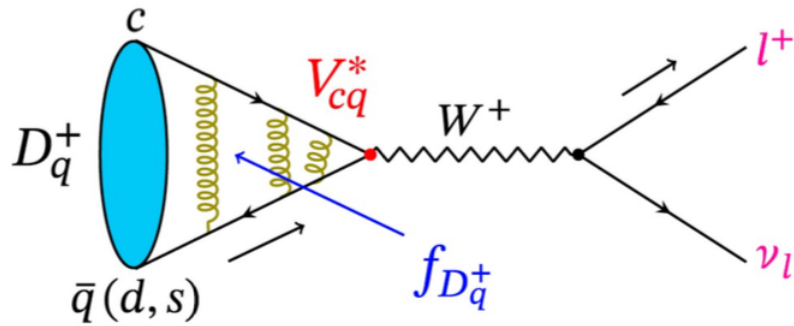
→ Get the information of $D \rightarrow VP, PP \dots$, where V and P denote vector and pseudoscalar, respectively.

Absolute BF's measurement

→ Test theoretical calculations of these BF's and benefit the understanding of the quark SU(3) flavor symmetry and CP violation.

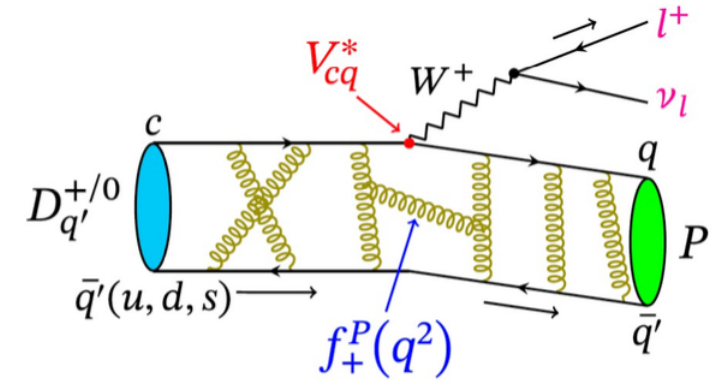
Why charm meson physics?

Pure leptonic decay



$$\Gamma(D_{(s)}^+ \rightarrow l^+ \nu_l) \propto |f_{D_s^+}|^2 \cdot |V_{cd(s)}|^2$$

Semi-leptonic decay



$$\Gamma(D_{(s)} \rightarrow Pl^+ \nu_l) \propto |f_+(q^2)|^2 \cdot |V_{cd(s)}|^2$$

Ideal bridge to access the strong and weak effects between quarks

- ✓ $|V_{cs}|$ and $|V_{cd}|$ → Test CKM matrix unitarity
- ✓ Decay Constant $f_{D_{(s)}^+}$ and form factor f_+ → Calibrate LQCD
- ✓ Branching fraction (BF) ratio → Test lepton flavor universality



Data sets

- 20 fb^{-1} @ $\sqrt{s} = 3.773 \text{ GeV}$

$$e^+e^- \rightarrow \psi(3770) \rightarrow D\bar{D}$$

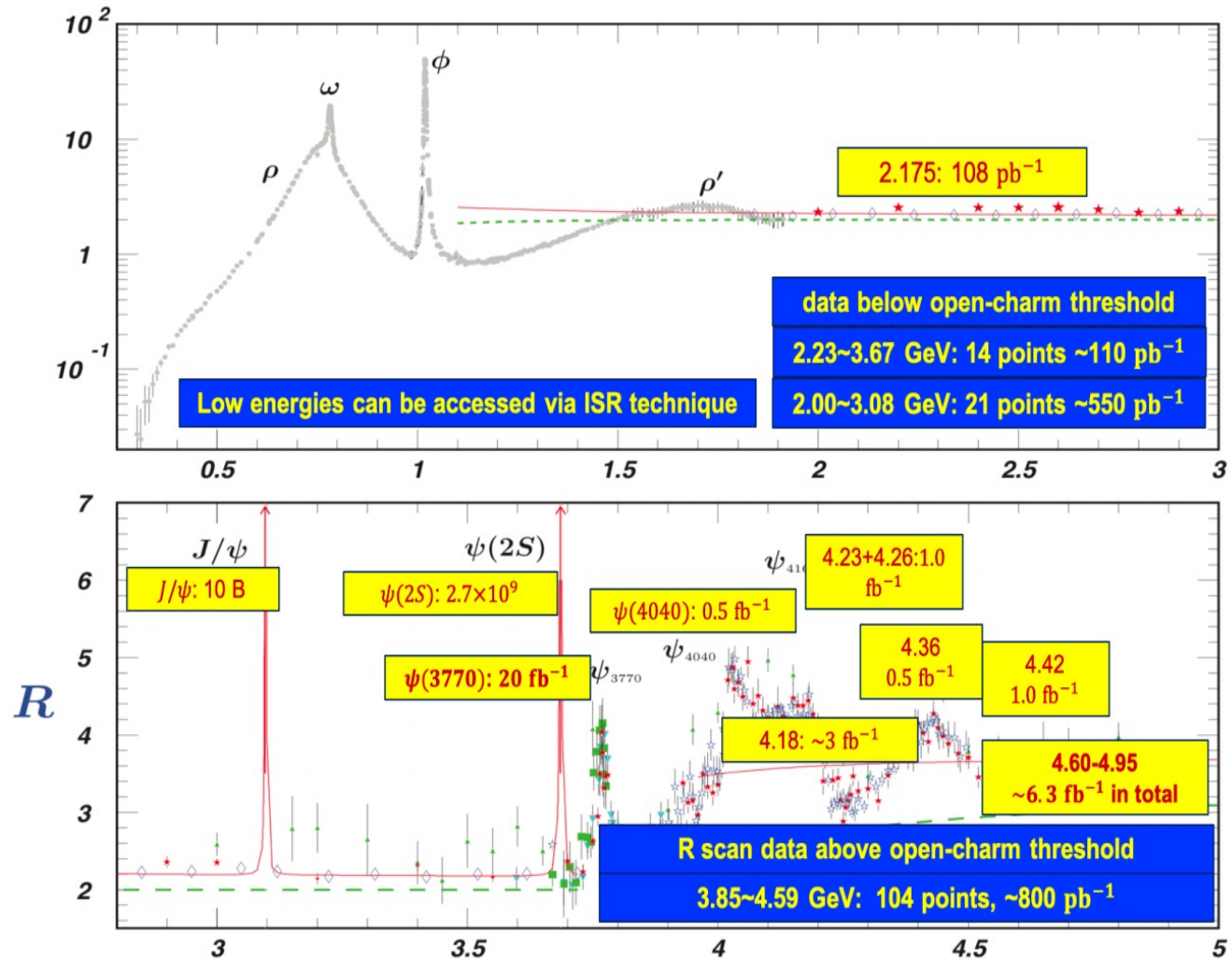
2.93 fb^{-1} in 2010-2011

- 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$

$$e^+e^- \rightarrow D_s^{*+} D_s^-$$

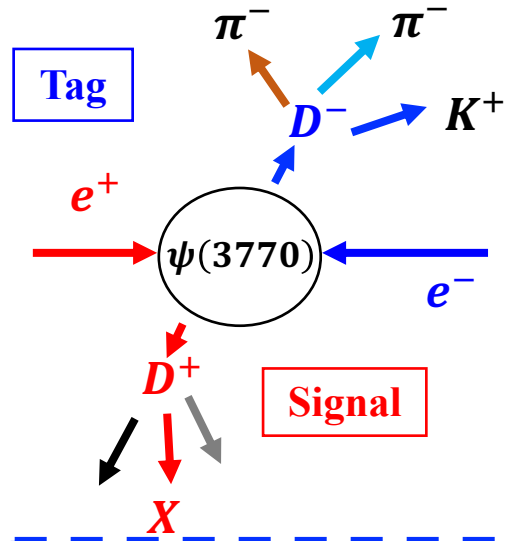
- 10.6 fb^{-1} @ $\sqrt{s} = 4.237 - 4.699 \text{ GeV}$

$$e^+e^- \rightarrow D_s^{*+} D_s^{*-}$$





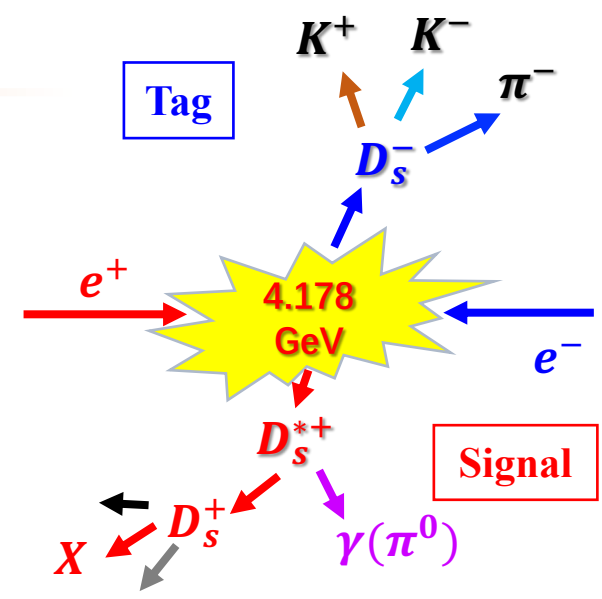
Analysis technique



Charge conjugated processes are implied

The signal branching fraction:

$$B_{\text{sig}} = \frac{N_{DT}^{\text{signal}}}{N_{D(s)}^{\text{ST}} \times \bar{\epsilon}_{\text{sig}}}$$



- **Single tag (ST):**
fully reconstruct one D^-

$$\Delta E = E_{D^-} - E_{\text{beam}}$$

$$M_{\text{BC}} = \sqrt{E_{\text{beam}}^2 - |\vec{p}_{D^-}|^2}$$

PRL124(2020)231801

- **Double tag (DT):**
in the recoil ST $D_{(s)}^-$,
analyze the signal $D_{(s)}^+$

$$M^2 = E_{\text{miss}}^2 - |\vec{p}_{\text{miss}}|^2$$

$$E_{\text{miss}} = E_{\text{cm}} - \sqrt{|\vec{p}_{D_{(s)}^-}|^2 + M_{D_{(s)}^-}^2} - E_X$$

$$\vec{p}_{\text{miss}} = -\vec{p}_{D_{(s)}^-} - \vec{p}_X$$

$$U_{\text{miss}} = E_{\text{miss}} - |\vec{p}_{\text{miss}}|$$

or other variables

- **Single tag (ST):**
fully reconstruct one D_s^-

$$M_{\text{rec}} = \sqrt{\left(E_{\text{cm}} - \sqrt{|\vec{p}_{D_s^-}|^2 + m_{D_s^-}^2}\right)^2 - |\vec{p}_{D_s^-}|^2}$$

PRD104(2021)052009 PRD 104 (2021) 012003



- Closed source / hand coded
 - **Tensor formulism:** most of charm decays. [$D^+ \rightarrow K_S^0 \pi^+ \pi^0 \pi^0$: [JHEP09, 077\(2023\)](#)]
 - Helicity formulism: [$e^+ e^- \rightarrow \omega \pi^+ \pi^-$: [JHEP08,159\(2023\)](#)]
- [GPUPWA](#):
 - First PWA tool based on GPU
 - Used in many PWA of light mesons: [$J/\psi \rightarrow \gamma \eta \eta$: [PRD87, 092009\(2013\)](#)]; [$J/\psi \rightarrow \gamma \eta \eta'$: [PRD106, 072012\(2022\)](#)]
- [FDC-PWA](#):
 - Feynman Diagram Calculation
 - Used in some baryon final states [$\psi' \rightarrow p \bar{p} \eta$: [PRD88, 032010\(2013\)](#)]; [$e^+ e^- \rightarrow p K^- \bar{\Lambda}$: [PRL131, 151901 \(2023\)](#)]
- **[TF-PWA](#)**:
 - TensorFlow-based, configurable, GPU acceleration, AD
 - as an example: [$\Lambda_c^+ \rightarrow \Lambda \pi^+ \pi^0$: [JHEP12, 033\(2022\)](#)]
- Other tools:
 - [Amptools](#): [$\chi_{c1} \rightarrow \eta \pi^+ \pi^-$: [PRD95,032002\(2017\)](#)]
 - [PAWIAN](#): [$e^+ e^- \rightarrow \phi K^+ K^-$: [PRD108, 032004 \(2023\)](#)]
 - [ComPWA](#): [$D^0 \rightarrow K_S K^+ K^-$: [arXiv:2006.02800](#)]

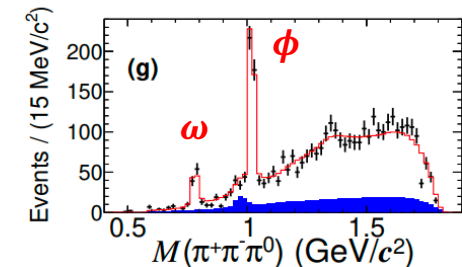
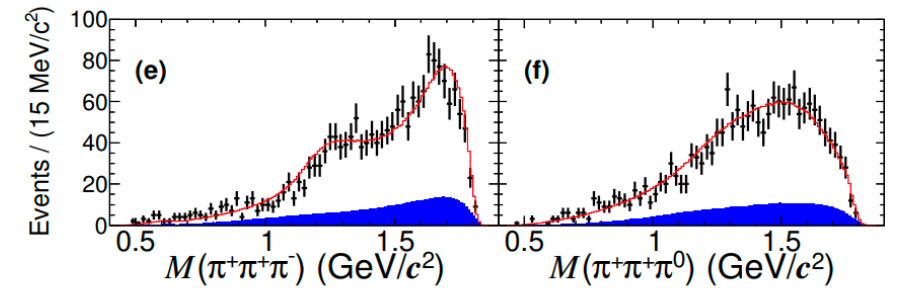
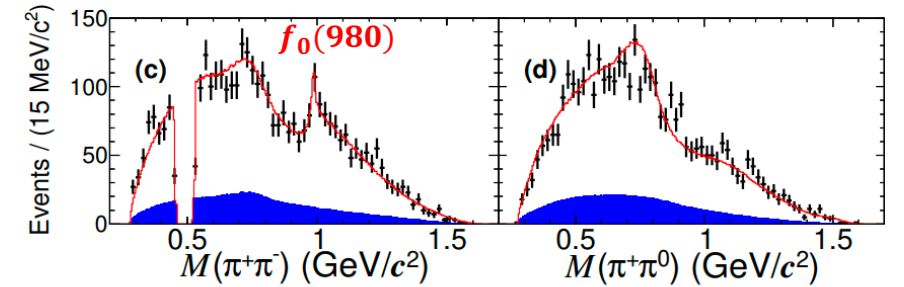
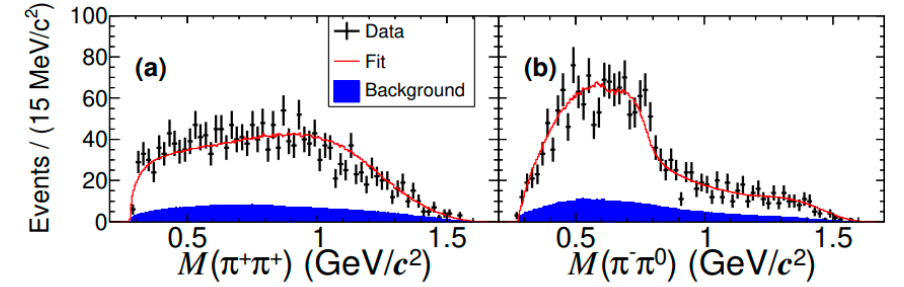
Hadronic decays



$D_S^+ \rightarrow \pi^+ \pi^+ \pi^- \pi^0$

- $e^+e^- \rightarrow D_S^{*+} D_S^-$ with 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$
- $D_S^+ \rightarrow f_0(980) \rho^+$: mainly a W -external emission channel.
- $D_S^+ \rightarrow \omega \pi^+$: solely via the W -annihilation process.

Component	Phase (rad)	FF (%)	BF (10^{-3})
$f_0(1370)\rho^+$	0.0(fixed)	$24.9 \pm 3.8 \pm 2.1$	$5.08 \pm 0.80 \pm 0.43$
$f_0(980)\rho^+$	$3.99 \pm 0.13 \pm 0.07$	$12.6 \pm 2.1 \pm 1.0$	$2.57 \pm 0.44 \pm 0.20$
$f_2(1270)\rho^+$	$1.11 \pm 0.10 \pm 0.10$	$9.5 \pm 1.7 \pm 0.6$	$1.94 \pm 0.36 \pm 0.12$
$(\rho^+ \rho^0)_S$	$1.10 \pm 0.18 \pm 0.10$	$3.5 \pm 1.2 \pm 0.6$	$0.71 \pm 0.25 \pm 0.12$
$(\rho(1450)^+ \rho^0)_S$	$0.43 \pm 0.18 \pm 0.17$	$4.6 \pm 1.3 \pm 0.8$	$0.94 \pm 0.27 \pm 0.16$
$(\rho^+ \rho(1450)^0)_P$	$4.58 \pm 0.16 \pm 0.09$	$8.6 \pm 1.3 \pm 0.4$	$1.75 \pm 0.27 \pm 0.08$
$\phi((\rho\pi) \rightarrow \pi^+ \pi^- \pi^0)\pi^+$	$2.90 \pm 0.15 \pm 0.18$	$24.9 \pm 1.2 \pm 0.4$	$5.08 \pm 0.32 \pm 0.10$
$\omega((\rho\pi) \rightarrow \pi^+ \pi^- \pi^0)\pi^+$	$3.22 \pm 0.21 \pm 0.09$	$6.9 \pm 0.8 \pm 0.3$	$1.41 \pm 0.17 \pm 0.06$
$a_1^+(\rho^0 \pi^+)_S \pi^0$	$3.78 \pm 0.16 \pm 0.12$	$12.5 \pm 1.6 \pm 1.0$	$2.55 \pm 0.34 \pm 0.20$
$a_1^0((\rho\pi)_S \rightarrow \pi^+ \pi^- \pi^0)\pi^+$	$4.82 \pm 0.15 \pm 0.12$	$6.3 \pm 1.9 \pm 1.2$	$1.29 \pm 0.39 \pm 0.24$
$\pi(1300)^0((\rho\pi)_P \rightarrow \pi^+ \pi^- \pi^0)\pi^+$	$2.22 \pm 0.14 \pm 0.08$	$11.7 \pm 2.3 \pm 2.2$	$2.39 \pm 0.48 \pm 0.45$



$$B_{D_S^+ \rightarrow \pi^+ \pi^+ \pi^- \pi^0 |_{\text{non-}\eta}} = (2.04 \pm 0.08_{\text{stat}} \pm 0.05_{\text{syst}}) \%$$

First time

$$\frac{B_{\phi \rightarrow \pi^+ \pi^- \pi^0}}{B_{\phi \rightarrow K^+ K^-}} = (0.230 \pm 0.014_{\text{stat}} \pm 0.010_{\text{syst}}) \%$$

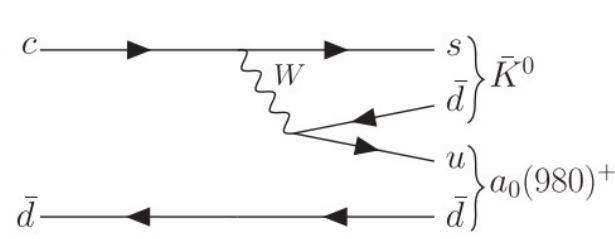
> 4 σ deviation from PDG value



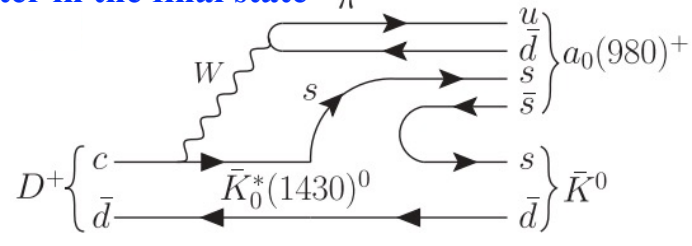
$D^+ \rightarrow K_S^0 \pi^+ \eta$

- $e^+e^- \rightarrow \psi(3770) \rightarrow D\bar{D}$ with 2.93 fb^{-1} @ 3.773 GeV
- $D^+ \rightarrow K_S^0 a_0(980)$ is the only W -annihilation-free decay among D to $a_0(980)$ pseudoscalar, which is the ideal decay in extracting the contributions of the W -emission amplitudes involving $a_0(980)$ and to study the final-state interactions.

Internal W -emission



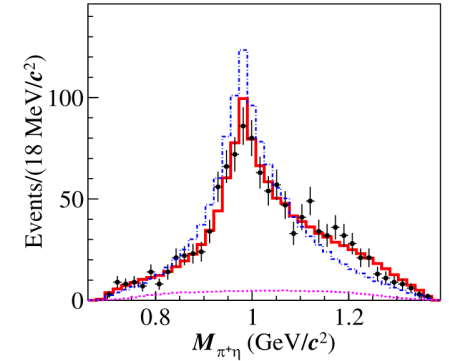
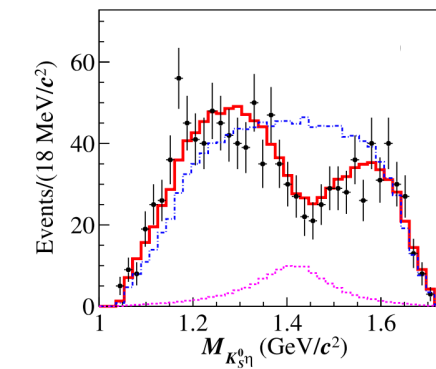
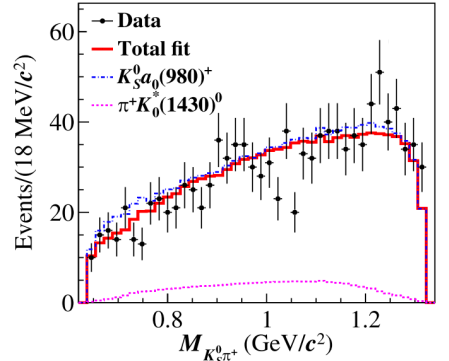
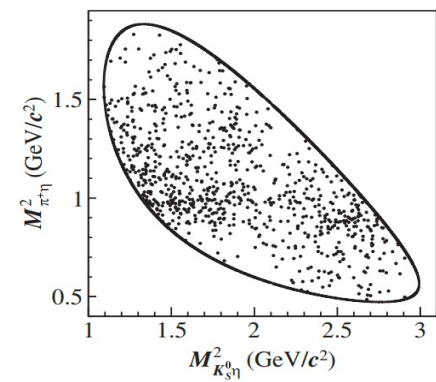
Rescatter in the final state π^+



$$B_{D^+ \rightarrow K_S^0 \pi^+ \eta} = (1.27 \pm 0.04_{\text{stat}} \pm 0.03_{\text{syst}})\%$$

$$B_{D^+ \rightarrow K_S^0 a_0(980)^+ (\rightarrow \pi^+ \eta)} = (1.33 \pm 0.05_{\text{stat}} \pm 0.04_{\text{syst}})\%$$

$$B_{D^+ \rightarrow \pi^+ \bar{K}_0^*(1430)^0 (\rightarrow K_S^0 \eta)} = (0.14 \pm 0.03_{\text{stat}} \pm 0.01_{\text{syst}})\%$$



Amplitude	Phase ϕ (rad)	FF (%)	Significance
$D^+ \rightarrow K_S^0 a_0(980)^+$	0.0 (fixed)	$105.00 \pm 0.94 \pm 1.04 \pm 0.07$	$>10\sigma$
$D^+ \rightarrow \bar{K}_0^*(1430)^0 \pi^+$	$2.58 \pm 0.06 \pm 0.09 \pm 0.01$	$10.83 \pm 1.50 \pm 1.27 \pm 0.08$	$>10\sigma$

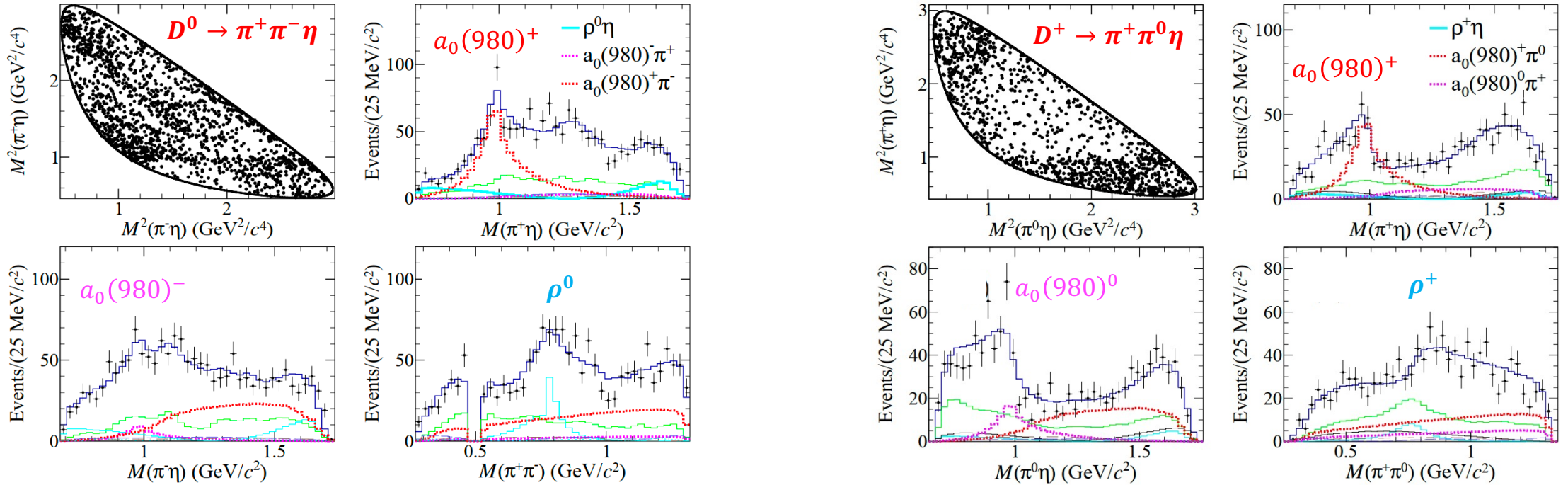
Destructive interference



$D \rightarrow \pi \pi \eta$

- $e^+e^- \rightarrow \psi(3770) \rightarrow D\bar{D}$ with 7.9 fb^{-1} @ 3.773 GeV
- The measurement of the branching fractions for $D \rightarrow \pi a_0(980)$ and the corresponding relative ratio can constrain the size and phase of the amplitude of the W -annihilation (W -exchange) process.

arXiv:2404.09219 [hep-ex]



$$\frac{FF_{D^0 \rightarrow \pi^- a_0(980)^+}}{FF_{D^0 \rightarrow \pi^+ a_0(980)^-}} = 7.5_{-0.8}^{+2.5} \pm 1.7_{\text{stat}} \pm 1.7_{\text{syst}} \quad \text{Significantly higher than unity.}$$

$$\frac{FF_{D^+ \rightarrow \pi^0 a_0(980)^+}}{FF_{D^+ \rightarrow \pi^+ a_0(980)^0}} = 2.6 \pm 0.6_{\text{stat}} \pm 0.3_{\text{syst}}$$

$$B_{D^0 \rightarrow \pi^+ \pi^- \eta} = (1.24 \pm 0.04_{\text{stat}} \pm 0.03_{\text{syst}}) \times 10^{-3}$$

$$B_{D^+ \rightarrow \pi^+ \pi^0 \eta} = (2.18 \pm 0.12_{\text{stat}} \pm 0.05_{\text{syst}}) \times 10^{-3}$$



Polarization of $D^0 \rightarrow \omega\phi$

PRL128(2022)011803



- The polarization in the $D^0 \rightarrow VV$ decay is sensitive to the V-A structure of electroweak interactions in the SM, spin correlations, and the final state interactions, among other effects.

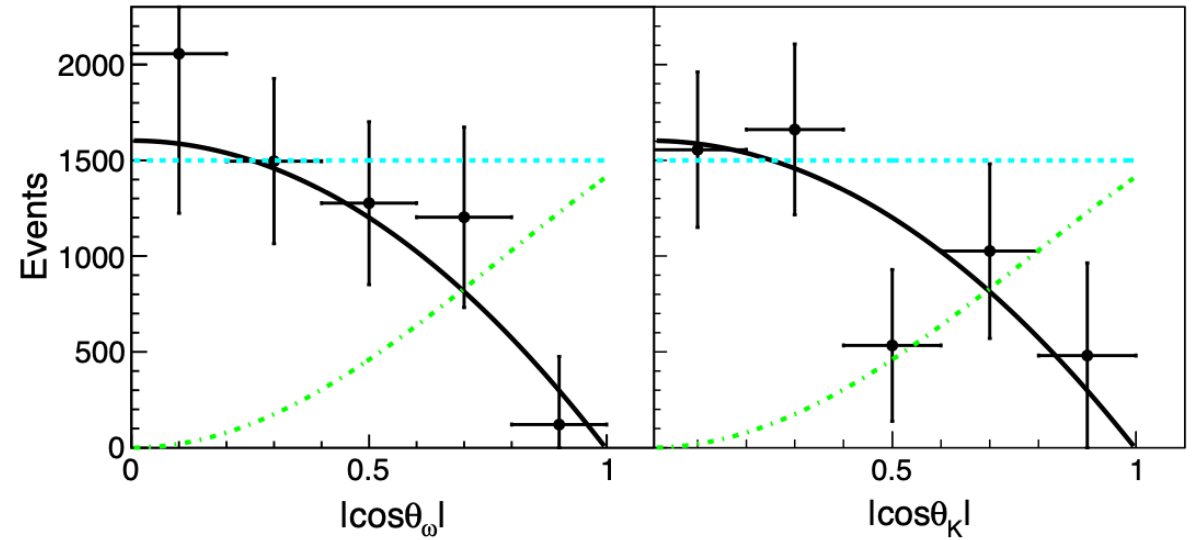
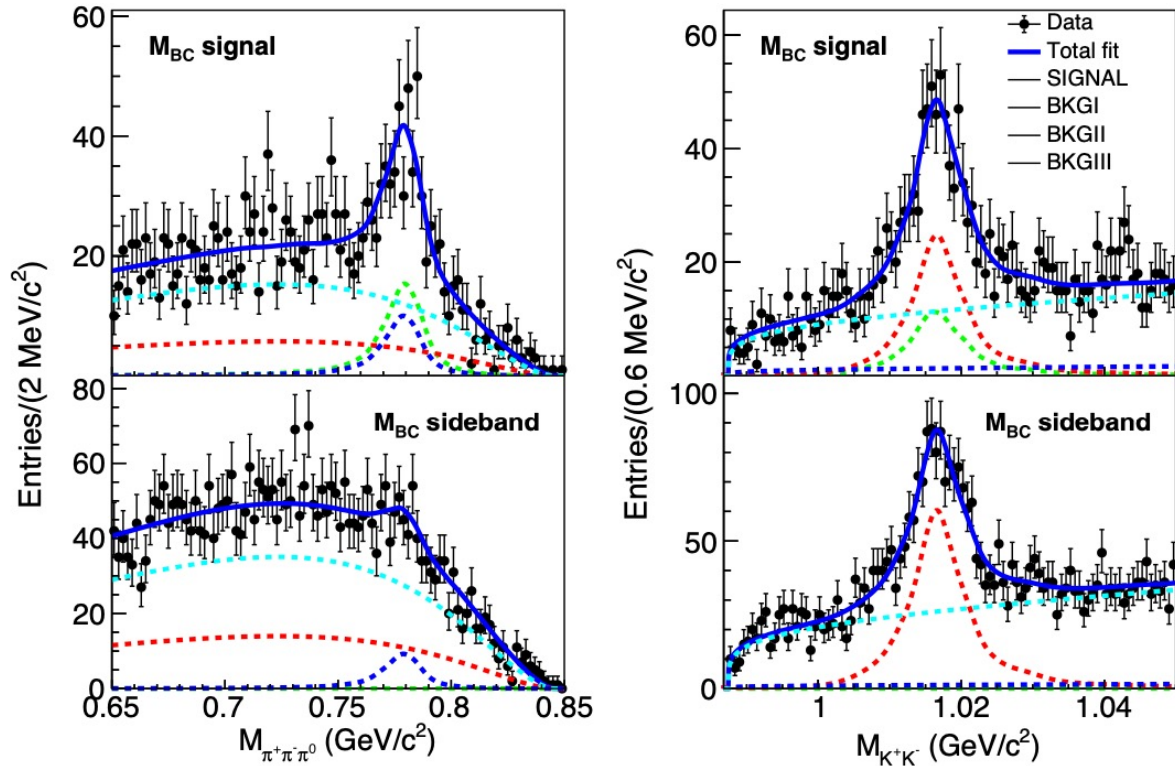
- 2-D unbinned maximum likelihood fit $N_{sig} = 195.9 \pm 29.1$

$$B(D^0 \rightarrow \omega\phi) = (6.48 \pm 0.96 \pm 0.40) \times 10^{-4}$$

$f_L < 0.24$ @ 95% CL

PHSP

Longitudinal polarization



ω and ϕ are transversely polarized

(Semi-) leptonic decays

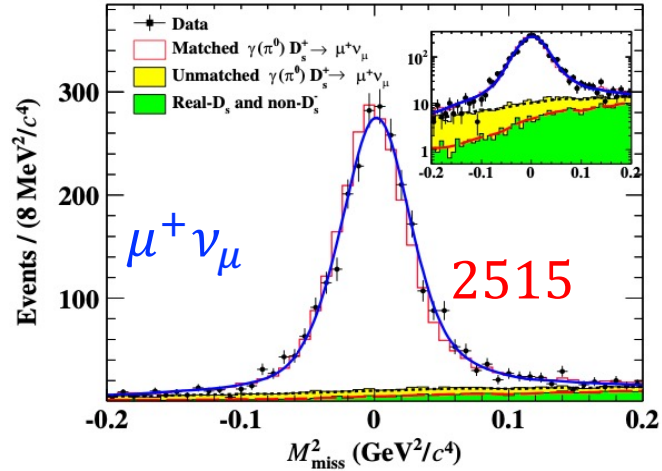


$D_s^+ \rightarrow \ell^+ \nu_\ell \ (\ell = \mu, \tau)$



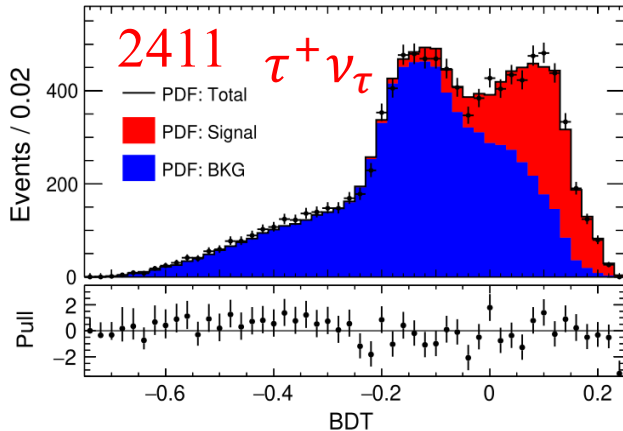
- $\mu^+ \nu_\mu$, muon counter

@4.128 – 4.226 GeV PRD108(2023)112001



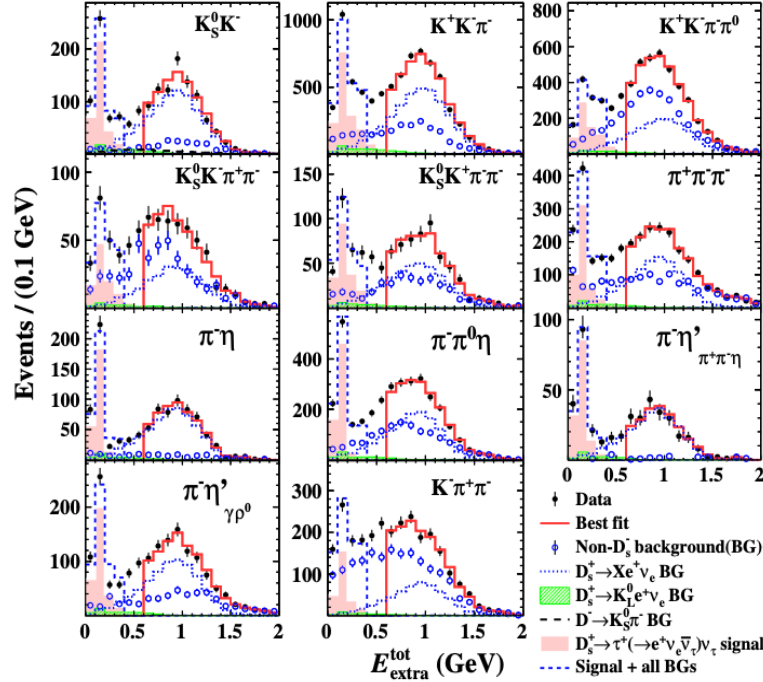
- $\tau^+ \rightarrow \pi^+ \bar{\nu}_\tau$ multivariate analysis

@4.128 – 4.226 GeV PRD108(2023)092014



- $\tau^+ \rightarrow e^+ \nu_e \bar{\nu}_\tau$ 4940

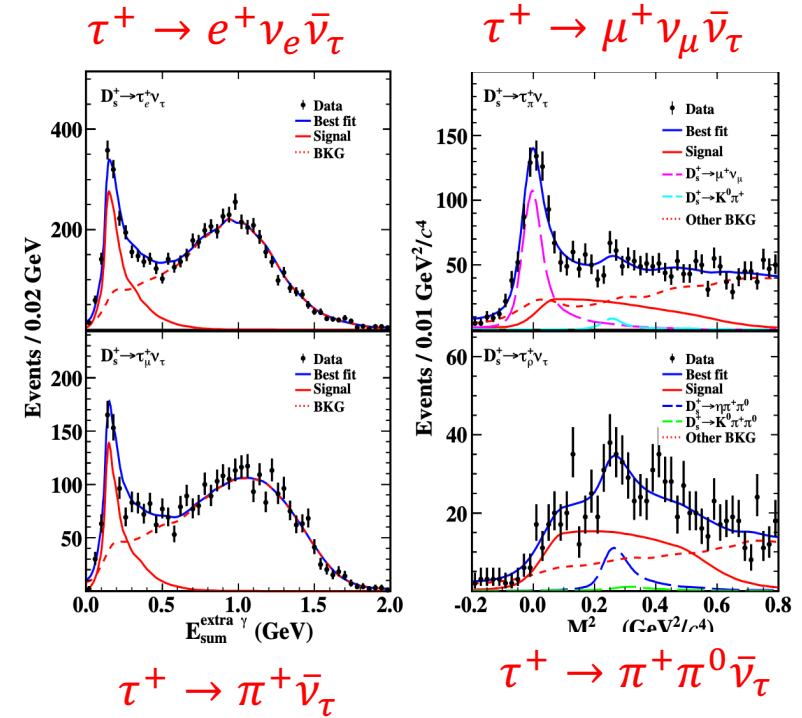
@4.178 – 4.226 GeV PRL127(2021)171801



- $\tau^+ \rightarrow e^+ \nu_e \bar{\nu}_\tau$ (17.82 ± 0.04)%
- $\tau^+ \rightarrow \mu^+ \nu_\mu \bar{\nu}_\tau$ (17.39 ± 0.04)%
- $\tau^+ \rightarrow \pi^+ \pi^0 \bar{\nu}_\tau$ (25.49 ± 0.09)%
- $\tau^+ \rightarrow \pi^+ \bar{\nu}_\tau$ (10.82 ± 0.05)%

- $\tau^+ \nu_\tau$ 2845

@4.237 – 4.699 GeV PRD110(2024)052002



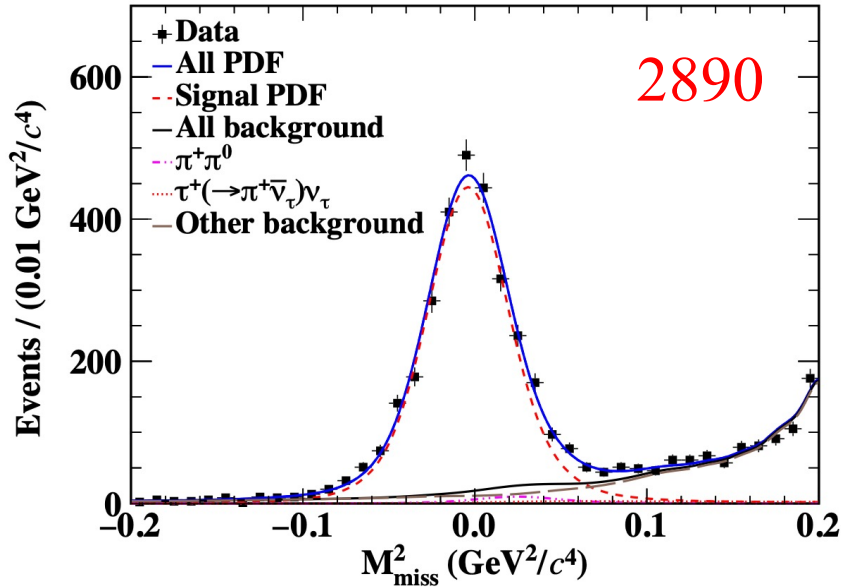
Simultaneous fit with the same BF



$D^+ \rightarrow \ell^+ \nu_\ell \ (\ell = \mu, \tau)$

- $\mu^+ \nu_\mu$, muon counter

20.3 fb⁻¹@3.773 GeV arXiv:2410.07626



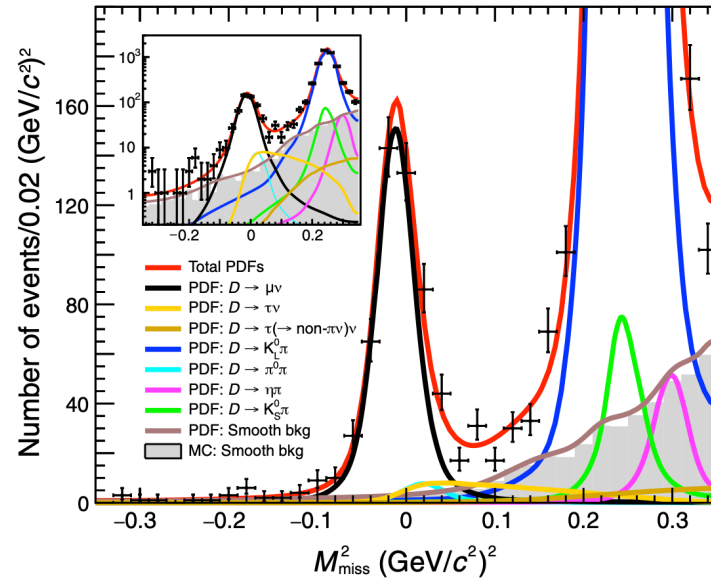
$$B(D^+ \rightarrow \mu^+ \nu_\mu) = (3.981 \pm 0.079 \pm 0.040) \times 10^{-4}$$

$f_{D^+|V_{cd}}$ The most precise to date

$$= (47.53 \pm 0.48_{\text{stat}} \pm 0.24_{\text{syst}} \pm 0.12_{\text{input}}) \text{ MeV}$$

- $\tau^+ \nu_\tau$ 2.93 fb⁻¹@3.773 GeV

μ -like: $E_{\text{EMC}} \leq 0.3 \text{ GeV}$



$$B(D^+ \rightarrow \tau^+ \nu_\tau) = (1.20 \pm 0.24 \pm 0.12) \times 10^{-3}$$

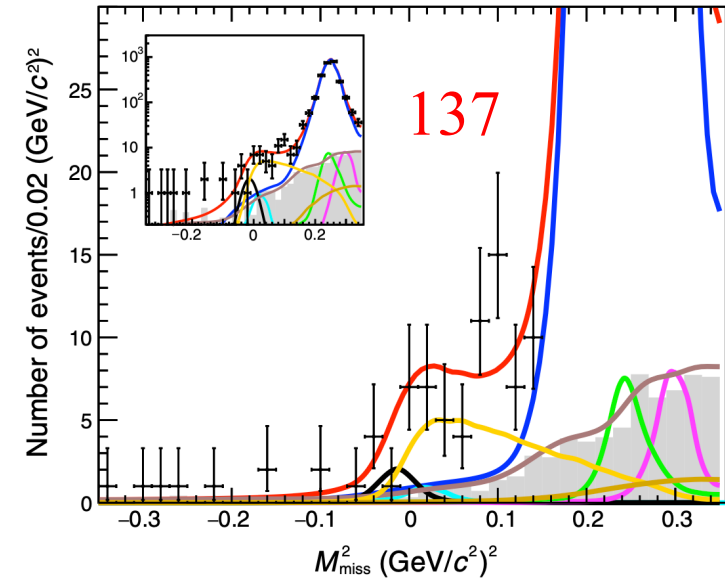
$$R_{\tau/\mu} = \frac{\bar{\Gamma}(D^+ \rightarrow \tau^+ \nu_\tau)}{\bar{\Gamma}(D^+ \rightarrow \mu^+ \nu_\mu)} = 3.02 \pm 0.68$$

SM prediction: 2.67 ± 0.01

No LFU violation

PRL123(2019)211802

π -like: $E_{\text{EMC}} > 0.3 \text{ GeV}$



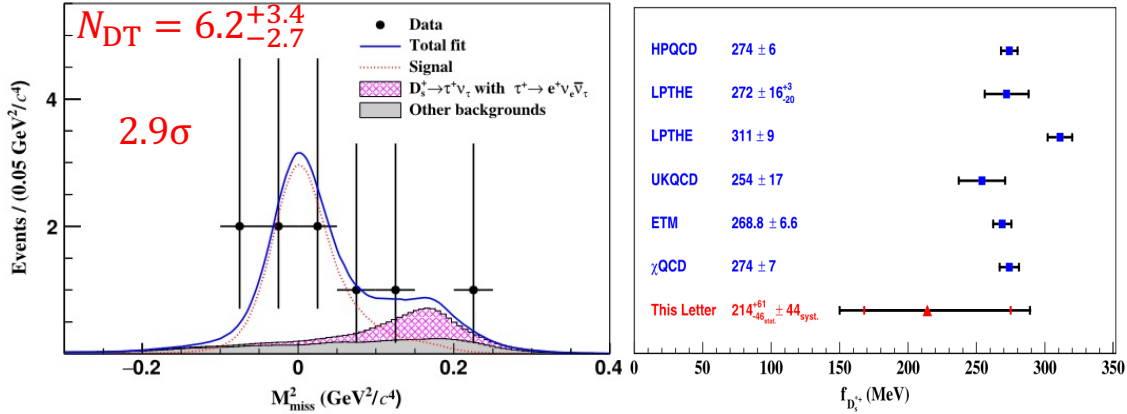


$D_{(s)}^{*+} \rightarrow \ell^+ \nu_\ell$ ($\ell = e, \mu$) Charmed vector meson



- $D_S^{*+} \rightarrow e^+ \nu_e$ @4.128 – 4.226 GeV

PRL131(2023)141802



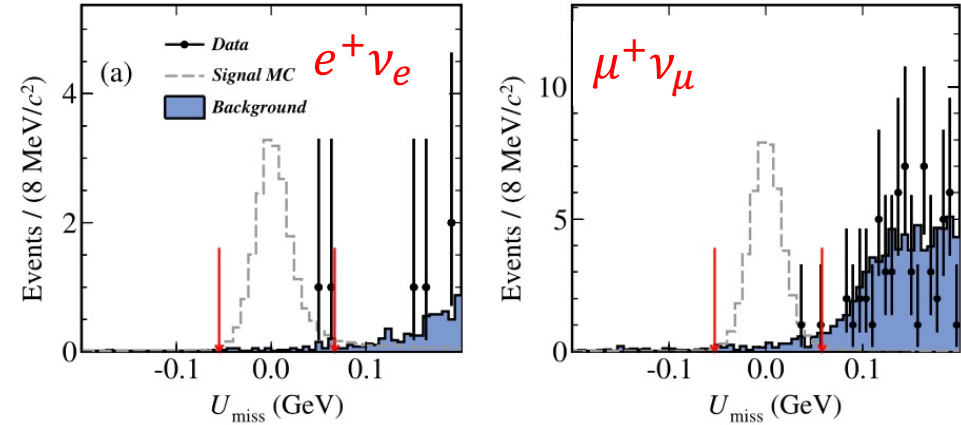
$$B(D_S^{*+} \rightarrow e^+ \nu_e) = (2.1_{-0.9}^{+1.2} \pm 0.2) \times 10^{-5}$$

$$f_{D_S^{*+}} |V_{cs}| = (208_{-45}^{+59} \pm 43) \text{ MeV}$$

$$f_{D_S^{*+}} = (214_{-46}^{+61} \pm 44) \text{ MeV} \quad \text{For the first time}$$

- $D^{*+} \rightarrow \ell^+ \nu_\ell$ ($\ell = e, \mu$) @4.178 – 4.226 GeV

PRD110(2024)012003



No significant signal. Profile likelihood method.

$$B(D^{*+} \rightarrow e^+ \nu_e) < 1.1 \times 10^{-5} \text{ @ 90\% C. L.}$$

$$B(D^{*+} \rightarrow \mu^+ \nu_\mu) < 4.3 \times 10^{-6} \text{ @ 90\% C. L.}$$

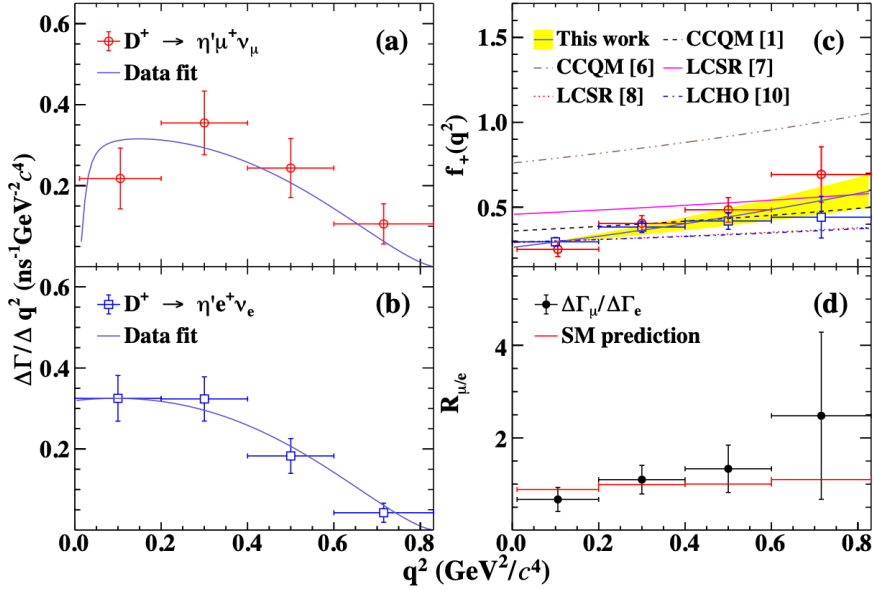


$D_{(s)}^+ \rightarrow P \ell^+ \nu_\ell$

$e^+e^- \rightarrow \psi(3770) \rightarrow D\bar{D}$ with 20.3 fb^{-1} @ 3.773 GeV

• $D^+ \rightarrow \eta' \ell^+ \nu_\ell$

arXiv:2410.08603



$$B(D^+ \rightarrow \eta' \mu^+ \nu_\mu) = (1.92 \pm 0.28_{\text{stat}} \pm 0.08_{\text{syst}}) \times 10^{-4}. \quad 8.6\sigma$$

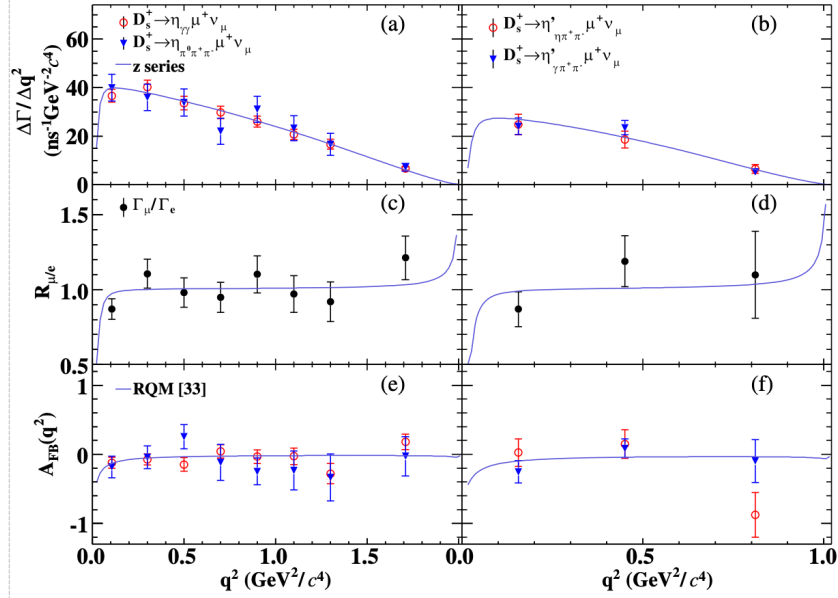
$$B(D^+ \rightarrow \eta' e^+ \nu_e) = (1.79 \pm 0.19_{\text{stat}} \pm 0.07_{\text{syst}}) \times 10^{-4}.$$

$$f_+^{\eta'}(0) |V_{cd}| = (5.92 \pm 0.56_{\text{stat}} \pm 0.13_{\text{syst}}) \times 10^{-2}.$$

The $\eta - \eta'$ mixing angle in the quark flavor basis is determined to be $\phi_P = (39.8 \pm 0.8_{\text{stat}} \pm 0.3_{\text{syst}})^\circ$

$e^+e^- \rightarrow D_s^* D_s^-$ with 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$

• $D_s^+ \rightarrow \eta^{(\prime)} \mu^+ \nu_\mu$ PRL132(2024)091802



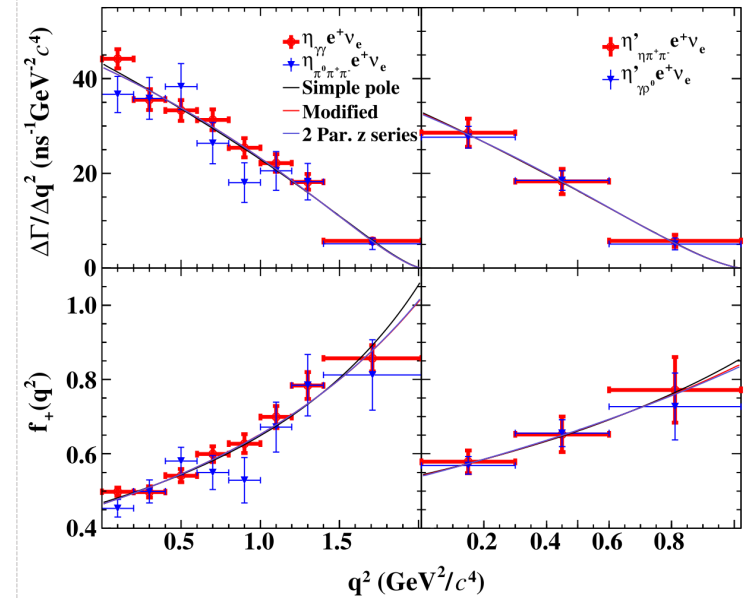
$$B(D_s^+ \rightarrow \eta' \mu^+ \nu_\mu) = (0.801 \pm 0.055_{\text{stat}} \pm 0.028_{\text{syst}}) \% \quad > 10\sigma$$

$$B(D_s^+ \rightarrow \eta \mu^+ \nu_\mu) = (2.235 \pm 0.051_{\text{stat}} \pm 0.052_{\text{syst}}) \%.$$

$$f_{+,0}^{\eta}(0) |V_{cs}| = 0.452 \pm 0.010_{\text{stat}} \pm 0.007_{\text{syst}}.$$

$$f_{+,0}^{\eta'}(0) |V_{cs}| = 0.504 \pm 0.037_{\text{stat}} \pm 0.012_{\text{syst}}.$$

• $D_s^+ \rightarrow \eta^{(\prime)} e^+ \nu_e$ PRD108(2023)092003



$$B(D_s^+ \rightarrow \eta' e^+ \nu_e) = (0.810 \pm 0.038_{\text{stat}} \pm 0.024_{\text{syst}}) \%.$$

$$B(D_s^+ \rightarrow \eta e^+ \nu_e) = (2.255 \pm 0.039_{\text{stat}} \pm 0.051_{\text{syst}}) \%.$$

$$f_+^{\eta}(0) |V_{cs}| = 0.4519 \pm 0.0071_{\text{stat}} \pm 0.0065_{\text{syst}}.$$

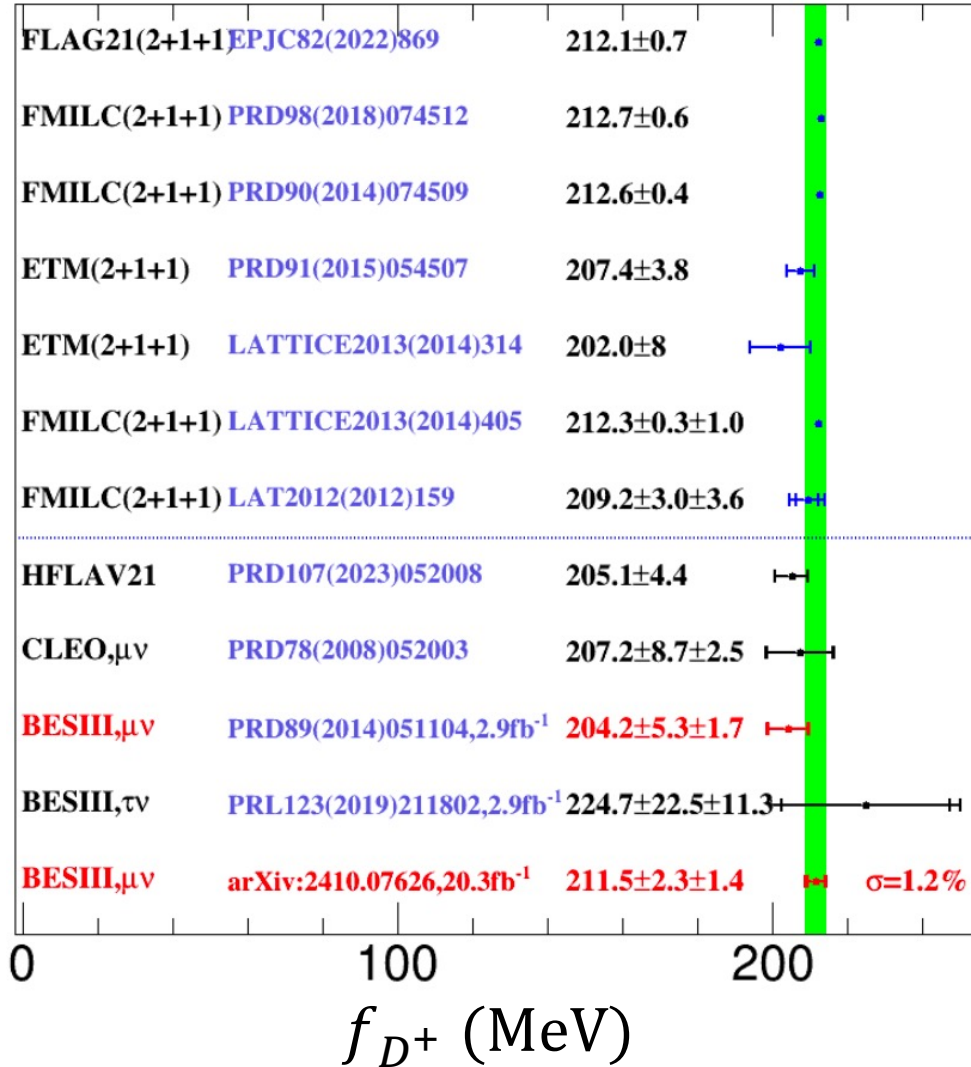
$$f_+^{\eta'}(0) |V_{cs}| = 0.525 \pm 0.024_{\text{stat}} \pm 0.009_{\text{syst}}.$$



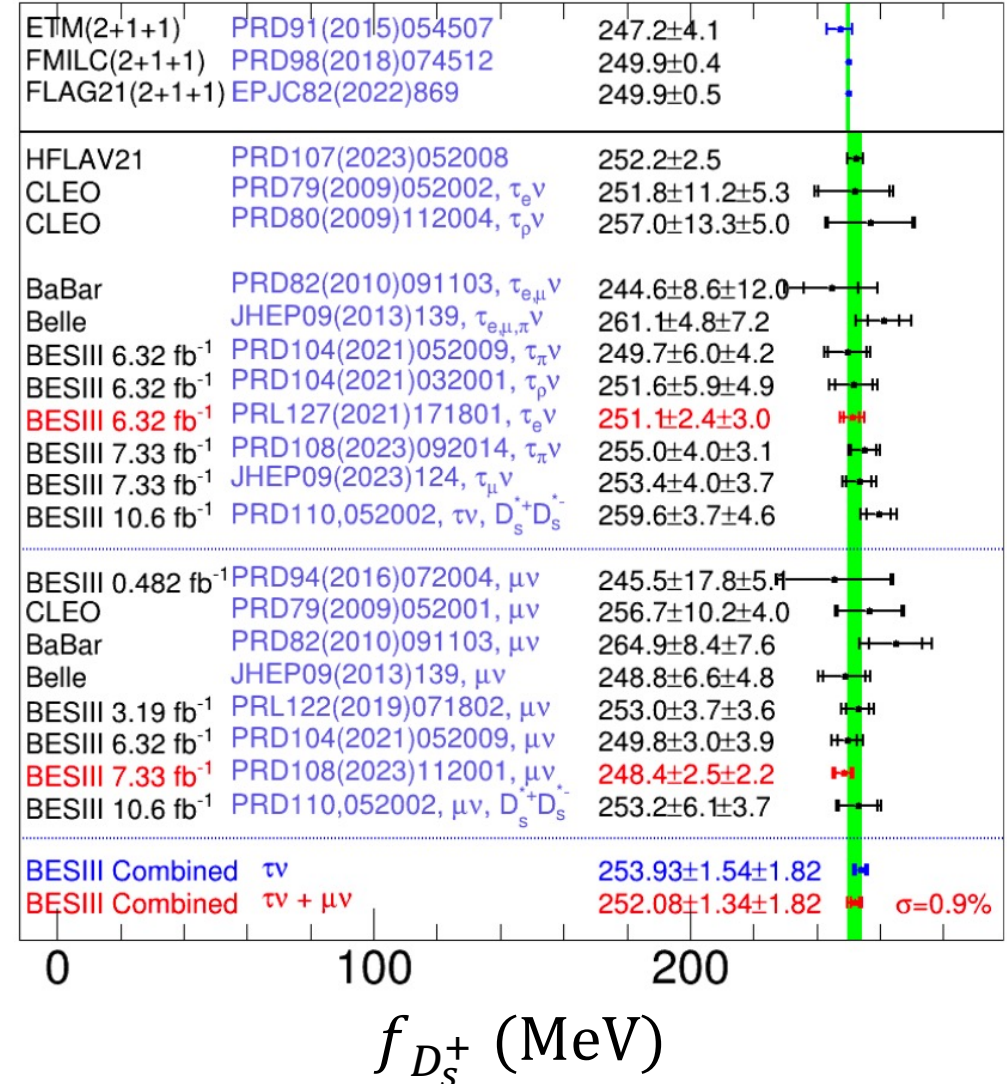
Comparison of the $D_{(s)}^+$ decay constant



- Input $|V_{cd}| = 0.22487 \pm 0.00068$ from CKM global fit



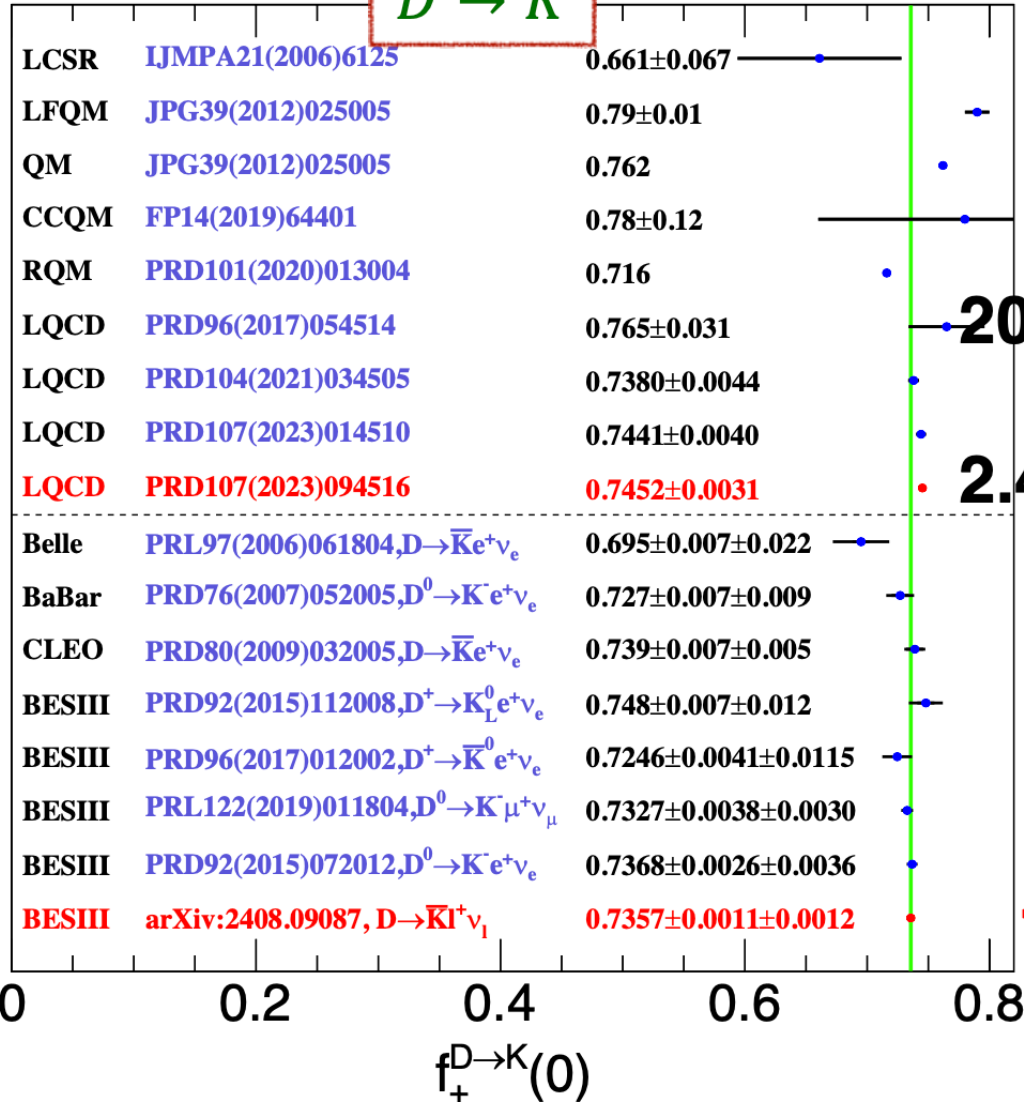
- Input $|V_{cs}| = 0.97320 \pm 0.00011$ MeV from CKM global fit





Comparison of form factor $D \rightarrow K(\pi)\ell^+ \nu_\ell$

$D \rightarrow K$

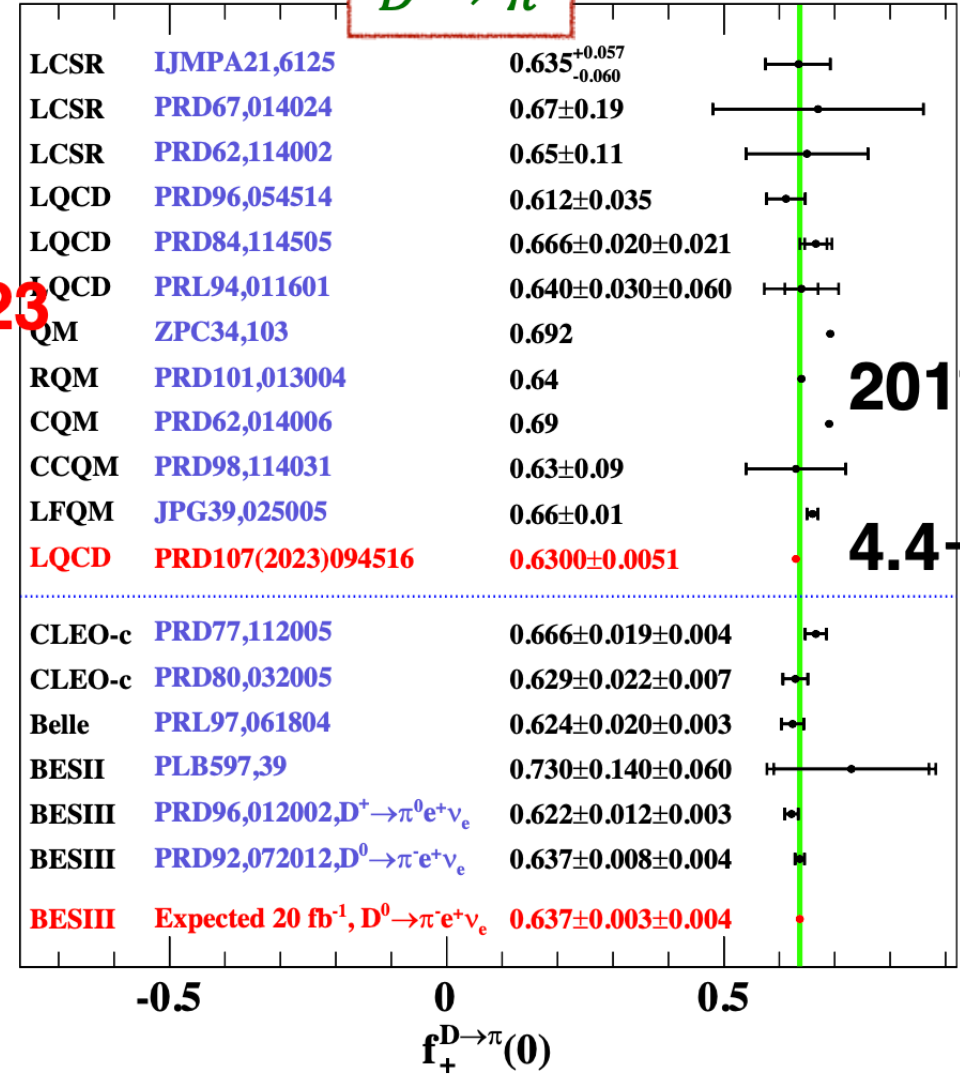


2010 \rightarrow 2023

2.4 \rightarrow 0.4%

< 0.3%

$D \rightarrow \pi$



2011 \rightarrow 2023

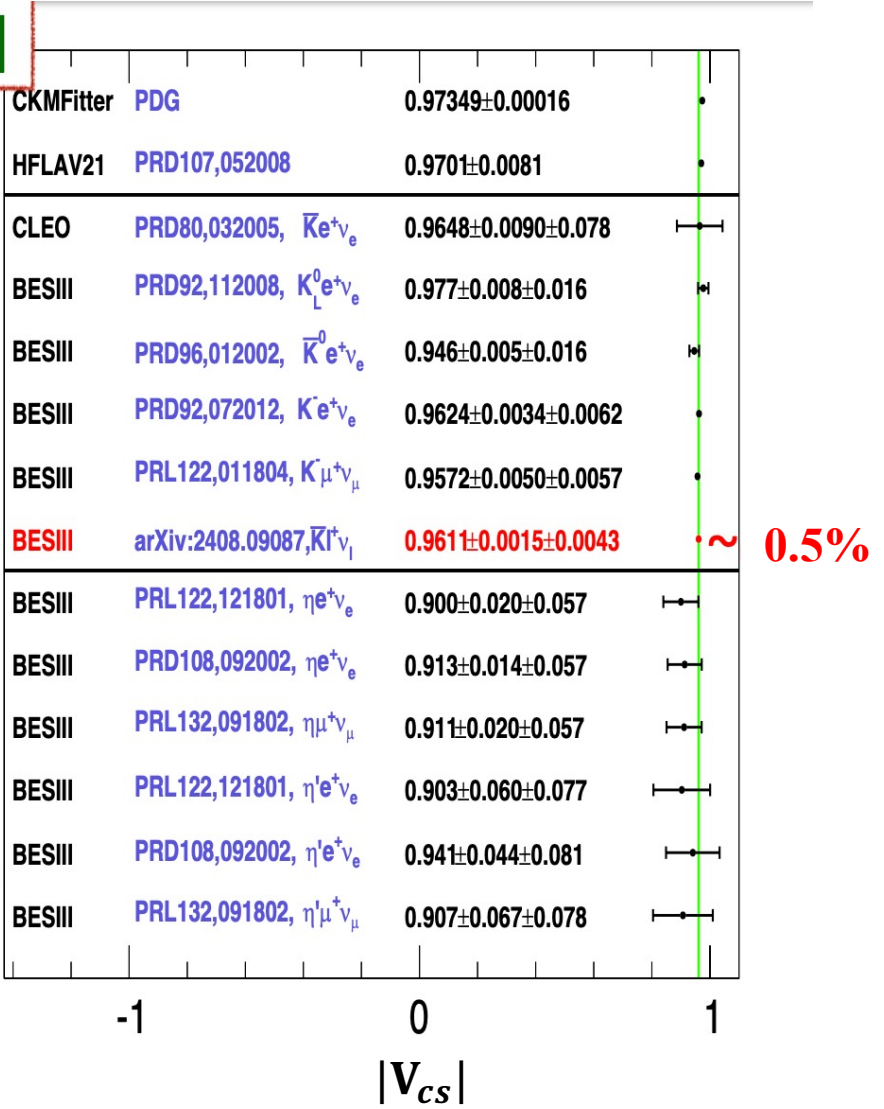
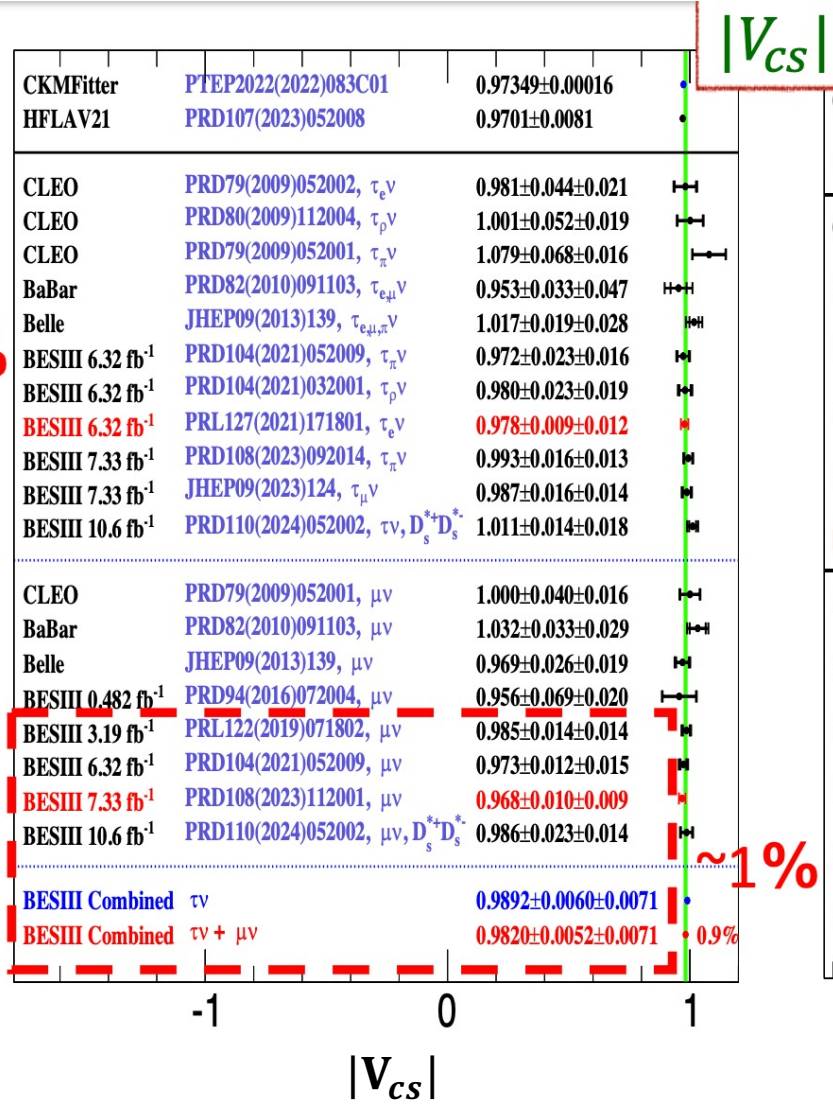
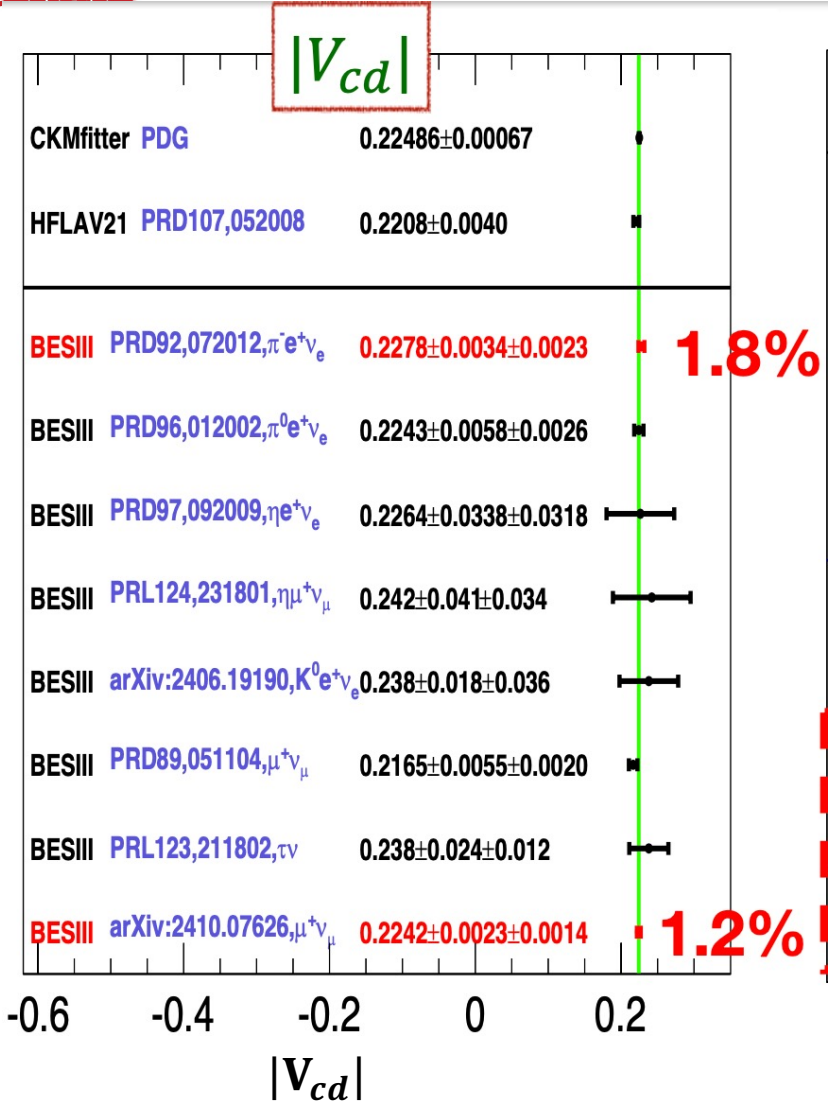
4.4 \rightarrow 0.8%

1.4%
< 1%

Experimental precision is comparable to the latest QCD result.



Comparison of $|V_{cd}(s)|$



pure leptonic D_s^+ decay

Semi-leptonic D_s^+ decay

Rare decays



$D_s^+ \rightarrow h^+(h^0)e^+e^-$

$e^+e^- \rightarrow D_s^+ D_s^-$ with 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$

- In the SM, the $D_s^+ \rightarrow h^+(h^0)e^+e^-$ involve both short-distance (SD) and long-distance (LD) contribution.
- SD: FCNC, suppressed by GIM mechanism.
- LD: through a radiated photon or an intermediate meson decaying to dileptons, can enhance the BFs to 10^{-6} .

$$B(D_s^+ \rightarrow \pi^+ \phi(\rightarrow e^+e^-)) = (1.17_{-0.21}^{+0.23} \pm 0.03_{\text{syst}}) \times 10^{-5} 7.8\sigma$$

$$B(D_s^+ \rightarrow \rho^+ \phi(\rightarrow e^+e^-)) = (2.44_{-0.62}^{+0.67} \pm 0.16_{\text{syst}}) \times 10^{-5} 4.4\sigma$$

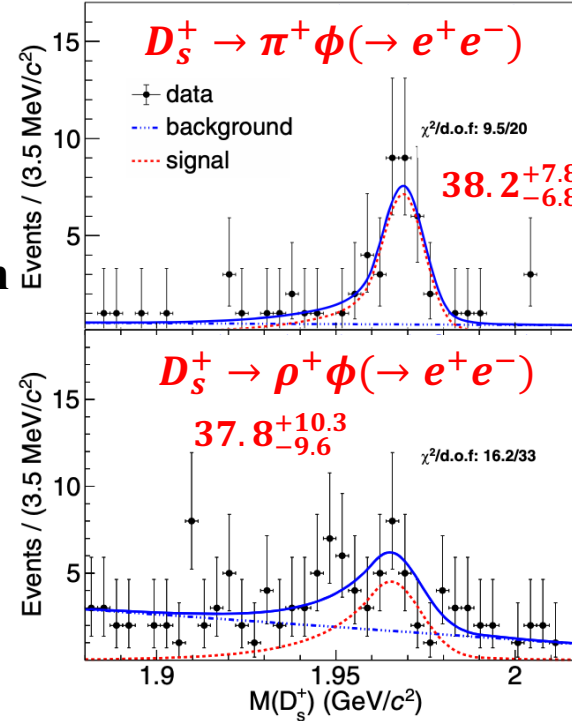
$$B(D_s^+ \rightarrow \pi^+ \pi^0 e^+e^-) < 7.0 \times 10^{-5} @ 90\% \text{ C. L.}$$

$$B(D_s^+ \rightarrow K^+ \pi^0 e^+e^-) < 7.1 \times 10^{-5} @ 90\% \text{ C. L.}$$

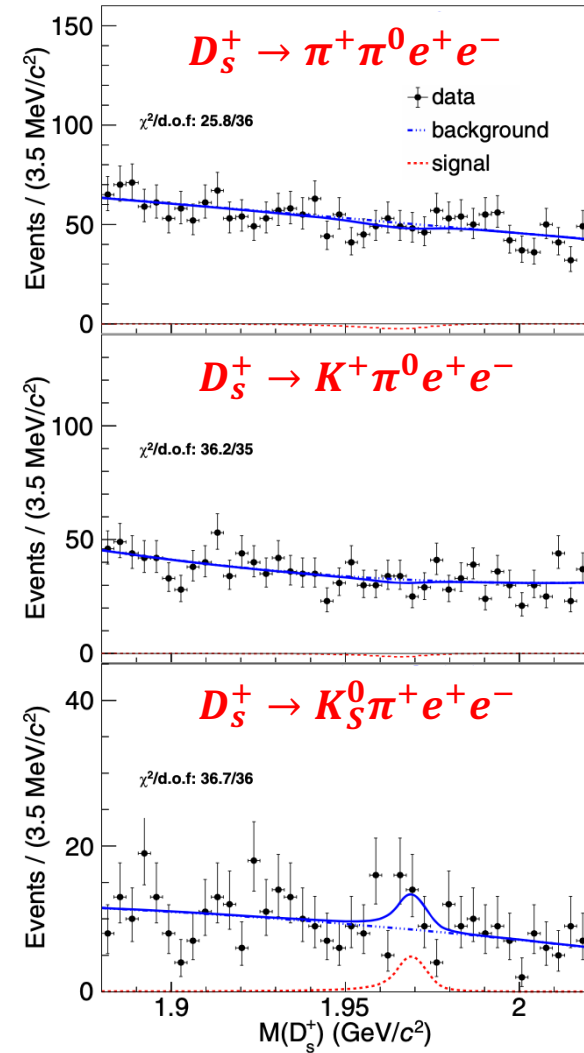
$$B(D_s^+ \rightarrow K_S^0 \pi^+ e^+e^-) < 8.1 \times 10^{-5} @ 90\% \text{ C. L.}$$

Out of ϕ mass window to reject potential LD effect.

First upper limits on the BF's



PRL133(2024)121801





$D_s^+ \rightarrow \gamma \rho(770)^+$

$e^+e^- \rightarrow D_s^{*+}D_s^-$ with 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$

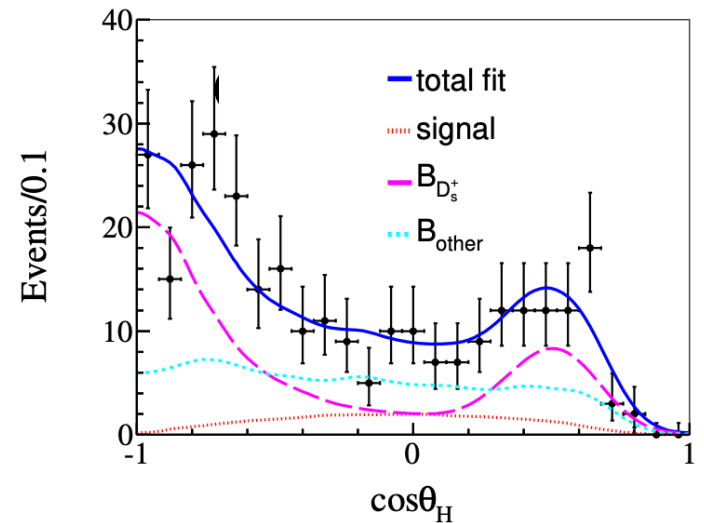
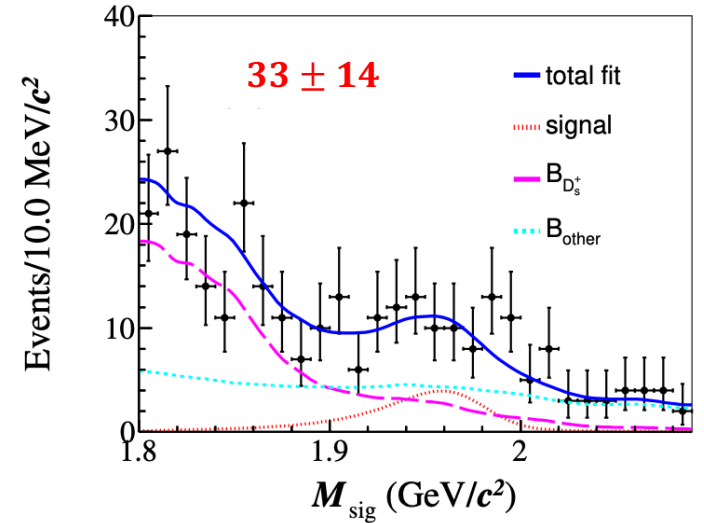
- In the rare radiative decays of charmed mesons, the short distance interaction of the $c \rightarrow u\gamma$ process make a negligible contribution.
- The long distance, non-perturbative processes dominate these decays, potentially enhancing BF's up to 10^{-3} .
- To test the QCD-based calculation of the long-distance dynamics.

$$B(D_s^+ \rightarrow \gamma \rho(770)^+) = (2.2 \pm 0.9_{\text{stat}} \pm 0.2_{\text{syst}}) \times 10^{-4}$$

3.8 σ smaller than the theoretical calculation of weak annihilation.

$$B(D_s^+ \rightarrow \gamma \rho(770)^+) < 6.1 \times 10^{-4} \text{ @90\% C.L.}$$

arXiv:2408.03980 [hep-ex]





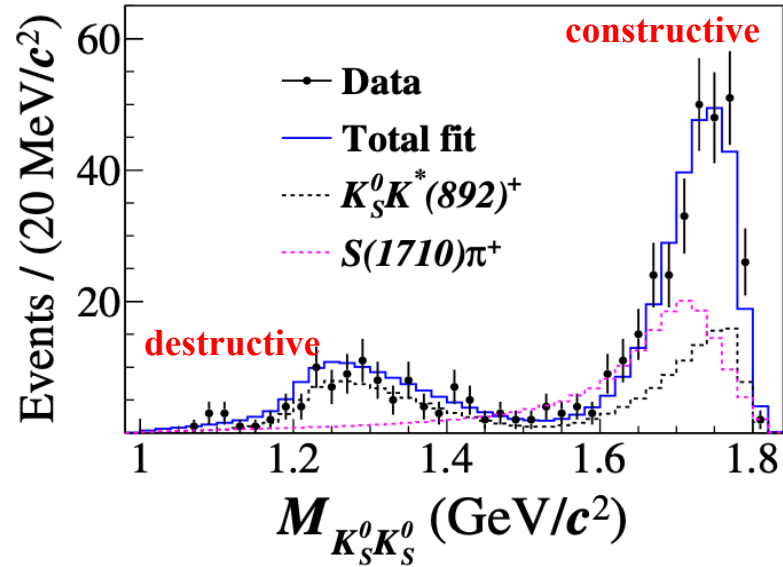
Summary

- **Based on 20 fb^{-1} $\psi(3770)$ data and 7.33 fb^{-1} @ $\sqrt{s} = 4.128 - 4.226 \text{ GeV}$ data, BESIII have reported many important results on the charmed mesons: the amplitude analysis for the hadronic decays, dynamic studies for the semi-leptonic decays, precision measurements on the pure leptonic decays, the rare decays, DCS decays, and so on.**
- **More precise measurements for about $D^{0(+)}$ decays are coming soon based on the 20 fb^{-1} $\psi(3770)$ data.**

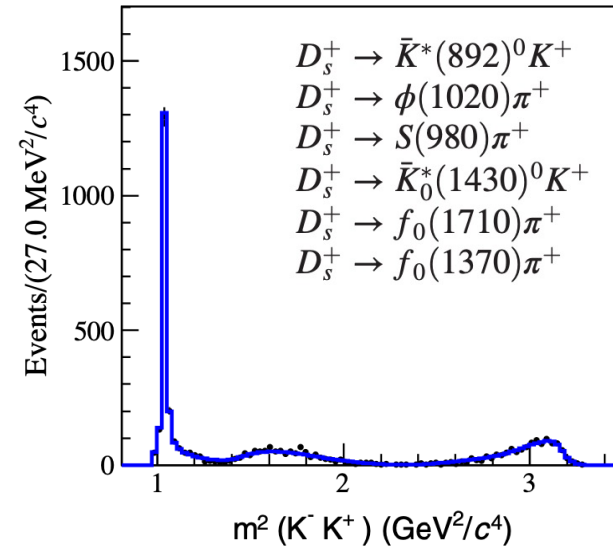
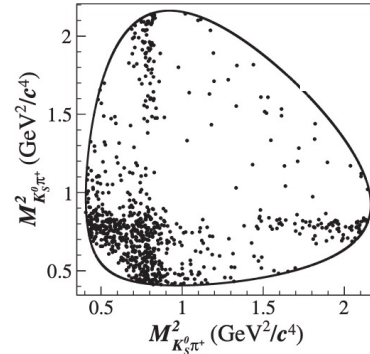
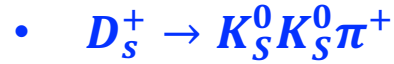
Backup



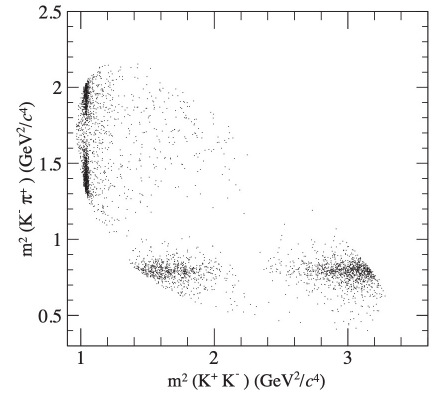
Observation of $a_0(1817)$ in D_s^+ decay



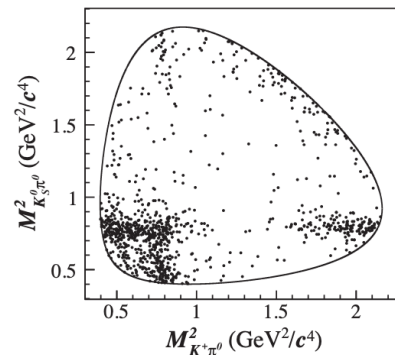
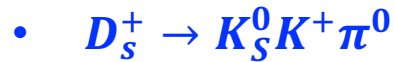
PRD105(2022)L051103



PRD104(2021)012016



PRL129(2022)182001



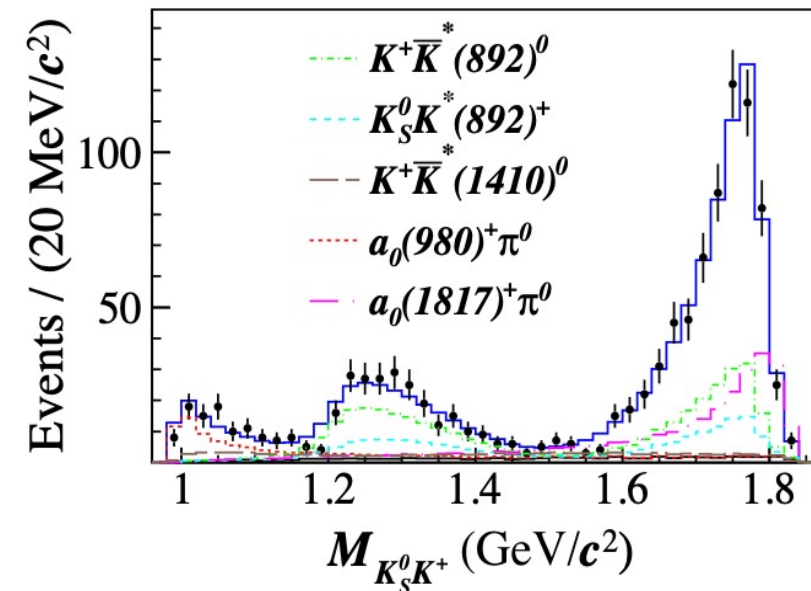
$$M(a_0(1817)^+) = 1.817 \pm 0.008 \pm 0.020 \text{ GeV}/c^2$$

$$\Gamma(a_0(1817)^+) = 0.097 \pm 0.022 \pm 0.015 \text{ GeV}/c$$

$$B(D_s^+ \rightarrow a_0(1817)^+ \pi^0) = (3.44 \pm 0.52 \pm 0.32) \times 10^{-3},$$

Stat. significance: $> 10\sigma$

$a_0(1817)^+$: the isovector partner of $f_0(1710)$ or X(1812)?

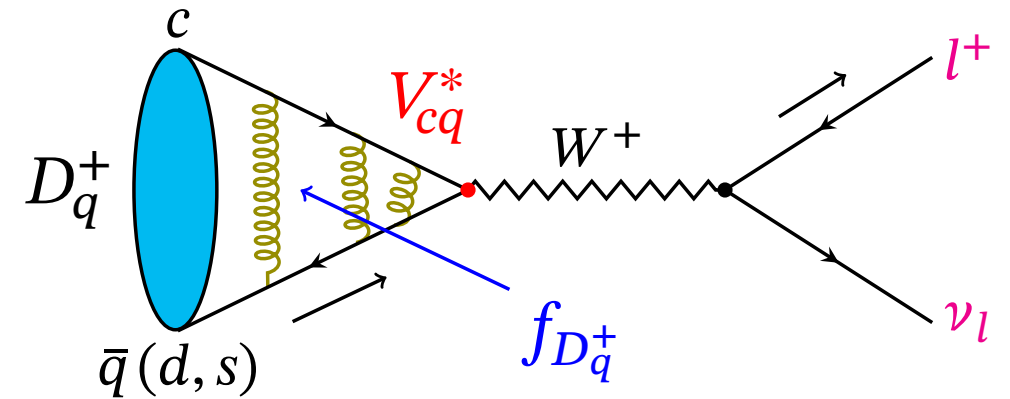


Pure leptonic $D_{(s)}^+$ decays

- Axial current (nonperturbative)

$$\langle 0 | \bar{s} \gamma_\mu \gamma_5 c | D_s^+ \rangle = i p_\mu f_{D_{(s)}^+}$$

($f_{D_{(s)}^+}$: $D_{(s)}^+$ decay constant)



- $V_{cd(s)}$: CKM matrix element, fundamental SM parameter, measured from experiments.

- In the SM, ignoring radiative corrections, the decay width:

$$\Gamma = \frac{G_F^2}{8\pi} f_{D_{(s)}^+}^2 |V_{cd(s)}|^2 m_\ell^2 m_{D_{(s)}^+} \left(1 - \frac{m_\ell^2}{m_{D_{(s)}^+}^2} \right)^2 \quad (\ell = e, \mu, \tau)$$

➤ **Branching fraction:** $Br(D_{(s)}^+ \rightarrow \ell^+ \nu_\ell) = \Gamma \cdot \tau_{D_{(s)}^+}$ ($\tau_{D_{(s)}^+}$: lifetime)

$$Br(D_{(s)}^+ \rightarrow \ell^+ \nu_\ell) \propto \left[f_{D_{(s)}^+} \cdot |V_{cd(s)}| \right]^2$$

Input $f_{D_{(s)}^+}$ from LQCD calculations

Input $|V_{cd(s)}|$ from SM global fit

Extract $|V_{cd(s)}|$ (Direct measurements)

Extract $f_{D_{(s)}^+}$

□ The decay constants $f_{D_{(s)}^+}$: help to calibrate the LQCD calculations;

□ The CKM matrix elements $|V_{cd(s)}|$: help to test the unitarity of CKM.

On the other side, although the $f_0(1710)$ is well established in the PDG [4], its quark structure is still unclear. The two most possible candidates are a glueball or a vector-vector, i.e., $K^*\bar{K}^*$, molecule [1,5,6]. If the $f_0(1710)$ is a $K^*\bar{K}^*$ molecule, there should be an isovector partner $a_0(1710)$, just as the nearly degenerate $f_0(980)$ and $a_0(980)$ for the case of $K\bar{K}$ molecules. If it is a glueball, a second state as pure glueballs must be isoscalar. Therefore, whether $a_0(1710)$ exists is crucial to understand the picture of $f_0(1710)$.

Both $a_0(980)$ and $f_0(980)$ decays to $K^+ K^-$. Impossible to separate them here

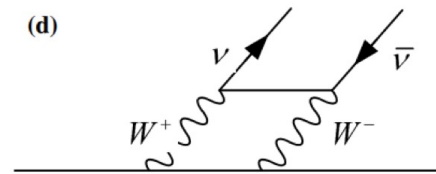
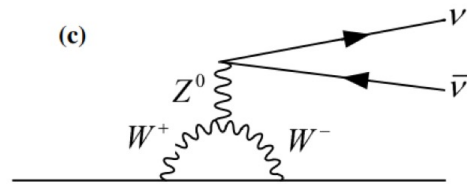
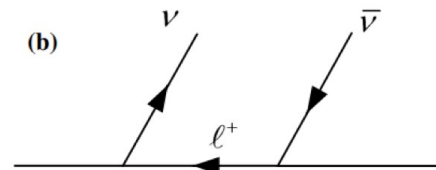
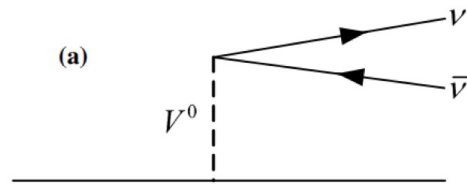
Isospin configurations:

$$a_0(980) \quad I=1 \rightarrow (|K^+ K^- \rangle - |K^0 \overline{K^0} \rangle)$$

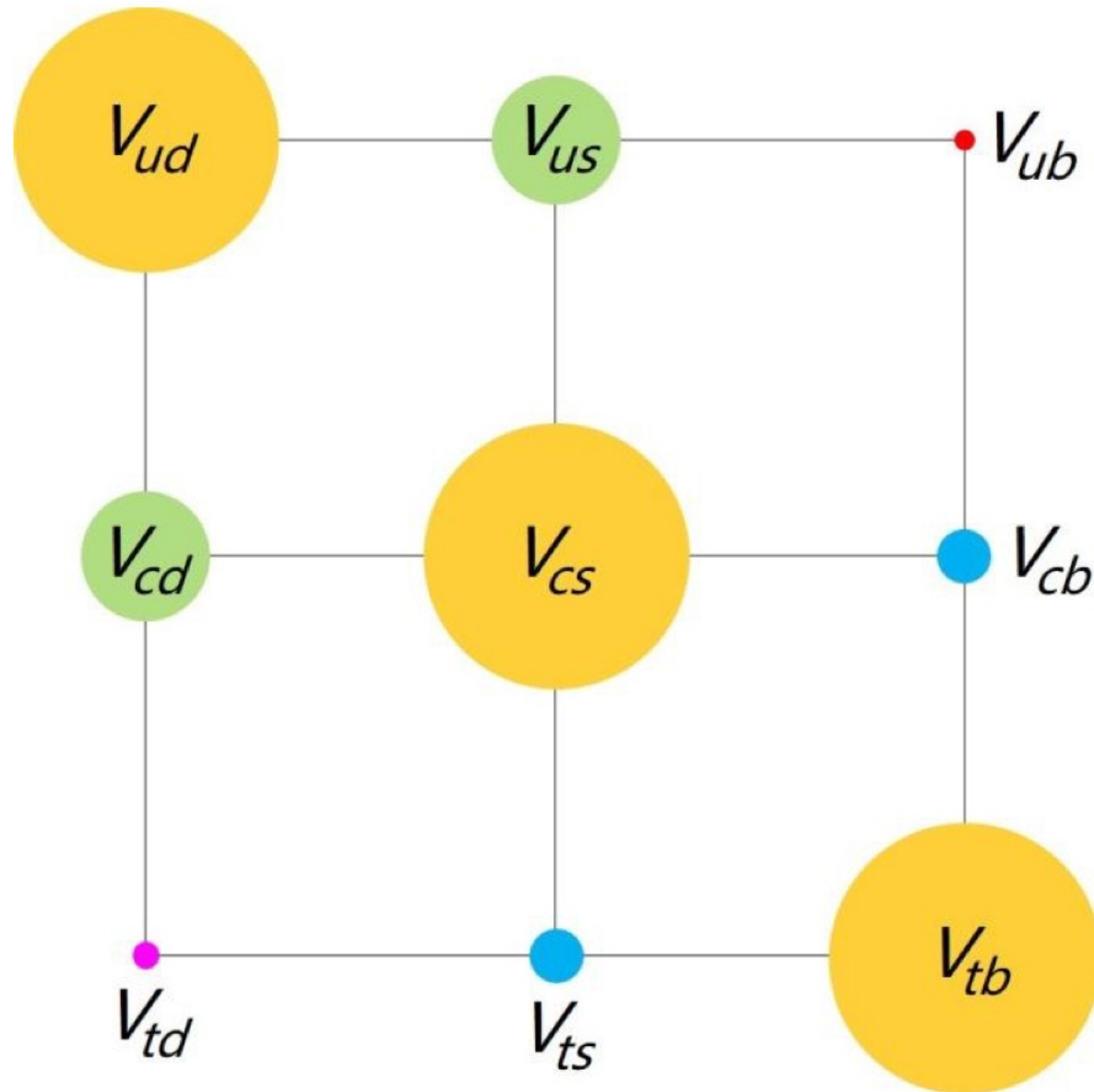
$$f_0(980) \quad I=0 \rightarrow (|K^+ K^- \rangle + |K^0 \overline{K^0} \rangle)$$

The comparison of $K^+ K^-$ and $K_S^0 K_S^0$ spectrum will reveal more information!

➤ Search for the flavor changing neutral current decay $D^0 \rightarrow \pi^0 \nu \bar{\nu}$



LD and SD contributions (a, c, d) to neutral D decays (a, c, d) and charged D decays (a, b, c, d).



Why charm meson physics?

Quantum correlation $D^0\bar{D}^0$: $|\psi(3770)\rangle \rightarrow \frac{1}{\sqrt{2}}(|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle)$

➔ Provide direct access to the $D^0 - \bar{D}^0$ strong-phase difference

- ✓ Important input in CKM γ measurement
- ✓ Precise test of perturbative QCD calculations in charm decays, mixing and CPV

Different methods depending on the final states of D decays

- **GLW** : D decaying to CP eigenstates
- **ADS** : D decaying to CF/DCS eigenstates
- **GGSZ** : D decaying to self-conjugate eigenstates

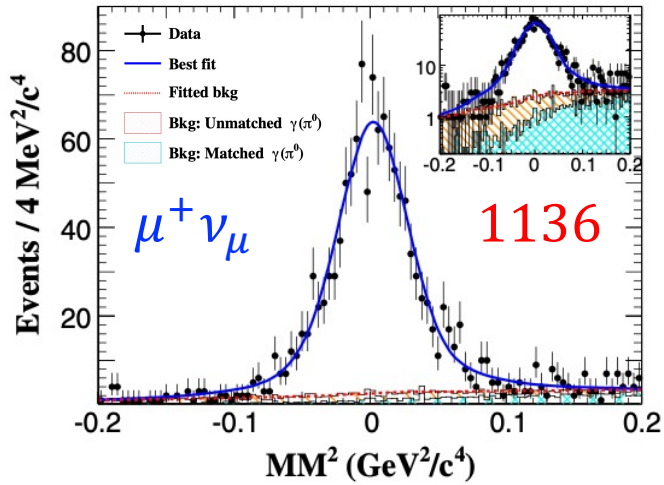
Flavour	$K^\pm \pi^\mp \pi^\mp \pi^-, K^\pm \pi^\mp \pi^0, K^\pm \pi^\mp, \dots$
CP -even	$K^+ K^-, \pi^+ \pi^-, \pi^0 \pi^0, K_S^0 \pi^0 \pi^0, K_L^0 \pi^0, K_L^0 \omega, \pi^+ \pi^- \pi^{0\dagger}$
CP -odd	$K_S^0 \pi^0, K_S^0 \eta, K_S^0 \omega, K_S^0 \eta', K_S^0 \phi, K_L^0 \pi^0 \pi^0$
Self-conjugate	$K_S^0 \pi^+ \pi^-, K_S^0 K^+ K^-, \dots$



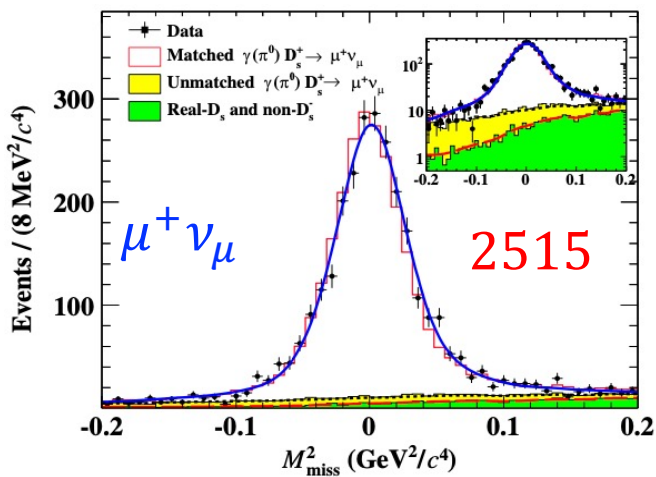
$D_s^+ \rightarrow \ell^+ \nu_\ell$ ($\ell = \mu, \tau$)

- $\mu^+ \nu_\mu$, muon counter

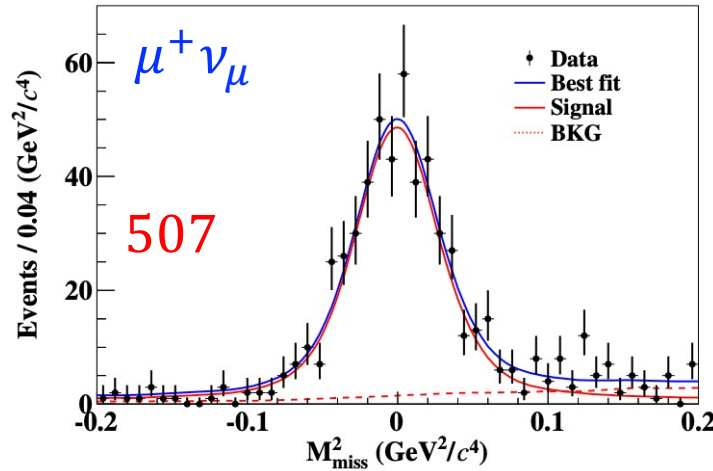
@4.178 GeV PRL122(2019)071802



@4.128 – 4.226 GeV PRD108(2023)112001

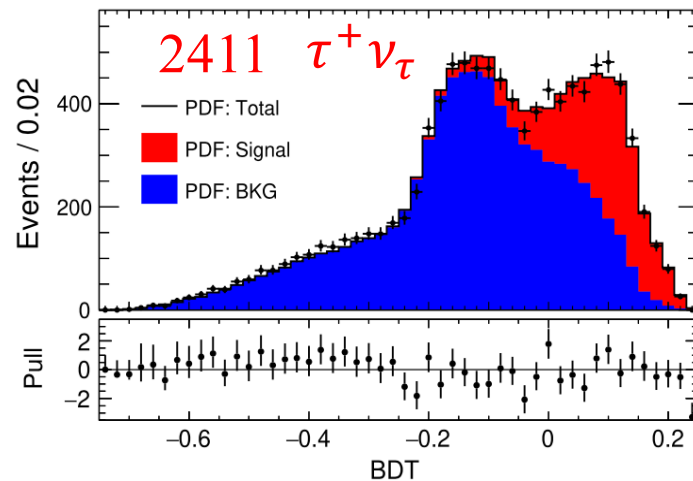


@4.237 – 4.699 GeV arXiv:2407.11727 [hep-ex]



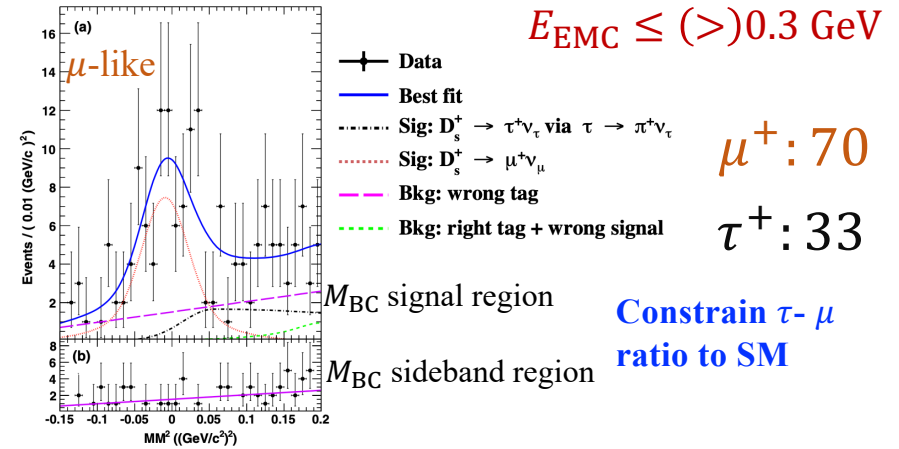
- $\tau^+ \rightarrow \pi^+ \bar{\nu}_\tau$ multivariate analysis

@4.128 – 4.226 GeV PRD108(2023)092014

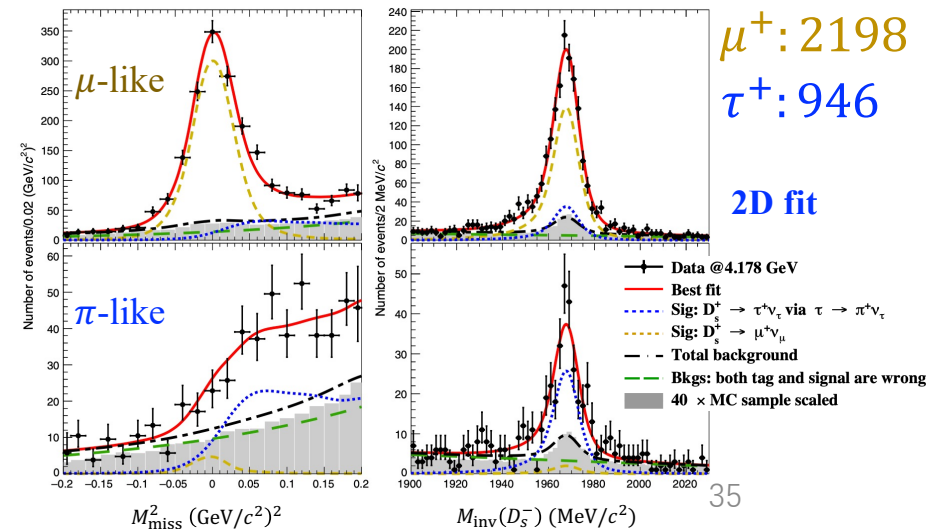


- $\mu^+ \nu_\mu, \tau^+ (\rightarrow \pi^+ \bar{\nu}_\tau) \nu_\tau$

@4.009 GeV PRD94(2016)072004 $\mu(\pi)$ -like:



@4.178 – 4.226 GeV PRD104(2021)052009





$$D_s^+ \rightarrow \tau^+ \nu_\tau$$

$e^+e^- \rightarrow D_s^{*+}D_s^{*-}$ @ 4.237 – 4.699 GeV

Simultaneous fit with the same BF

PRD110(2024)052002

E_{cm} (GeV)	\mathcal{L} (pb ⁻¹)	M_{BC} (GeV/c ²)	N_{ST}
4.237	530.3	(2.107, 2.117)	6477±163
4.246	593.9	(2.107, 2.118)	11944±246
4.260	828.4	(2.107, 2.118)	21550±320
4.270	531.1	(2.107, 2.118)	13319±244
4.280	175.7	(2.106, 2.119)	4063±152
4.290	502.4	(2.106, 2.119)	9316±221
4.310-4.315	546.3	(2.106, 2.119)	5758±228
4.400	507.8	(2.106, 2.119)	1855±87
4.420	1090.7	(2.106, 2.121)	14890±443
4.440	569.9	(2.106, 2.121)	9699±443
4.470-4.699	4768.3	(2.104, 2.123)	25156±762
Sum			124027±1121

$$N_{DT} = 2845 \pm 83$$

$$B_{D_s^+ \rightarrow \tau^+ \nu_\tau} = (5.60 \pm 0.16 \pm 0.20)\%$$

