

Doubly heavy tetraquark states with various configurations

Wei-Lin Wu (吴伟霖)

Peking University

Apr 13th, 2025

Based on: Phys.Rev.D 110 (2024) 9, 094041

Coauthors: Yao Ma, Yan-ke Chen, Lu Meng, Shi-Lin Zhu

- 1 Introduction
- 2 Theoretical Framework
- 3 Doubly heavy tetraquark states
- 4 Summary

Introduction

Experimentally,

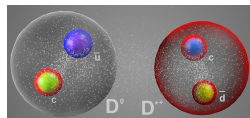
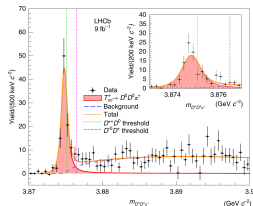
- The first doubly charmed tetraquark $T_{cc}(3875)^+$ discovered in the $D^0 D^0 \pi^+$ channel [LHCb:2021vvq]

$$m_{T_{cc}^+} - (m_{D^{*+}} + m_{D^0}) \sim -300 \text{ keV}$$

- $X(3872) \rightarrow XYZ$ states
 $T_{cc}(3875)^+ \rightarrow$ doubly heavy exotic states?

Theoretically,

- $T_{cc}^+(3875)$: natural interpretation as the $D^{*+} D^0$ molecular state
- Other doubly heavy tetraquarks: compact tetraquark / hadronic molecule



Quark Model + Effective Few-Body Methods:

A possible framework for unified descriptions

- 1 Introduction
- 2 Theoretical Framework**
- 3 Doubly heavy tetraquark states
- 4 Summary

Constituent Quark Model

- Success in conventional hadrons → **nontrivial** extension to multiquark
 - Richer color structure

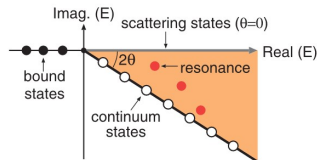
$$\mathbf{3}_q \otimes \mathbf{3}_q \otimes \bar{\mathbf{3}}_{\bar{q}} \otimes \bar{\mathbf{3}}_{\bar{q}} \rightarrow \mathbf{1}_{[qq]\bar{\mathbf{3}}[\bar{q}\bar{q}]\mathbf{3}} \oplus \mathbf{1}_{[qq]\mathbf{6}[\bar{q}\bar{q}]\bar{\mathbf{6}}}$$
 - Complicated few-body calculations
- Do not make *a priori* assumptions about clustering of quarks → **various** types of multiquark states in a **unified** framework
- AL1 model: one-gluon-exchange + linear confinement [Silvestre-Brac:1996myf]

$$V_{ij} = -\frac{3}{16} \lambda_i \cdot \lambda_j \left(-\frac{\kappa}{r_{ij}} + \lambda r_{ij} - \Lambda + \frac{8\pi\kappa'}{3m_i m_j} \frac{\exp(-r_{ij}^2/r_0^2)}{\pi^{3/2} r_0^3} \mathbf{s}_i \cdot \mathbf{s}_j \right)$$

- Theoretical uncertainties $\sim$ tens of MeV

Complex Scaling Method

- Resonant states: poles at $E = m - i\frac{\Gamma}{2}$, not square-integrable
- Previous works: states above thresholds with real energies — possible false states
- Complex scaling method: transform the asymptotic behavior of wave functions by analytical continuation
[Aguilar:1971ve, Balslev:1971vb]



$$U(\theta)\mathbf{r} = \mathbf{r}e^{i\theta} \quad U(\theta)\mathbf{p} = \mathbf{p}e^{-i\theta}$$

$$H(\theta) = \sum_{i=1}^4 \left(m_i + \frac{p_i^2 e^{-2i\theta}}{2m_i} \right) + \sum_{i<j=1}^4 V_{ij}(r_{ij}e^{i\theta})$$

$$H(\theta)\Psi(\theta) = E(\theta)\Psi(\theta)$$

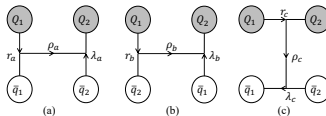
Solve bound, resonant and continuum states simultaneously!

Wave Function

- Color-spin wave functions: complete set

$$\chi_{\bar{3}_c \otimes 3_c}^{s_1, s_2, S} = \left[(Q_1 Q_2)_{\bar{3}_c}^{s_1} (\bar{q}_1 \bar{q}_2)_{3_c}^{s_2} \right]_{1_c}^S \quad \chi_{6_c \otimes \bar{6}_c}^{s_1, s_2, S} = \left[(Q_1 Q_2)_{6_c}^{s_1} (\bar{q}_1 \bar{q}_2)_{\bar{6}_c}^{s_2} \right]_{1_c}^S$$

- Spatial wave functions: Gaussian expansion method [Hiyama:2003cu]

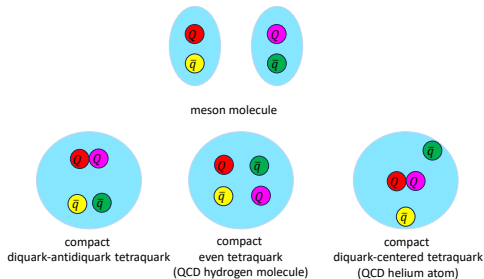


$$\Phi_{n_1, n_2, n_3}^{(\text{jac})} = \phi_{n_1}(r_{\text{jac}}) \phi_{n_2}(\lambda_{\text{jac}}) \phi_{n_3}(\rho_{\text{jac}}) \quad \phi_{n_i}(r) = N_{n_i} e^{-\nu_{n_i} r^2}$$

- dimeson** (a,b) and **diquark-antidiquark** (c) configurations
- consistently describe **molecular**, **compact** states and their couplings
- $\sim 10^4$ bases to solve four-body Schrödinger equation with strong correlations in color, spin and spatial d.o.f.

Spatial Structures

- Meson molecules & compact tetraquarks: different internal structures and binding mechanisms
- 3 typical types of compact tetraquarks expected in $QQ\bar{q}\bar{q}$
- Identification by rms radius



$$r_{ij}^{\text{rms}} \equiv \text{Re} \left[\sqrt{\frac{\langle \Psi_{\text{nA}}(\theta) | r_{ij}^2 e^{2i\theta} | \Psi_{\text{nA}}(\theta) \rangle}{\langle \Psi_{\text{nA}}(\theta) | \Psi_{\text{nA}}(\theta) \rangle}} \right]$$

A novel definition to avoid ambiguities from antisymmetrization

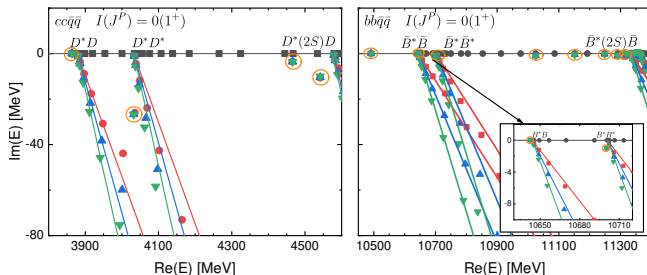
- 1 Introduction
- 2 Theoretical Framework
- 3 Doubly heavy tetraquark states**
- 4 Summary

Bound states

System	$I(J^P)$	M	ΔE	$\chi_{\bar{3}_c \otimes 3_c}$	$\chi_{6_c \otimes \bar{6}_c}$	$r_{Q_1 \bar{q}_1}^{\text{rms}}$	$r_{Q_2 \bar{q}_2}^{\text{rms}}$	$r_{Q_1 \bar{q}_2}^{\text{rms}}$	$r_{Q_2 \bar{q}_1}^{\text{rms}}$	$r_{Q_1 Q_2}^{\text{rms}}$	$r_{\bar{q}_1 \bar{q}_2}^{\text{rms}}$	Configuration
$cc\bar{q}\bar{q}$	$0(1^+)$	3864	-14	58%	42%	0.71	0.64	1.13	1.16	1.02	1.22	M. (D^*D)
$bb\bar{q}\bar{q}$	$0(1^+)$	10642	-1	33%	67%	0.66	0.63	2.06	2.07	1.98	2.15	M. ($\bar{B}^*\bar{B}$)
$bc\bar{q}\bar{q}$	$0(2^+)$	7363	-3	27%	73%	0.66	0.70	1.95	1.97	1.86	2.05	M. ($\bar{B}^*D^*$)
$bc\bar{q}\bar{q}$	$0(0^+)$	7129	-26	48%	52%	0.64	0.64	0.91	0.95	0.76	1.03	C.E.
	$0(1^+)$	7185	-27	60%	40%	0.67	0.66	0.88	0.93	0.71	1.00	C.E.
$bb\bar{q}\bar{q}$	$0(1^+)$	10491	-153	97%	3%	0.68	0.67	0.70	0.71	0.33	0.78	C.DC.
$bb\bar{s}\bar{q}$	$\frac{1}{2}(1^+)$	10647	-64	91%	9%	0.56	0.67	0.71	0.61	0.36	0.76	C.DC.

- D^*D molecule with $\Delta E = -14$ MeV as candidate for $T_{cc}(3875)^+$
- Bound states with various configurations
 - Molecular (M.) shallow bound state
 - Compact even tetraquark (C.E.)
 - Compact diquark-centered tetraquark (C.DC.) deeply bound state
- No isovector bound state $\rightarrow$ importance of “good” antidiquark $(\bar{q}\bar{q})_{3_c}^{S=0, I=0}$

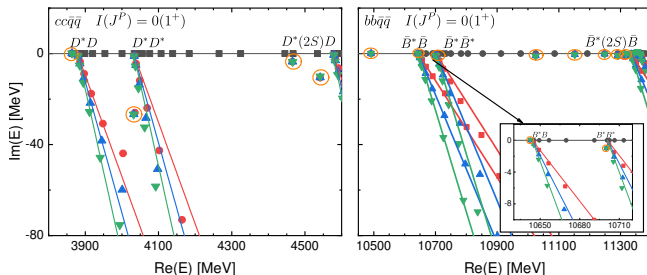
Resonant states: isoscalar $QQ\bar{q}\bar{q}$



- Lowest resonant states near the D^*D^* and $\bar{B}^*\bar{B}^*$ thresholds
- First T_{bb} resonance as $\bar{B}^*\bar{B}^*$ molecule

$M - i\Gamma/2$	$\chi_{3_c \otimes 3_c}$	$\chi_{6_c \otimes 6_c}$	$r_{Q_1\bar{q}_1}^{\text{rms}}$	$r_{Q_2\bar{q}_2}^{\text{rms}}$	$r_{Q_1\bar{q}_2}^{\text{rms}}$	$r_{Q_2\bar{q}_1}^{\text{rms}}$	$r_{Q_1Q_2}^{\text{rms}}$	$r_{\bar{q}_1\bar{q}_2}^{\text{rms}}$	Config.
10491	97%	3%	0.68	0.67	0.70	0.71	0.33	0.78	C.DC.
10642	33%	67%	0.66	0.63	2.06	2.07	1.98	2.15	M. ($\bar{B}^*\bar{B}$)
10700 - 1i	44%	56%	0.67	0.67	1.96	1.96	1.88	2.02	M. ($\bar{B}^*\bar{B}^*$)
11025 - 1i	98%	2%	1.08	1.07	1.08	1.08	0.33	0.83	C.DC.

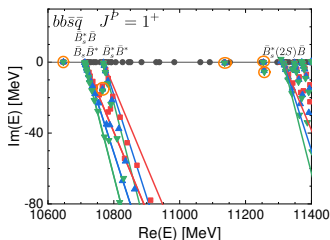
Resonant states: isoscalar $QQ\bar{q}\bar{q}$



- Second T_{bb} resonance as radial excitation of the deeply bound T_{bb} state
- More resonant states in higher energy region and $bc\bar{q}\bar{q}$ system

$M - i\Gamma/2$	$\chi_{3_c \otimes 3_c}$	$\chi_{6_c \otimes 6_c}$	$r_{Q_1 \bar{q}_1}^{\text{rms}}$	$r_{Q_2 \bar{q}_2}^{\text{rms}}$	$r_{Q_1 \bar{q}_2}^{\text{rms}}$	$r_{Q_2 \bar{q}_1}^{\text{rms}}$	$r_{Q_1 Q_2}^{\text{rms}}$	$r_{\bar{q}_1 \bar{q}_2}^{\text{rms}}$	Config.
10491	97%	3%	0.68	0.67	0.70	0.71	0.33	0.78	C.DC.
10642	33%	67%	0.66	0.63	2.06	2.07	1.98	2.15	M. ($\bar{B}^* \bar{B}$)
10700 - 1i	44%	56%	0.67	0.67	1.96	1.96	1.88	2.02	M. ($\bar{B}^* \bar{B}^*$)
11025 - 1i	98%	2%	1.08	1.07	1.08	1.08	0.33	0.83	C.DC.

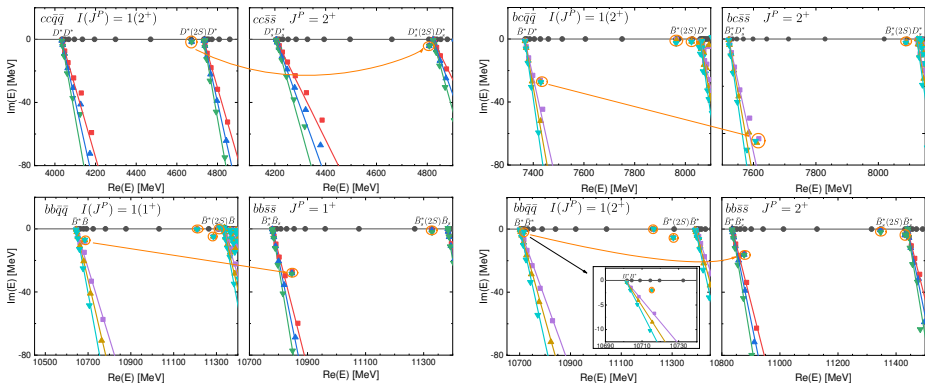
Resonant states: $QQ\bar{s}\bar{q}$



$M - i\Gamma/2$	$\chi_{\bar{3}_c \otimes \bar{3}_c}$	$\chi_{6_c \otimes \bar{6}_c}$	$r_{Q_1 \bar{s}}^{\text{rms}}$	$r_{Q_2 \bar{q}}^{\text{rms}}$	$r_{Q_1 \bar{q}}^{\text{rms}}$	$r_{Q_2 \bar{s}}^{\text{rms}}$	$r_{Q_1 Q_2}^{\text{rms}}$	$r_{s \bar{q}}^{\text{rms}}$	Config.
10647	91%	9%	0.56	0.67	0.71	0.61	0.36	0.76	C.DC.
10766 - 16i	92%	8%	0.58	0.71	0.74	0.63	0.36	0.85	C.DC.

- Lowest $T_{bb\bar{s}}$ resonant state near $\bar{B}_s^* \bar{B}^*$ threshold
- Similar to $bb\bar{q}\bar{q}$ system: the resonant state is the radial excitation in the light d.o.f of the deeply bound state

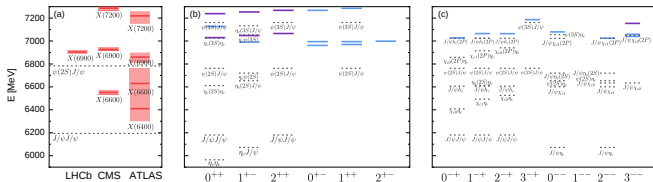
Resonant states: isovector $QQ\bar{q}\bar{q}$ and $QQ\bar{s}\bar{s}$



- A series of compact tetraquark resonant states
- Resemblance between energy spectra of isovector $QQ\bar{q}\bar{q}$ and $QQ\bar{s}\bar{s}$ systems $\rightarrow$ same internal symmetries
- Mass differences between $QQ\bar{q}\bar{q}$ resonances and their strange partners ~ 150 MeV

Other systems

- Successfully describe $T_{cs0}^*(2870)^0$ [Chen:2023syh], $X(6900)$ and $X(7200)$ [Wu:2024euj, Wu:2024hrv]
- Both 0^{++} and 2^{++} candidates for $X(6900)$ and $X(7200)$ are found
- All $cc\bar{c}\bar{c}$ resonances are compact tetraquark states



		BW ₁	BW ₂	BW ₃
Interference	m (MeV)	$6593^{+15}_{-14} \pm 25$	$6847^{+10}_{-10} \pm 15$	$7173^{+9}_{-10} \pm 13$
(Run 2+Run 3)	Γ (MeV)	$446^{+66}_{-54} \pm 87$	$135^{+16}_{-14} \pm 14$	$73^{+18}_{-15} \pm 10$
Interference	m (MeV)	$6638^{+43}_{-38} \pm 31$	$6847^{+44}_{-28} \pm 20$	$7134^{+48}_{-25} \pm 15$
(Run 2 [12])	Γ (MeV)	$440^{+230}_{-200} \pm 110$	$191^{+66}_{-49} \pm 25$	$97^{+40}_{-29} \pm 26$

	J^{PC}	0^{++}	0^{+-}	2^{++}	2^{+-}
Quark Model	m (MeV)	6978	7167	7017	7204
	Γ (MeV)	72	38	78	58

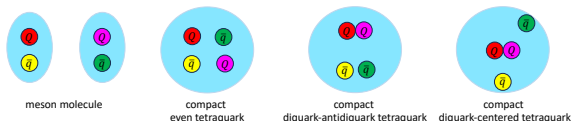
Open questions for $cc\bar{c}\bar{c}$:

- If experimental states are determined to be 2^{++} , where are 0^{++} states?
- Where is $X(6600)$? It's too wide to be calculated.

- 1 Introduction
- 2 Theoretical Framework
- 3 Doubly heavy tetraquark states
- 4 Summary**

Summary

- Framework: QM + CSM + GEM
- Various types of tetraquark states in a unified framework



- A series of doubly heavy tetraquark bound and resonant states
 - D^*D molecular shallow bound state as candidate for T_{cc}^+ (3875)
 - Bound states in $bb\bar{q}\bar{q}$, $bc\bar{q}\bar{q}$, $bbs\bar{q}$ systems
 - Resonant states near D^*D^* and $\bar{B}^*\bar{B}^*$ thresholds, $\dots$
- The capability and limit of QM are to be explored with comprehensive theoretical calculations.

Thanks for your listening

Backup: RMS Radii

- Molecular or compact states: rms radii
- Ambiguities from antisymmetrization:

$$\begin{aligned}
 \Psi &= \mathcal{A}\Psi_{\text{nA}} = \mathcal{A} \sum_{s_1 \geq s_2} [(Q_1 \bar{q}_1)_{1_c}^{s_1} (Q_2 \bar{q}_2)_{1_c}^{s_2}]_{1_c}^S \otimes |\psi_1^{s_1 s_2}\rangle \\
 &= \sum_{s_1 \geq s_2} \left([(Q_1 \bar{q}_1)_{1_c}^{s_1} (Q_2 \bar{q}_2)_{1_c}^{s_2}]_{1_c}^S \otimes |\psi_1^{s_1 s_2}\rangle + [(Q_1 \bar{q}_1)_{1_c}^{s_2} (Q_2 \bar{q}_2)_{1_c}^{s_1}]_{1_c}^S \otimes |\psi_2^{s_1 s_2}\rangle \right. \\
 &\quad \left. + [(Q_1 \bar{q}_2)_{1_c}^{s_1} (Q_2 \bar{q}_1)_{1_c}^{s_2}]_{1_c}^S \otimes |\psi_3^{s_1 s_2}\rangle + [(Q_1 \bar{q}_2)_{1_c}^{s_2} (Q_2 \bar{q}_1)_{1_c}^{s_1}]_{1_c}^S \otimes |\psi_4^{s_1 s_2}\rangle \right)
 \end{aligned}$$

The quarks belong to both constituent mesons, $\langle \Psi | r_{Q\bar{q}}^2 | \Psi \rangle$ cannot reflect the size of the constituent meson.

- Define rms radii using Ψ_{nA}

$$r_{ij}^{\text{rms}} \equiv \text{Re} \left[\sqrt{\frac{\langle \Psi_{\text{nA}}(\theta) | r_{ij}^2 e^{2i\theta} | \Psi_{\text{nA}}(\theta) \rangle}{\langle \Psi_{\text{nA}}(\theta) | \Psi_{\text{nA}}(\theta) \rangle}} \right]$$