

ω meson from lattice QCD

燕浩波

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Based on:

1. Yan *et al.*, PRL 133 (2024) 211906, **Editors' Suggestion**
2. Yan *et al.*, PRD 111 (2025) 014503
3. Yan *et al.*, arXiv 2504.xxxx on OpTion
4. Yan *et al.*, arXiv 250x.xxxx on coupled channel $D_0^*(2300)$

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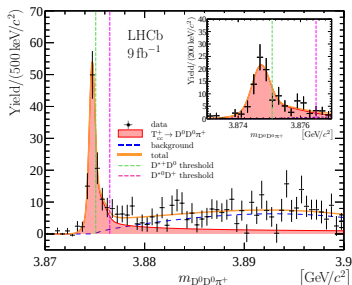
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Introduction

- Many resonances decay into three-body states
- LHCb discovered a tetraquark candidate $T_{cc}(3875) \rightarrow DD^* \rightarrow DD\pi$ in 2022¹



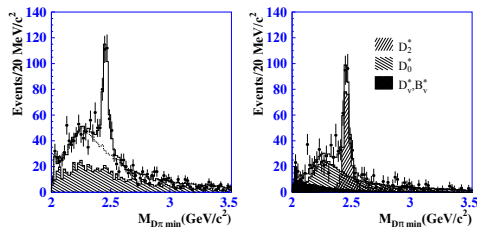
- Previous lattice studies lifted the pion mass and study the two-body DD^* scattering² and found the corresponding pole
- The real **three-body scattering** problem relies on the two-body $D\pi$ scattering amplitudes as an input
- Lattice QCD is the only way to calculate observables at low energy from the first principle
- We need to know how to solve 1. two-body and 2. three-body problems

¹LHCb collaboration, NP 18 (2022) 751

²Padmanath *et al.*, PRL 129 (2022) 032002; Chen *et al.*, PLB 833 (2022) 137391; Lyu *et al.*, PRL 131 (2023) 161901; Collins *et al.*, PRD 109 (2024) 094509

Introduction 1: Two-body: $D_0^* \rightarrow D\pi$

- A broad resonance $D_0^* \rightarrow D\pi$ was found by Belle collaboration³ in 2004



- Puzzles on $D_0^*(2300)$

- $m_{D_0^*} \approx m_{D_{s0}^*}$
- $m_{D_0^*}^{\text{ChPT}} \ll m_{D_0^*}^{\text{exp.}}$

could be explained by UChPT where the two-pole structure is proposed⁴

- Should be tested on the lattice at a series of m_π ⁵

³Satpathy *et al.*, PRB 159 (2003) 553.

⁴For e.g., Albaladejo *et al.*, PLB 767 (2017) 465.

⁵See previous lattice works in [Moir:2016srx, Gayer:2021xzv]

- Turn off the weak and electromagnetic interactions,

$$\mathcal{L}_{\text{QCD}} = \bar{\psi} (i(\gamma^\mu D_\mu) - m) \psi - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu} \rightarrow \text{discretizing}$$

- Two-point function is calculated by

$$\langle \Omega | O(t) O^\dagger(0) | \Omega \rangle = \frac{1}{Z} \int \mathcal{D}[\psi, \bar{\psi}] \mathcal{D}[U] e^{-S_F[\psi, \bar{\psi}, U] - S_G[U]} (O(t) O^\dagger(0)) [\psi, \bar{\psi}, U]$$

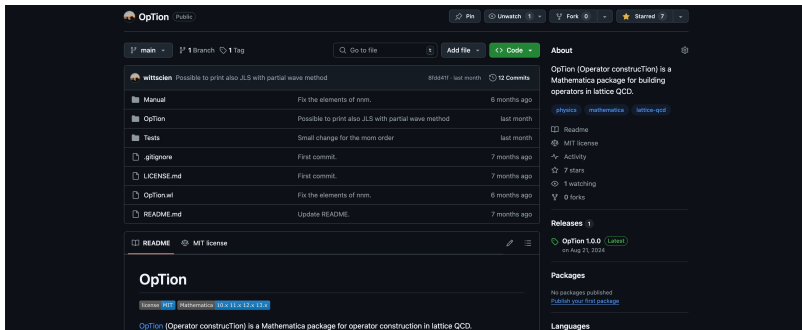
- On the other hand,

$$\begin{aligned} \langle \Omega | O(t) O^\dagger(0) | \Omega \rangle &= \langle \Omega | O(t) \sum_n \frac{1}{2E_n} | E_n \rangle \langle E_n | O^\dagger(0) | \Omega \rangle \\ &= \frac{1}{2E_1} |\langle \Omega | O(t) | E_1 \rangle|^2 e^{-E_1 t} + \text{excited states} \end{aligned}$$

- Towers of finite-volume energy levels E_n encode the interactions

Short introduction to lattice QCD – Operator construction

- Now, how to build these operators O , so that we create $D\pi$ from the vacuum?
- Symmetry broken $O(3) \rightarrow O_h$ or even smaller group
- A package **OpTion**⁶ (**O**perator **con**struc**T**ion) has been developed to construct general N -hadron operators. It is widely used and has become a routine tool in the CLQCD collaboration



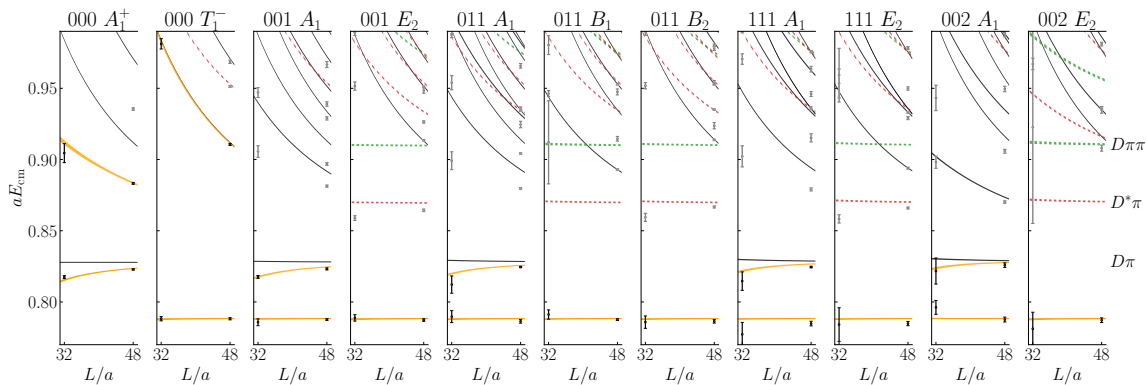
- e.g. $D(+p_z)\pi(-p_z) - D(-p_z)\pi(+p_z)$ transforms as the T_1^- irrep

⁶<https://github.com/wittscien/OpTion>

configuration	volume	a/fm	m_π/MeV	N_{cfgs}
C48P14	$48^3 \times 96$	0.10530(18)	135.5(1.6)	259
F32P21	$32^3 \times 64$	0.07746(18)	210.9(2.2)	459
F48P21	$48^3 \times 96$	0.07746(18)	207.2(1.1)	222
F32P30	$32^3 \times 96$	0.07746(18)	303.2(1.3)	567
F48P30	$48^3 \times 96$	0.07746(18)	303.4(9)	201
H48P32	$48^3 \times 144$	0.05187(26)	317.2(0.9)	274

- CLQCD ensembles with $N_f = 2 + 1$ Clover fermions [CLQCD, 2024]
- 4 different pion masses to track the chiral behavior
- 3 lattice spacings to estimate the discretization error

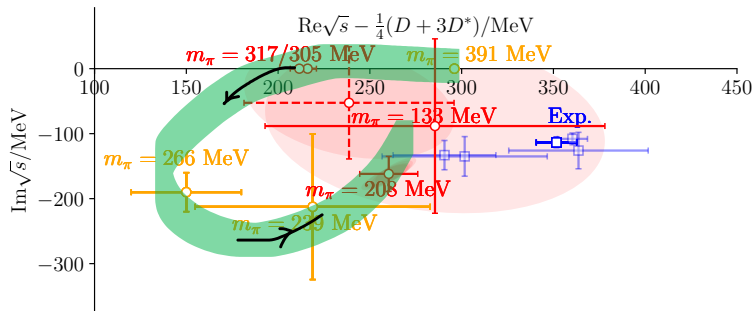
Finite-volume spectra at $m_\pi \approx 208$ MeV



- Strong attraction in S -wave
- The levels are interpreted by quantization conditions: $f(\text{energy levels, amplitude}) = 0$

Pole positions of $D_0^*(2300)$ as a function of m_π

- Using the scattering amplitude, we perform an analytic continuation to find poles



- Only one pole is found $\rightarrow$ two-pole structure is to be found in coupled channel scattering
- An clear trend for the motion of **one** D_0^* pole is identified

bound state $\rightarrow$ virtual state $\rightarrow$ resonance

Introduction 2: Three-body: $\omega \rightarrow \pi\pi\pi$



- Now we have the two-body input – we want to approach the three-body channel
- However, 3bd is too much more complicated (only one lattice calculation considered resonance in the history^a)
- We take a challenge on the ω meson
- crucial for T_{cc}^b , Roper, ...

Why is this problem difficult?

- Three-body problem with resonances in two-body problem
 - ▶ $\pi\pi\pi \rightarrow \omega$
 - ▶ $\pi\pi \rightarrow \rho$
- Computationally most complicated among all isospin channel
- Complicated quantization conditions

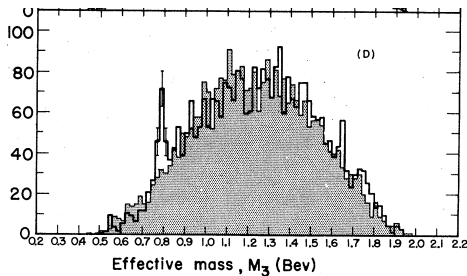
^aMai:2021nul

^bHansen:2024ffk

ω : the first neutral vector meson (1961)¹



PRESS/TV CONFERENCE ON DISCOVERY OF OMEGA MESON
Berkeley, August 31, 1961
Maglic, Alvarez, Rosenfeld & Stevenson, Phys. Rev. Lett. September 1, 1961

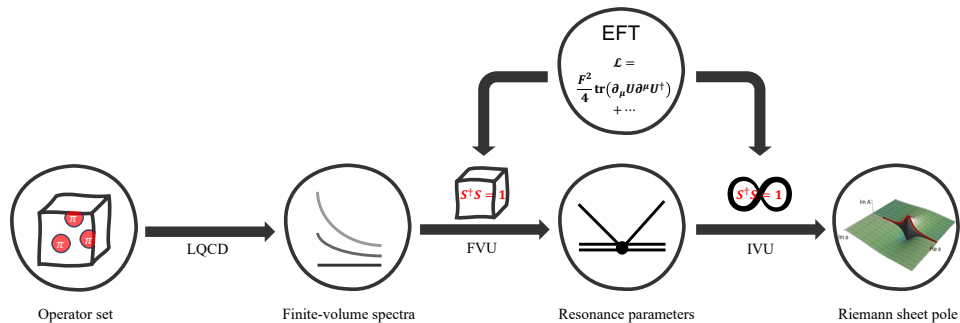


- The lightest hadron decaying into three particles: $\omega \rightarrow 3\pi$
- Phenomenologically interesting, for instance in the $N - N$ interaction, muon $g - 2$, dark matter...⁸

¹Maglic *et al.* (1961).

⁸Sakurai (1960); Erkelenz (1974); Brown and Jackson (1976); Barkov *et al.*, 1985; Connell *et al.* (1997); Bazavov *et al.* (2021).

The methodology

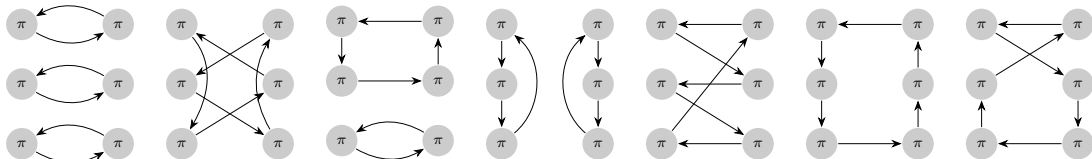


- $\omega, \rho\pi, \pi\pi\pi$ operator from OpTion²
- 2bd and 3bd spectra
- Develop FVU and the pertinent EFT to parametrize the three-body force
- Solve the integral equations and search for the ω poles

²<https://github.com/wittscien/OpTion>

Contraction & quantization conditions

- Insanely many diagrams (202 for only $\pi\pi\pi \rightarrow \pi\pi\pi$, only 9 for the two-body problem)
- The topologies for $\pi\pi\pi \rightarrow \pi\pi\pi$



- Quantization conditions: **FVU** (Finite-Volume Unitarity)

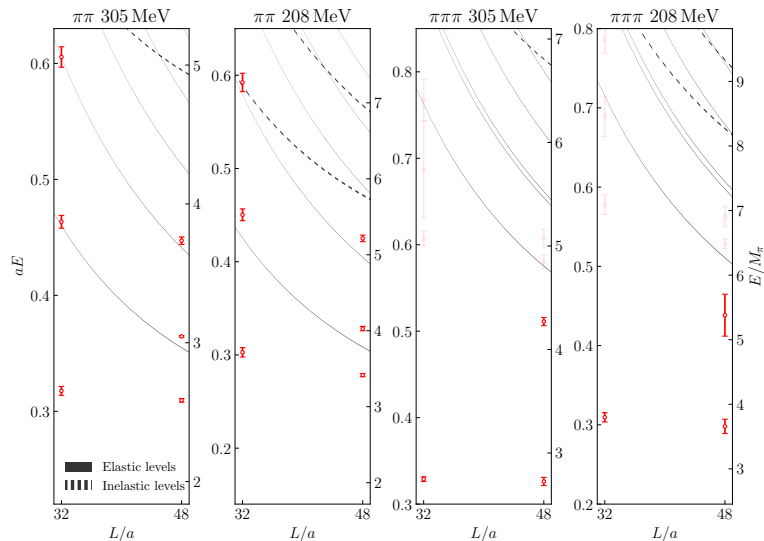
$$\begin{cases} \tilde{K}^{-1}(\sigma) - \Sigma^{FV}(\sigma) = 0 \\ \det[(\tilde{K}^{-1}(s) - \Sigma^{FV}(s))E_L - (\tilde{B}(s) + \tilde{C}(s))] = 0 \end{cases}$$

- Combined fit** for $\pi\pi$ and $\pi\pi\pi$ spectra for all M_π 's (**EFT4**)³

$$\left\{ \begin{array}{l} \text{[Vertex diagram]} \sim \frac{\sigma - M_\rho^2}{g^2} \\ \text{[Loop diagram]} \sim \frac{s(M_\rho^2 - \sigma_q + 6g^2 f_\pi^2)(M_\rho^2 - \sigma_p + 6g^2 f_\pi^2)}{g^2 f_\pi^6 (s - M_\omega^2)} \end{array} \right.$$

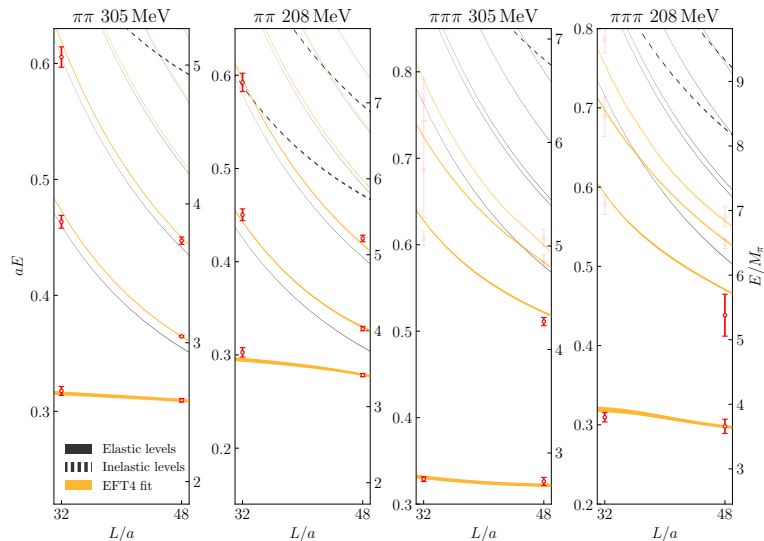
³For a review of the EFT form, see Meißner (1988).

Finite-volume spectra revisited



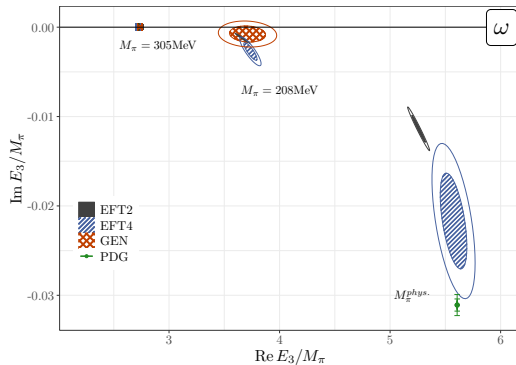
- Strong attraction in both the $\pi\pi$ and $\pi\pi\pi$ channels
- Restricted to be below the $\omega(1420)$ region, but high-lying energies are also well-predicted

Finite-volume spectra revisited



- Strong attraction in both the $\pi\pi$ and $\pi\pi\pi$ channels
- Restricted to be below the $\omega(1420)$ region, but high-lying energies are also well-predicted

Pole positions of $\omega(782)$ as a function of m_π



- Solve the integral equation
- ω is a **bound state** at $M_\pi \approx 305$ MeV and a **resonance** at $M_\pi \approx 208$ MeV
- Extrapolate to the **physical pion mass**, the poles agree well with the PDG values⁴

⁴See related discussions in Hoferichter (2023), Hoferichter (2019), Hoid (2020), Colangelo (2022), Colangelo (2018).

Summary

For $D_0^*(2300)$,

- An clear trend for the motion of one D_0^* pole is identified

bound state $\rightarrow$ virtual state $\rightarrow$ resonance

- Coupled $D\pi-D\eta-D_s\bar{K}$ analysis is expected soon
- The package OpTion is developed

For $\omega(782)$,

- First-ever determination of the ω -meson pole from lattice QCD
- Development of the FVU, matching EFT and FVU
- Paved the way to study heavier three-hadron resonances
- The ρ and ω pole positions at the physical point

$$\sqrt{s_\rho} = (748.9(10.0) - i63.5(1.8)) \text{ MeV}$$

$$\sqrt{s_\omega} = (778.0(11.2) - i3.0(5)) \text{ MeV}$$

Thank you!

@Maxim Mai and Marco Garofalo

