

# The Workshop on Grand Unified Theory, Phenomenology and Cosmology (GUTPC 2025)

## Leptogenesis via symmetry non-restoration

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**Wei Chao**

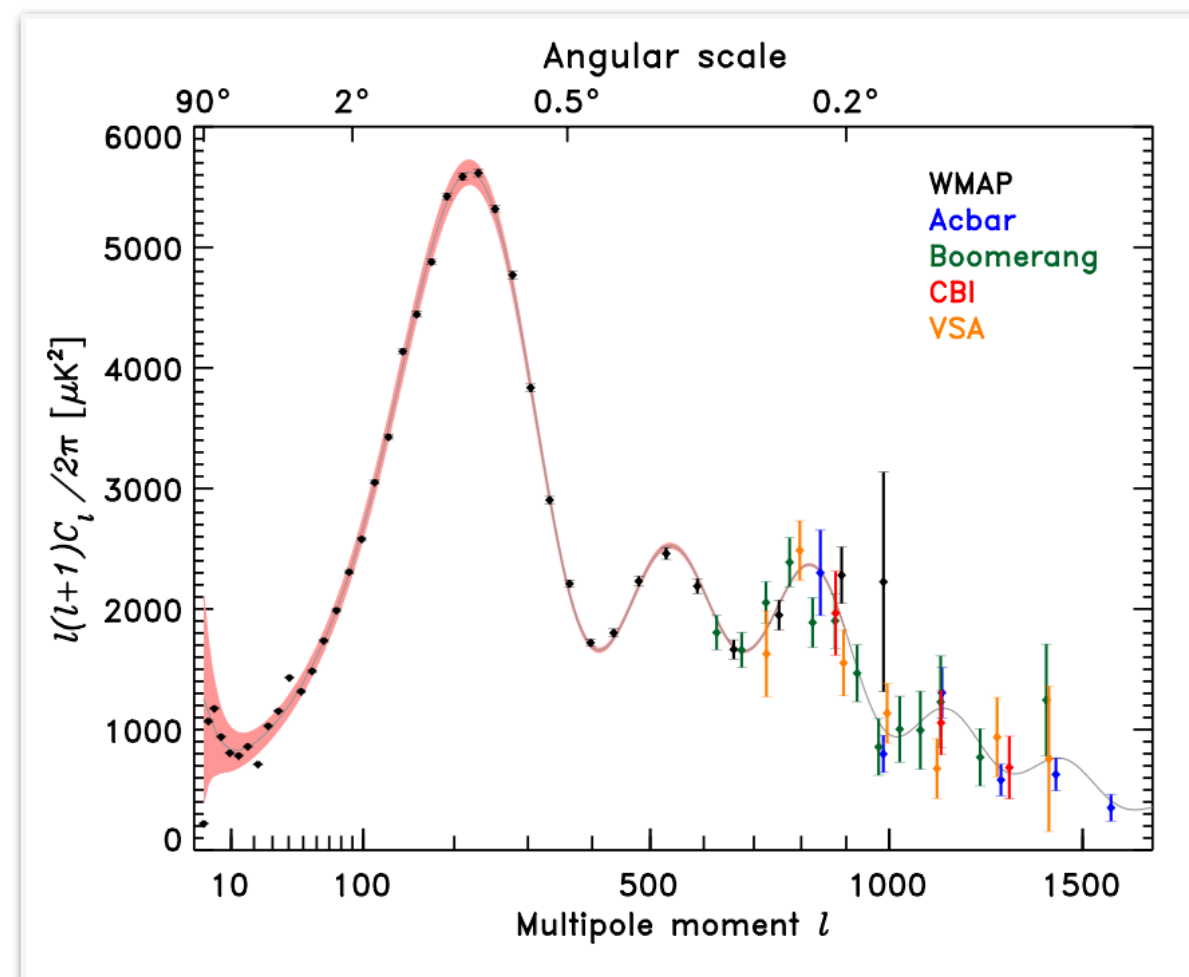
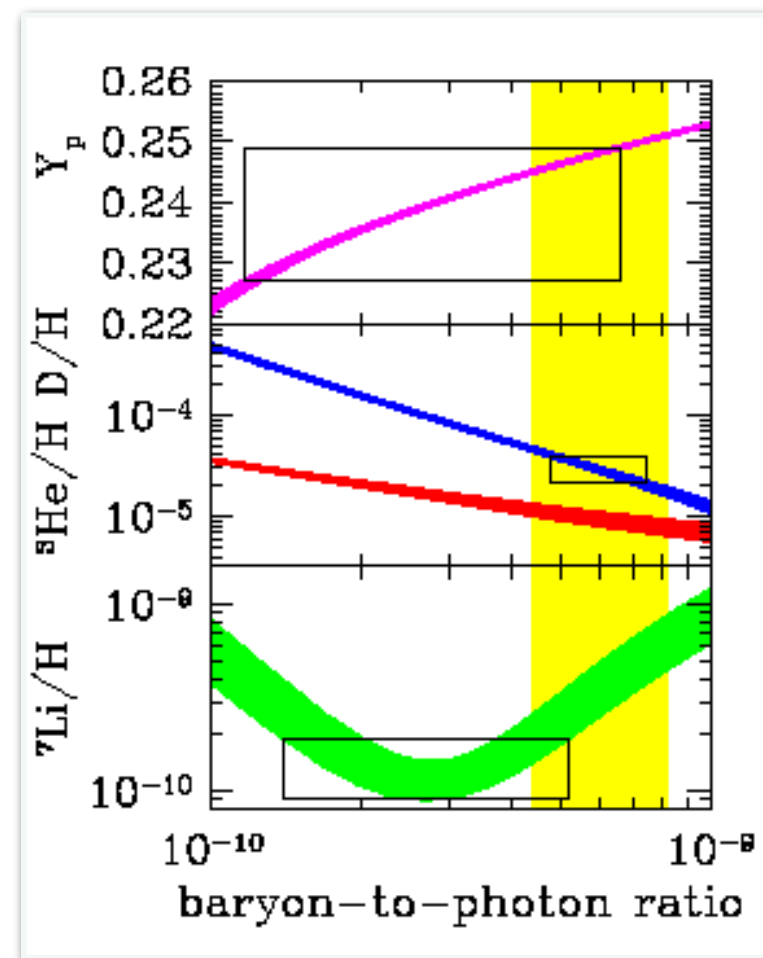
**Department of Physics and Astrophysics, Beijing Normal University**

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# New physics—The Baryon asymmetry

## Matter-antimatter asymmetry

- ✳ No galaxy made by ant-baryon is observed
- ✳ Baryon asymmetry is measured by the Planck.

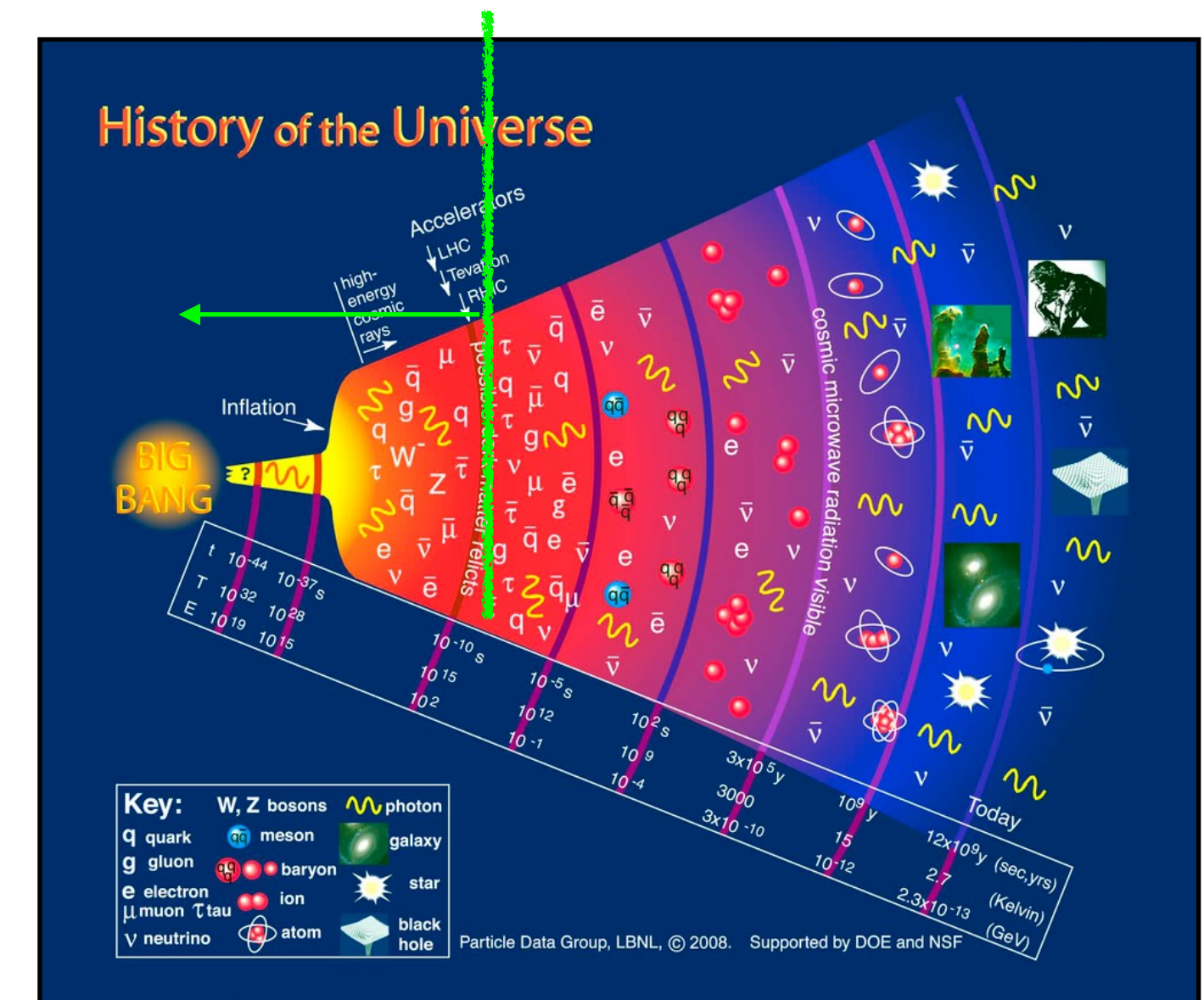


Baryon asymmetry:  $Y_B = \frac{\rho_B}{s} = (8.59 \pm 0.11) \times 10^{-11}$  Planck

## Baryogenesis



- ★ Baryon number violating
- ★ C&CP violation
- ★ Departure from equilibrium



- Leptogenesis
- Electroweak Baryogenesis
- GUT Baryogenesis
- Affleck-Dine Baryogenesis
- Post-sphaleron baryogenesis

# History and development: Leptogenesis

## Type-I seesaw mechanism

Neutrino physics

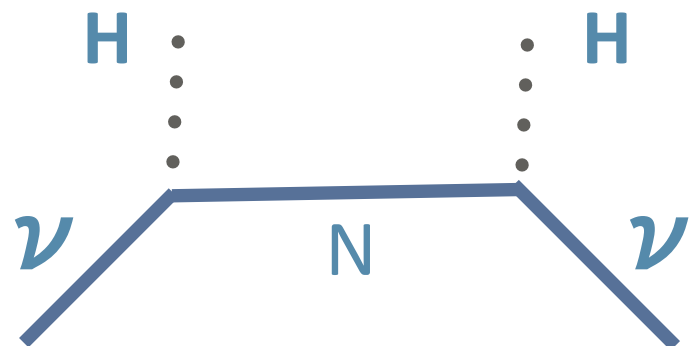
Leptogenesis via seesaw

BAU

Lagrangian:

$$-\mathcal{L} = Y_\nu \bar{\ell}_L \tilde{H} N_R + \frac{1}{2} \overline{N_R^C} M_R N_R + \text{h.c.}$$

Seesaw Mechanism



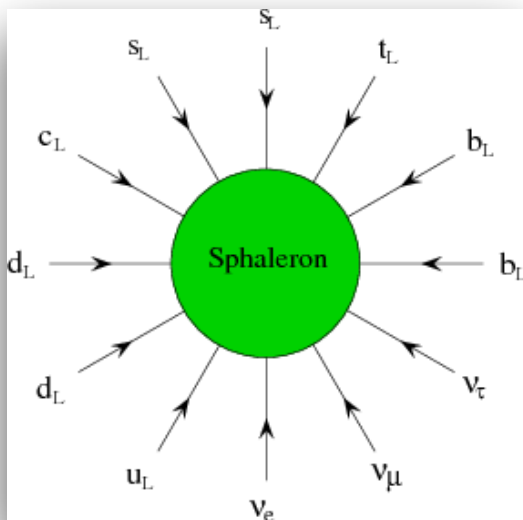
Elegant but losing testability at colliders!

Neutrino mass:

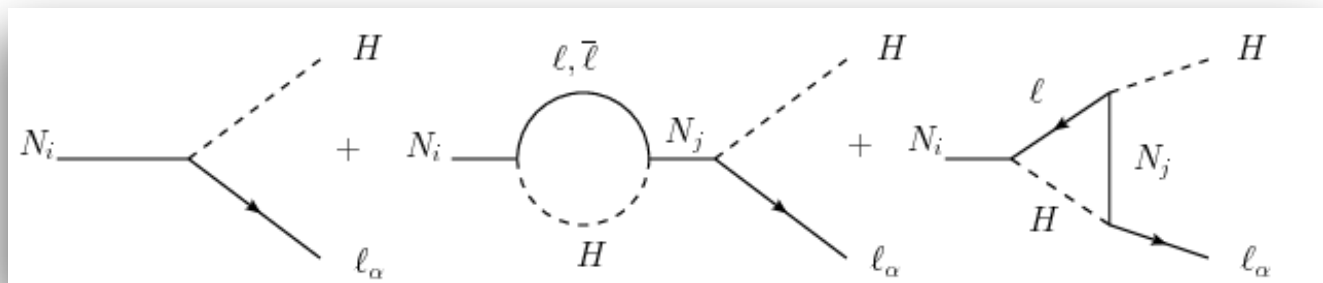
$$M_\nu = -M_D M_R^{-1} M_D^T$$

## Leptogenesis

BV



CPV

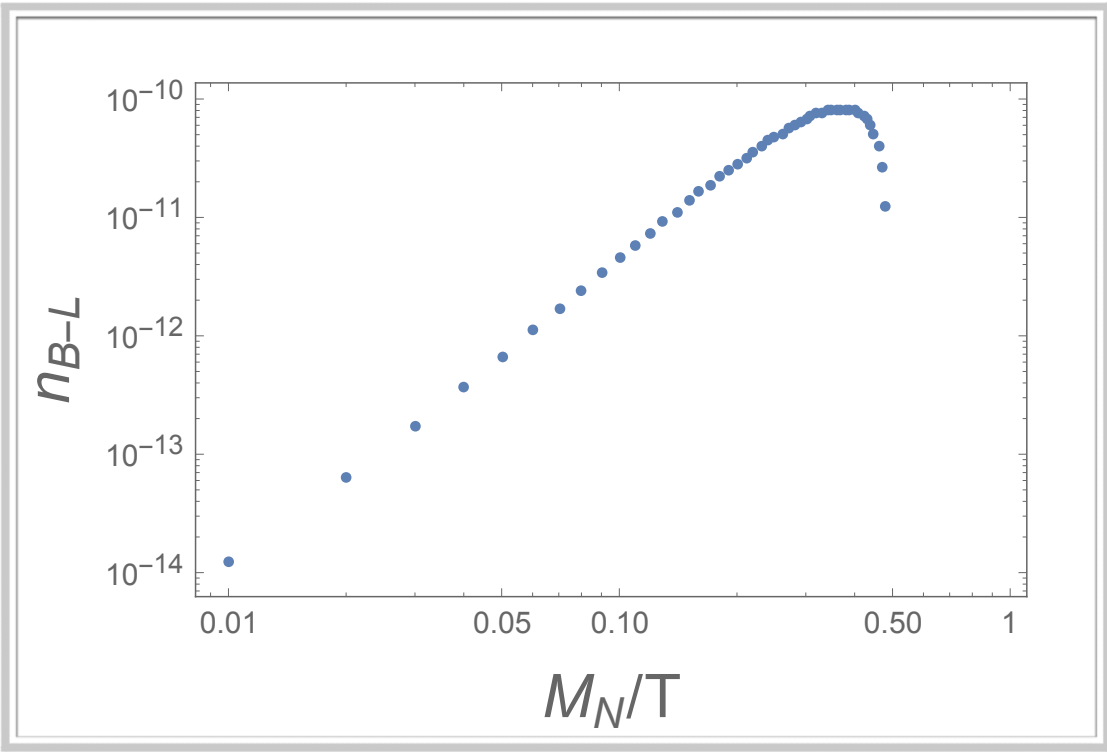


Out of EQ

Heavy neutrinos decouple

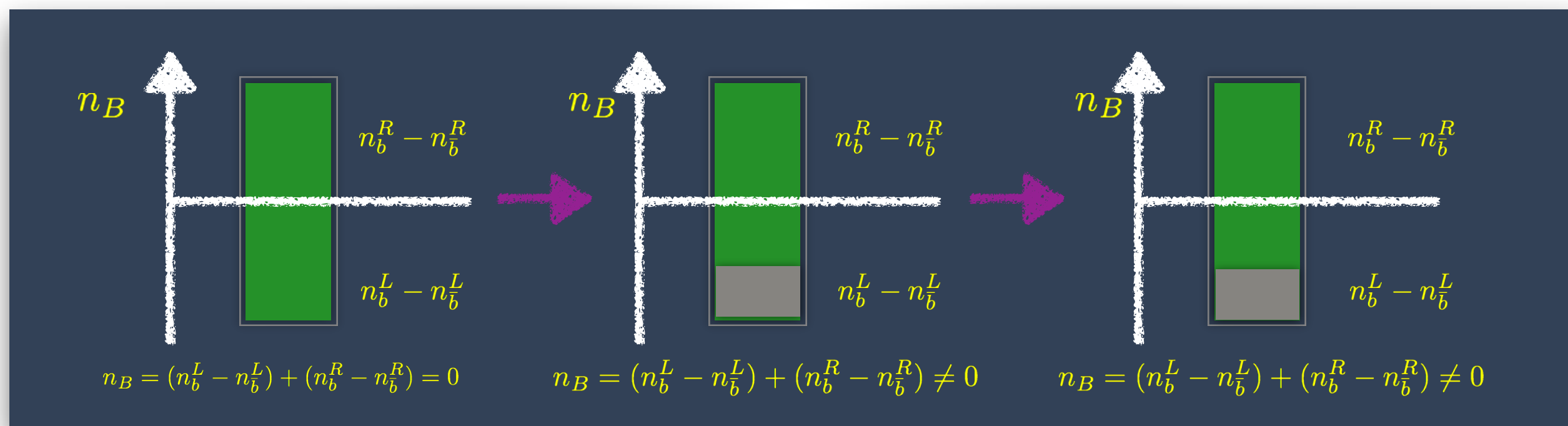
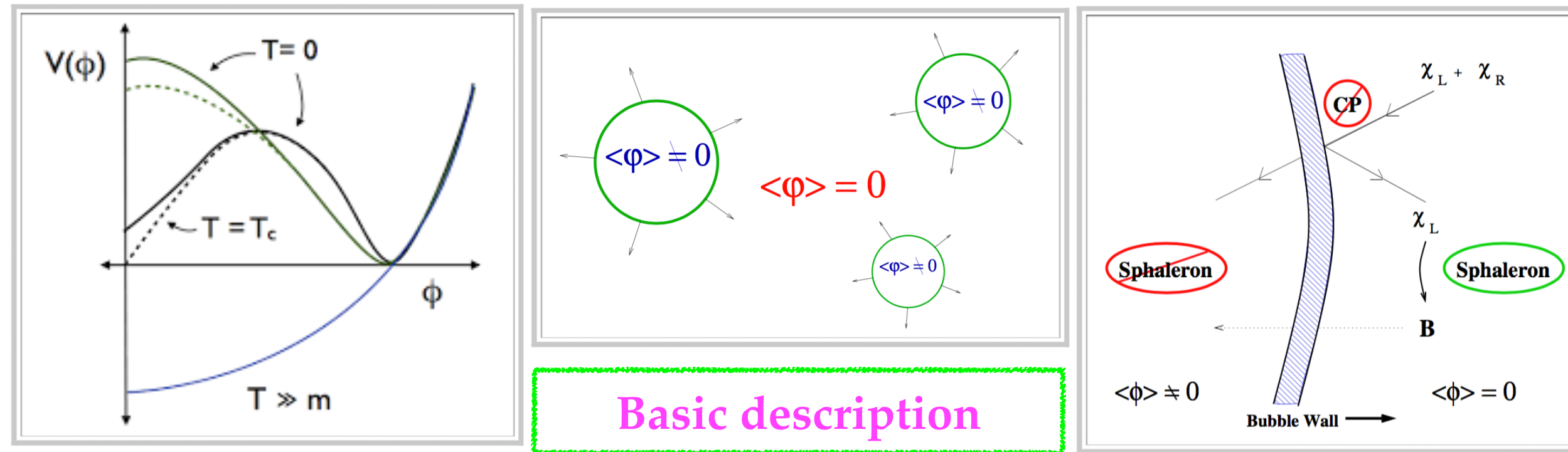
Boltzmann EQ

$$\frac{dN}{dz} = -D(N - N^{\text{eq}})$$
$$\frac{dN_{B-L}}{dz} = -\epsilon D(N - N^{\text{eq}}) + W_{\text{ID}} N_{B-L}$$



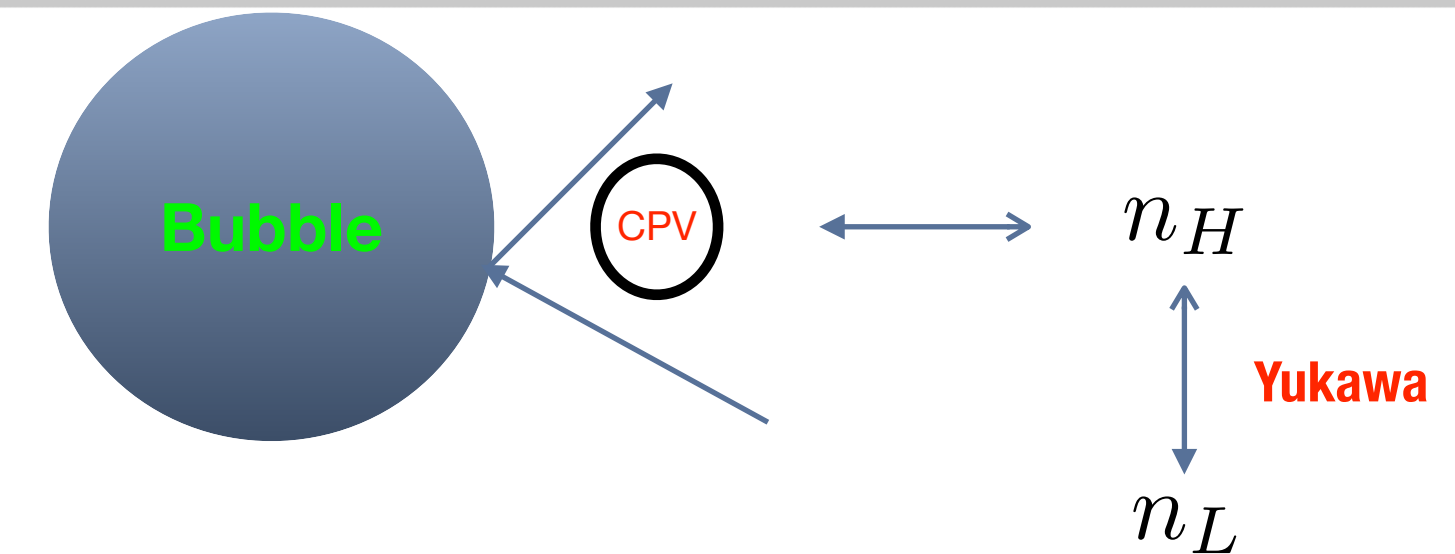
# History and development: EW Baryogenesis

## EWBG



## Transport equations

$$\frac{\partial n}{\partial t} + \nabla \cdot j(x) = - \int d^3z \int_{-\infty}^{x_0} dz^0 \text{Tr}[\Sigma^>(x, z) S^<(z, x) - S^>(x, z) \Sigma^<(z, x) + S^<(x, z) \Sigma^>(z, x) - \Sigma^<(x, z) S^>(z, x)]$$



$$\begin{aligned} \partial_\mu \psi_\mu &= +\Gamma_\psi^+ \left( \frac{\chi}{k_\chi} + \frac{\psi}{k_\psi} \right) + \Gamma_\psi^- \left( \frac{\chi}{k_\chi} - \frac{\psi}{k_\psi} \right) + \left( \sum_i \Gamma_{y_i} \right) \left( \frac{\chi}{k_\chi} - \frac{H}{k_H} - \frac{\psi}{k_\psi} \right) + S_{\text{CP}}^\psi \\ \partial_\mu \chi_\mu &= -\Gamma_\psi^+ \left( \frac{\chi}{k_\chi} + \frac{\psi}{k_\psi} \right) - \Gamma_\psi^- \left( \frac{\chi}{k_\chi} - \frac{\psi}{k_\psi} \right) - \left( \sum_i \Gamma_{y_i} \right) \left( \frac{\chi}{k_\chi} - \frac{H}{k_H} - \frac{\psi}{k_\psi} \right) - S_{\text{CP}}^\psi \\ \partial_\mu H_\mu &= \Gamma_{Y_t} \left( \frac{T}{k_T} - \frac{H}{k_H} - \frac{Q}{k_Q} \right) + \left( \sum_i \Gamma_{y_i} \right) \left( \frac{\chi}{k_\chi} - \frac{H}{k_H} - \frac{\psi}{k_\psi} \right) - \Gamma_h \frac{H}{k_H}, \end{aligned}$$

# History and development: Afleck-Dine

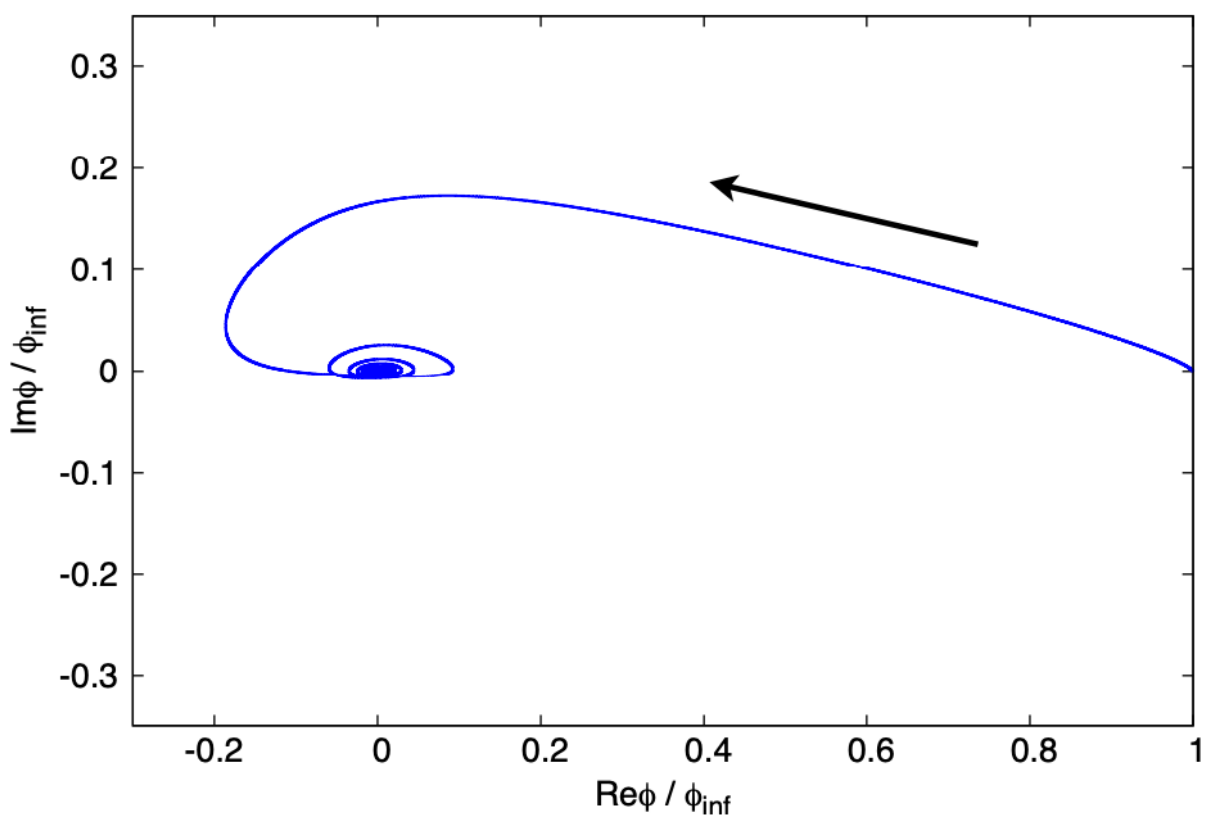
## Affleck-Dine Mechanism

Scalars carrying  
non-zero U(1)  
charges

Flat directions  
(AD fields)

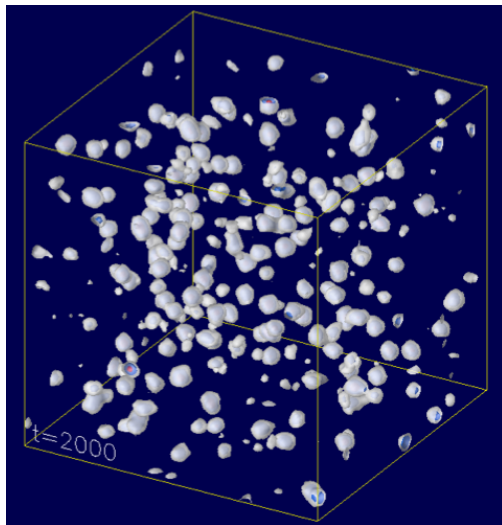
Lifting the potential  
via B/L violation  
operators

$$V = (m^2 - cH^2) |\phi|^2 + \lambda |\phi|^4 + \left( \frac{\phi^n}{M^{n-4}} + \text{h.c.} \right)$$



$$\dot{n}_{B,L} + 3Hn_{B,L} = 2\beta \text{Im} \left[ \frac{\partial V}{\partial \phi} \phi \right]$$

## Q-ball formation (Non-topological soliton in scalar field theory)



Oscillation of AD field



Q-ball formation



Long lived Q-ball



DM candidate

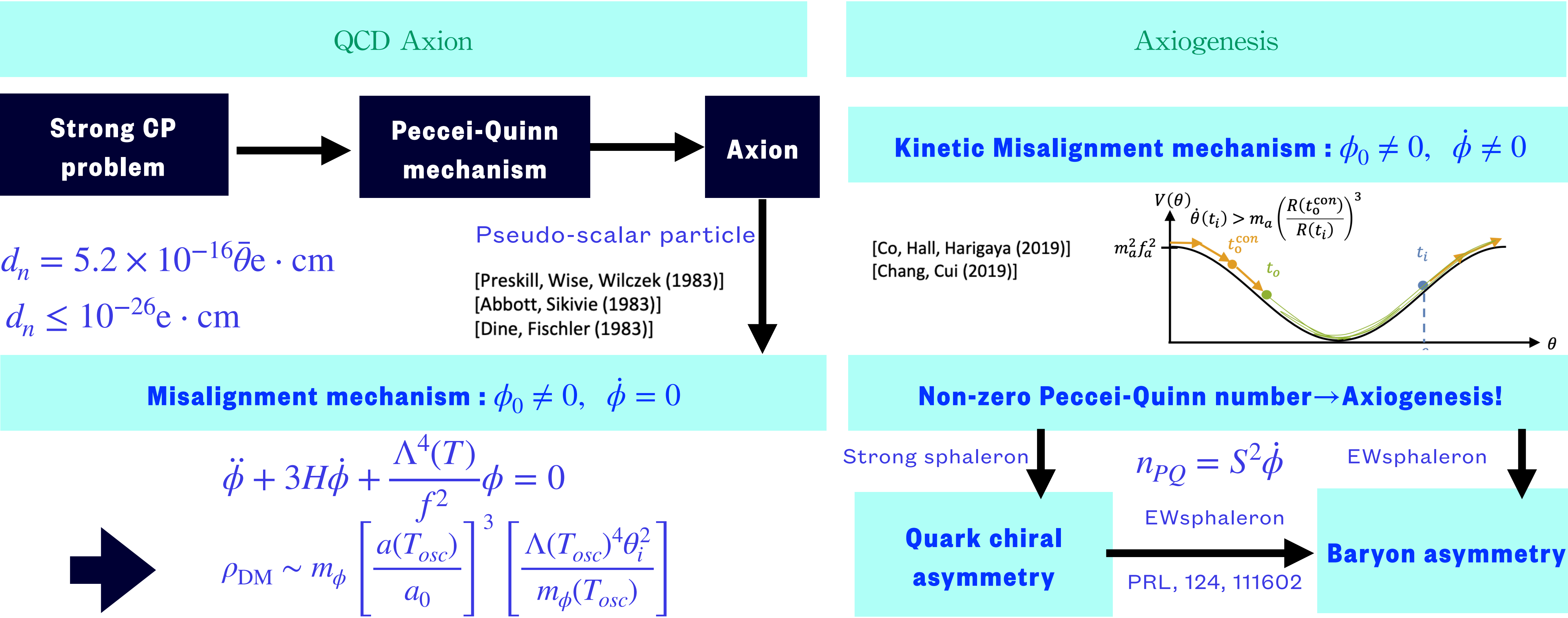


Evaporation



BAU when sphaleron  
erase is irrelevant

# History and development: Axionogenesis



# Main point of this talk about Leptogenesis

**Traditional  
Leptogenesis  
mechanism:**

- **There must be primordial B-L violation**
- **There must exist right-handed neutrinos**

Eogenesis via the High-scale Electroweak Symmetry Restoration

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2412.03902

<sup>1</sup>*Key Laboratory of Multi-scale Spin Physics, Ministry of Education,  
Beijing Normal University, Beijing 100875, China*

<sup>2</sup>*Center of Advanced Quantum Studies, School of Physics and Astronomy,  
Beijing Normal University, Beijing, 100875, China*

In this paper, we propose a novel electron-assisted Baryogenesis scenario that does not require explicit B-L violation, which is essential for the traditional Leptogenesis mechanism. This scenario

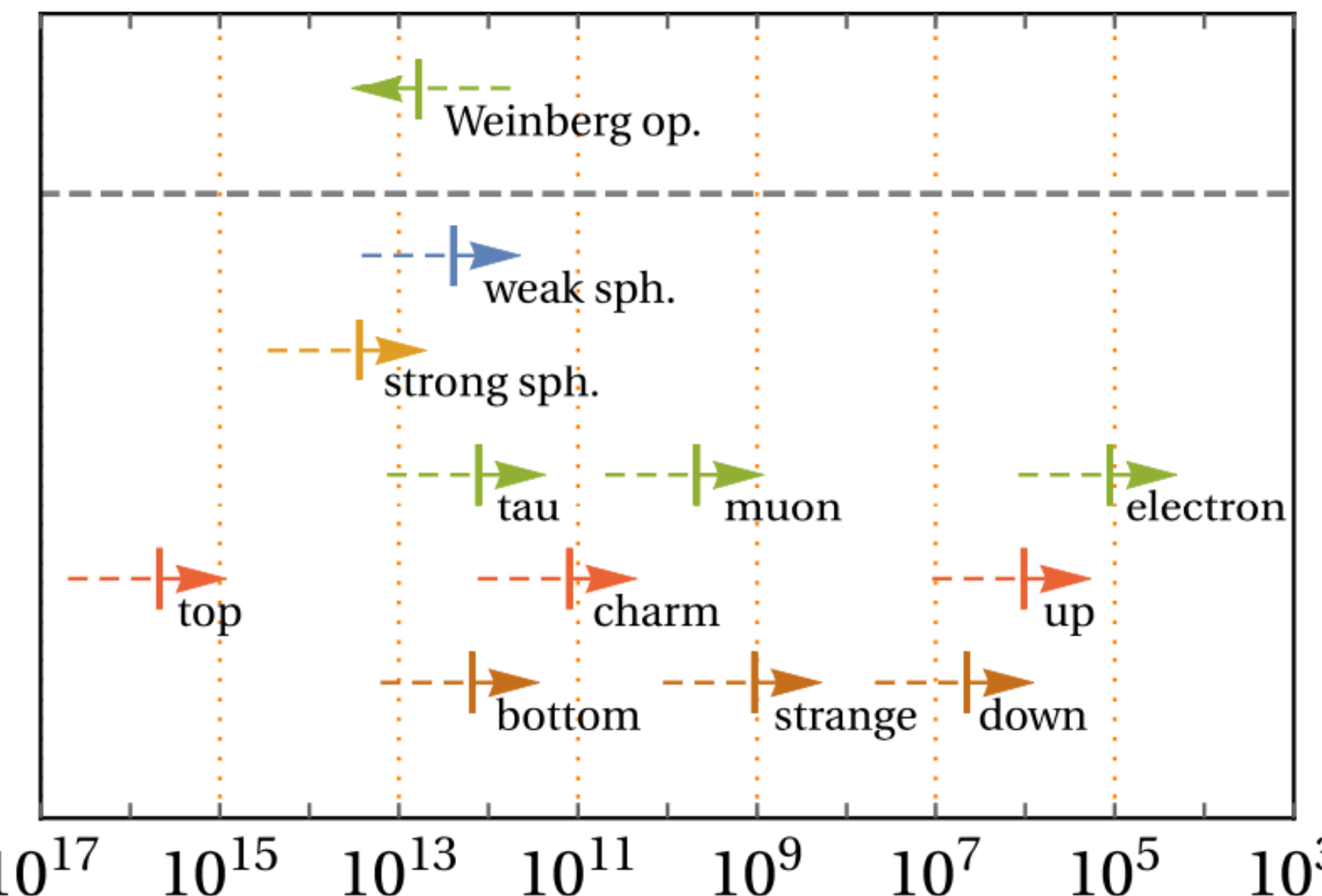
**Not necessary!**

**Sphaleron quenches before electron  
Yukawa interaction entering  
thermal equilibrium in the early  
universe!**

# Parameters for various interactions

- Key Point: Sphaleron may quench before the electron Yukawa interaction entering the equilibrium**

| Interaction         | Weinberg                                | WS                                   | SS                                   | $Y_e$                | $Y_\mu$                | $Y_\tau$                |
|---------------------|-----------------------------------------|--------------------------------------|--------------------------------------|----------------------|------------------------|-------------------------|
| $\Gamma_\alpha/T^4$ | $\kappa_W \frac{m_\nu^2 T^2}{v_{EW}^4}$ | $\frac{1}{2} \kappa_{WS} \alpha_2^5$ | $\frac{1}{2} \kappa_{SS} \alpha_3^5$ | $\kappa_{Y_e} y_e^2$ | $\kappa_{Y_e} y_\mu^2$ | $\kappa_{Y_e} y_\tau^2$ |
| $T_\alpha$ [GeV]    | $6.0 \times 10^{12}$                    | $2.5 \times 10^{12}$                 | $2.8 \times 10^{13}$                 | $1.1 \times 10^5$    | $4.7 \times 10^9$      | $1.3 \times 10^{12}$    |
| Interaction         | $Y_u$                                   | $Y_c$                                | $Y_t$                                | $Y_d$                | $Y_s$                  | $Y_b$                   |
| $\Gamma_\alpha/T^4$ | $\kappa_{Y_u} y_u^2$                    | $\kappa_{Y_u} y_c^2$                 | $\kappa_{Y_t} y_t^2$                 | $\kappa_{Y_d} y_d^2$ | $\kappa_{Y_d} y_s^2$   | $\kappa_{Y_d} y_b^2$    |
| $T_\alpha$ [GeV]    | $1.0 \times 10^6$                       | $1.2 \times 10^{11}$                 | $4.7 \times 10^{15}$                 | $4.5 \times 10^6$    | $1.1 \times 10^9$      | $1.5 \times 10^{12}$    |



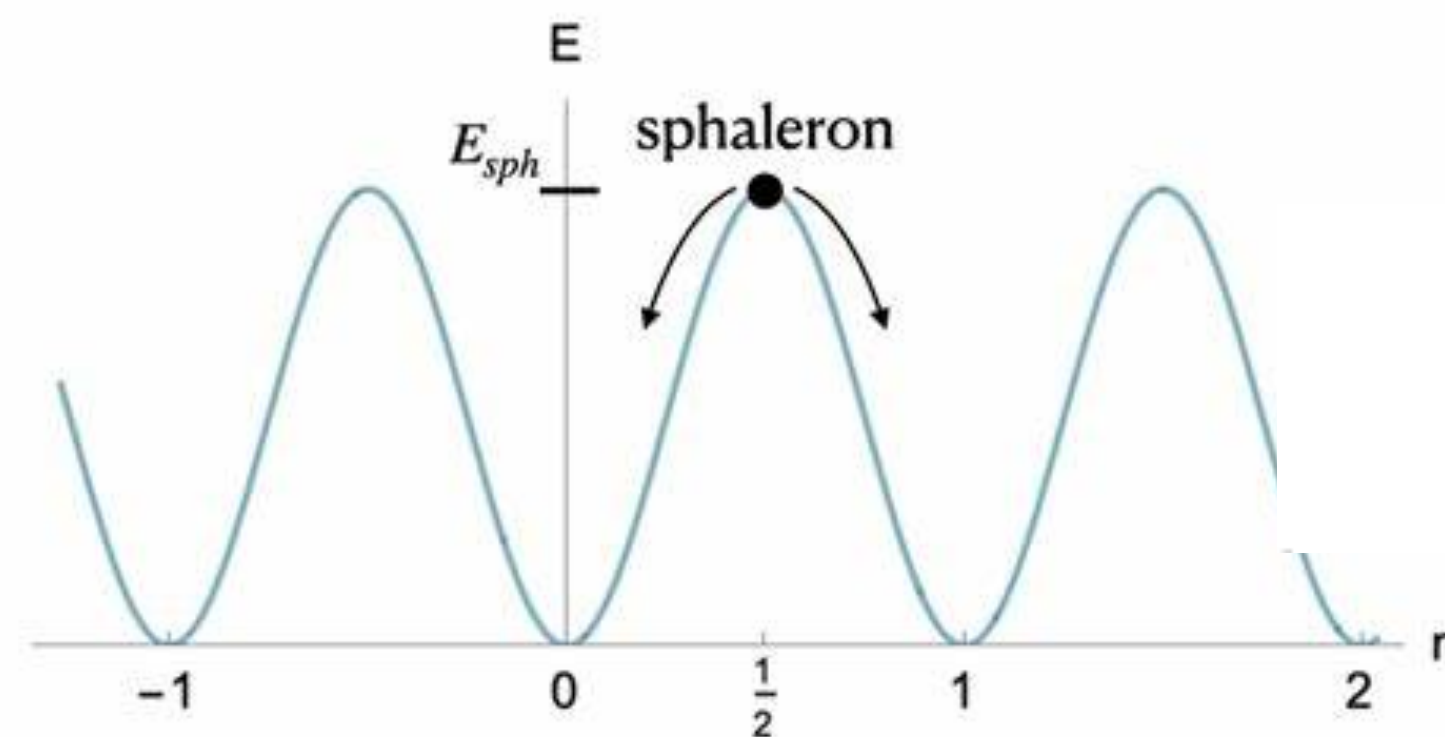
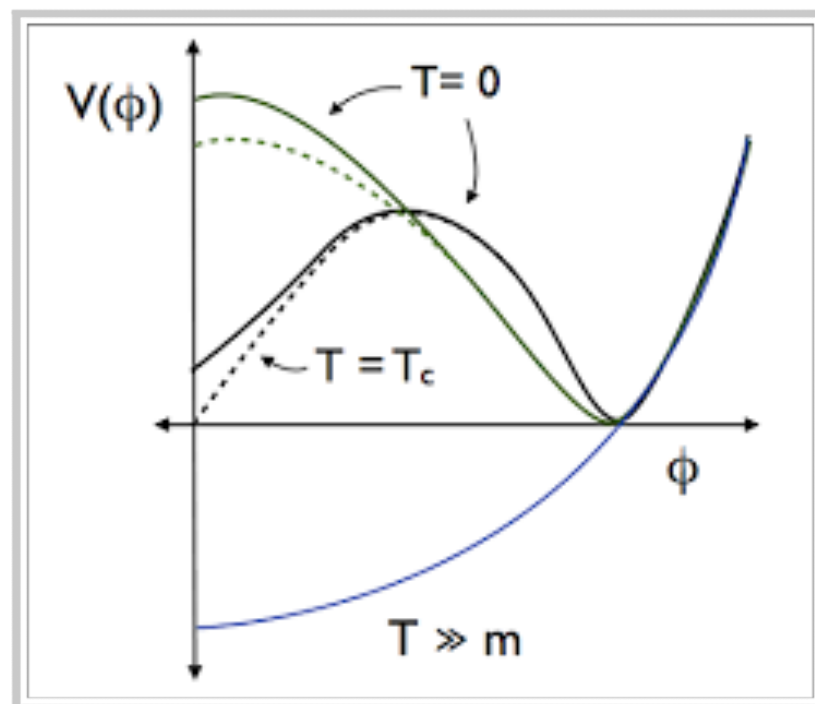
Valerie Domck, Yohei Ema, Kyohei Mukai and Masaki Yamada, JHEP 08(2020)096

# Sphaleron

- **Typical temperature:**

Sphaleron quench temperature:  $T = 130 \text{ GeV}$

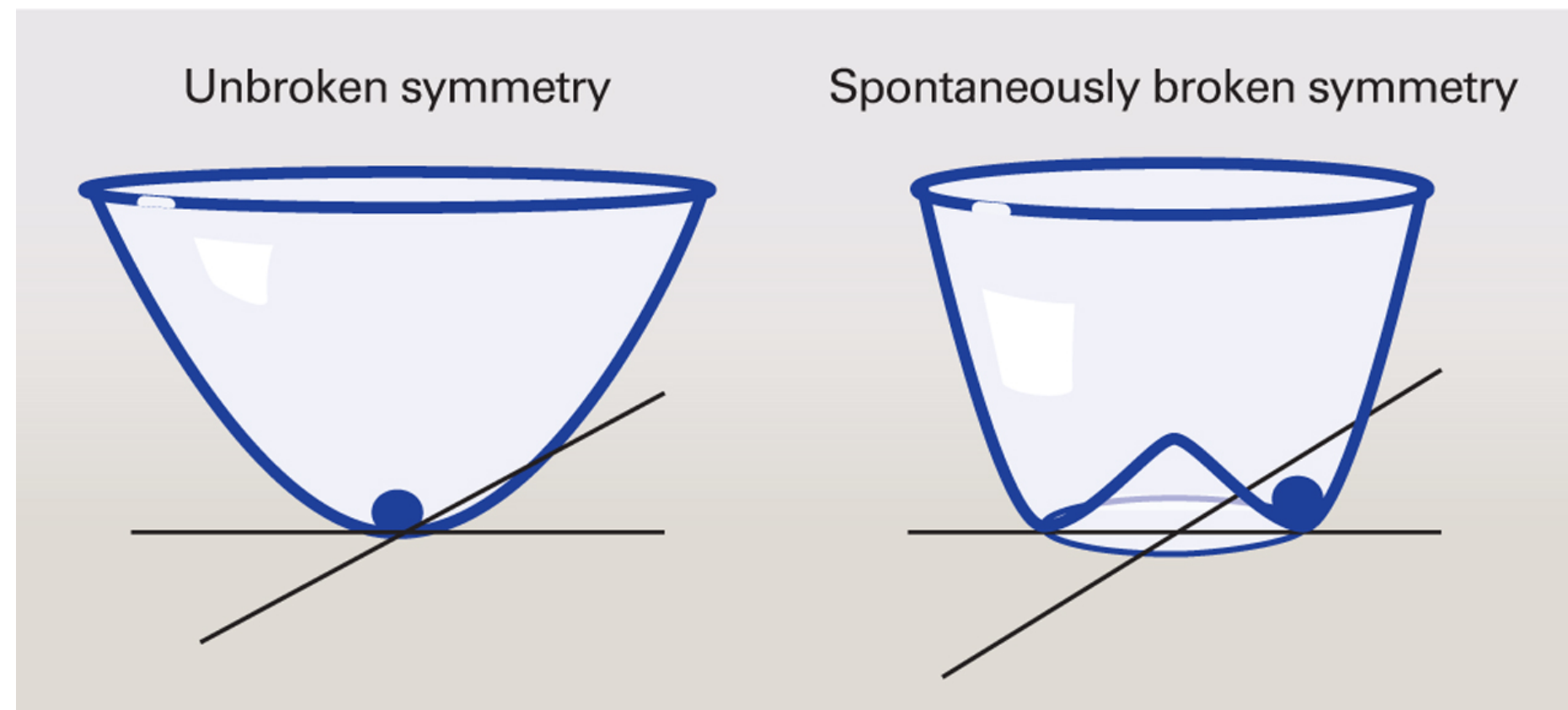
Electroweak symmetry restored temperature:  $T = 160 \text{ GeV}$



$$\Gamma_{\text{sph}}^{\text{brok}}(T) = \kappa_{\text{brok}} \alpha_W^4 T^4 \exp\left(-\frac{E_{\text{sph}}}{T}\right) \quad (4)$$

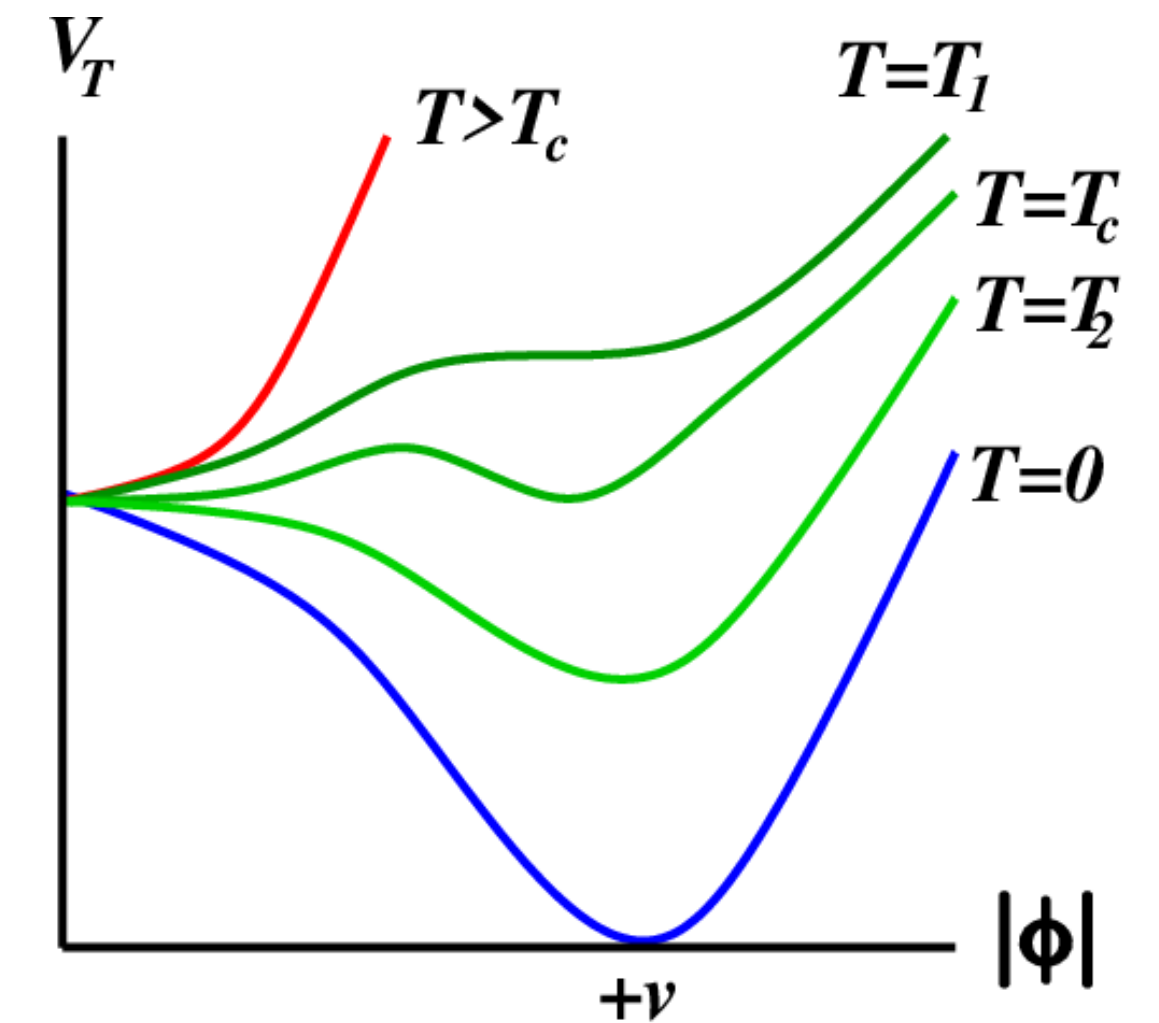
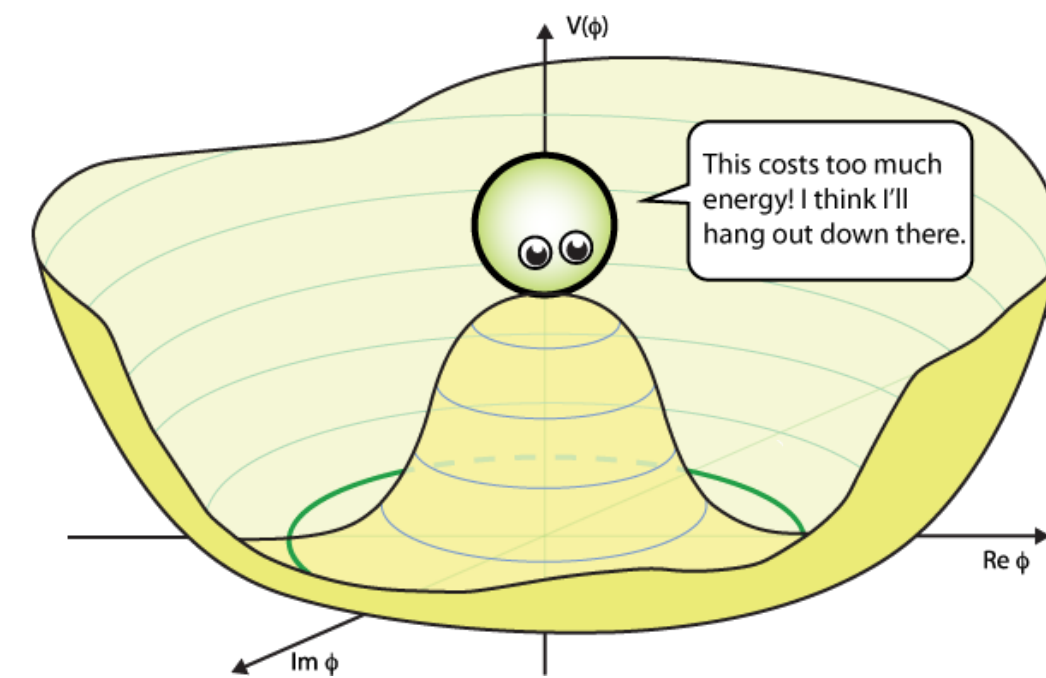
# Symmetry at high temperature

- EW symmetry non-restoration!**



$$V = -\frac{\mu^2}{2}h^2 + \frac{\lambda}{4}h^4$$

Tree-level



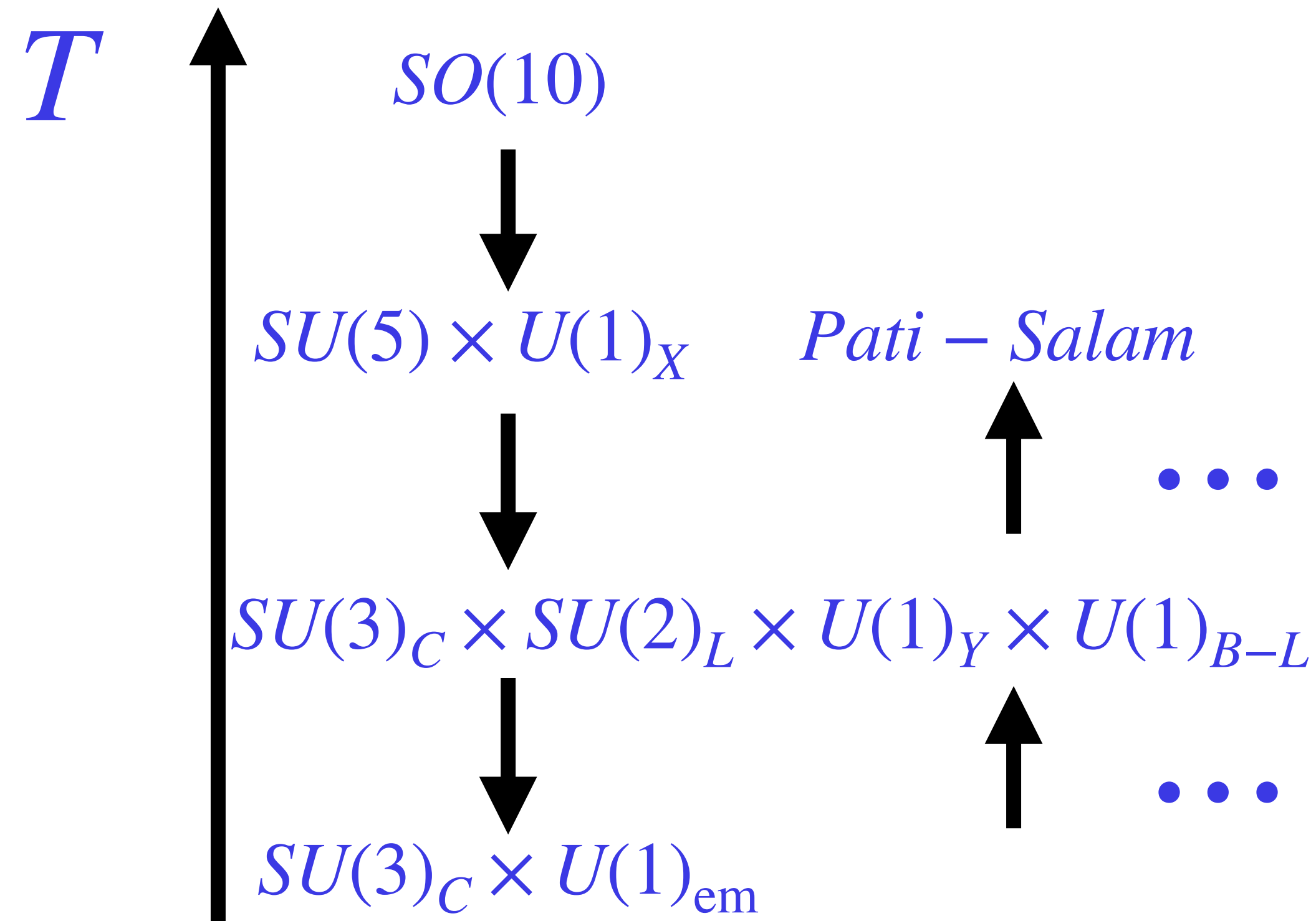
$$\Delta V_1(\phi, T) = \sum_F \frac{g_F T^4}{2\pi^2} \left[ \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} (\beta m_F)^2 K_2(\beta m_F n) \right] - \sum_B \frac{g_B T^4}{2\pi^2} \left[ \sum_{n=1}^{\infty} \frac{1}{n^2} (\beta m_B)^2 K_2(\beta m_B n) \right].$$

Thermal corrections

# Symmetry restoration at high T?

Conventional symmetry breaking sequence

Question: Must all symmetries be restored at high T?



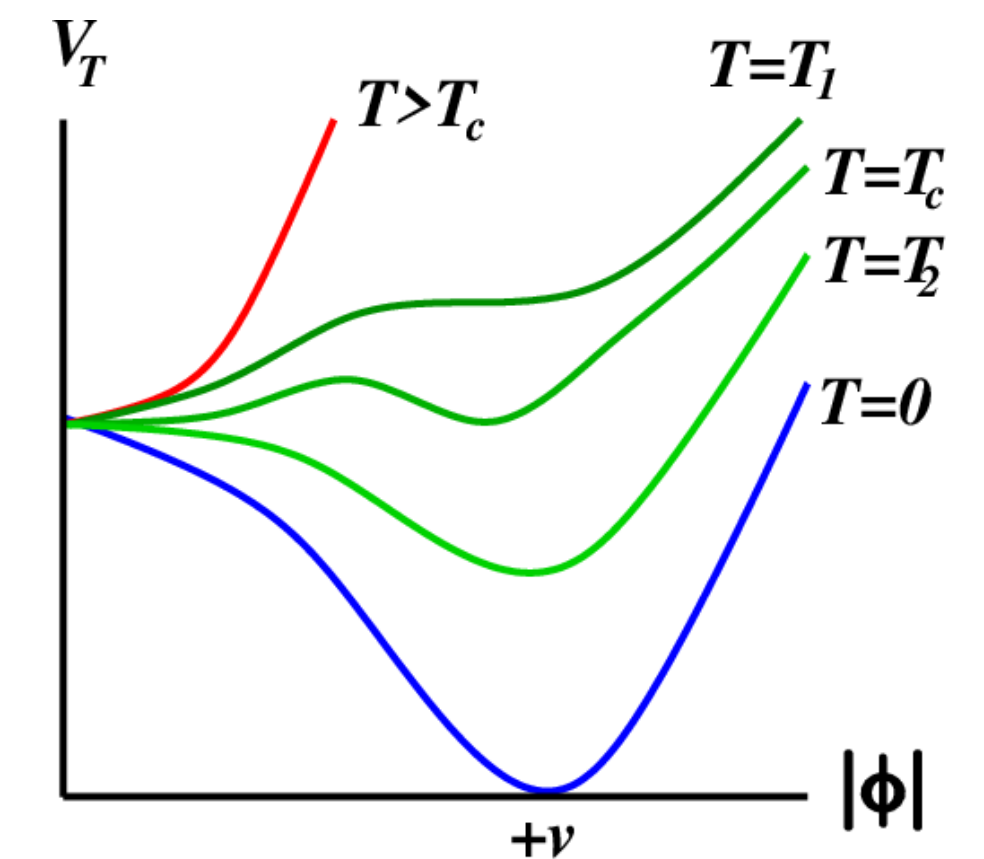
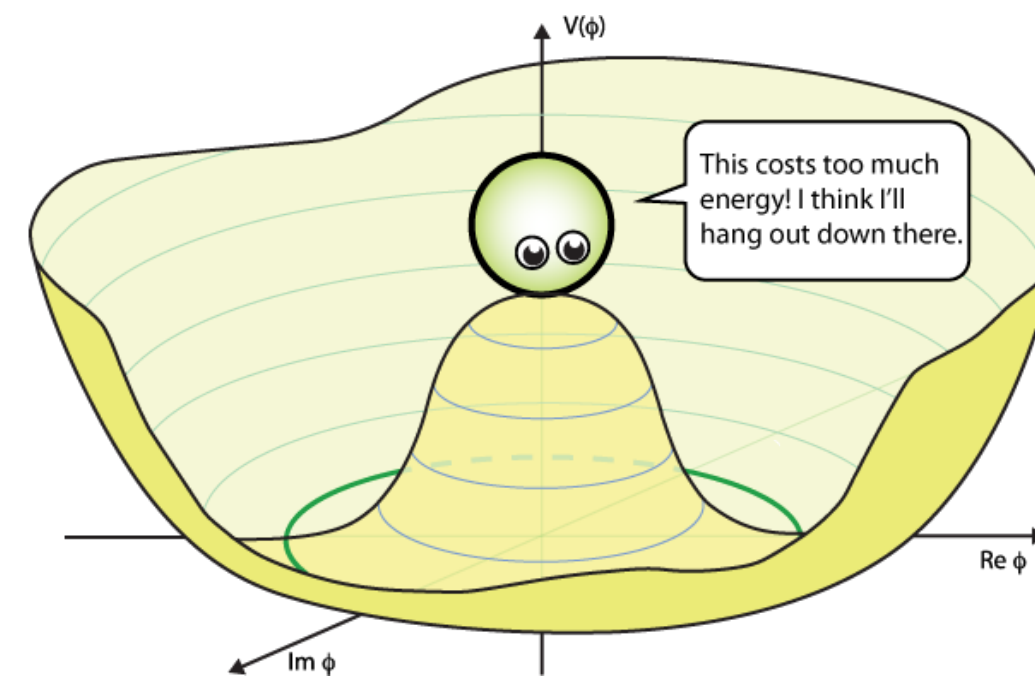
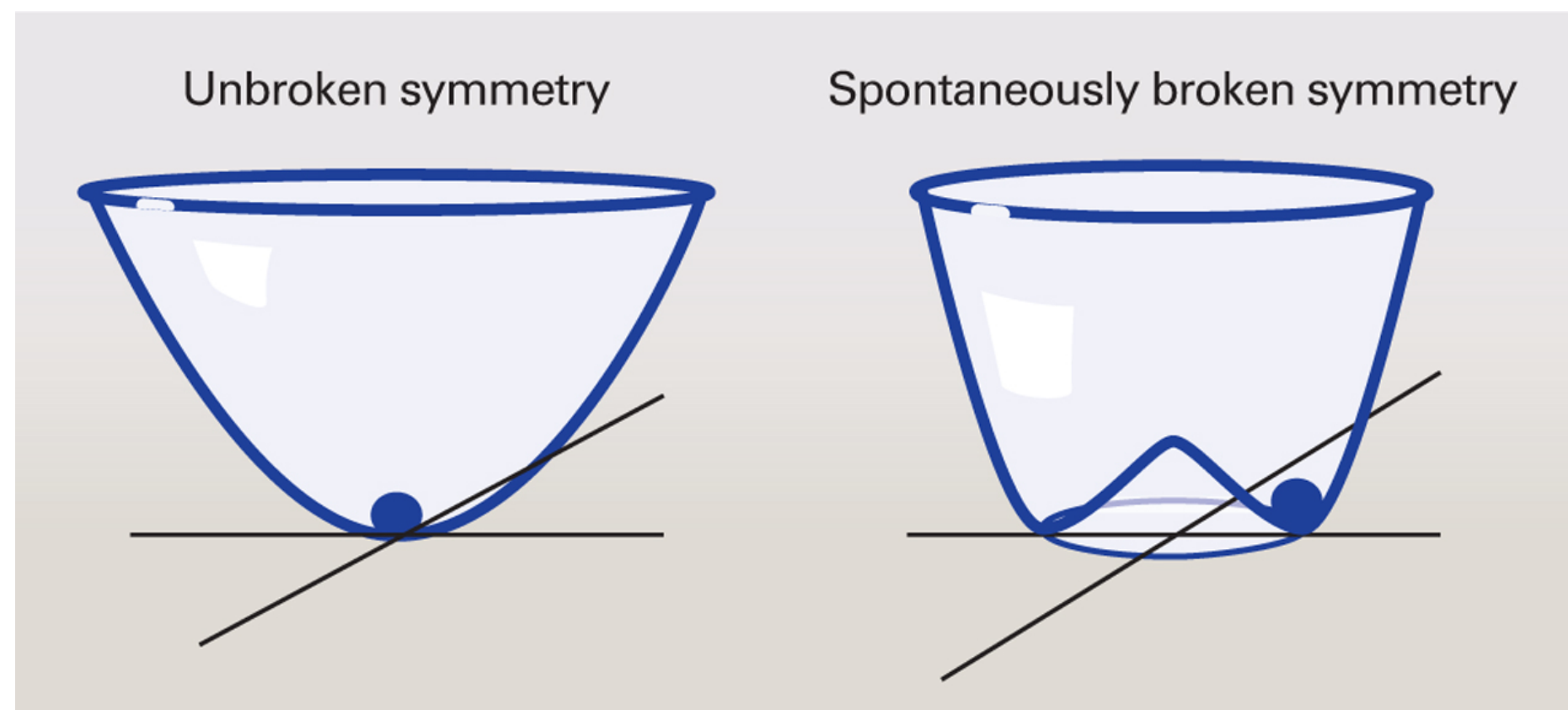
Maybe not!

- \*  $O(n) \times O(n)$ : Weinberg (1974)
- \*  $SU(5)$ : Dvali, Mohapatra, Senjanovic ('79, 80'S, 90'S)
- \* Cline, Moore, Servant et al (1999)
- \* EM: Langacker & Pi (1980)
- \*  $SU(3)$ : Patel, Ramsey-Musolf, Wise (2013)

# Symmetry non-restoration

- EW symmetry non-restoration?

Maybe!



$$V = -\frac{\mu^2}{2}h^2 + \frac{\lambda}{4}h^4$$

Tree-level

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} + N_s \frac{\lambda_{hs}}{12} \right) \quad (3)$$

BSM corrections can be negative!

# Color symmetry breaking at high T?

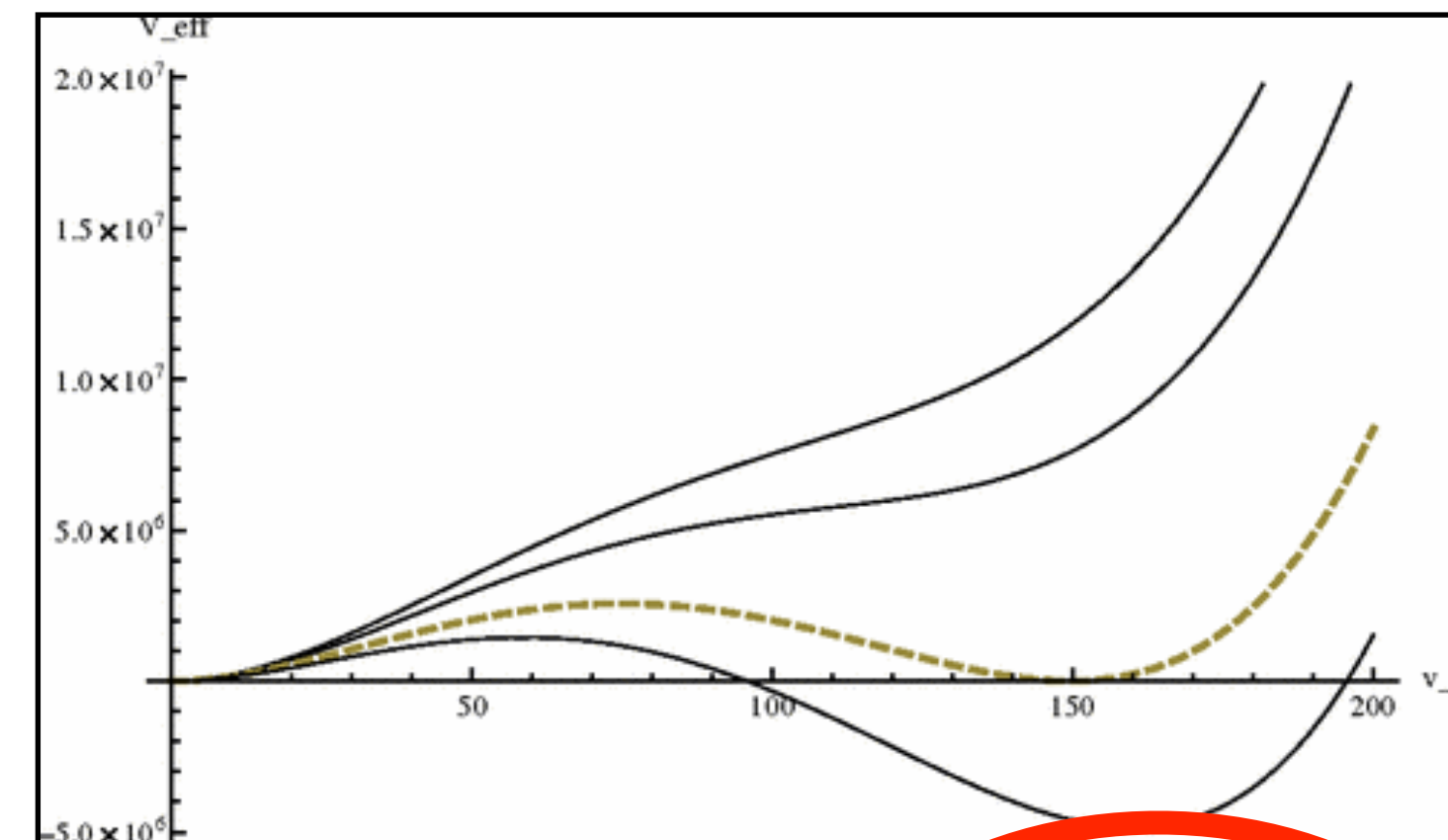
Key: symmetry non-storiation

• P. Meade and H. Ramani, *Phys. Rev. Lett.* 122, 041802

$$\delta V_h(T) \sim \frac{1}{2} c_h T^2 h^2 = \frac{1}{2} \left[ \frac{3g^2 + g'^2}{16} + \frac{\lambda}{2} + \frac{y_t^2}{4} + \zeta \right] T^2 h^2 \rightarrow$$

Higgs gets large VEV at high T

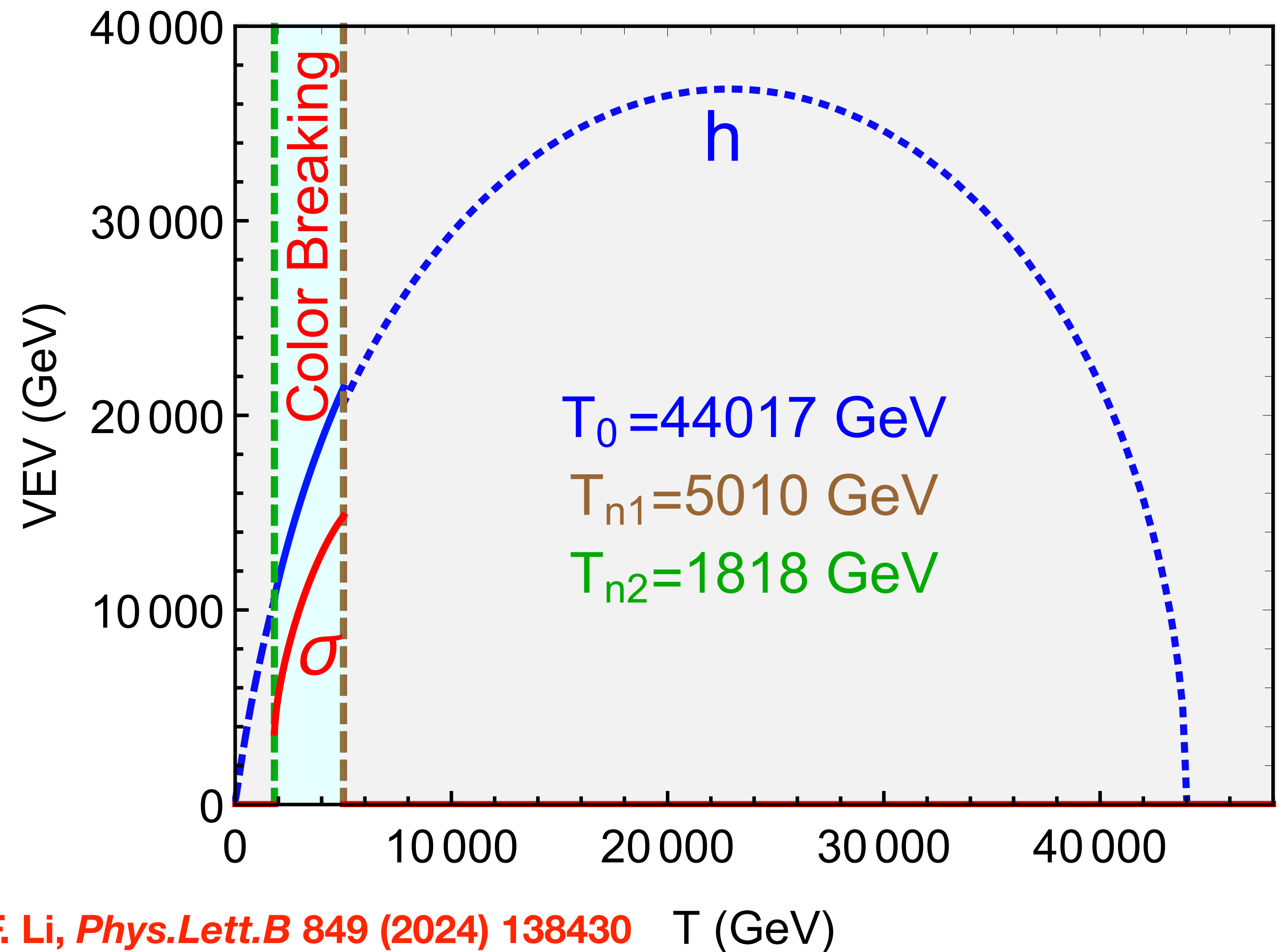
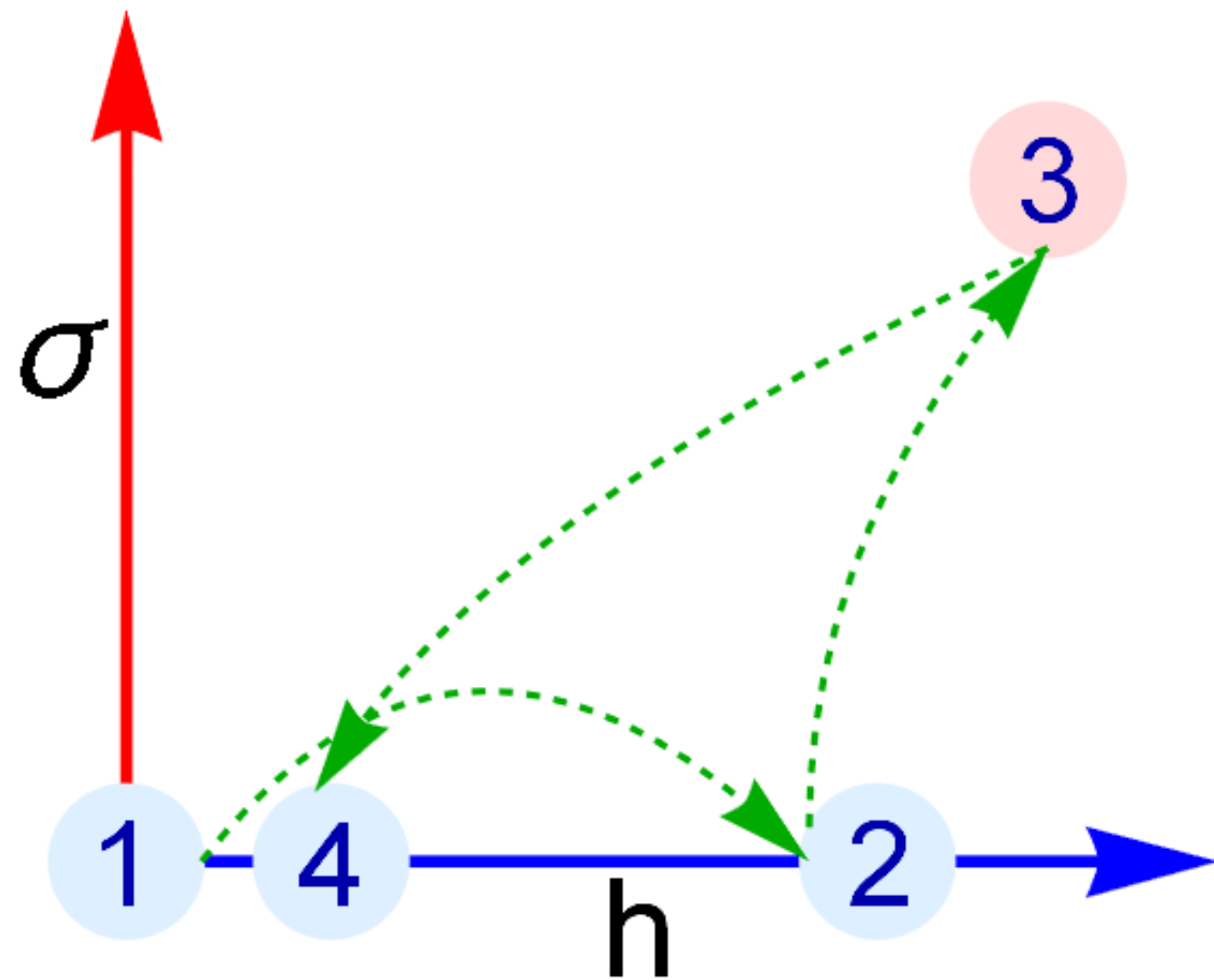
$$V_{\text{eff}} = \mu_c^2 \Delta^\dagger \Delta - \lambda_{hc} H^\dagger H \Delta^\dagger \Delta + \lambda_c (\Delta^\dagger \Delta)^2 + V_{\text{SNR}}$$



$$V(h, \sigma, T) \supset -\frac{\mu_{hT}^2}{2} h^2 + \frac{\mu_{cT}^2}{2} \sigma^2 - \frac{9 + 2\sqrt{3}}{72\pi} T g_s^3 \sigma^3 + \frac{\lambda_h}{4} h^4 + \frac{\lambda_c}{4} \sigma^4 - \frac{\lambda_{hc}}{4} h^2 \sigma^2,$$

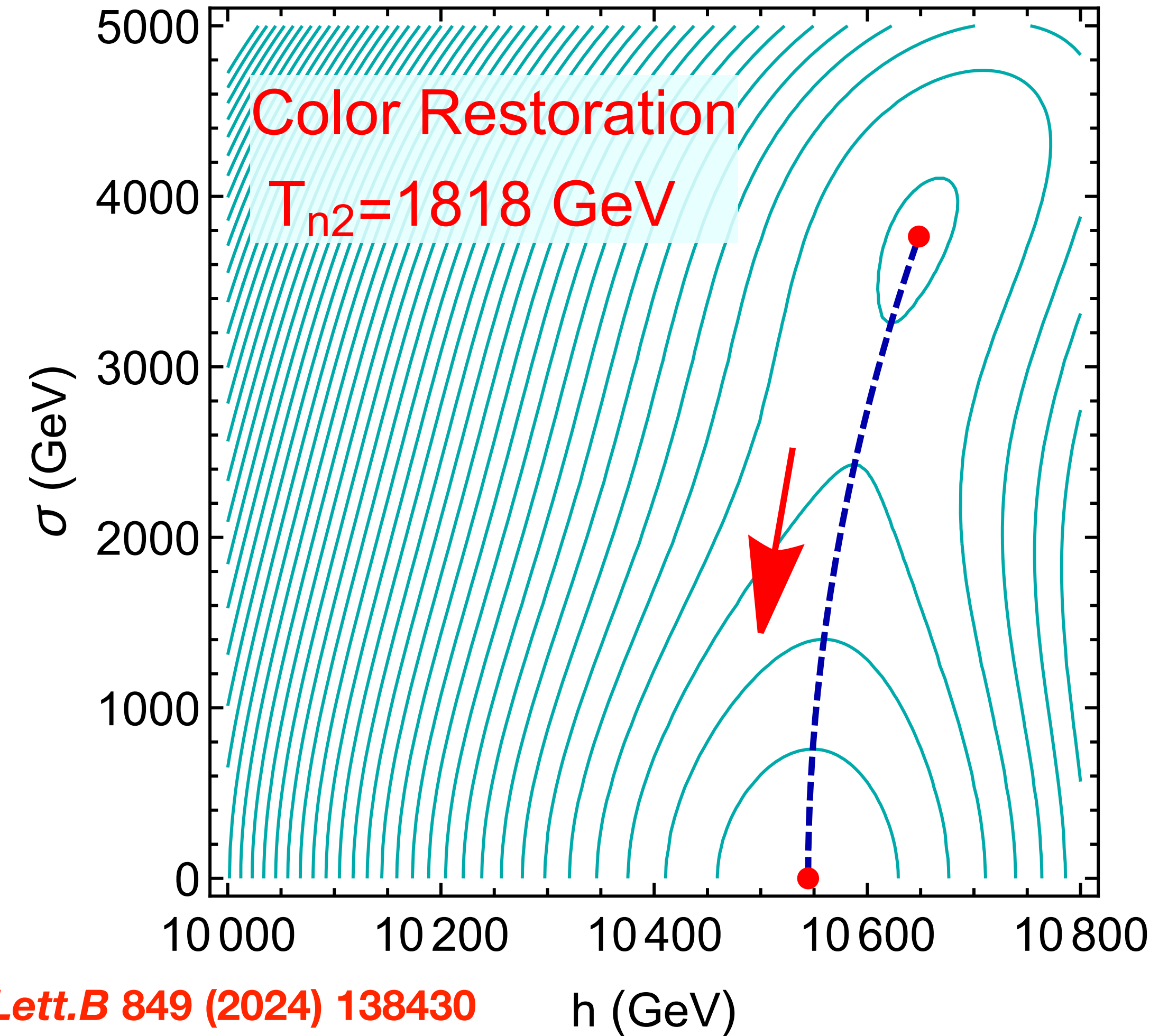
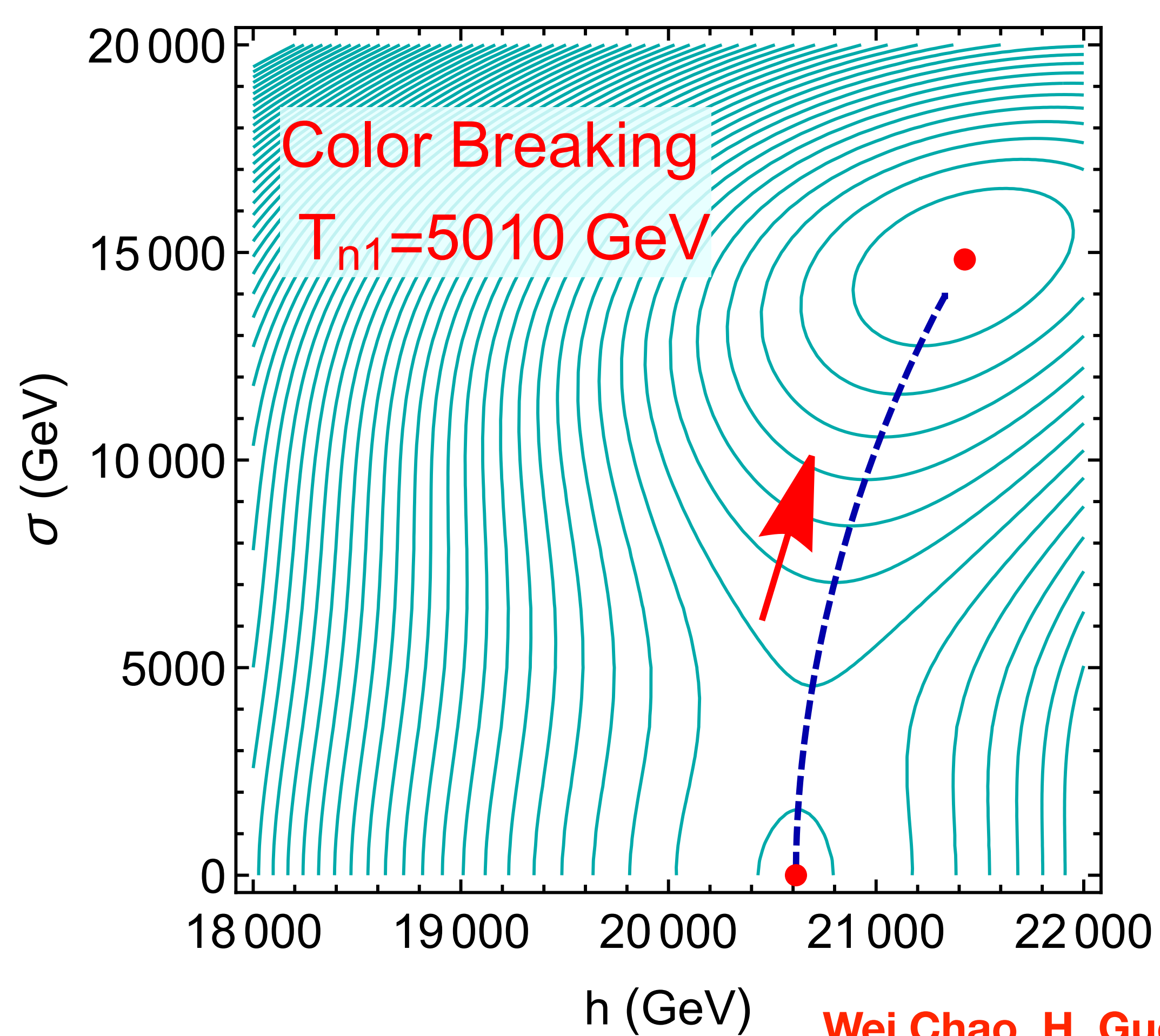
Wei Chao, H. Guo, X. F. Li, *Phys. Lett. B* 849 (2024) 138430

# Color symmetry breaking at high T?



Wei Chao, H. Guo, X. F. Li, *Phys.Lett.B* 849 (2024) 138430

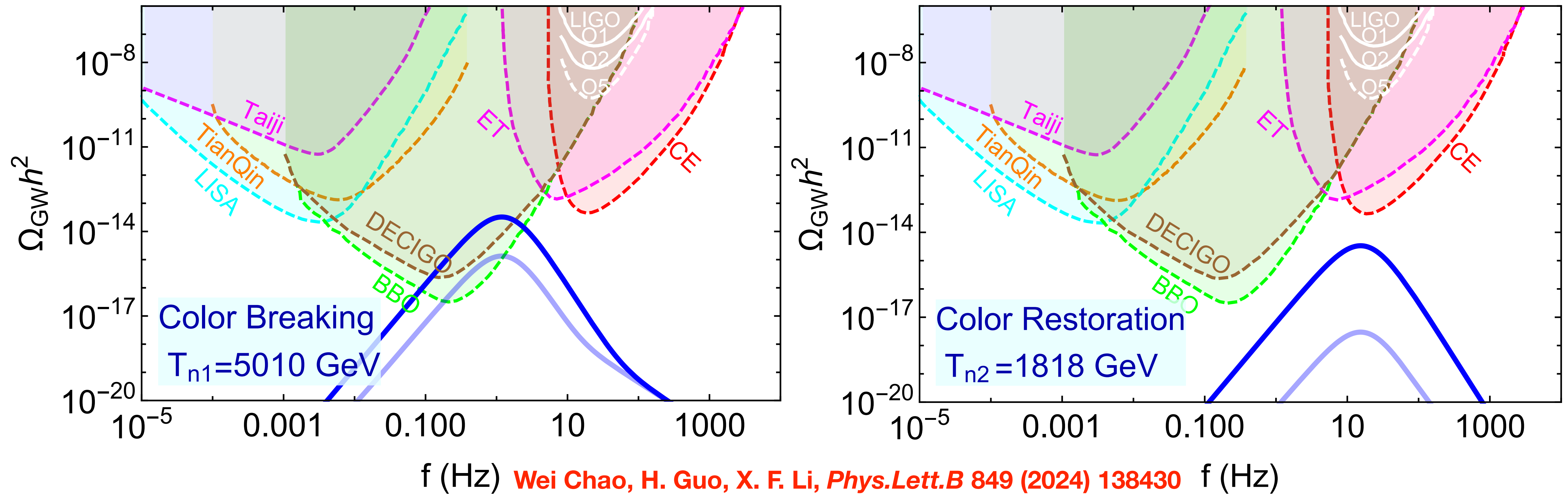
# Phase diagrams



Wei Chao, H. Guo, X. F. Li, *Phys.Lett.B* 849 (2024) 138430

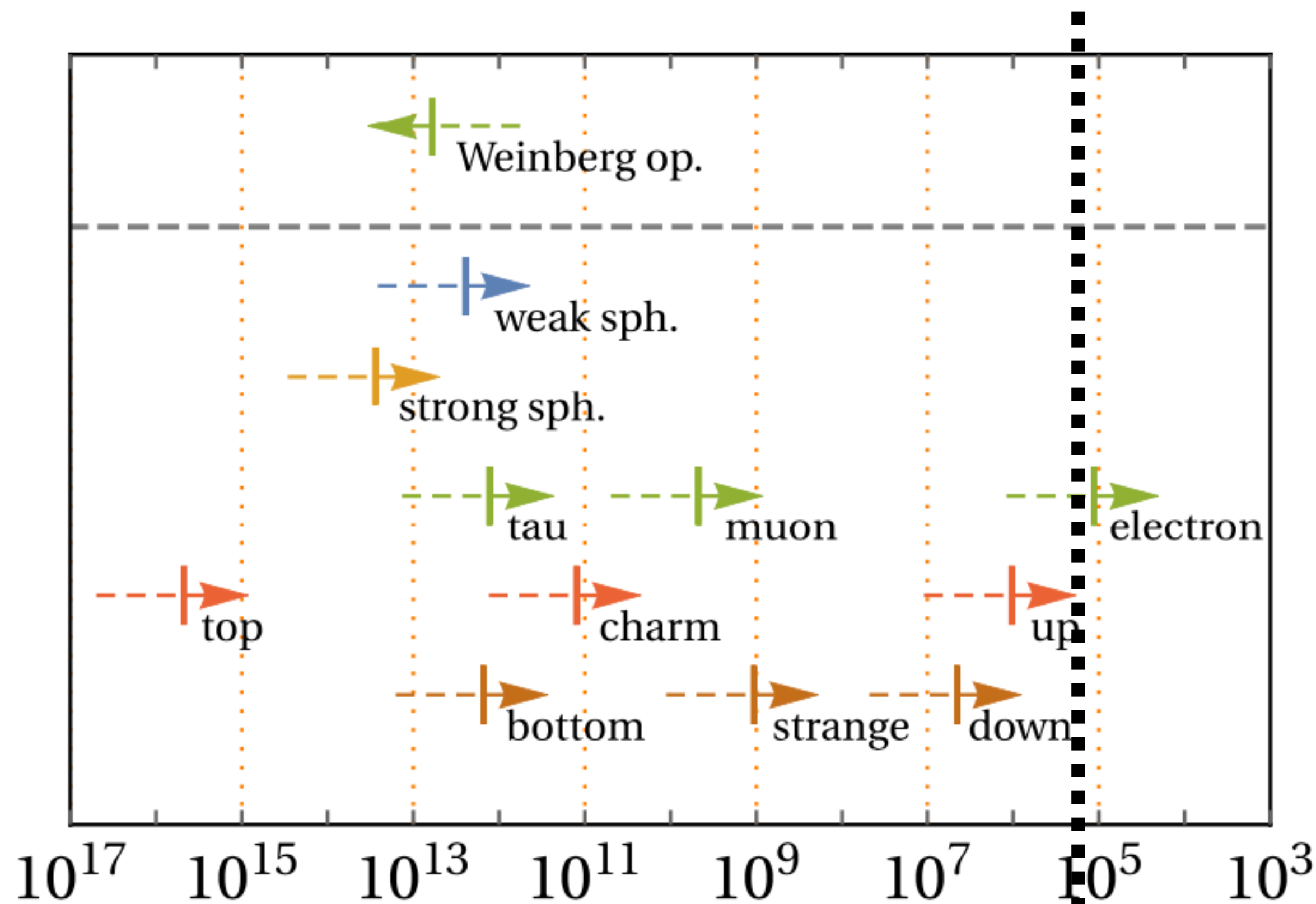
# Gravitational waves

Two peak frequencies!



# Symmetry non-restoration

- **EW symmetry non-restoration? The answer is Yes!**



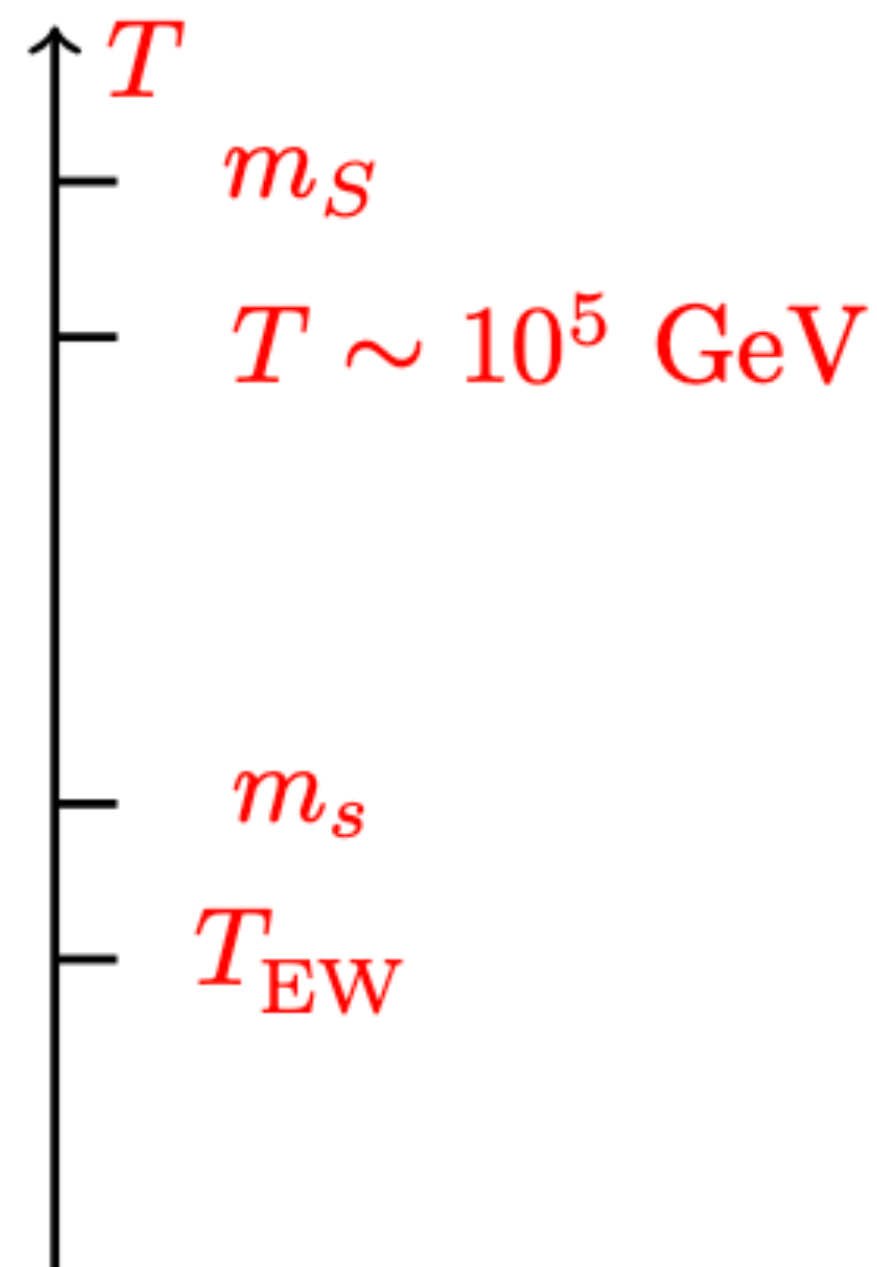
Sphaleron quenches at  $10^5$  GeV.

\* Left-handed Electron asymmetry can be transported to the BAU via sphaleron.

\* No B-L violation is needed!

# Symmetry non-restoration

- **Basic set up**



$$\Gamma_{\text{sph}}^{\text{brok}}(T) = \kappa_{\text{brok}} \alpha_W^4 T^4 \exp\left(-\frac{E_{\text{sph}}}{T}\right) \quad E_{\text{sph}} > T \log\left(\frac{\kappa_{\text{brok}} \alpha_W^4 M_P}{3T}\right)$$

$$E_{\text{sph}} = \frac{4\pi v}{g} \int_0^\infty d\xi \left[ 4(f')^2 + \frac{8}{\xi^2} f^2(1-f)^2 + \frac{\xi^2}{2} (h')^2 + h^2(1-f)^2 + \frac{\xi^2}{16} \sigma^2(h^2-1)^2 \right]$$

$$\rightarrow N_s \lambda_{hs} < -4.82$$

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# Leptogenesis

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- **Scenario-1 : Axion inflation Leptogenesis**
- **Scenario-2 : Eogenesis**

# Axion-inflation baryogenesis

- **U(1)<sub>R</sub>: Gauged U(1) symmetry for right-handed fermions**

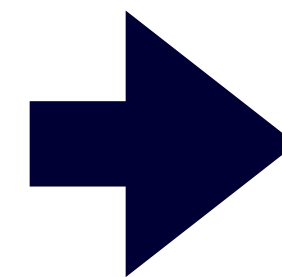
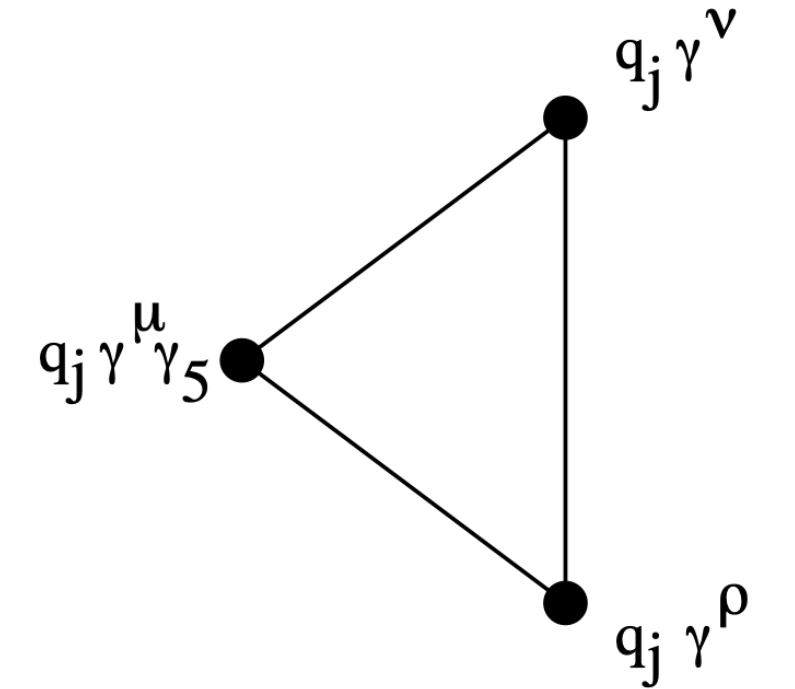
| scenario | symmetries          | $Q_L$          | $\ell_L$ | $U_R$          | $D_R$          | $E_R$ | $N_R$ | $H$ | $\psi_L$ | $\psi_R$ | $\chi_L$ | $\chi_R$ | $\eta_L$ | $\eta_R$ |
|----------|---------------------|----------------|----------|----------------|----------------|-------|-------|-----|----------|----------|----------|----------|----------|----------|
| (i)      | U(1) <sub>B-L</sub> | $+\frac{1}{3}$ | -1       | $+\frac{1}{3}$ | $+\frac{1}{3}$ | -1    | -1    | 0   | ×        | ×        | ×        | ×        | ×        | ×        |
| (ii)     | U(1) <sub>R</sub>   | 0              | 0        | -1             | +1             | -1    | +1    | 1   | ×        | ×        | ×        | ×        | ×        | ×        |
| (iii)    | U(1) <sub>B</sub>   | $+\frac{1}{3}$ | 0        | $+\frac{1}{3}$ | $+\frac{1}{3}$ | 0     | 0     | 0   | -1       | +2       | +2       | -1       | +2       | -1       |
| (iv)     | U(1) <sub>L</sub>   | 0              | +1       | 0              | 0              | +1    | +1    | 0   | -1       | +2       | +2       | -1       | +2       | -1       |

$$\partial_\mu \left( j_{B,u}^\mu \right) = \frac{1}{16\pi^2} \left( -\frac{4}{9} g'^2 F \widetilde{F} - g_R^2 F' \widetilde{F}' \right)$$

$$\partial_\mu \left( j_{B,d}^\mu \right) = \frac{1}{16\pi^2} \left( -\frac{1}{9} g'^2 F \widetilde{F} - g_R^2 F' \widetilde{F}' \right)$$

$$\partial_\mu \left( j_{L,E}^\mu \right) = \frac{1}{16\pi^2} \left( -g'^2 F \widetilde{F} - g_R^2 F' \widetilde{F}' \right)$$

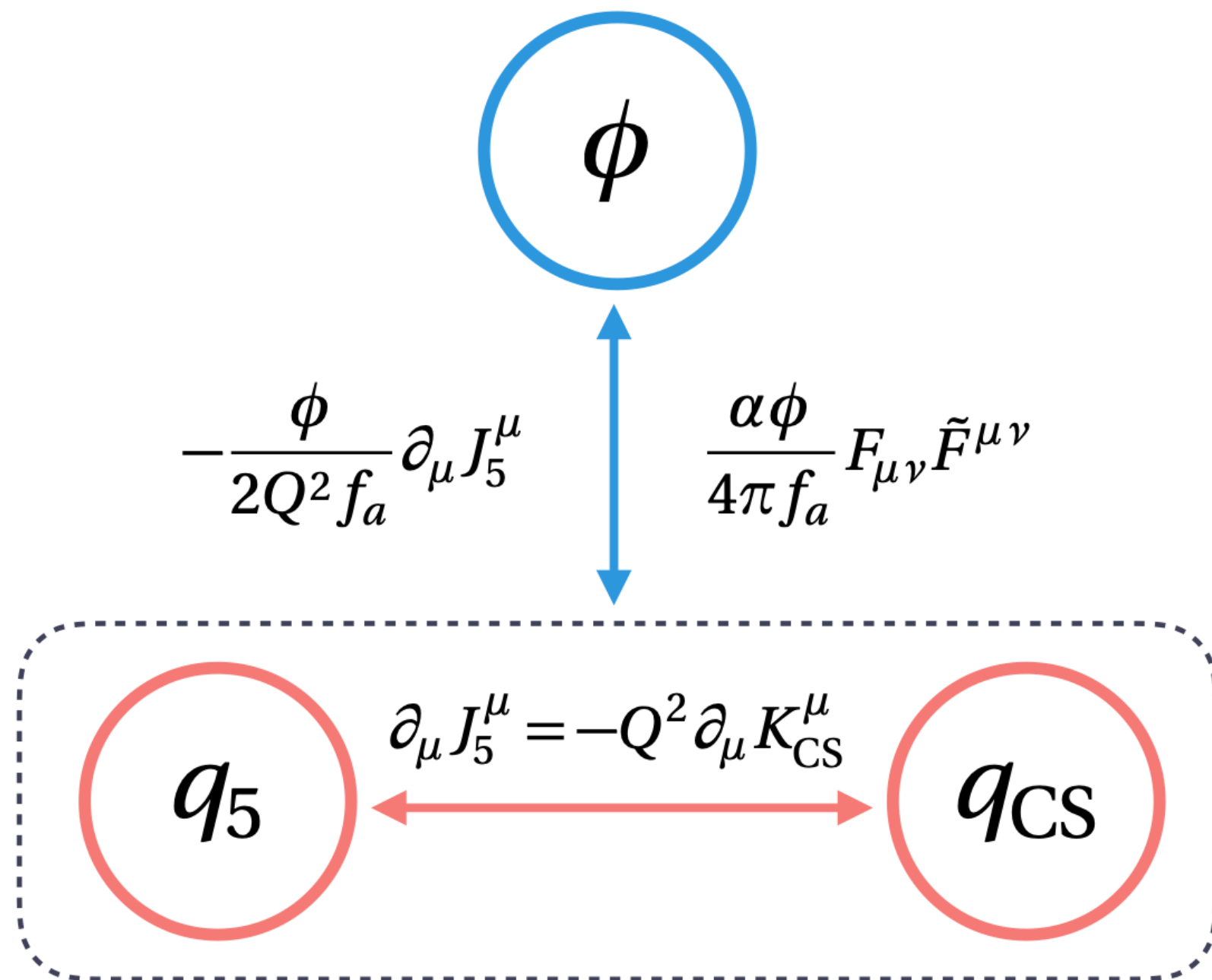
$$\partial_\mu \left( j_{L,N}^\mu \right) = \frac{1}{16\pi^2} \left( -g_R^2 F' \widetilde{F}' \right)$$



$$\partial_\mu \left( j_B^\mu - j_{L'}^\mu - j_N^\mu - j_e^\mu \right) = 0$$

# Axion-inflation baryogenesis

- **Axion inflation**  $\rightarrow q_{CS} \rightarrow q_5$



- **Chern Simons number:**

$$n_{CS} \equiv \frac{1}{(2\pi)^2} \mathcal{K}(\xi) a^3 H^3$$

$$= \frac{1}{(2\pi)^2} \sum_{\lambda=\pm} \lambda e^{i\kappa_\lambda \pi} \int \tilde{\tau}^3 d \ln \tilde{\tau} W_{\kappa_\lambda, \mu}^*(-2i\tilde{\tau}) W_{\lambda_\sigma, \mu}(-2i\tilde{\tau}) a^3 H^3$$

- **Chiral fermion asymmetry during reheating**

$$n_{f,\sigma} = -\epsilon_\sigma N_{f,\sigma} \frac{g_X^2}{8\pi^2 a^3} n_{CS} = -\epsilon_i N_i \frac{g_X^2}{2(2\pi)^4} H^3 \mathcal{K}(\xi)$$

# Axion-inflation baryogenesis

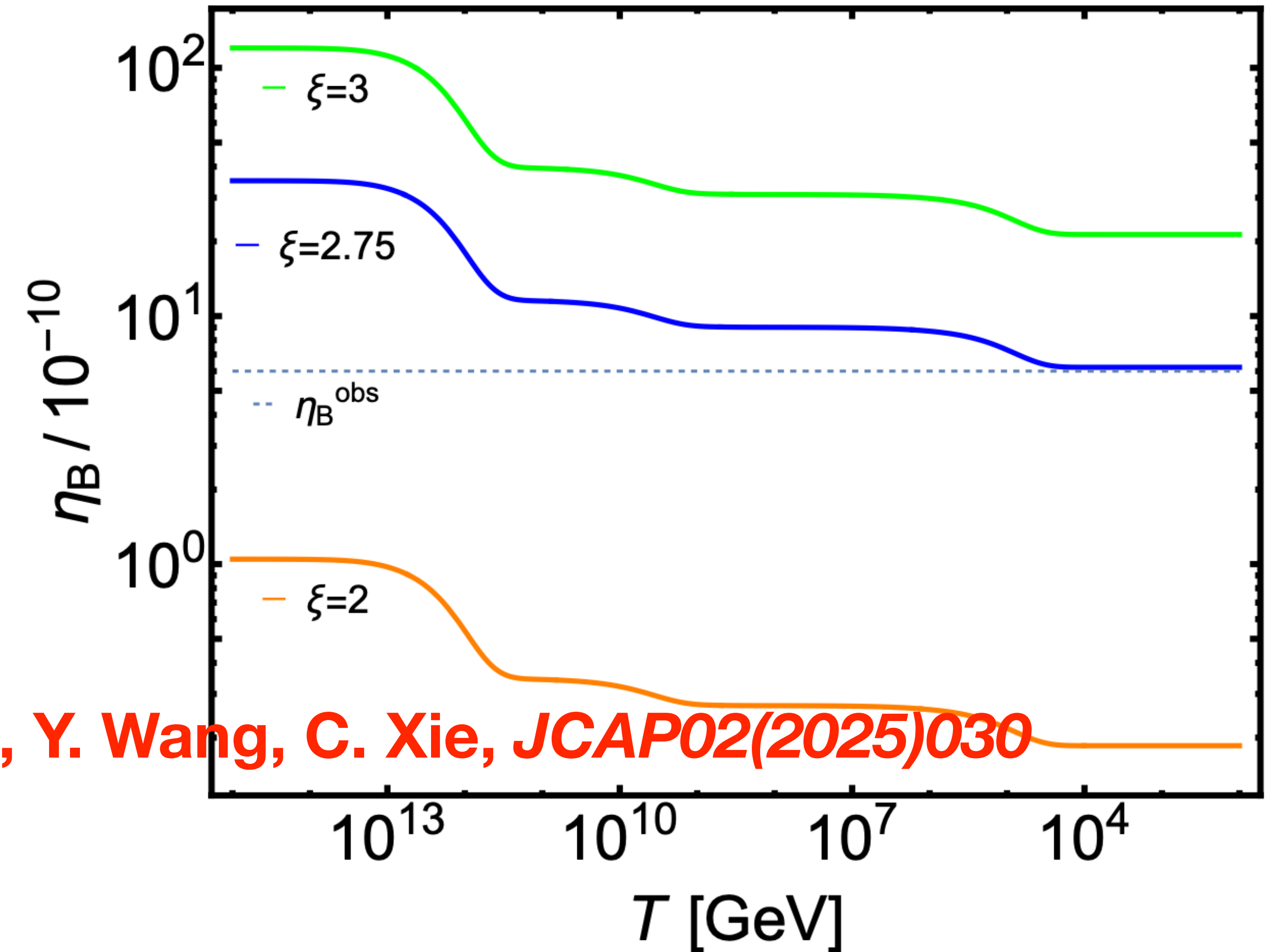
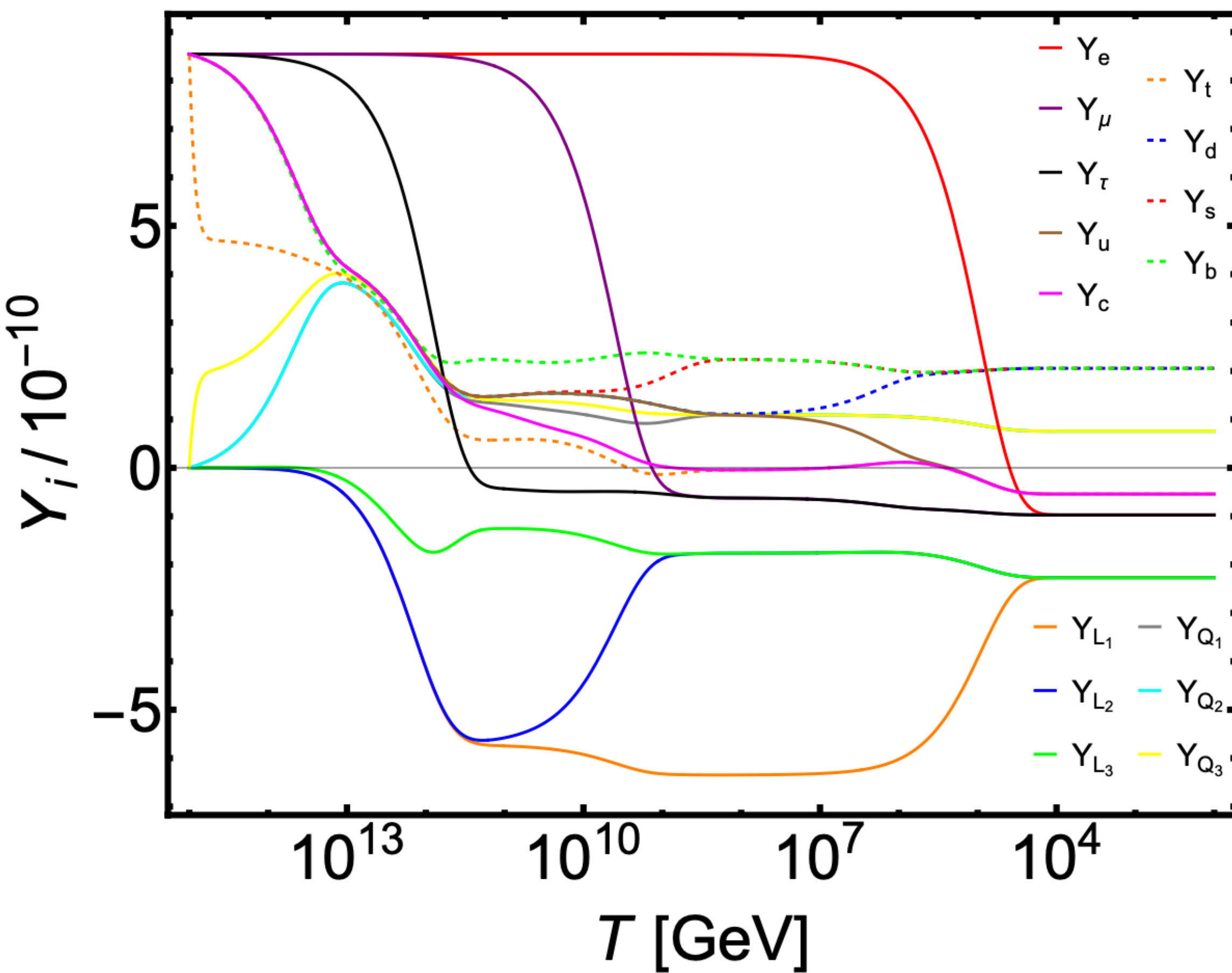
For the case in which neutrinos are Majorana particle

$$Y_B = \frac{28}{79} Y_{B-L'} = \frac{28}{79} \frac{n_e}{s} \bigg|_{T=T_{\text{FO}}}$$

- **Concrete value depends on the input of the axion inflation parameters!**

# Axion-inflation baryogenesis

If neutrinos are Dirac particles!



Chao, Y. Wang, C. Xie, JCAP02(2025)030

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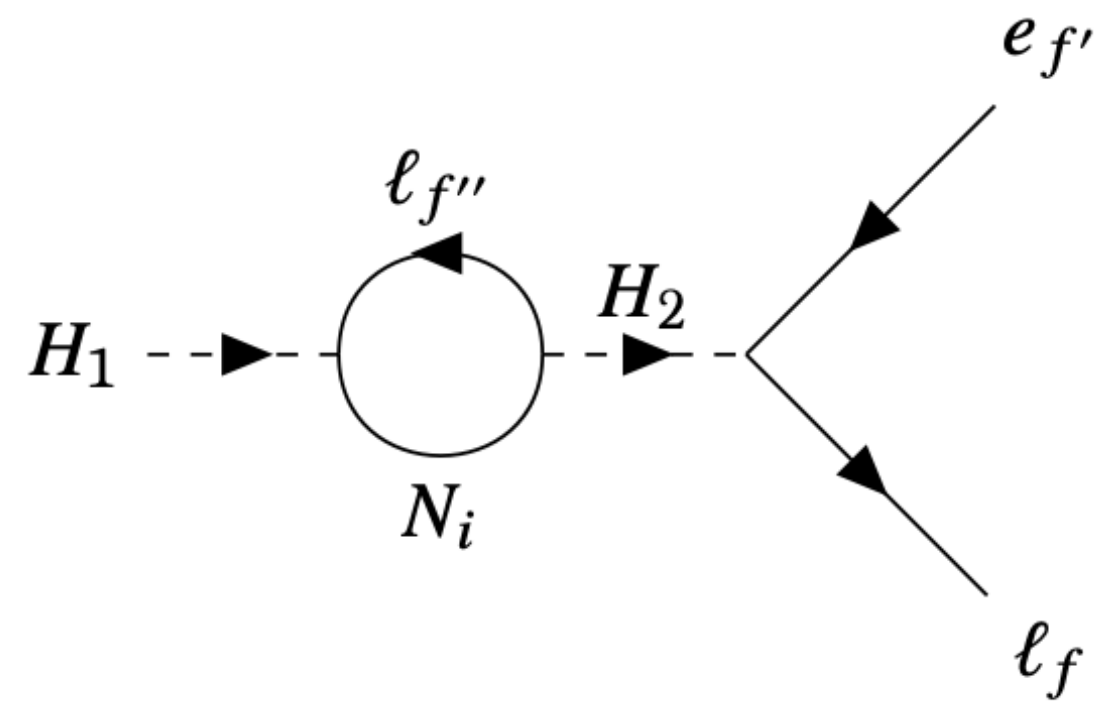
# Leptogenesis

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- Scenario-1 : Axion inflation Leptogenesis
- **Scenario-2 : Eogenesis**

# Eogenesis

- **Higgs doublet decay into chiral electrons with CP violation**



$$\varepsilon = \frac{1}{8\pi} \frac{\text{Im} \left[ Y_{fg}^1 Y_{gf}^{2\dagger} \text{tr}(Y^{1\dagger} Y^2) \right]}{\text{tr}(Y^{1\dagger} Y^1)} f \left( \frac{M_{\Phi_2}}{M_{\Phi_1}} \right)$$

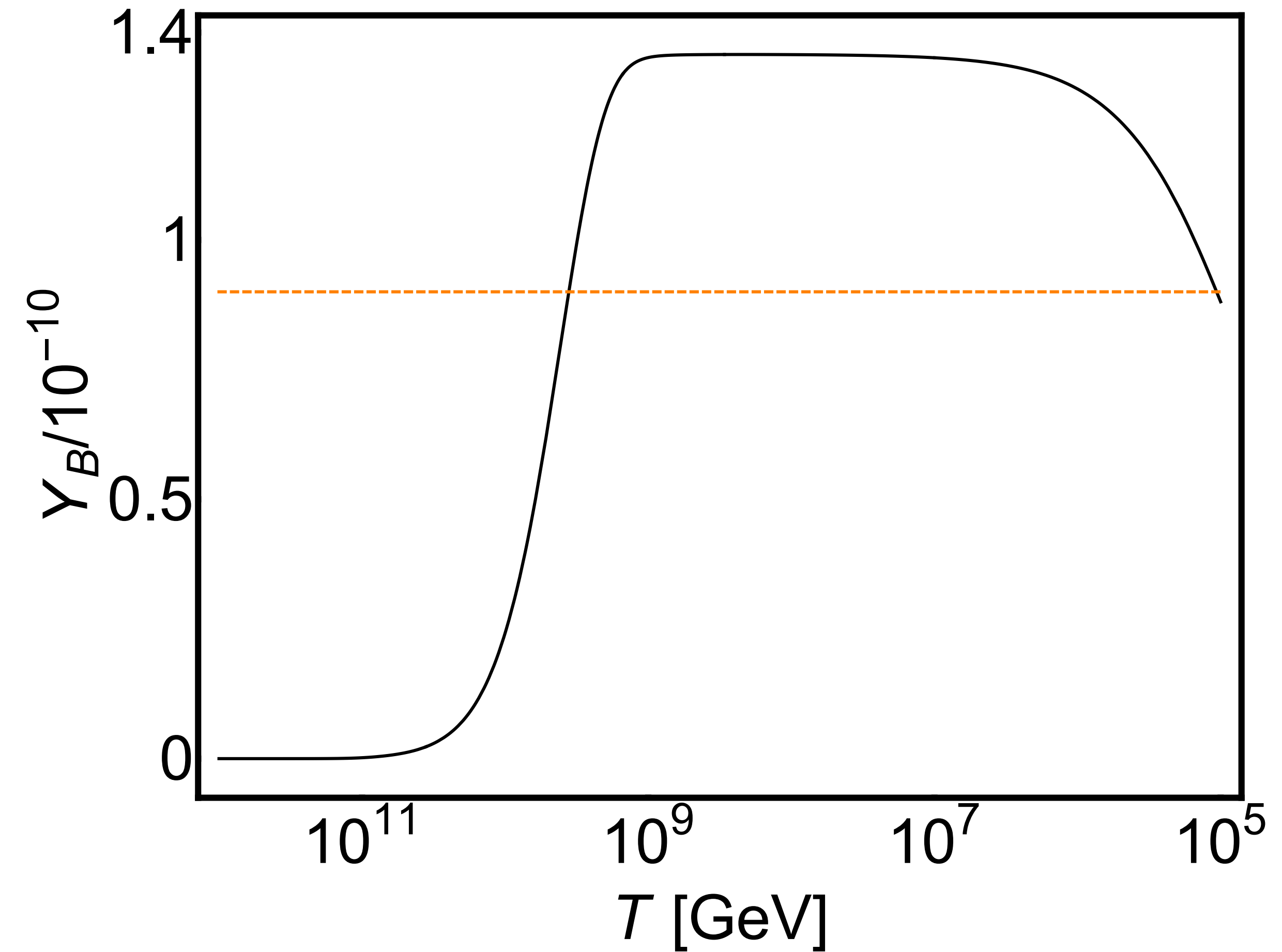
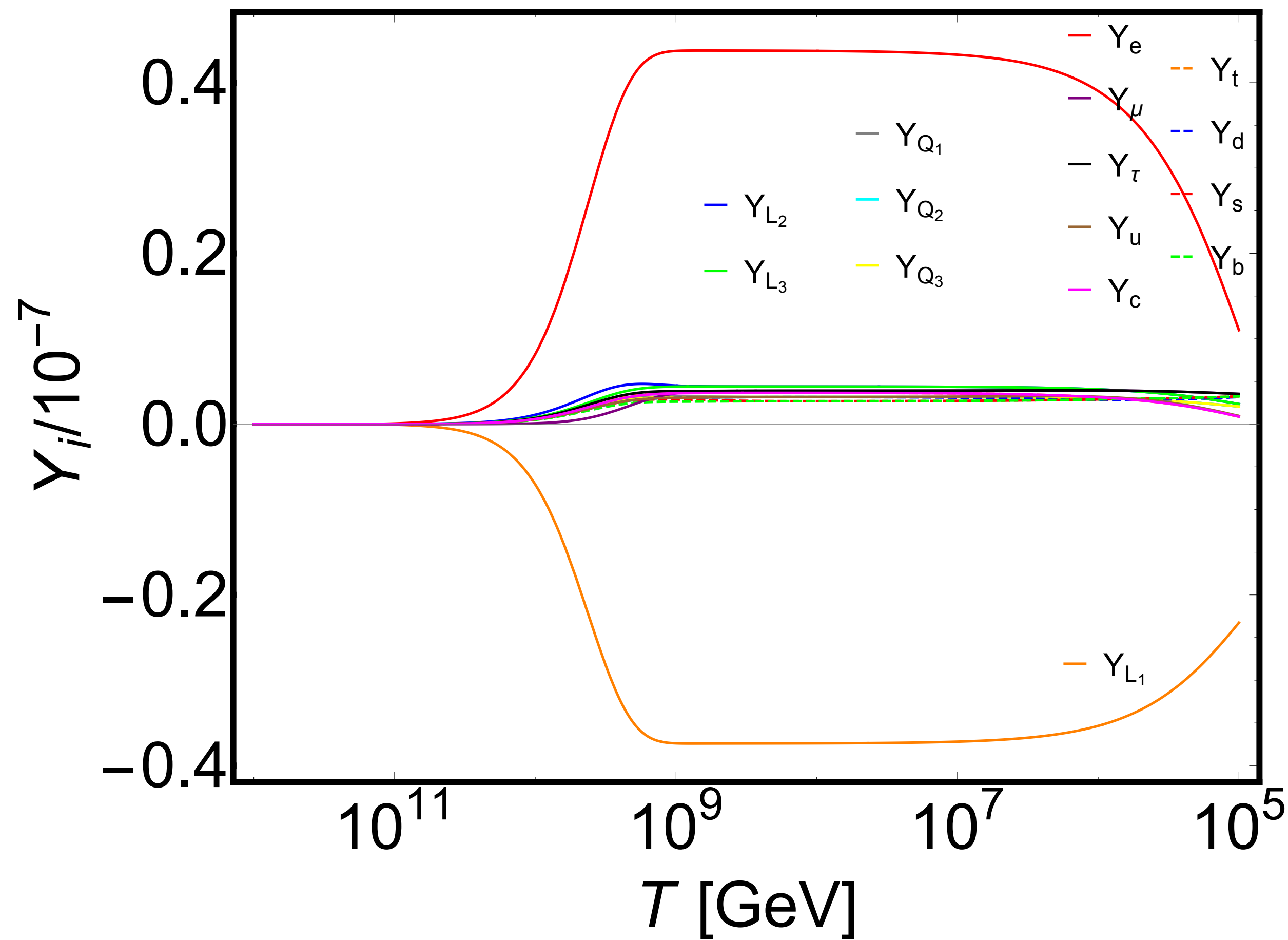
- **Chiral asymmetries are generated in the first generation.**
- **No primordial B-L violation is generated!**
- **No matter right-handed neutrinos exist or not!**

# Eogenesis

- **Transport equations**

$$\begin{aligned} -\frac{dY_{\Sigma}}{d \ln T} &= -\frac{\gamma_D}{H} \left[ Y_{\Sigma} - Y_{\Sigma}^{\text{EQ}} \right] \\ -\frac{d}{d \ln T} \left( \frac{\mu_{\Phi}}{T} \right) &= -2 \frac{\gamma_D}{H} \left( \frac{\mu_{\Phi}}{T} - \frac{4}{3} \frac{\mu_{L_k}}{T} \right) \\ \gamma_D &= \frac{K_1(z)}{K_2(z)} \Gamma_{\Phi} \\ -\frac{d}{d \ln T} \left( \frac{\mu_{L_k}}{T} \right) &= -\frac{1}{g_{L_k}} \frac{\gamma_{WS}}{H} \left[ \sum_{i=1}^3 \left( \frac{\mu_{L_i}}{T} + 3 \frac{\mu_{Q_i}}{T} \right) \right] \\ &\quad -\frac{1}{g_{L_k}} \frac{\gamma_{Y_{E_k}}}{H} \left( -\frac{\mu_{E_k}}{T} + \frac{\mu_{L_k}}{T} - \frac{\mu_H}{T} \right) \\ &\quad + \frac{1}{g_{L_k}} \frac{\gamma_D}{H} \epsilon \left[ \frac{4\pi^2 g_*^S}{15} Y_{\Sigma} - \frac{48\zeta(3)}{\pi^2} \right] \\ &\quad -\frac{2}{g_{L_k}} g_{\Phi} \frac{\gamma_D}{H} \left( \frac{4}{3} \frac{\mu_{L_k}}{T} - \frac{\mu_{\Phi}}{T} \right) \end{aligned}$$

# Eogenesis



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# Conclusion

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A new type of Leptogenesis is proposed, which is dubbed as **Eogenesis**

- No primordial B-L is needed
- No dependence on the nature of neutrinos

**Thank you for your attention!**