



Silk damping in induced gravitational waves: a novel probe for new physics

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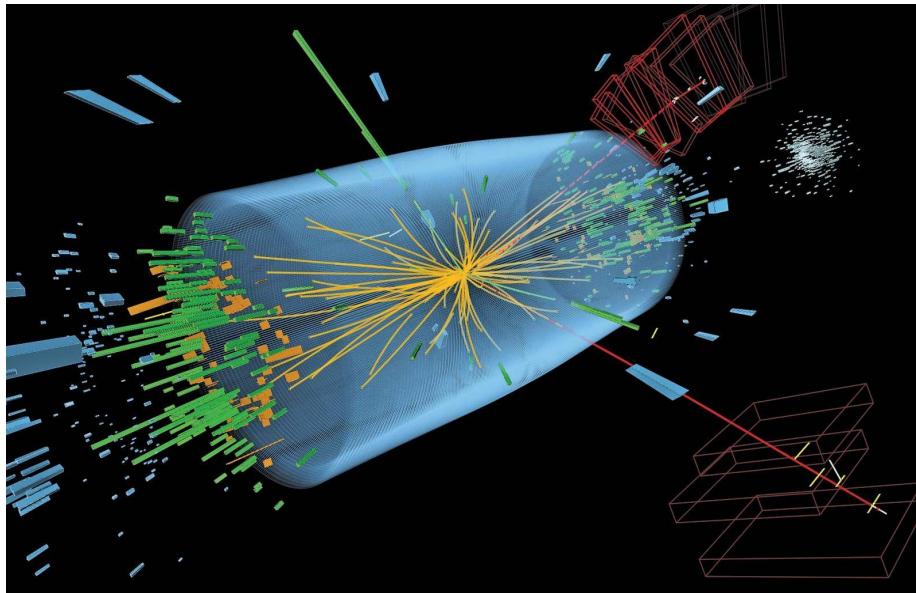
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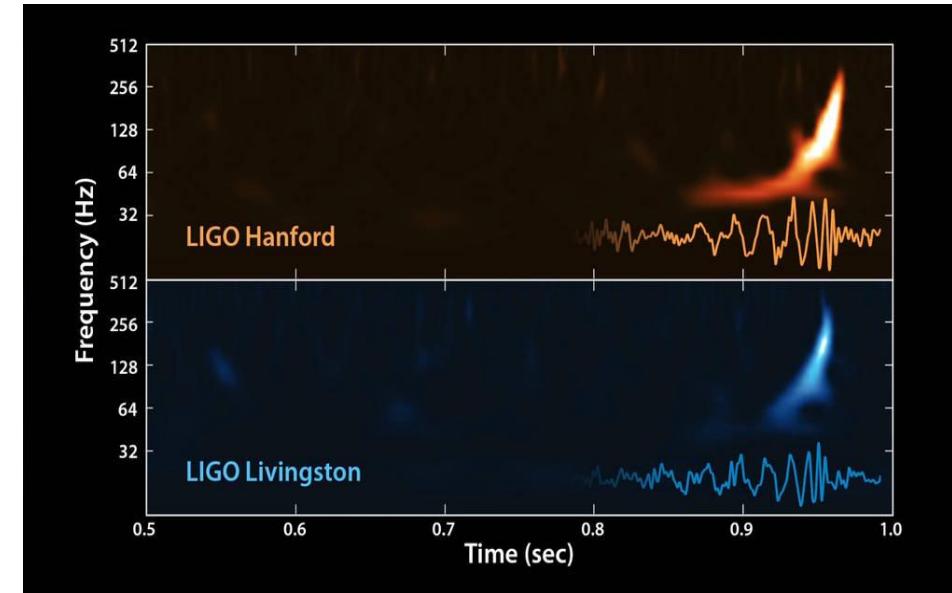
Based on: **YHY** and Sai Wang: *Sci.China Phys.Mech.Astron.* 68 (2025) 210412, arXiv: 2405.02960

Two breakthroughs in 21st century



Higgs bosons (2012)

additional symmetry breakings?

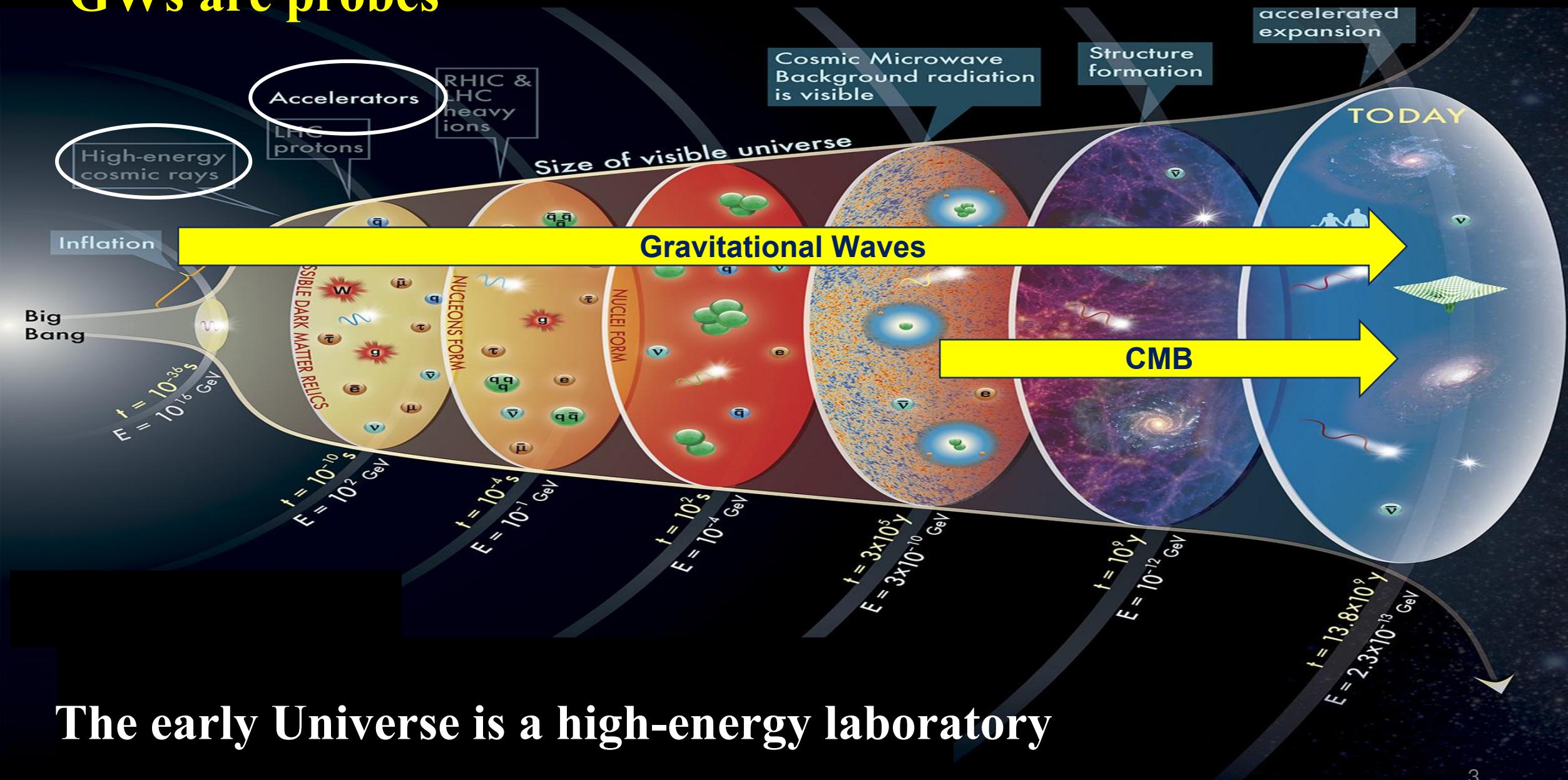


Gravitational waves (2016)

GW astronomy & GW cosmology

GWs provide new opportunities to detect fundamental particle physics!

GWs are probes

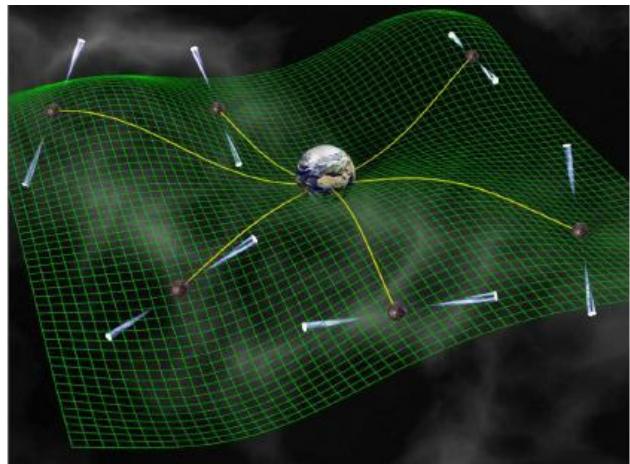


GWs in the early Universe

- Primordial GWs
- Preheating GWs
- GWs from first-order phase transition
- GWs from cosmological topological defects
- Induced GWs
-

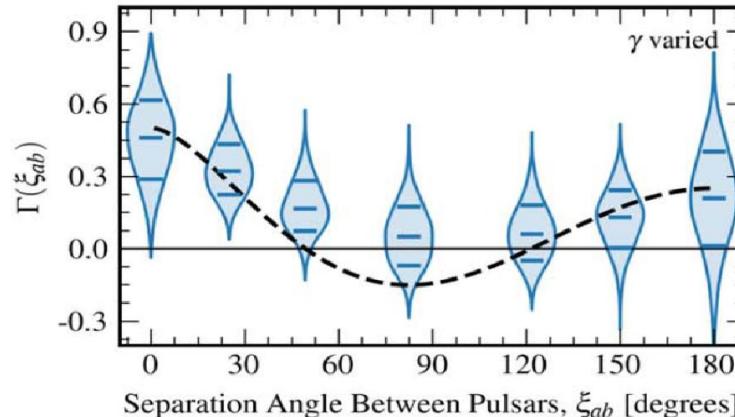
Evidence for GWB

PTA experiments

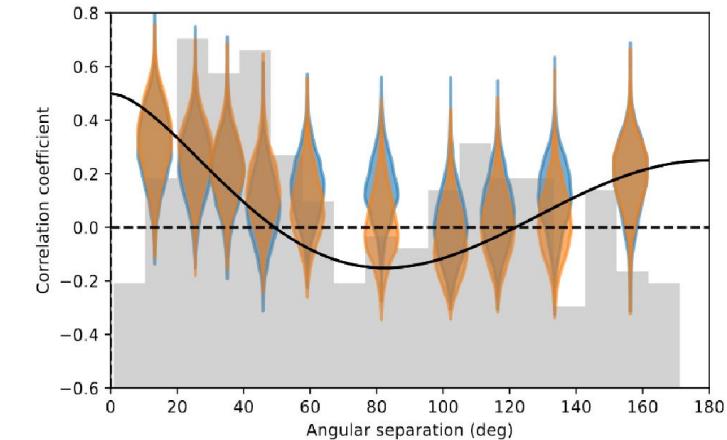


- ApJL 951 (2023) L8: NANOGrav*
A&A 678 (2023) A50: EPTA, InPTA
ApJL 951 (2023) L6: PPTA
RAA 23 (2023) 075024: CPTA

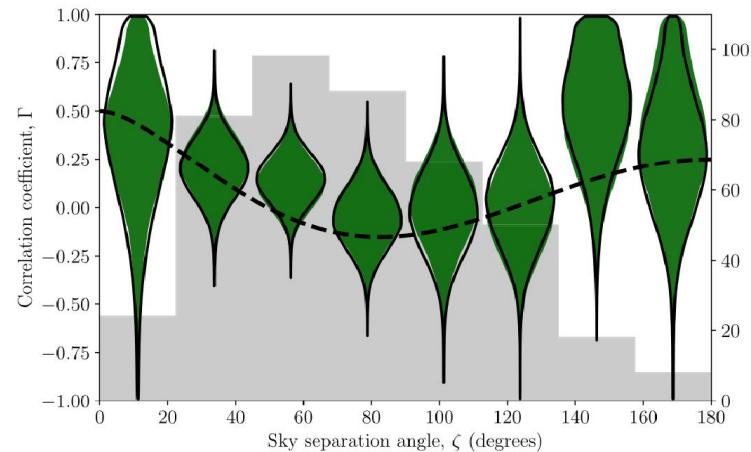
NANOGrav



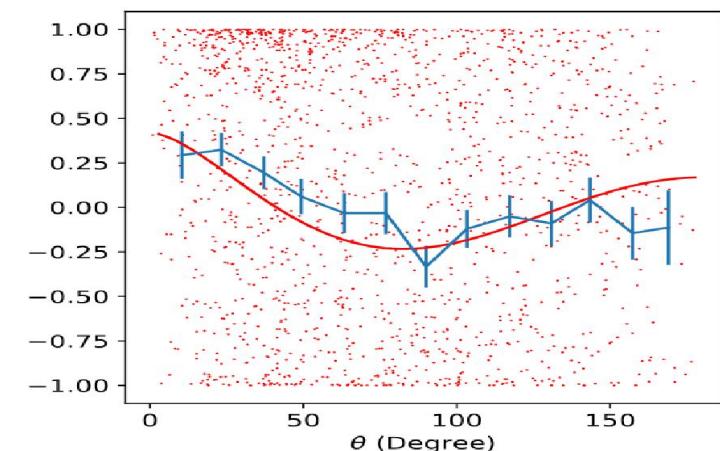
EPTA



PPTA



CPTA



Induced GWs

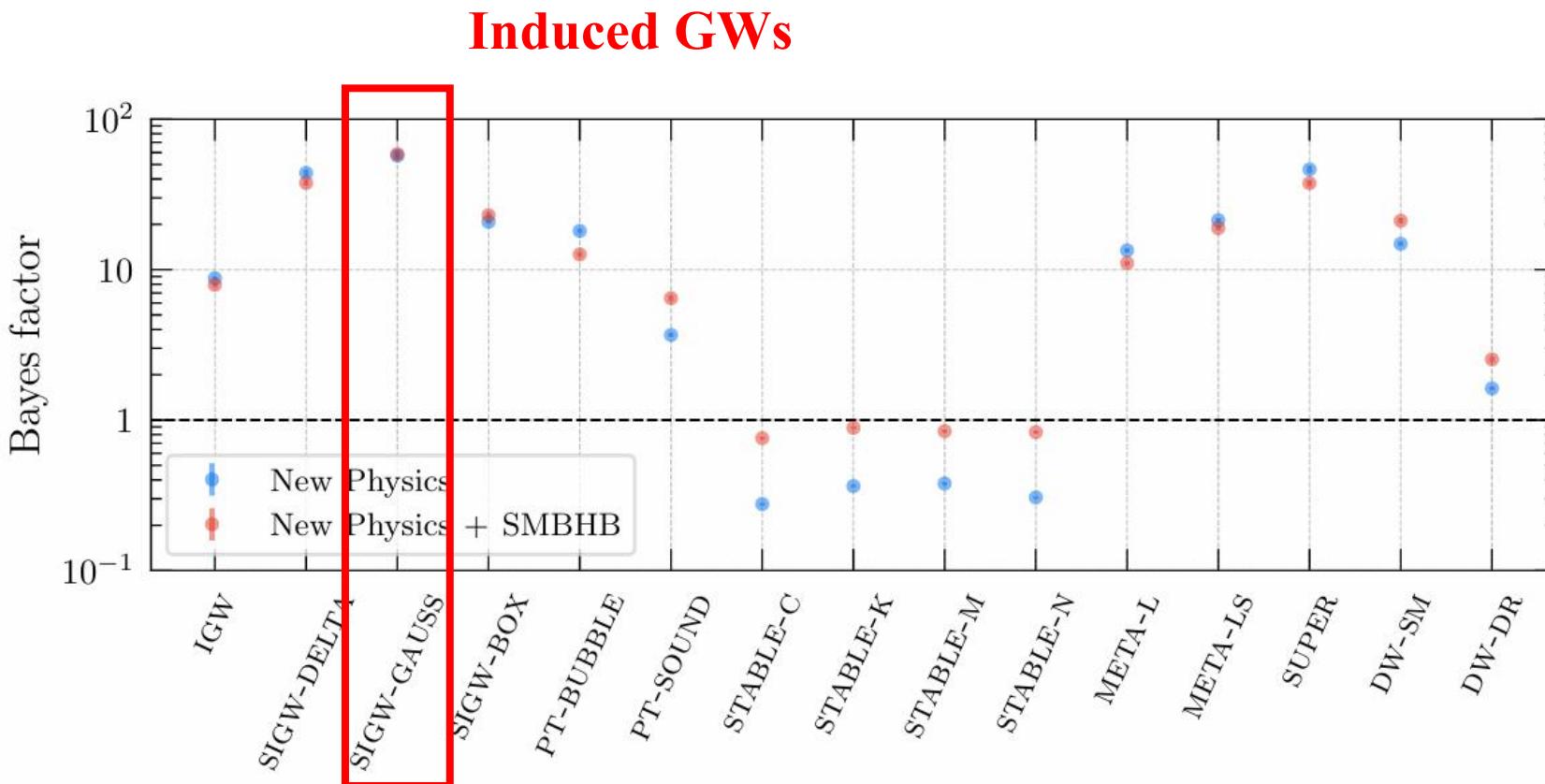


Figure 2. Bayes factors for the model comparisons between the new-physics interpretations of the signal considered in this work and the interpretation in terms of SMBHBs alone. Blue points are for the new physics alone, and red points are for the new physics in combination with the SMBHB signal. We also plot the error bars of all Bayes factors, which we obtain following the bootstrapping method outlined in Section 3.2. In most cases, however, these error bars are small and not visible.

Induced GWs

- **Spacetime metric:** $ds^2 = a^2 \left\{ -(1 + 2\phi^{(1)}) d\tau^2 + \left[(1 - 2\phi^{(1)}) \delta_{ij} + \frac{1}{2} h_{ij}^{(2)} \right] dx^i dx^j \right\}$

- **Induced GWs:** $h''_{\lambda, \mathbf{k}} + 2\mathcal{H}h'_{\lambda, \mathbf{k}} - \nabla^2 h_{\lambda, \mathbf{k}} = 4\mathcal{S}_{\lambda, \mathbf{k}}$ $h^{(2)} \sim \phi^{(1)} \phi^{(1)}$

$$\mathcal{S}_{\lambda, \mathbf{k}} = \int \frac{d^3 \mathbf{q}}{(2\pi)^{3/2}} \epsilon_{\mathbf{q}, ij}^\lambda q^i q^j \left\{ 2\phi_{\mathbf{k}-\mathbf{q}} \phi_{\mathbf{q}} + \frac{4}{3(1+w)\mathcal{H}^2} [\phi'_{\mathbf{k}-\mathbf{q}} + \mathcal{H}\phi_{\mathbf{k}-\mathbf{q}}] [\phi'_{\mathbf{q}} + \mathcal{H}\phi_{\mathbf{q}}] \right\}$$

- **Energy-density fraction spectrum:**

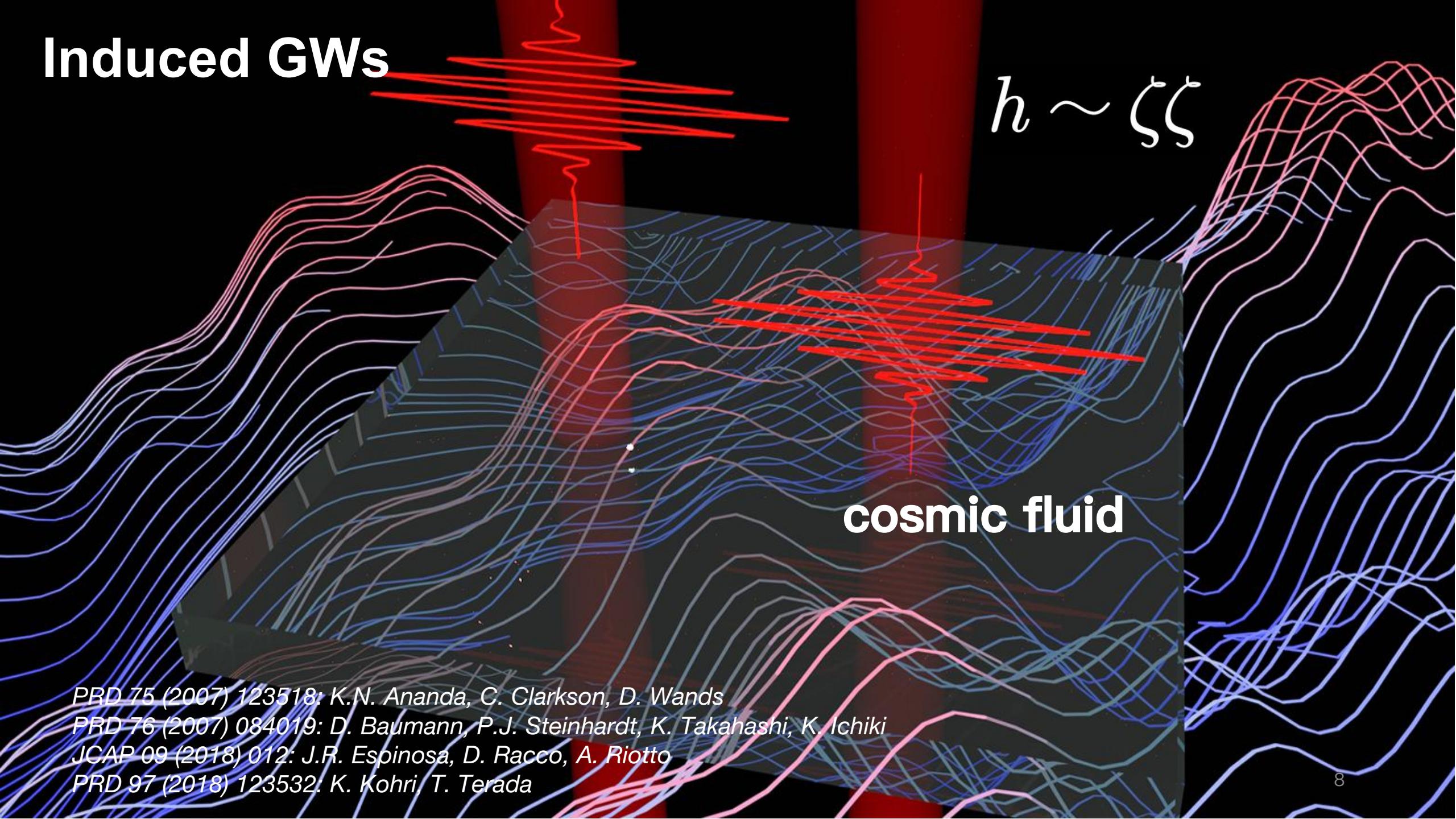
$$\Omega_{\text{gw}}(k) = \frac{k^2}{6a^2 H^2} \int_0^\infty du \int_{|1-u|}^{|1+u|} dv \left[\frac{4v^2 - (1+v^2-u^2)^2}{4uv} \right]^2 \times \frac{1}{I^2(u, v, k, \tau)} \mathcal{P}_\zeta(uk) \mathcal{P}_\zeta(vk)$$

$\phi_{\mathbf{k}}^{(1)}(\tau) = \Phi(k, \tau) |\zeta_{\mathbf{k}}$

classical evolution
after inflation

primordial curvature
perturbations

Induced GWs



PRD 75 (2007) 123518: K.N. Ananda, C. Clarkson, D. Wands

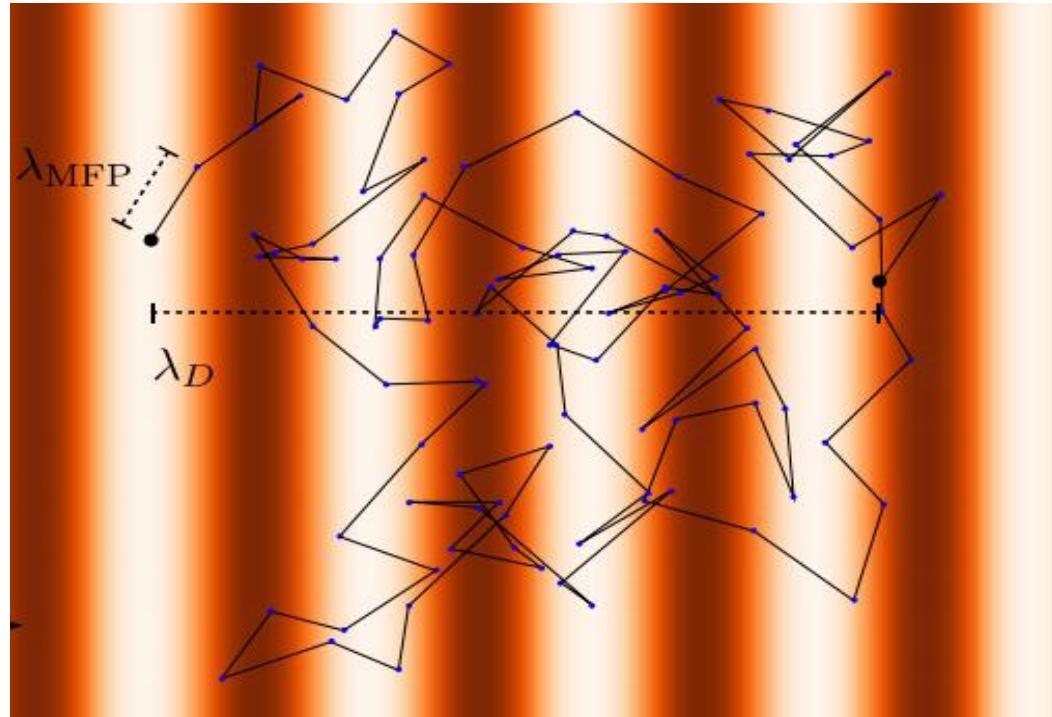
PRD 76 (2007) 084019: D. Baumann, P.J. Steinhardt, K. Takahashi, K. Ichiki

JCAP 09 (2018) 012: J.R. Espinosa, D. Racco, A. Riotto

PRD 97 (2018) 123532: K. Kohri, T. Terada

Silk damping

Cosmic fluid

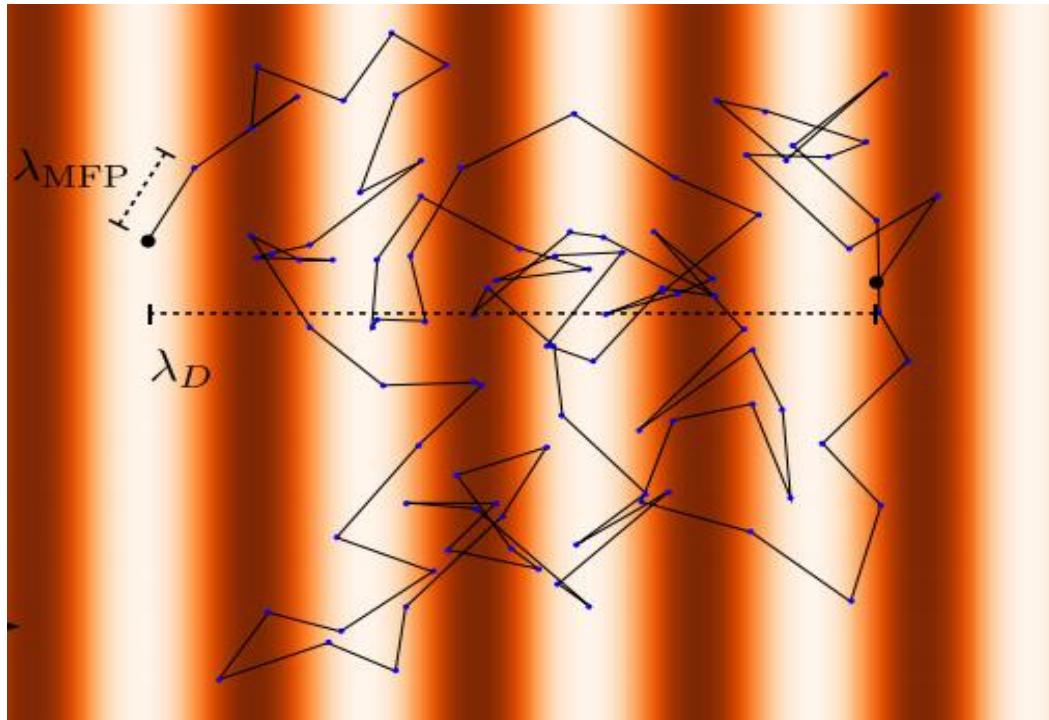


perturbation scale $k^{-1} \ll$ diffusion scale k_D^{-1}

Diffusion erases perturbations !

Silk damping in the CMB

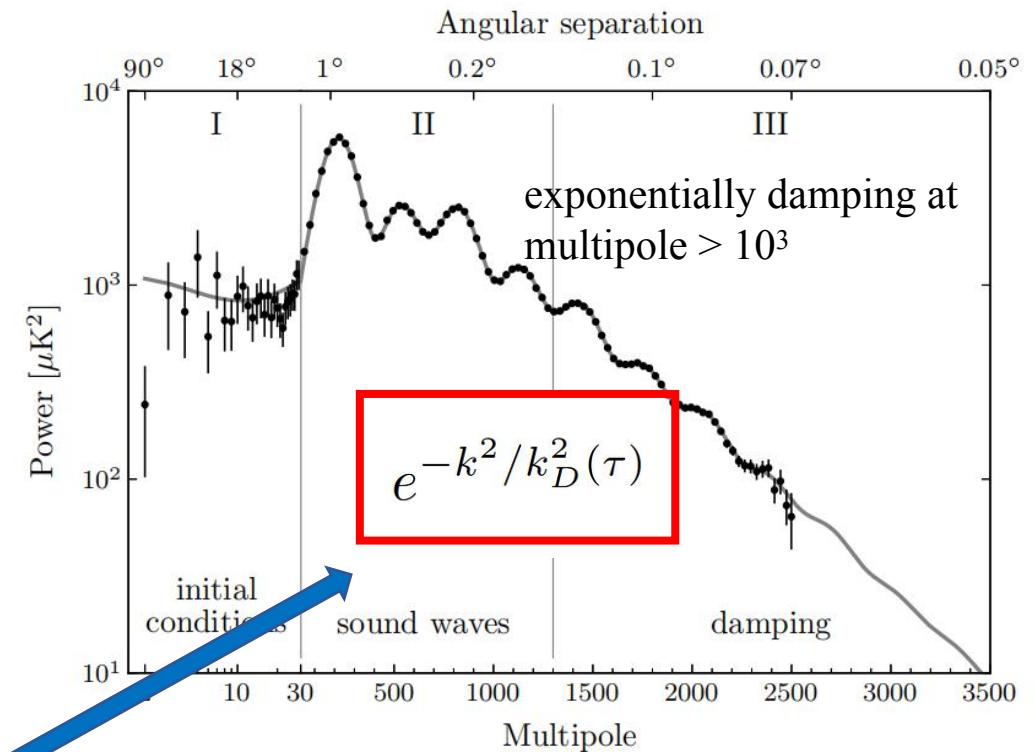
Photon–baryon plasma



perturbation scale $k^{-1} \ll$ diffusion scale k_D^{-1}

Diffusion erases perturbations !

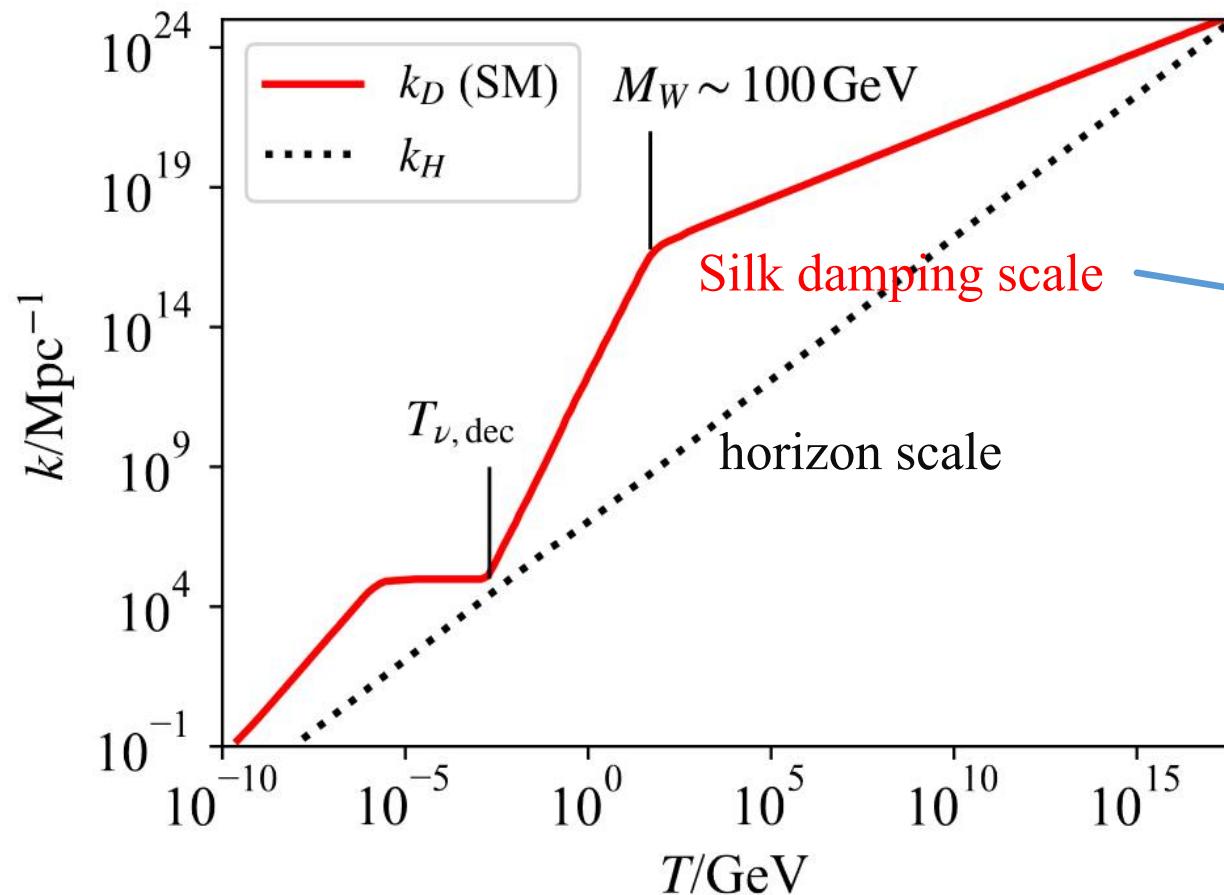
CMB anisotropy power spectrum



“Silk damping”

Nature 215 (1967) 5106: J. Silk

Silk damping scale



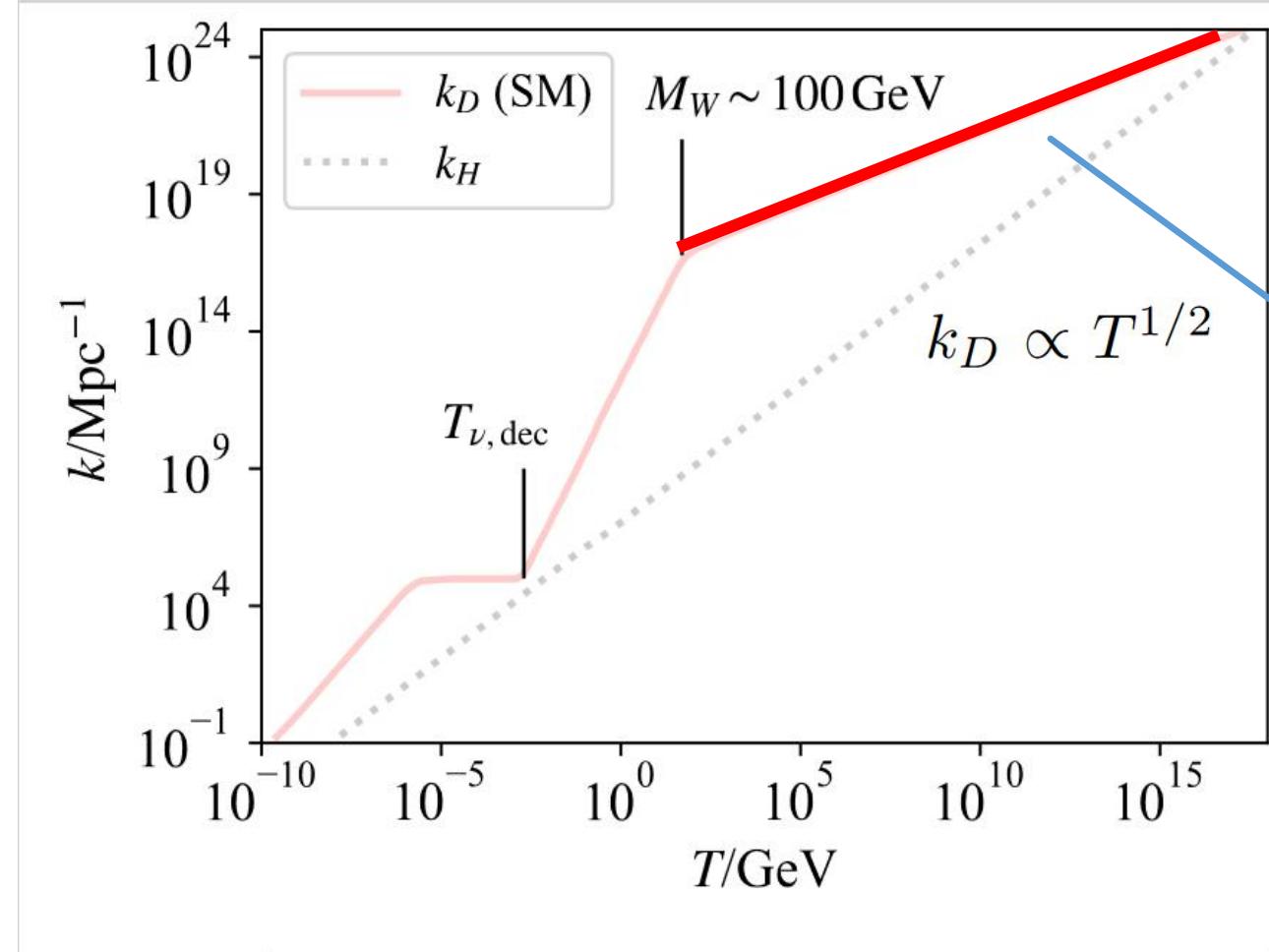
- **Microscopic origin:**

particle interaction in cosmic fluid

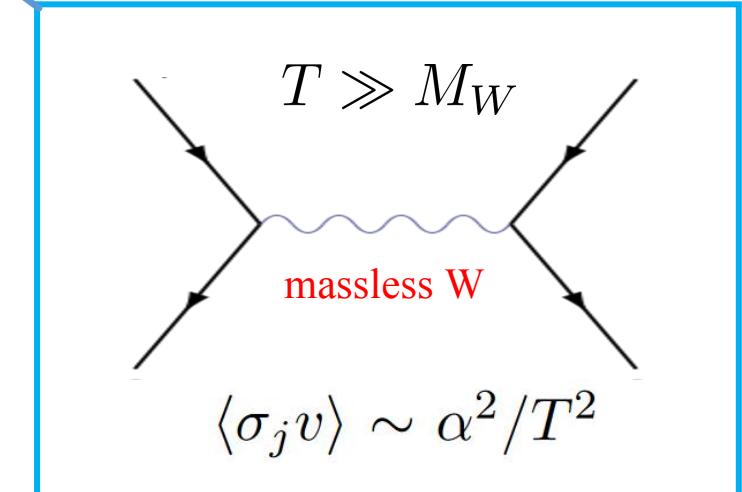
$$k_D^{-2}(\tau) \simeq \int_0^\tau d\tau' \frac{2\eta(\tau')}{3a(\rho+p)}$$
$$\eta(\tau) \simeq \sum_{j=\nu,X,\dots} \frac{4}{15} \rho_j t_j \Theta(\tau_{j,dec} - \tau)$$
$$t_j = 1/(n_j \langle \sigma_j v \rangle)$$

thermally-averaged
cross section

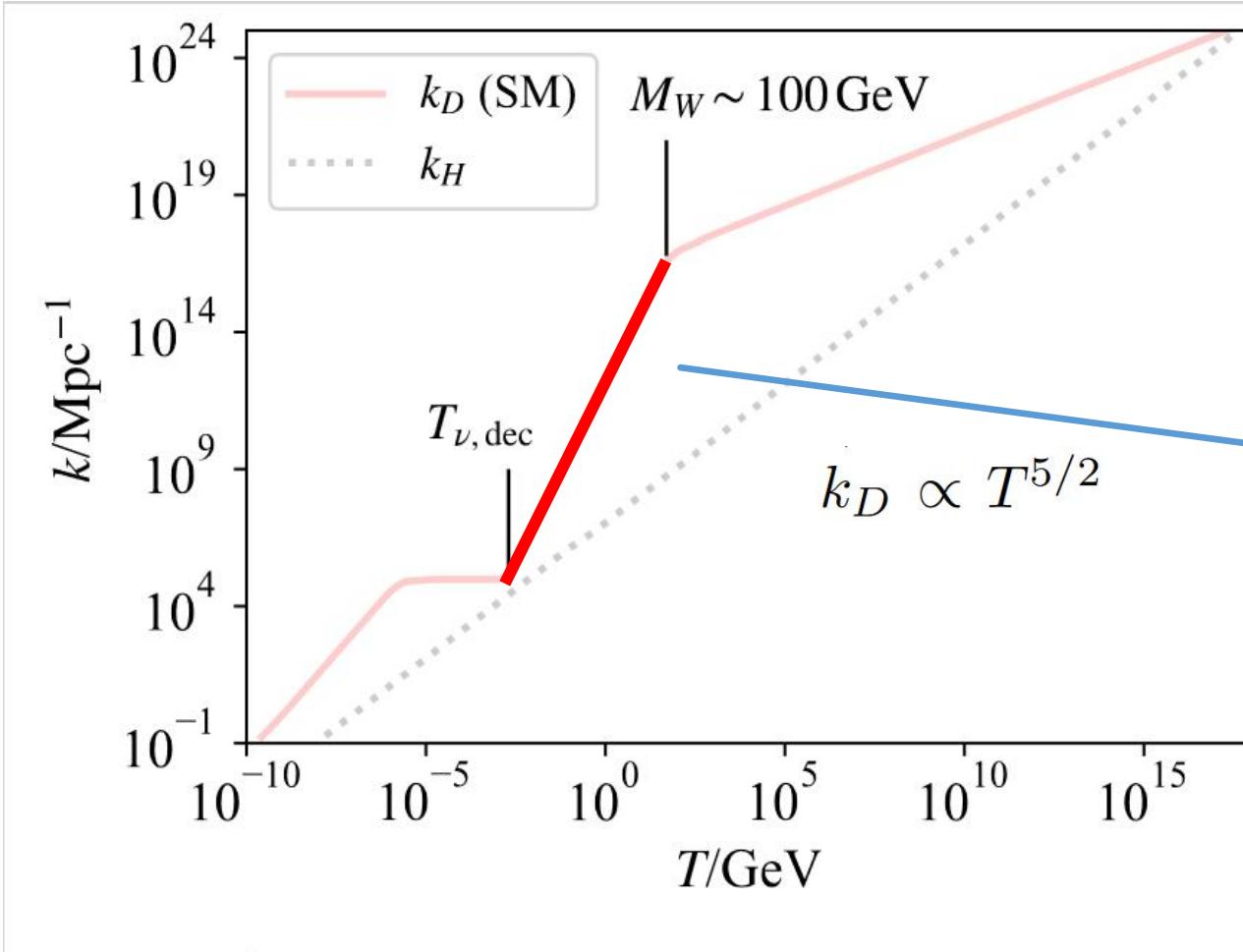
Silk damping scale



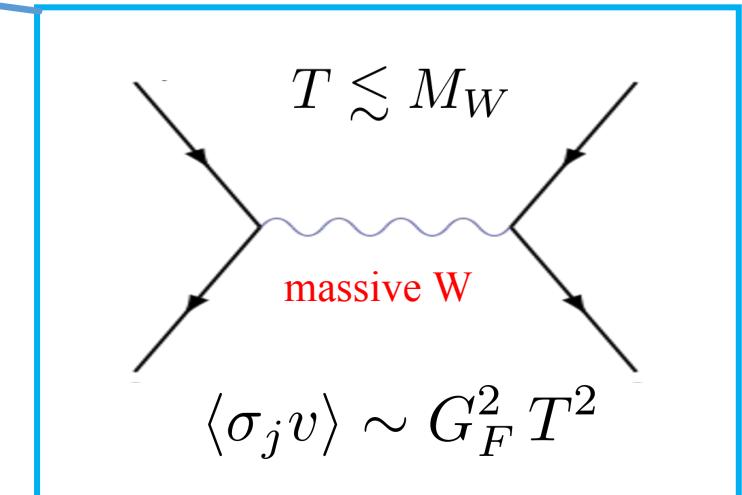
- **Microscopic origin:**
particle interaction in cosmic fluid
- **Within the SM:**
dominated by neutrinos



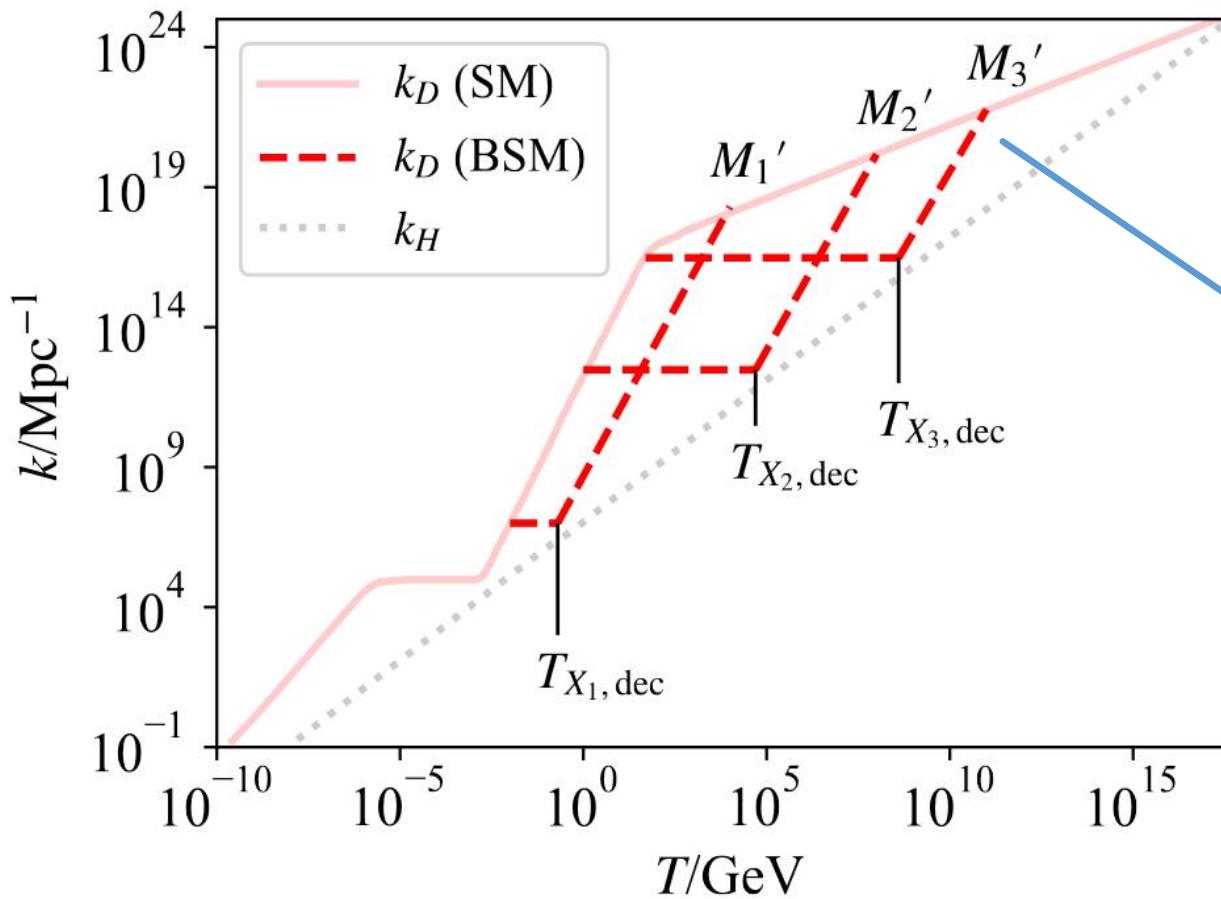
Silk damping scale



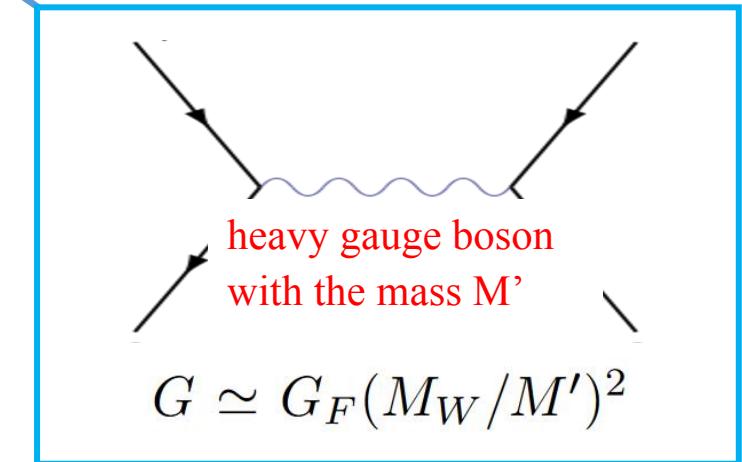
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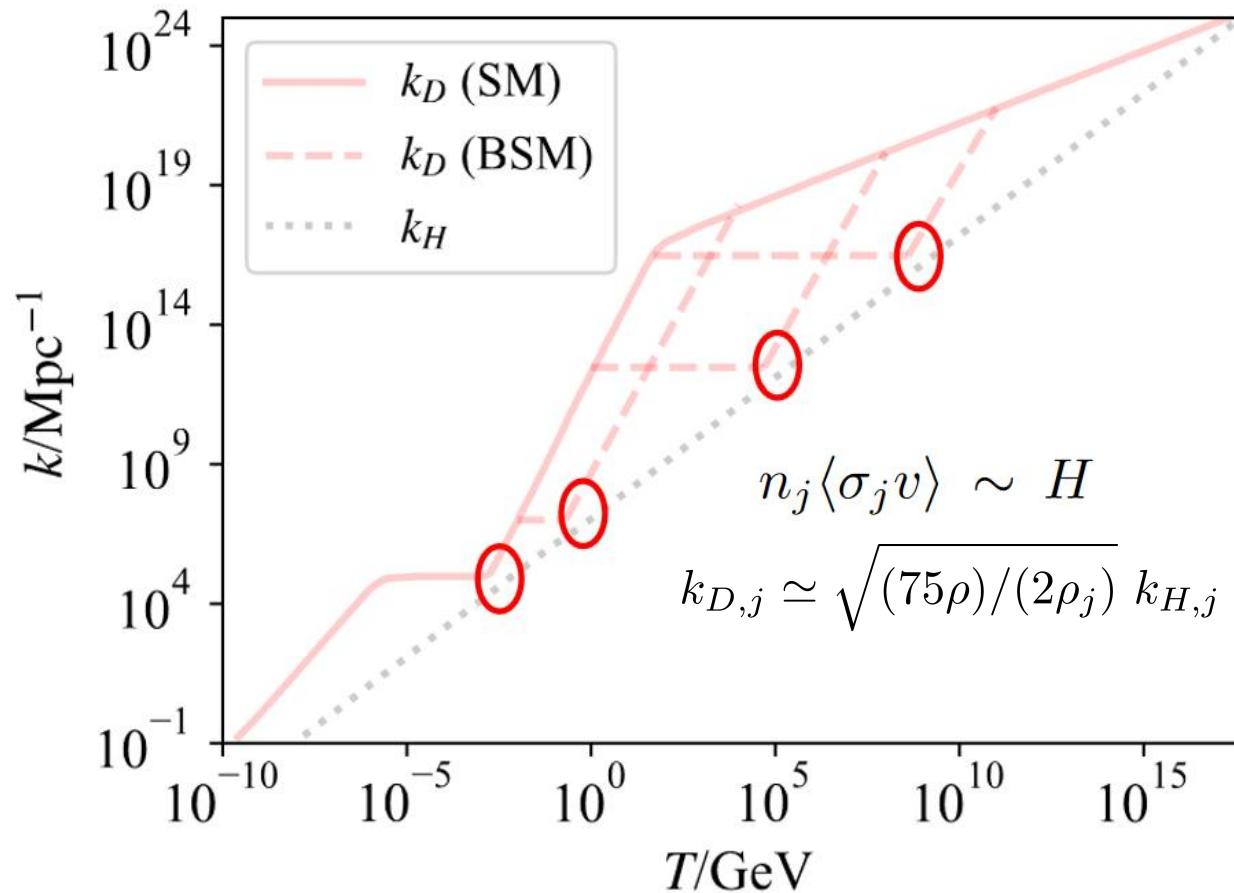
Silk damping scale



- **Microscopic origin:**
particle interaction in cosmic fluid
- **Beyond the SM:**
dominated by weakest-interacting particles



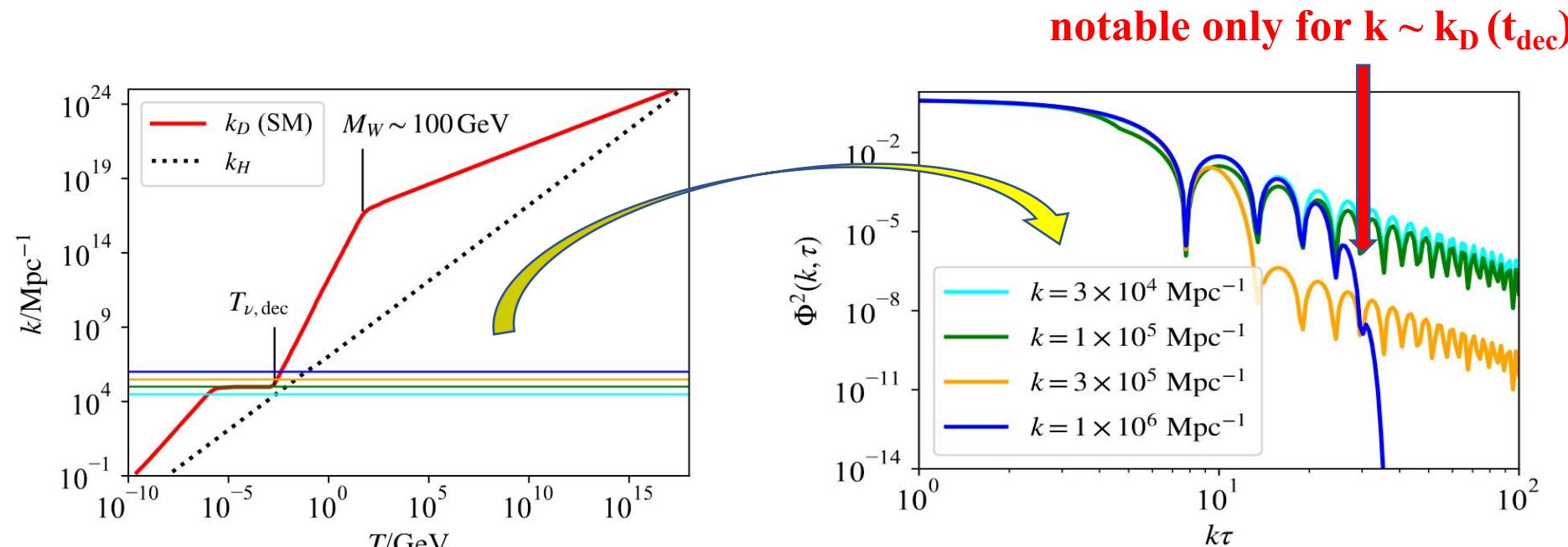
Silk damping scale



- **Microscopic origin:**
particle interaction in cosmic fluid
 - **Within the SM:**
dominated by neutrinos
 - **Beyond the SM:**
dominated by weakest-interacting particles
 - **At particle decoupling:**
Silk damping scale \sim horizon scale
- important for induced GWs

Silk damping in induced GWs

SCPMA 68 (2025) 210412: YHY and S. Wang



$$\phi_{\mathbf{k}}^{(1)}(\tau) = \Phi(k, \tau) \zeta_{\mathbf{k}}$$

perfect fluid, RD

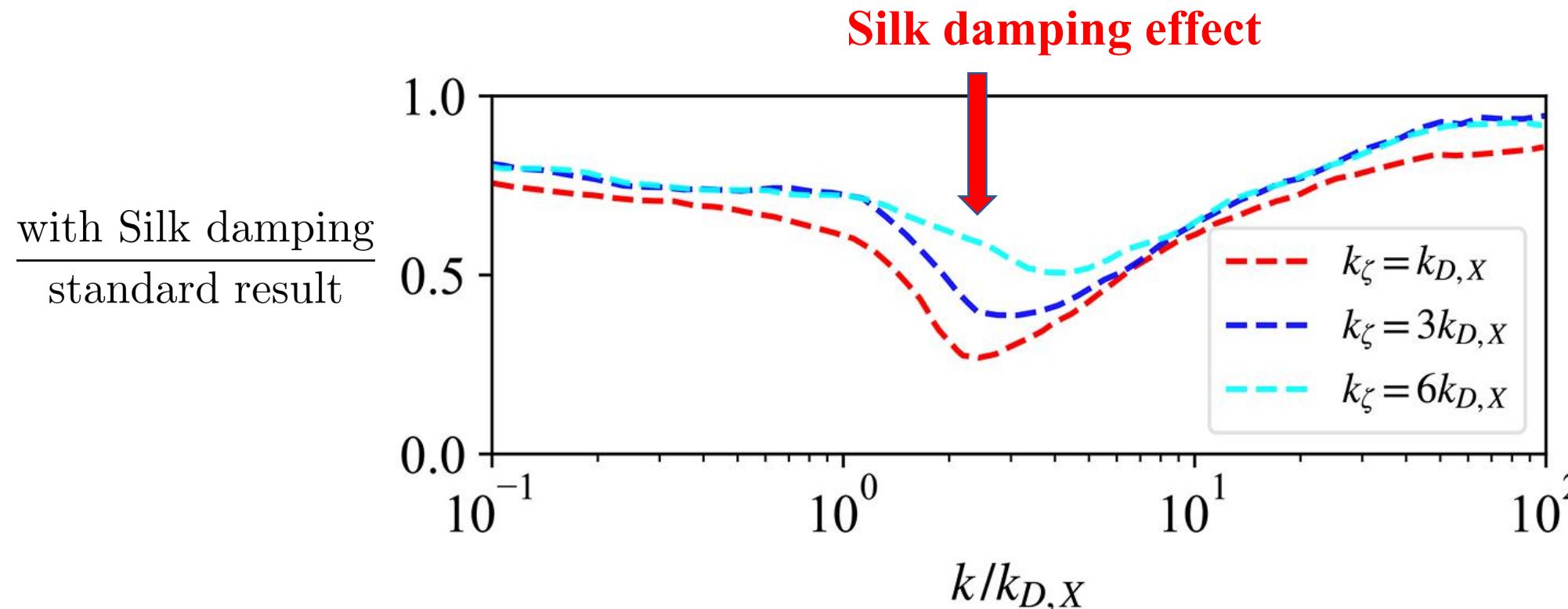
Silk damping

$$\Phi(k, \tau) \simeq \frac{9}{(k\tau)^2} \left(\frac{\sqrt{3}}{k\tau} \sin \frac{k\tau}{\sqrt{3}} - \cos \frac{k\tau}{\sqrt{3}} \right) \times e^{-k^2/k_D^2(\tau)}$$

$$\Omega_{\text{gw}}(k) = \frac{k^2}{6a^2H^2} \int_0^\infty du \int_{|1-u|}^{|1+u|} dv \left[\frac{4v^2 - (1+v^2-u^2)^2}{4uv} \right]^2 \times \frac{1}{I^2(u, v, k, \tau)} \mathcal{P}_\zeta(uk) \mathcal{P}_\zeta(vk)$$

Results

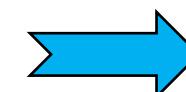
SCPMA 68 (2025) 210412: YHY and S. Wang



Silk damping can suppress
the GW spectrum by $\sim 70\%$
at the scale $k \sim k_D(t_{dec,X})$



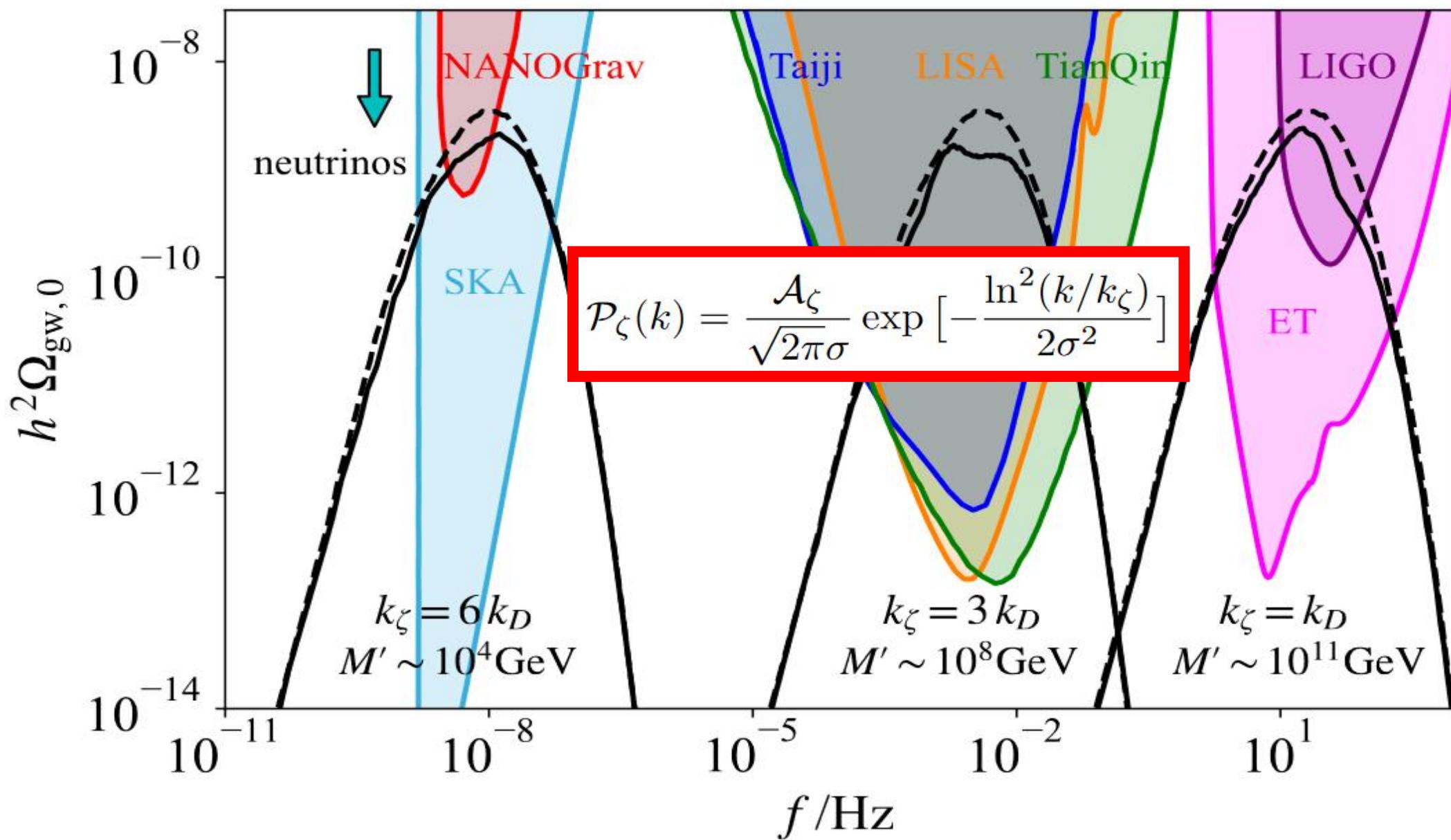
decoupling time of
weakly-interacting
particles



particle
interaction

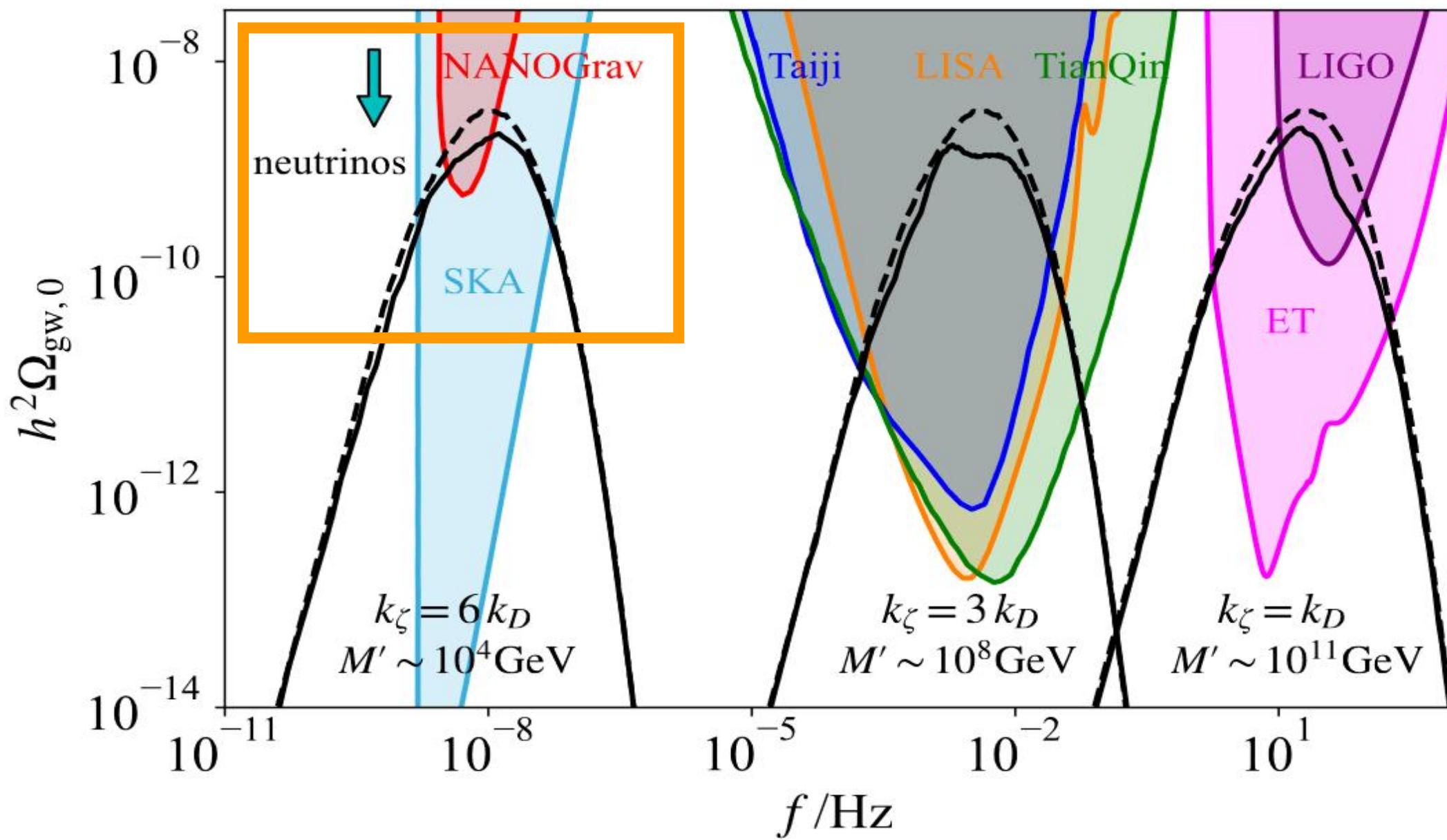
Observation

SCPMA 68 (2025) 210412: YHY and S. Wang



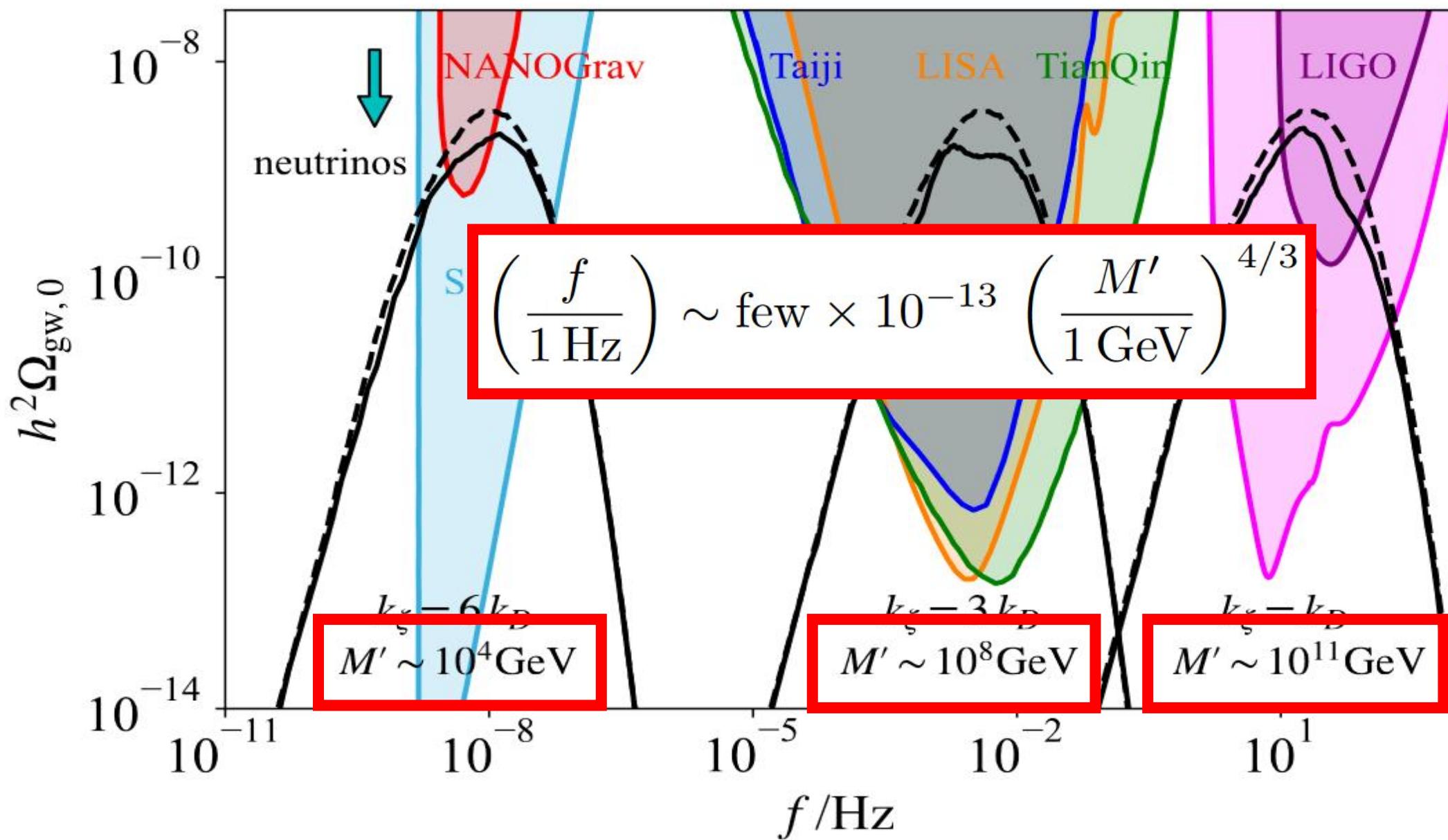
Observation (within the SM)

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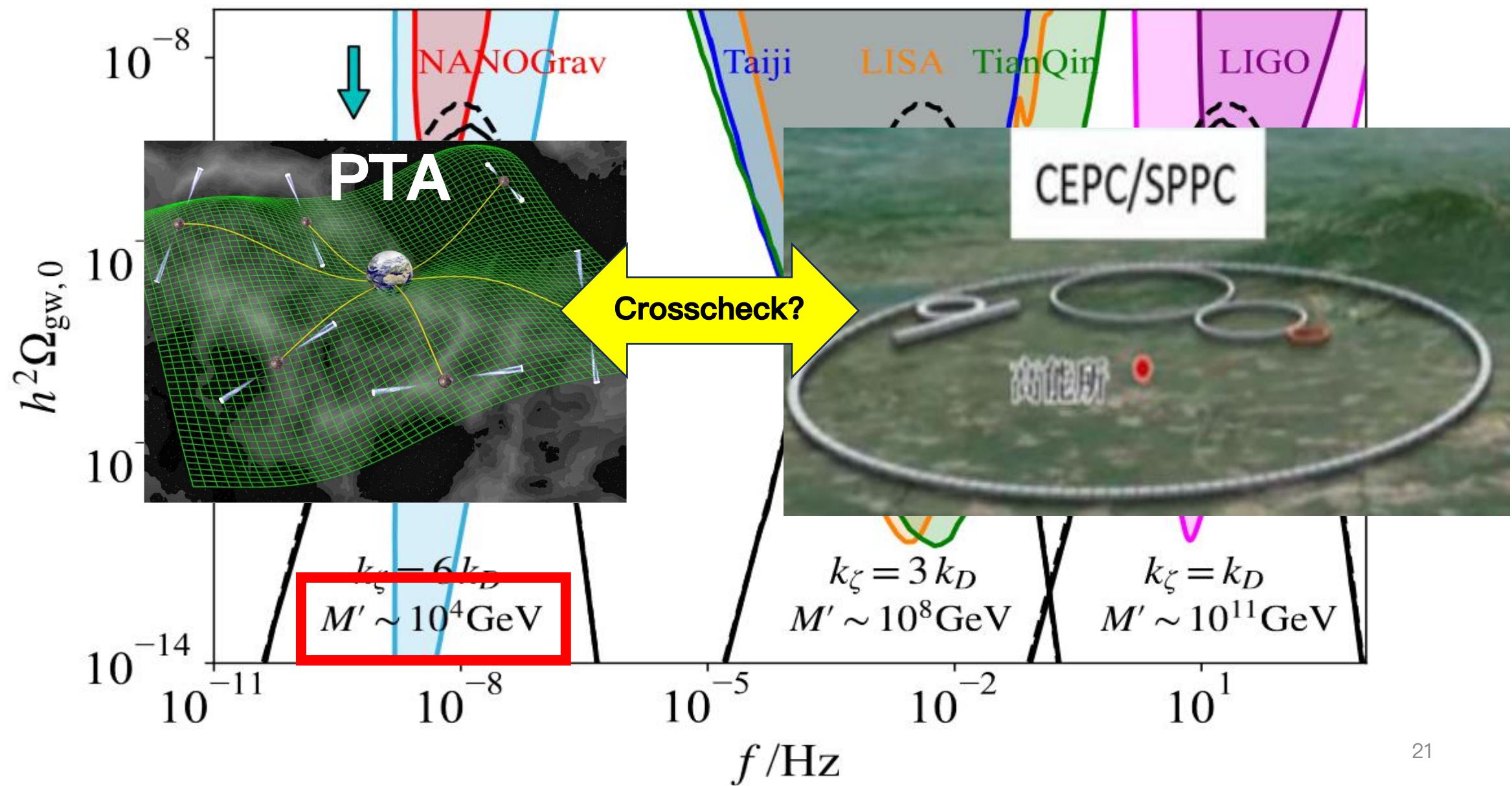
Observation (beyond the SM)

SCPMA 68 (2025) 210412: YHY and S. Wang



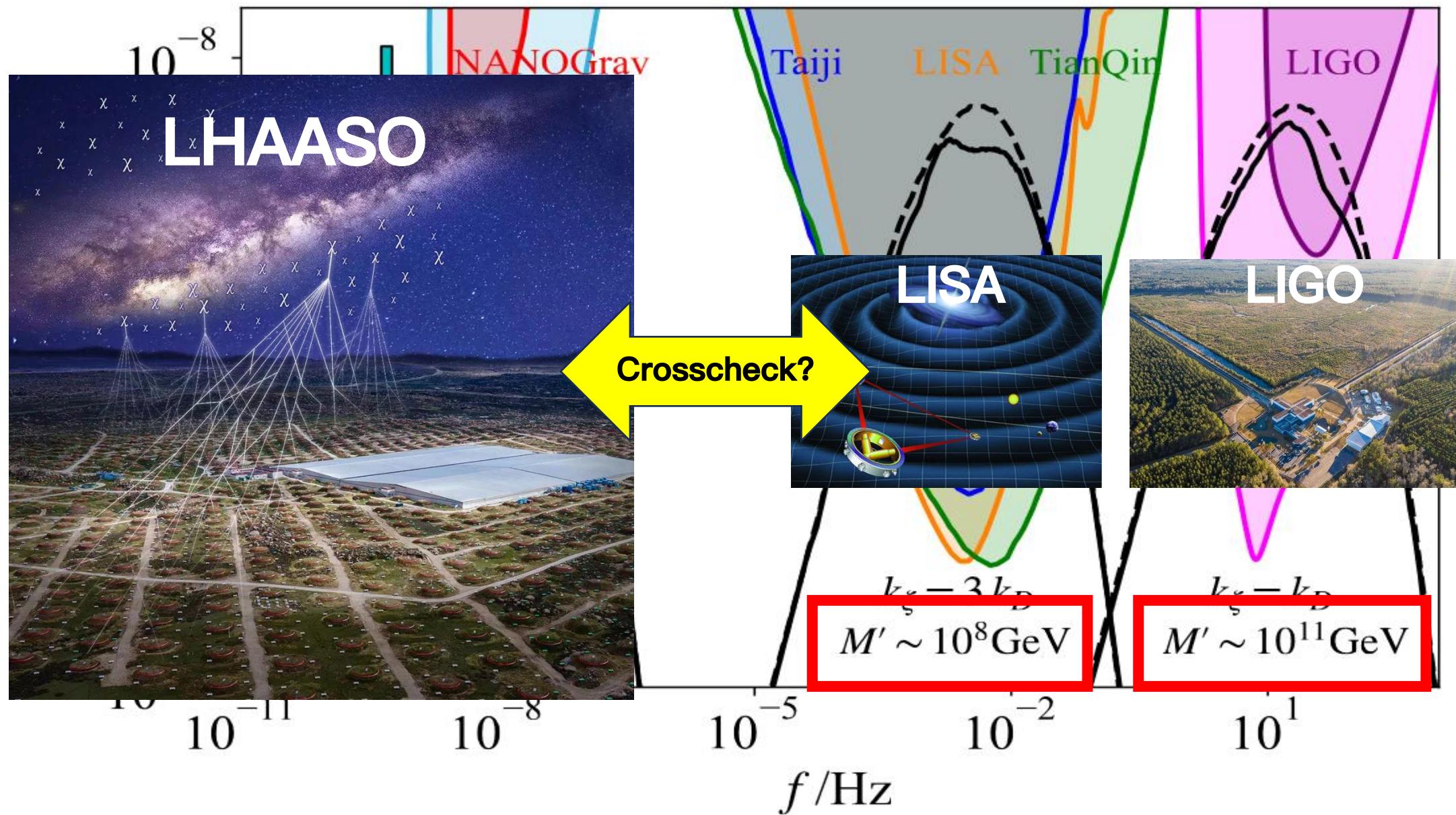
Future detection

SCPMA 68 (2025) 210412: YHY and S. Wang



Future detection

SCPMA 68 (2025) 210412: YHY and S. Wang



Conclusion

Particle interaction

STANDARD MODEL OF ELEMENTARY PARTICLES					
QUARKS	UP	CHARM	TOP	GLUON	HIGGS BOSON
	mass 2.3 MeV/c ² charge 2/3 spin 1/2 u	1.275 GeV/c ² 2/3 c	173.07 GeV/c ² 2/3 t	0 0 1 g	0 0 1 H
LEPTONS	DOWN	STRANGE	BOTTOM	PHOTON	Z BOSON
	4.8 MeV/c ² -1/3 1/2 d	95 MeV/c ² -1/3 1/2 s	4.18 GeV/c ² -1/3 1/2 b	0 0 1 γ	91.2 GeV/c ² 0 1 Z
GAUGE BOSONS	ELECTRON	MUON	TAU	W BOSON	W BOSON
	0.511 MeV/c ² -1 1/2 e	105.7 MeV/c ² -1 1/2 μ	1,777 GeV/c ² -1 1/2 τ	80.4 GeV/c ² 0 1 W	80.4 GeV/c ² ±1 1 W
new physics?					

Dissipation in cosmic fluid

shear viscosity bulk viscosity

$$\Delta T_{ij} = -\eta \left(\frac{\partial u_i}{\partial x^j} + \frac{\partial u_j}{\partial x^i} - \frac{2}{3} \delta_{ij} \nabla \cdot \mathbf{u} \right) - \xi \delta_{ij} \nabla \cdot \mathbf{u},$$

$$\Delta T_{0i} = -\chi \left(\frac{\partial T}{\partial x^i} + T \frac{\partial u_i}{\partial t} \right),$$

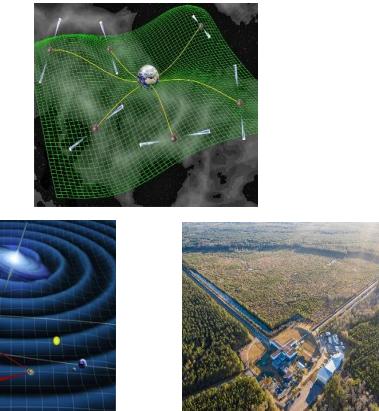
$$\Delta T_{00} = 0 \quad \text{thermal conduction}$$

We find a basic relation

$$\left(\frac{f}{1 \text{ Hz}} \right) \sim \text{few} \times 10^{-13} \left(\frac{M'}{1 \text{ GeV}} \right)^{4/3}$$



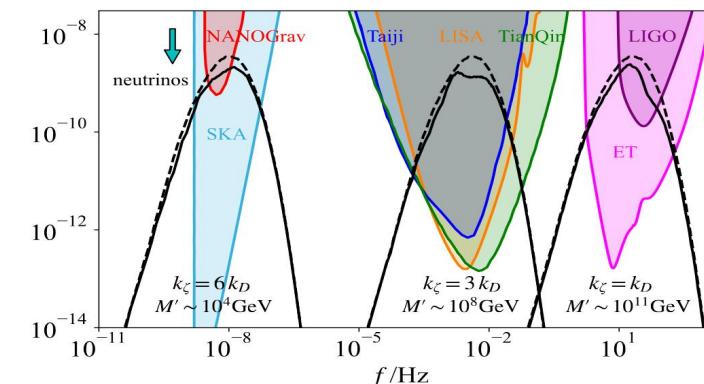
GW observation



Silk damping



Induced GWs



Conclusion

- Silk damping is an essential yet unexplored effect for GW observations.
- Silk damping notably suppresses the spectrum of induced GWs at the frequencies related to the decoupling of weakly-interacting particles.
- Within the SM, Silk damping caused by neutrinos will help to determine the origin of the GWB reported by PTA.
- Beyond the SM, Silk damping in induced GWs opens a new road to detect new physics, especially for those at extremely high energy scales.

Thanks!