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Effective Field Theories from GUT Scale to Nuclear Probes

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Based on CQS, J.H. Yu, arxiv:2504.XXXX

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Effective Field Theory

Our dream is to understand the most fundamental (UV) theory. On the other hand, we can always do computation without knowing the exact theory.

For example

- $v \ll c \rightarrow$ no need of Special Relativity
- $x \gg \lambda \sim h/p \rightarrow$ no need of Quantum Mechanics
- $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, h_{\mu\nu} \ll 1 \rightarrow$ no need of General Relativity

Consider a QFT with a fundamental high energy scale M but interested in observables as energy $E \ll M$.

Choose a cutoff $\Lambda < M$ and divide all quantum fields

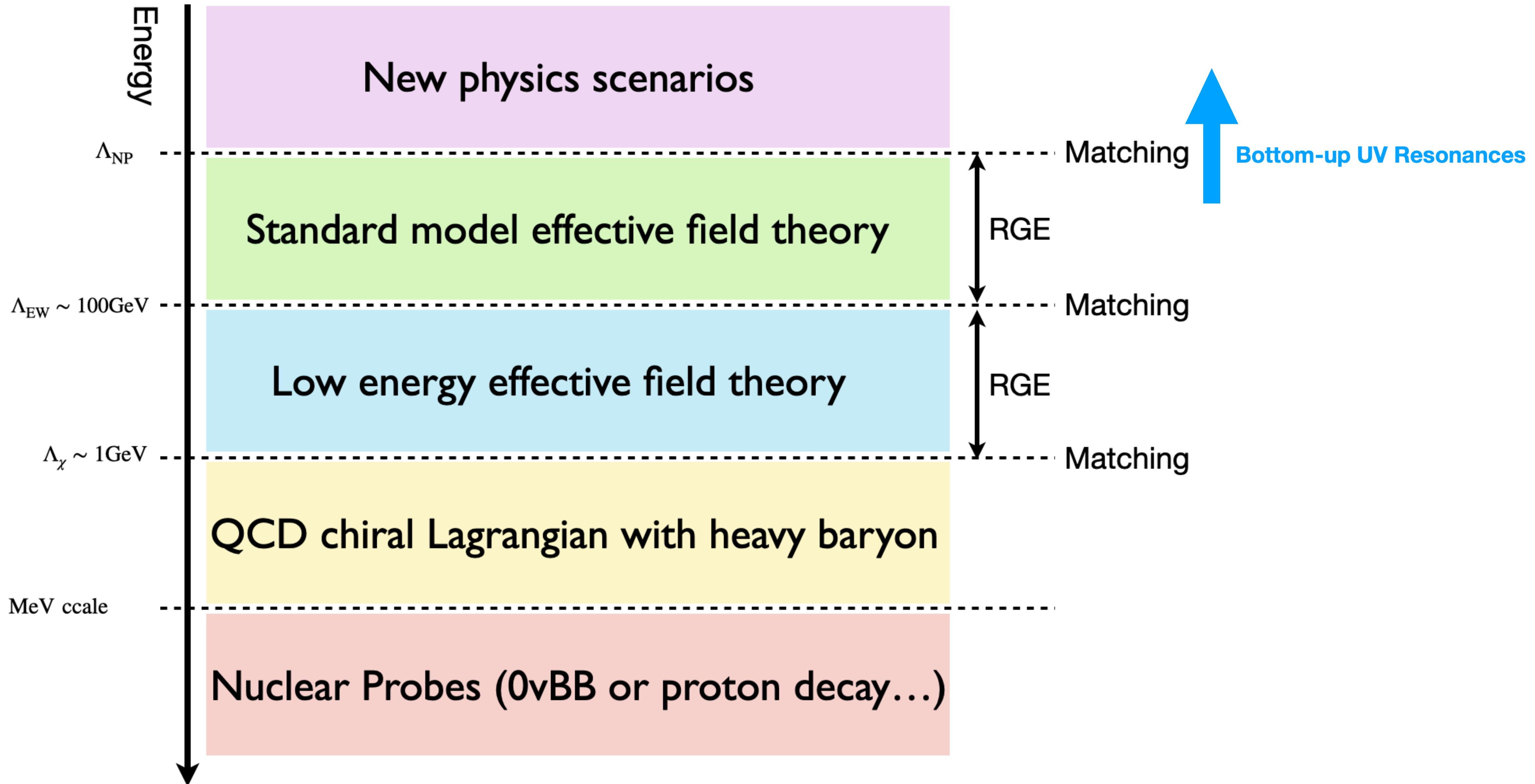
$$\varphi = \varphi_H + \varphi_L \quad (\varphi_H : \omega > \Lambda \quad \varphi_L : \omega < \Lambda)$$

Integrate-out

$$\int \mathcal{D}\varphi_H e^{i \int \mathcal{L}(\varphi_L, \varphi_H)} = e^{i \int \mathcal{L}_{eff}(\varphi_L)}$$

$$\mathcal{L}_{eff}(\varphi_L) = \mathcal{L}_{d \leq 4} + \sum_i \frac{O_i}{\Lambda^{dim(O_i)-4}}$$

Effective Field Theory



Grand Unified Theory

$\mathcal{G}_{\text{GUT}} : SU(5) \text{ or } SO(10) \dots$

Symmetry breaking



Unify into a single force



$\mathcal{G}_{\text{SM}} = SU(3)_C \times SU(2)_L \times U(1)_Y$

Different breaking chains

$\mathcal{G}_{\text{PS}} = SU(4)_C \times SU(2)_L \times SU(2)_R$

$\mathcal{G}_{421} = SU(4)_C \times SU(2)_L \times U(1)_{B-L}$

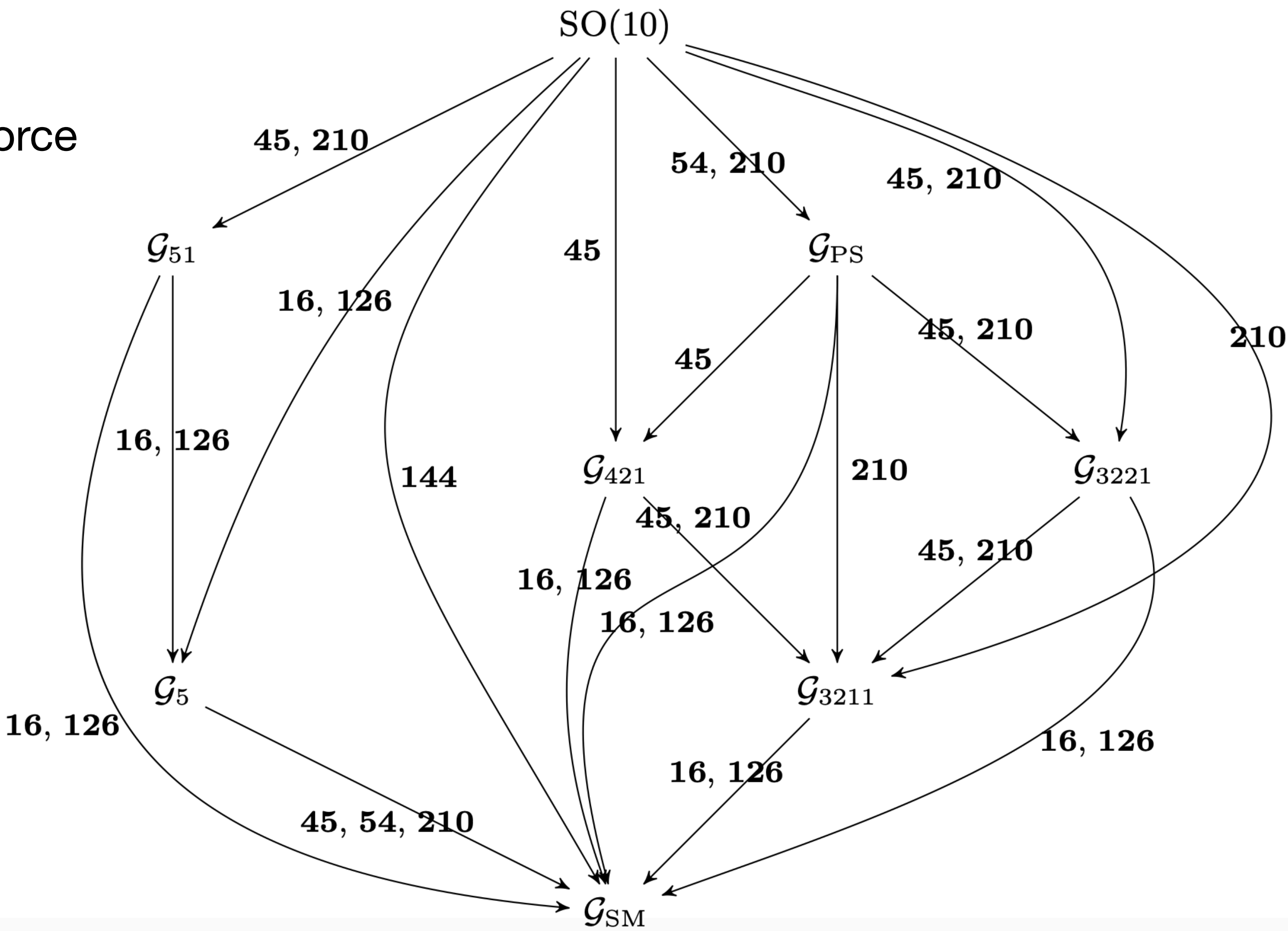
$\mathcal{G}_{3221} = SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$

$\mathcal{G}_{3211} = SU(3)_C \times SU(2)_L \times U(1)_R \times U(1)_{B-L}$

$\mathcal{G}_{51} = SU(5) \times U(1)_X$

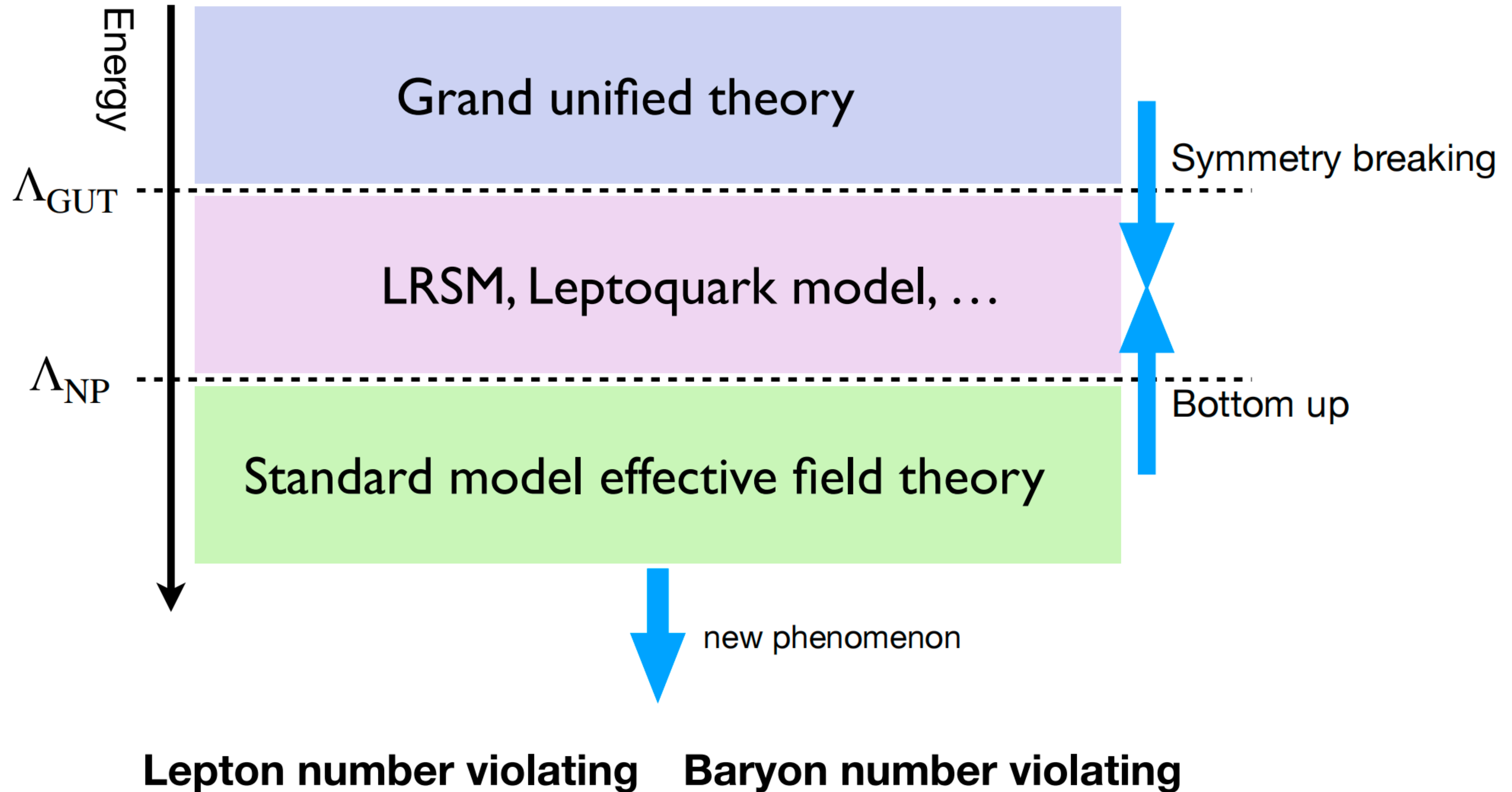
$\frac{1}{10}(-X + 4Y) = B - L,$

$Y = B - L + T_R^3,$



M.Pernow (2021)

SMEFT and GUT

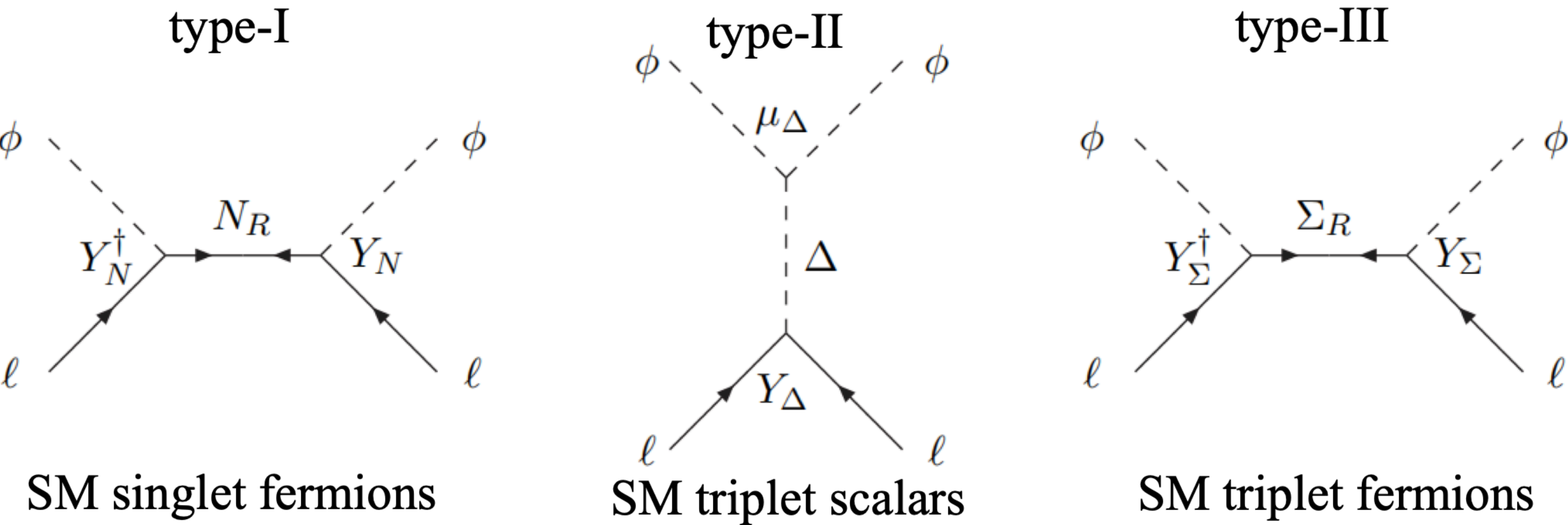


Lepton Number Violating

Dimension-5

$$Q_{\nu\nu} = \varepsilon_{jk}\varepsilon_{mn}\varphi^j\varphi^m(l_p^k)^TC l_r^n$$

Corresponding three types seesaw



The breaking chain in SO(10)

$SO(10)$	\rightarrow	\mathcal{G}_{PS}	\rightarrow	\mathcal{G}_{SM}
$\mathbf{45}_F$	\rightarrow	$(\mathbf{1}, \mathbf{1}, \mathbf{3})$	\rightarrow	$F_1(\mathbf{1}, \mathbf{1})_0$
$\mathbf{45}_F$	\rightarrow	$(\mathbf{1}, \mathbf{3}, \mathbf{1})$	\rightarrow	$F_5(\mathbf{1}, \mathbf{3})_0$
$\mathbf{54}_S$	\rightarrow	$(\mathbf{1}, \mathbf{3}, \mathbf{3})$	\rightarrow	$S_6(\mathbf{1}, \mathbf{3})_1$

$SO(10)$	\rightarrow	\mathcal{G}_{51}	\rightarrow	\mathcal{G}_{SM}
$\mathbf{45}_F$	\rightarrow	$\mathbf{24}_0$	\rightarrow	$F_1(\mathbf{1}, \mathbf{1})_0$
$\mathbf{45}_F$	\rightarrow	$\mathbf{24}_0$	\rightarrow	$F_5(\mathbf{1}, \mathbf{3})_0$
$\mathbf{54}_S$	\rightarrow	$\mathbf{15}_4$	\rightarrow	$S_6(\mathbf{1}, \mathbf{3})_1$

Lepton Number Violating

Dimension-7

Only L -violating			
Type: $\psi^2 H^4$		Type: $\psi^2 H^3 D$	
\mathcal{O}_{LH}	$\epsilon^{ik}\epsilon^{jl}(\ell_i^T C \ell_j) H_k H_l (H^\dagger H)$	\mathcal{O}_{LeHD}	$\epsilon^{ij}\epsilon^{kl}(\ell_i^T C \gamma^\mu e) H_j H_k (i D_\mu H_l)$
Type: $\psi^2 H^2 D^2$		Type: $\psi^2 H^2 X$	
\mathcal{O}_{LDH1}	$\epsilon^{ij}\epsilon^{kl}(\ell_i^T C D_\mu \ell_j)(H_k D^\mu H_l)$	\mathcal{O}_{LHW}	$\epsilon^{ik}(\epsilon \tau^I)^{jl}(\ell_i^T C i \sigma^{\mu\nu} \ell_j) H_k H_l W_{\mu\nu}^I$
\mathcal{O}_{LDH2}	$\epsilon^{ik}\epsilon^{jl}(\ell_i^T C D_\mu \ell_j)(H_k D^\mu H_l)$	\mathcal{O}_{LHB}	$\epsilon^{ik}\epsilon^{jl}(\ell_i^T C i \sigma^{\mu\nu} \ell_j) H_k H_l B_{\mu\nu}$
Type: $\psi^4 D$		Type: $\psi^4 H$	
\mathcal{O}_{duLLD}	$\epsilon^{ij}(\bar{d}^a \gamma^\mu u_a)(\ell_i^T C i D_\mu \ell_j)$	\mathcal{O}_{eLLLH} \mathcal{O}_{dLQLH1} \mathcal{O}_{dLQLH2} \mathcal{O}_{dLueH} \mathcal{O}_{QuLLH}	$\epsilon^{ij}\epsilon^{kl}(\bar{e} \ell_i)(\ell_j^T C \ell_k) H_l$ $\epsilon^{ij}\epsilon^{kl}(\bar{d}^a \ell_i)(q_{aj}^T C \ell_k) H_l$ $\epsilon^{ik}\epsilon^{jl}(\bar{d}^a \ell_i)(q_{aj}^T C \ell_k) H_l$ $\epsilon^{ij}(\bar{d}^a \ell_i)(u_a^T C e) H_j$ $\epsilon^{ij}(\bar{q}^{ak} u_a)(\ell_k^T C \ell_i) H_j$

$$\mathcal{L}_{eff} = \frac{1}{\Lambda} \mathcal{O}^{(5)} + \frac{1}{\Lambda^3} \mathcal{O}^{(7)} + \dots$$

The dimension-7 operators are suppressed, the UV models which contribute to dimension-7 operators should exclude three types seesaw.

Lepton Number Violating

The UV models contribute to dimension-7 operators without seesaw.

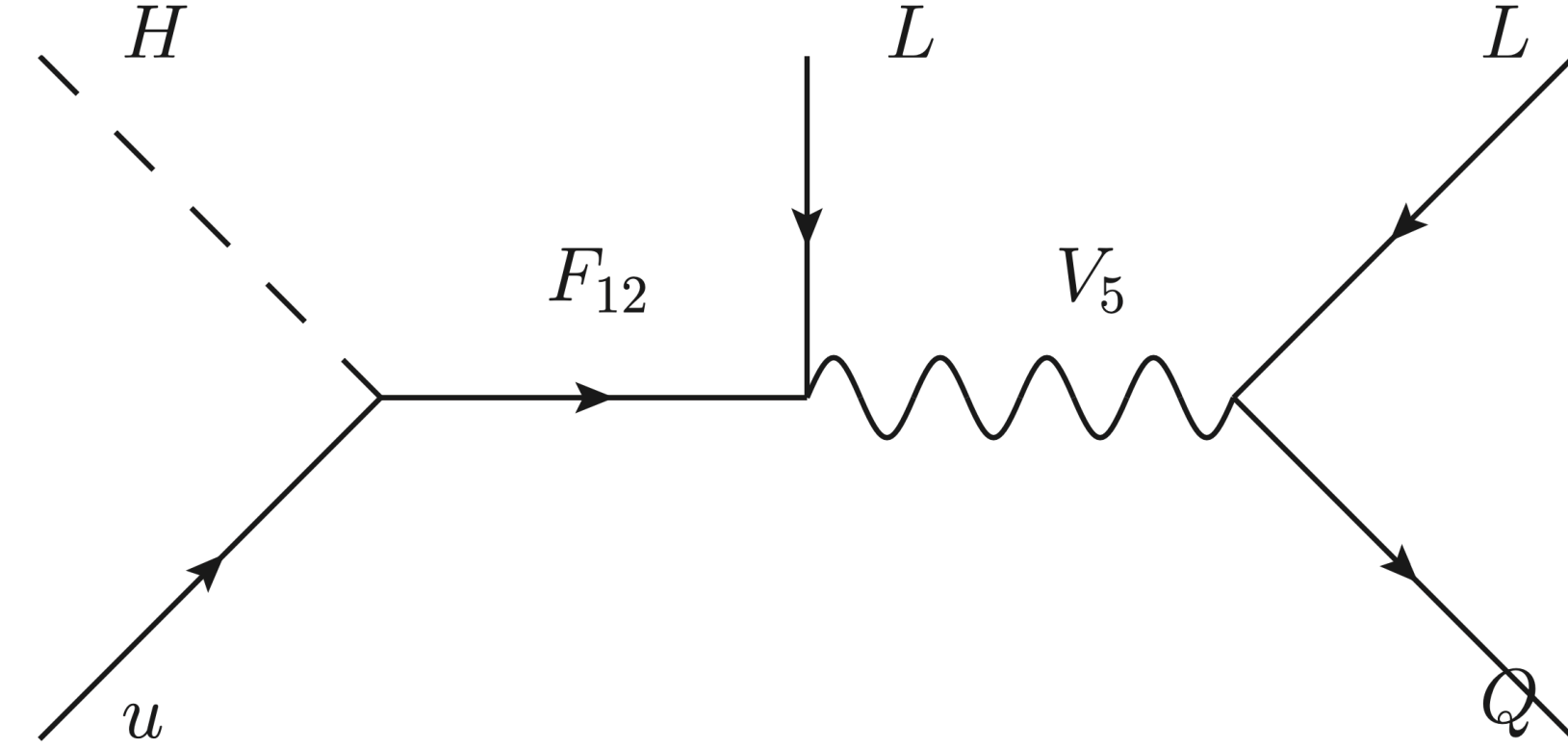
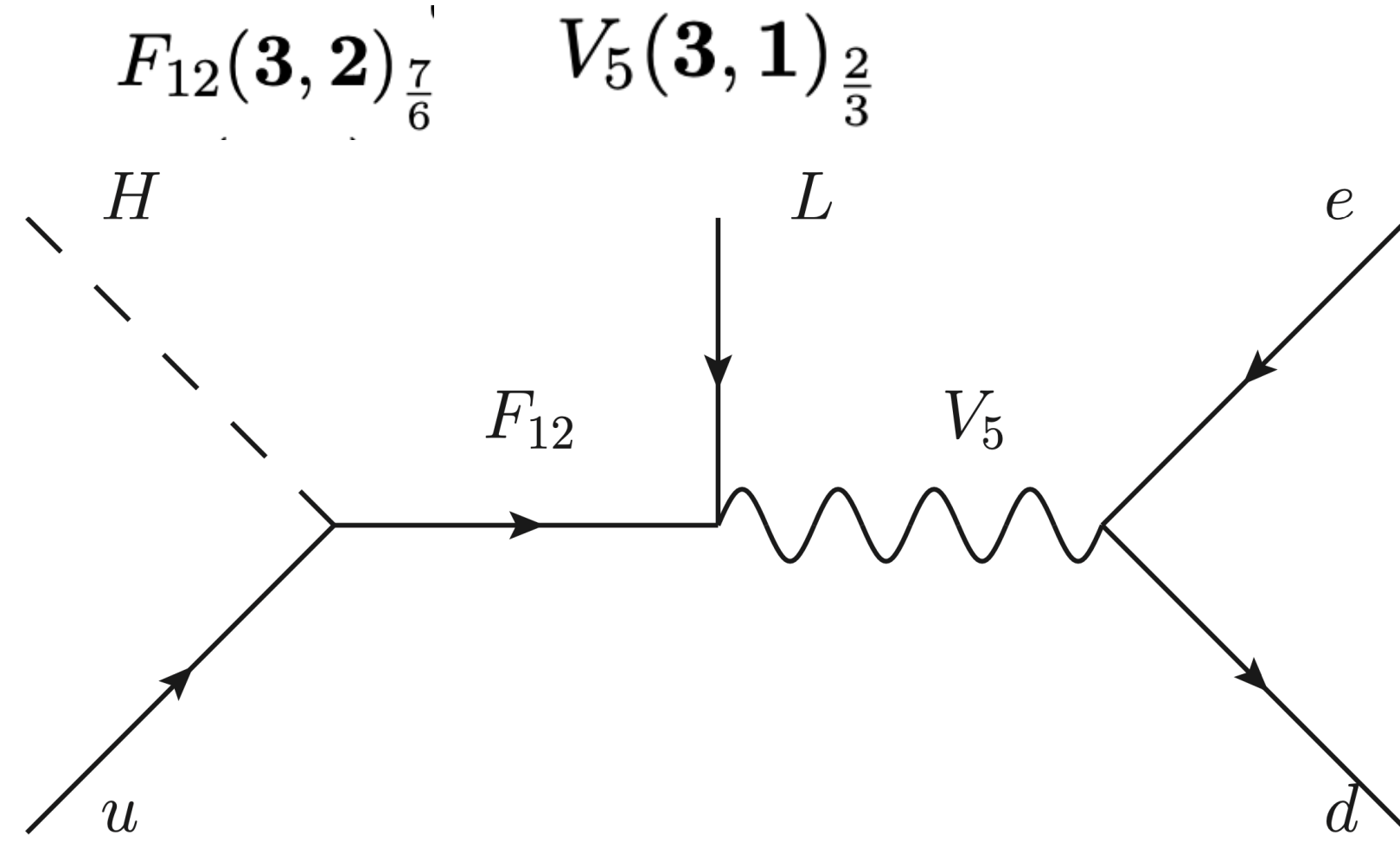
(S_8, F_6)	\rightarrow	\mathcal{O}_{LH}
$(F_3, V_2), (V_2, V_3)$	\rightarrow	\mathcal{O}_{LeHD}
$(S_2, S_4), (S_2, F_4)$	\rightarrow	\mathcal{O}_{eLLH}
(S_{12}, F_{14})	\rightarrow	\mathcal{O}_{dLQLH1}
$(S_2, S_4), (S_2, F_9),$ $(S_2, F_{10}), (S_{12}, F_9), (S_{12}, F_{14})$	\rightarrow	\mathcal{O}_{dLQLH2}
$(S_{12}, F_3), (S_{12}, F_{12}),$ $(F_3, V_2), (F_{10}, V_3),$ $(F_{12}, V_3), (F_{12}, V_5), (V_2, V_3)$	\rightarrow	\mathcal{O}_{dLueH}
$(S_2, S_4), (S_2, F_8),$ $(S_2, F_{12}), (F_{12}, V_5), (F_{12}, V_9)$	\rightarrow	\mathcal{O}_{QuLLH}

The breaking chain in SO(10).

$SO(10)$	\rightarrow	\mathcal{G}_{51}	\rightarrow	\mathcal{G}_{SM}	$SO(10)$	\rightarrow	\mathcal{G}_{PS}	\rightarrow	\mathcal{G}_{SM}
$\mathbf{10}_F$	\rightarrow	$\mathbf{5}_2$	\rightarrow	$F_3(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$	$\mathbf{10}_F$	\rightarrow	$(\mathbf{1}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_3(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$
$\mathbf{210}_F$	\rightarrow	$\mathbf{40}_{-4}$	\rightarrow	$F_4(\mathbf{1}, \mathbf{2})_{\frac{3}{2}}$	$\mathbf{210}_F$	\rightarrow	$(\overline{\mathbf{10}}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_4(\mathbf{1}, \mathbf{2})_{\frac{3}{2}}$
$\mathbf{54}_F$	\rightarrow	$\mathbf{15}_4$	\rightarrow	$F_6(\mathbf{1}, \mathbf{3})_1$	$\mathbf{54}_F$	\rightarrow	$(\mathbf{1}, \mathbf{3}, \mathbf{3})$	\rightarrow	$F_6(\mathbf{1}, \mathbf{3})_1$
$\mathbf{10}_F$	\rightarrow	$\mathbf{5}_2$	\rightarrow	$F_8(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$	$\mathbf{10}_F$	\rightarrow	$(\mathbf{6}, \mathbf{1}, \mathbf{1})$	\rightarrow	$F_8(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
$\mathbf{45}_F$	\rightarrow	$\overline{\mathbf{10}}_{-4}$	\rightarrow	$F_9(\mathbf{3}, \mathbf{1})_{\frac{2}{3}}$	$\mathbf{45}_F$	\rightarrow	$(\mathbf{15}, \mathbf{1}, \mathbf{1})$	\rightarrow	$F_9(\mathbf{3}, \mathbf{1})_{\frac{2}{3}}$
$\mathbf{45}_F$	\rightarrow	$\mathbf{24}_0$	\rightarrow	$F_{10}(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$	$\mathbf{45}_F$	\rightarrow	$(\mathbf{6}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_{10}(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$
$\overline{\mathbf{126}}_F$	\rightarrow	$\overline{\mathbf{45}}_{-2}$	\rightarrow	$F_{12}(\mathbf{3}, \mathbf{2})_{\frac{7}{6}}$	$\overline{\mathbf{126}}_F$	\rightarrow	$(\mathbf{15}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_{12}(\mathbf{3}, \mathbf{2})_{\frac{7}{6}}$
$\mathbf{120}_F$	\rightarrow	$\mathbf{45}_2$	\rightarrow	$F_{13}(\mathbf{3}, \mathbf{3})_{-\frac{1}{3}}$	$\mathbf{120}_F$	\rightarrow	$(\mathbf{6}, \mathbf{3}, \mathbf{1})$	\rightarrow	$F_{13}(\mathbf{3}, \mathbf{3})_{-\frac{1}{3}}$
$\mathbf{210}_F$	\rightarrow	$\mathbf{40}_{-4}$	\rightarrow	$F_{14}(\mathbf{3}, \mathbf{3})_{\frac{2}{3}}$	$\mathbf{210}_F$	\rightarrow	$(\mathbf{15}, \mathbf{3}, \mathbf{1})$	\rightarrow	$F_{14}(\mathbf{3}, \mathbf{3})_{\frac{2}{3}}$
$\mathbf{16}_S$	\rightarrow	$\mathbf{10}_{-1}$	\rightarrow	$S_2(\mathbf{1}, \mathbf{1})_1$	$\mathbf{16}_S$	\rightarrow	$(\overline{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	\rightarrow	$S_2(\mathbf{1}, \mathbf{1})_1$
$\mathbf{10}_S$	\rightarrow	$\mathbf{5}_2$	\rightarrow	$S_4(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$	$\mathbf{10}_S$	\rightarrow	$(\mathbf{1}, \mathbf{2}, \mathbf{2})$	\rightarrow	$S_4(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$
$\mathbf{16}_S$	\rightarrow	$\mathbf{10}_{-1}$	\rightarrow	$S_{12}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$	$\mathbf{16}_S$	\rightarrow	$(\mathbf{4}, \mathbf{2}, \mathbf{1})$	\rightarrow	$S_{12}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
$\mathbf{16}_V$	\rightarrow	$\mathbf{10}_{-1}$	\rightarrow	$V_2(\mathbf{1}, \mathbf{1})_1$	$\mathbf{16}_V$	\rightarrow	$(\overline{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	\rightarrow	$V_2(\mathbf{1}, \mathbf{1})_1$
$\mathbf{210}_V$	\rightarrow	$\mathbf{40}_{-4}$	\rightarrow	$V_3(\mathbf{1}, \mathbf{2})_{\frac{3}{2}}$	$\mathbf{210}_V$	\rightarrow	$(\overline{\mathbf{10}}, \mathbf{2}, \mathbf{2})$	\rightarrow	$V_3(\mathbf{1}, \mathbf{2})_{\frac{3}{2}}$
$\mathbf{45}_V$	\rightarrow	$\overline{\mathbf{10}}_{-4}$	\rightarrow	$V_5(\mathbf{3}, \mathbf{1})_{\frac{2}{3}}$	$\mathbf{45}_V$	\rightarrow	$(\mathbf{15}, \mathbf{1}, \mathbf{1})$	\rightarrow	$V_5(\mathbf{3}, \mathbf{1})_{\frac{2}{3}}$
$\mathbf{210}_V$	\rightarrow	$\mathbf{40}_{-4}$	\rightarrow	$V_9(\mathbf{3}, \mathbf{3})_{\frac{2}{3}}$	$\mathbf{210}_V$	\rightarrow	$(\mathbf{15}, \mathbf{3}, \mathbf{1})$	\rightarrow	$V_9(\mathbf{3}, \mathbf{3})_{\frac{2}{3}}$

Lepton Number Violating

For example



Interaction

$$\begin{aligned} \mathcal{L}_{F_{12}V_5} = & -2\mathcal{D}_{F_{12}LH^\dagger u^\dagger}^{pr} [(\bar{u}_r)^a (F_{12p})_{ai}] H^{\dagger i} \\ & -2\mathcal{D}_{d^\dagger e V_5}^{rsp} [(\bar{d}_r)^a \gamma_\mu e_s] (V_{5p})_a^\mu + 2\mathcal{D}_{L^\dagger Q V_5^\dagger}^{srp} [(\bar{l}_s)^i \gamma_\mu (q_r)_{ai}] (V_{5p}^\dagger)^{a\mu} \\ & + 2\mathcal{D}_{F_{12R}^\dagger L^\dagger V_5}^{psr} \epsilon_{ij} \delta_a^b [(\bar{l}_s)^j \gamma_\mu (F_{12p})^{ai}] (V_{5r})_b^\mu \end{aligned}$$

Wilson coefficient

$$\begin{aligned} \mathcal{C}_{dLueH}^{f_1 f_2 f_4 f_5} &= \sum_{p_1 p_2} \frac{16 \mathcal{D}_{d^\dagger e V_5}^{f_1 f_4 p_1} \mathcal{D}_{F_{12}LH^\dagger u^\dagger}^{p_2 * f_5} \mathcal{D}_{F_{12R}^\dagger L^\dagger V_5}^{p_2 f_2 p_1 *}}{M_{F_{12}} M_{V_5}^2} \\ \mathcal{C}_{QuLLH}^{f_1 f_2 f_4 f_5} &= \sum_{p_1 p_2} \frac{16 \mathcal{D}_{F_{12}LH^\dagger u^\dagger}^{p_1 * f_5} \mathcal{D}_{F_{12R}^\dagger L^\dagger V_5}^{p_1 f_2 p_2 *}}{M_{F_{12}} M_{V_5}^2} \mathcal{D}_{L^\dagger Q V_5^\dagger}^{f_1 f_4 p_2} \end{aligned}$$

Effective Lagrangian

$$\mathcal{L}_{eff} = \mathcal{C}_{dLueH}^{f_1 f_2 f_4 f_5} \times \epsilon^{ij} (\bar{d}_{f_1}^a \ell_{if_4}) (u_{af_5}^T C e_{f_2}) H_j + \mathcal{C}_{QuLLH}^{f_1 f_2 f_4 f_5} \times \epsilon^{ij} (\bar{q}_{f_4}^{ak} u_{af_5}) (\ell_{kf_1}^T C \ell_{if_2}) H_j$$

Baryon Number Violating

Dimension-6 $\Delta(B - L) = 0$

Type: <i>B</i> -violating	
\mathcal{O}_{duq}	$\epsilon^{abc}\epsilon^{jk}(d_a^T C u_b)(q_{cj}^T C \ell_k)$
\mathcal{O}_{qqu}	$\epsilon^{abc}\epsilon_{jk}(q_{aj}^T C q_{bk})(u_c^T C e)$
\mathcal{O}_{qqq}	$\epsilon^{abc}\epsilon_{jn}\epsilon_{km}(q_{aj}^T C q_{bk})(q_{cn}^T C \ell_m)$
\mathcal{O}_{duu}	$\epsilon^{abc}(d_a^T C u_b)(u_c^T C e)$

The breaking chain in SO(10).

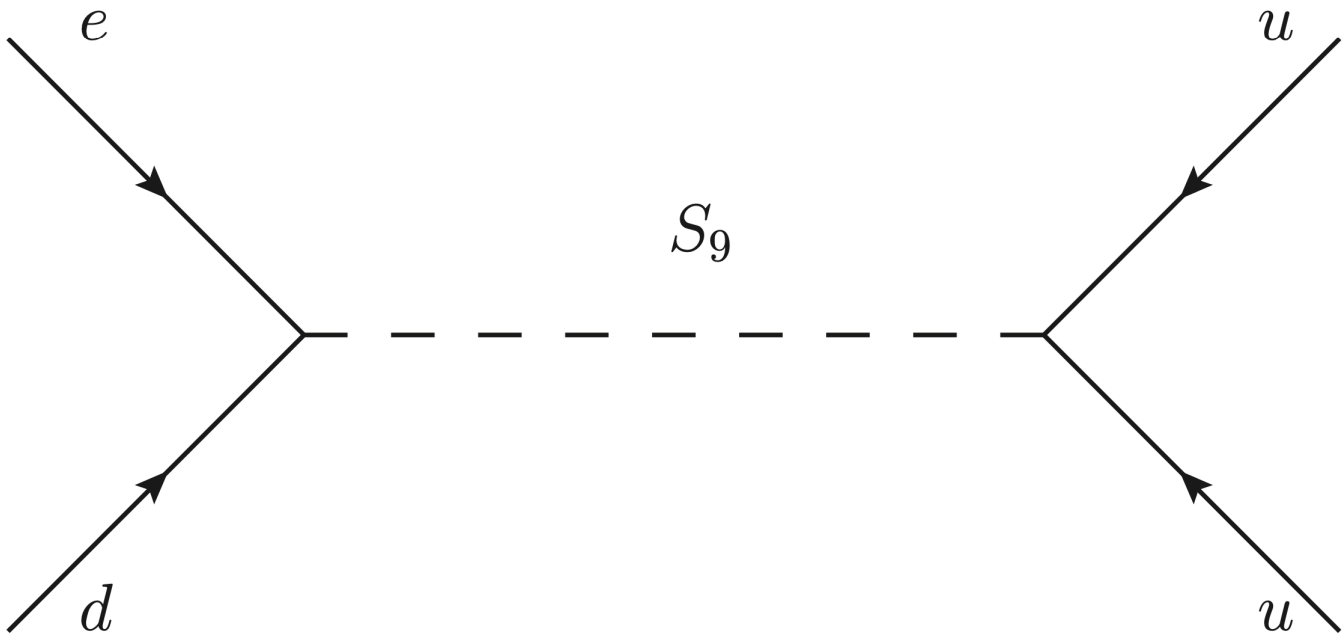
$SO(10)$	\rightarrow	\mathcal{G}_{PS}	\rightarrow	\mathcal{G}_{SM}
$\mathbf{45}_V$	\rightarrow	$(\mathbf{6}, \mathbf{2}, \mathbf{2})$	\rightarrow	$V_7(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$
$\mathbf{16}_V$	\rightarrow	$(\mathbf{4}, \mathbf{2}, \mathbf{1})$	\rightarrow	$V_8(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
$\mathbf{120}_S$	\rightarrow	$(\mathbf{6}, \mathbf{1}, \mathbf{3})$	\rightarrow	$S_9(\mathbf{3}, \mathbf{1})_{-\frac{4}{3}}$
$\mathbf{10}_S$	\rightarrow	$(\mathbf{6}, \mathbf{1}, \mathbf{1})$	\rightarrow	$S_{10}(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
$\mathbf{120}_S$	\rightarrow	$(\mathbf{6}, \mathbf{3}, \mathbf{1})$	\rightarrow	$S_{14}(\mathbf{3}, \mathbf{3})_{-\frac{1}{3}}$

$SO(10)$	\rightarrow	\mathcal{G}_{51}	\rightarrow	\mathcal{G}_{SM}
$\mathbf{45}_V$	\rightarrow	$\mathbf{24}_0$	\rightarrow	$V_7(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$
$\mathbf{16}_V$	\rightarrow	$\mathbf{10}_{-1}$	\rightarrow	$V_8(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
$\mathbf{120}_S$	\rightarrow	$\overline{\mathbf{45}}_{-2}$	\rightarrow	$S_9(\mathbf{3}, \mathbf{1})_{-\frac{4}{3}}$
$\mathbf{10}_S$	\rightarrow	$\mathbf{5}_2$	\rightarrow	$S_{10}(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
$\mathbf{120}_S$	\rightarrow	$\mathbf{45}_2$	\rightarrow	$S_{14}(\mathbf{3}, \mathbf{3})_{-\frac{1}{3}}$

The UV models contribute to dimension-6 operators.

S_{10}, V_7, V_8	\rightarrow	\mathcal{O}_{duq}
S_{10}, V_7	\rightarrow	\mathcal{O}_{qqu}
S_{10}, S_{14}	\rightarrow	\mathcal{O}_{qqq}
S_9, S_{10}	\rightarrow	\mathcal{O}_{duu}

For example



$$\mathcal{L}_{S_9} = -2\mathcal{D}_{d^\dagger e^\dagger S_9}^{rsp} [(\bar{d}_r)^a C \bar{e}_s] (S_{9p})_a - 4\mathcal{D}_{S_9^\dagger u^\dagger u^\dagger}^{prs} \epsilon^{abc} [(\bar{u}_r)^b C (\bar{u}_s)^c] (S_{9p}^\dagger)^a,$$

Baryon Number Violating

Dimension-7 $\Delta(B - L) = 2$

<i>B</i> -violating			
Type: $\psi^4 D$		Type: $\psi^4 H$	
\mathcal{O}_{LQddD}	$\epsilon^{abc}(\bar{\ell}^i \gamma^\mu q_{ai})(d_b^T C i D_\mu d_c)$	\mathcal{O}_{LdudH}	$\epsilon^{abc} \epsilon_{ij}(\bar{\ell}^i d_a)(u_b^T C d_c) H^{*j}$
\mathcal{O}_{edddD}	$\epsilon^{abc}(\bar{e} \gamma^\mu d_a)(d_b^T C i D_\mu d_c)$	\mathcal{O}_{LdddH}	$\epsilon^{abc}(\bar{\ell}^i d_a)(d_b^T C d_c) H_i$
		\mathcal{O}_{eQddH}	$-\epsilon^{abc}(\bar{e} Q_{ai})(d_b^T C d_c) H^{*i}$
		\mathcal{O}_{LdQQH}	$-\epsilon^{abc}(\bar{\ell}^k d_a)(q_{bk}^T C q_{ci}) H^{*i}$

The breaking chain in SO(10).

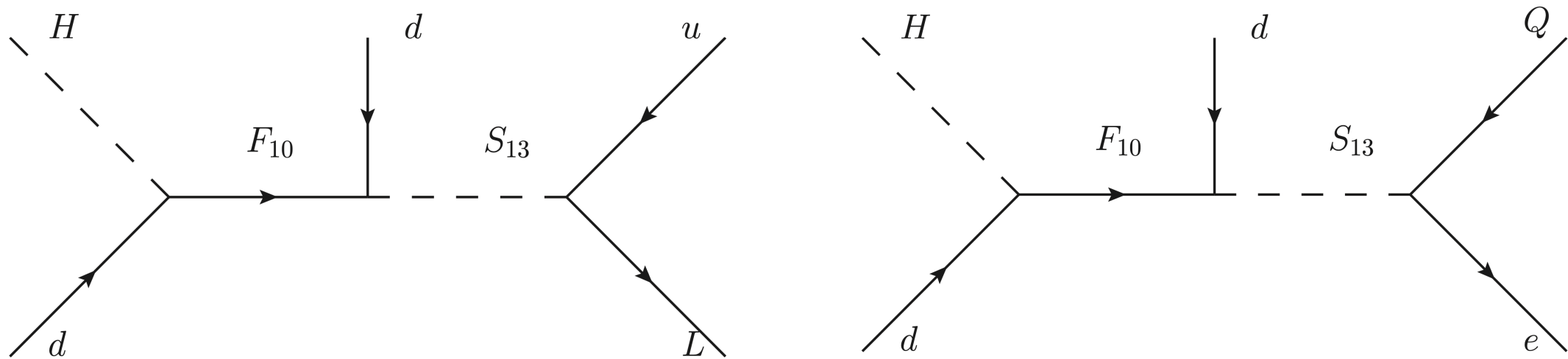
$SO(10)$	\rightarrow	\mathcal{G}_{51}	\rightarrow	\mathcal{G}_{SM}
45_F	\rightarrow	24_0	\rightarrow	$F_1(\mathbf{1}, \mathbf{1})_0$
16_F	\rightarrow	$\overline{10}_{-1}$	\rightarrow	$F_2(\mathbf{1}, \mathbf{1})_1$
10_F	\rightarrow	5_2	\rightarrow	$F_3(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$
10_F	\rightarrow	5_2	\rightarrow	$F_8(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
45_F	\rightarrow	24_0	\rightarrow	$F_{10}(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$
16_F	\rightarrow	10_{-1}	\rightarrow	$F_{11}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
16_S	\rightarrow	5_2	\rightarrow	$S_{11}(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
16_S	\rightarrow	10_{-1}	\rightarrow	$S_{12}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
$\overline{126}_S$	\rightarrow	$\overline{45}_{-2}$	\rightarrow	$S_{13}(\mathbf{3}, \mathbf{2})_{\frac{7}{6}}$

$SO(10)$	\rightarrow	\mathcal{G}_{PS}	\rightarrow	\mathcal{G}_{SM}
45_F	\rightarrow	$(\mathbf{1}, \mathbf{1}, \mathbf{3})$	\rightarrow	$F_1(\mathbf{1}, \mathbf{1})_0$
16_F	\rightarrow	$(\overline{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	\rightarrow	$F_2(\mathbf{1}, \mathbf{1})_1$
10_F	\rightarrow	$(\mathbf{1}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_3(\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$
10_F	\rightarrow	$(\mathbf{6}, \mathbf{1}, \mathbf{1})$	\rightarrow	$F_8(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
45_F	\rightarrow	$(\mathbf{6}, \mathbf{2}, \mathbf{2})$	\rightarrow	$F_{10}(\mathbf{3}, \mathbf{2})_{-\frac{5}{6}}$
16_F	\rightarrow	$(\mathbf{4}, \mathbf{2}, \mathbf{1})$	\rightarrow	$F_{11}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
16_S	\rightarrow	$(\mathbf{6}, \mathbf{1}, \mathbf{1})$	\rightarrow	$S_{11}(\mathbf{3}, \mathbf{1})_{-\frac{1}{3}}$
16_S	\rightarrow	$(\mathbf{4}, \mathbf{2}, \mathbf{1})$	\rightarrow	$S_{12}(\mathbf{3}, \mathbf{2})_{\frac{1}{6}}$
$\overline{126}_S$	\rightarrow	$(\mathbf{15}, \mathbf{2}, \mathbf{2})$	\rightarrow	$S_{13}(\mathbf{3}, \mathbf{2})_{\frac{7}{6}}$

The UV models contribute to dimension-7 operators which do not contribute to dimension-6 operators.

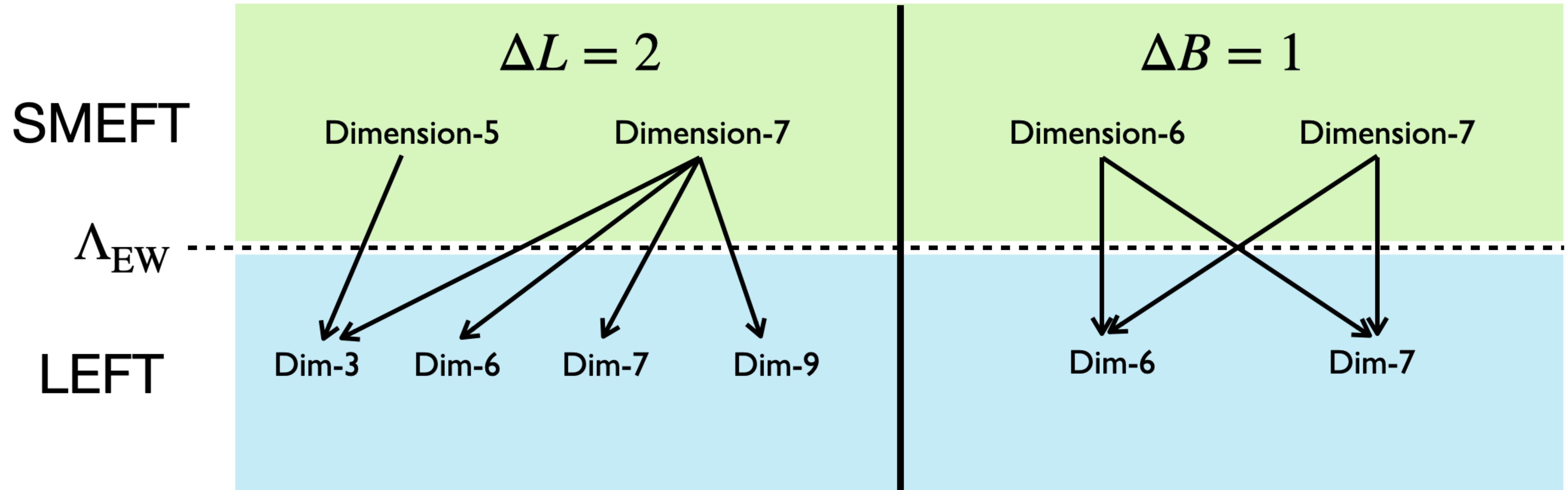
$(S_{11}, S_{12}), (S_{11}, F_2), (S_{11}, F_{11})$	\rightarrow	\mathcal{O}_{LdddH}
$(S_{11}, S_{13}), (S_{11}, F_1), (S_{11}, F_{11}), (S_{13}, F_{10})$	\rightarrow	\mathcal{O}_{LdudH}
$(S_{11}, S_{13}), (S_{11}, F_3), (S_{11}, F_8), (S_{13}, F_{10})$	\rightarrow	\mathcal{O}_{eQddH}

For example



$$\begin{aligned} \mathcal{L}_{S_{13}F_{10}} = & -2\mathcal{D}_{e^\dagger Q S_{13}^\dagger}^{rsp} [\bar{e}_r(q_s)_{ai}](S_{13p}^\dagger)^{ai} - 2\mathcal{D}_{LS_{13}u^\dagger}^{rps} \epsilon^{ji}[(\bar{u}_s)^a(l_r)_j](S_{13p})_{ai} \\ & - 2\mathcal{D}_{d^\dagger F_{10}LH}^{rp} \epsilon^{ij}[(\bar{d}_r)^a(F_{10p})_{ai}]H_j - 4\mathcal{D}_{d^\dagger F_{10R}S_{13}^\dagger}^{spr} \epsilon_{ij} \epsilon^{cab}[(\bar{d}_s)^c(F_{10p})^{ai}](S_{13r}^\dagger)^{bj} \end{aligned}$$

Matching between SMEFT and LEFT



The particles in LEFT : $(u_R, u_L, d_L, d_R, e_R, e_L, \nu_L)$, $n_u = 2, n_d = 3, n_e = 3, n_\nu = 3$.

The involving LEFT operator form : $(\bar{q}q\bar{l}l), (\bar{q}q\bar{l}lD), (\bar{q}q\bar{q}ql), (qqql), (qqq\bar{l}), (qqqlD), (qqq\bar{l}D)$, $q = (u, d, s)^T$

Matching between LEFT and ChPT

The QCD Lagrangian with external source

$$\mathcal{L} = \mathcal{L}_{\text{QCD}}^0 + \mathcal{L}_{\text{ext}} ,$$

$$\mathcal{L}_{\text{ext}} = \bar{q} \gamma^\mu (v_\mu + \gamma^5 a_\mu) q - \bar{q} (s - i \gamma^5 p) q + \bar{q} \sigma_{\mu\nu} \bar{t}^{\mu\nu} q ,$$

$$SU(3)_L \otimes SU(3)_R \quad \text{chiral symmetry}$$

symmetry breaking



The chiral Lagrangian

$$\mathcal{L}_2 = \frac{F_0^2}{4} \text{Tr} \left[D_\mu U (D^\mu U)^\dagger \right] + \frac{F_0^2}{4} \text{Tr} \left(\chi U^\dagger + U \chi^\dagger \right)$$

$$\mathcal{L}_{MB}^{(1)} = \langle \bar{B} (i \gamma^\mu D_\mu - M_0) B \rangle + \frac{D}{2} \langle \bar{B} \gamma_\mu \gamma_5 \{ u^\mu, B \} \rangle + \frac{F}{2} \langle \bar{B} \gamma_\mu \gamma_5 [u^\mu, B] \rangle .$$

$$SU(3)_V \quad \text{symmetry}$$

Dimension-6 and dimension-7 LEFT LNV operators

$$(\bar{q} \Gamma q)(l \Gamma l^c)$$

symmetry breaking



The chiral Lagrangian

$$(B + U + \dots)(l \Gamma l^c)$$

Dimension-9 LEFT LNV operators

$$(\bar{q} \Gamma q)(\bar{q} \Gamma q)(l \Gamma l^c)$$

$$SU(3)_V \quad \text{symmetry}$$

$$SU(3)_L \otimes SU(3)_R \quad \text{chiral symmetry}$$

$$U(x) = \exp \left(i \frac{\phi(x)}{F_0} \right), \quad \phi(x) = \begin{pmatrix} \pi^0 + \frac{1}{\sqrt{3}} \eta & \sqrt{2} \pi^+ & \sqrt{2} K^+ \\ \sqrt{2} \pi^- & -\pi^0 + \frac{1}{\sqrt{3}} \eta & \sqrt{2} K^0 \\ \sqrt{2} K^- & \sqrt{2} \bar{K}^0 & -\frac{2}{\sqrt{3}} \eta \end{pmatrix} \quad B = \begin{pmatrix} \frac{\Sigma^0}{\sqrt{2}} + \frac{\Lambda}{\sqrt{6}} & \Sigma^+ & p \\ \Sigma^- & -\frac{\Sigma^0}{\sqrt{2}} + \frac{\Lambda}{\sqrt{6}} & n \\ \Xi^- & \Xi^0 & -\frac{2\Lambda}{\sqrt{6}} \end{pmatrix}$$

Matching between LEFT and ChPT

Quark bilinear with spurion

$$(\bar{q}_L \Gamma \Sigma_L q_L), \quad (\bar{q}_R \Gamma \Sigma_R q_R), \quad (\bar{q}_R \Gamma \Sigma q_L), \quad (\bar{q}_L \Gamma \Sigma^\dagger q_R),$$

	$\bar{q}q$	$\bar{q}\gamma^5 q$	$\bar{q}\gamma^\mu q$	$\bar{q}\gamma^5 \gamma^\mu q$	$\bar{q}\sigma^{\mu\nu} q$	$\bar{q}\overleftrightarrow{\partial}^\mu q$	$\bar{q}\gamma^5 \overleftrightarrow{\partial}^\mu q$	$\bar{q}\gamma^\mu \overleftrightarrow{\partial}^\nu q$	$\bar{q}\gamma^5 \gamma^\mu \overleftrightarrow{\partial}^\nu q$
C	+	+	-	+	-	-	-	+	-
P	+	-	+	-	+	+	-	+	-

	$\bar{B}B$	$\bar{B}\gamma^5 B$	$\bar{B}\gamma^\mu B$	$\bar{B}\gamma^5 \gamma^\mu B$	$\bar{B}\sigma^{\mu\nu} B$	$\bar{B}\overleftrightarrow{D}^\mu B$
C	+	+	-	+	-	-
P	+	-	+	-	+	+
	$\bar{B}\gamma^5 \overleftrightarrow{D}^\mu B$	$\bar{B}\gamma^\mu \overleftrightarrow{D}^\nu B$	$\bar{B}\gamma^5 \gamma^\mu \overleftrightarrow{D}^\nu B$	u_μ	Σ_\pm	Q_\pm
C	-	+	-	+	+	\mp
P	-	+	-	-	\pm	\pm

$$\Sigma_\pm = u \left(\Sigma^\dagger \pm \Sigma \right) u \pm u^\dagger \left(\Sigma^\dagger \pm \Sigma \right) u^\dagger,$$

$$Q_\mp = u^\dagger \left(\Sigma_R \pm \Sigma_L \right) u \mp u \left(\Sigma_R \pm \Sigma_L \right) u^\dagger,$$

$$u_\mu = i(u^\dagger \partial_\mu u - u \partial_\mu u^\dagger),$$

$$\nabla_\mu X = \partial_\mu X + [\Gamma_\mu, X], \quad \Gamma_\mu = \frac{1}{2}(u^\dagger \partial_\mu u + u \partial_\mu u^\dagger).$$

$$SU(3)_L \otimes SU(3)_R$$

$$\begin{pmatrix} q_L \\ q_R \\ \bar{q}_L \\ \bar{q}_R \\ \Sigma \\ \Sigma^\dagger \\ \Sigma_L \\ \Sigma_R \end{pmatrix} \rightarrow \begin{pmatrix} Lq_L \\ Rq_R \\ \bar{q}_L L^\dagger \\ \bar{q}_R R^\dagger \\ R\Sigma L^\dagger \\ L\Sigma^\dagger R^\dagger \\ L\Sigma_L L^\dagger \\ R\Sigma_R R^\dagger \end{pmatrix}, \quad L \in SU(3)_L, R \in SU(3)_R.$$

$$u \rightarrow RuV^\dagger = VuL^\dagger \quad u = \sqrt{U}$$

$$\begin{pmatrix} u_\mu \\ \Sigma_- \\ \Sigma_+ \\ Q_- \\ Q_+ \\ B \end{pmatrix} \rightarrow \begin{pmatrix} Vu_\mu V^\dagger \\ V\Sigma_- V^\dagger \\ V\Sigma_+ V^\dagger \\ VQ_- V^\dagger \\ VQ_+ V^\dagger \\ VB V^\dagger \end{pmatrix}, \quad V \in SU(3)_V.$$

Matching between LEFT and ChPT

Quark bilinear with two spurion

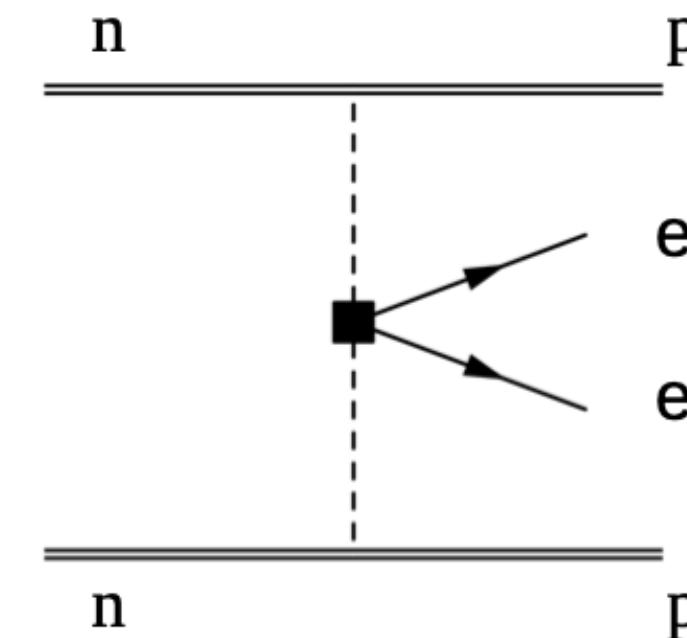
$$\begin{array}{ll}
 (\bar{q}_L \Gamma \Sigma_L q_L)(\bar{q}_L \Gamma \Sigma_L q_L), & (\bar{q}_R \Gamma \Sigma_R q_R)(\bar{q}_R \Gamma \Sigma_R q_R), \\
 (\bar{q}_L \Gamma \Sigma_L q_L)(\bar{q}_R \Gamma \Sigma_R q_R), & (\bar{q}_L \Gamma \Sigma^\dagger q_R)(\bar{q}_L \Gamma \Sigma q_R), \\
 (\bar{q}_L \Gamma \Sigma^\dagger q_R)(\bar{q}_R \Gamma \Sigma^\dagger q_L), & (\bar{q}_R \Gamma \Sigma q_L)(\bar{q}_R \Gamma \Sigma q_L), \\
 (\bar{q}_L \Gamma \Sigma_L q_L)(\bar{q}_R \Gamma \Sigma q_L), & (\bar{q}_L \Gamma \Sigma_L q_L)(\bar{q}_L \Gamma \Sigma^\dagger q_R), \\
 (\bar{q}_R \Gamma \Sigma_R q_R)(\bar{q}_R \Gamma \Sigma q_L), & (\bar{q}_R \Gamma \Sigma_R q_R)(\bar{q}_L \Gamma \Sigma^\dagger q_R).
 \end{array}
 \quad \longrightarrow \quad
 \begin{array}{ll}
 Q_{1+} = u^\dagger(\Sigma_{1R} - \Sigma_{1L})u + u(\Sigma_{1R} - \Sigma_{1L})u^\dagger, & \Sigma_{1+} = u^\dagger(\Sigma_1 + \Sigma_1^\dagger)u^\dagger + u(\Sigma_1 + \Sigma_1^\dagger)u, \\
 Q_{1-} = u^\dagger(\Sigma_{1R} + \Sigma_{1L})u - u(\Sigma_{1R} + \Sigma_{1L})u^\dagger, & \Sigma_{1-} = u^\dagger(\Sigma_1 - \Sigma_1^\dagger)u^\dagger - u(\Sigma_1 - \Sigma_1^\dagger)u, \\
 Q_{2+} = u^\dagger(\Sigma_{2R} - \Sigma_{2L})u + u(\Sigma_{2R} - \Sigma_{2L})u^\dagger, & \Sigma_{2+} = u^\dagger(\Sigma_2 + \Sigma_2^\dagger)u^\dagger + u(\Sigma_2 + \Sigma_2^\dagger)u, \\
 Q_{2-} = u^\dagger(\Sigma_{2R} + \Sigma_{2L})u - u(\Sigma_{2R} + \Sigma_{2L})u^\dagger. & \Sigma_{2-} = u^\dagger(\Sigma_2 - \Sigma_2^\dagger)u^\dagger - u(\Sigma_2 - \Sigma_2^\dagger)u.
 \end{array}$$

$0\nu\beta\beta$ operators

$$\mathcal{L}_{\Delta L=2}^{(9)} = \frac{1}{v^5} \sum_i \left[\left(C_{iR}^{(9)} \bar{e}_R C \bar{e}_R^T + C_{iL}^{(9)} \bar{e}_L C \bar{e}_L^T \right) O_i^{(9)} + C_i^{(9)} \bar{e} \gamma^\mu \gamma^5 C \bar{e}^T O_i^{\mu(9)} \right]$$

$$\begin{array}{ll}
 O_{1+}^{(9)} = \bar{q}_L^\alpha \gamma_\mu \tau^a q_L^\alpha \bar{q}_L^\beta \gamma^\mu \tau^a q_L^\beta + \bar{q}_R^\alpha \gamma_\mu \tau^a q_R^\alpha \bar{q}_R^\beta \gamma^\mu \tau^a q_R^\beta & O_{1+}^{(9)} \rightarrow 2\langle Q_{a-} Q_{a-} \rangle + 2\langle Q_{a+} Q_{a+} \rangle = 0, \\
 O_{2-}^{(9)} = \bar{q}_R^\alpha \gamma_\mu \tau^a q_R^\alpha \bar{q}_R^\beta \gamma^\mu \tau^a q_R^\beta - \bar{q}_L^\alpha \gamma_\mu \tau^a q_L^\alpha \bar{q}_L^\beta \gamma^\mu \tau^a q_L^\beta & O_{2-}^{(9)} \rightarrow -2\langle Q_{a-} Q_{a+} \rangle - 2\langle Q_{a+} Q_{a-} \rangle = 0, \\
 O_{7+}^{(9)} = \bar{q}_L^\alpha \gamma_\mu \tau^a q_L^\alpha \bar{q}_R^\beta \gamma^\mu \tau^a q_R^\beta & O_{7+}^{(9)}, O_{8+}^{(9)} \rightarrow \langle Q_{a-} Q_{a-} \rangle - \langle Q_{a+} Q_{a+} \rangle - \langle Q_{a-} Q_{a+} \rangle + \langle Q_{a+} Q_{a-} \rangle = 4\langle U^\dagger \tau^\pm U \tau^\pm \rangle, \\
 O_{8+}^{(9)} = \bar{q}_L^\alpha \gamma_\mu \tau^a q_L^\beta \bar{q}_R^\beta \gamma^\mu \tau^a q_R^\alpha
 \end{array}$$

The vertex $\pi\pi ee$ can be obtained for $0\nu\beta\beta$.



Matching between LEFT and ChPT

Dimension-6 and dimension-7 LEFT BNV operators

$$(q_L q_L q_L) \ell, (q_R q_R q_R) \ell, (q_L q_L q_R) \ell, (q_R q_R q_L) \ell$$

Three spurions can be added

$$(T_L q_L T_L q_L T_L q_L) \ell, (T_R q_R T_R q_R T_R q_R) \ell, (T_L q_L T_R q_R T_R q_R) \ell, (T_L q_L T_L q_L T_R q_R) \ell,$$

$$\begin{pmatrix} q_L \\ q_R \\ T_L \\ T_R \end{pmatrix} \rightarrow \begin{pmatrix} L q_L \\ R q_R \\ T_L L^\dagger \\ T_R R^\dagger \end{pmatrix}, \quad L \in SU(3)_L, R \in SU(3)_R.$$



$$\begin{aligned} (T_L u) &\rightarrow (T_L u) V^\dagger, \\ (T_R u^\dagger) &\rightarrow (T_R u^\dagger) V^\dagger, \\ B &\rightarrow V u V^\dagger, \\ u^\mu &\rightarrow V u^\mu V^\dagger, \end{aligned}$$

Type $BT_L T_L T_L$ LO

$$\begin{aligned} &\epsilon^{acd} (\lambda^A)_d^b (T_L u)_a (T_L u)_b (T_L u)_c (B^{AT} CL), \\ &\epsilon^{bcd} (\lambda^A)_d^a (T_L u)_a (T_L u)_b (T_L u)_c (B^{AT} CL), \end{aligned}$$

NLO

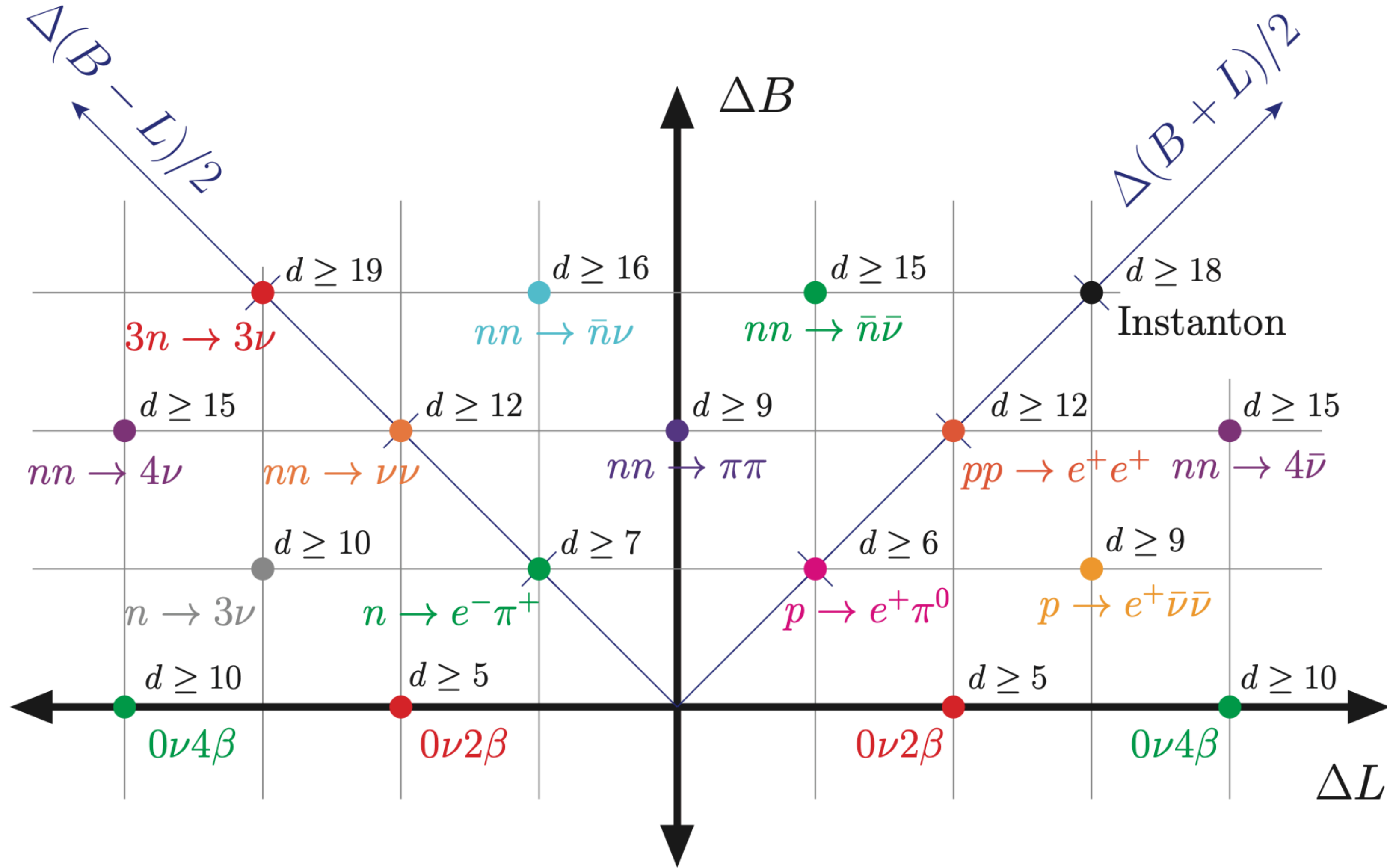
$$\begin{aligned} &\epsilon^{abc} (T_L u)_a (T_L u)_b (T_L u)_c (B^{AT} C \gamma_\mu L) u_\mu^A, \\ &f_{ABC} \epsilon^{bcd} (\lambda^A)_d^a (T_L u)_a (T_L u)_b (T_L u)_c (B^{BT} C \gamma_\mu L) u_\mu^C, \\ &d_{ABC} \epsilon^{bcd} (\lambda^A)_d^a (T_L u)_a (T_L u)_b (T_L u)_c (B^{BT} C \gamma_\mu L) u_\mu^C, \end{aligned}$$

The corresponding chiral Lagrangian for $(u_L u_L d_L L)$ can be expand as

$$\begin{aligned} &\epsilon^{acd} (\lambda^A)_d^b (T_L u)_a (T_L u)_b (T_L u)_c (B^{AT} CL) \\ &= \epsilon^{123} (\lambda^A)_3^b (T_L u)_1 (T_L u)_b (T_L u)_2 (B^{AT} CL) + \epsilon^{231} (\lambda^A)_1^b (T_L u)_2 (T_L u)_b (T_L u)_3 (B^{AT} CL) \\ &\quad + \epsilon^{312} (\lambda^A)_2^b (T_L u)_3 (T_L u)_b (T_L u)_1 (B^{AT} CL) + \epsilon^{132} (\lambda^A)_2^b (T_L u)_1 (T_L u)_b (T_L u)_3 (B^{AT} CL) \\ &\quad + \epsilon^{213} (\lambda^A)_3^b (T_L u)_2 (T_L u)_b (T_L u)_1 (B^{AT} CL) + \epsilon^{321} (\lambda^A)_1^b (T_L u)_3 (T_L u)_b (T_L u)_2 (B^{AT} CL) \\ &= p^T CL + (\pi^0 + \frac{1}{\sqrt{3}}\eta, \sqrt{2}\pi^+, \sqrt{2}K^+) \cdot (p, n, -\frac{2}{\sqrt{6}}\Lambda)^T CL + (\pi^0 + \frac{1}{\sqrt{3}}\eta) p^T CL + (-\pi^0 + \frac{1}{\sqrt{3}}\eta) p^T CL \\ &\quad - \sqrt{2}K^0 \Sigma^{+T} CL - \sqrt{2}K^- (\frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda)^T CL + \dots \\ &\quad + (\pi^0 + \frac{1}{\sqrt{3}}\eta)(-\pi^0 + \frac{1}{\sqrt{3}})(\pi^0 + \frac{1}{\sqrt{3}}\eta, \sqrt{2}\pi^+, \sqrt{2}K^+) \cdot (p, n, -\frac{2}{\sqrt{6}})^T \\ &\quad + 2\pi^+ K^0 (\pi^0 + \frac{1}{\sqrt{3}}\eta, \sqrt{2}\pi^+, \sqrt{2}K^+) (\frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda, \Sigma^-, \Xi^-)^T + \dots \end{aligned}$$

Contribute to the processes

$$\begin{array}{ll} p \rightarrow \pi^0 + \pi^0 + e^+, & p \rightarrow \pi^+ + \pi^- + e^+, \\ p \rightarrow \pi^0 + e^+, & p \rightarrow \eta + \pi^0 + e^+, \\ p \rightarrow \eta + e^+, & n \rightarrow \pi^0 + \pi^- + e^+, \\ n \rightarrow \pi^- + e^+, & n \rightarrow \eta + \pi^- + e^+, \end{array}$$



Heeck, Takhistov 1910.07647

Summary

$\Delta L = 2$

$\Delta B = 1$

$(\overline{126}_F, 45_V)$

Grand unified theory

$(45_F, \overline{126}_S)$

$[F_{12}(\mathbf{3}, 2)_{\frac{7}{6}}, V_5(\mathbf{3}, 1)_{\frac{2}{3}}]$

LRSM, Leptoquark model, ...

$[F_{10}(\mathbf{3}, 2)_{-\frac{5}{6}}, S_{13}(\mathbf{3}, 2)_{\frac{7}{6}}]$

$\epsilon^{ij}(\overline{d}_{f_1}^a \ell_{if_4})(u_{af_5}^T C e_{f_2}) H_j$
 $\epsilon^{ij}(\overline{q}_{f_4}^{ak} u_{af_5})(\ell_{kf_1}^T C \ell_{if_2}) H_j$

Standard model effective field theory

$-\epsilon^{abc}(\overline{e} Q_{ai})(d_b^T C d_c) H^{*i}$
 $-\epsilon^{abc}(\overline{\ell}^k d_a)(q_{bk}^T C q_{ci}) H^{*i}$

$\overline{u}_R \gamma^\mu d_R \overline{e}_{R,i} \gamma_\mu C \overline{\nu}_{L,j}^T$
 $\overline{u}_R d_L \overline{e}_{L,i} C \overline{\nu}_{L,j}^T$

Low energy effective field theory

$\epsilon_{\alpha\beta\gamma}(d_{Rp}^{\alpha T} C d_{Rr}^\beta)(\overline{e}_{Rs} d_{Lt}^\gamma)$
 $\epsilon_{\alpha\beta\gamma}(u_{Rp}^{\alpha T} C d_{Rr}^\beta)(\overline{\nu}_{Ls} d_{Rt}^\gamma)$

$\langle \overline{B} Q + \gamma^\mu B \rangle (\overline{e}_R \gamma_\mu C \overline{\nu}_L^T)$
 $\langle \overline{B} \Sigma + B \rangle (\overline{e}_L C \overline{\nu}_L^T)$

QCD chiral Lagrangian with heavy baryon

$\epsilon^{acd}(\lambda^A)_d^b (T_L u)_a (T_L u)_b (T_L u)_c (\overline{e}_R B_L^A),$
 $\epsilon^{bcd}(\lambda^A)_d^a (T_L u)_a (T_L u)_b (T_L u)_c (\overline{\nu}_L B_R),$

$n \rightarrow p + e + \nu^c$

Nuclear Probes (0vBB or proton decay...)

$p \rightarrow \pi^+ + \nu$

Thanks !