



# Recent results on the charmonium and X(6900) productions in pp collisions at the LHC energies from the PACIAE model

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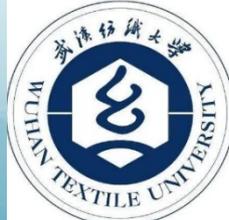
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# Outline

1. Background and motivations
2. The PACIAE and DCPC models
3. Results and discussions
4. Summary



# Background and Motivation

- $J/\psi$  is the lightest vector charmonium meson.
- The suppression of  $J/\psi$  production was proposed as a signal of QGP formation in HI collisions.
- The  $J/\psi$  production could also be suppressed due to the CNM effects, such as modifications of nuclear PDFs.
- In order to disentangle the hot and cold medium effects, it is necessary to understand the  $J/\psi$  production in pp collisions where the initial state effects are absent.



# Background and Motivation

- The  $J/\psi$  production was extensively investigated at colliders such as the Tevatron, RHIC and LHC.
- The ALICE collaboration had published the **inclusive  $J/\psi$**  production in the fwd- and mid-rapidity regions in pp collisions at  $\sqrt{s} = 2.76, 5.02, 7, 8$  and 13 TeV.

Inclusive  $J/\psi$  {  
  direct  $J/\psi$   
  feed-down from heavier charmonium decay  
  beauty hadrons weak decay

[Phys. Lett. B 718, 295 \(2012\)](#),  
[JHEP 10, 084 \(2019\)](#),  
[Phys. Lett. B 704, 442 \(2011\)](#),  
[Eur. Phys. J. C 74, 2974 \(2014\)](#),  
[Eur. Phys. J. C 76, 184 \(2016\)](#),  
[Eur. Phys. J. C 81, 1121 \(2021\)](#),  
[Eur. Phys. J. C 77, 392 \(2017\)](#).



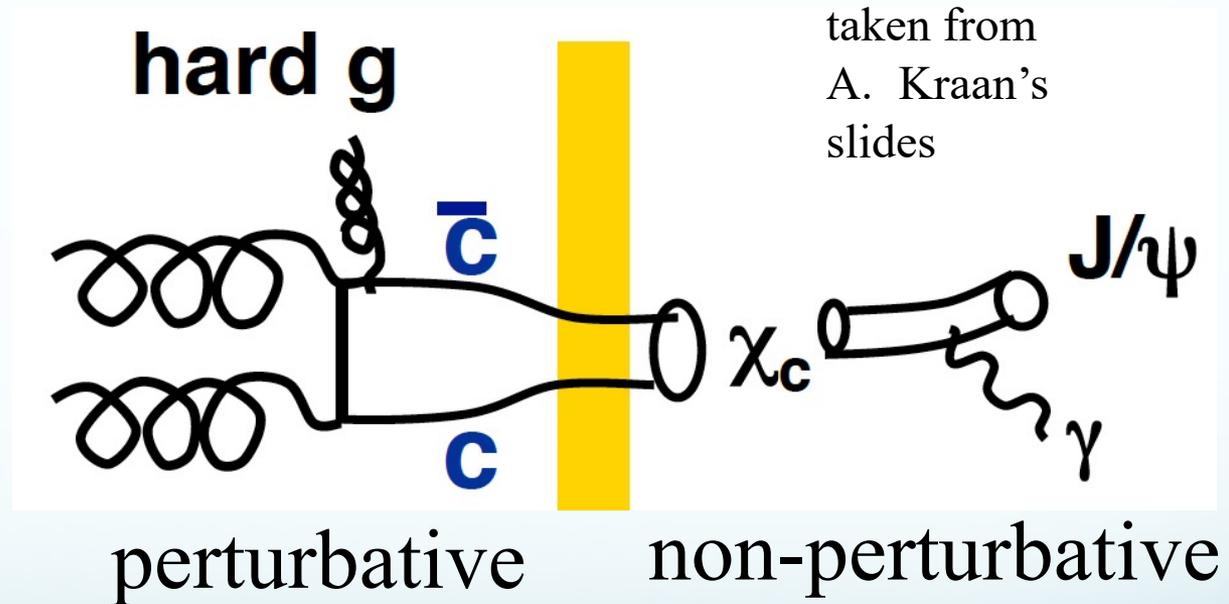
# Background and Motivation

- The ALICE collaboration had also published multiplicity dependence of  $J/\psi$  production in pp collisions at  $\sqrt{s} = 5.02, 7$  and 13 TeV. [Phys. Lett. B 712, 165–175 \(2012\)](#), [Phys. Lett. B 810, 135758 \(2020\)](#), [JHEP 06, 015 \(2022\)](#)
- Several theoretical approaches, such as CSM, COM and CEM have been utilized to describe the experimental data.
- They differ mostly in the treatment of non-perturbative evolution of the  $c\bar{c}$  pair into the bound state  $J/\psi$ .

# Background and Motivation

- Color Singlet Model [Phys. Lett. B 102, 364 \(1981\)](#).

Color  
Singlet

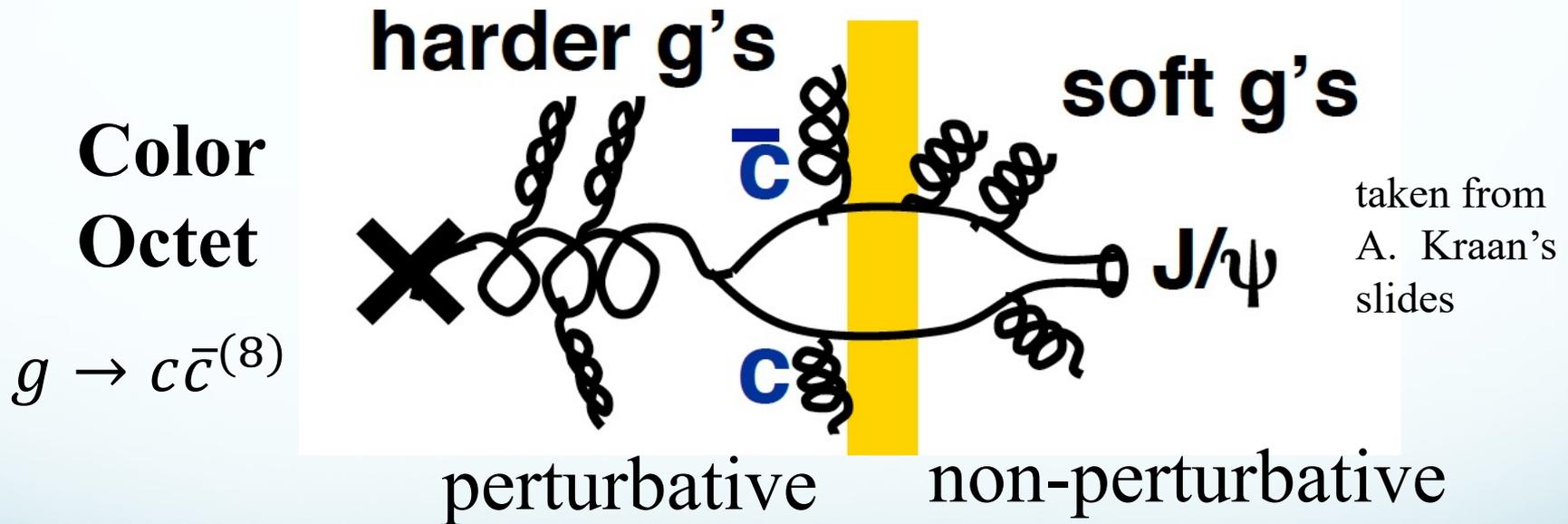


Color singlet  $\rightarrow$  No gluon radiation

$J/\psi$  produced in isolation

# Background and Motivation

- Color Octet Model [Phys. Rev. D 51, 1125 \(1995\)](#)



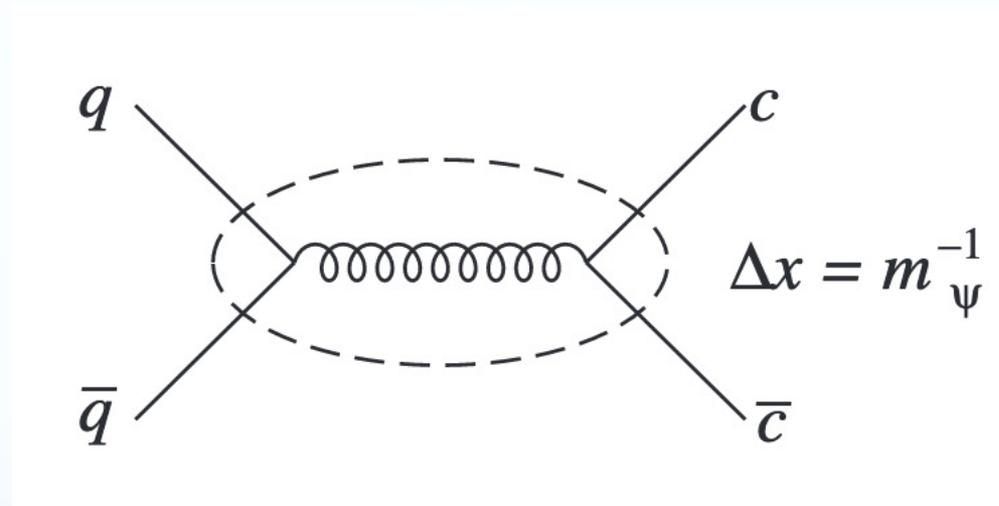
Color Octet  $\rightarrow$  shower expected

$J/\psi$  produced in shower

# Background and Motivation

- Color Evaporation Model [Phys. Lett. B 390, 323-328 \(1997\)](#)

## Color Evaporation

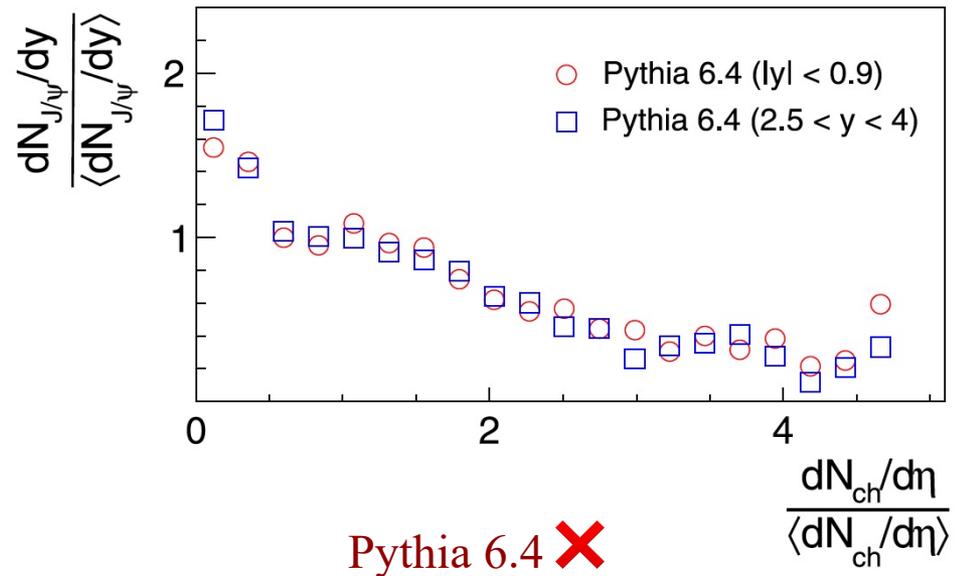
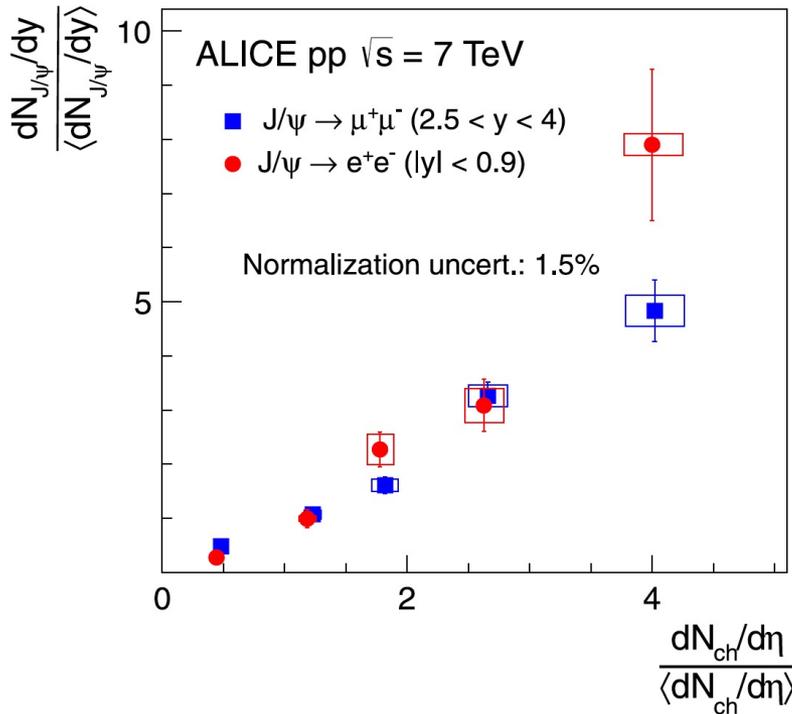


$q\bar{q}$  annihilation into  $c\bar{c}$ , similar to the Drell-Yan process

There is an infinite time for soft gluons to readjust the color of the  $c\bar{c}$  pair before it appears as an asymptotic  $J/\psi$ .

# Background and Motivation

- The  $J/\psi$  production was also investigated by Monte Carlo simulations [Phys. Lett. B 712 \(2012\) 165–175](#)

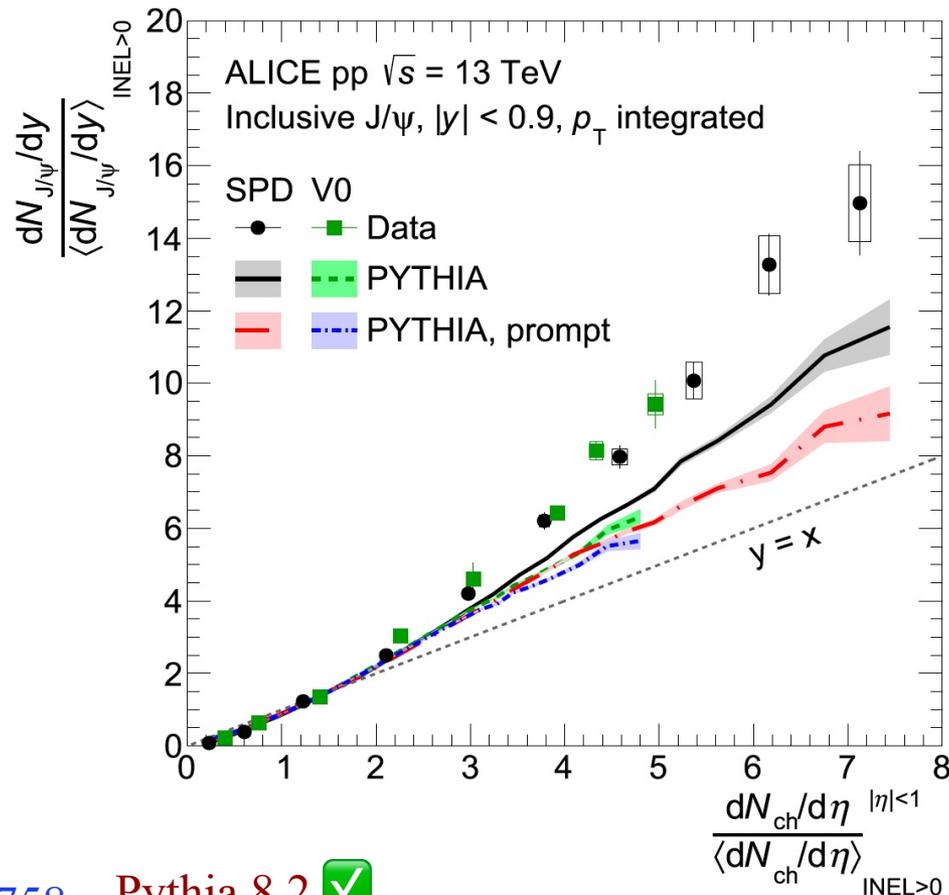


# Background and Motivation

- The  $J/\psi$  production was also investigated by Monte Carlo

simulations

Different treatment  
of MPI in PYTHIA 6  
and PYTHIA 8





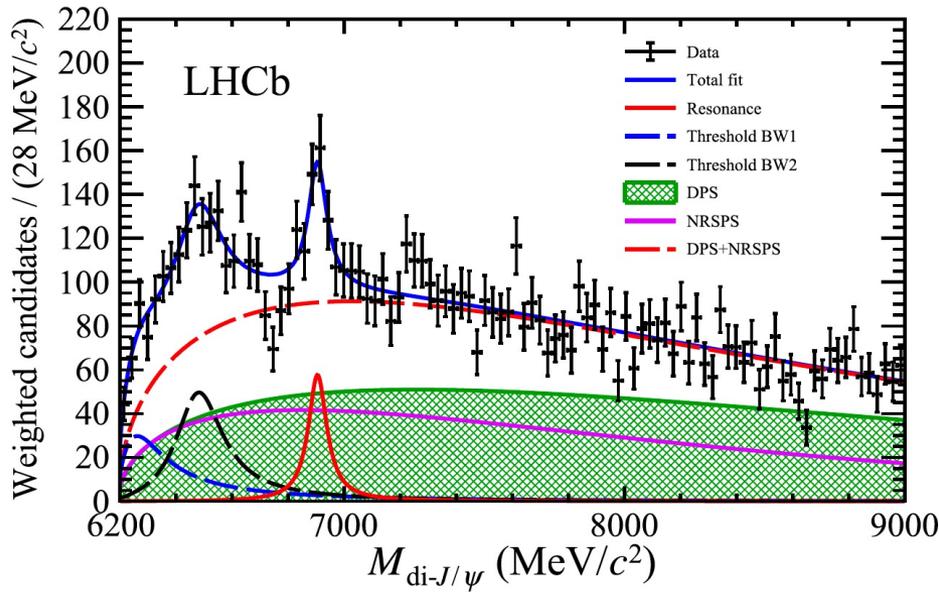
# Background and Motivation

- In pp collisions, a pair of  $J/\psi$  mesons can be produced in two separate interactions of gluons or quarks, named **double-parton scattering** (DPS) or in a **single-parton scattering** (SPS).
- The SPS process includes both resonant production via **intermediate states**, which could be **tetraquarks**, and non-resonant production (NRSPS).
- Within the DPS process, the two  $J/\psi$  mesons are usually assumed to be **produced independently**.
- The fully tetraquark states are searched for by the LHCb collaboration in the di-  $J/\psi$  invariant mass spectra in pp collisions at  $\sqrt{s} = 7, 8$  and 13 TeV.

# Background and Motivation

- Experimental results of X(6900)

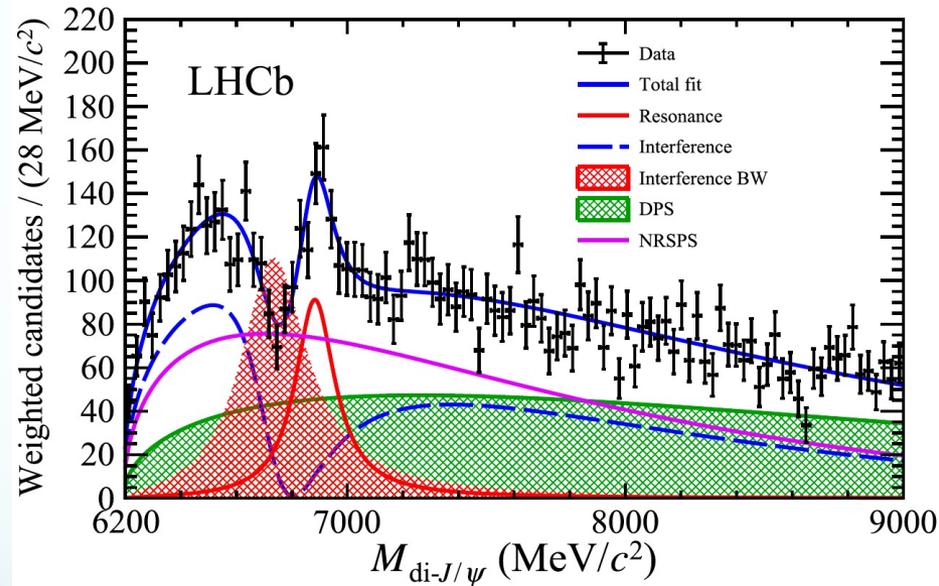
Science Bulletin 65, 1983 (2020)



$$m[X(6900)] = 6905 \pm 11 \pm 7 \text{ MeV}/c^2$$

$$\Gamma[X(6900)] = 80 \pm 19 \pm 33 \text{ MeV}$$

w/o considering interference



$$m[X(6900)] = 6886 \pm 11 \pm 11 \text{ MeV}/c^2$$

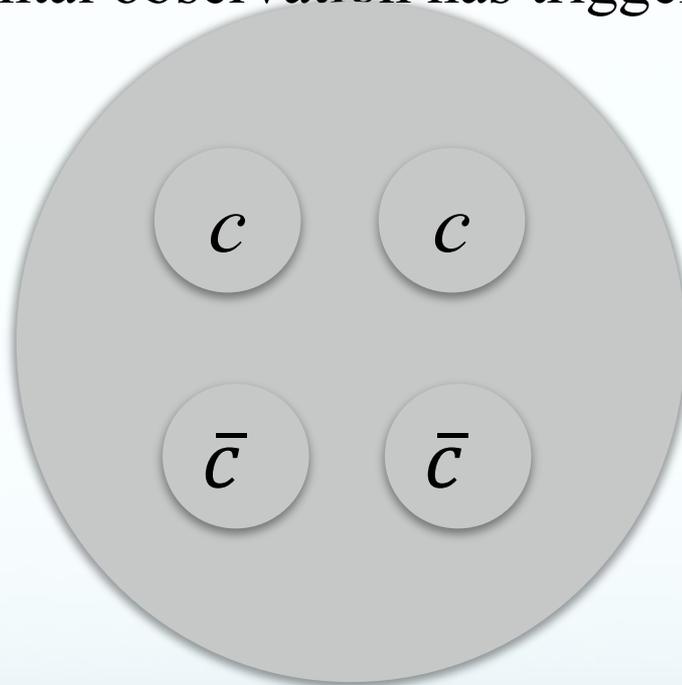
$$\Gamma[X(6900)] = 168 \pm 33 \pm 69 \text{ MeV}$$

w/ considering interference

# Background and Motivation

- The experimental observation has triggered tremendous studies.

Tetraquark  
 $c\bar{c}c\bar{c}$



X.-Z. Weng, X.-L. Chen, W.-Z. Deng, and S.-L. Zhu, Phys. Rev. D 103, 034001 (2021).

J.-Z. Wang, D.-Y. Chen, X. Liu, and T. Matsuki, Phys. Rev. D 103, L071503 (2021).

R.-H. Wu, Y.-S. Zuo, C.-Y. Wang, C. Meng, Y.-Q. Ma, and K.-T. Chao, JHEP 11, 023 (2022).

M.-S. Liu, F.-X. Liu, X.-H. Zhong, and Q. Zhao, Phys. Rev. D 109, 076017 (2024).

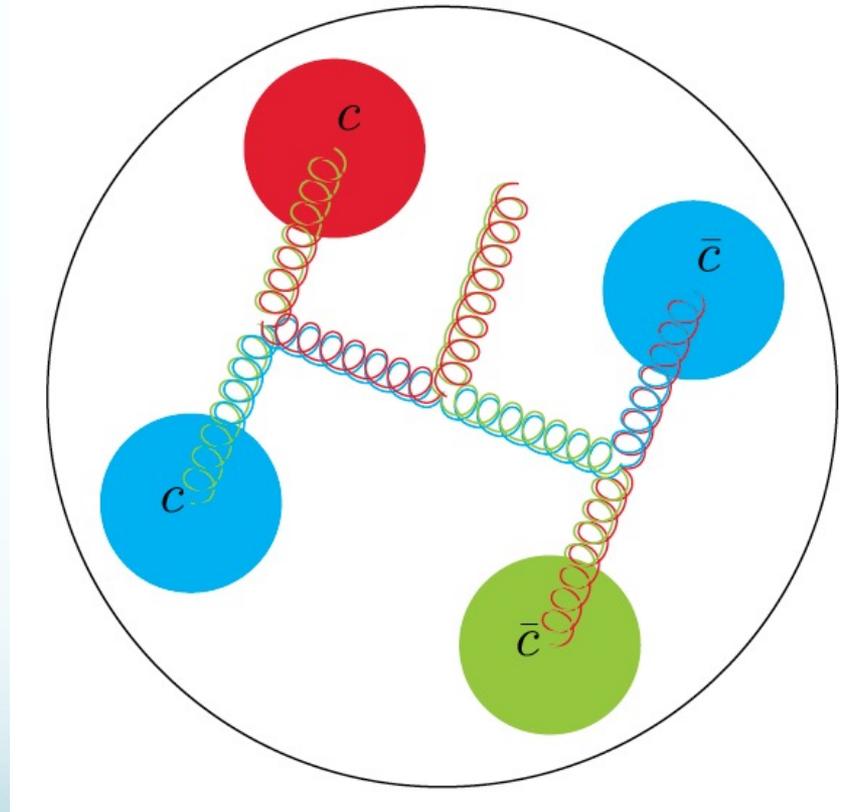
Z.-G. Wang, Nucl. Phys. B 985, 115983 (2022).

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# Background and Motivation

- The experimental observation has triggered tremendous studies.

Tetracharm  
hybrid



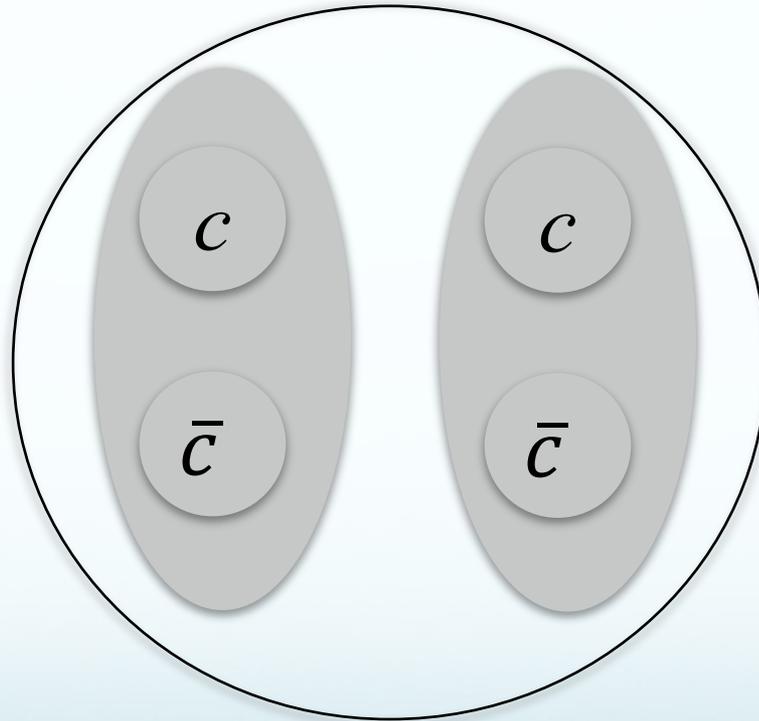
B.-D. Wan, C.-F. Qiao, Phys. Lett. B 817, 136339 (2021).

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# Background and Motivation

- The experimental observation has triggered tremendous studies.

Molecular state  
 $J/\psi J/\psi$



X.-K. Dong, V. Baru, F.-K. Guo, C. Hanhart, and A. Nefediev, Phys. Rev. Lett. 126, 132001 (2021)  
Y. Lu, C. Chen, K.-G. Kang, G.-Y. Qin, and H.-Q. Zheng, Phys. Rev. D 107, 094006 (2023).

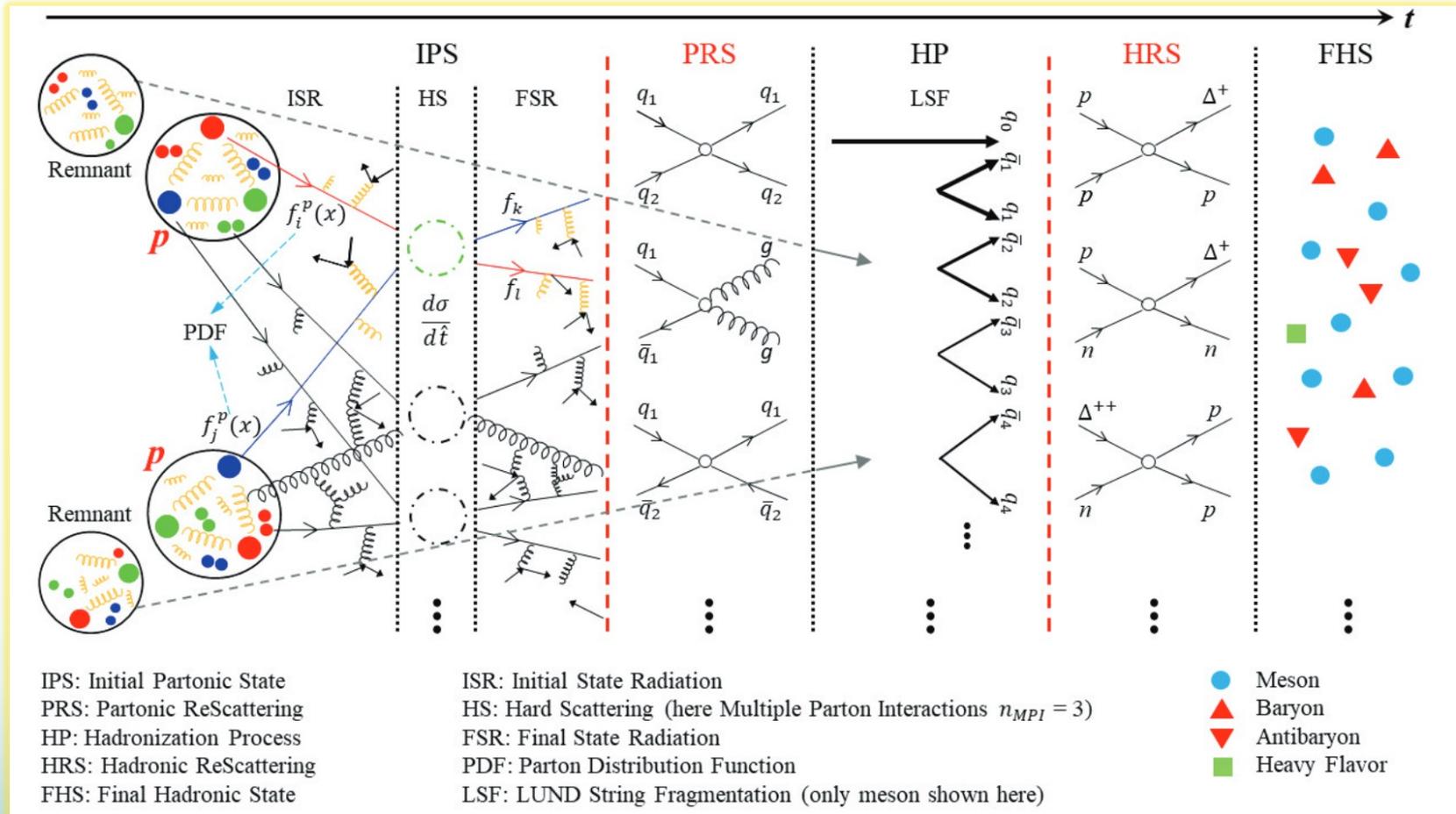
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# Background and Motivation

- In this work, we use a parton and hadron cascade model PACIAE 4.0 to investigate the  $J/\psi$  production in pp collisions at  $\sqrt{s} = 5.02, 7,$  and 13 TeV.
- We have also investigated the multiplicity dependence of  $J/\psi$  production in pp collisions at  $\sqrt{s} = 5.02, 7$  and 13 TeV.
- Finally, we have investigated the productions of X(6900) in the configuration of the tetraquark state and the double  $J/\psi$  molecular state.

# The PACIAE Model



PACIAE4.0 is based on PYTHIA 8.3 but further considers the partonic rescattering before hadronization and the hadronic rescattering after hadronization.

A.-K. Lei et al., Computer Physics Communications 310, 109520 (2025).

# The DCPC Model

- In the quantum statistical mechanics inspired dynamically constrained phase-space coalescence (DCPC) model, the yield of  $N$ -particle cluster in six-dimensional phase space can be estimated by [Phys. Rev. C 85, 024907 \(2012\)](#).

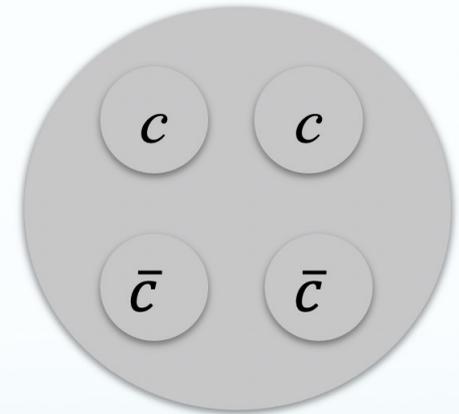
$$Y_{cluster} = \int \cdots \int \frac{d\vec{q}_1 d\vec{p}_1 \cdots d\vec{q}_N d\vec{p}_N}{h^{3N}}$$

- For the  $X(6900)$  tetraquark state, its yield reads

$$Y_{tetraquark} = \int \cdots \int \delta_{1234} \frac{d\vec{q}_1 d\vec{p}_1 d\vec{q}_2 d\vec{p}_2 d\vec{q}_3 d\vec{p}_3 d\vec{q}_4 d\vec{p}_4}{h^{12}}$$

$$\delta_{1234} = \begin{cases} 1. & \text{if } 1 = c, 2 = \bar{c}, 3 = c', 4 = \bar{c}' \\ & m_{X(6900)} - \Delta m < m_{inv} < m_{X(6900)} + \Delta m, \quad R \leq 1fm \\ 0. & \text{otherwise} \end{cases}$$

$\Delta m = m_{X(6900)}/10$ , ( $6.2 < m_{inv} < 7.4 \text{ GeV}/c^2$ ),  $R$  is the radius of the cluster.



# The DCPC Model

- In the quantum statistical mechanics inspired dynamically constrained phase-space coalescence (DCPC) model, the yield of  $N$ -particle cluster in six-dimensional phase space can be estimated by

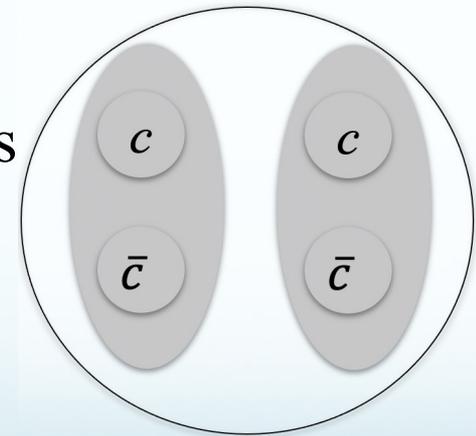
$$Y_{cluster} = \int \dots \int \frac{d\vec{q}_1 d\vec{p}_1 \dots d\vec{q}_N d\vec{p}_N}{h^{3N}}$$

- For the  $X(6900)$  double  $J/\psi$  molecular state, its yield reads as

$$Y_{glueball} = \frac{1}{2!} \int \int \delta_{12} \frac{d\vec{q}_1 d\vec{p}_1 d\vec{q}_2 d\vec{p}_2}{h^6}$$

$$\delta_{12} = \begin{cases} 1. & \text{if } 1 = J/\psi, 2 = J/\psi' \\ & m_{X(6900)} - \Delta m < m_{inv} < m_{X(6900)} - \Delta m, R > 1fm \\ 0. & \text{otherwise} \end{cases}$$

$\Delta m = m_{X(6900)}/10$ , ( $6.2 < m_{inv} < 7.4 \text{ GeV}/c^2$ ),  $R$  is the radius of the cluster.



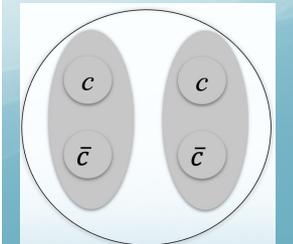
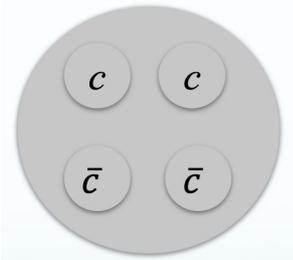
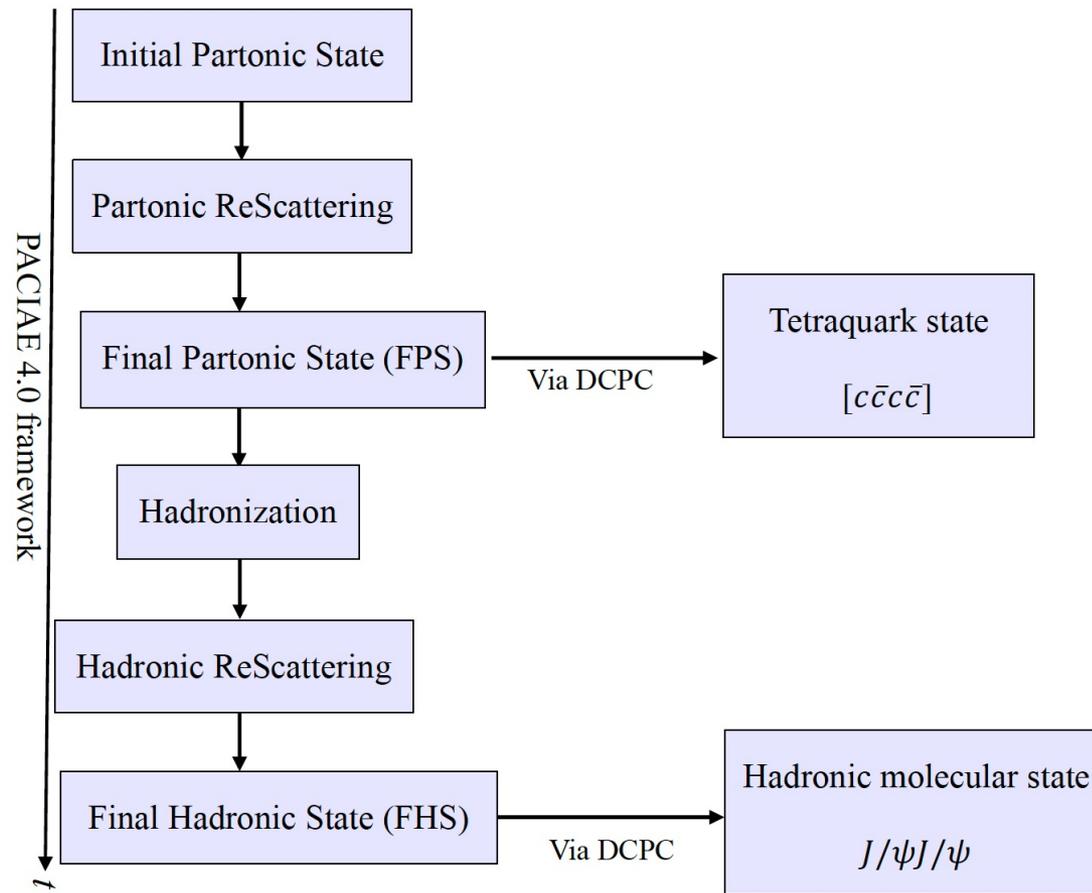


# The DCPC Model

- Using the DCPC model, we have investigated the production of the glueball-like particle  $X(2370)$ , the exotic states  $X(3872)$  and  $G(3900)$  in  $e^+e^-$  and  $pp$  collisions. See the followings for details.
- $X(2370)$ , [Phys. Rev. D 110, 054046 \(2024\)](#)
- $X(3872)$ , [Phys. Rev. C 110, 014910 \(2024\)](#)
- $G(3900)$ , [arXiv: 2502.16822](#)

# The DCPC Model

- The tetraquark state and double  $J/\psi$  hadronic molecular state of  $X(6900)$  are assumed to be coalesced via DCPC in PFS and HFS, respectively.





# Results and Discussions

- In PACIAE 4.0, different charmonium production mechanisms are implemented.
- 1. In the perturbative scattering processes, leading order NRQCD channels via color-singlet and color-octet pre-resonant states are included.
- For color-octet state, one additional gluon is emitted in the transition to the physical color-singlet state.



# Results and Discussions

- In PACIAE 4.0, **different production mechanisms** for the charmonium productions are implemented.
- 2. Charmonium can be produced from the **cluster collapse mechanism**.
- If  $c$  and  $\bar{c}$  are close in phase space the potential energy in the string is **sufficient to create a light quark-antiquark pair**, and instead the two heavy quarks bin into a charmonium bound state.
- 3. In the so-called **non-prompt** charmonium production, charmonia can be produced from the weak decay of a hadron containing a beauty quark.

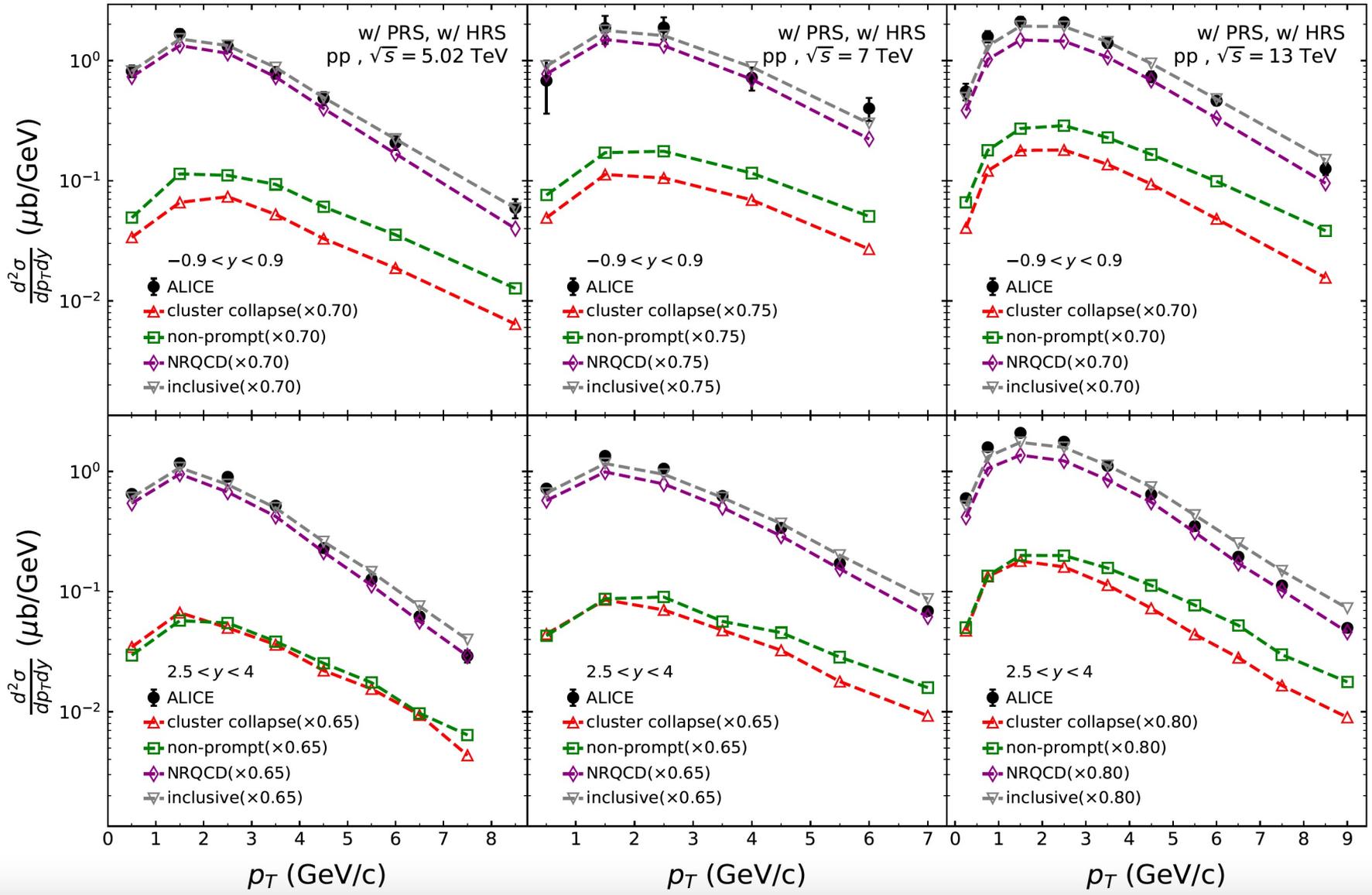
# Results and Discussions

- For the  $J/\psi$  hadronic rescattering, besides the elastic scattering, so far the following inelastic processes are considered.
- So far, the scattering of  $J/\psi$  with partons is not considered in PACIAE4.0 .

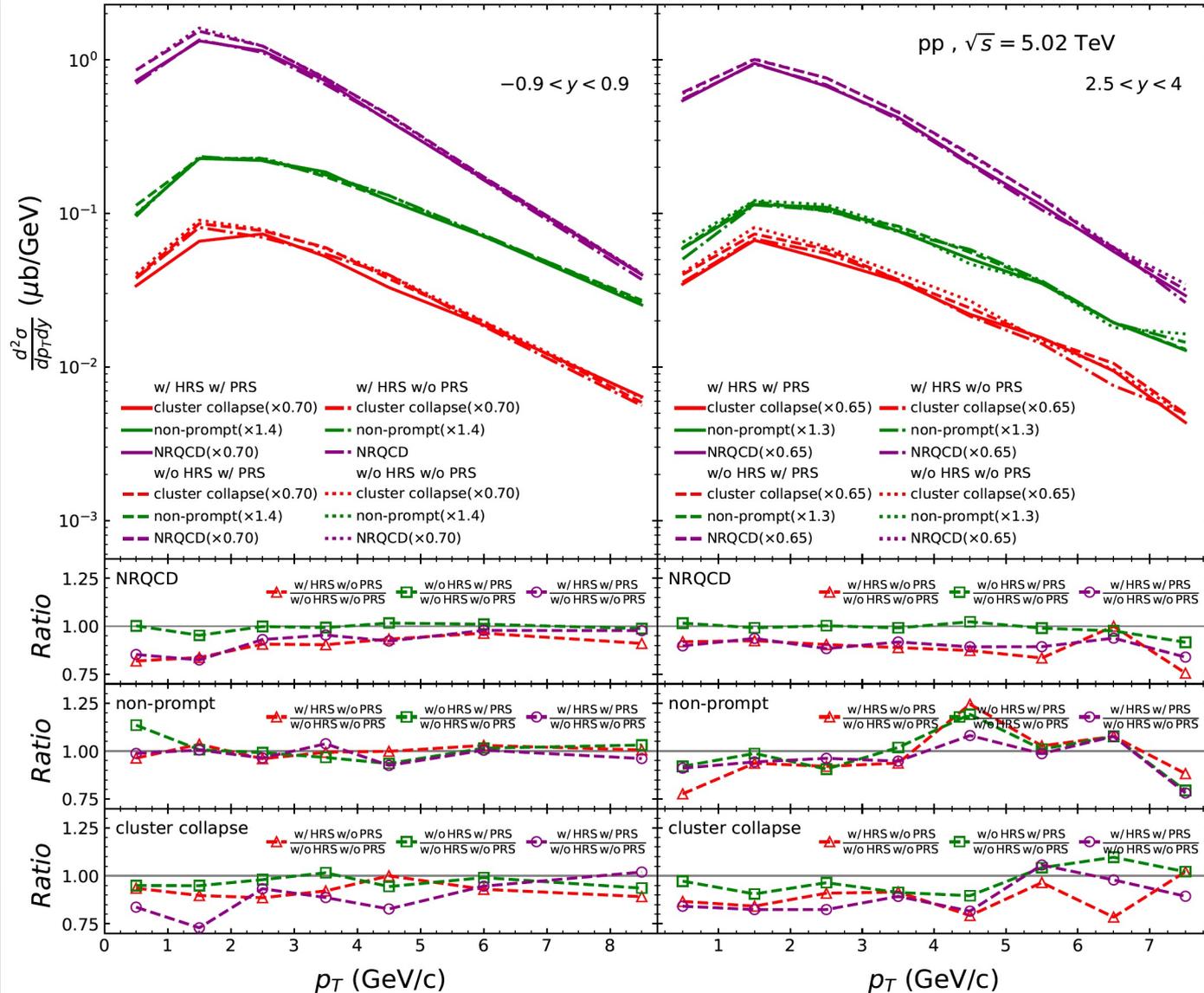
$$\begin{aligned} J/\psi + n &\rightarrow \Lambda_c^+ + D^-, & J/\psi + n &\rightarrow \Sigma_c^+ + D^-, \\ J/\psi + n &\rightarrow \Sigma_c^0 + \bar{D}^0, & J/\psi + p &\rightarrow \Lambda_c^+ + \bar{D}^0, \\ J/\psi + p &\rightarrow \Sigma_c^+ + \bar{D}^0, & J/\psi + p &\rightarrow \Sigma_c^{++} + D^-, \\ J/\psi + \pi^+ &\rightarrow D^+ + \bar{D}^{*0}, & J/\psi + \pi^- &\rightarrow D^0 + D^{*-}, \\ J/\psi + \pi^0 &\rightarrow D^0 + \bar{D}^{*0}, & J/\psi + \pi^0 &\rightarrow D^+ + D^{*-}, \\ J/\psi + \rho^+ &\rightarrow D^+ + \bar{D}^0, & J/\psi + \rho^- &\rightarrow D^0 + D^-, \\ J/\psi + \rho^0 &\rightarrow D^0 + \bar{D}^0, & J/\psi + \rho^0 &\rightarrow D^+ + D^-. \end{aligned}$$



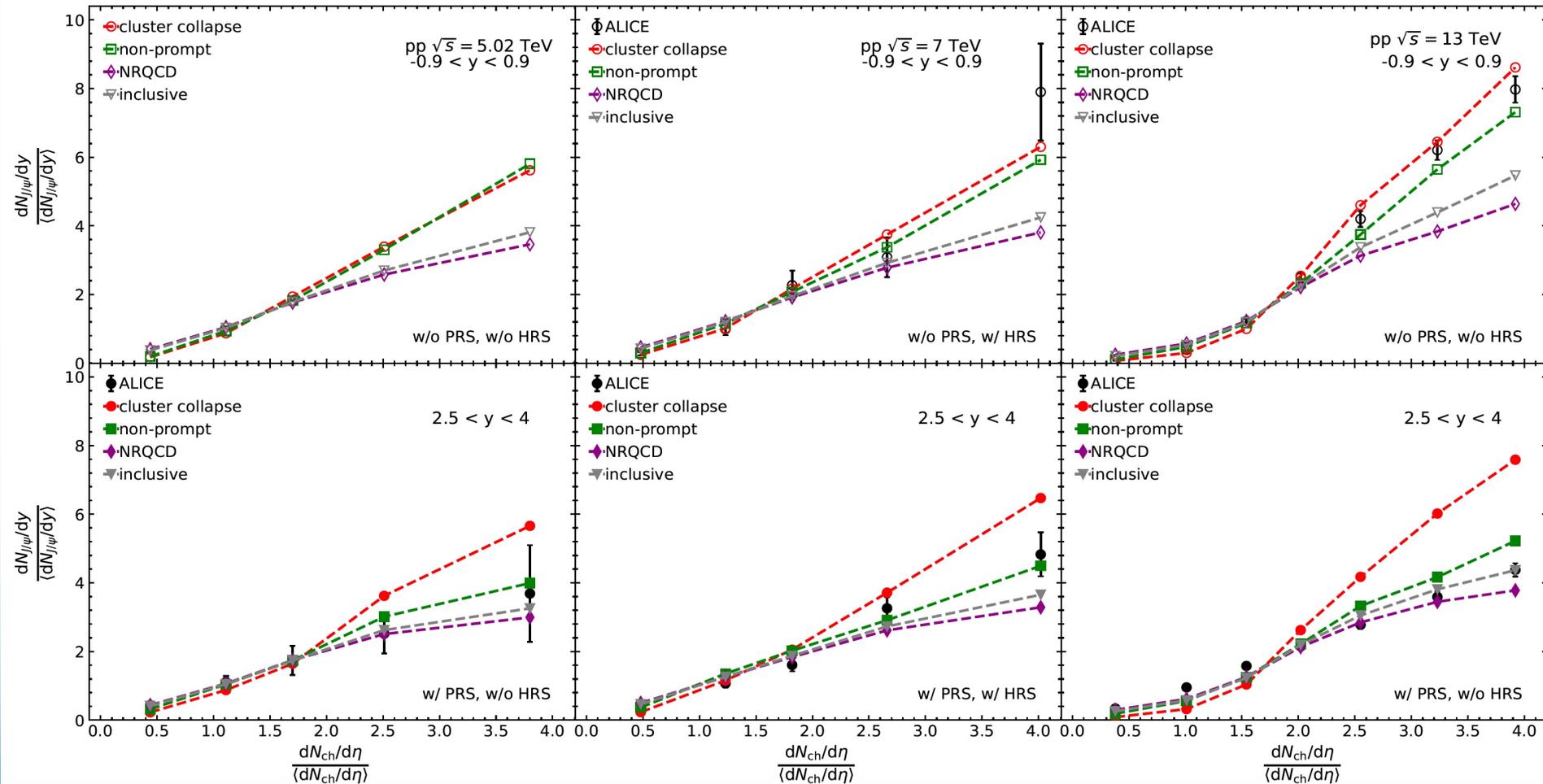
# Results and Discussions



# Results and Discussions

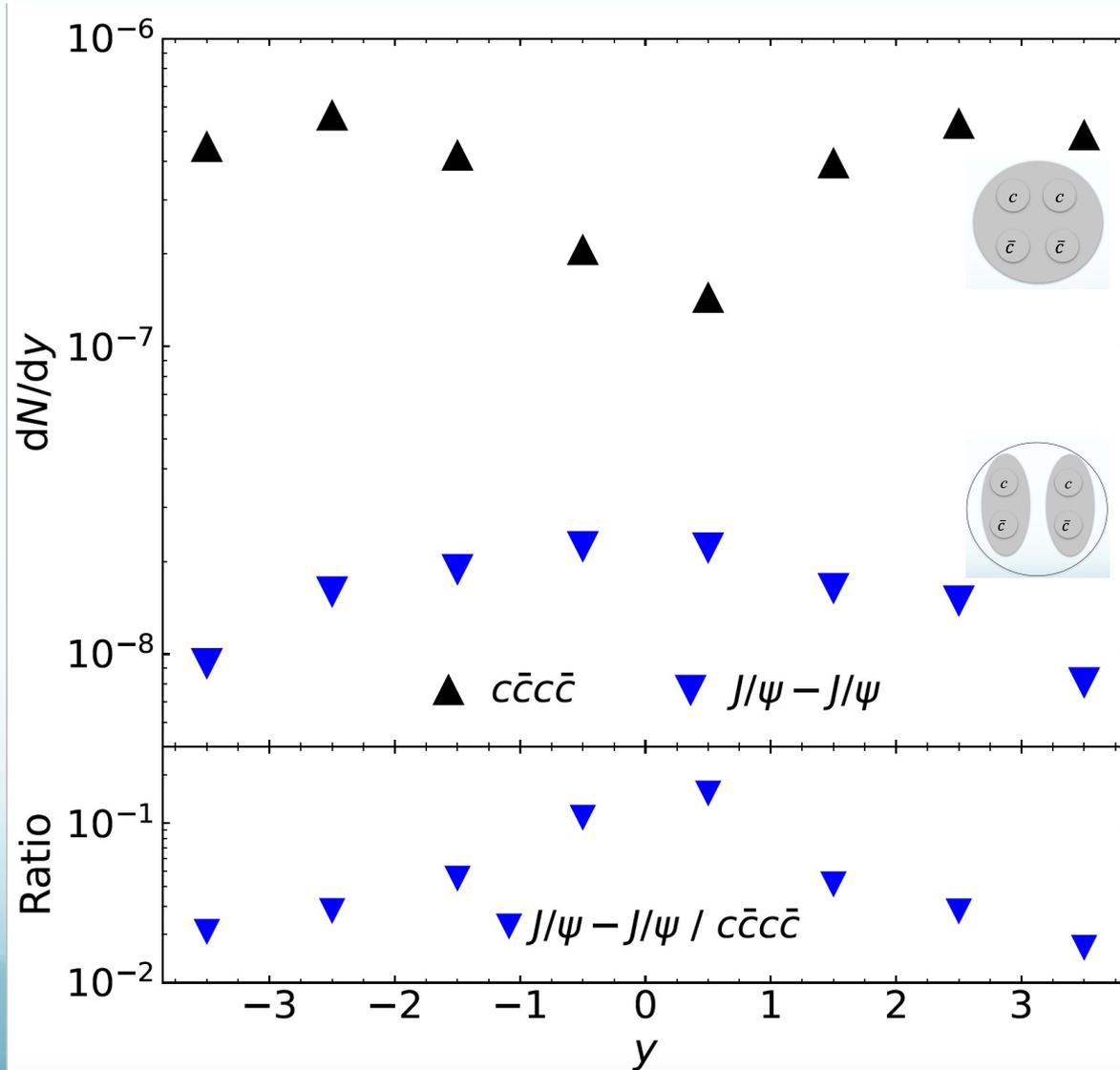


# Results and Discussions

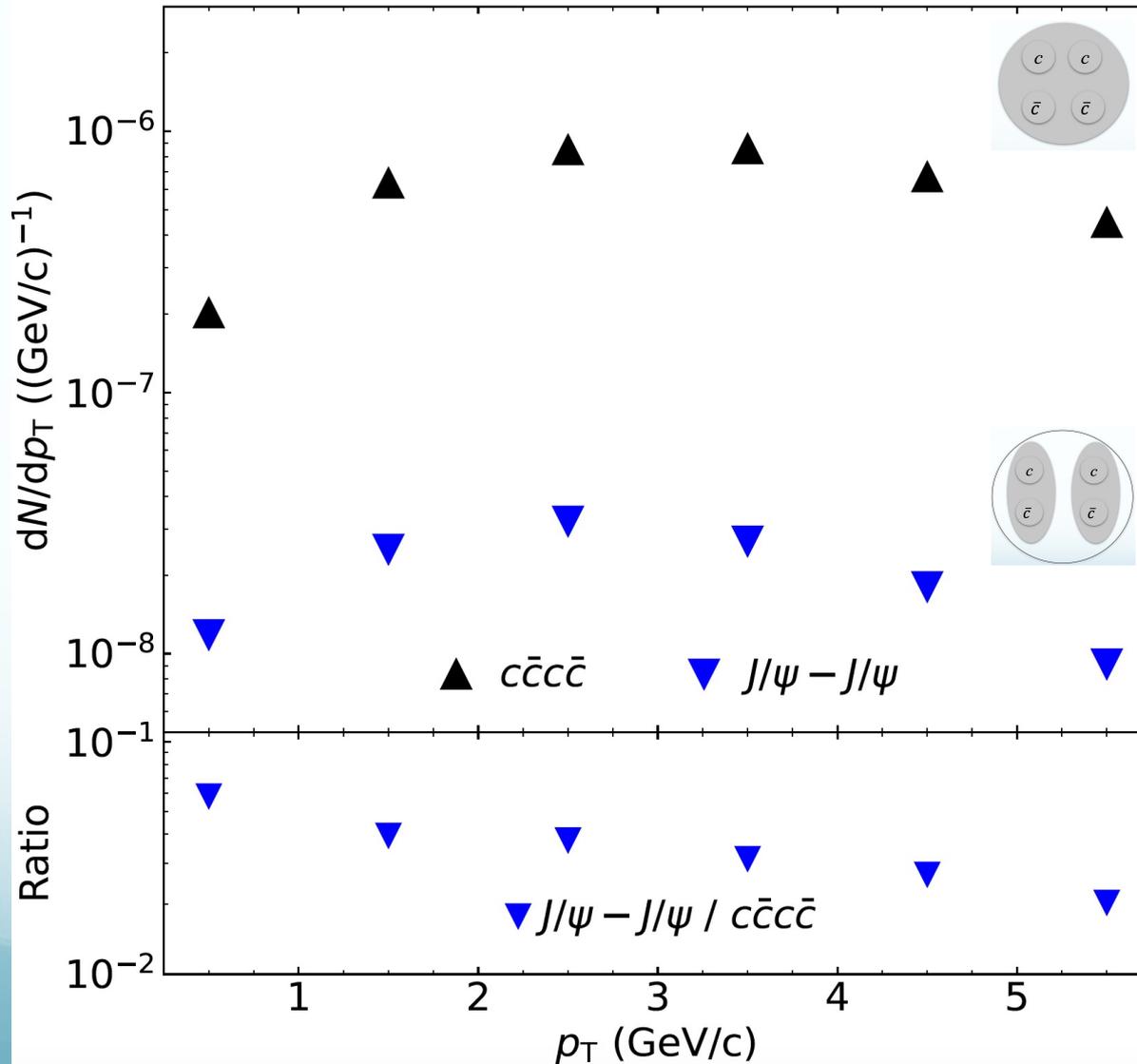


$\langle dN_{J/\psi} \rangle$  and  $\langle dN_{ch} \rangle$  are the average  $J/\psi$  yield and the charged particle multiplicity in minimum bias events.

# Results and Discussions



# Results and Discussions





# Conclusions

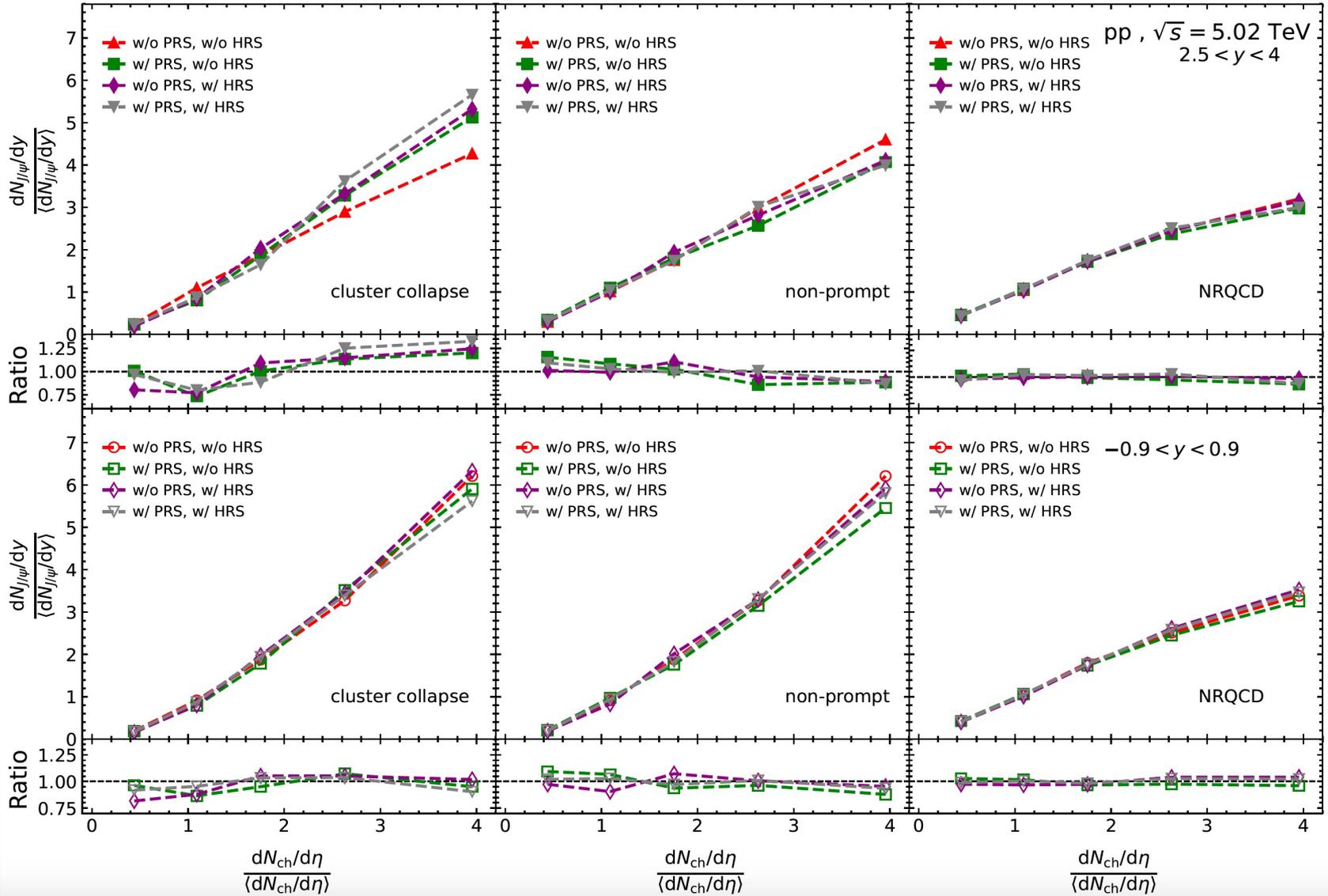
- We have investigated the  $J/\psi$  productions in pp collisions at  $\sqrt{s} = 5.02, 7$  and 13 TeV with PACIAE 4.0.
- We have also investigated the multiplicity dependence of  $J/\psi$  production in pp collisions at  $\sqrt{s} = 5.02, 7$  and 13 TeV.
- The contributions of the  $J/\psi$  productions from the NRQCD, cluster collapse and the weak decay from the B hadrons are analyzed in detail.
- The effect of the partonic and hadronic rescattering effects on the  $J/\psi$  production has been presented.
- Discrepancy between tetraquark- and molecule-state in the rapidity and  $p_T$  distributions can be served as a criterion for identifying the configuration of X(6900).



# Back up

- If MPI were mainly affecting processes involving only light quarks and gluons, as implemented e.g. in PYTHIA 6.4, processes like  $J/\psi$  and open heavy flavor production should not be influenced and their rates are expected to be independent of the overall event multiplicity.
- However, at the high center-of-mass energies reached at the LHC, there might be a substantial contribution of MPI on a harder scale which can also induce a correlation between the yield of quarkonia and the total charged particle multiplicity.

# Back up





# Back up

$$\frac{dN_{J/\psi}}{\langle dN_{J/\psi} \rangle} = \frac{N_{J/\psi}^i}{N_{J/\psi}^{tot}} \cdot \frac{N_{evt}^{tot}}{N_{evt}^i}$$

$N_{J/\psi}^i$  and  $N_{evt}^i$  are the number of  $J/\psi$  and number of events in  $i^{\text{th}}$  multiplicity bin, respectively.  $N_{J/\psi}^{tot}$  and  $N_{evt}^{tot}$  are the total number of  $J/\psi$  produced and total number of minimum bias events, respectively.