



河南省科学院
HENAN ACADEMY OF SCIENCES



Energy dependence of transverse momentum fluctuations in Au+Au collisions from a multiphase transport model

Physical Review C 111, 024911 (2025)

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In collaboration with Jinhui Chen², Chunjian Zhang²

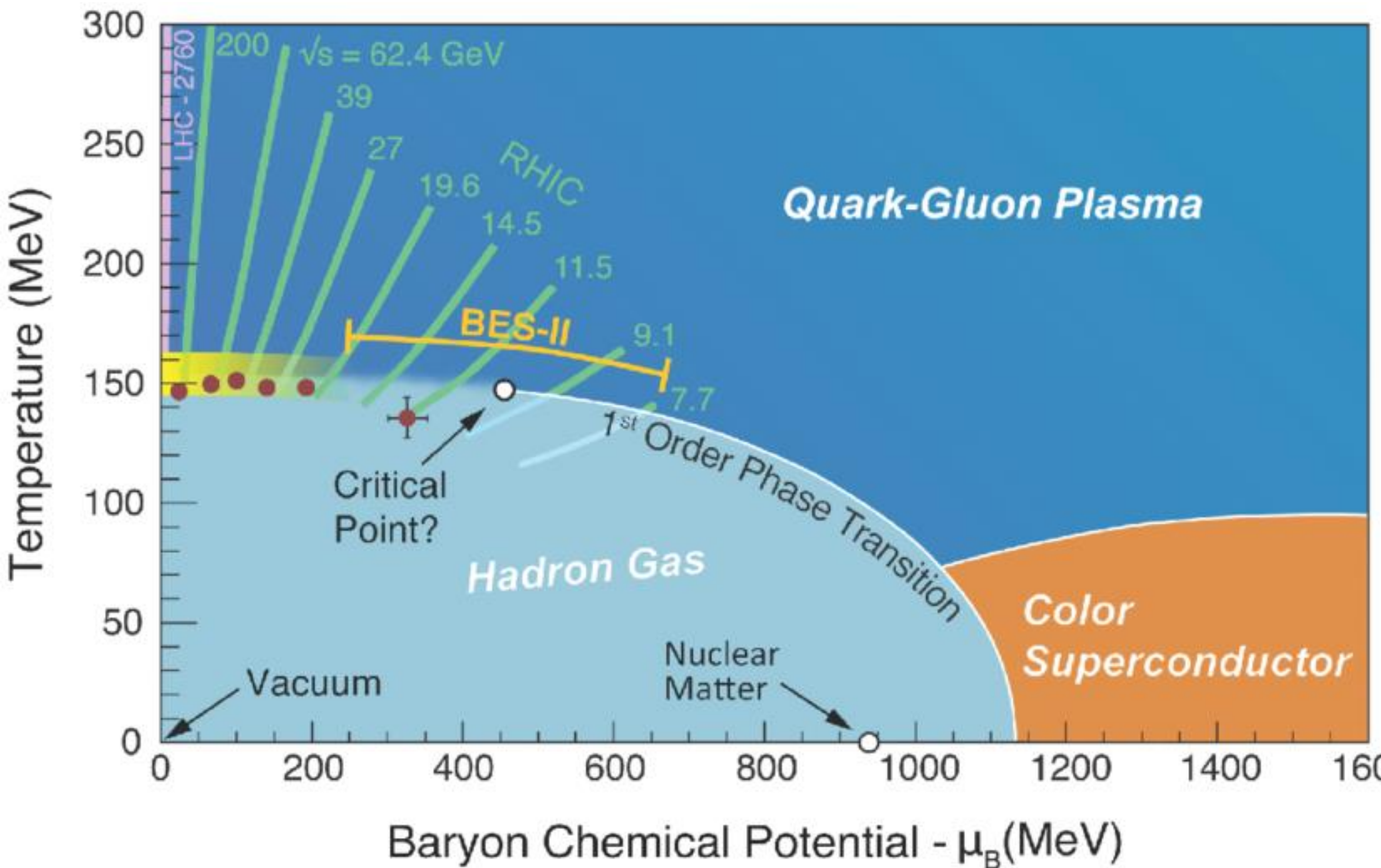
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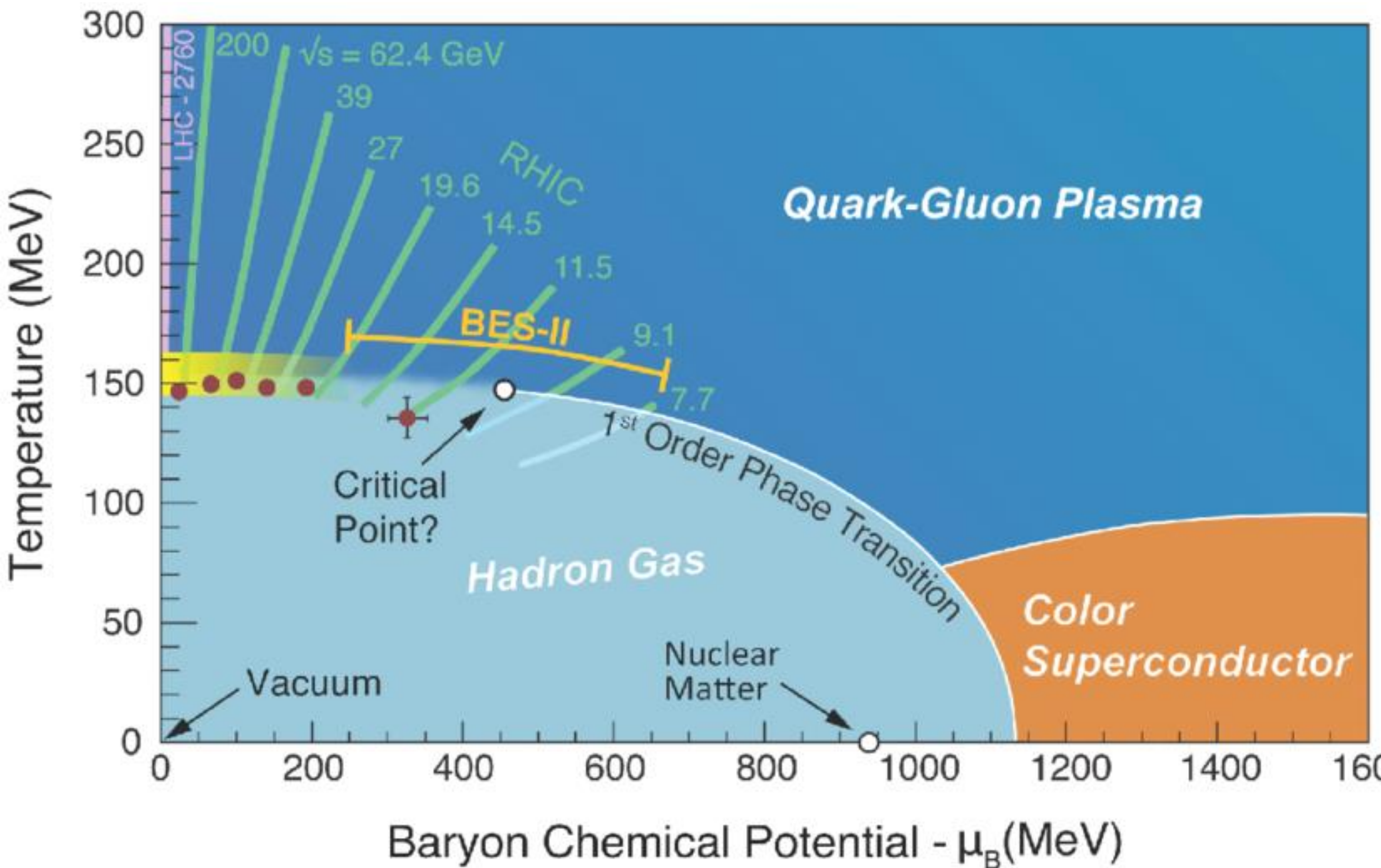
The evidence of Quark-Gluon Plasma (QGP) existence: searching critical point and 1st order boundary.



Sensitive observables :

- Collectivity: long-range correlation.
- Chirality: polarization.
- Criticality: net proton/charge/strange fluctuation, light nucleus yield ratios fluctuation.

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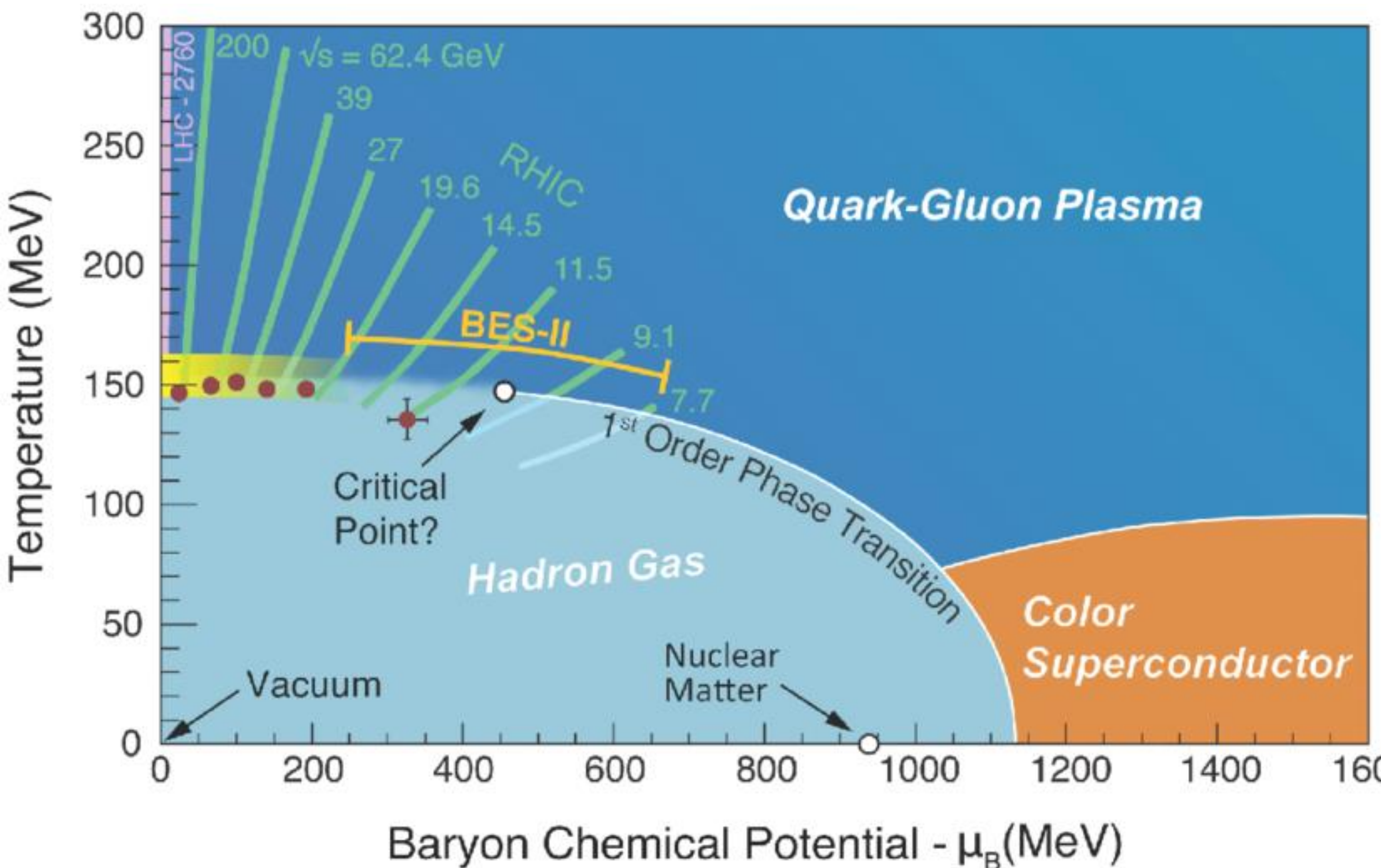


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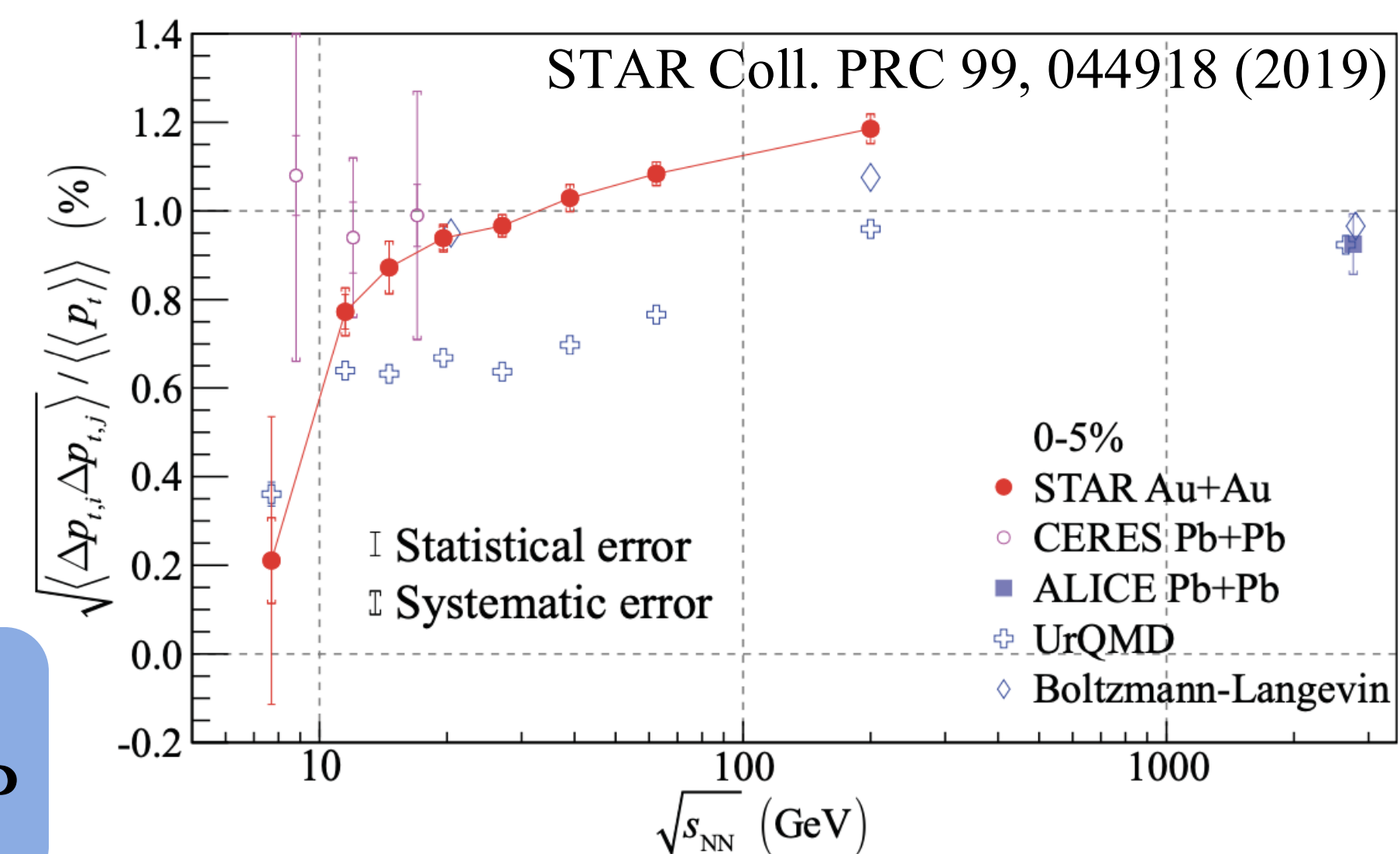
EbE $\langle p_T \rangle$ fluctuations

The evidence of Quark-Gluon Plasma (QGP) existence: searching critical point and 1st order boundary.



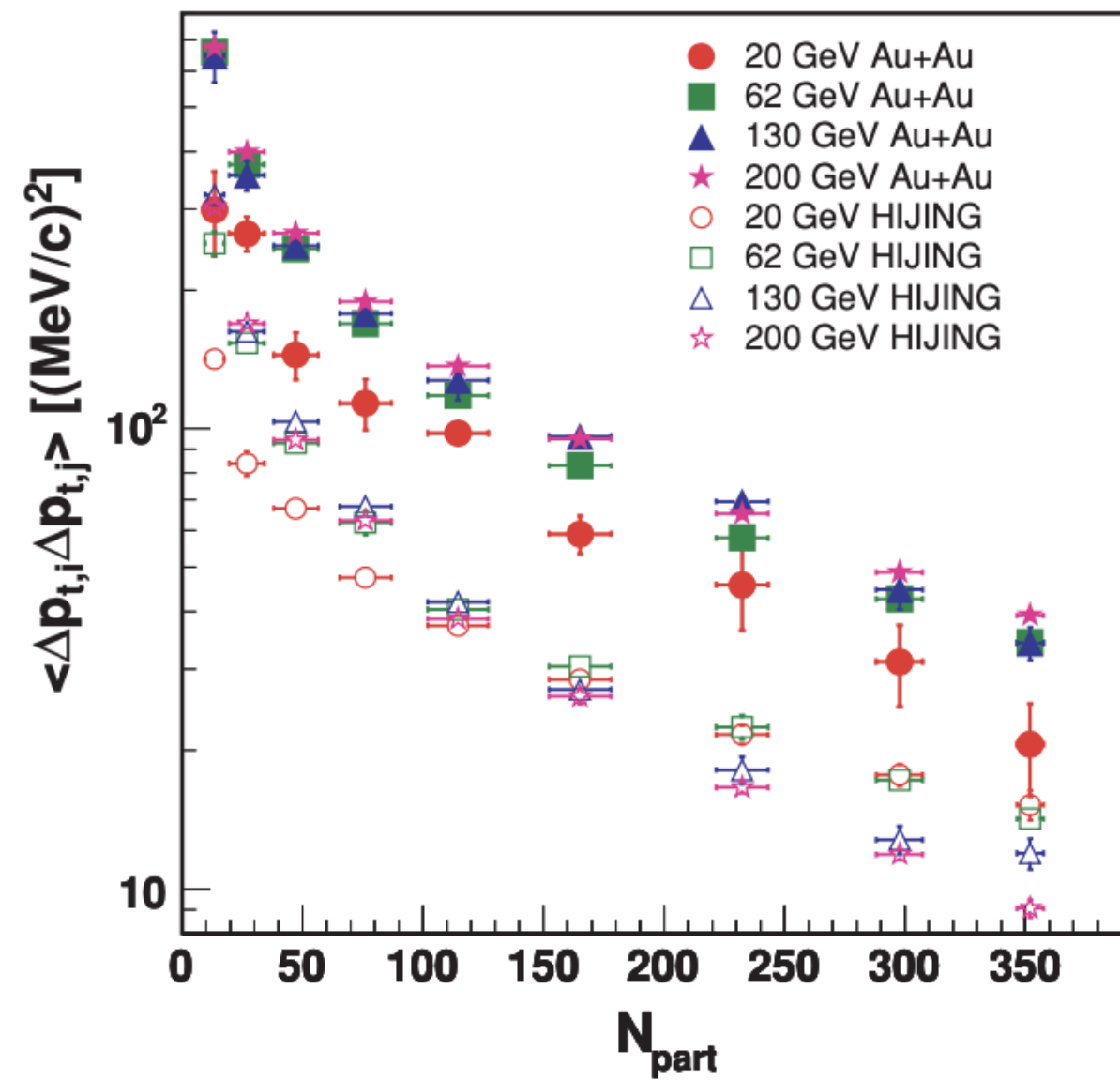
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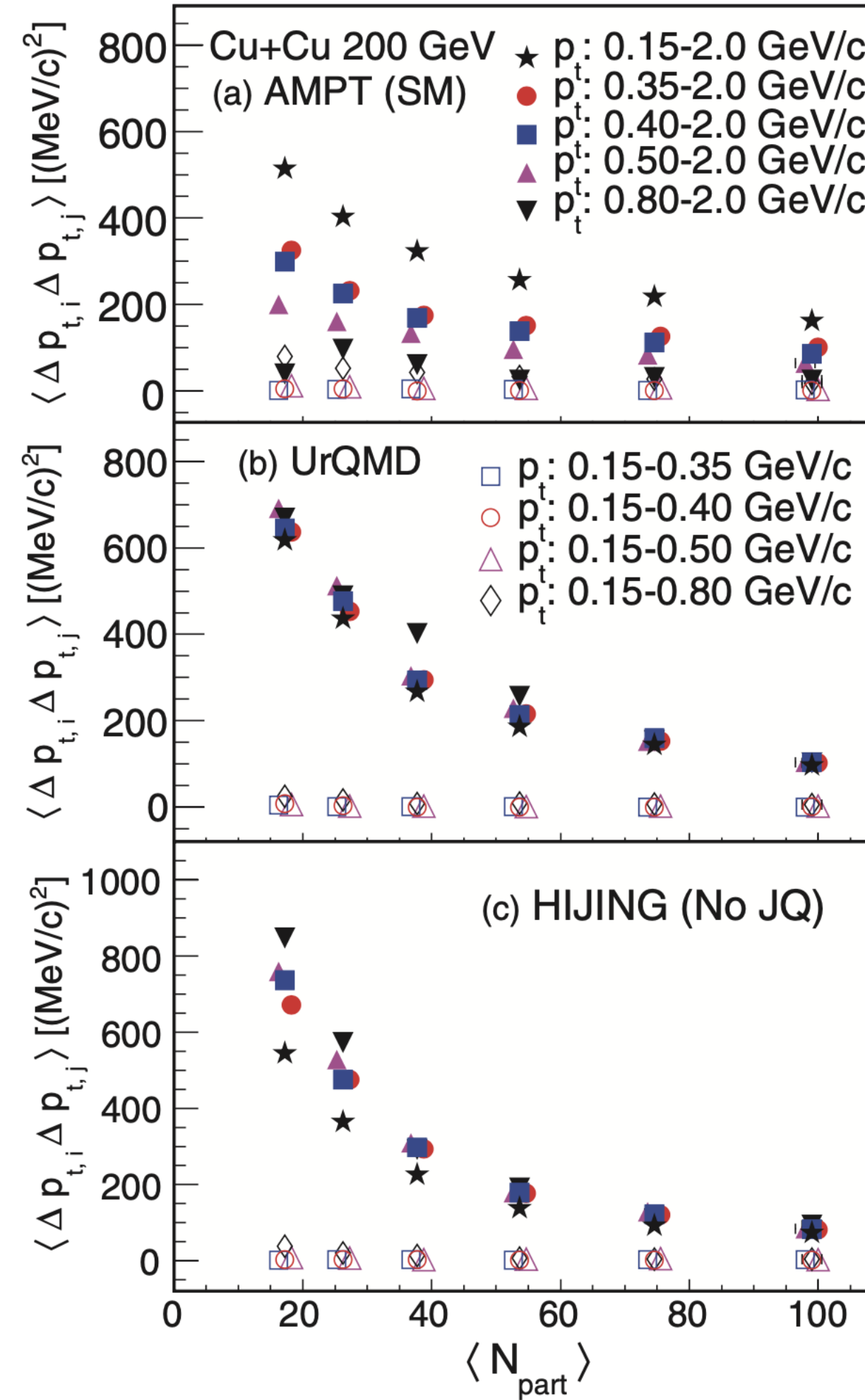


A non-monotonic behavior of $\langle p_T \rangle$ fluctuations as a function of centrality or incident energy was suggested as one of the possible signals of the QGP

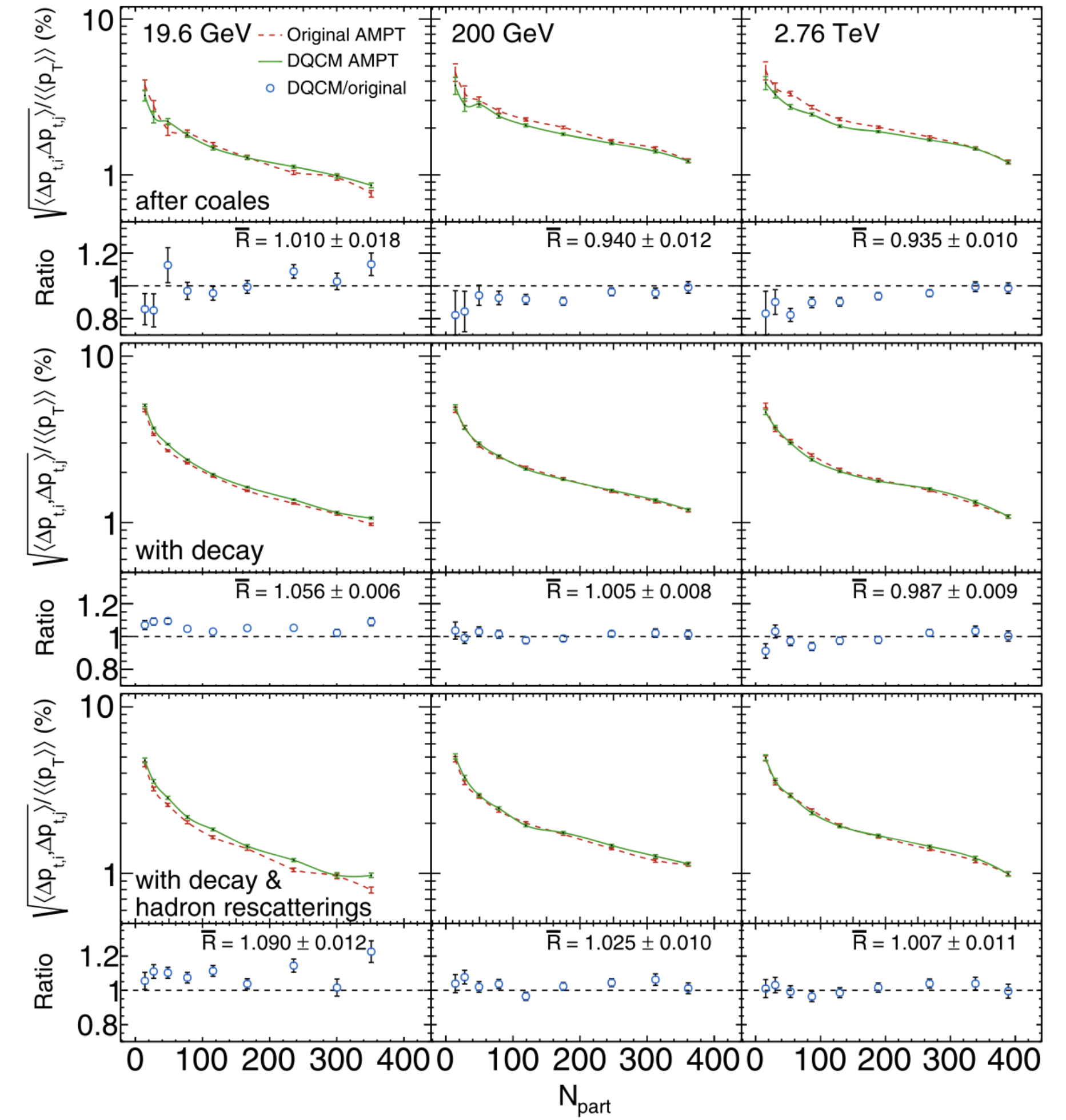
PRC 72, 044902 (2005)



PHYSICAL REVIEW C 87, 064902 (2013)



Zhenyu Xu et al 2020 J. Phys. G: Nucl. Part. Phys. 47 125102



The higher-order $\langle p_T \rangle$ fluctuations need to be systematically explored considering the BES program.

$\langle p_T \rangle$ trend : AMPT optimization

1) vs centrality, 2) vs incident energy

Two key parameters in hadronization mechanisms in the AMPT model:

a_L, b_L : which are inversely related to $\langle p_T \rangle^2$ of hadrons.

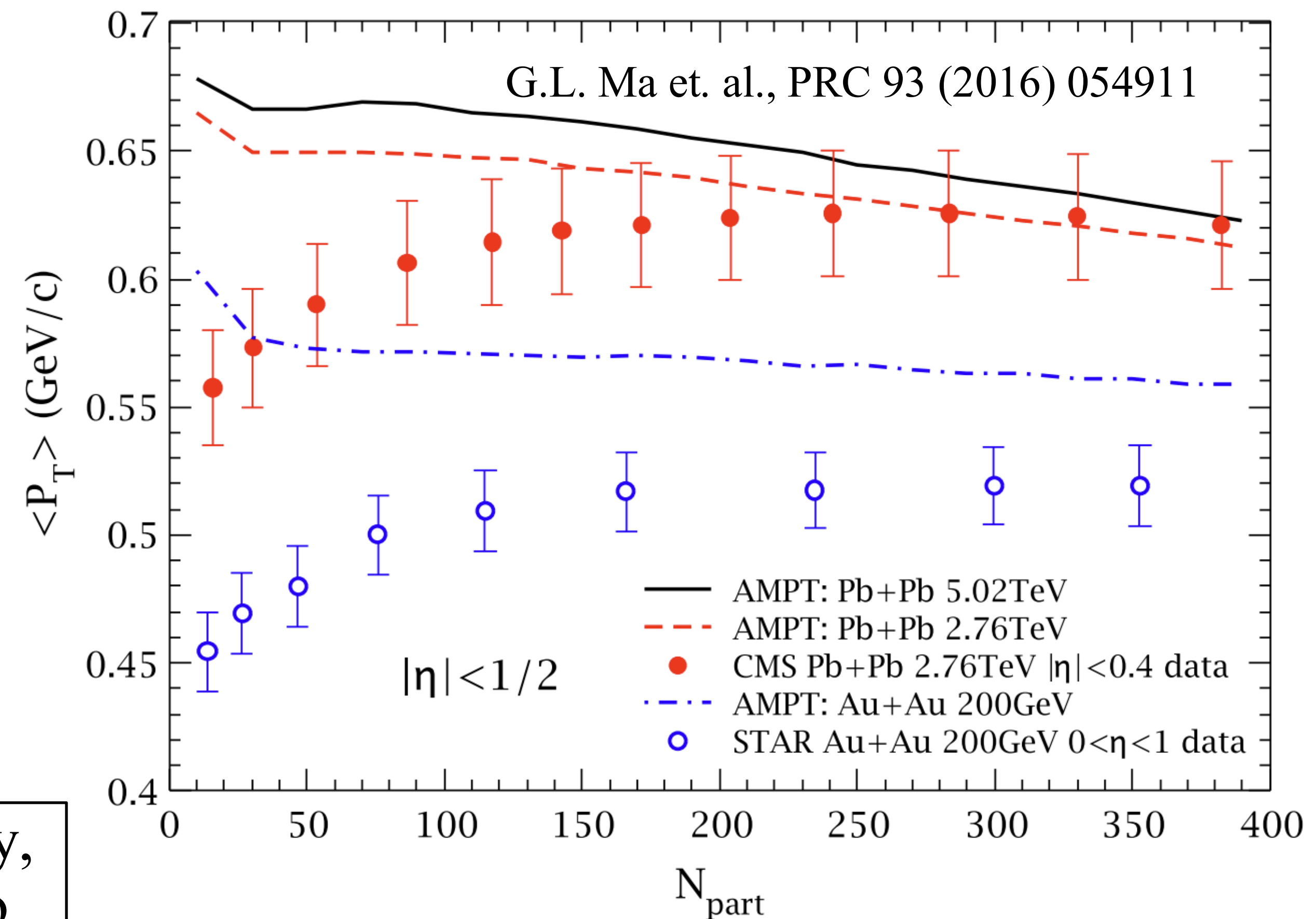
$$\langle p_T^2 \rangle = \frac{1}{b_L(2 + a_L)}$$

Traditionally, a_L, b_L are constant values, e.g.

RHIC@200 GeV: $a_L = 0.55, b_L = 0.15$

LHC@2.75/5.02 TeV: $a_L = 0.3, b_L = 0.15$

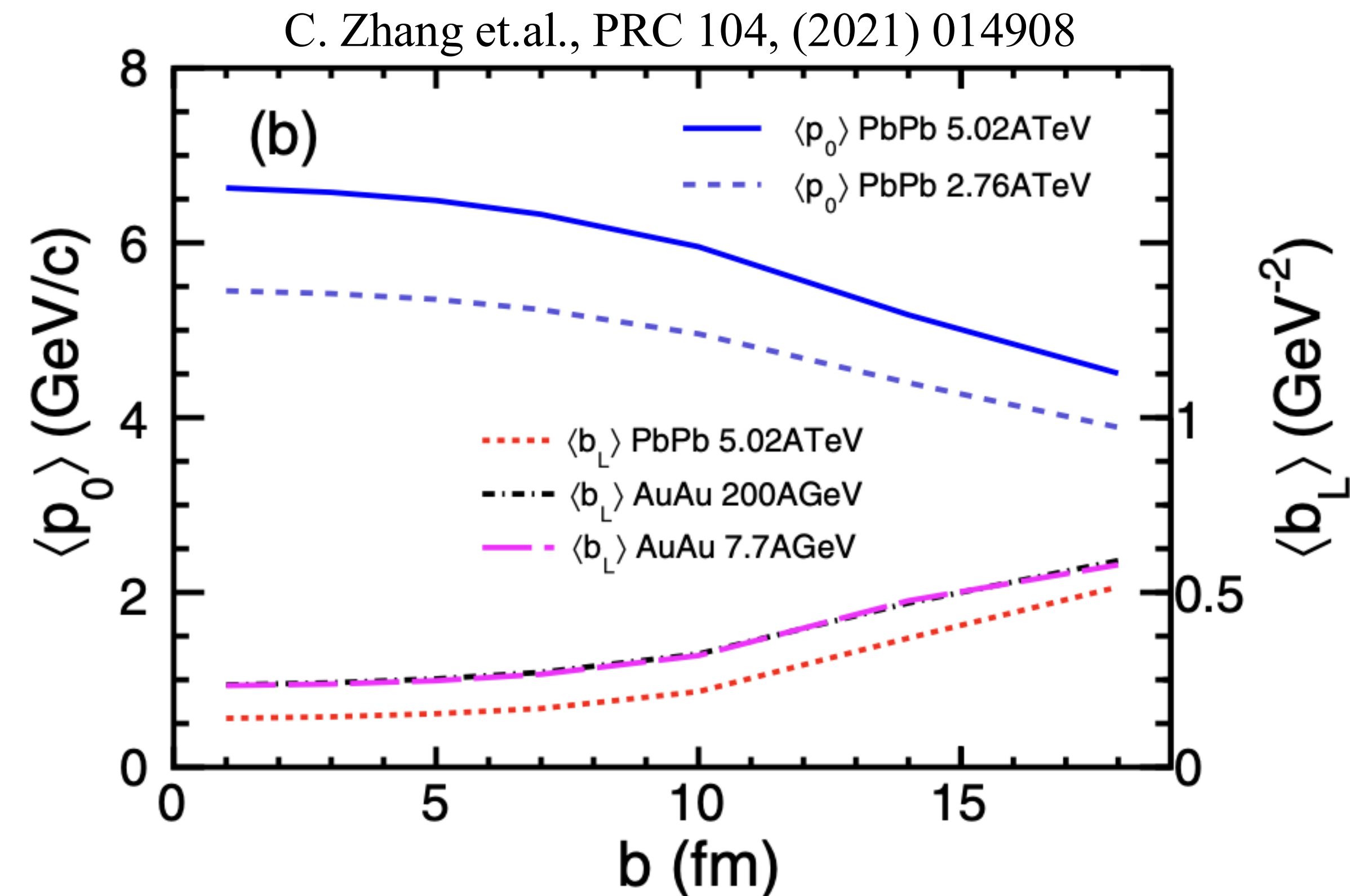
The $\langle p_T \rangle$ (cent.) is inconsistent with the data. Actually, the $\langle p_T \rangle$ is expected to increase with centrality due to higher initial temperature in more central collisions.



$\langle p_T \rangle$ trend : AMPT optimization

A possible solution:

Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.



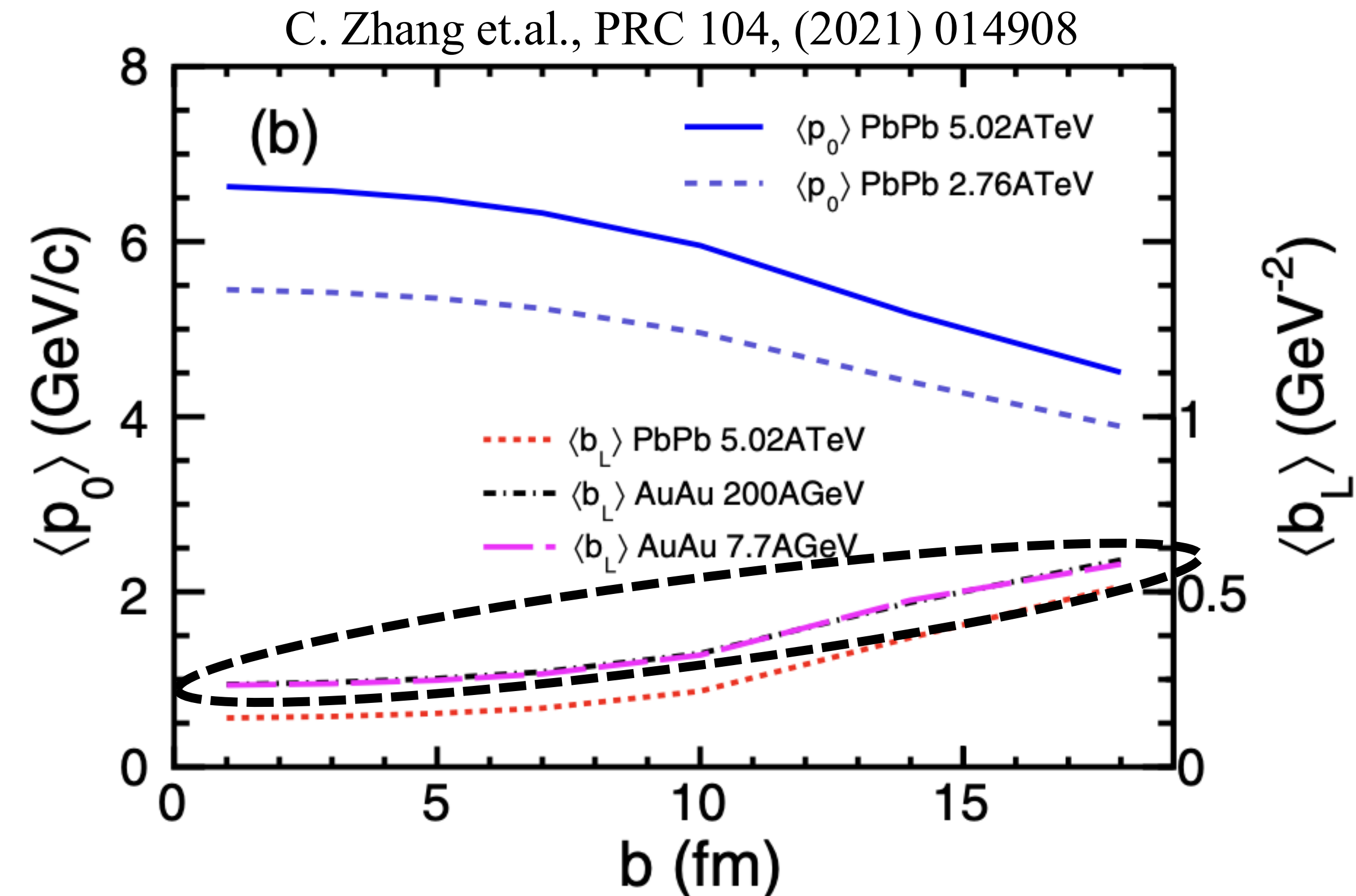
At 200 GeV and 7.7 GeV:

C. Zhang's work reveals b_L has approximate linear dependence on the impact parameter.

$\langle p_T \rangle$ trend : AMPT optimization

A possible solution:

Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.



Based on the published work, we adopt:

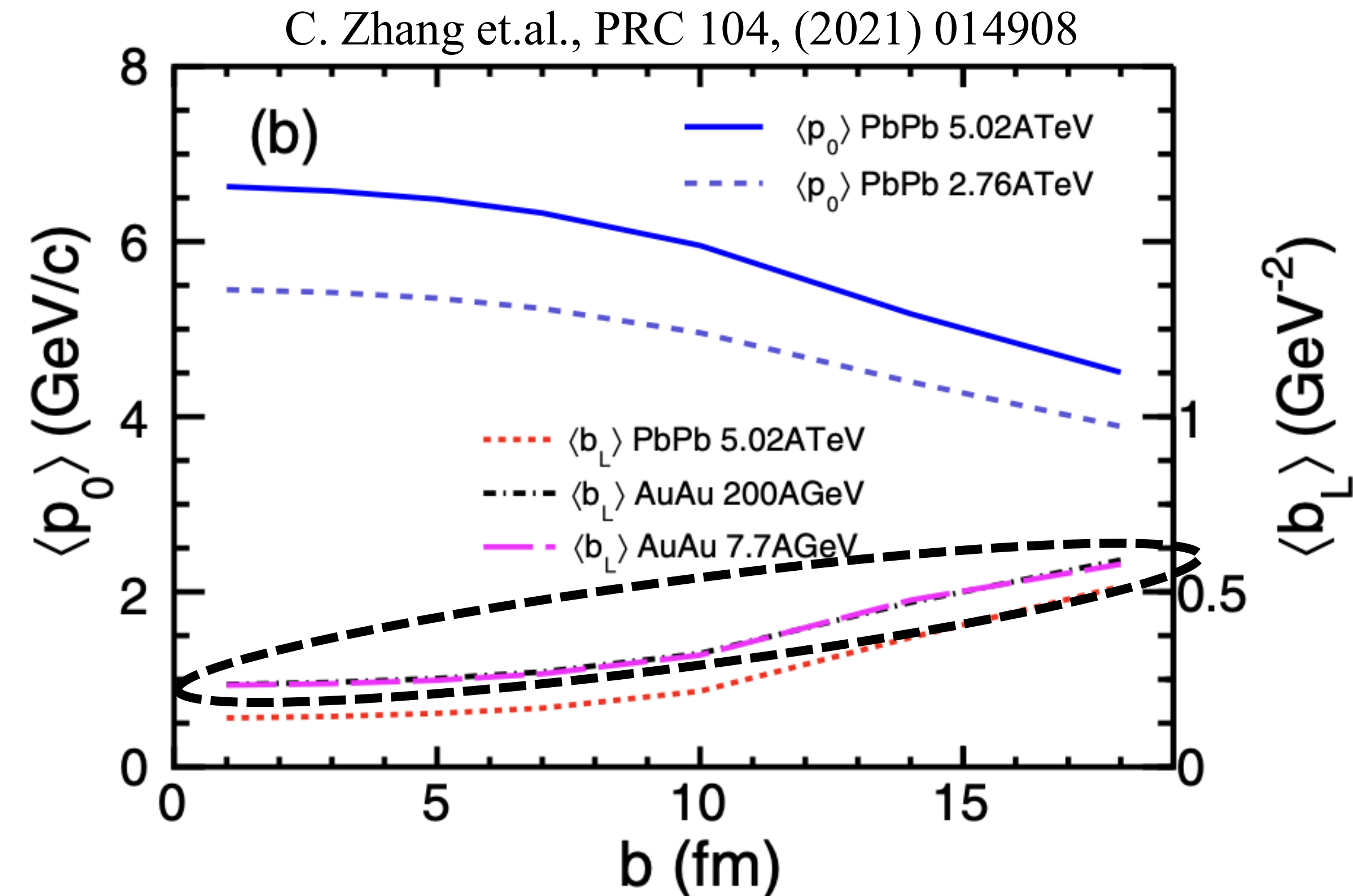
a) b_L has approximate linear dependence on the b .

→ Distribution of $\langle p_T \rangle$ (cent.) have a reasonable trend.

$\langle p_T \rangle$ trend : AMPT optimization

A possible solution:

Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.



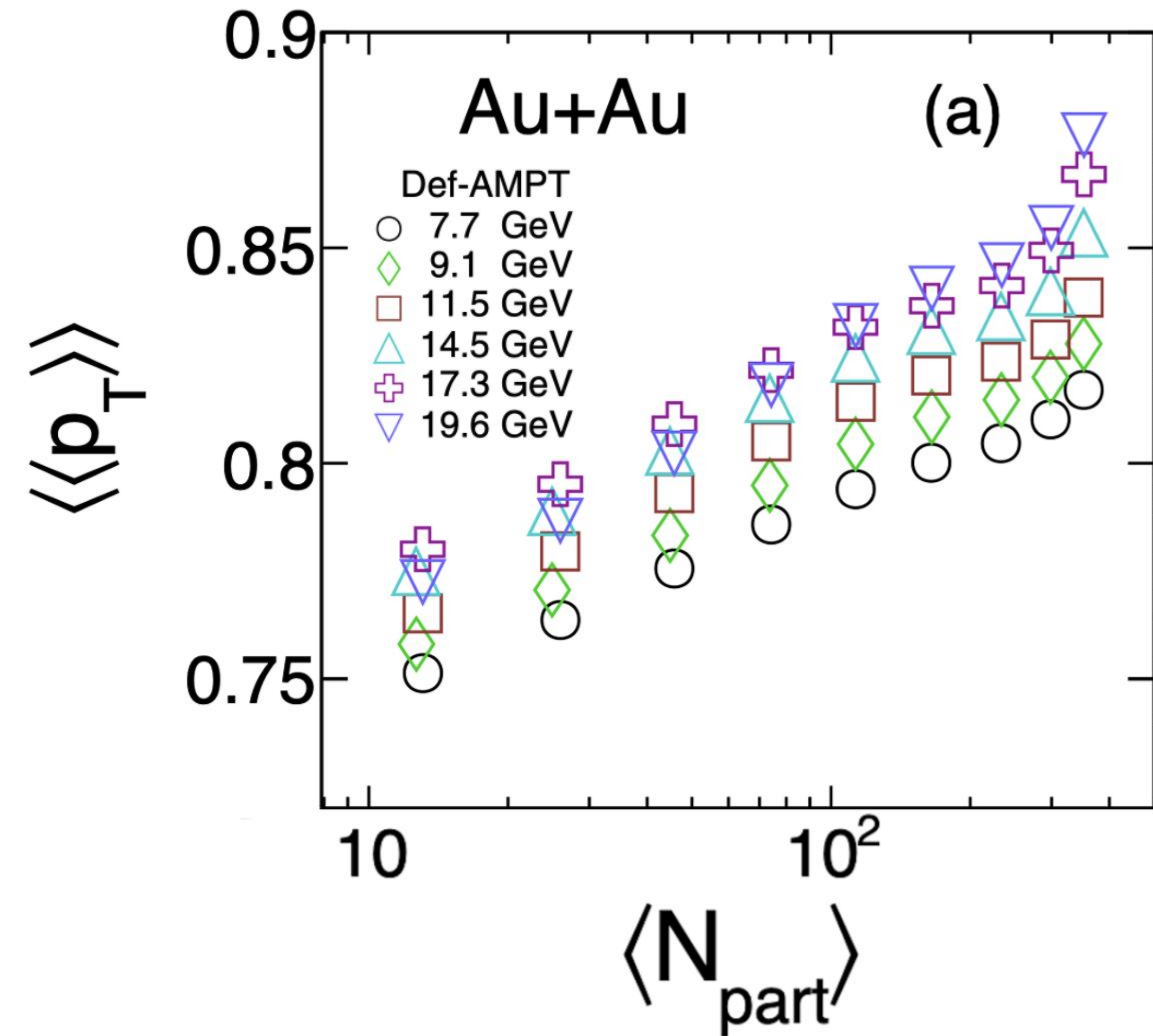
Based on the published work, we adopt:

a) b_L has approximate linear dependence on the b .

→ Distribution of $\langle p_T \rangle$ (cent.) have a reasonable trend.

b) b_L is systematically increased with E .

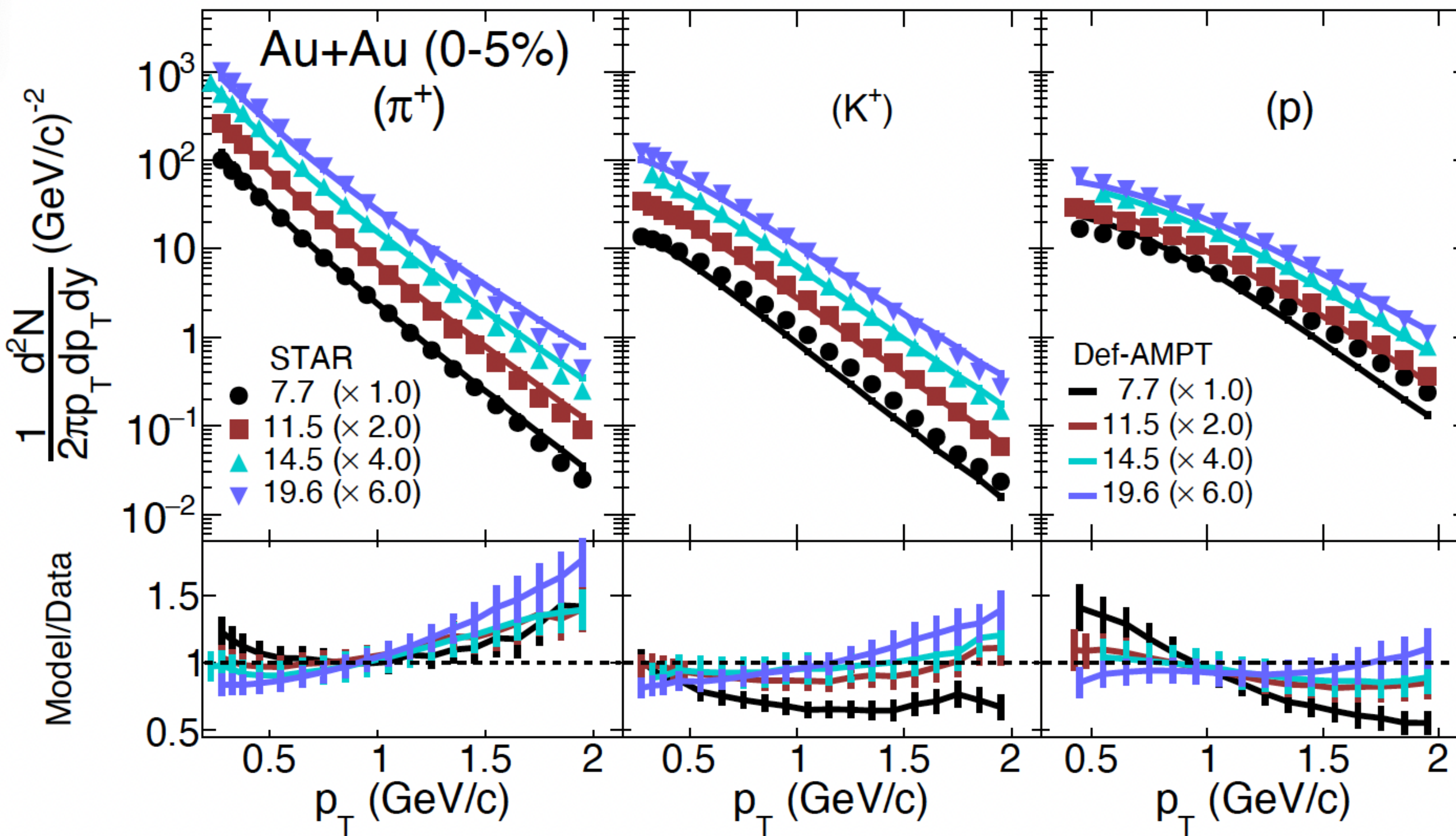
→ Distribution of $\langle p_T \rangle$ (cent.) is higher at higher incident energy due to enhanced collective effects.



Combine the optimization (a) and (b), the $\langle\langle p_T \rangle\rangle$ distributions from the enhanced AMPT exhibit:

- ✓ A reasonable centrality dependence, verified the expectation that higher initial temperatures in more central collisions.
- ✓ A significant energy dependence is observed across beam energy scan, with larger $\langle\langle p_T \rangle\rangle$ values at higher incident energy due to enhanced collective effects.

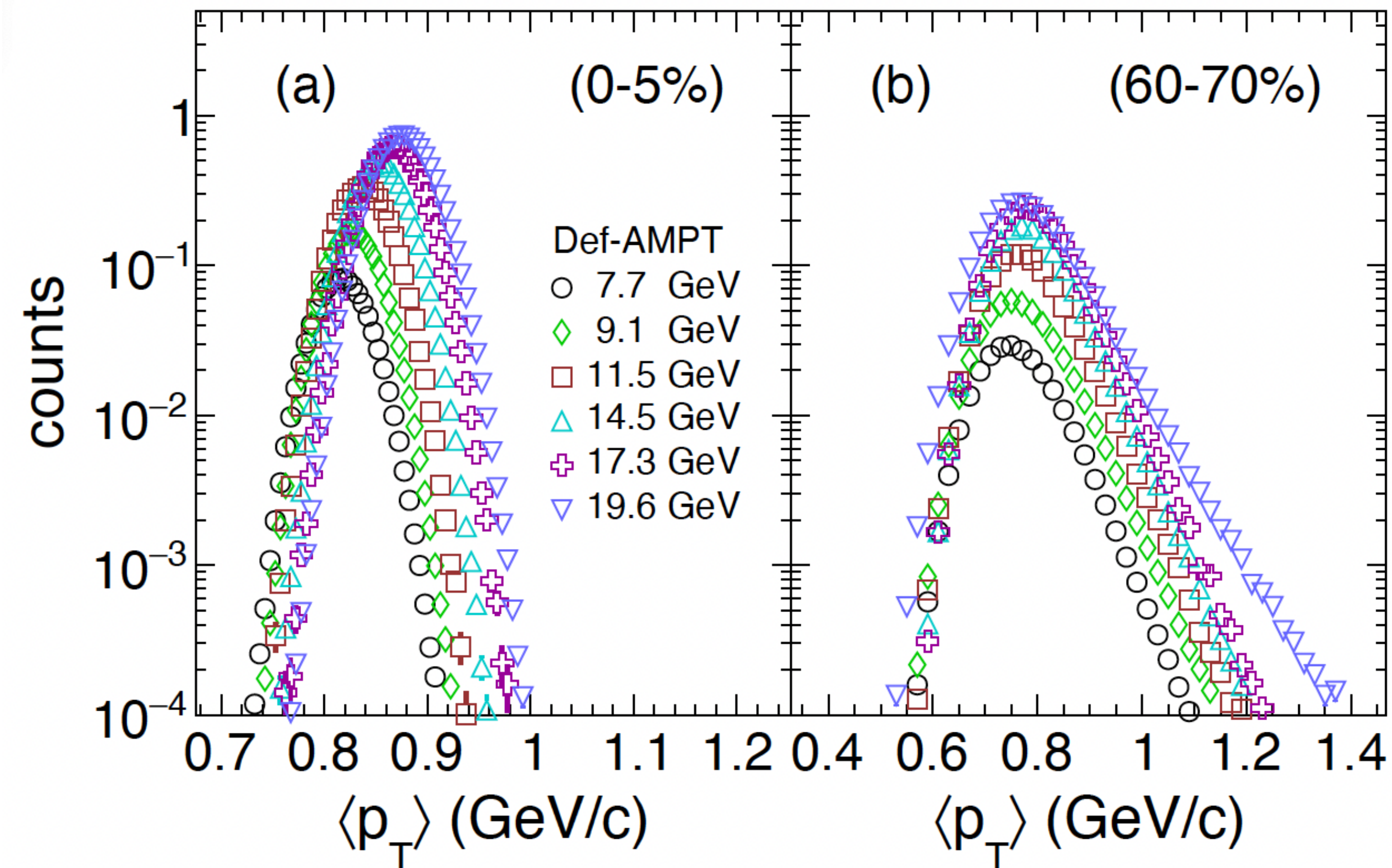
The AMPT validity



The improved AMPT model qualitatively produces the experimental data across a broad p_T spectrum, spanning a variety of incident energy.

Event-wise $\langle p_T \rangle$ distribution

The $\langle p_T \rangle$ fluctuations can be straightforwardly studied by its event-wise distributions.

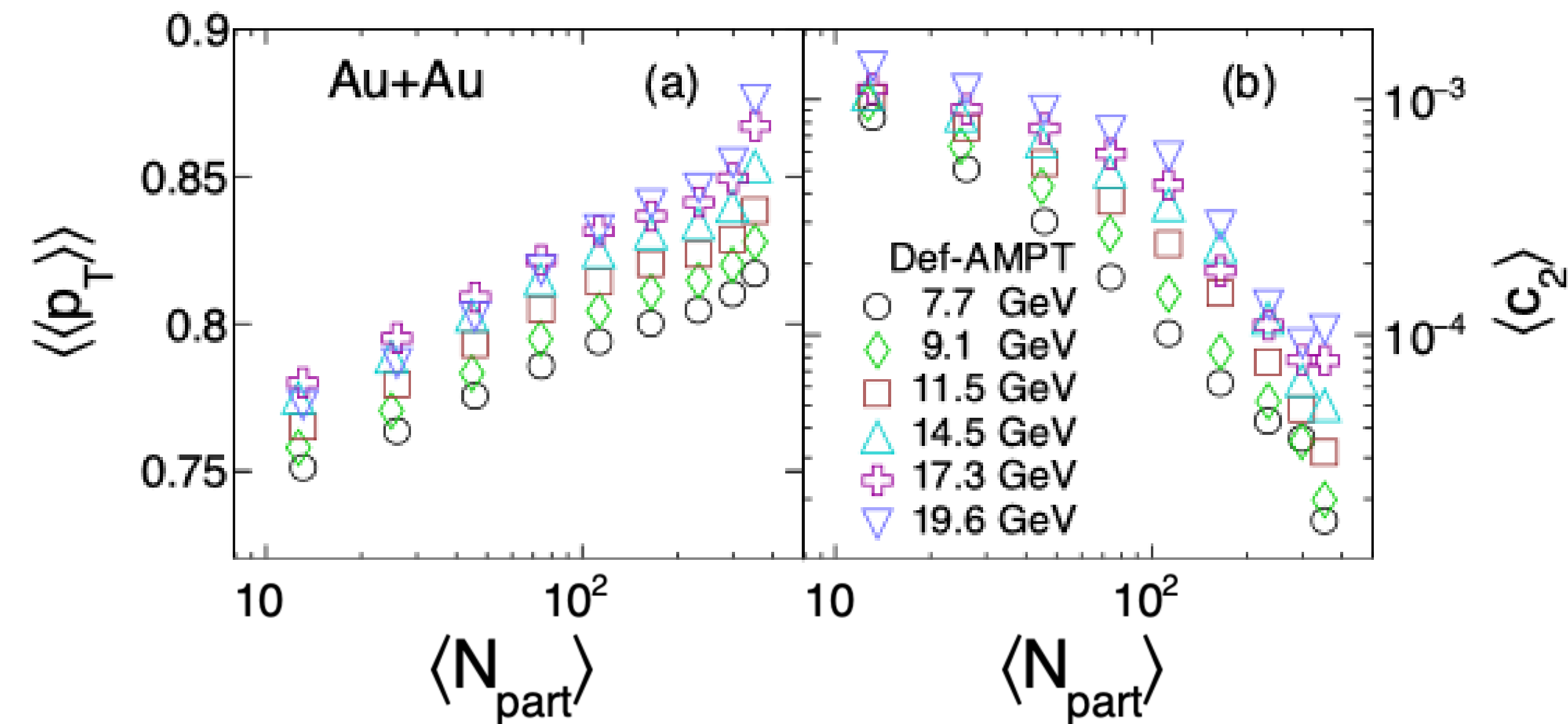


The peripheral collisions (60-70%) compared to central collisions (0-5%) :

- the $\langle p_T \rangle$ distributions exhibit greater variances, indicating enhanced fluctuations.
- the $\langle p_T \rangle$ distributions show a significant rightward tail, suggesting positive skewness.

Second-order p_T cumulants

Variance($\langle c_2 \rangle$)



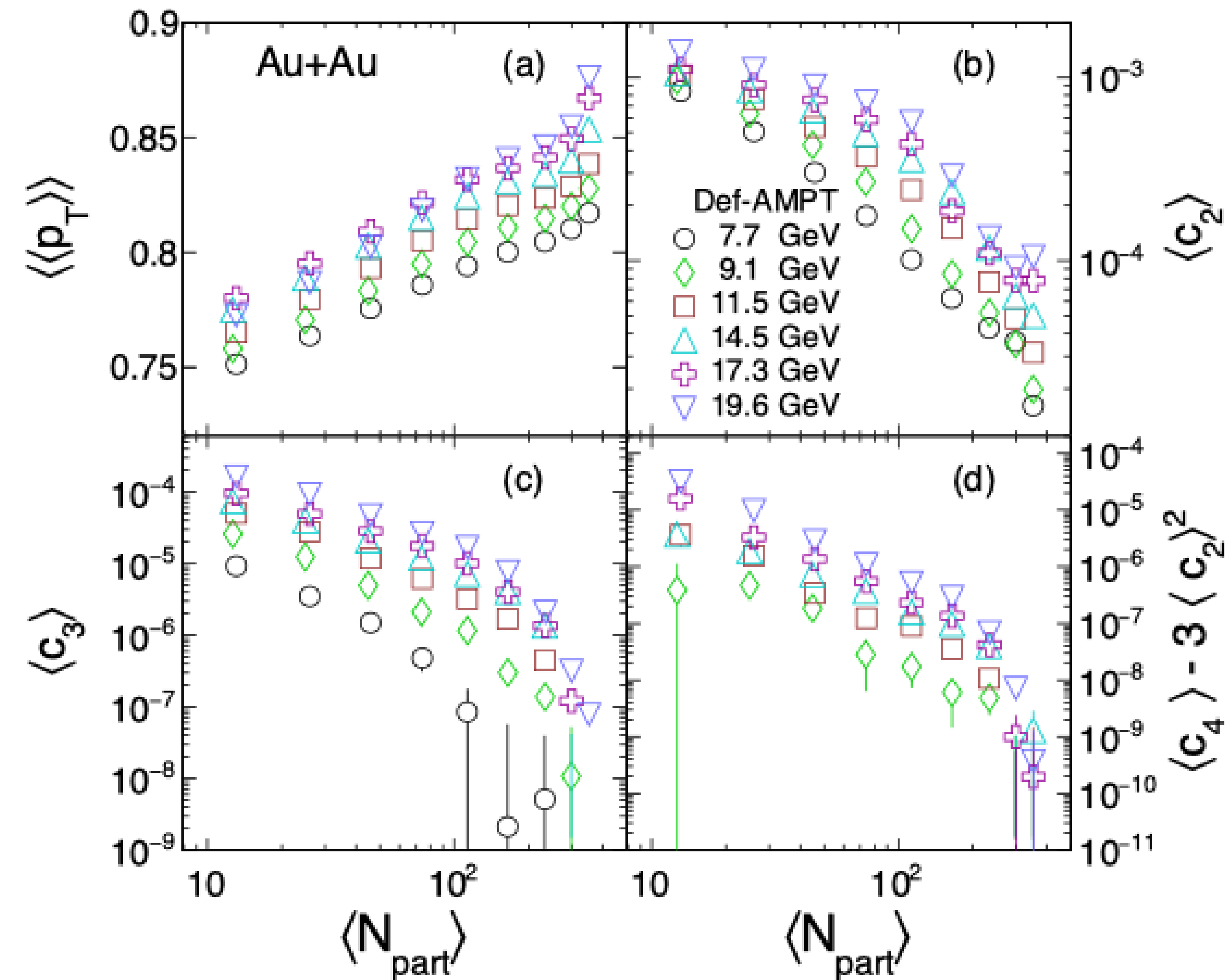
$\langle c_2 \rangle$ (cent.) vs energy

An inverse dependence on centrality is conserved across all energies.

May from a reduction in particle-pair correlations if they are dominated by particles originating from the same NN collisions.

Higher-order p_T cumulants

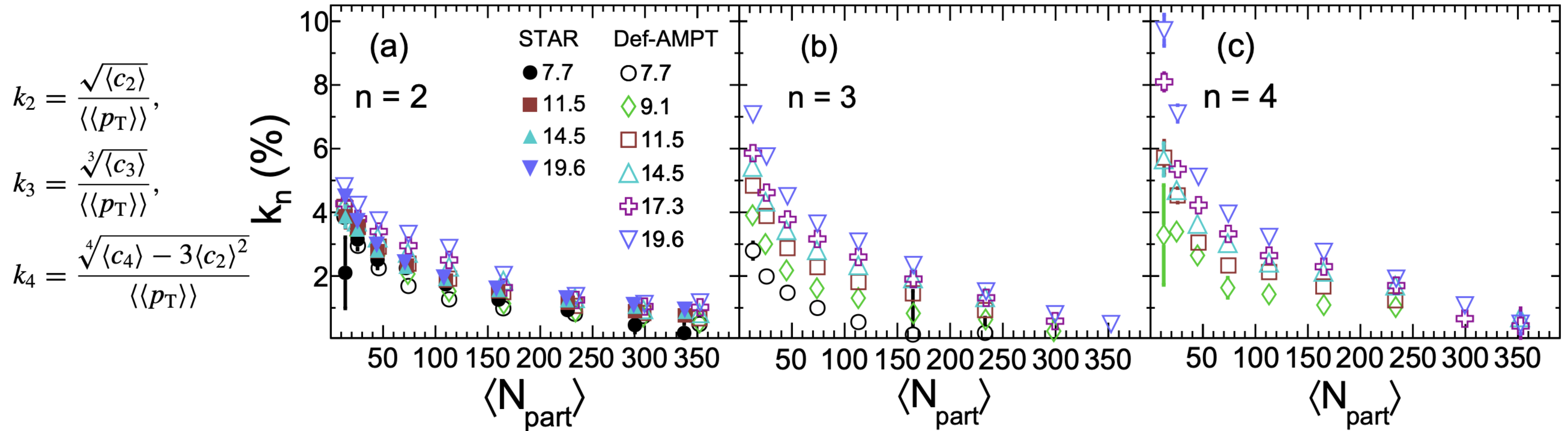
Skewness ($\langle c_3 \rangle$), Kurtosis ($\langle c_4 \rangle - 3\langle c_2 \rangle^2$)



Share significant dependence on cent. and energy:
their magnitudes decrease by more than one order of
magnitude with increasing centrality classes.

Scaled p_T cumulants

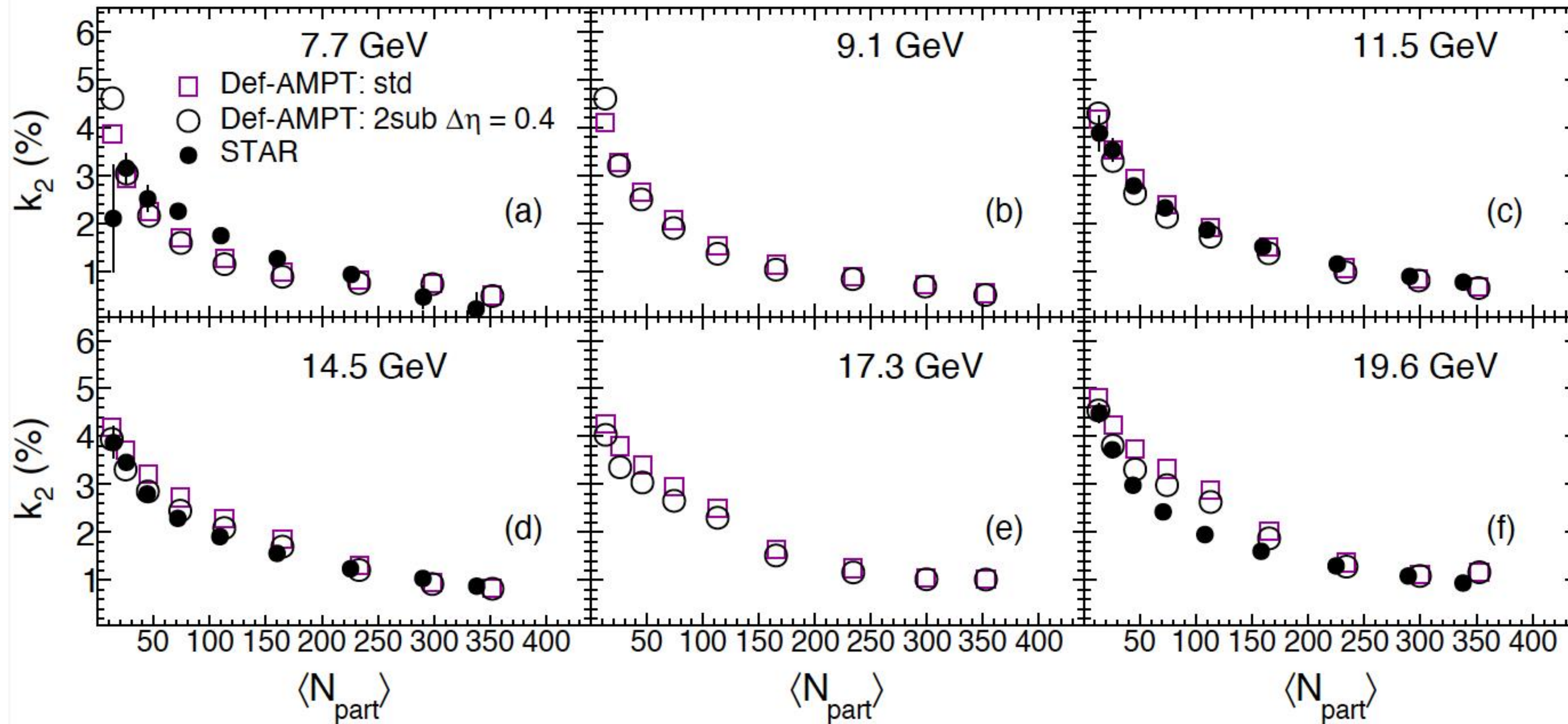
To mitigate the influence on the cumulants in $\langle\langle p_T \rangle\rangle$ with centrality or incident energy.



- k_2 qualitatively produce the trends of STAR experiment.
- k_n ($n=2,3,4$) exhibits significant centrality and energy dependence;
- k_n ($n=2,3,4$) exhibit an approximate power-law behavior.

Scaled p_T cumulants: the short-range correlations

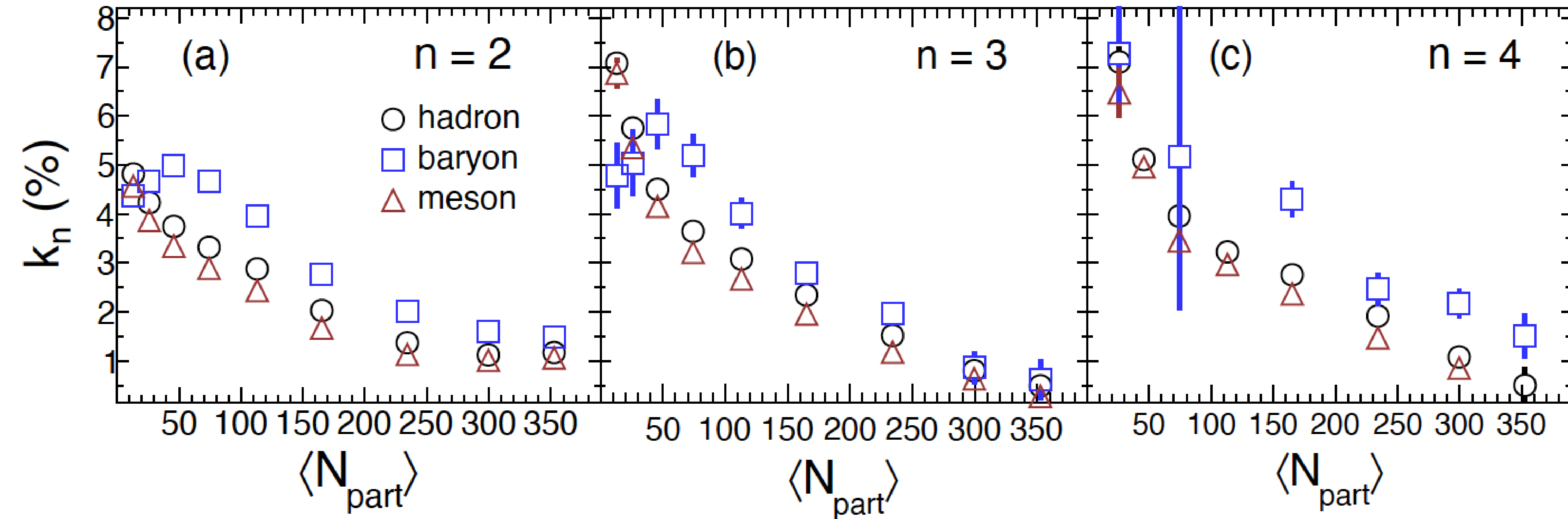
The scaled p_T cumulants by a two-subevent method: to minimize the impact of short-range correlations.



All values from two-subevent method are slightly suppressed compared to the standard method.

Scaled p_T cumulants: baryon Vs meson

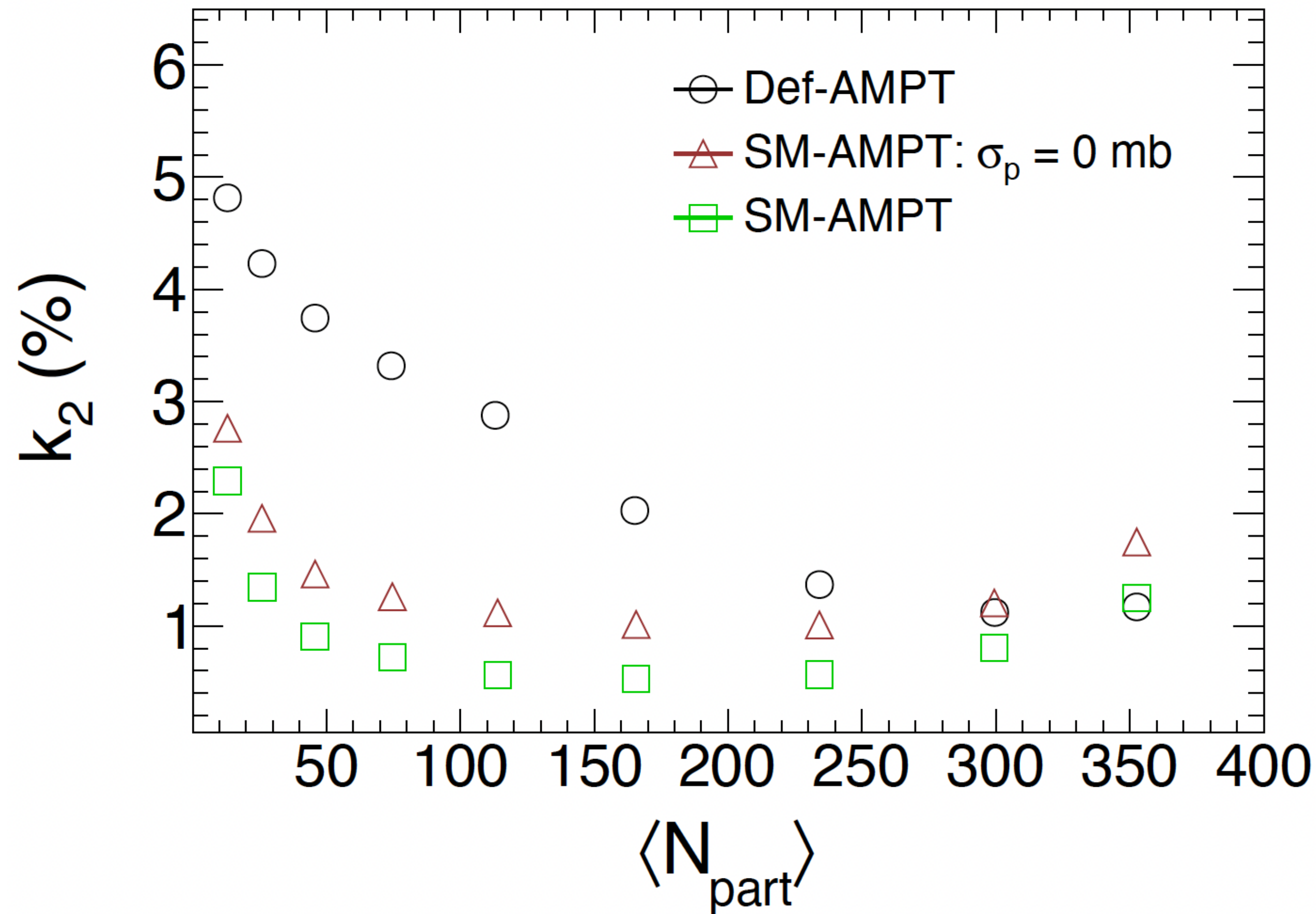
To further explore the radial flow mechanism.



Baryons exhibit more pronounced fluctuations in all scaled variance, skewness and kurtosis compared to mesons.

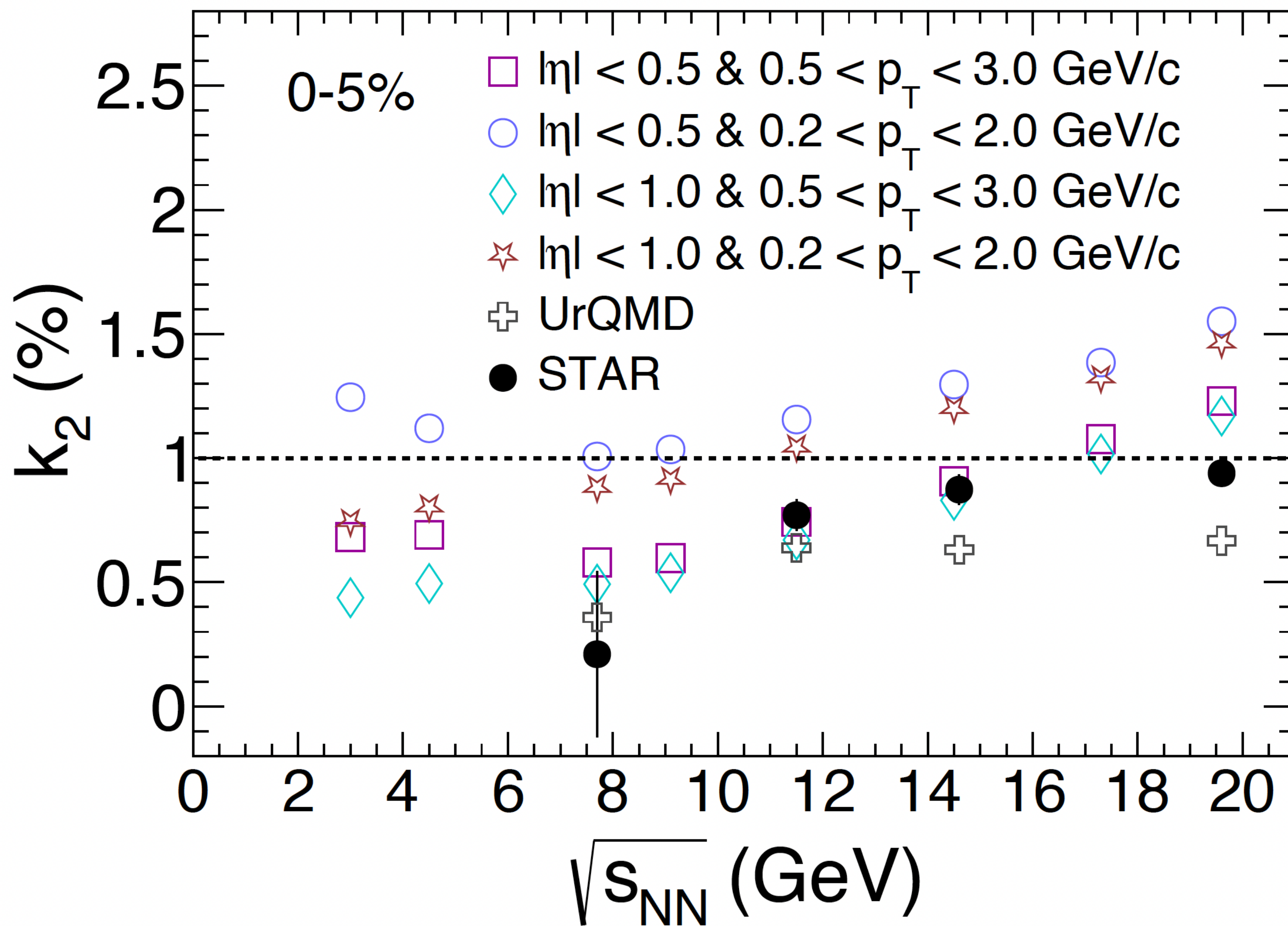
→ this behavior might be attributed to the effects of radial flow.

Scaled p_T cumulants: SM- Vs default-AMPT



- The default configuration exhibits more pronounced $\langle p_T \rangle$ fluctuations than SM version with and without partonic interactions.
→ likely different hadronization mechanisms.
- Partonic evolution slightly suppress the EbE fluctuations.

Scaled p_T cumulants: vs energy



- k_2 is more sensitive to p_T variations than η acceptance.
- $0.5 < p_T < 3.0$ GeV: quantitatively consistent with STAR measurements.
- extend to 3.0 GeV, k_2 exhibits a significant abnormal increase, indicating an enhancement of dynamical correlations at lower collision energies.

- A comprehensive systematic study of higher-order dynamical p_T cumulants up to fourth order.
- Higher-order p_T cumulants with and without normalization exhibit a strong dependence on centrality across 3.0-19.6 GeV.
- Our finding provides variable references for the experimental measurements.
- The inner mechanism from acceptance or decorrelation need further to be explored.

Thank you !!!

Backup

$$c_n = \frac{\sum_{i_1 \neq \dots \neq i_n} w_{i_1} \cdots w_{i_n} (p_{T,i_1} - \langle\langle p_T \rangle\rangle) \cdots (p_{T,i_n} - \langle\langle p_T \rangle\rangle)}{\sum_{i_1 \neq \dots \neq i_n} w_{i_1} \cdots w_{i_n}}$$

$$p_{mk} = \sum_i w_i^k p_i^m / \sum_i w_i^k, \quad \tau_k = \frac{\sum_i w_i^{k+1}}{(\sum_i w_i)^{k+1}}$$

$$\begin{aligned} c_2 &= \frac{\bar{p}_{11}^2 - \tau_1 \bar{p}_{22}}{1 - \tau_1}, \\ c_3 &= \frac{\bar{p}_{11}^3 - 3\tau_1 \bar{p}_{22} \bar{p}_{11} + 2\tau_2 \bar{p}_{33}}{1 - 3\tau_1 + 2\tau_2}, \\ c_4 &= \frac{\bar{p}_{11}^4 - 6\tau_1 \bar{p}_{22} \bar{p}_{11}^2 + 3\tau_1^2 \bar{p}_{22}^2 + 8\tau_2 \bar{p}_{33} \bar{p}_{11} - 6\tau_3 \bar{p}_{44}}{1 - 6\tau_1 + 3\tau_1^2 + 8\tau_2 - 6\tau_3} \end{aligned} \quad \langle p_T \rangle \text{ cumulant}$$

S. Bhatta, C, Zhang, and J. Jia, PRC 105, 024904 (2022)