# Shear and Bulk Viscosities of Gluon Plasma across the Transition Temperature from Lattice QCD

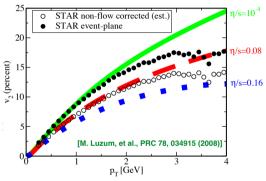
### Cheng Zhang (张成)<sup>1</sup> collaborators: Heng-Tong Ding<sup>1</sup>, Hai-Tao Shu<sup>1</sup>

<sup>1</sup>Institute of Particle Physics Central China Normal University



第二十届全国中高能核物理大会 Apr. 24-28, 2025, 上海

## Introduction



• inputs for tranport/hydro models

 η/s quantifies the dissipation processes in the hydrodynamics.

G. Denicol et al., PRC 80, 064901 (2008)

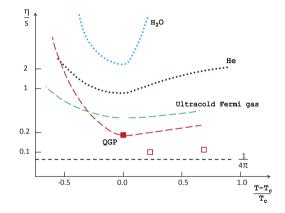
 Small η/s suggested by phenomenological interpretation of experimental data.

K. H. Ackermann et al., (STAR), PRL 86, 402 (2001)

 Extracting η/s from experiments needs accurate inputs: initial condition, EoS ...

U. Heinz et al., Annu. Rev. Nucl. Part. Sci. 63, 123 (2013)

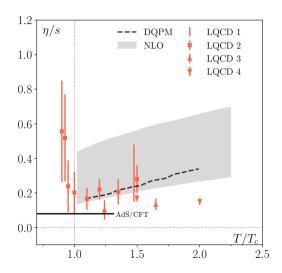
## Introduction



S.Cremonini et al., JHEP 1208 (2012) 167

- η/s of QGP is sensitive to the phase transitions.
- Determinations of viscosities require theoretical inputs.

# Determinations of viscosities from theory



1.5  $T_c \Rightarrow$  this work:  $0.76 \leq T/T_c \leq 2.25$ 

# Theoretical framework

$$\begin{split} G_{\rm shear}(\tau,\tau_{\rm F}) &= \frac{1}{10} \int \mathrm{d}^3 x \, \left\langle \pi_{ij}(0,\vec{0},\tau_{\rm F}) \, \pi_{ij}(\tau,\vec{x},\tau_{\rm F}) \right\rangle \\ \hline \mathbf{Renormalization} & a \to 0, \tau_{\rm F} \to 0 \\ \hline G(\tau) &= \int_0^\infty \frac{\mathrm{d}\omega}{\pi} \, \frac{\cosh[\omega(1/2T-\tau)]}{\sinh(\omega/2T)} \, \rho(\omega,T) \\ \hline \mathbf{Inverse \ problem} \\ \hline \mathbf{Spectra \ reconstruction} \, \rho(\omega,T) \\ \hline \mathbf{Kubo \ formula} \\ \hline \eta(T) &= \lim_{\omega \to 0} \, \frac{\rho_{\rm shear}(\omega,T)}{\omega} \end{split}$$

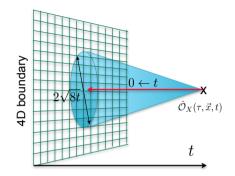
#### Key Points:

- Lattice computation of EMT correlators.
- Spectra reconstruction from the correlators.

#### Challenges:

- Severe UV fluctuations in the correlators.
- Theoretical uncertainty in the spectra reconstruction.

Noise reduction technique: gradient flow



Flow equation :

$$\begin{aligned} \frac{\mathrm{d}B_{\mu}(x,\tau_{\mathsf{F}})}{\mathrm{d}\tau_{\mathsf{F}}} &= D_{\nu}\,\mathcal{G}_{\nu\mu}(x,\tau_{\mathsf{F}})\\ \mathcal{B}_{\nu}(x,\tau_{\mathsf{F}}=0) &= \mathcal{A}_{\nu}(x) \end{aligned}$$

LO solution:

$$B_
u(x, au_{F}) \propto m{exp} \left(rac{-(x-y)^2}{\sqrt{8 au_{F}}^2/2}
ight) B_
u(y)$$

M. Lüscher, JHEP 08, 071 (2010)

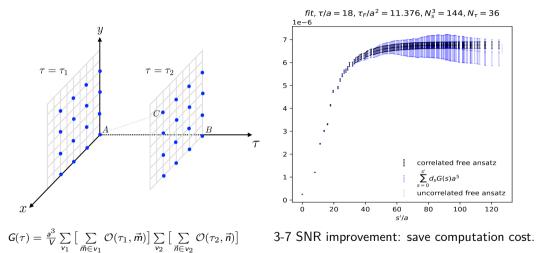
Smearing radius:  $\sqrt{8\tau_{\rm F}}$ . Advantage:

- ► The UV fluctuations strongly suppressed.
- Provide renormalization framework for EMT:

 $T_{\mu\nu}(\tau_F, \mathbf{x}) = c_1(\tau_F) U_{\mu\nu}(\tau_F, \mathbf{x}) + 4c_2(\tau_F) \delta_{\mu\nu} E(\tau_F, \mathbf{x})$ 

• Operator Product Expansion of  $G(\tau, \tau_{\rm F})$  in  $\tau_{\rm F}/\tau^2$ .

## Noise reduction technique: blocking fit



L. Altenkort et al. PRD 105, 094505 (2022)

## Lattice setup

Pure SU(3) Yang-Mills gauge theory:

# NSD

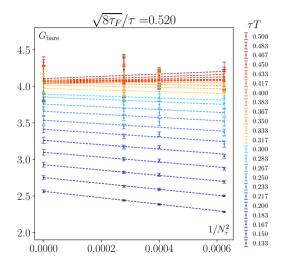
$T/T_c$	0.76			0.9			1.125			1.267			1.5			1.9			2.25		
$N_{\sigma}$	96	120	144	96	120	144	96	120	144	96	120	144	96	120	144	96	120	144	96	120	144
$N_{ au}$	40	50	60	40	50	60	32	40	48	24	30	36	24	30	36	16	20	24	16	20	24
#Conf.	5000		5000		5000		5000		5000		5000		5000								

#### Lattice spacing:

				7.3874		
<i>a</i> (fm)	0.02068	0.01746	0.01654	0.01397	0.01379	0.01164

H.-T. Ding, H.-T. Shu and CZ, work in progress

## Continuum extrapolation



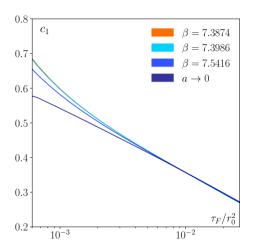
 $a \rightarrow 0$  at  $0.9 T_c$  in the shear channel.

The joint fit Ansatz:

$$G_{\text{bare}}(N_{\tau}) = G_{\text{bare}}^{\tau T}(a = 0) + \left(b + m_1 \cdot \tau T + \frac{m_2}{\tau T}\right) / N_{\tau}^2$$
$$G_{\text{bare}} = \frac{G^{\text{t.l.}}(\tau T, \tau_{\text{F}})}{G^{\text{norm}}(\tau T)}$$

The Ansatz describes the data well.

## Renormalization

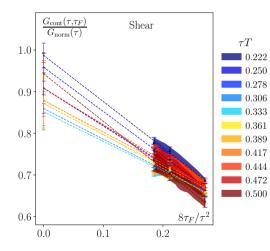


 $c_1 \& c_2$ : renormalization constants matching Gradient Flow scheme to  $\overline{\text{MS}}$  scheme.

$$\begin{aligned} T_{\mu\nu} \left( \tau_{\rm F}, x \right) &= c_1 \left( \tau_{\rm F} \right) U_{\mu\nu} \left( \tau_{\rm F}, x \right) + 4 c_2 \left( \tau_{\rm F} \right) \delta_{\mu\nu} E(\tau_{\rm F}, x) \\ c_1(\tau_{\rm F}) &= \frac{1}{g^2(\mu)} \sum_{n=0}^2 k_1^{(n)} (L(\mu, \tau_{\rm F})) \Big[ \frac{g^2(\mu)}{(4\pi)^2} \Big]^n \end{aligned}$$

To control lattice spacing effects in  $c_1$ :  $\tau_{\rm F}/r_0^2 \geq 9\times 10^{-4}$ 

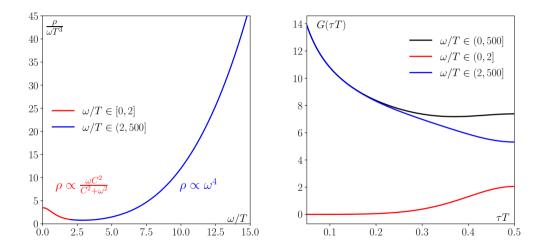
## Flow time extrapolation



$$\begin{split} \tau_{\mathsf{F}} &\rightarrow 0 \text{ extrapolation Ansatz:} \\ \mathcal{G}(\tau_{\mathsf{F}}/\tau^2, \tau T) = \mathcal{G}_{\tau_{\mathsf{F}}=0}^{\tau \, \mathsf{T}} + \left(b + m_1 \cdot \tau \, T + \frac{m_2}{\tau \, T}\right) \cdot \tau_{\mathsf{F}}/\tau^2 \\ \mathsf{Flow time window:} \\ &\sqrt{8\tau_{\mathsf{F}}}/\tau \in [0.43, \, 0.52] \\ & \text{L. Altenkort et al. PRD 103, 114513 (2021)} \\ \mathsf{Avoid over smearing:} \\ &\sqrt{8\tau_{\mathsf{F}}} > \sqrt{2}a \end{split}$$

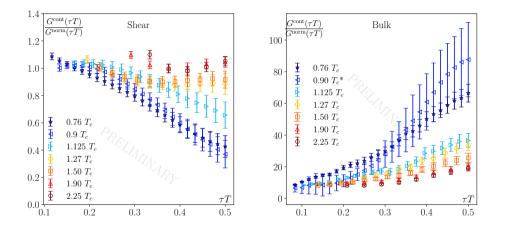
 $\tau_{\rm F} \rightarrow 0$  at  $1.5\,T_c.$ 

## Illustration of sensitivity of correlators to the transport peak



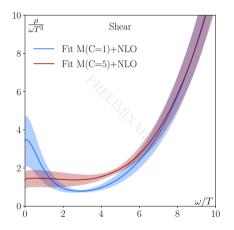
 $G(\tau T)$  at  $\tau T \sim 0.5$  are more sensitive to the transport peak.

## Normalized correlators in the continuum limit



Clear temperature dependencies for both channels. Negative slope for low temperatures in the shear channel.

## Reconstructed spectral functions in the shear channel at 1.5 $T_c$



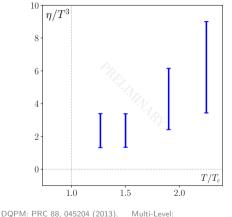
Spectral function at 1.5  $T_c$ .

$$\begin{split} G(\tau) &= \int_0^\infty \frac{d\omega}{\pi} \frac{\cosh[\omega(1/2T-\tau)]}{\sinh(\omega/2T)} \rho(\omega,T) \\ \frac{\rho(\omega)}{\omega T^3} &= \frac{A}{T^3} \frac{C^2}{C^2 + (\omega/T)^2} + B \frac{\rho_{\text{pert}}(\omega)}{\omega T^3} \\ \rho_{\text{pert}}(\omega) \propto (\omega/T)^4 \end{split}$$

Y. Zhu et al. JHEP 03, 002 (2013) (shear)M. Laine et al., JHEP 09, 084(2011 (bulk)

C = 1, sharp peak, long-lived excitation. C = 5, broad peak, short-lived excitation.

## Temperature dependencies of shear viscosity

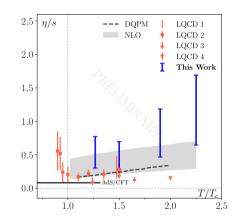


 DQFWI: FRC 06, 049204 (2013).
 Multit-Level:

 NLO: JHEP 03, 179 (2018).
 LQCD2: JHEP 04, 101 (2017).

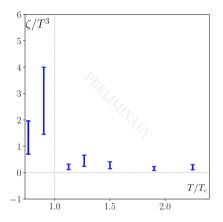
 Gradient Flow:
 LQCD3: PRD 76, 101701 (2007).

 LQCD1: PRD 108, 014503 (2023).
 LQCD4: PRD 98, 014512 (2018).

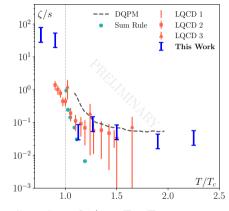


Mild increase with temperature in  $\eta/s$  at  $T \ge 1.27 T_c$ .  $\eta/s$  agrees with LQCD1 & NLO at  $T \ge 1.27 T_c$ .

## Temperature dependencies of bulk viscosity



DQPM: PRC 88, 045204 (2013) Sum Rule: JHEP 09, 093 (2008) LQCD1 (GF): PRD 108, 014503 (2023). a = 0.0117 fm LQCD2 (ML): PRD 98, 054515 (2018).  $a \ge 0.0253$  fm LQCD3 (ML): PRL 100, 162001 (2008).  $a \ge 0.0475$  fm



Smaller values of  $\zeta/s$  at  $T > T_c$ .  $\zeta/s$  agrees with LQCD1 & LQCD2 at  $T > T_c$ .

## Results

- ► Large and fine lattices are generated to extract the viscosities.
- High-precision EMT correlators are obtained via gradient flow and blocking methods.
- Temperature dependencies of η/s and ζ/s are investigated in SU(3) across the phase transition region.
- ▶  $\eta/s$  increases mildly with temperature at  $T \ge 1.27 T_c$ .
- ►  $\zeta/s(T < T_c) \gg \zeta/s(T > T_c)$ , and  $\zeta/s$  is most flat at  $T \ge 1.13T_c$ .
- ▶ The full QCD investigation (including dynamical quarks) is progressing.