



利用国产格点组态计算高精度 ρ 介子质量和宽度

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arXiv: 2502.03700

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Outline

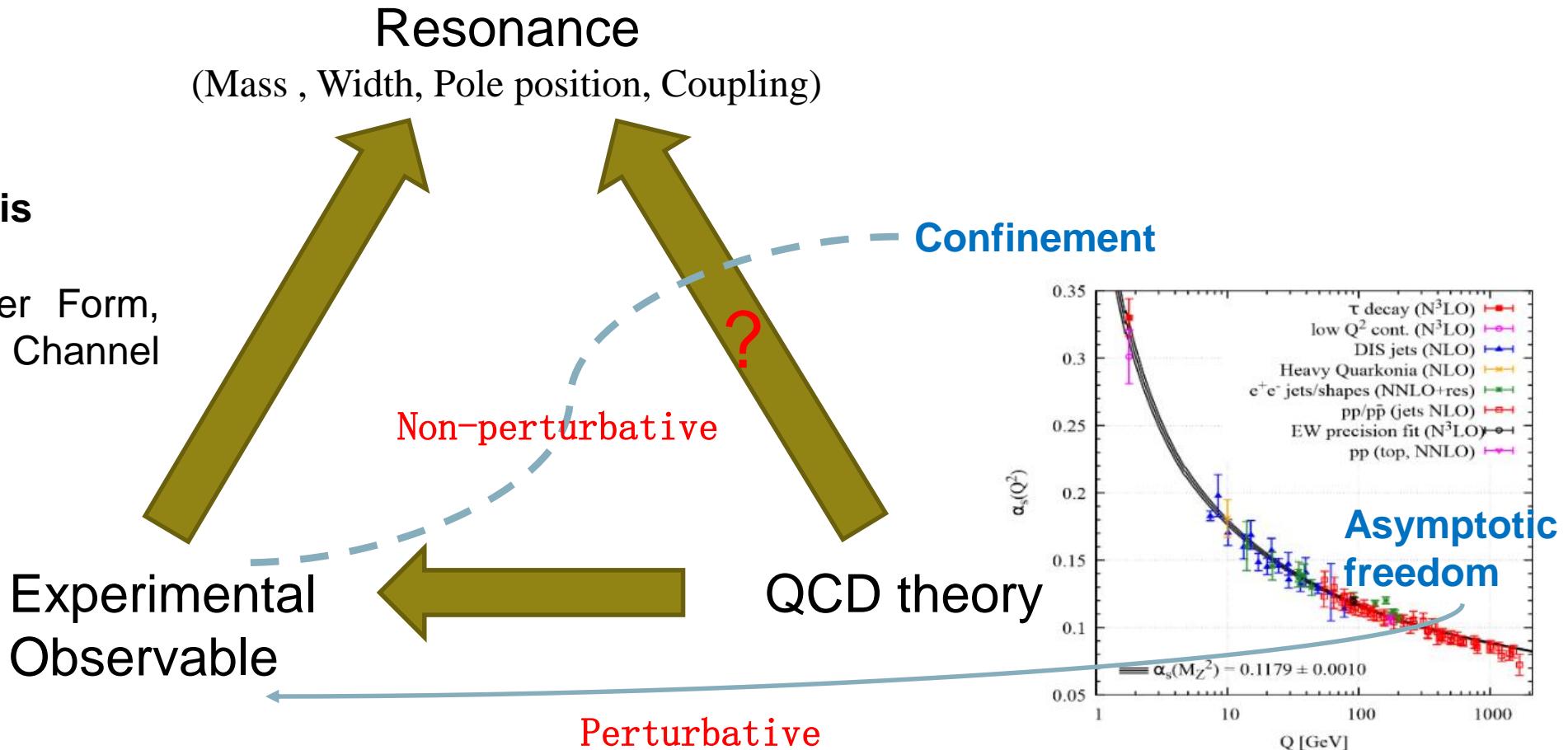
- Motivation
- Lattice QCD introduction
- ρ meson spectra on LQCD
- Summary

Motivation

Partial Wave Analysis

With various Model:

Such as Breit-Wigner Form,
Flatte Form, Coupled Channel
Form, and so on



Motivation: Spectrum & Scattering

- Spectrum: Table of the energy of the physical states, can reflect the inner structure information.
- Scattering: The line shapes of Cross section or T-matrix show the properties of the states and interaction.
- These observables can constrain the models.

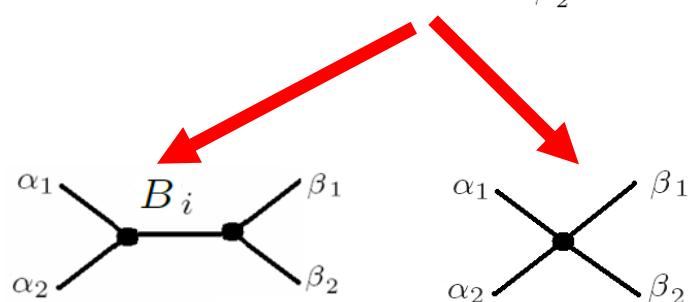
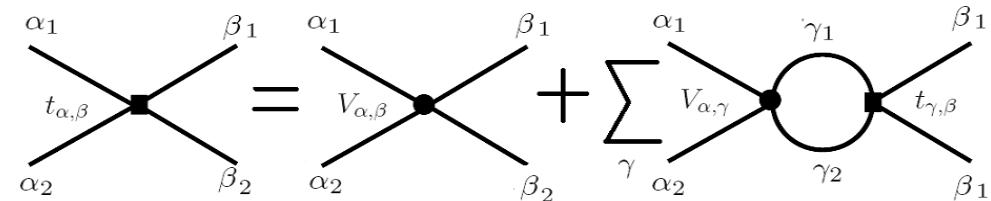


Powerful Parameters, Powerful Model

Motivation: Scattering

Lattice Spectrum !!!
The dimension of m_π

$$t_{\alpha,\beta}(k_\alpha, k_\beta, E) = V_{\alpha,\beta}(k_\alpha, k_\beta) + \sum_\gamma \int k_\gamma^2 dk_\gamma \frac{V_{\alpha,\gamma}(k_\alpha, k_\gamma) t_{\gamma,\beta}(k_\gamma, k_\beta, E)}{E - \sqrt{m_{\gamma 1}^2 + k_\gamma^2} - \sqrt{m_{\gamma 2}^2 + k_\gamma^2} + i\epsilon}$$



$$\frac{g_{i,\alpha}^*}{E - m_i} \frac{1}{E - m_i} g_{i,\beta}$$

$$v_{\alpha,\beta}$$

$$G_{i,\alpha} = g_{i,\alpha} + \sum_\beta \int \frac{k^2 dk \tilde{t}_{\alpha\beta} g_{j,\beta}}{E - \sqrt{m_{\beta 1}^2 + k^2} - \sqrt{m_{\beta 2}^2 + k^2} + i\epsilon}$$

$$\tilde{t}_{\alpha\beta} = v_{\alpha\beta} + \sum_\gamma \int \frac{k_\gamma^2 dk_\gamma v_{\alpha\gamma} \tilde{t}_{\gamma\beta}}{E - \sqrt{m_{\gamma 1}^2 + k_\gamma^2} - \sqrt{m_{\gamma 2}^2 + k_\gamma^2} + i\epsilon}$$

- $T^{-1}(E) \sim E - m_i - \Sigma(E)$
 $= E - m_i - \text{Re}(\Sigma(E)) - i\text{Im}(\Sigma(E))$
- $T^{-1}(E) \sim E - m_i - \text{Re}(\Sigma(E)) + i\Gamma(E)/2$
- $m_{phys} = m_i + \text{Re}(\Sigma(E))$
 $= m_0 + am_\pi^2 + \dots + \text{Re}(\Sigma(E))$

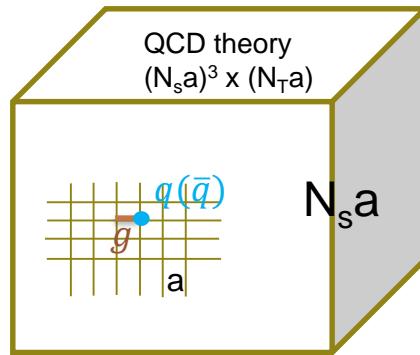
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Lattice QCD

1. QCD theory: on a box in the Euclid four space



2. $a \rightarrow$ UV cutoff, $N_s a \rightarrow$ Infrared truncation

3. Lattice QCD \rightarrow a model of statistical physics.

$$\langle O \rangle = \int D\phi O[\phi] P[\phi] \quad P[\phi] = \frac{1}{Z} e^{-S[\phi]} \quad Z = \int D\phi e^{-S[\phi]}$$

ϕ : field quantity, $S[\phi]$: Action, $O[\phi]$: physical quantity

4. Monte Carlo method

5. Three steps for Lattice QCD to real world

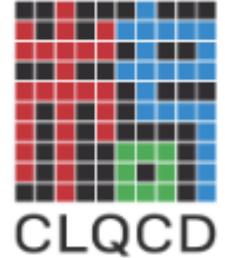
a, Configuration

b, Measurement $\sum_{(\vec{y}-\vec{x}) \in Z^3} e^{i\vec{p} \cdot (\vec{y}-\vec{x})} \langle T(\psi(t; \vec{y}), \psi^\dagger(t; \vec{y})) \rangle \sim \sum_{\Gamma, i} Z_i^\Gamma e^{-E_i^\Gamma t}$

c, Transformation

Build configuration for different pion mass

Cost!



Operator in quark and gluon level

Calculate Correlation function

Extract energy levels

HEFT

Luescher

Translate energy levels to observable

Model/EFT

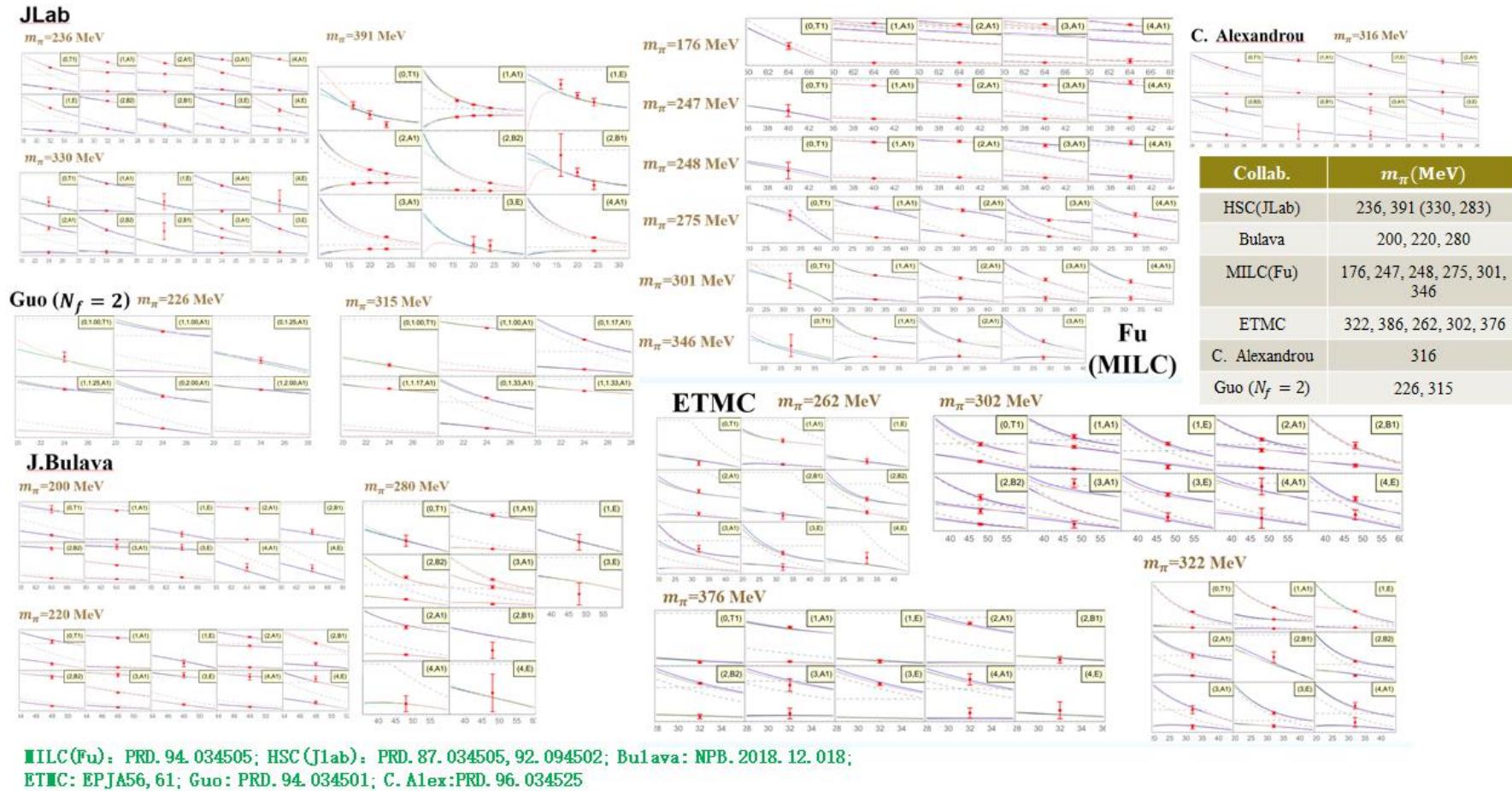
Extract physics parameters from observables



Motivation: π mass dependence

- The ρ is the best one!
- The largest lattice data of ρ .
- The lowest state can decay through strong interaction.
- We almost clearly understand ρ .

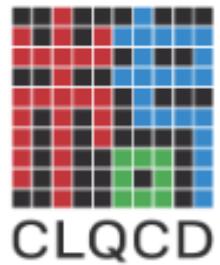
Kang Yu, D. Leinweber,
 A. Thomas, J.j. Wu
 Phys.Rev.D 109 (2024) 3, 034505



Motivation: π mass dependence

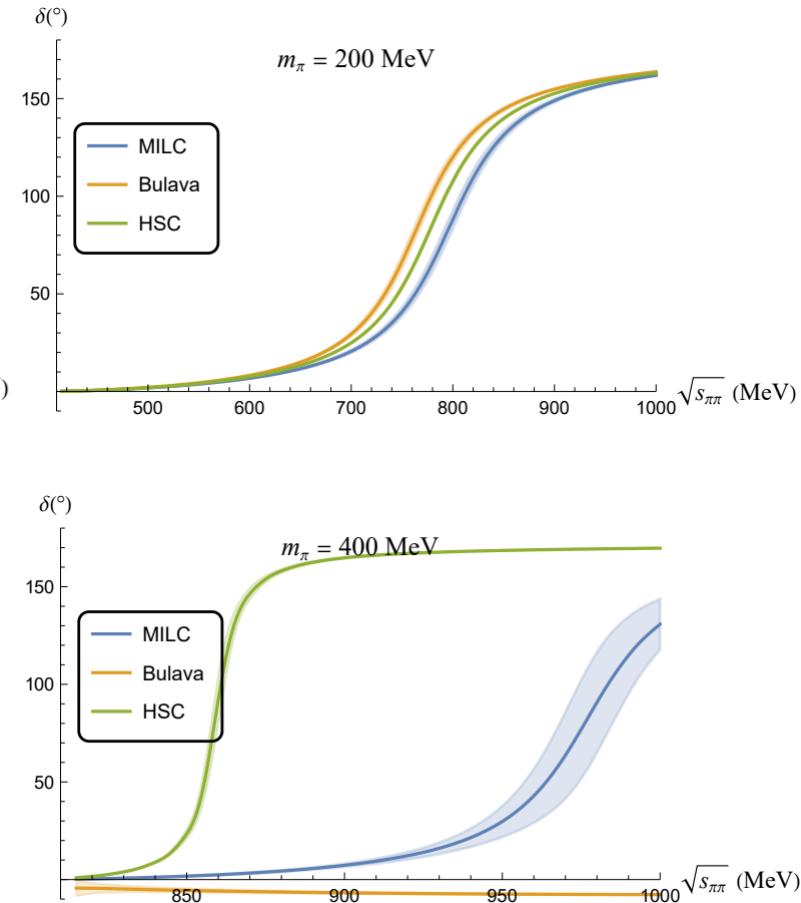
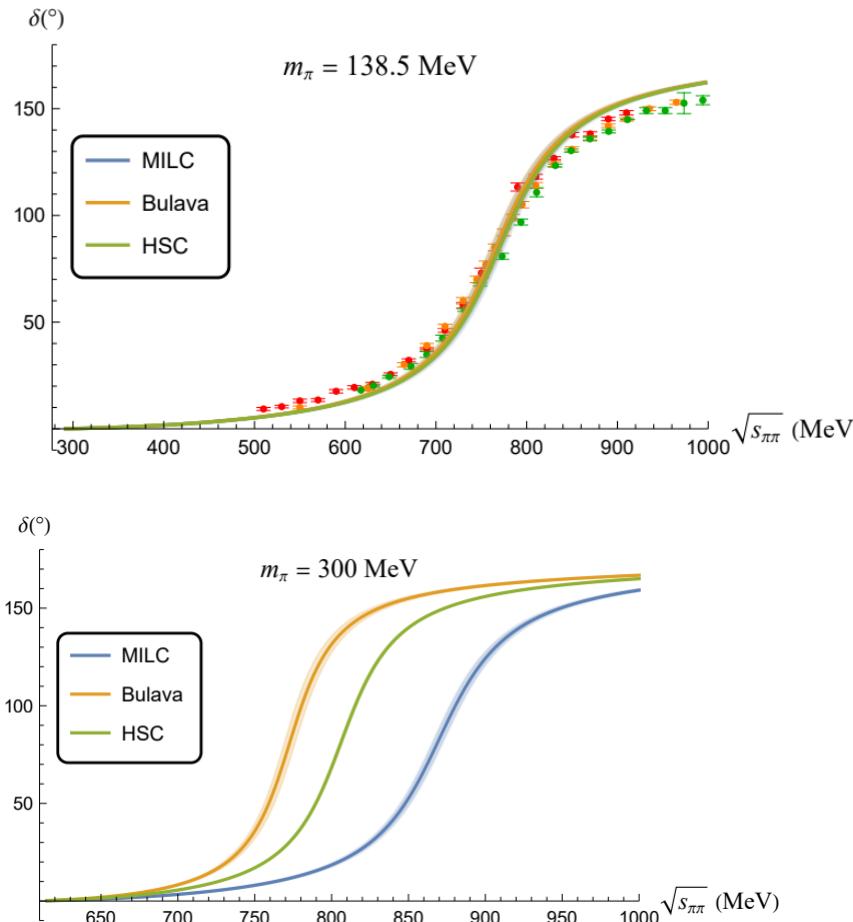
- The ρ is the best one!
- The largest lattice data of ρ .
- The lowest state can decay through strong interaction.
- We almost clearly understand ρ .

However



Lattice spacing effect,
Fermion action effect...

We need a systematic Lattice
data group to study ρ !



Kang Yu, D. Leinweber, A. Thomas, J.j. Wu

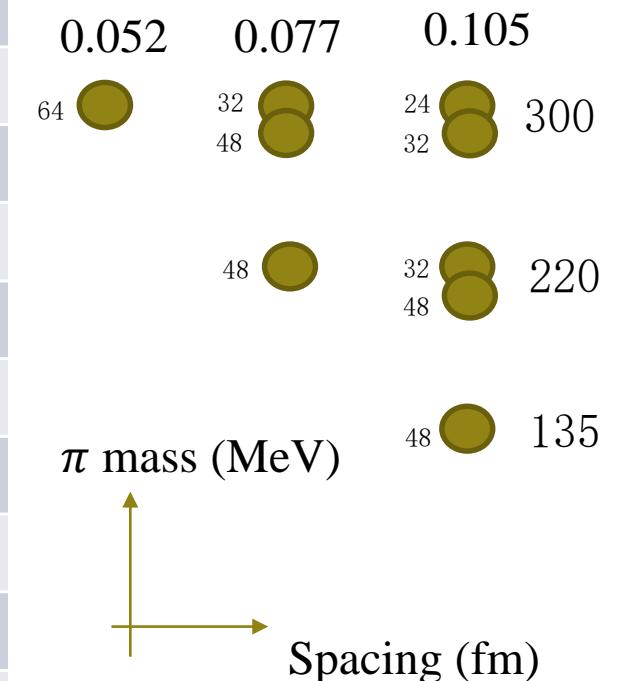


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ρ meson spectra on LQCD -- Configuration

| Name | Volume | Spacing | β | m_π/MeV | $m_\pi L$ |
|--------|-------------------|------------|---------|-------------|-----------|
| C24P34 | $24^3 \times 64$ | 0.10530 fm | 6.20 | 340 | 4.38 |
| C24P29 | $24^3 \times 72$ | | | 292 | 3.75 |
| C32P29 | $32^3 \times 64$ | | | 292 | 5.01 |
| C32P23 | $32^3 \times 64$ | | | 228 | 3.91 |
| C48P23 | $48^3 \times 96$ | | | 225 | 5.79 |
| C48P14 | $48^3 \times 96$ | | | 135 | 3.56 |
| F32P30 | $32^3 \times 96$ | 0.07746 fm | 6.41 | 303 | 3.81 |
| F48P30 | $48^3 \times 96$ | | | 303 | 5.72 |
| F32P21 | $32^3 \times 64$ | | | 210 | 2.67 |
| F48P21 | $48^3 \times 96$ | | | 207 | 3.91 |
| H48P32 | $48^3 \times 144$ | 0.05187 fm | 6.72 | 321 | 4.06 |
| H64P32 | $64^3 \times 128$ | | | 321 | 5.41 |



CLQCD, PRD109
(2024) 5, 054507

ρ meson spectra on LQCD -- Operator

| $[000]T_1^-$ | $[001]A_1$ | $[001]E_2$ | $[011]A_1$ |
|-----------------------------|----------------------------|----------------------------|---------------------------|
| $\rho_{[000]}$ | $\rho_{[001]}$ | $\rho_{[001]}$ | $\rho_{[011]}$ |
| $\pi_{[001]}\pi_{[00-1]}$ | $\pi_{[000]}\pi_{[001]}$ | $\pi_{[0-10]}\pi_{[011]}$ | $\pi_{[000]}\pi_{[011]}$ |
| $\pi_{[011]}\pi_{[0-1-1]}$ | $\pi_{[0-10]}\pi_{[011]}$ | $\pi_{[0-1-1]}\pi_{[111]}$ | $\pi_{[-100]}\pi_{[111]}$ |
| $\pi_{[111]}\pi_{[-1-1-1]}$ | $\pi_{[-1-10]}\pi_{[111]}$ | | $\pi_{[01-1]}\pi_{[002]}$ |
| | $\pi_{[00-1]}\pi_{[002]}$ | | |

| $[011]B_1$ | $[011]B_2$ | $[111]A_1$ | $[002]A_1$ |
|---------------------------|---------------------------|---------------------------|--------------------------|
| $\rho_{[011]}$ | $\rho_{[011]}$ | $\rho_{[111]}$ | $\rho_{[002]}$ |
| $\pi_{[010]}\pi_{[001]}$ | $\pi_{[-100]}\pi_{[111]}$ | $\pi_{[000]}\pi_{[111]}$ | $\pi_{[000]}\pi_{[002]}$ |
| $\pi_{[110]}\pi_{[-101]}$ | $\pi_{[110]}\pi_{[-101]}$ | $\pi_{[100]}\pi_{[011]}$ | |
| $\pi_{[0-11]}\pi_{[002]}$ | | $\pi_{[200]}\pi_{[-111]}$ | |

$$\rho^0 = \frac{1}{\sqrt{2}}(\bar{u}\Gamma^P u - \bar{d}\Gamma^P d)$$

$$\rho^{0\dagger} = \frac{1}{\sqrt{2}}(\bar{u}\bar{\Gamma}^P u - \bar{d}\bar{\Gamma}^P d)$$

$$\pi^+ = \bar{d}\Gamma^k u, \quad \pi^- = \bar{u}\Gamma^k d$$

$$\pi\pi = \frac{1}{\sqrt{2}}(\pi^+(k_1)\pi^-(k_2) - \pi^-(k_1)\pi^+(k_2))$$

$$= \frac{1}{\sqrt{2}}(\bar{d}\Gamma^{k_1} u \bar{u}\Gamma^{k_2} d - \bar{u}\Gamma^{k_1} d \bar{d}\Gamma^{k_2} u)$$

$$(\pi\pi)^\dagger = \frac{1}{\sqrt{2}}(\bar{d}\bar{\Gamma}^{k_2} u \bar{u}\bar{\Gamma}^{k_1} d - \bar{u}\bar{\Gamma}^{k_2} d \bar{d}\bar{\Gamma}^{k_1} u)$$



ρ meson spectra on LQCD -- correlation function

$$C_{ij} = \langle 0 | \mathcal{O}_i \mathcal{O}_j^\dagger | 0 \rangle = \sum_n Z_i^n Z_j^{n*} e^{-E_n t}$$

Computed based on wick contraction and quark propagator calculated based on the configuration.

To parameterize the correlation function !

Energy level in the finite volume

- Solve the generalized eigenvalue problem(G EVP)

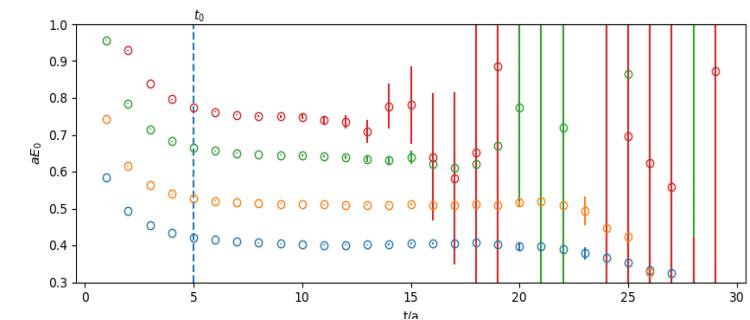
$$\mathcal{G}_{ij}(t_0 + dt) u_j^\alpha = e^{-m_\alpha dt} \mathcal{G}_{ij}(t_0) u_j^\alpha,$$

$$\mathcal{G}^\alpha(t) = v_i^\alpha \mathcal{G}_{ij}(t) u_j^\alpha,$$

$$[\mathcal{G}^{-1}(t_0) \mathcal{G}(t_0 + dt)]_{ij} u_j^\alpha = c^\alpha u_i^\alpha$$

$$E^\alpha(t) = \frac{1}{n} \log \frac{\mathcal{G}^\alpha(t)}{\mathcal{G}^\alpha(t+n)},$$

$$v_i^\alpha [\mathcal{G}(t_0 + dt) \mathcal{G}^{-1}(t_0)]_{ij} = c^\alpha v_j^\alpha,$$



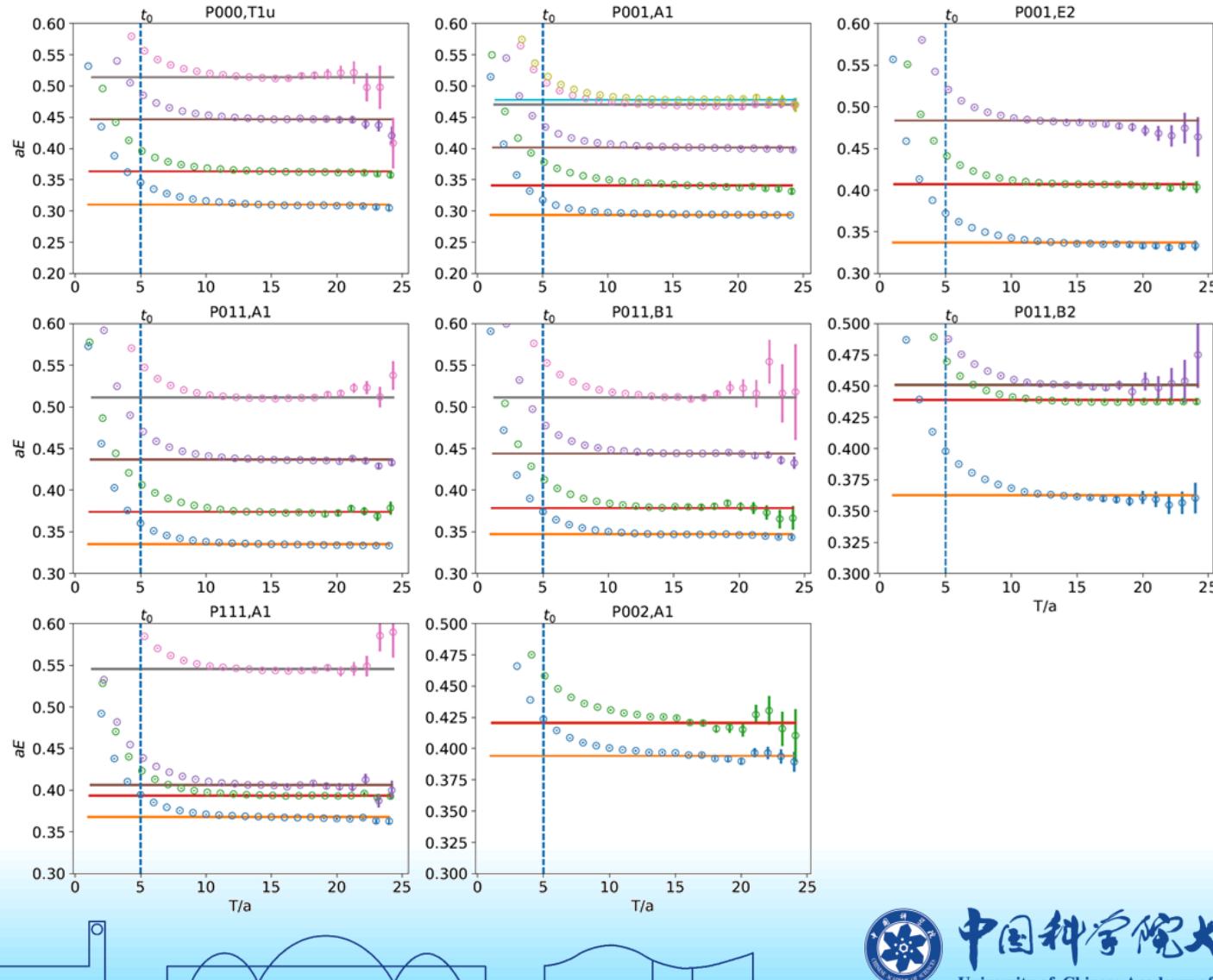
Kiratidis, et al., PRD 91, 094509 (2015)

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ρ meson spectra on LQCD -- Spectra



Configuration: F48P30

There are 8 different spectra with different total momentum and irreps, including following irreducible representations.:

$$p^2 = 0 : T_1^- ,$$

$$p^2 = 1 : A_1 \text{ and } E_2 ,$$

$$p^2 = 2 : A_1, B_1 \text{ and } B_2 ,$$

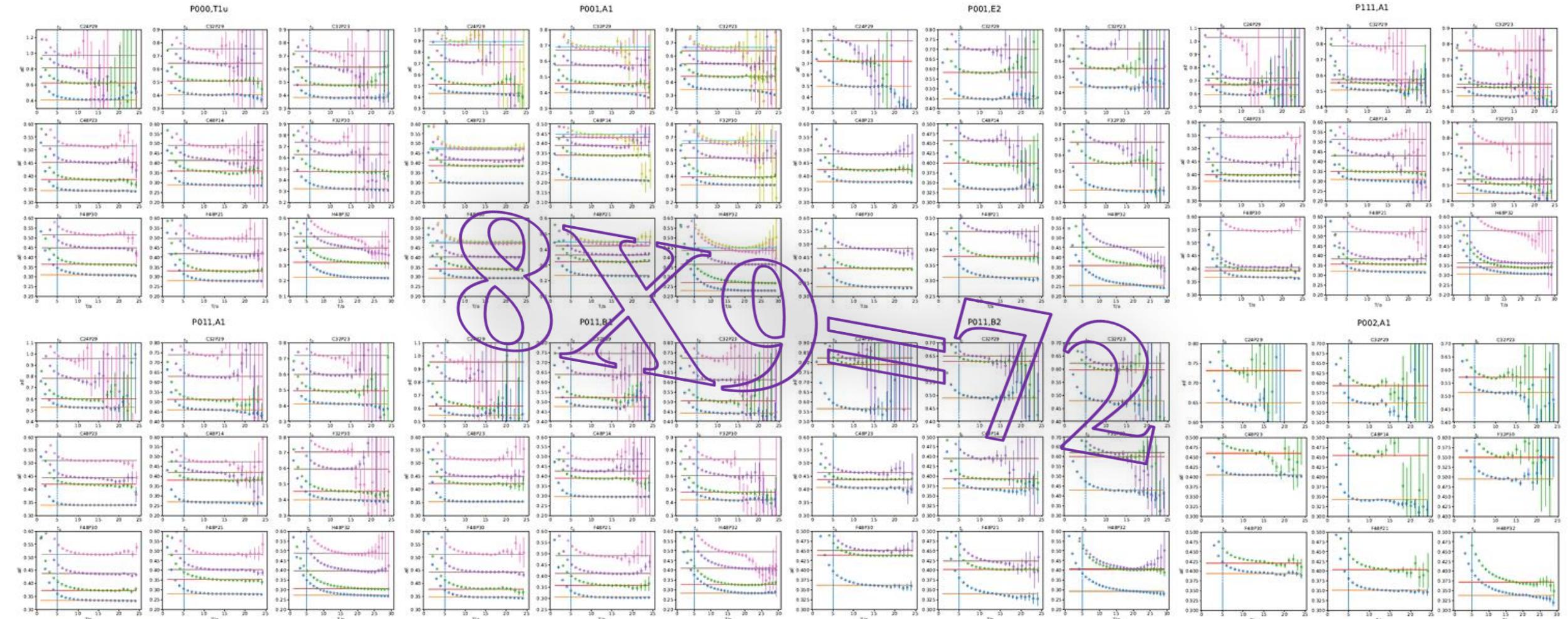
$$p^2 = 3 : A_1 ,$$

$$p^2 = 4 : A_1 .$$



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ρ meson spectra on LQCD -- Spectra



ρ meson spectra on LQCD -- Spectra

C32P29/C48P29

$m_\pi \sim 300$ MeV

C48P23/C32P23

~ 230 MeV

C48P14

~ 135 MeV

F48P30/F32P30

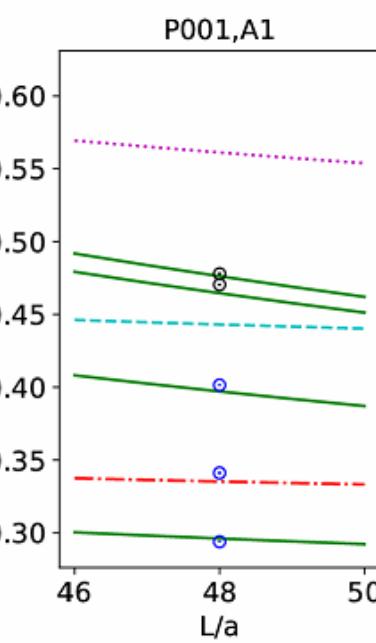
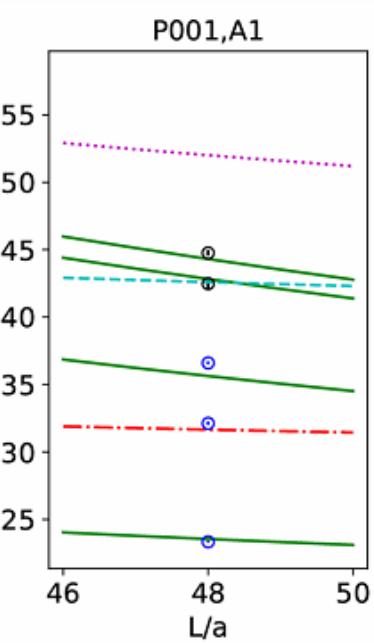
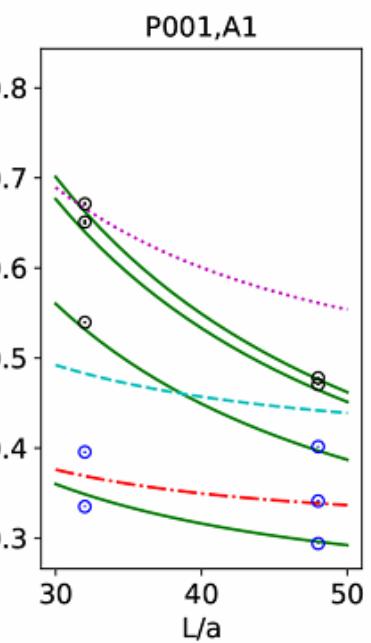
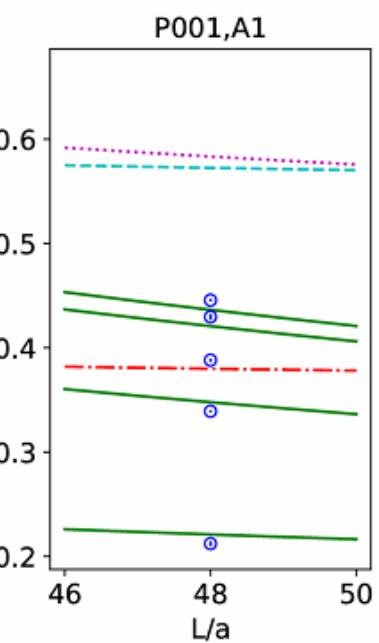
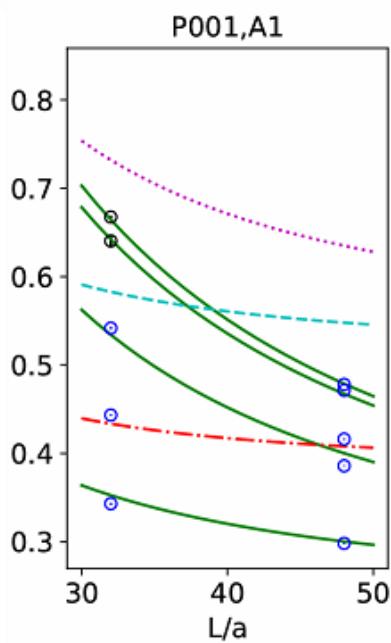
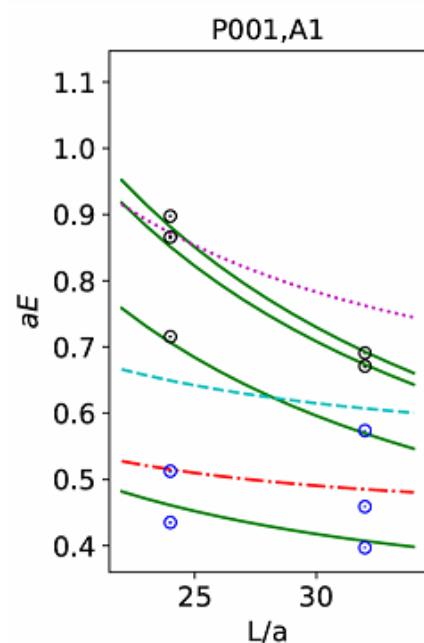
~ 300 MeV

F48P21

~ 220 MeV

H48P32

~ 300 MeV



$a \sim 0.105$ fm



$a \sim 0.077$ fm

$a \sim 0.052$ fm

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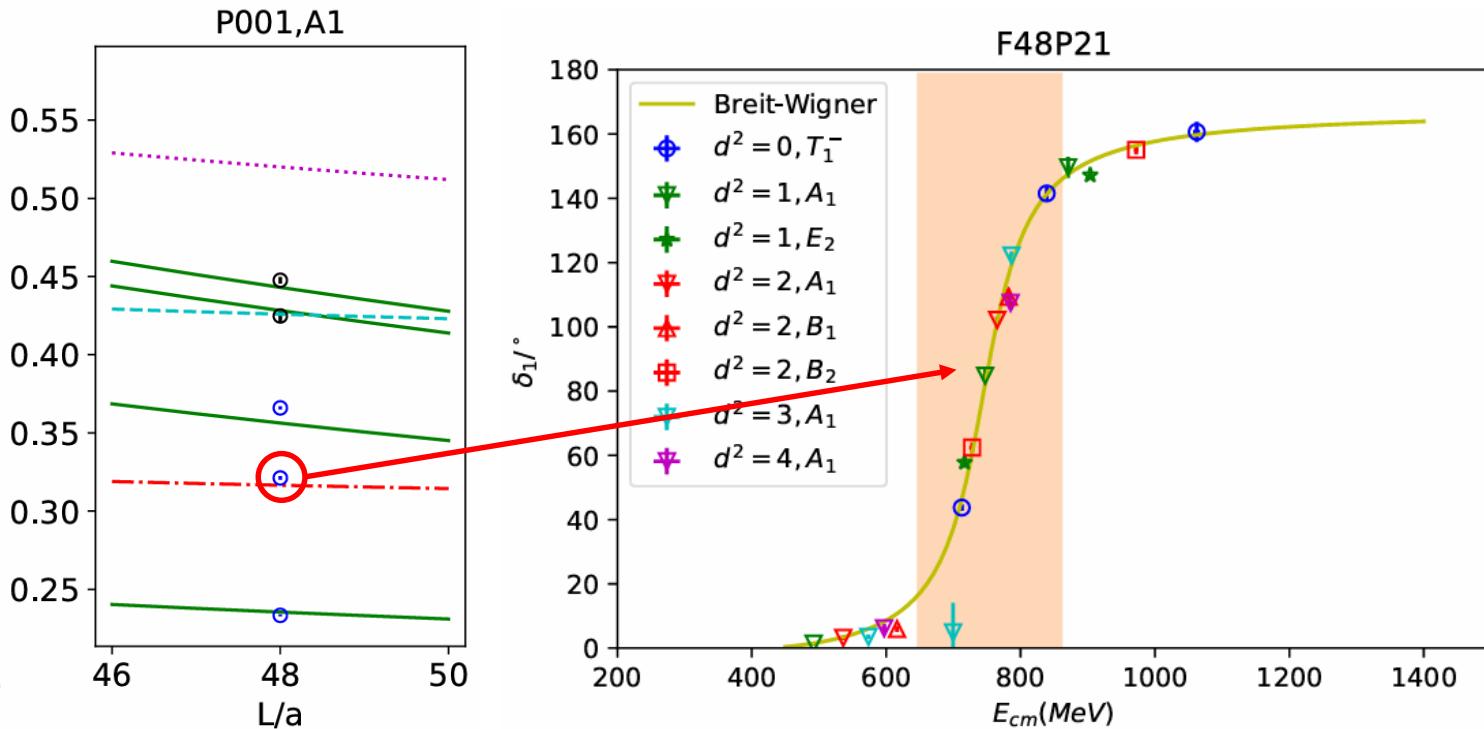


ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

| d^2 | Λ, μ | $M_{11,11}^{(\vec{P}, \Lambda, \mu)}$ |
|-------|----------------|---|
| 0 | $T_1^-, 2$ | $w_{0,0}$ |
| 1 | $A_1, 1$ | $w_{0,0} + 2w_{2,0}$ |
| 1 | $E_2, 1$ | $w_{0,0} - w_{2,0}$ |
| 1 | $E_2, 2$ | $w_{0,0} - w_{2,0}$ |
| 2 | $A_1, 1$ | $w_{0,0} + \frac{1}{2}w_{2,0} + i\sqrt{6}w_{2,1} - \frac{\sqrt{6}}{2}w_{2,2}$ |
| 2 | $B_1, 1$ | $w_{0,0} + \frac{1}{2}w_{2,0} - i\sqrt{6}w_{2,1} - \frac{\sqrt{6}}{2}w_{2,2}$ |
| 2 | $B_2, 1$ | $w_{0,0} - w_{2,0} + \sqrt{6}w_{2,2}$ |
| 3 | $A_1, 1$ | $w_{0,0} - i2\sqrt{6}w_{2,2}$ |

$$\det(M_{ln,l'n'}(k) - \delta_{ll'}\delta_{nn'}\cot(\delta_l)) = 0, \quad \delta_1 = \arccot M_{11,1}^{(\vec{P}, \Lambda, \mu)}$$

$$w_{j,s} = \frac{\mathcal{Z}_{js}(1, q^2)}{\pi^{3/2}\sqrt{2j+1}\gamma q^{j+1}}, \quad q = \frac{kL}{2\pi}, \quad \delta_{l \geq 3} \text{ are negligible}$$

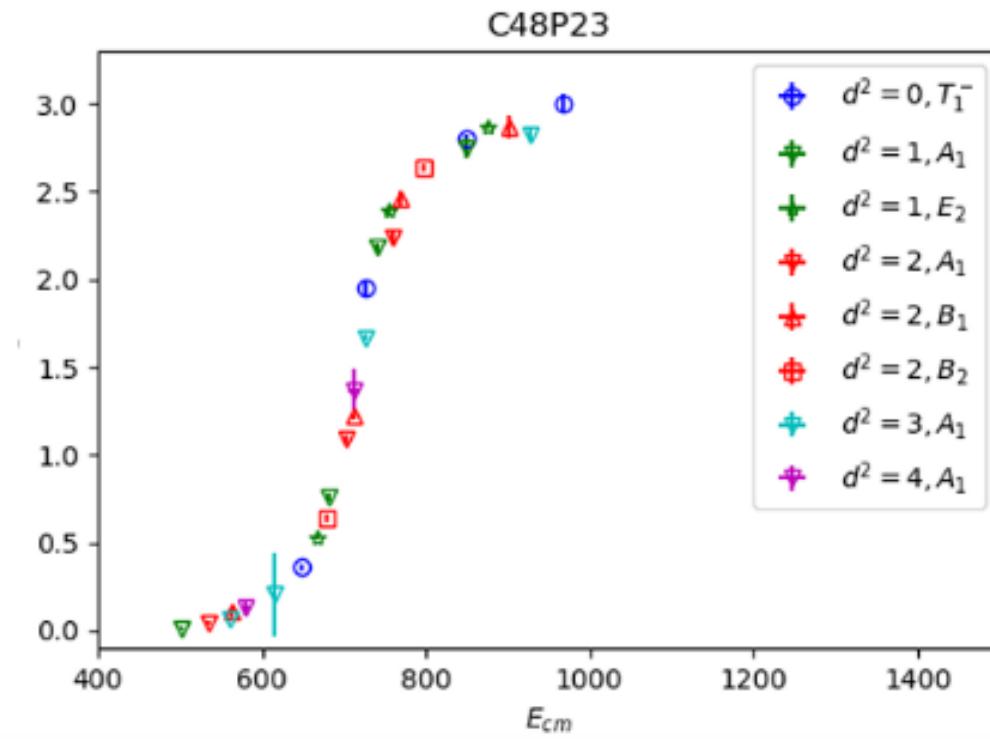
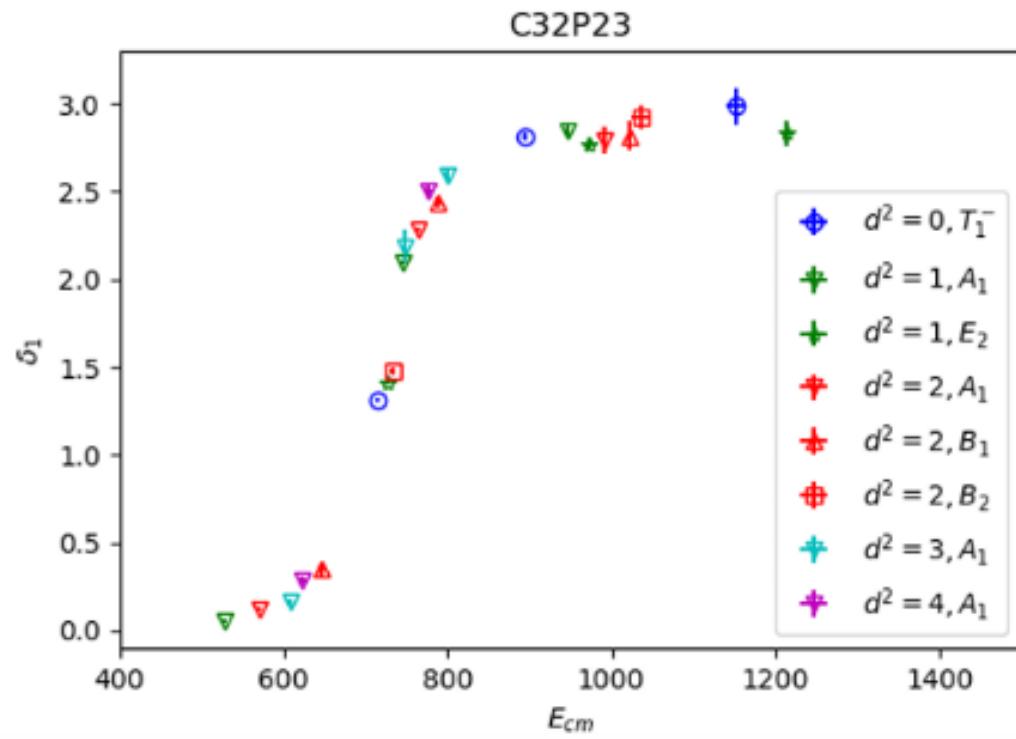


$$E_{cm} = \sqrt{E_{001}^2 - 4\pi^2/L^2} = \sqrt{(0.34/a)^2 - 4\pi^2/(48a)^2} = 771 \text{ MeV}$$



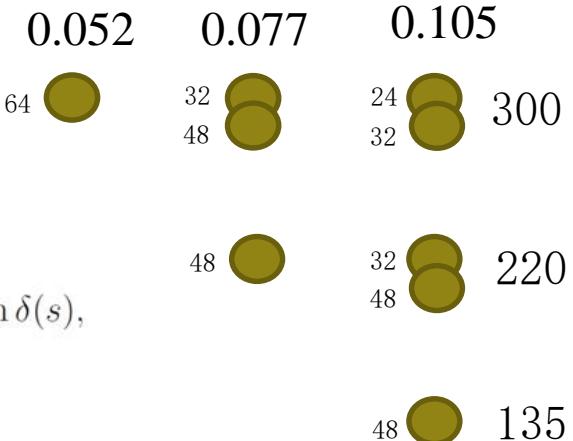
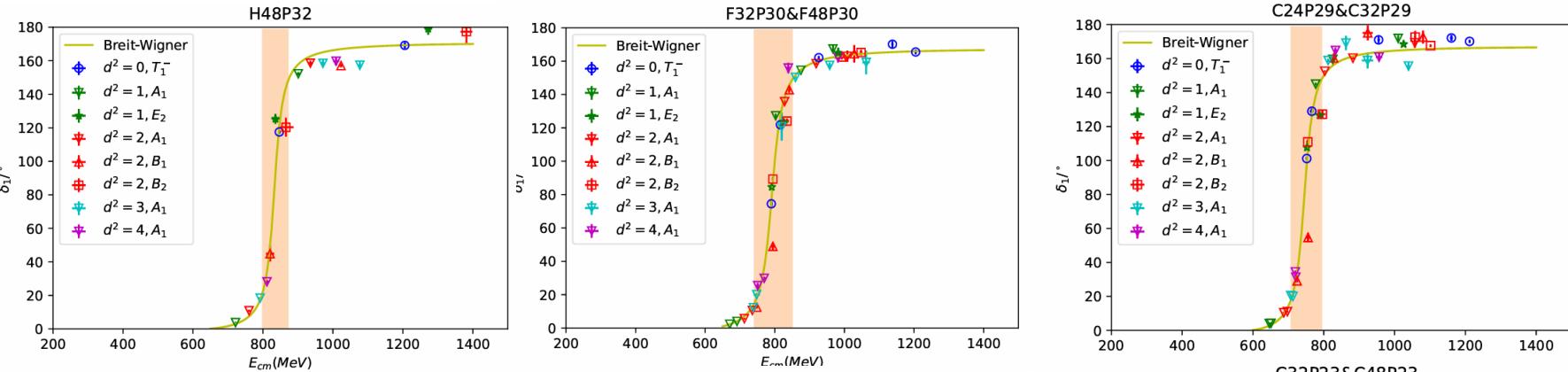
ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

| Name | Volume | Spacing | β | m_π/MeV | $m_\pi L$ |
|--------|------------------|--------------|---------|-------------|-----------|
| C32P23 | $32^3 \times 64$ | $0.10530 fm$ | 6.20 | 228 | 3.91 |
| C48P23 | $48^3 \times 96$ | | | 225 | 5.79 |



ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

| | m_ρ (MeV) | $g_{\rho\pi\pi}$ |
|---------------|----------------|------------------|
| C24P29/C32P29 | 749.0(6.0) | 5.859(79) |
| C32P23/C48P23 | 723.9(5.6) | 5.827(14) |
| C48P14 | 714.0(34.5) | 6.300(440) |
| F32P30/F48P30 | 795.5(4.2) | 5.775(58) |
| F48P21 | 754.2(3.0) | 5.954(38) |
| H48P32 | 834.8(7.4) | 4.916(45) |

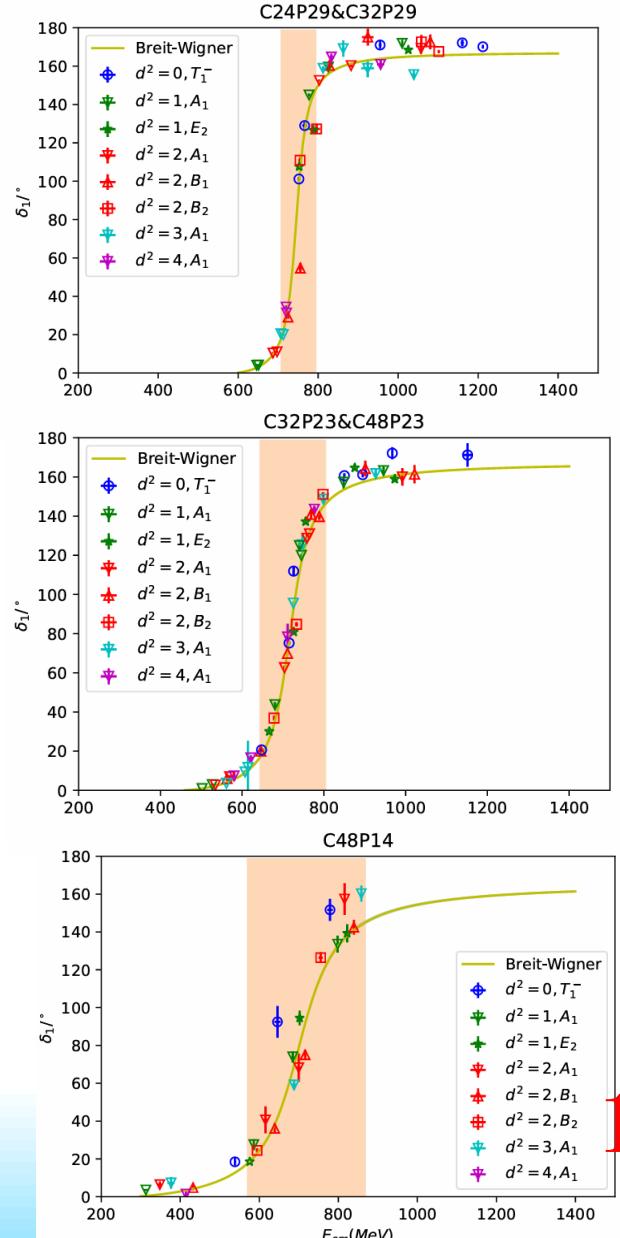
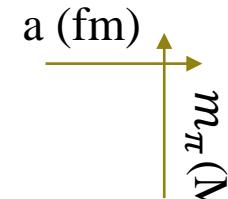


$$a_1 = \frac{-\sqrt{s}\Gamma(s)}{s - m_\rho^2 + i\sqrt{s}\Gamma(s)} = e^{i\delta(s)} \sin \delta(s),$$

$$\Gamma(s) = \frac{p^{*3}}{s} \frac{g_{\rho\pi\pi}^2}{6\pi}.$$



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ρ meson spectra on LQCD -- M & Γ of ρ meson

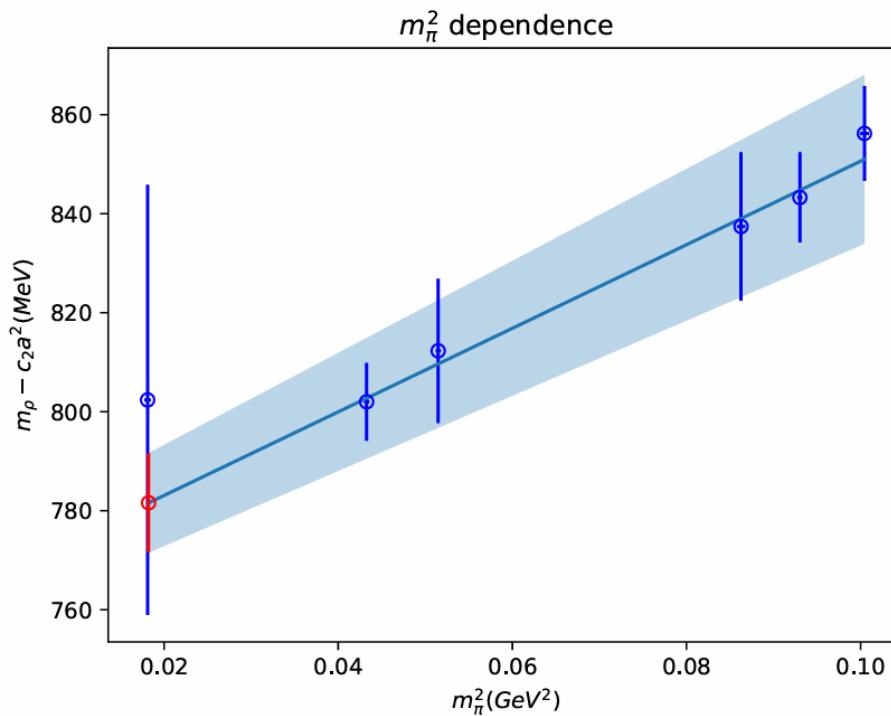
| | m_ρ (MeV) | $g_{\rho\pi\pi}$ |
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| H48P32 | 834.8(7.4) | 4.916(45) |

$$m_\rho = c_0 + c_1 m_\pi^2 + c_2 a^2 \quad \chi^2/d.o.f = 0.23$$

$$(c_0, c_1, c_2)^{\text{BW}} = (766.2(8.5), 0.84(9), -7.97(81)) \quad (\text{MeV}, \text{GeV}^{-1}, \text{GeV} \cdot \text{fm}^{-2})$$

$$g = \tilde{c}_0 + \tilde{c}_1 m_\pi^2 + \tilde{c}_2 a^2 \quad \chi^2/d.o.f = 0.25$$

$$(\tilde{c}_0, \tilde{c}_1, \tilde{c}_2)^{\text{BW}} = (5.74(67), -5.3(6.4), 46(48)) \quad (\text{None}, \text{GeV}^{-2}, \text{fm}^{-2})$$



Here we take $m_{\pi-phy} = m_{\pi^0} \sim 135$ MeV

$$m_\rho = 781.6(10.0) \text{ MeV}, \quad \Gamma_\rho(m_\rho) = 146.5(9.9) \text{ MeV}$$

$$Z_{\text{pole}} = 768.1(10.0) - i 70.5(4.9) \text{ MeV}$$

$$T_{\pi\pi \rightarrow \pi\pi}^{l=1}(z) = \frac{|V_{\rho\pi\pi}(\bar{k}(z))|^2}{z - m_\rho^B - \Sigma(z)},$$

$$\Sigma(z) = \int q^2 dq \frac{|V_{\rho\pi\pi}(q)|^2}{z - 2\sqrt{q^2 + m_\pi^2} + i\epsilon}.$$

HEFT
BW

$$\frac{-\sqrt{s}\Gamma(s)}{s - m_\rho^2 + i\sqrt{s}\Gamma(s)} = e^{i\delta(s)} \sin \delta(s)$$

$$\Gamma(s) = \frac{g_{\rho\pi\pi}^2}{6\pi s} \left(\frac{s}{4} - m_\pi^2 \right)^{3/2}$$



Introduction of HEFT

J. M. M. Hall etc. PRD 87(2013), 094510
 J.-j. Wu etc. PRC90 (2014), 055206
 Y. Li etc. PRD 101(2020), 114501
 PRD 103(2021), 094518

$$H = H_0 + H_I$$

$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha} |\alpha(k_{\alpha})\rangle \left[\sqrt{m_{\alpha 1}^2 + k_{\alpha}^2} + \sqrt{m_{\alpha 2}^2 + k_{\alpha}^2} \right] \langle \alpha(k_{\alpha})|$$

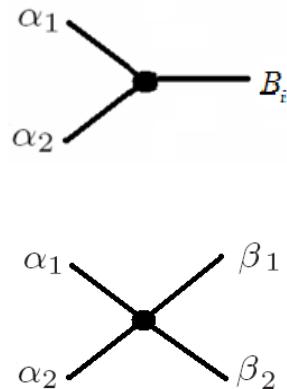
$|B_i\rangle$ bare state, bare mass m_i

$|\alpha(k_{\alpha})\rangle$ non-interaction channels

$$H_I = \hat{g} + \hat{v}$$

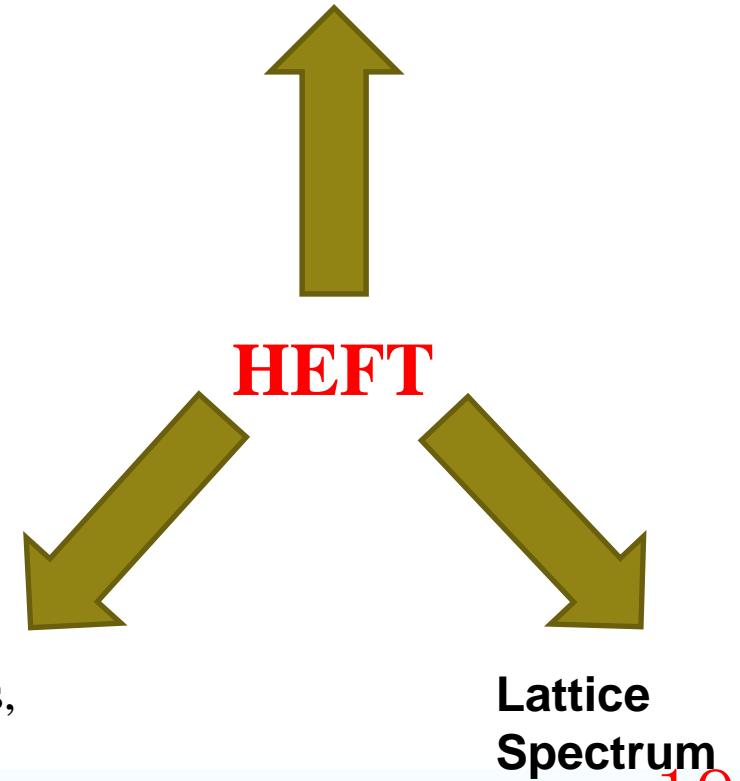
$$\hat{g} = \sum_{\alpha} \sum_{i=1,n} [|\alpha(k_{\alpha})\rangle g_{i,\alpha}^{+} \langle B_i| + |B_i\rangle g_{i,\alpha} \langle \alpha(k_{\alpha})|]$$

$$\hat{v} = \sum_{\alpha, \beta} |\alpha(k_{\alpha})\rangle v_{\alpha, \beta} \langle \beta(k_{\beta})|$$



T matrix
 (Phase Shifts,
 inelasticity)

Resonance
 (Mass , Width, Pole position, Coupling)



ρ meson spectra on LQCD -- Pole of T

$$T_{\pi\pi \rightarrow \pi\pi}^{l=1}(z) = \frac{|V_{\rho\pi\pi}(\bar{k}(z))|^2}{z - m_\rho^B - \Sigma(z)} \quad e^{i\delta(s)} \sin \delta(s) = -\rho(s) T_{\pi\pi \rightarrow \pi\pi}^{l=1}(\sqrt{s}) \quad \rho(s) = \frac{\pi}{4} \left(s \left(\frac{s}{4} - m_\pi^2 \right) \right)^{1/2}$$

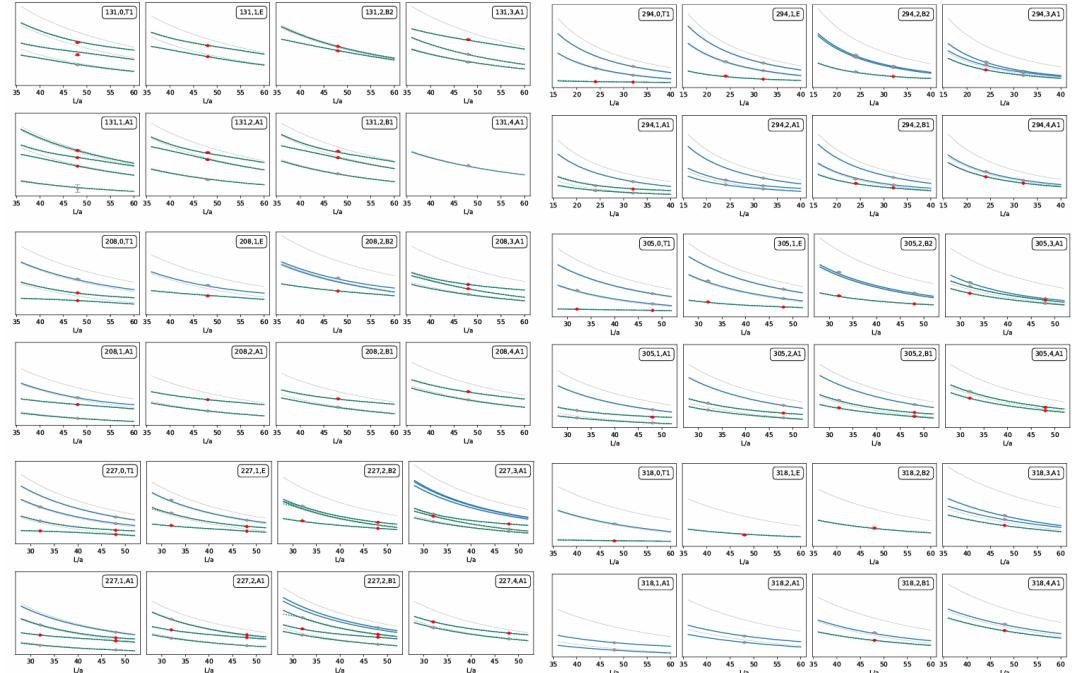
$$\Sigma(z) = \int q^2 dq \frac{|V_{\rho\pi\pi}(q)|^2}{z - 2\sqrt{q^2 + m_\pi^2} + i\epsilon}. \quad V_{\rho\pi\pi}(k) = \frac{g_{\rho\pi\pi} k}{\sqrt{m_\rho^B} \sqrt{k^2 + m_\pi^2}} \left(\frac{\Lambda_{\rho\pi\pi}^2}{k^2 + \Lambda_{\rho\pi\pi}^2} \right)^2$$

$$\Sigma_{\omega\pi}(z) = \int q^2 dq \frac{|V_{\rho\omega\pi}(q)|^2}{z - \sqrt{q^2 + m_\pi^2} - \sqrt{q^2 + m_\omega^2} + i\epsilon} \quad V_{\rho\omega\pi}(k) = \frac{g_{\rho\omega\pi} k \sqrt{m_\rho^B}}{\sqrt{2E_\pi(k) E_\omega(k)}} \left(\frac{\Lambda_{\rho\omega\pi}^2 - \mu_\pi^2}{k^2 + \Lambda_{\rho\omega\pi}^2} \right)^2$$

- Five free parameters, $m_\rho^B, g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}, g_{\rho\omega\pi}, \Lambda_{\rho\omega\pi}$
- Here we take $g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$ as constant for different m_π and a .
- $g_{\rho\omega\pi}=18 \text{ GeV}^{-1}$, $\Lambda_{\rho\omega\pi}=1 \text{ GeV}$ by $\omega \rightarrow 3\pi$
- But m_ρ^B is pion mass dependence.

Scheme A, $g_{\rho\pi\pi} = 7.0 \pm 0.1$, $\Lambda_{\rho\pi\pi} = 0.9 \pm 0.1 \text{ GeV}$ without $\omega\pi$ channel.

Scheme B, $g_{\rho\pi\pi} = 7.4 \pm 0.1$, $\Lambda_{\rho\pi\pi} = 1.0 \pm 0.1 \text{ GeV}$ with $\omega\pi$ channel.



ρ meson spectra on LQCD -- Pole of T

$$T_{\pi\pi \rightarrow \pi\pi}^{l=1}(z) = \frac{|V_{\rho\pi\pi}(\bar{k}(z))|^2}{z - m_\rho^B - \Sigma(z)} \quad e^{i\delta(s)} \sin \delta(s) = -\rho(s) T_{\pi\pi \rightarrow \pi\pi}^{l=1}(\sqrt{s}) \quad \rho(s) = \frac{\pi}{4} \left(s \left(\frac{s}{4} - m_\pi^2 \right) \right)^{1/2}$$

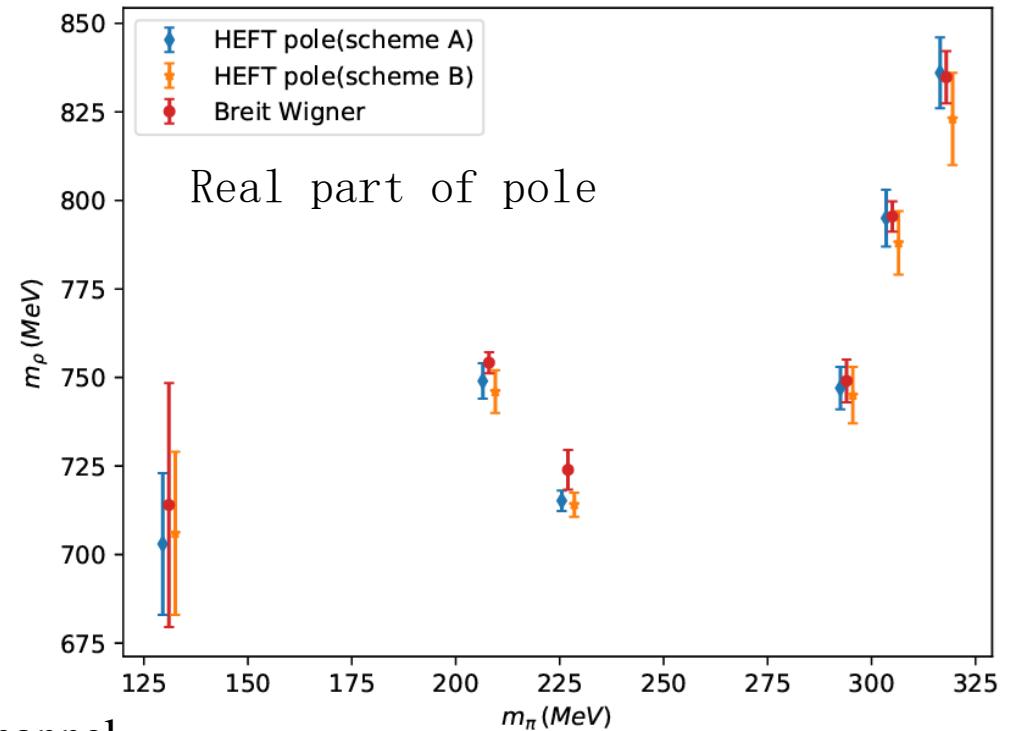
$$\Sigma(z) = \int q^2 dq \frac{|V_{\rho\pi\pi}(q)|^2}{z - 2\sqrt{q^2 + m_\pi^2} + i\epsilon}. \quad V_{\rho\pi\pi}(k) = \frac{g_{\rho\pi\pi} k}{\sqrt{m_\rho^B} \sqrt{k^2 + m_\pi^2}} \left(\frac{\Lambda_{\rho\pi\pi}^2}{k^2 + \Lambda_{\rho\pi\pi}^2} \right)^2$$

$$\Sigma_{\omega\pi}(z) = \int q^2 dq \frac{|V_{\rho\omega\pi}(q)|^2}{z - \sqrt{q^2 + m_\pi^2} - \sqrt{q^2 + m_\omega^2} + i\epsilon} \quad V_{\rho\omega\pi}(k) = \frac{g_{\rho\omega\pi} k \sqrt{m_\rho^B}}{\sqrt{2E_\pi(k) E_\omega(k)}} \left(\frac{\Lambda_{\rho\omega\pi}^2 - \mu_\pi^2}{k^2 + \Lambda_{\rho\omega\pi}^2} \right)^2$$

- Five free parameters, $m_\rho^B, g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}, g_{\rho\omega\pi}, \Lambda_{\rho\omega\pi}$
- Here we take $g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$ as constant for different m_π and a .
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ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

Only m_ρ^B is pion mass dependence.

$$m_\rho^B(m_\pi, a) = c_0 + c_1 m_\pi^2 + c_2 a^2$$

$$(c_0, c_1, c_2)^A = (817.0(13.0), 0.85(12), -7.61(89))$$

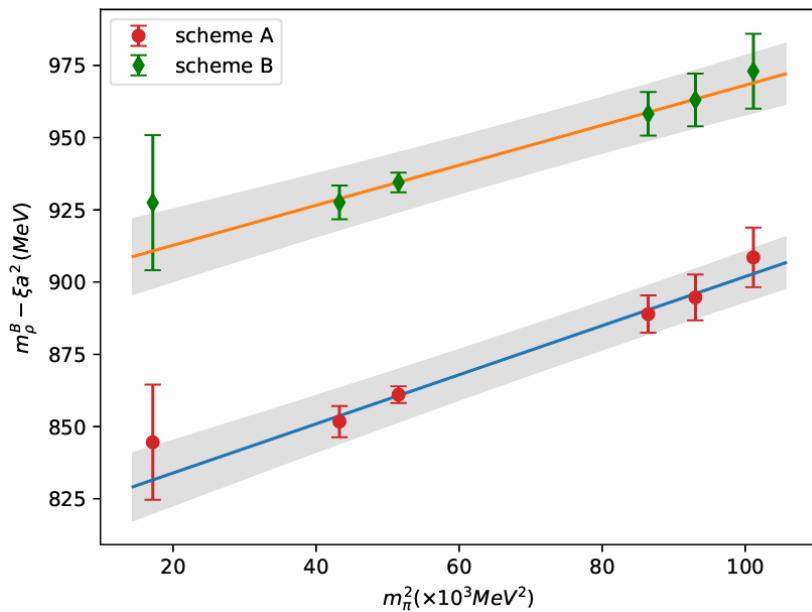
$$(c_0, c_1, c_2)^B = (840.0(11.2), 0.72(11), -7.91(86))$$

$$m_\rho^{\text{pole,ext(A)}} = 787.0(15.0) - i 59.0(6.0) \text{ MeV}$$

$$m_\rho^{\text{pole,ext(B)}} = 777.0(15.0) - i 60.0(6.0) \text{ MeV}$$

$\rho(770)$ T-MATRIX POLE \sqrt{s}

$(761 - 765) - i(71 - 74)$ MeV



at $\delta = 90$ degree

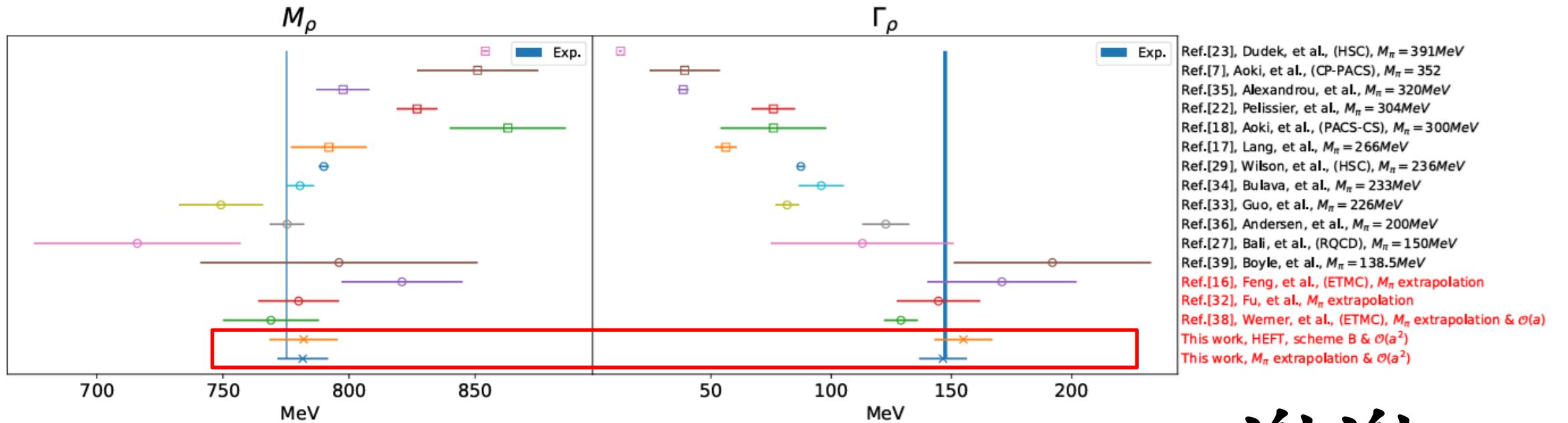
$$m_\rho = 782.0(13.5) \text{ MeV}, \quad \Gamma_\rho(m_\rho) = 155.0(12.0) \text{ MeV}$$

$$m_\rho = 781.6(10.0) \text{ MeV}, \quad \Gamma_\rho(m_\rho) = 146.5(9.9) \text{ MeV}$$

$$Z_{\text{pole}} = 768.1(10.0) - i 70.5(4.9) \text{ MeV}$$

Summary

We using CLQCD configuration measure the rho meson mass and width at the **physical pion mass** and **continuum limit** with the highest precision up to now!



谢谢

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