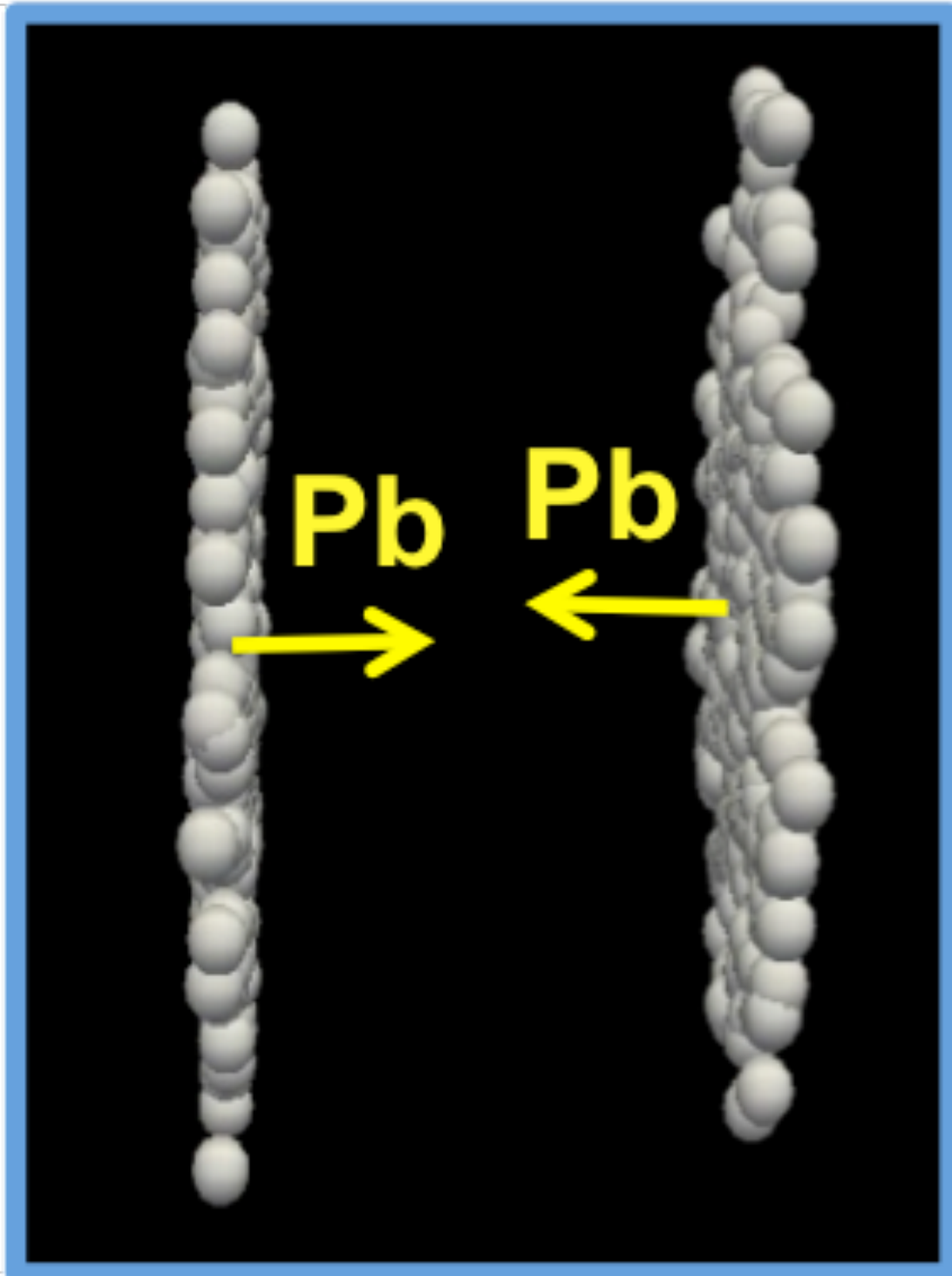


Quantum thermalization and hydrodynamization
in high energy collisions

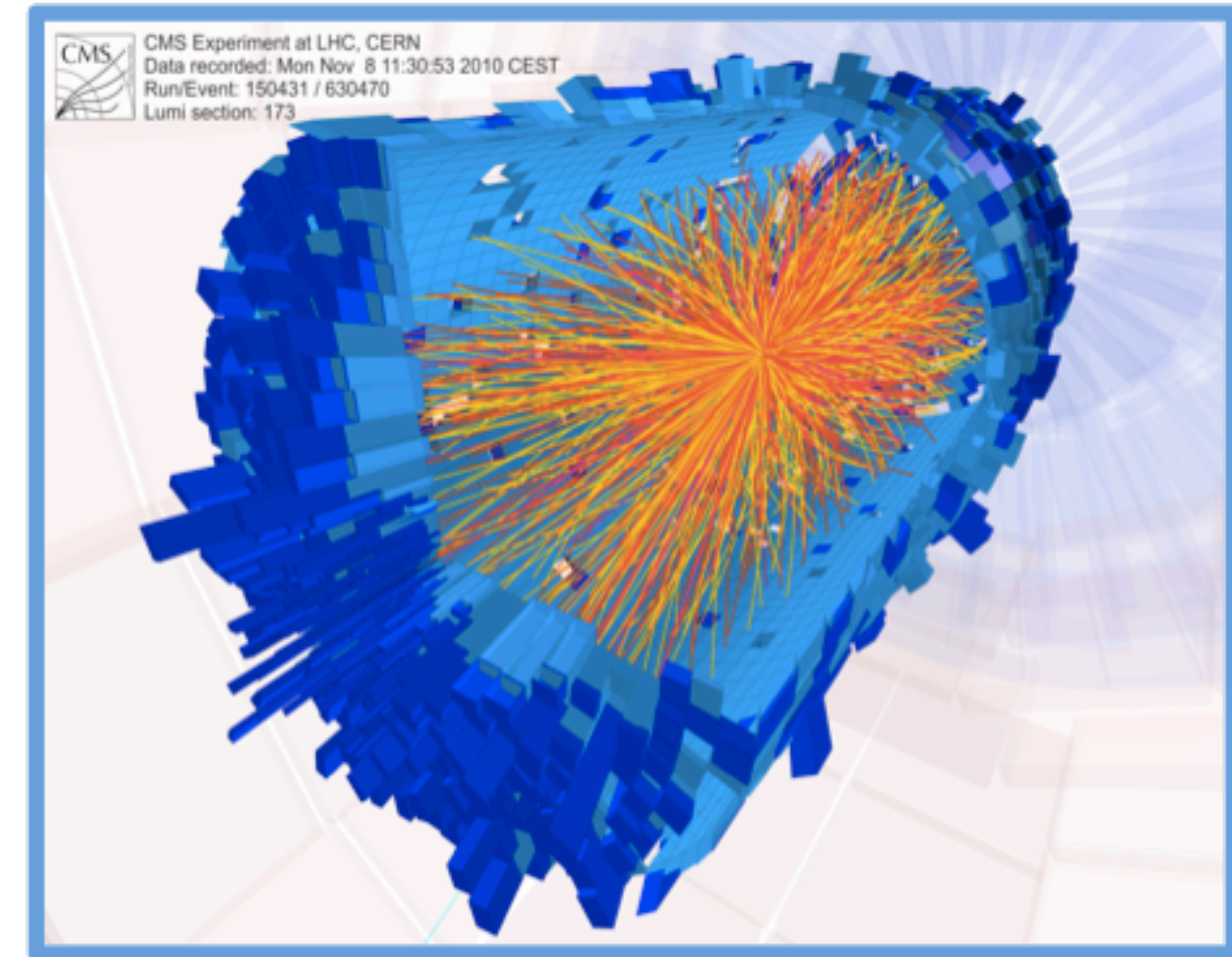
施舒哲，清华大学

w/ Shile Chen (陈诗乐) & Li Yan (严力), 2412.00662.
w/ Haiyang Shao (邵海洋) & Shile Chen (陈诗乐), in preparation.

Pre-reaction



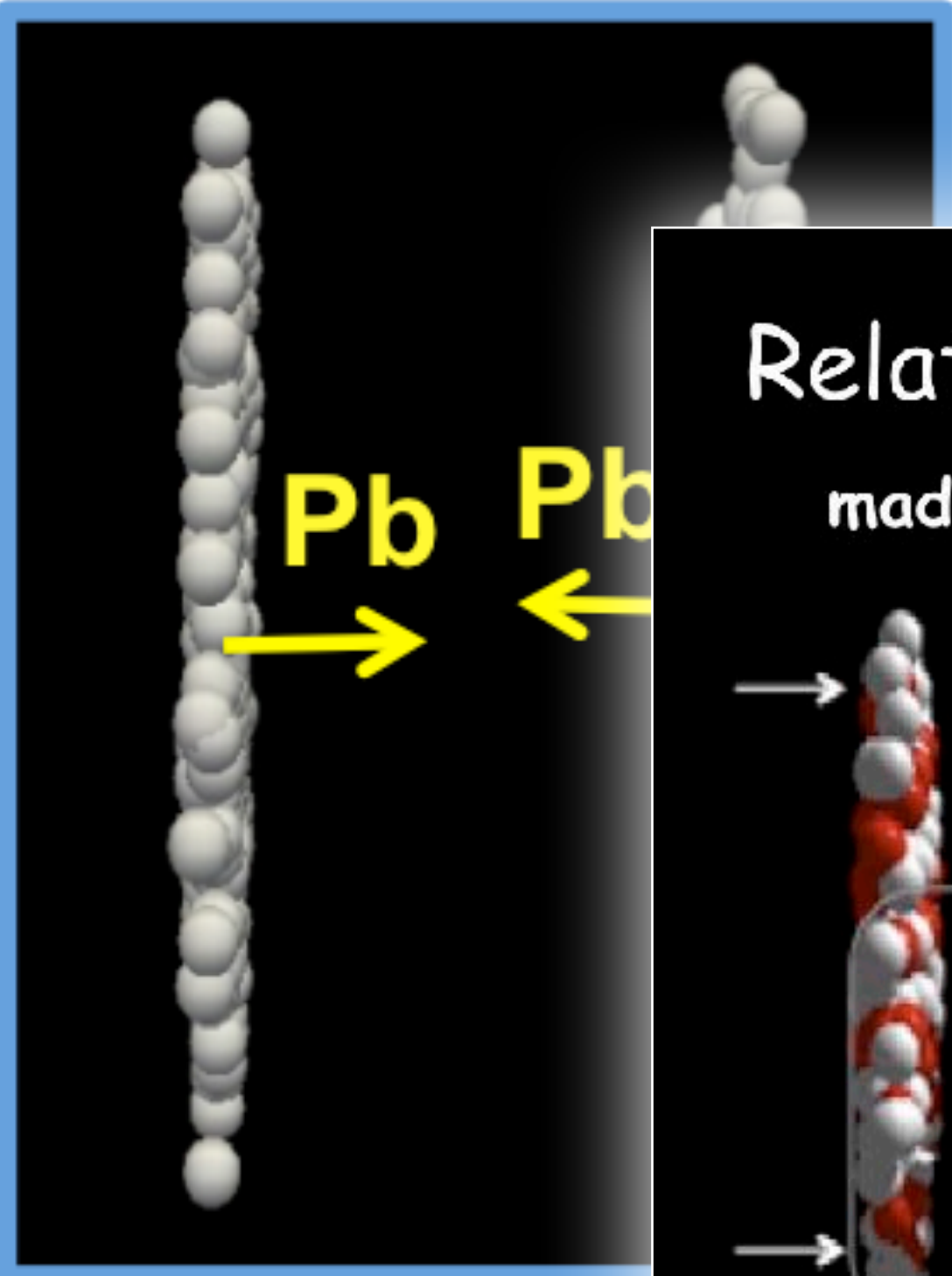
Detection



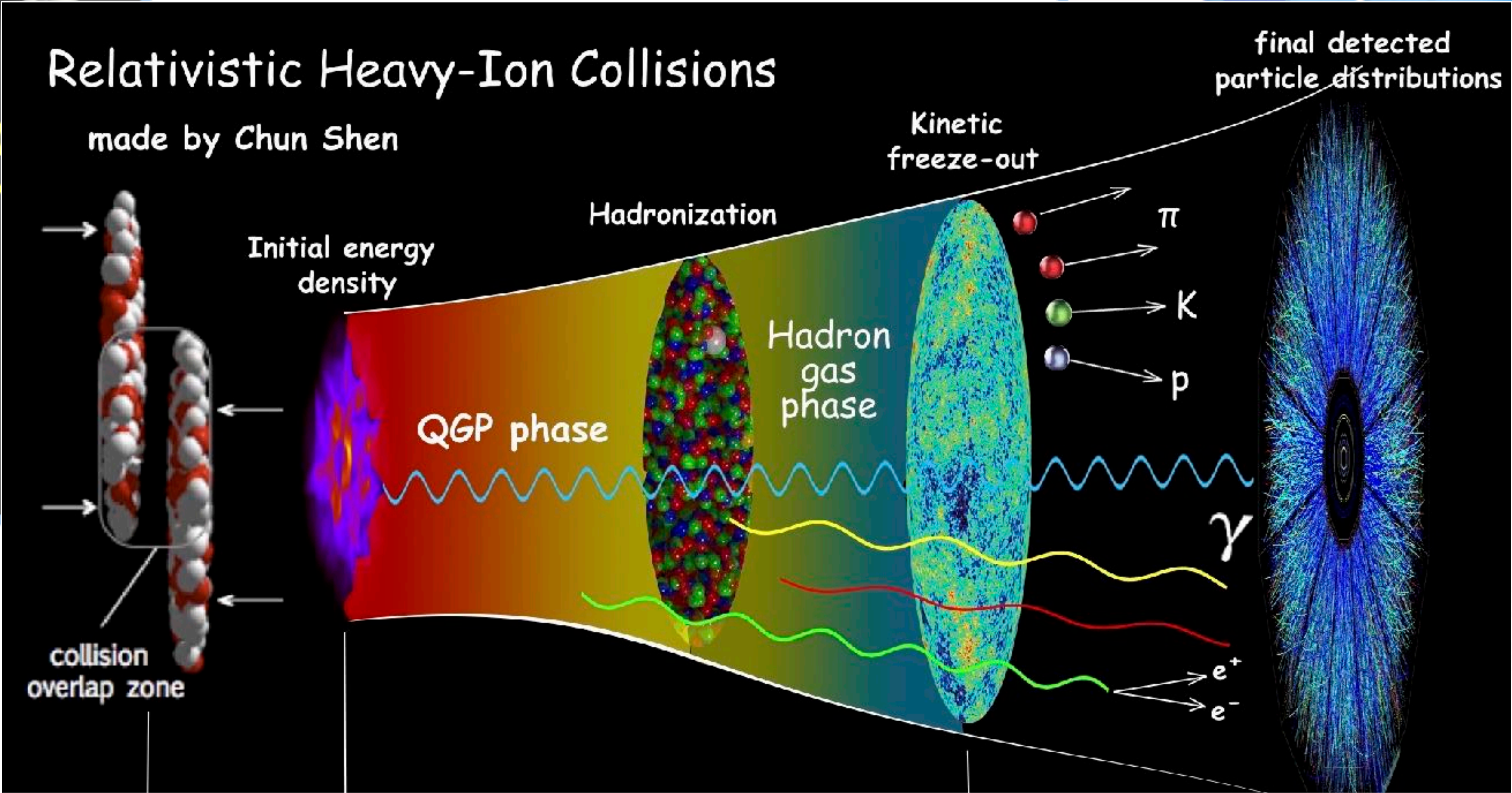
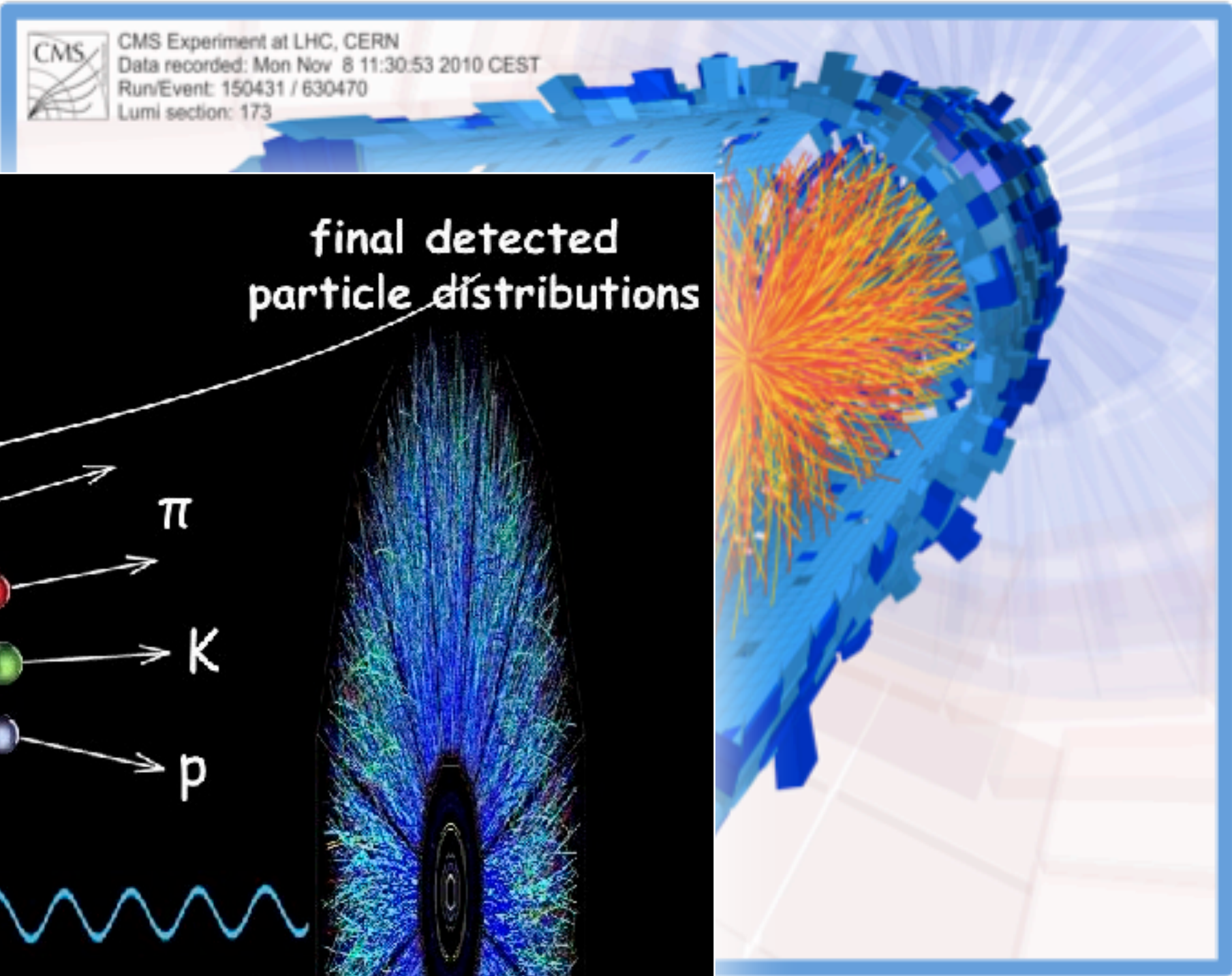
final state particles

$$N \sim 10^3-4$$

Pre-reaction



Detection



particles
3-4

Hydrodynamic attractors

- focus on properties of *hydrodynamic equations*
- *common behavior* developed very rapidly

rapidity distribution? *Shile Chen's talk tomorrow 09:00, Parallel I*

Hydrodynamic attractors

- focus on properties of *hydrodynamic equations*
- *common behavior* developed very rapidly

rapidity distribution? *Shile Chen's talk tomorrow 09:00, Parallel I*

Thermalization in phase space distribution

- solve Boltzmann equations, compared with hydro
- usually *linearized* Boltzmann, e.g., Relaxation Time Approximation

non-linear collision kernel? multiparticle correlations? *Shuai Lu's poster #223*

Quark-Gluon Plasma:

A strongly coupled, quantum, (quasi)many-body system!

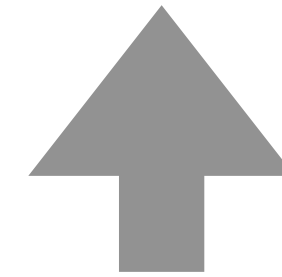
Ideally, full quantum simulation of QCD in 3+1D

Practically, strong coupling QED in 1+1D

– confinement, chiral condensate

1+1D Schwinger model

$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$



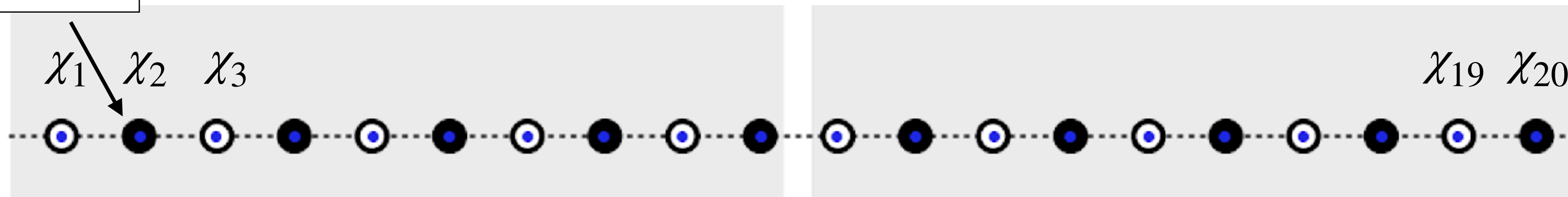
E : electric field
 A : electric potential
 $\psi, \bar{\psi}$: fermion field

$$L(t) = \int \left(-\frac{F^{\mu\nu}F_{\mu\nu}}{4} + \bar{\psi}(i\gamma^\mu \partial_\mu - g\gamma^\mu A_\mu - m)\psi \right) dx.$$

1+1D Schwinger model

$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$

$|0\rangle$:empty
 $|1\rangle$:occupied



$$\{\psi_a(x), \psi_b^\dagger(y)\} = \delta_{a,b} \delta(x-y)$$

field operators
 represented by
 matrices

$$\{\chi_n^\dagger, \chi_m\} = \delta_{nm}, \quad \{\chi_n^\dagger, \chi_m^\dagger\} = \{\chi_n, \chi_m\} = 0.$$

Jordan-Wigner

$$\chi_n = \frac{X_n - iY_n}{2} \prod_{m=1}^{n-1} (-iZ_m)$$

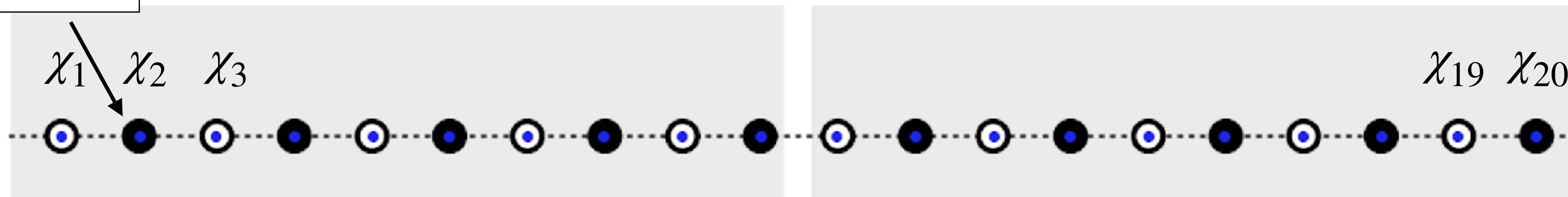
$$X_n \equiv I \otimes \cdots \otimes I \otimes \sigma_x \otimes I \otimes \cdots \otimes I$$

n^{th}

1+1D Schwinger model

$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$

$|0\rangle$:empty
 $|1\rangle$:occupied



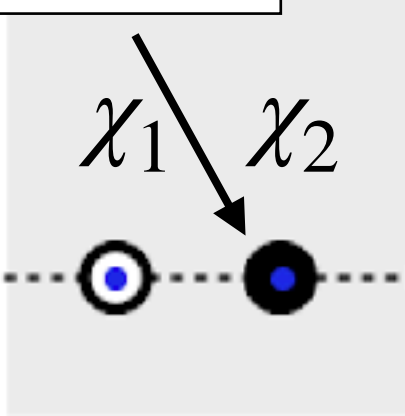
$$\{\psi_a(x), \psi_b^\dagger(y)\} = \delta_{a,b}\delta(x-y)$$

field operators
 represented by
 matrices

$$\{\chi_n^\dagger, \chi_m\} = \delta_{nm}, \quad \{\chi_n^\dagger, \chi_m^\dagger\} = \{\chi_n, \chi_m\} = 0.$$

$$H = \frac{1}{4a} \sum_{n=1}^{N-1} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{a g^2}{2} \sum_{n=1}^{N-1} L_n^2.$$

$|0\rangle$:empty
 $|1\rangle$:occupied



Eigenstates of H : vacuum and (quasi)-particles

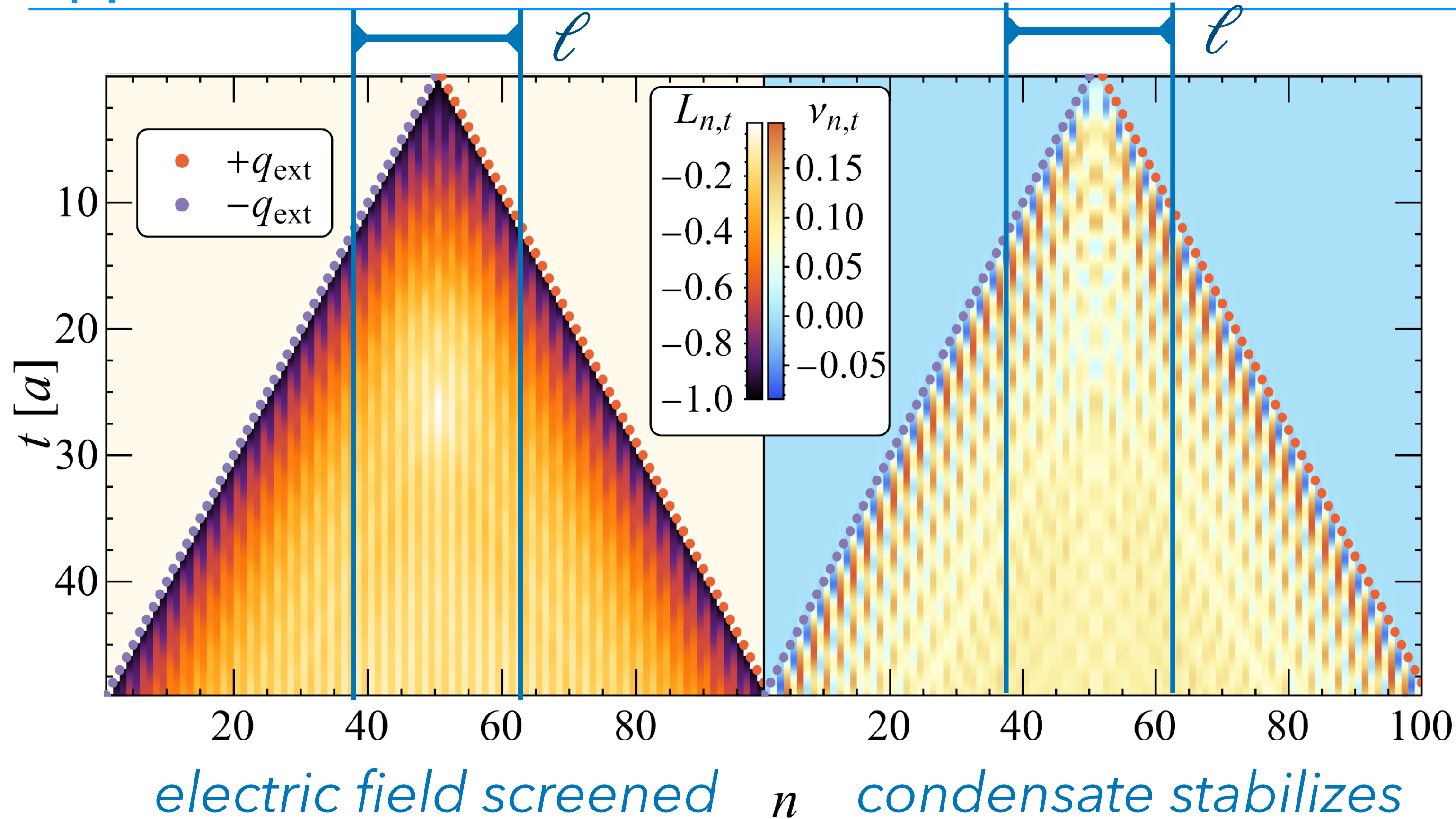
Time evolution:

$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i H |\psi(t)\rangle$$

$$O(t) = \langle \psi(t) | \hat{O} | \psi(t) \rangle$$

$$H = \frac{1}{4a} \sum_{n=1} (\Lambda_n \Lambda_{n+1} + I_n I_{n+1}) + \frac{1}{2} \sum_{n=1} (-1)^n Z_n + \frac{1}{2} \sum_{n=1} L_n .$$

$\delta_{a,b} \delta(x - y)$
d operators
resented by
rices
 $\{\chi_n, \chi_m\} = 0 .$

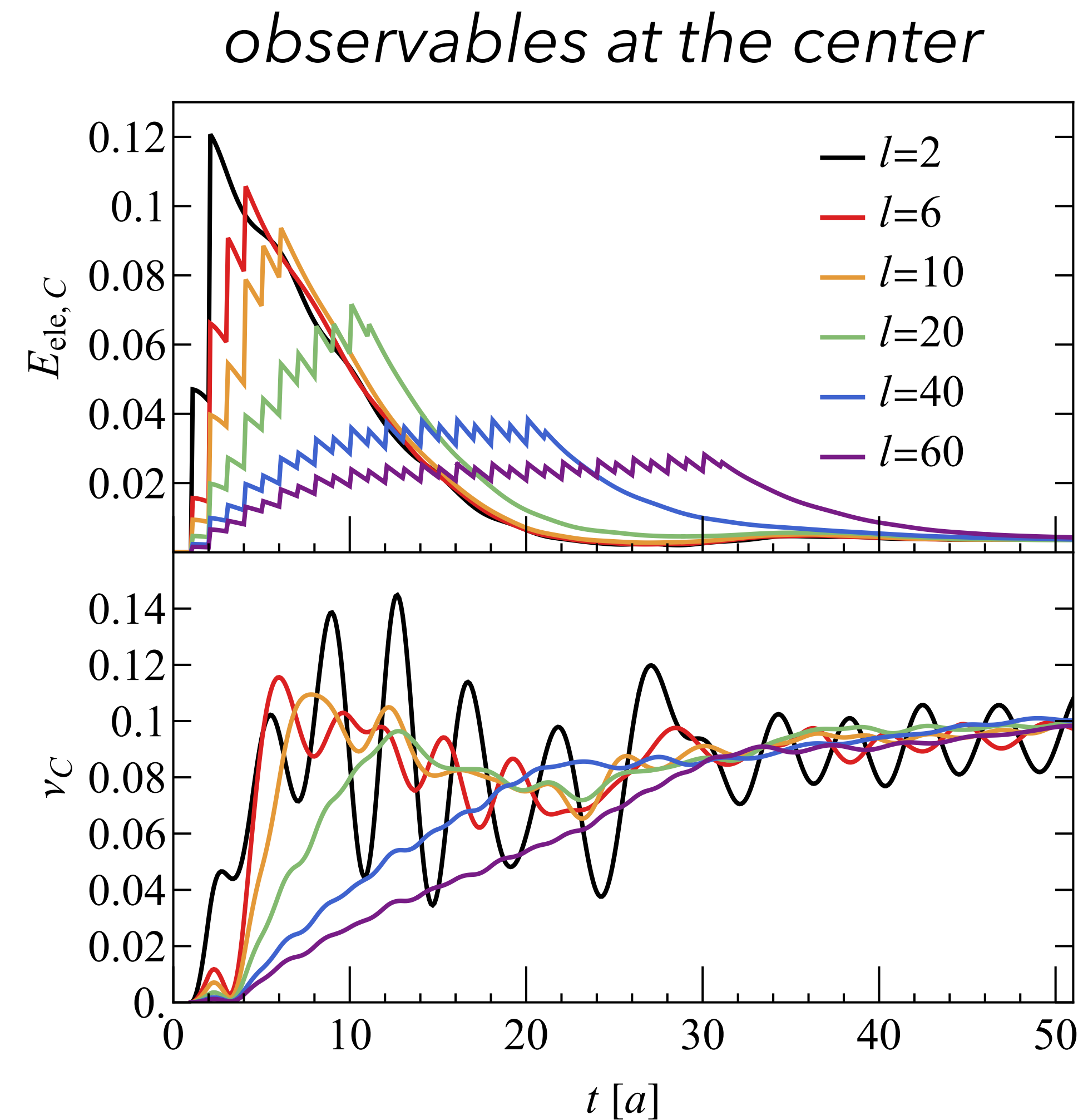


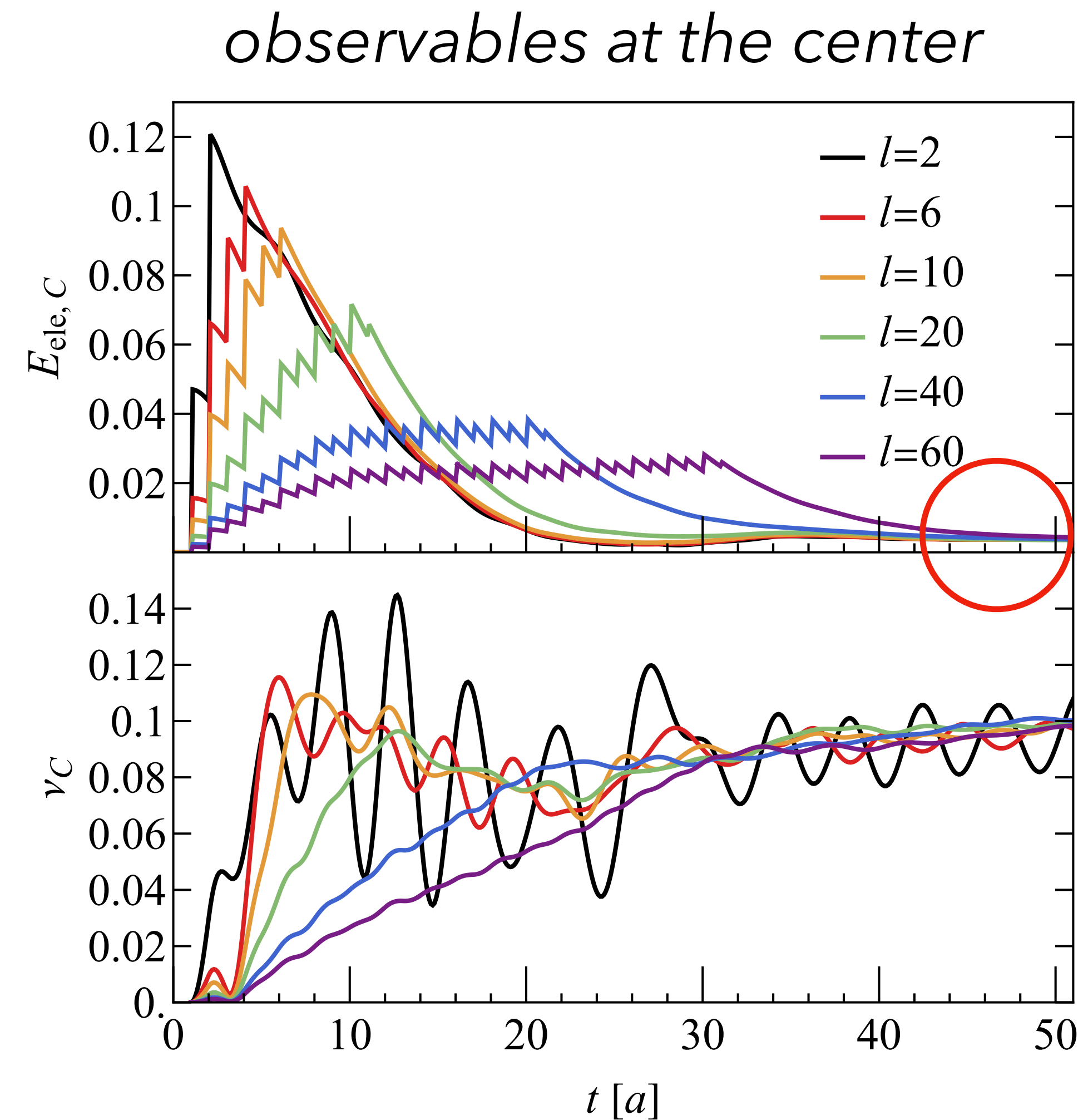
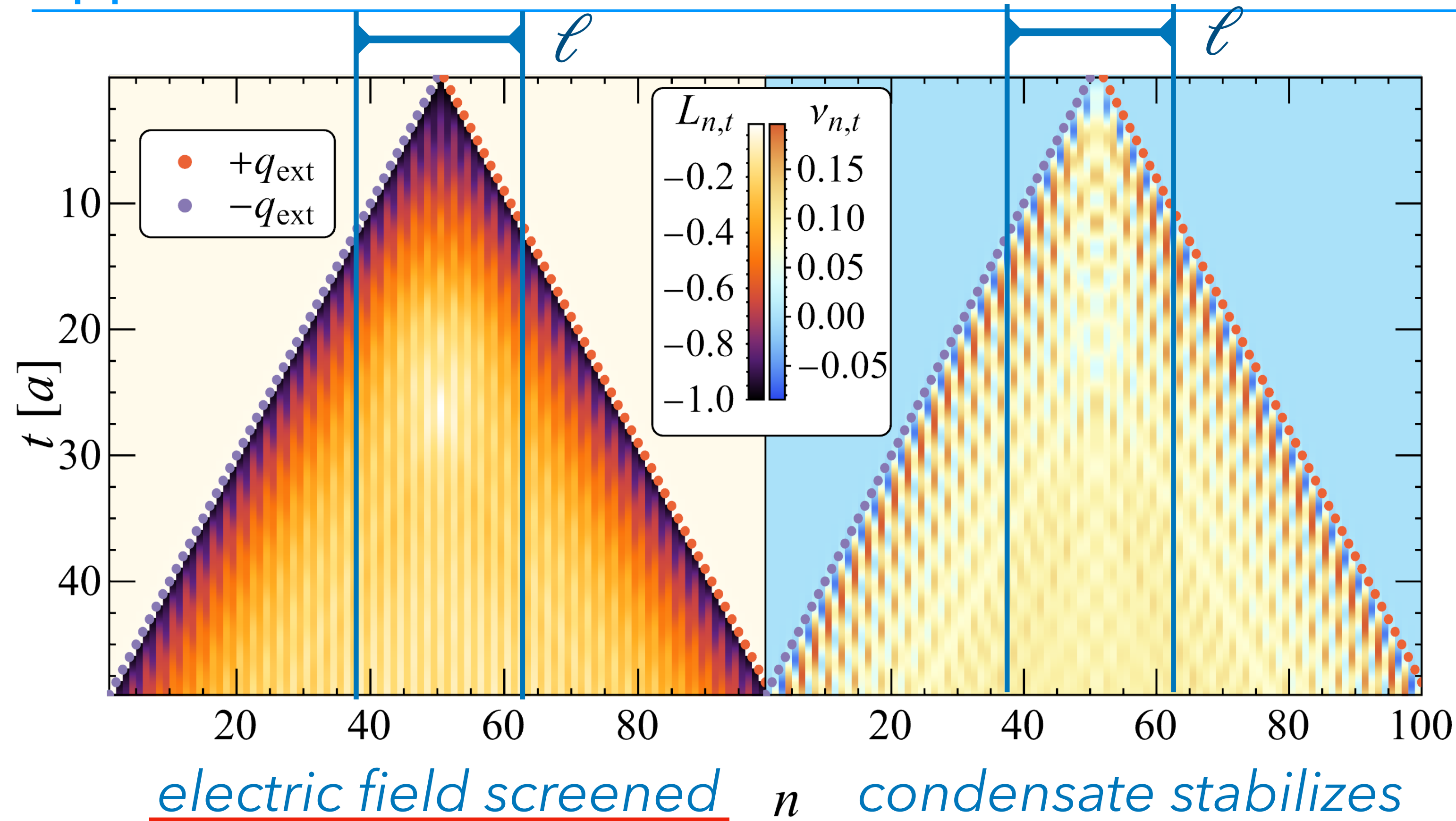
$$H_{\text{Schwinger}} |\psi(t=0)\rangle = E_0 |\psi(t=0)\rangle$$

$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i \left(H_{\text{Schwinger}} + H_{\text{source}}(t) \right) |\psi(t)\rangle$$

w/ Florio, Frenklakh, Ikeda, Kharzeev, Korepin, Yu,

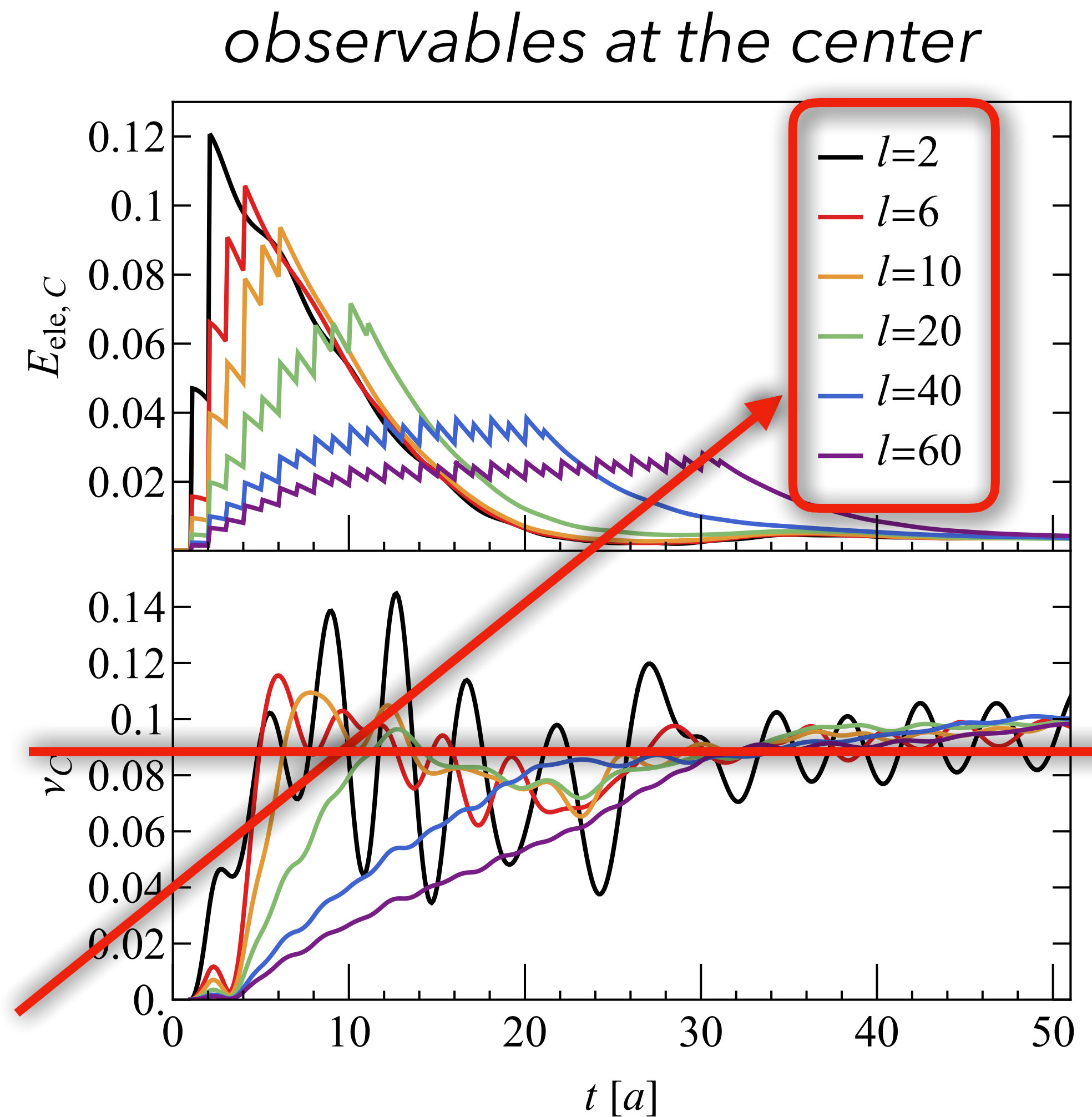
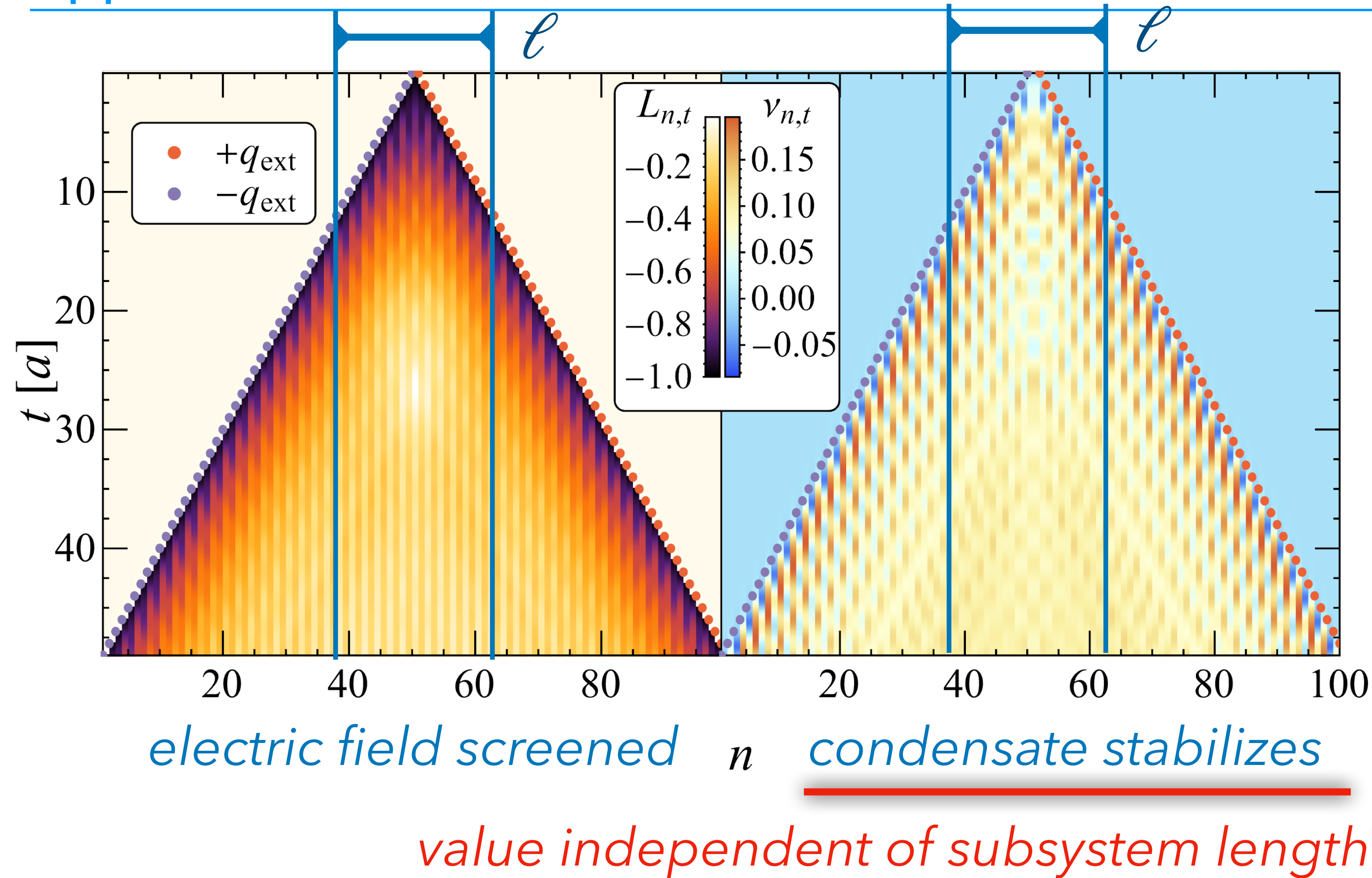
PhysRevLett.131.021902; PhysRevD.110. 094029; in progress





w/ Florio, Frenklakh, Ikeda, Kharzeev, Korepin, Yu,

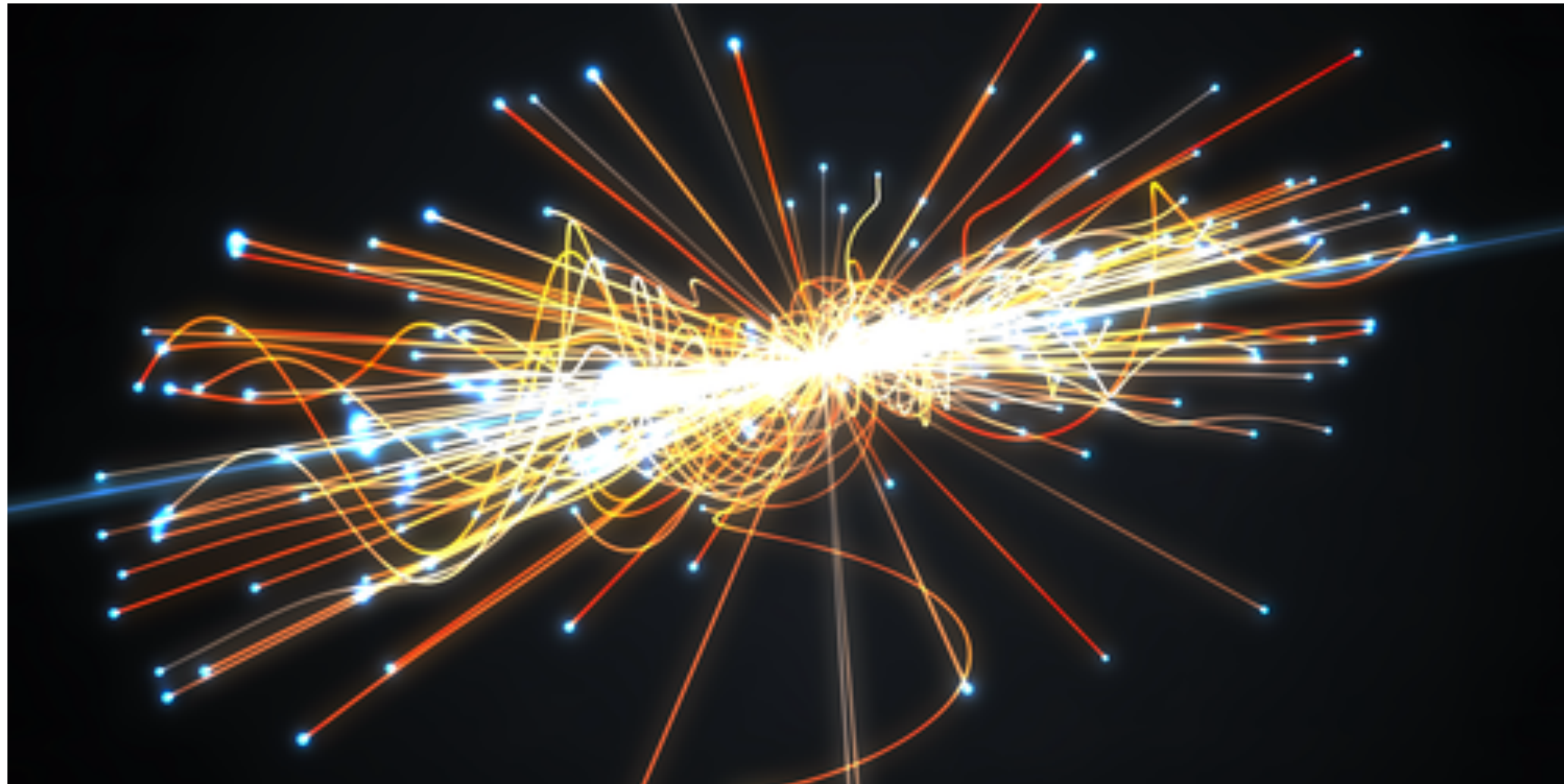
PhysRevLett.131.021902; PhysRevD.110. 094029; in progress



w/ Florio, Frenklakh, Ikeda, Kharzeev, Korepin, Yu,

PhysRevLett.131.021902; PhysRevD.110. 094029; in progress





$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$

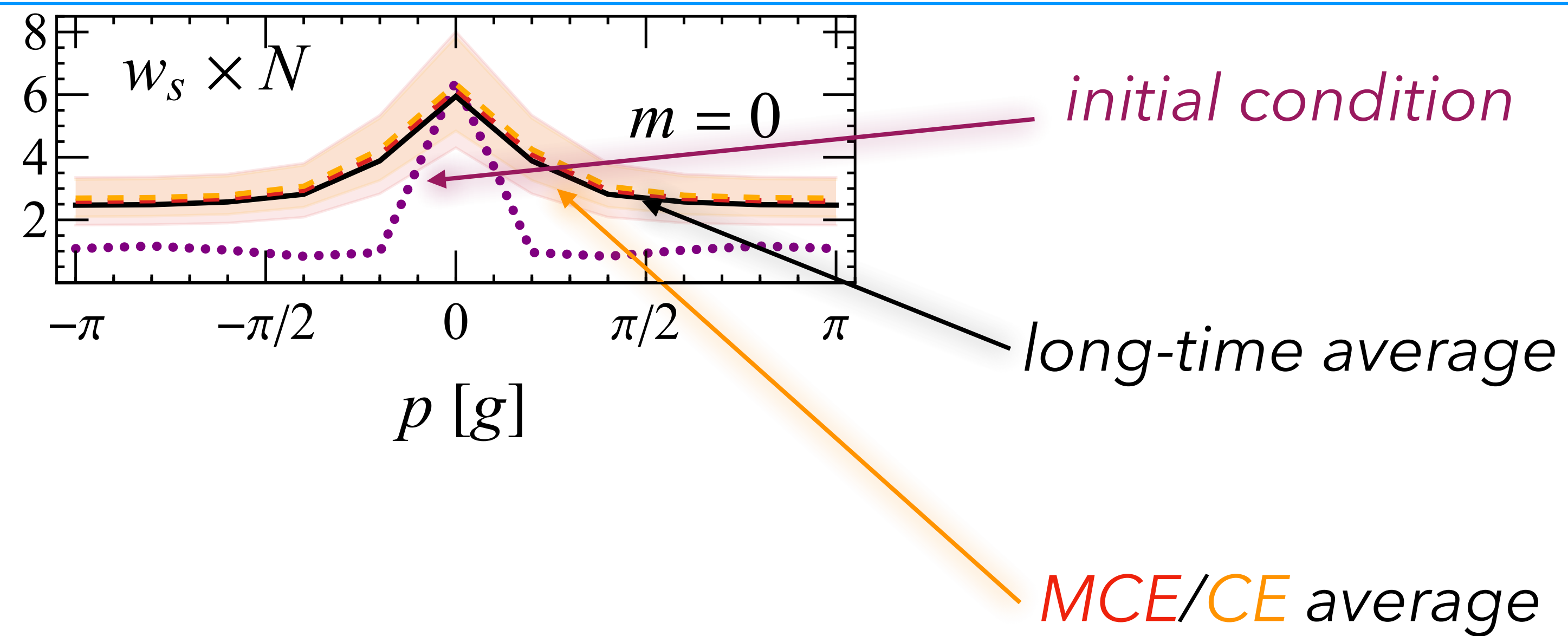
$$|\psi(t)\rangle = \exp(-i H t) |\psi(0)\rangle$$

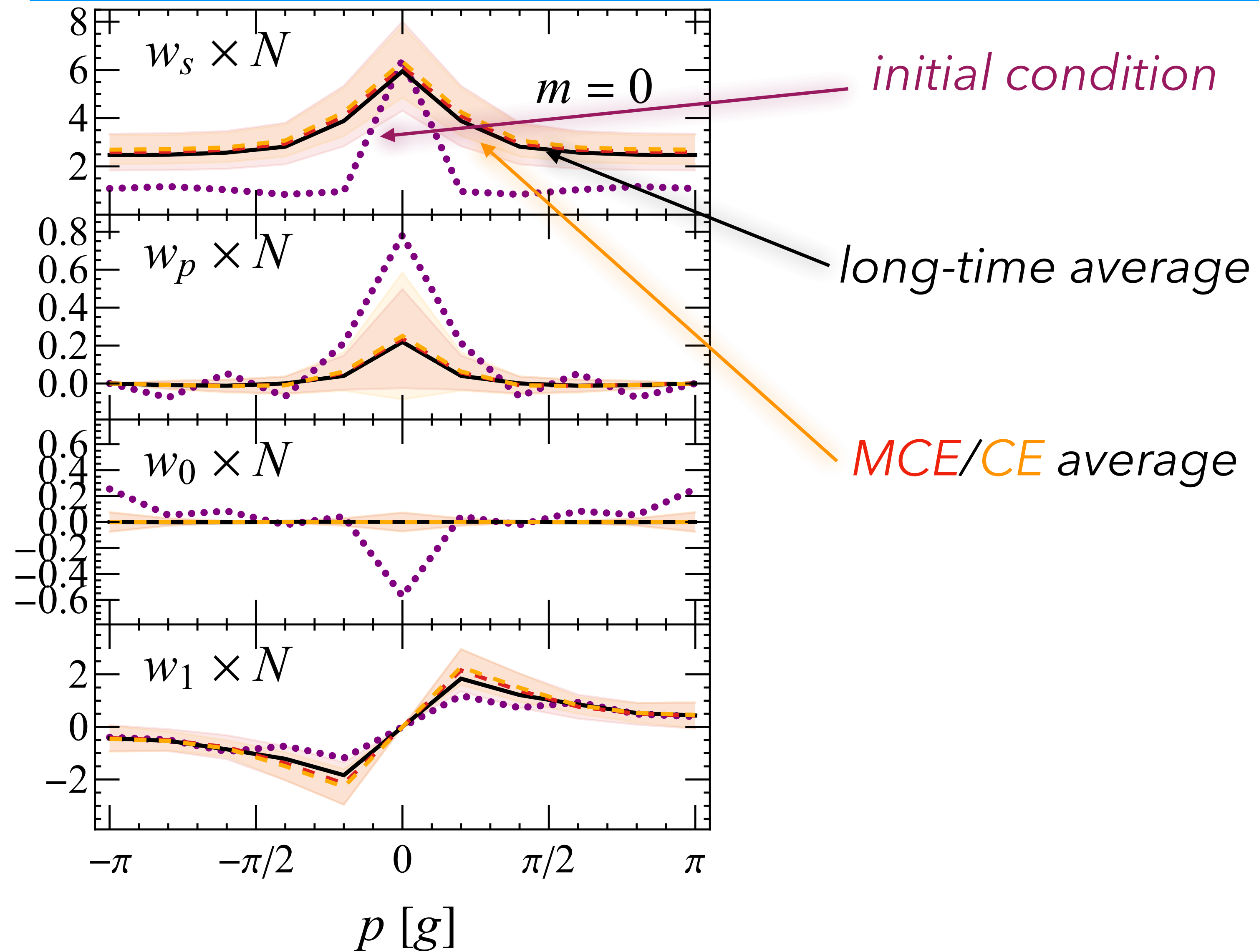
$$W(t) = \langle \psi(t) | \hat{W} | \psi(t) \rangle$$

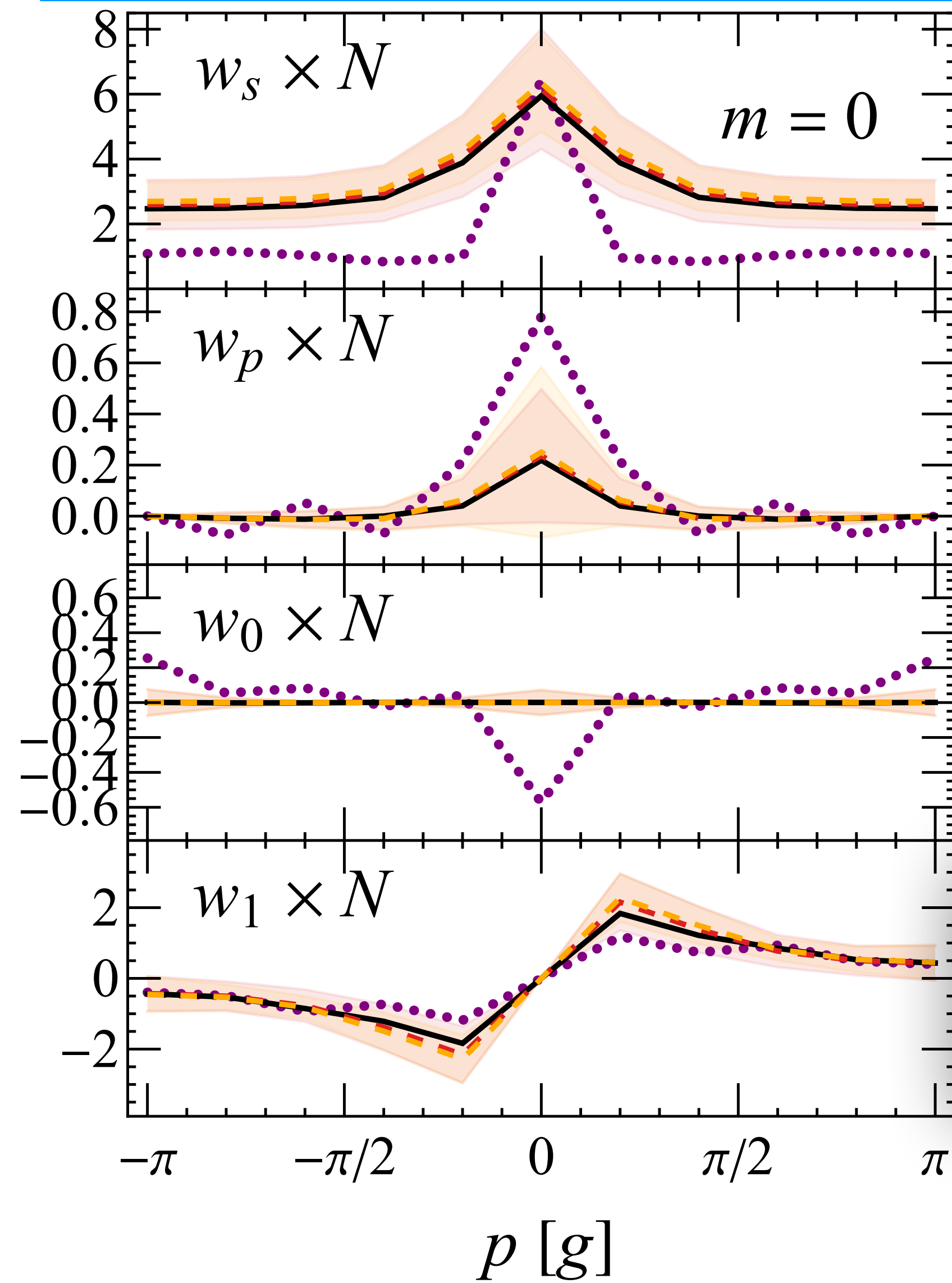
Thermalization in an isolated quantum system?

Quantum distribution function: Wigner function

$$\hat{W}_{\alpha\beta}(t, z, p) = \int \bar{\psi}_{\alpha}(z_{+}) U(z_{+}, z_{-}) \psi_{\beta}(z_{-}) e^{i \frac{py}{\hbar}} dy$$







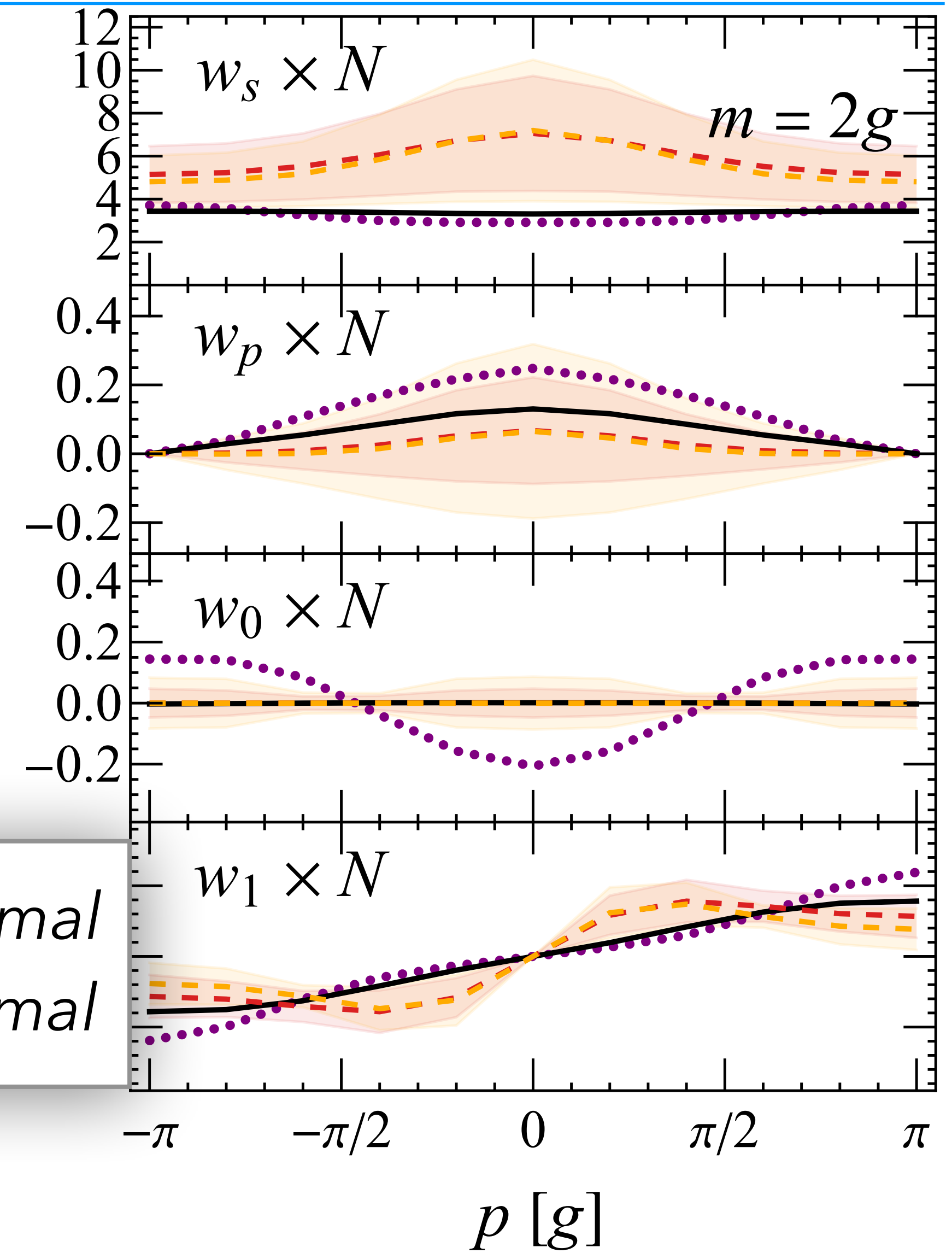
strong coupling

initial condition

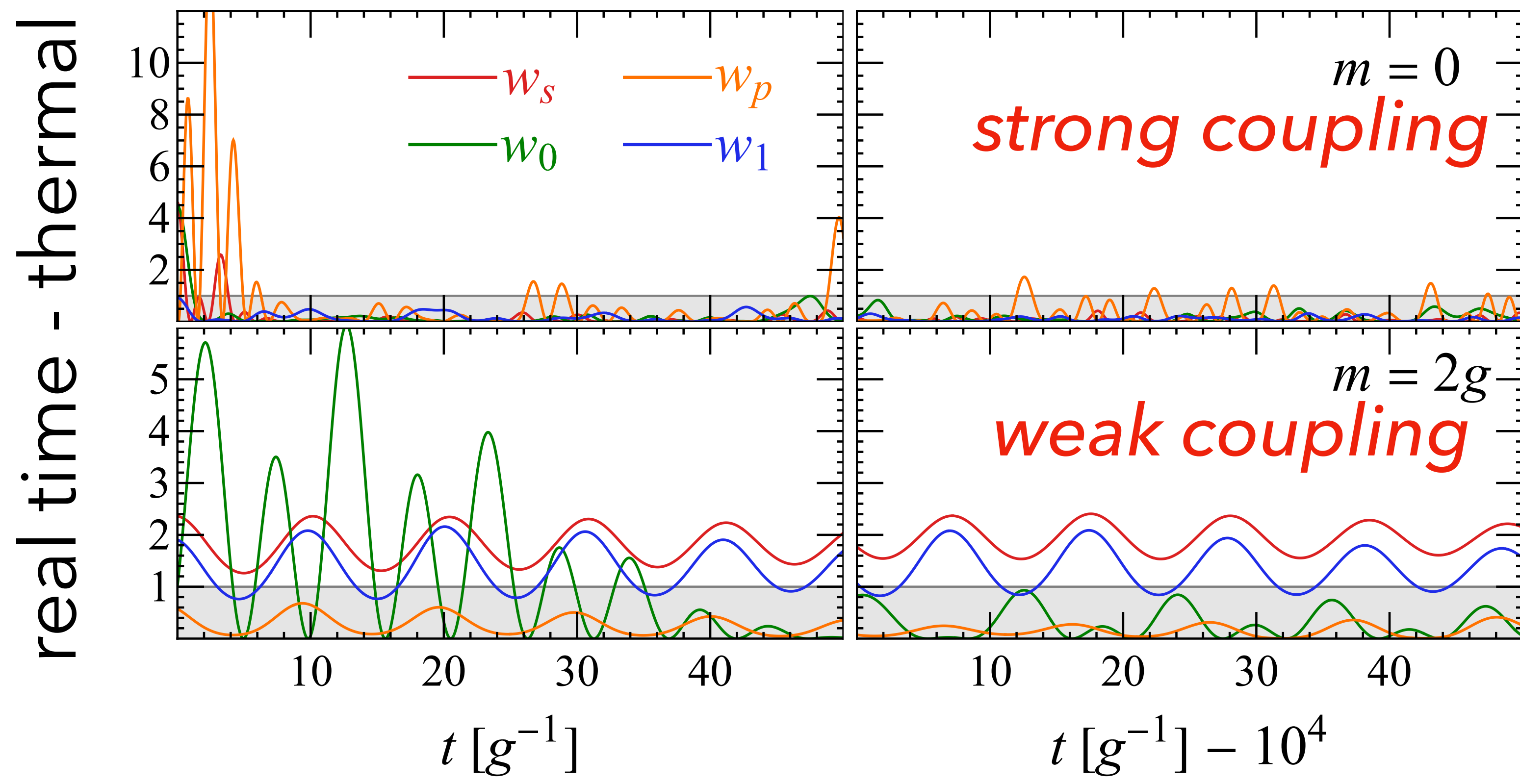
long-time average

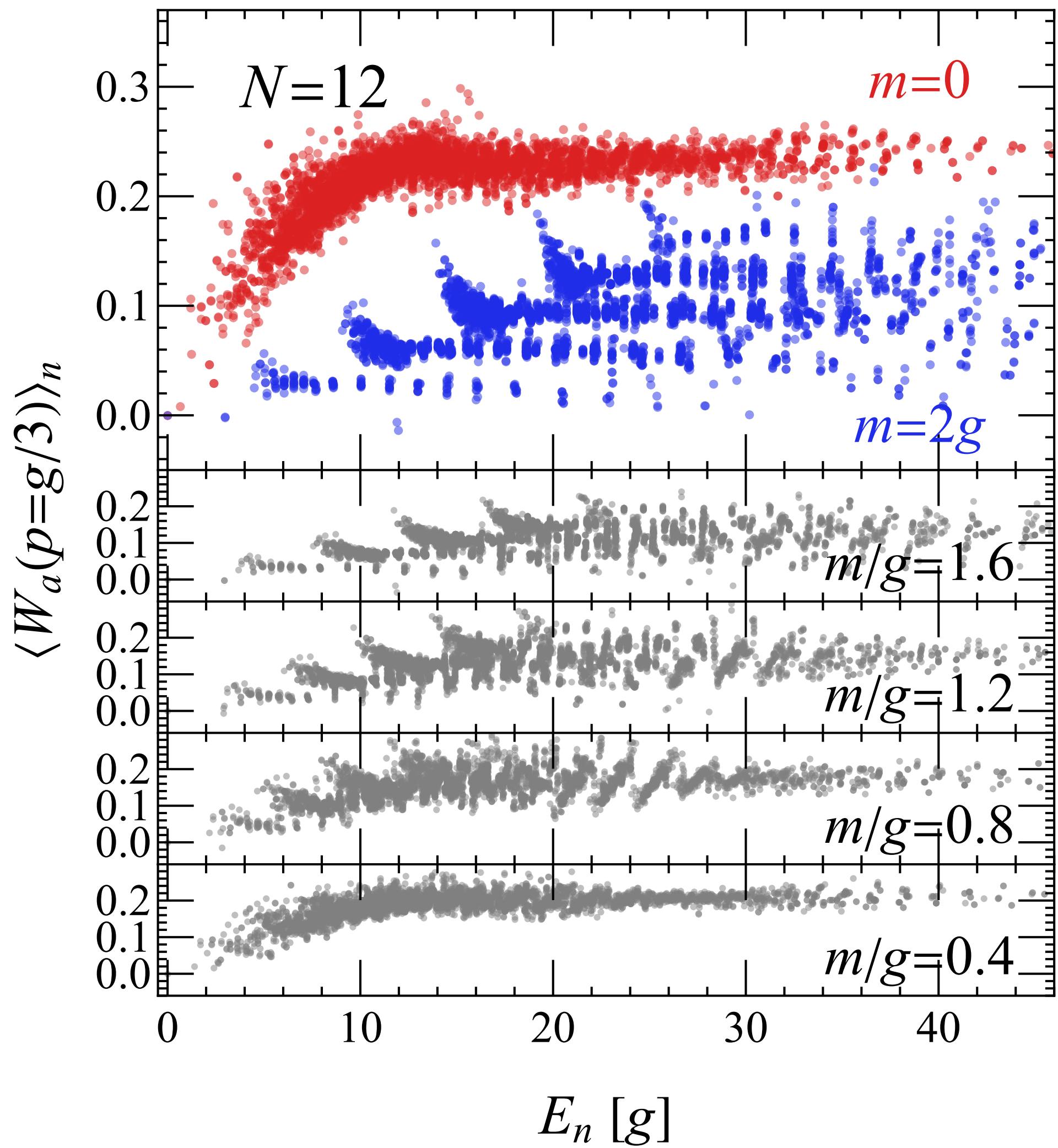
MCE/CE average

strong coupling: LTA = thermal
weak coupling: LTA $\neq$ thermal

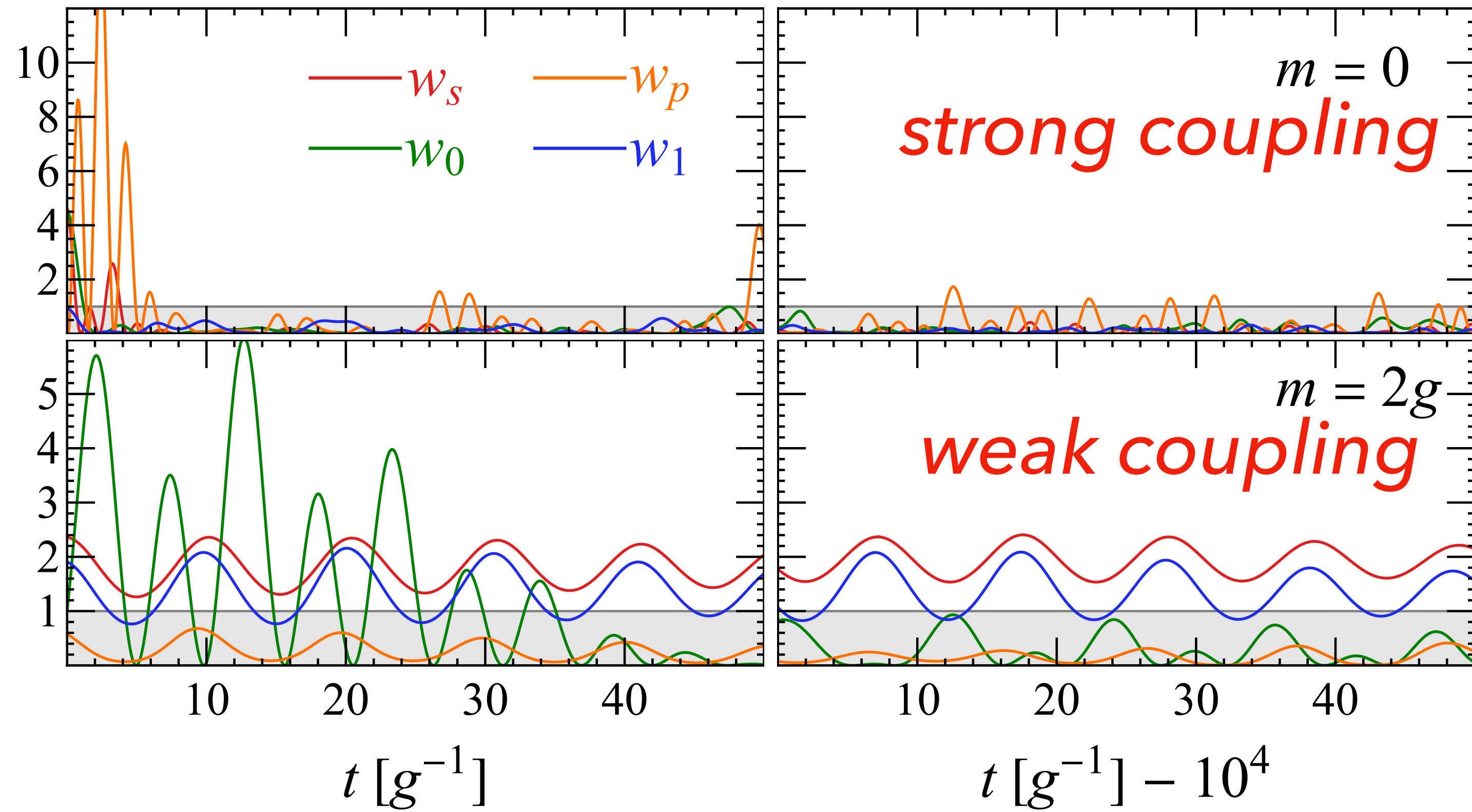


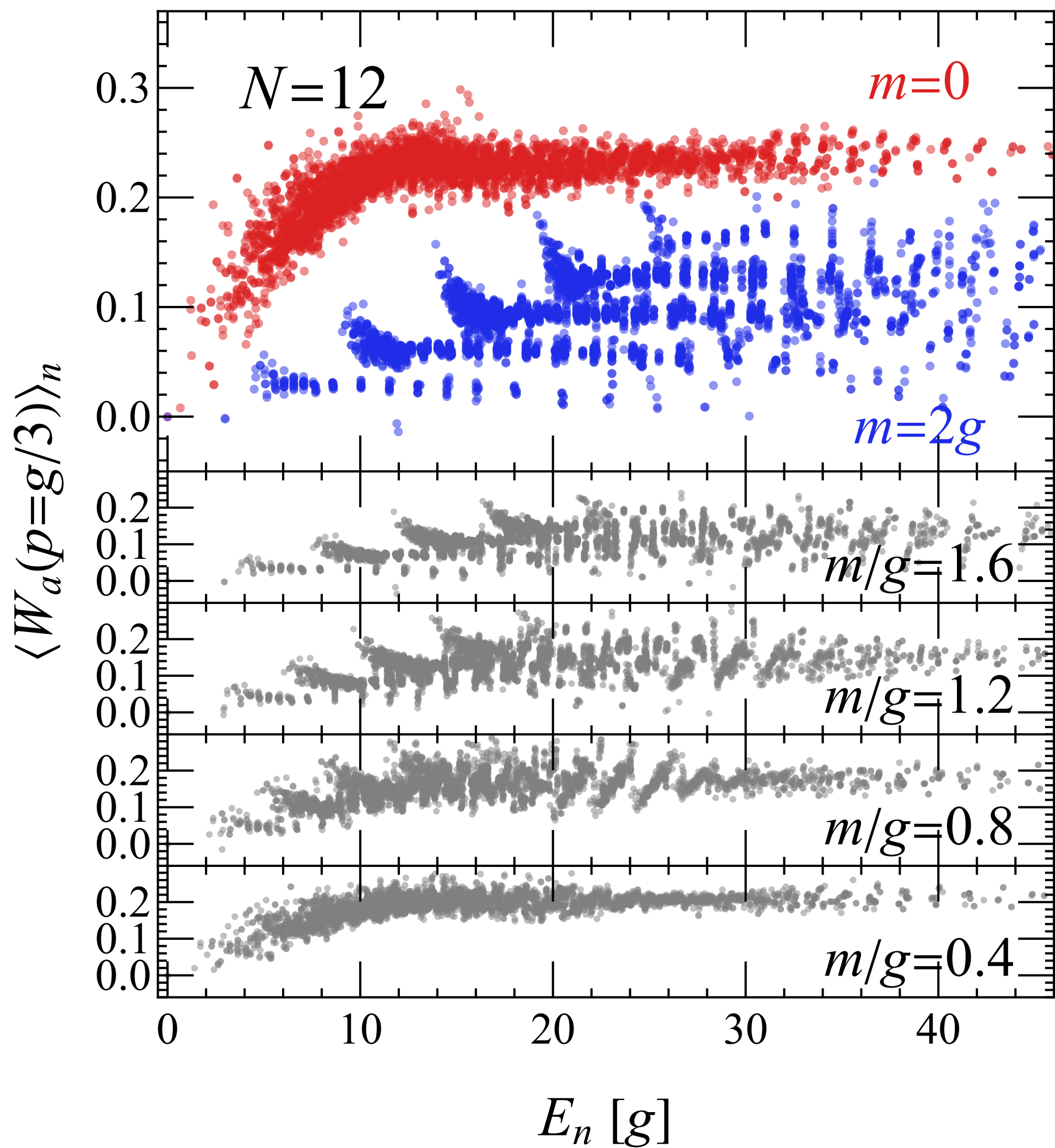
weak coupling





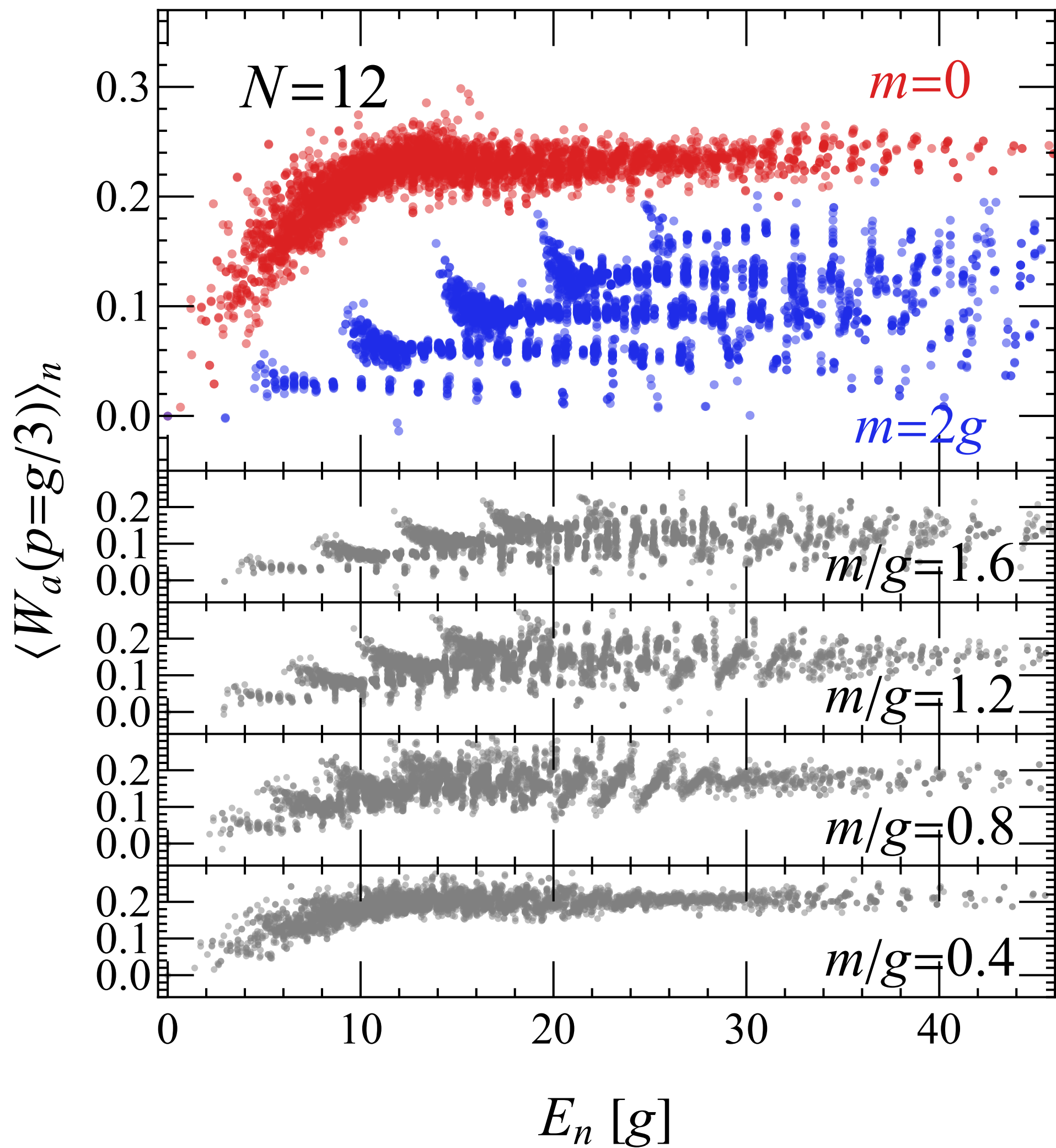
real time - thermal





eigenstate thermalization hypothesis

$$\langle n | \hat{O} | n \rangle \approx f_O(E_n)$$



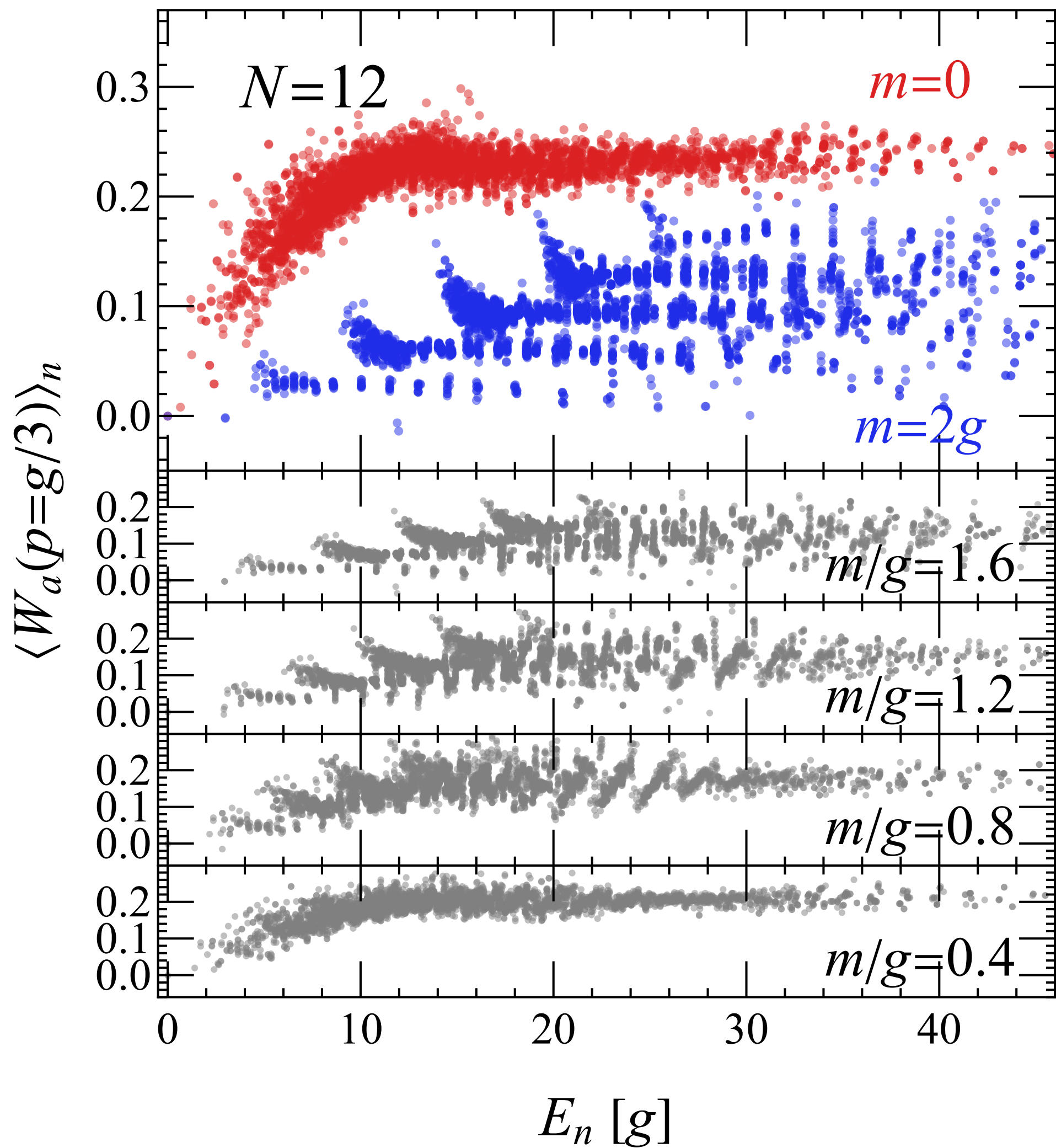
eigenstate thermalization hypothesis

$$\langle n | \hat{O} | n \rangle \approx f_O(E_n)$$

$$\sum_n p_n \langle n | \hat{O} | n \rangle \approx f_O(\sum_n p_n E_n)$$

general pure state $|c_n|^2$

thermal $e^{-\beta E_n}/Z$



eigenstate thermalization hypothesis

$$\langle n | \hat{O} | n \rangle \approx f_O(E_n)$$

$$\sum_n p_n \langle n | \hat{O} | n \rangle \approx f_O(\sum_n p_n E_n)$$

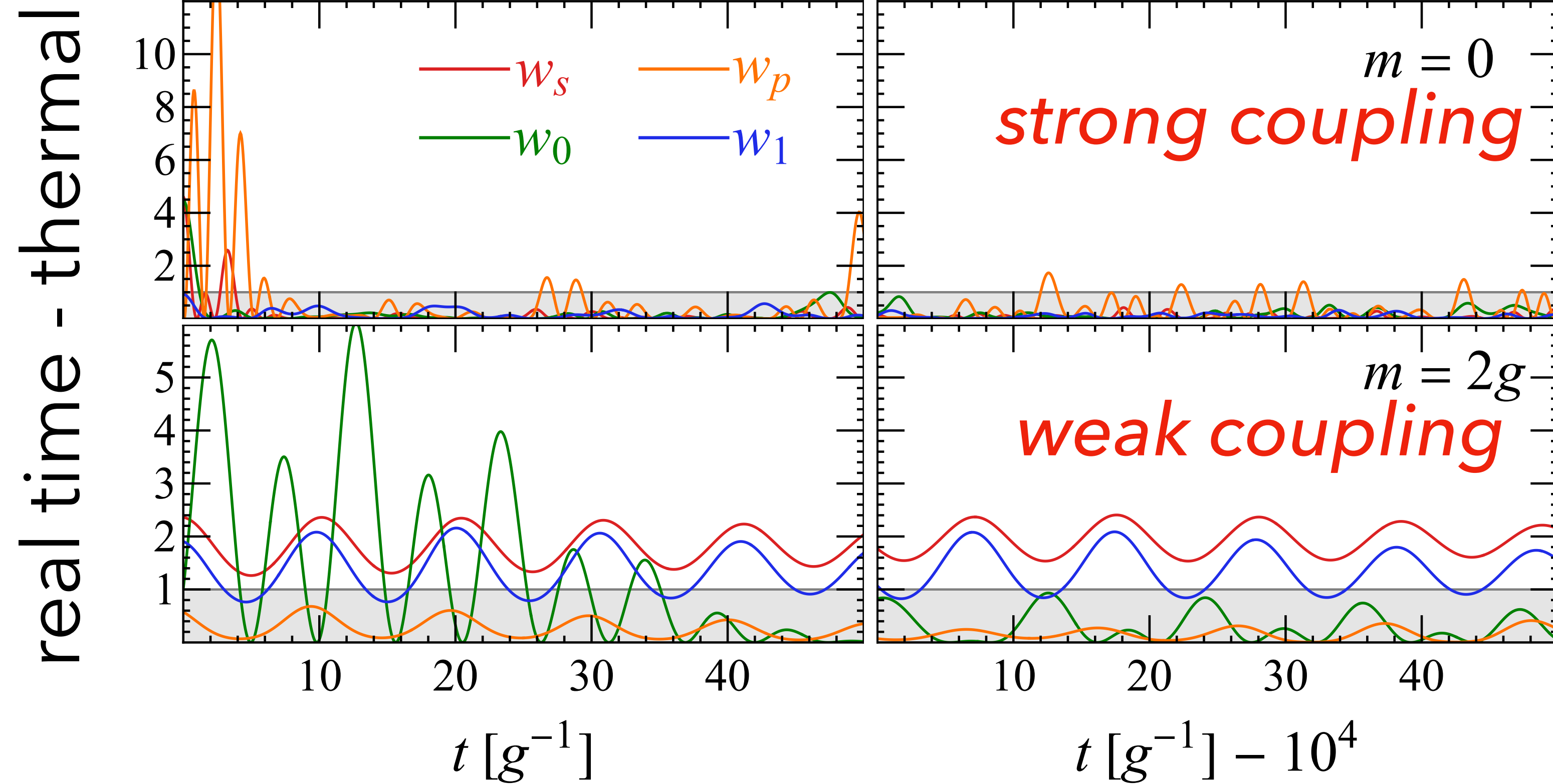
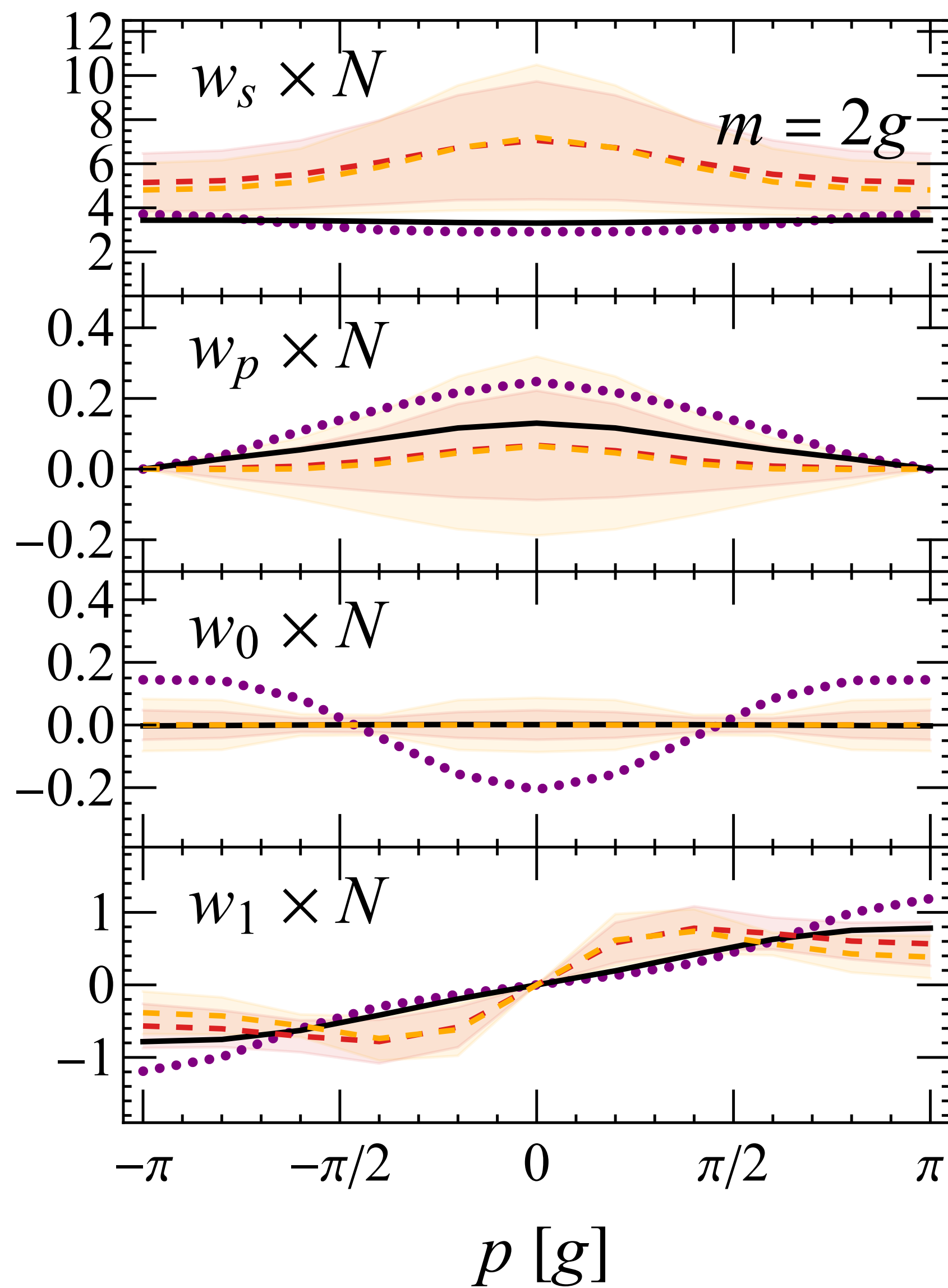
general pure state

thermal

$$|c_n|^2 \approx e^{-\beta E_n} / Z$$

$$\langle O \rangle_{PS} \approx \langle O \rangle_{th}$$

if $\langle E \rangle_{PS} = \langle E \rangle_{th}$



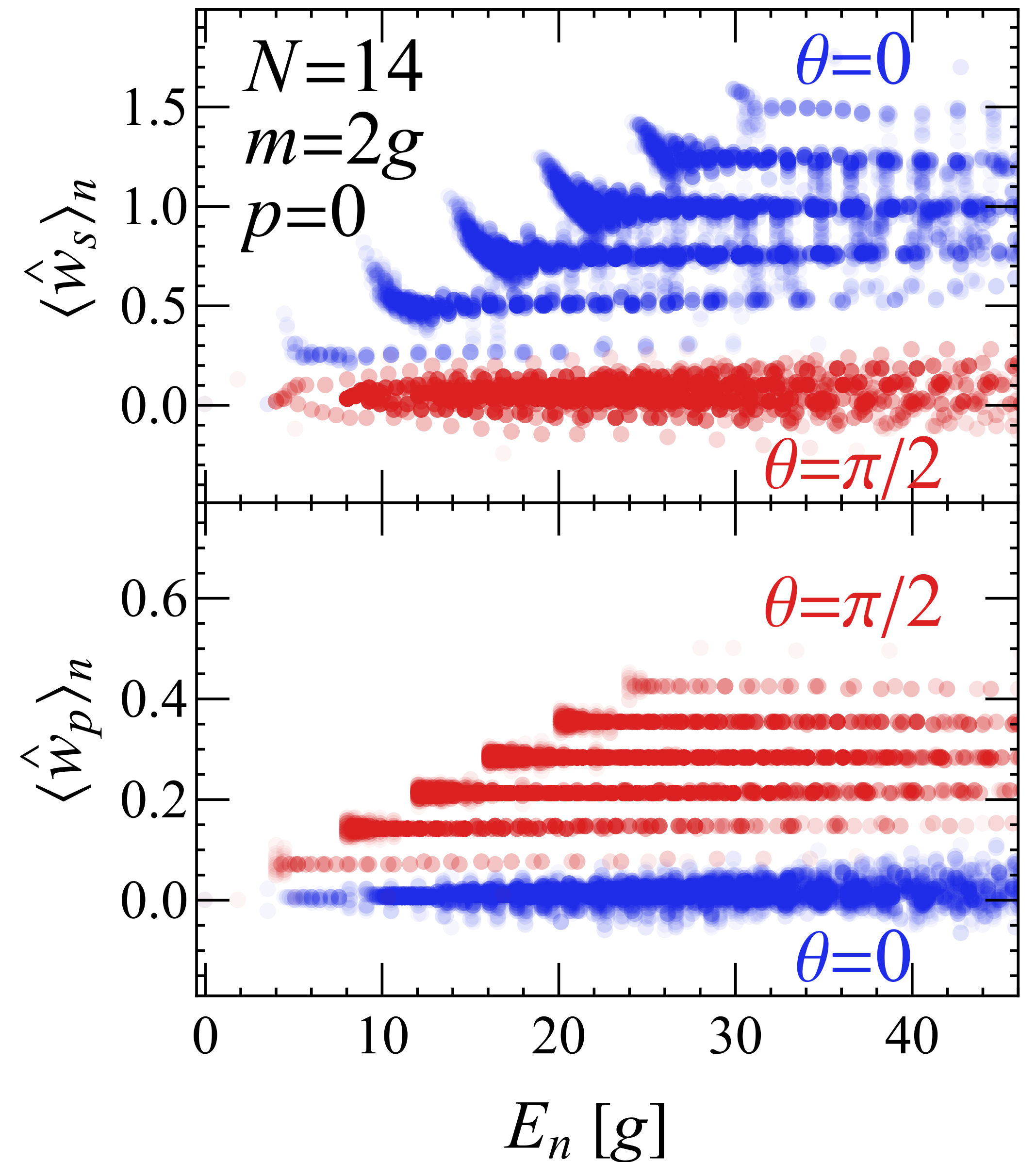
$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - \underline{m e^{i\theta} \gamma^5}) \psi \right) dx$$

$\theta = \frac{\pi}{2}$
 $\theta = 0$

$m \bar{\psi} \gamma^5 \psi$
 $m \bar{\psi} \psi$

$$w_s(x, p) \sim \int_{x-y} e^{ip(x-y)} \bar{\psi}(x) \psi(y)$$

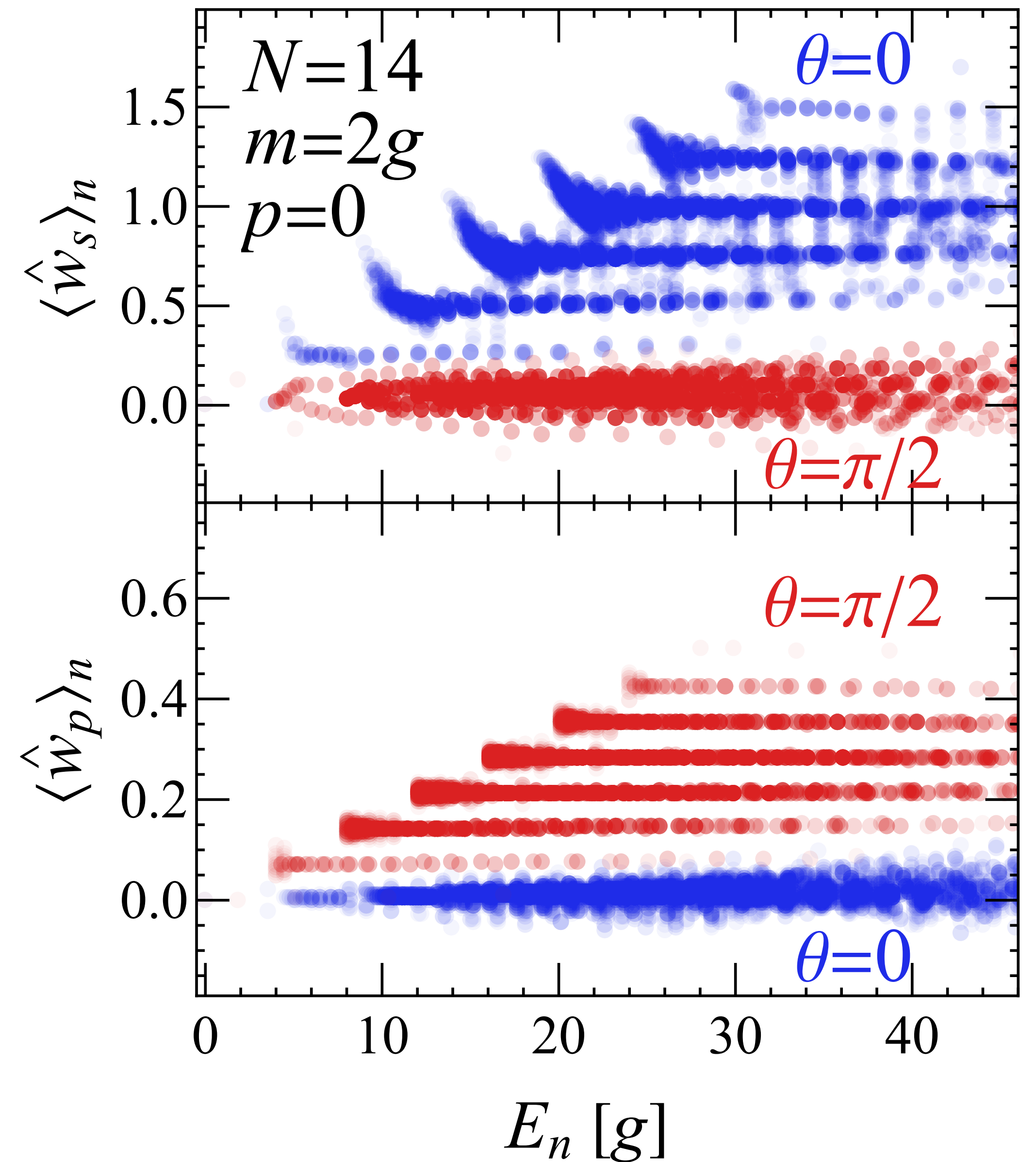
$$w_p(x, p) \sim \int_{x-y} e^{ip(x-y)} \bar{\psi}(x) \gamma^5 \psi(y)$$

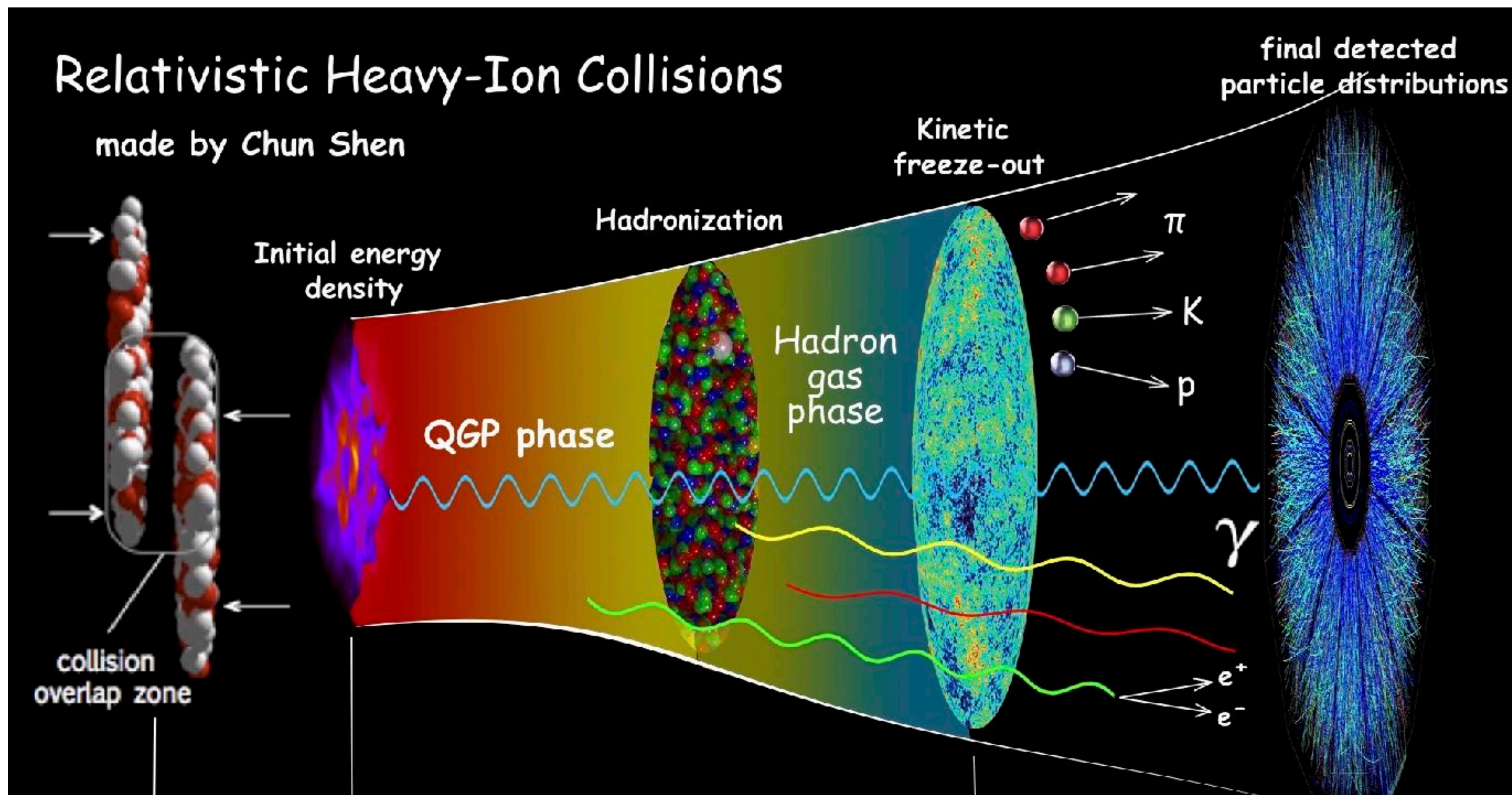




Shile Chen (陈诗乐)

w/ Shile Chen and Li Yan, 2412.00662

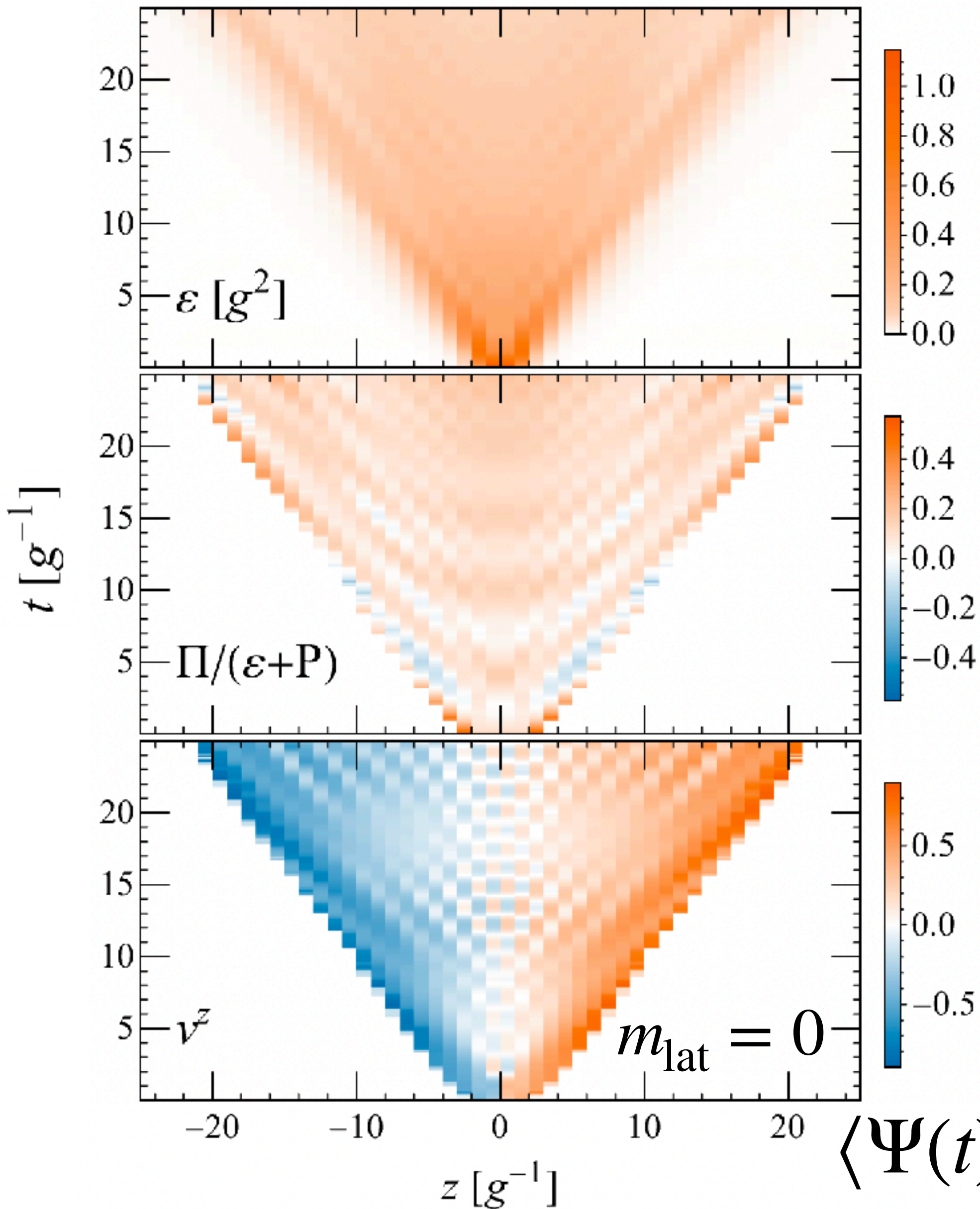




Emergence of Hydrodynamics?

$$\mathcal{L} \Rightarrow \hat{T}^{\mu\nu}$$

$$\langle \Psi(t) | \hat{T}^{\mu\nu}(x) | \Psi(t) \rangle = (\varepsilon + P + \Pi) u^\mu u^\nu - (P + \Pi) g^{\mu\nu}$$

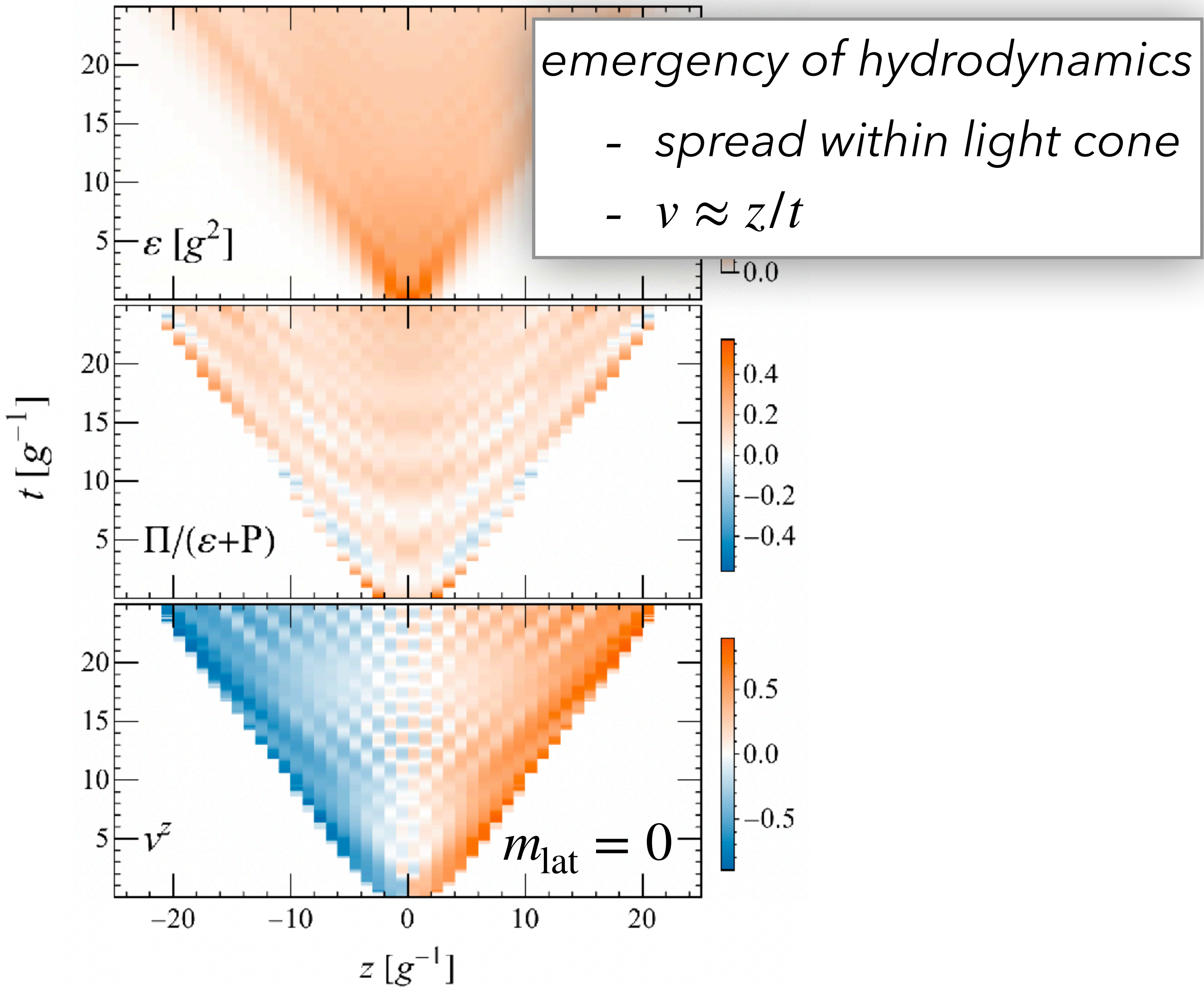


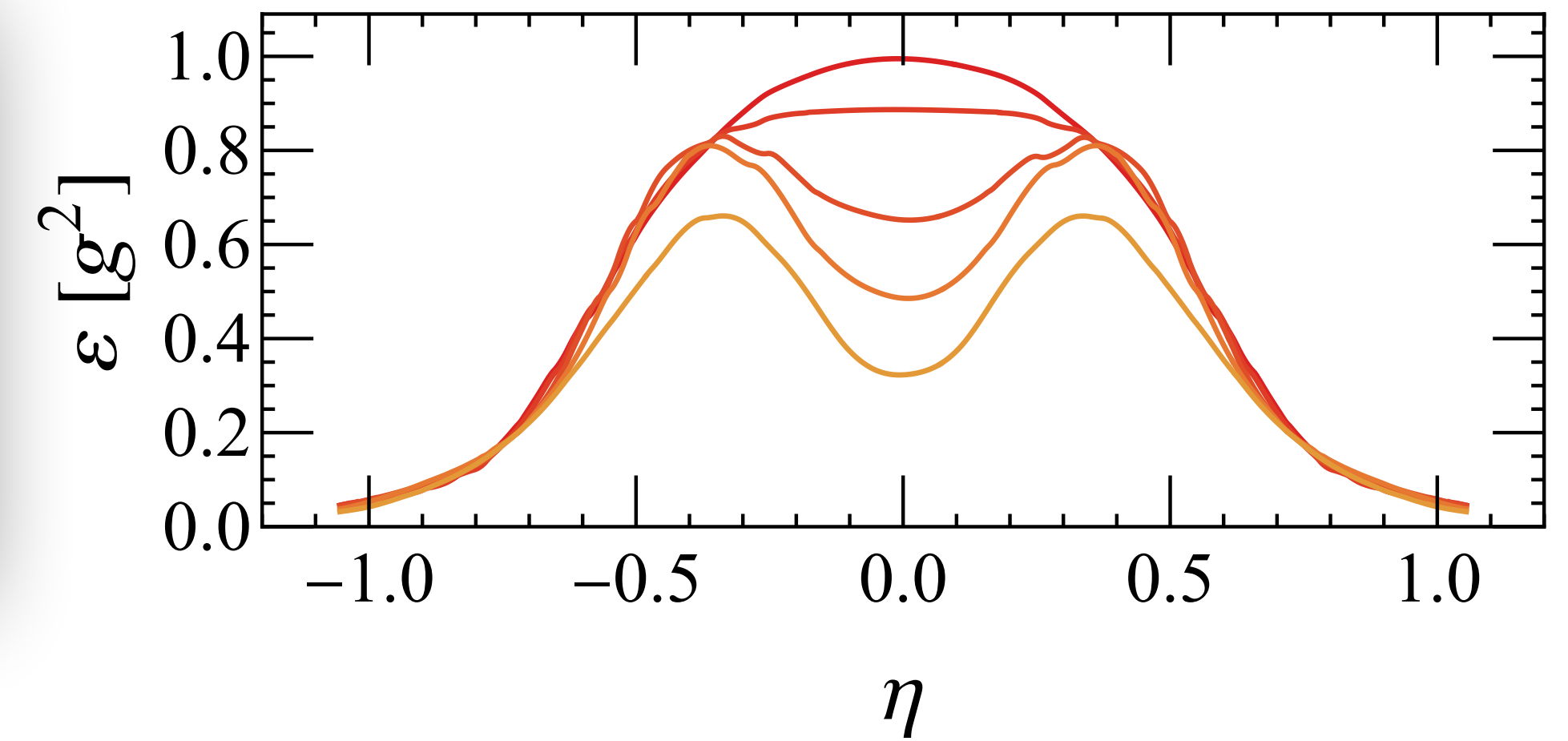
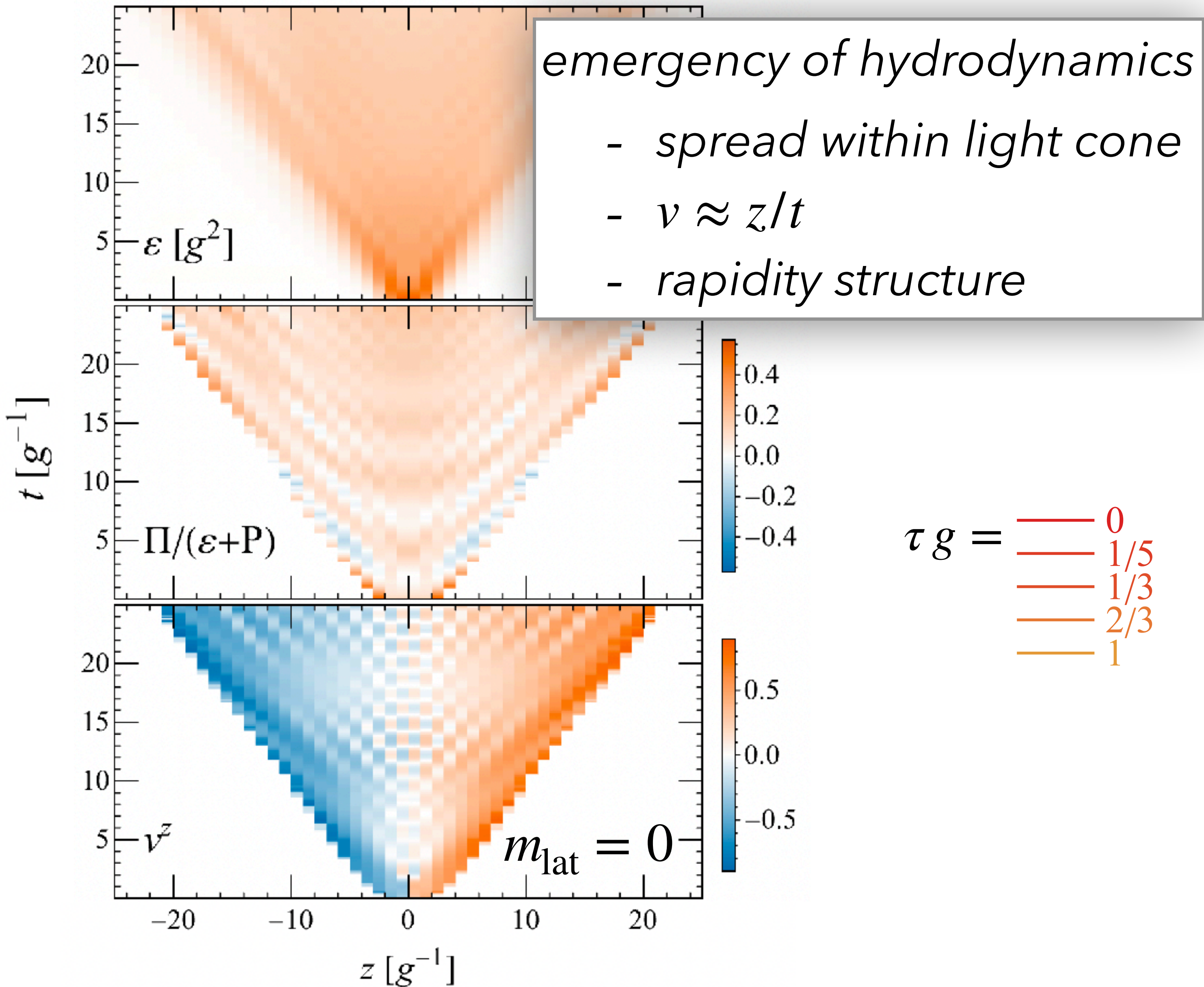
initial state:

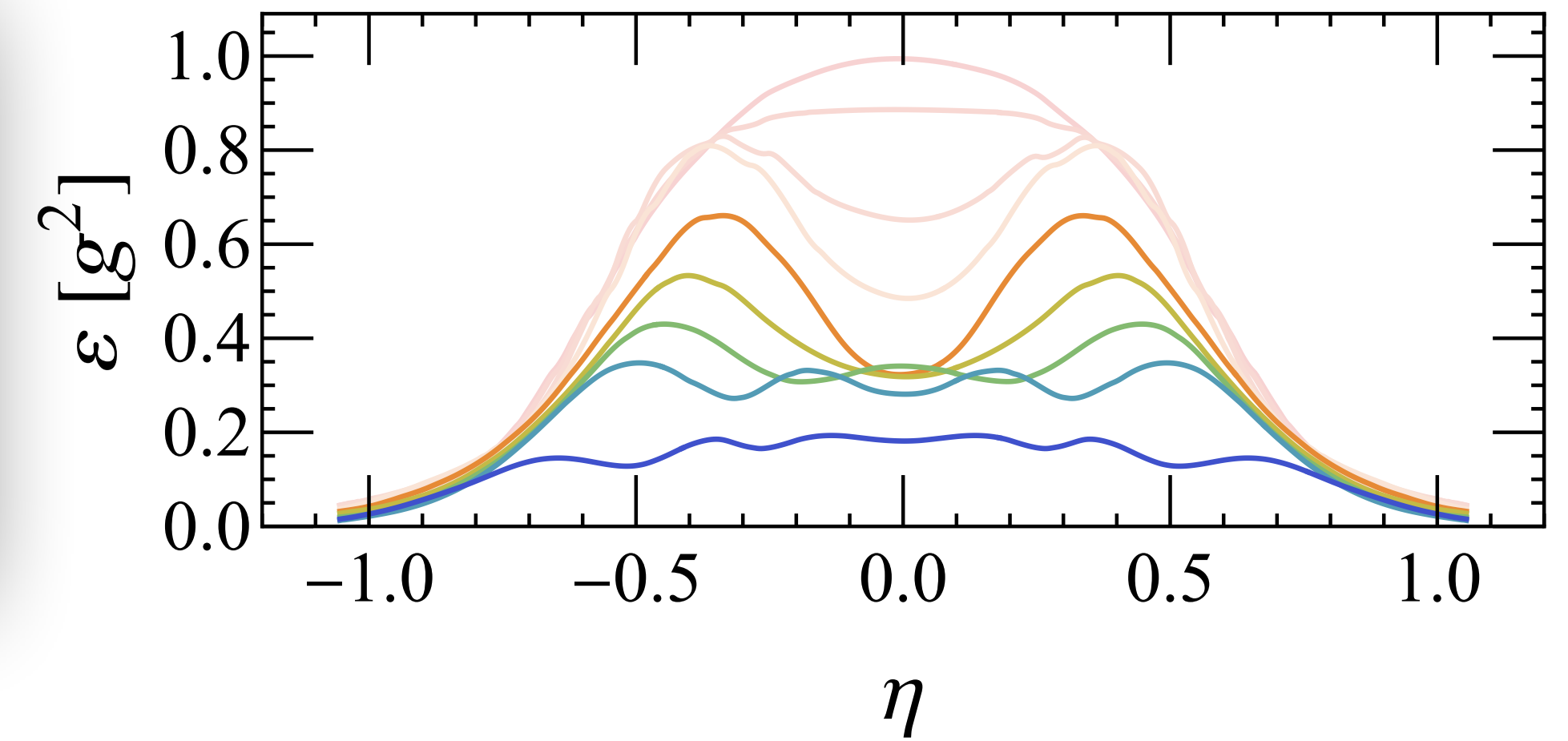
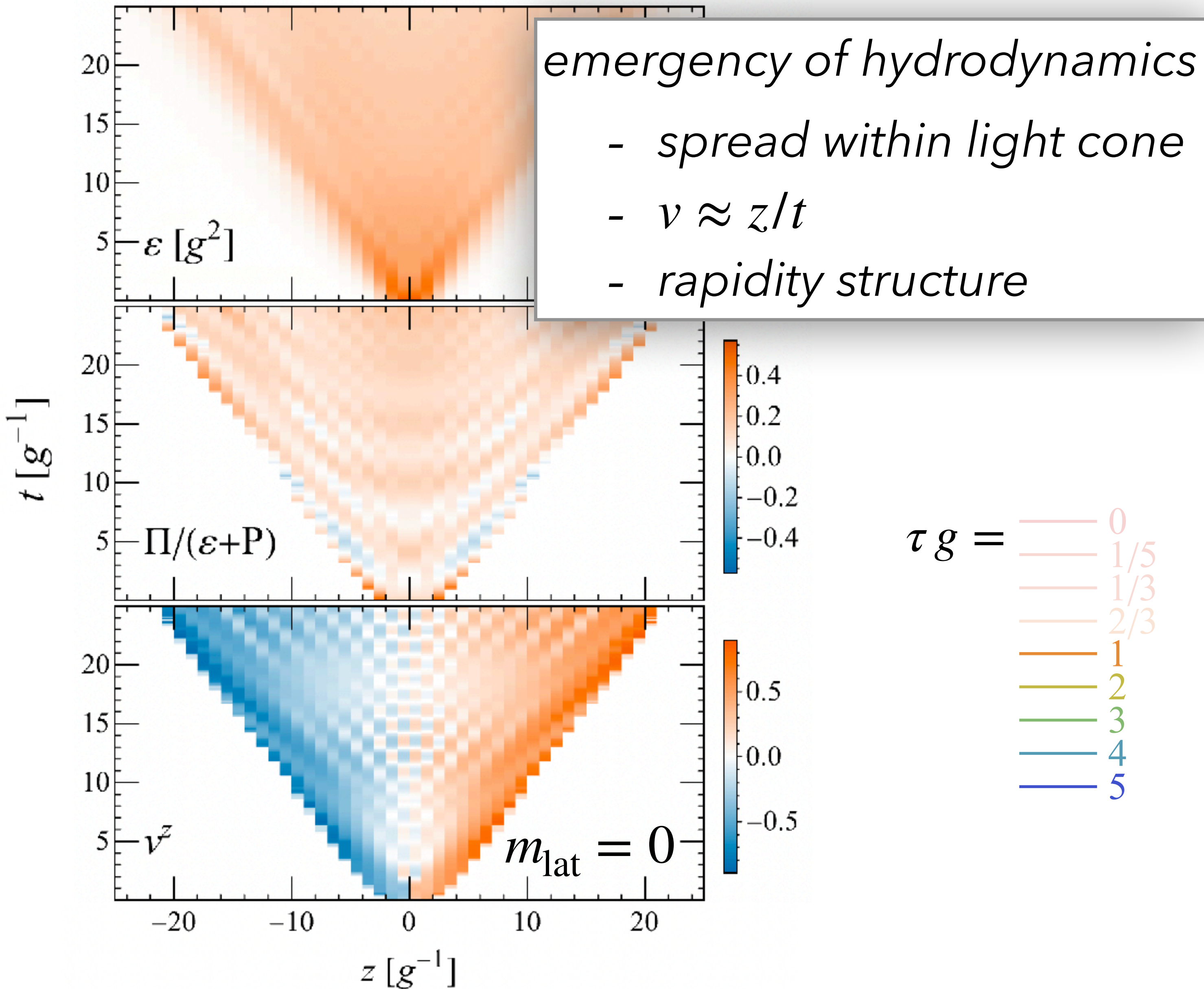
vacuum + [excitation @ center]

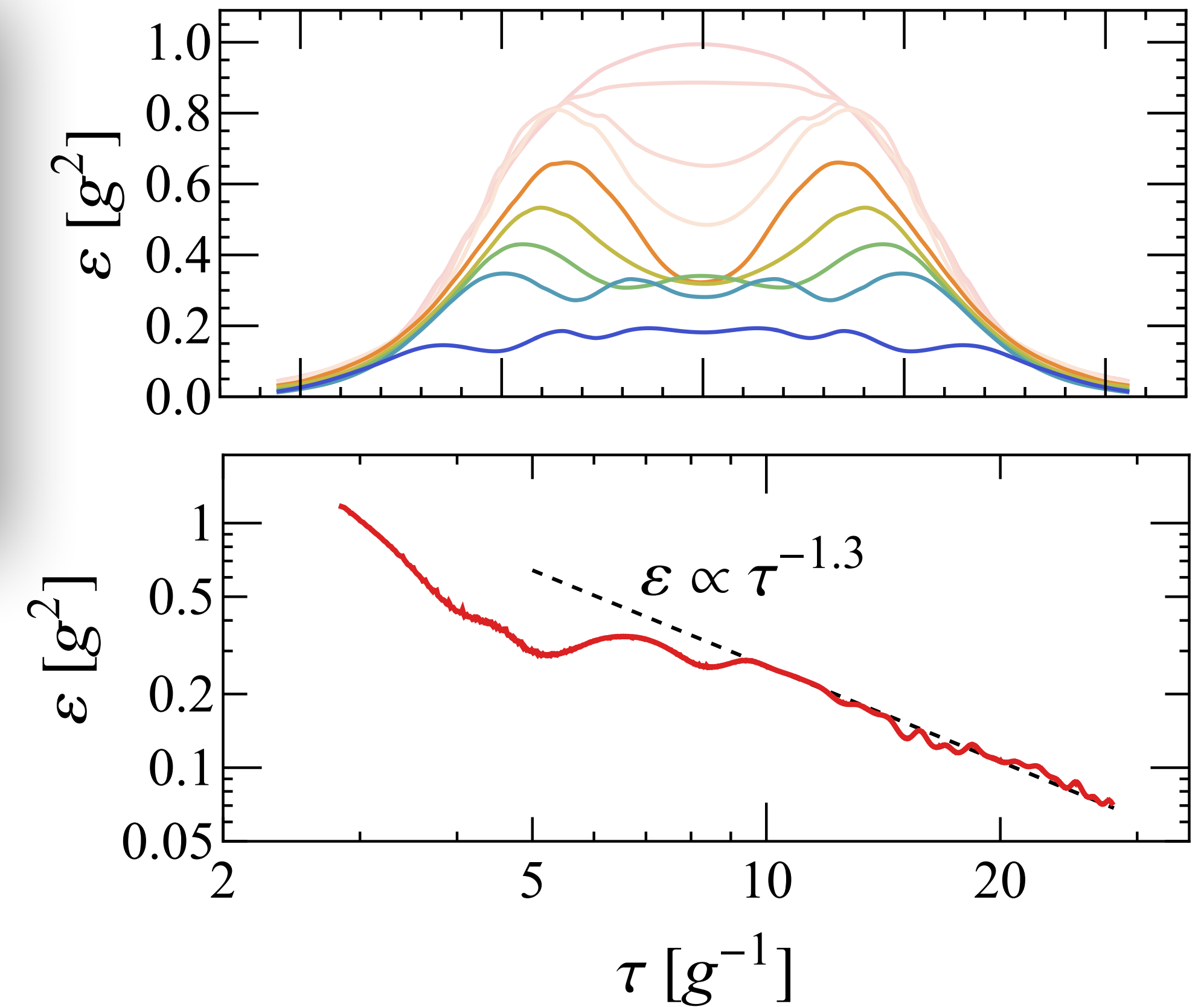
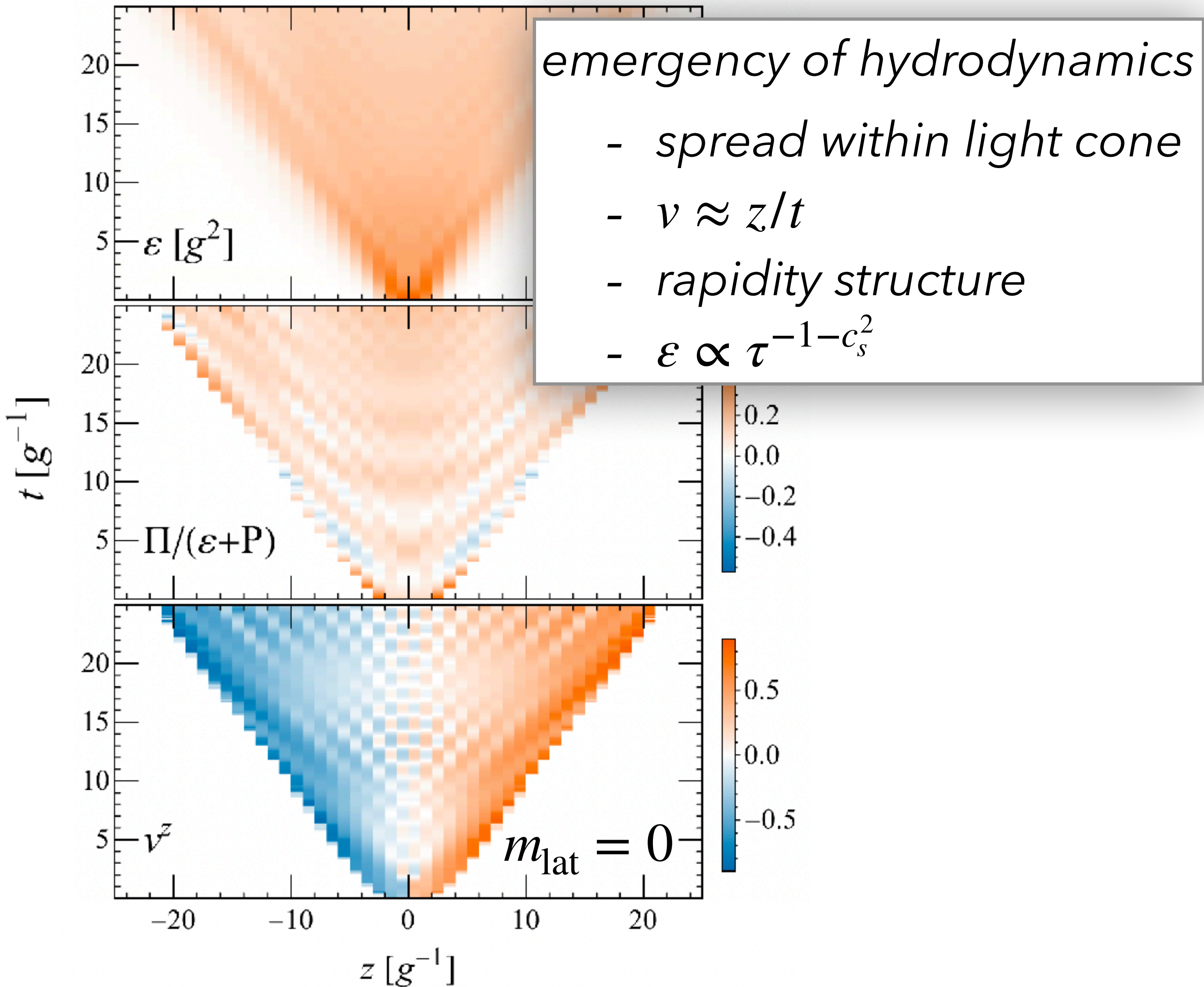


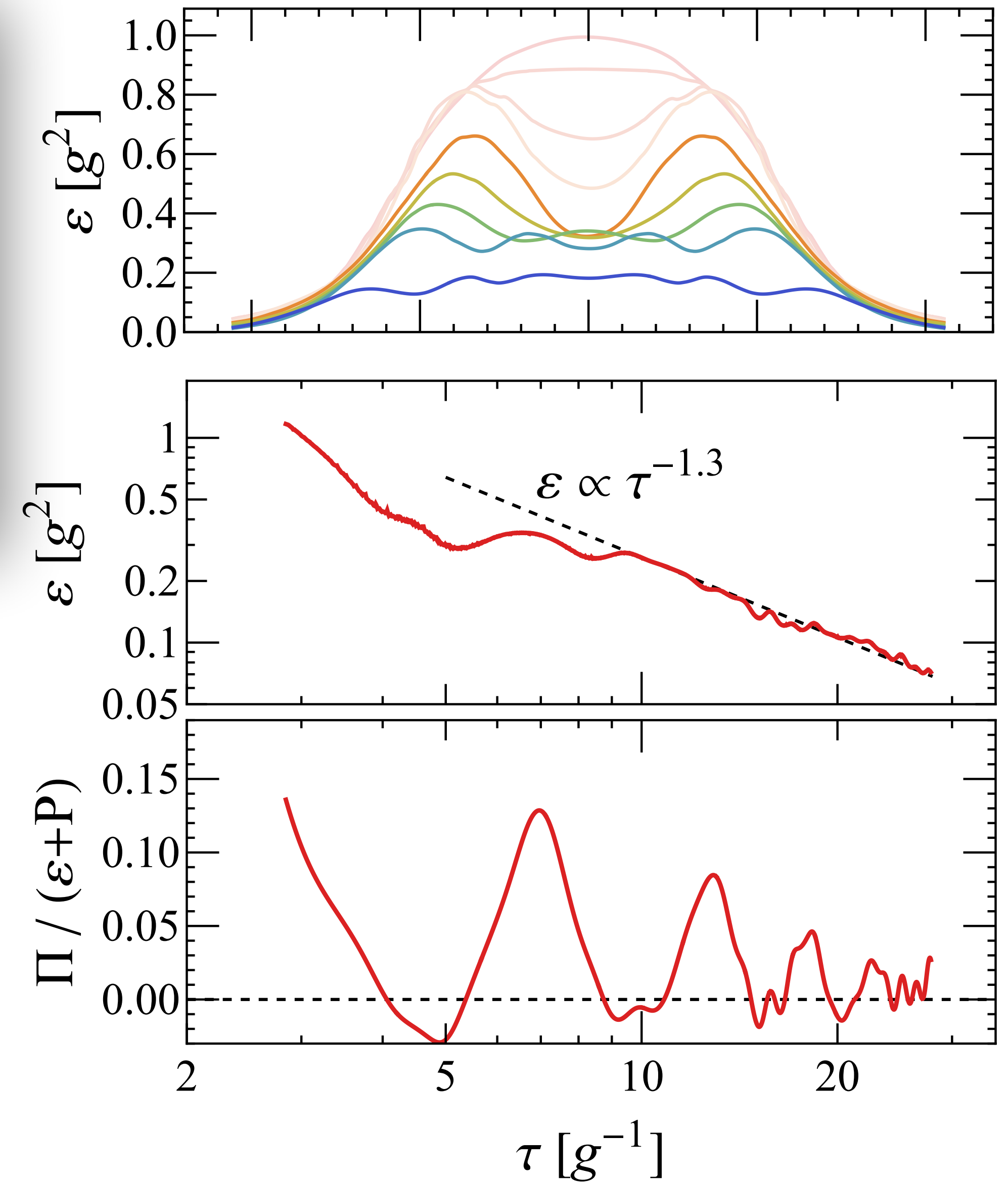
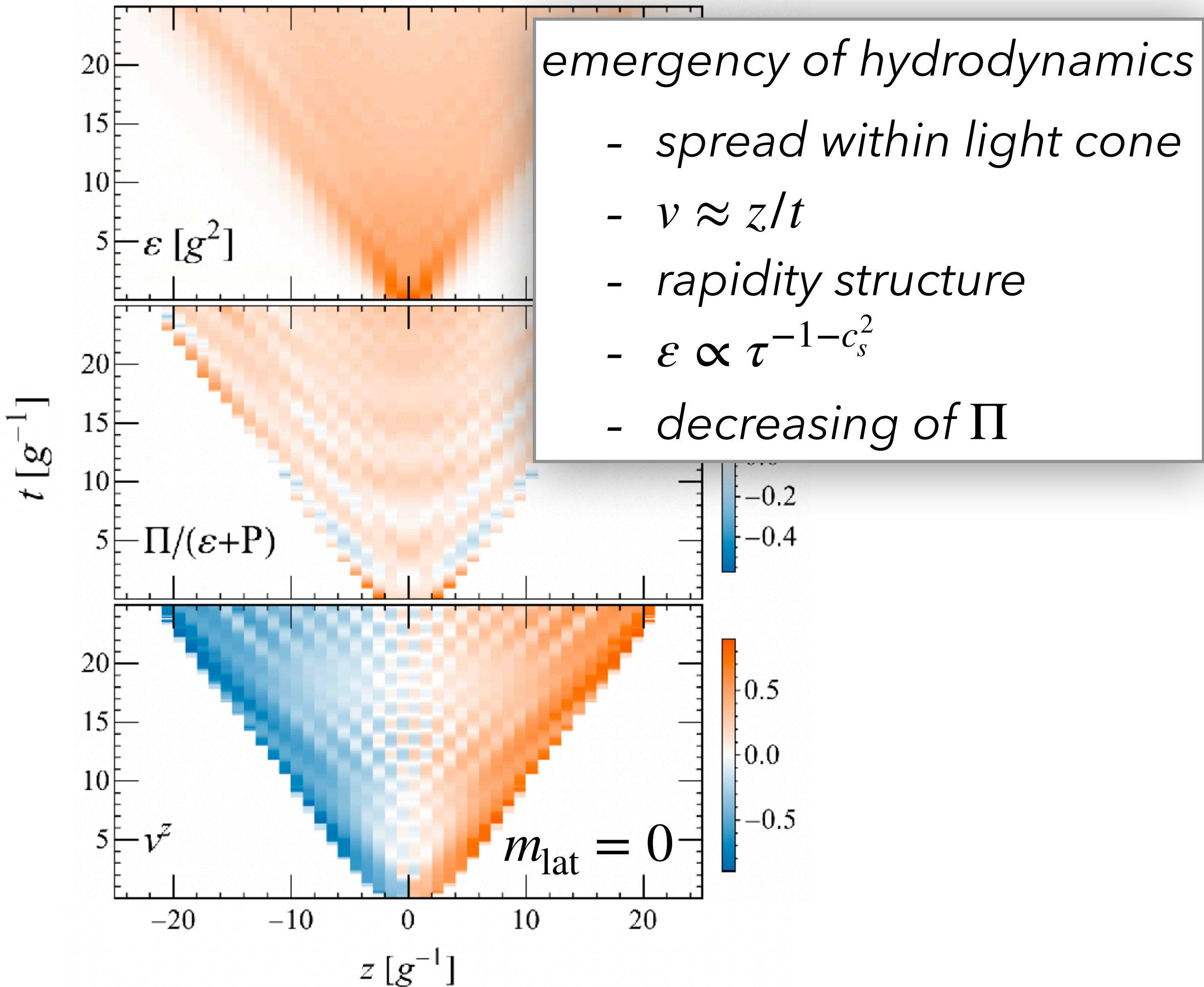
$$\langle \Psi(t) | \hat{T}^{\mu\nu}(x) | \Psi(t) \rangle = (\varepsilon + P + \Pi) u^\mu u^\nu - (P + \Pi) g^{\mu\nu}$$

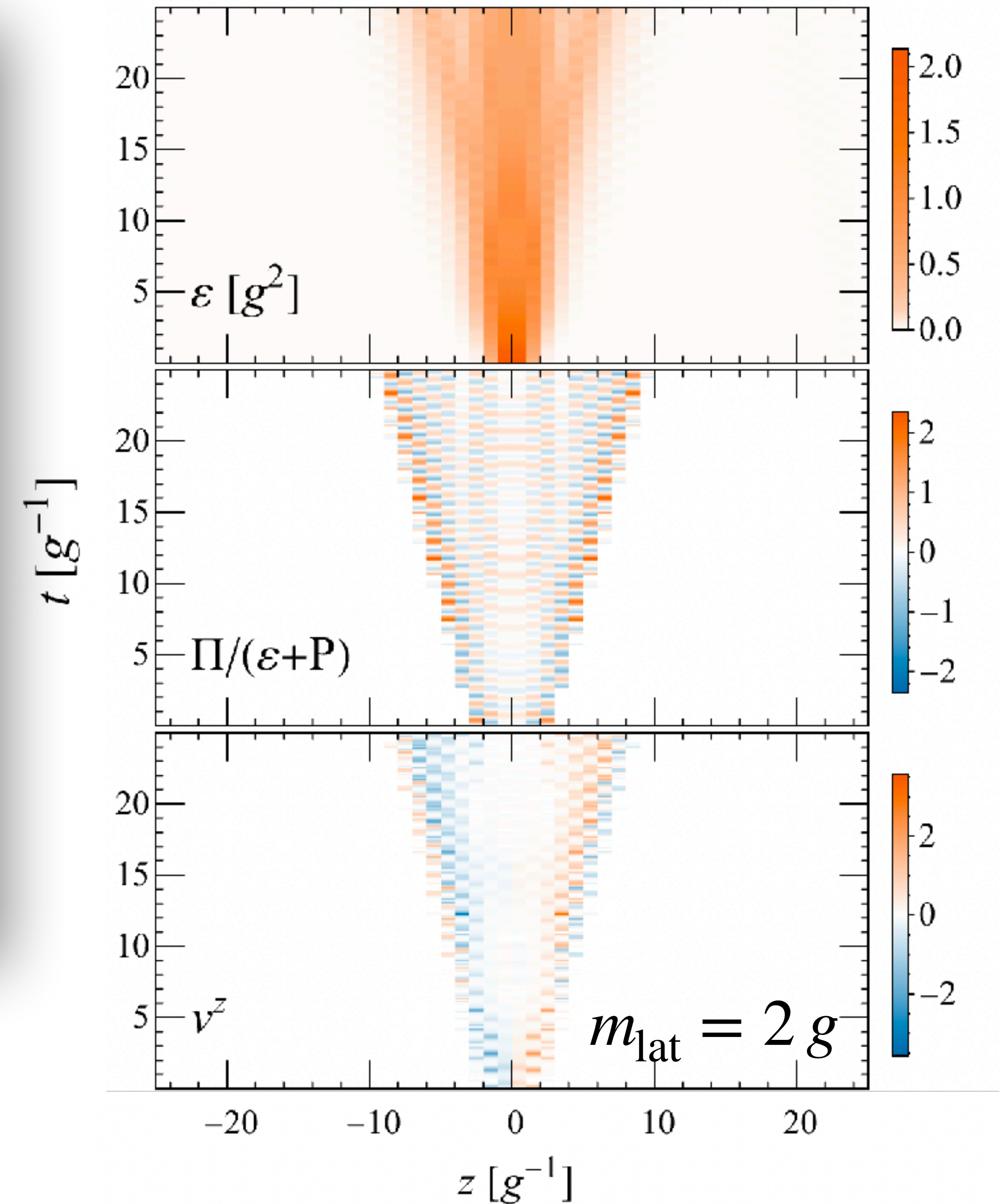
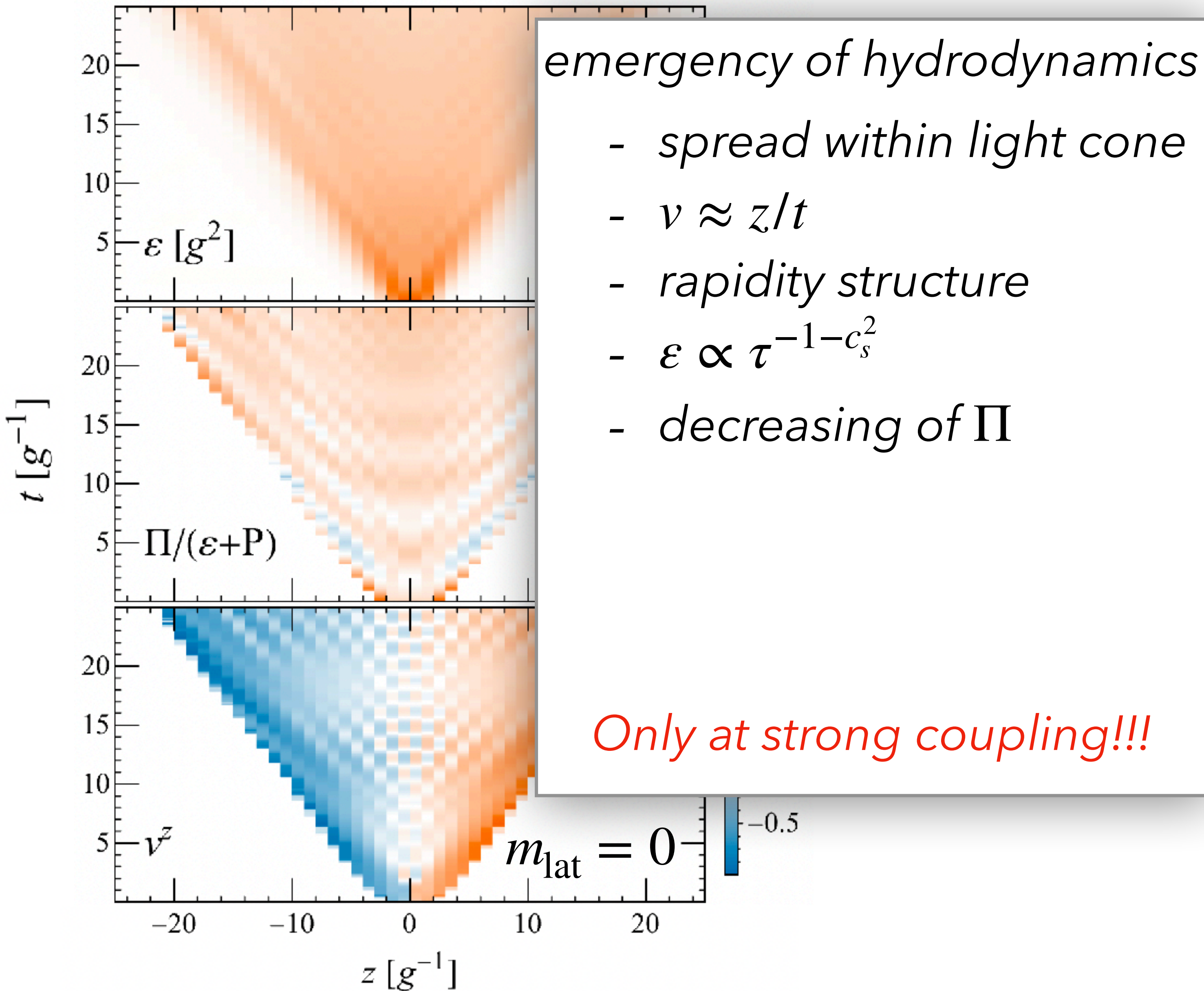


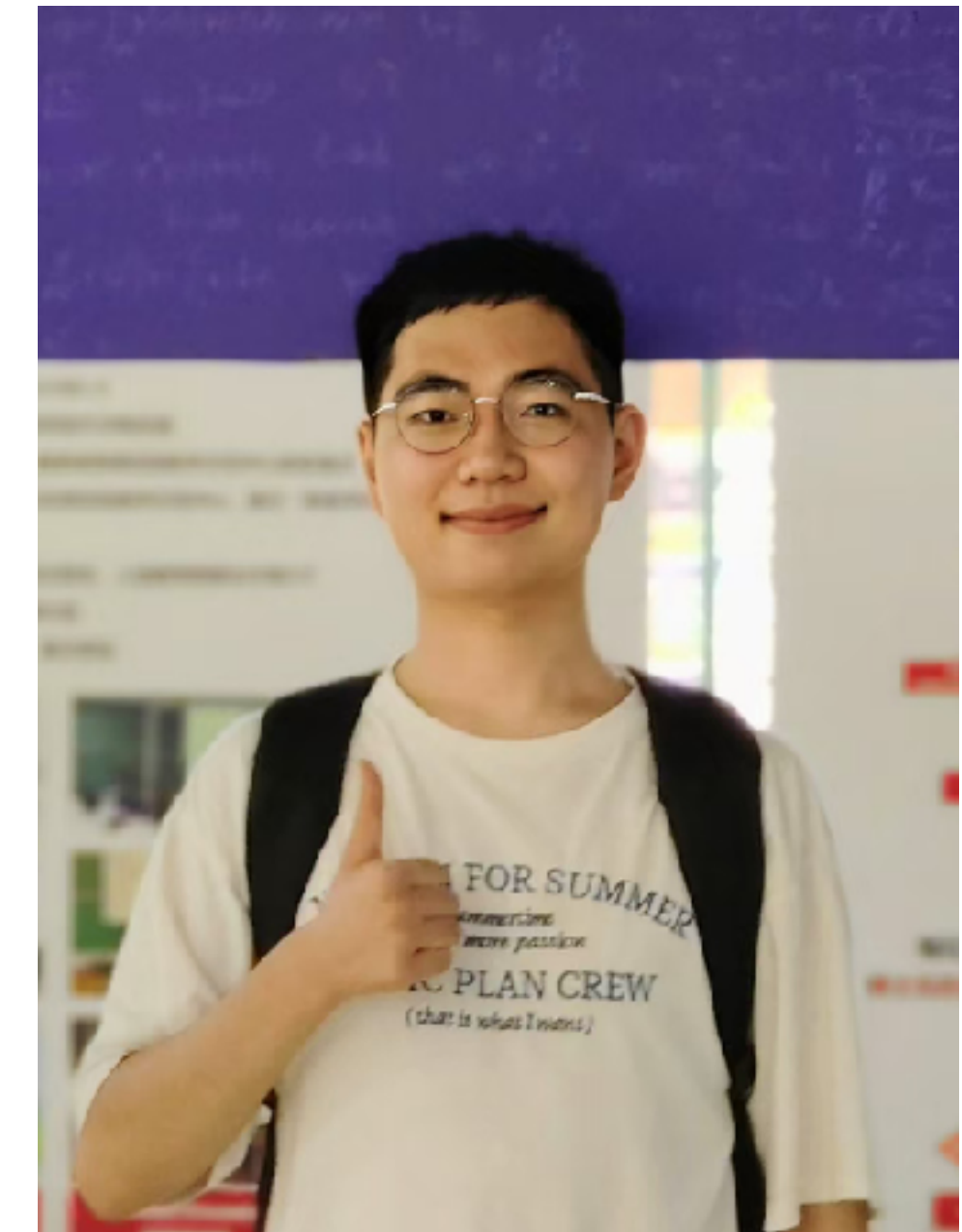
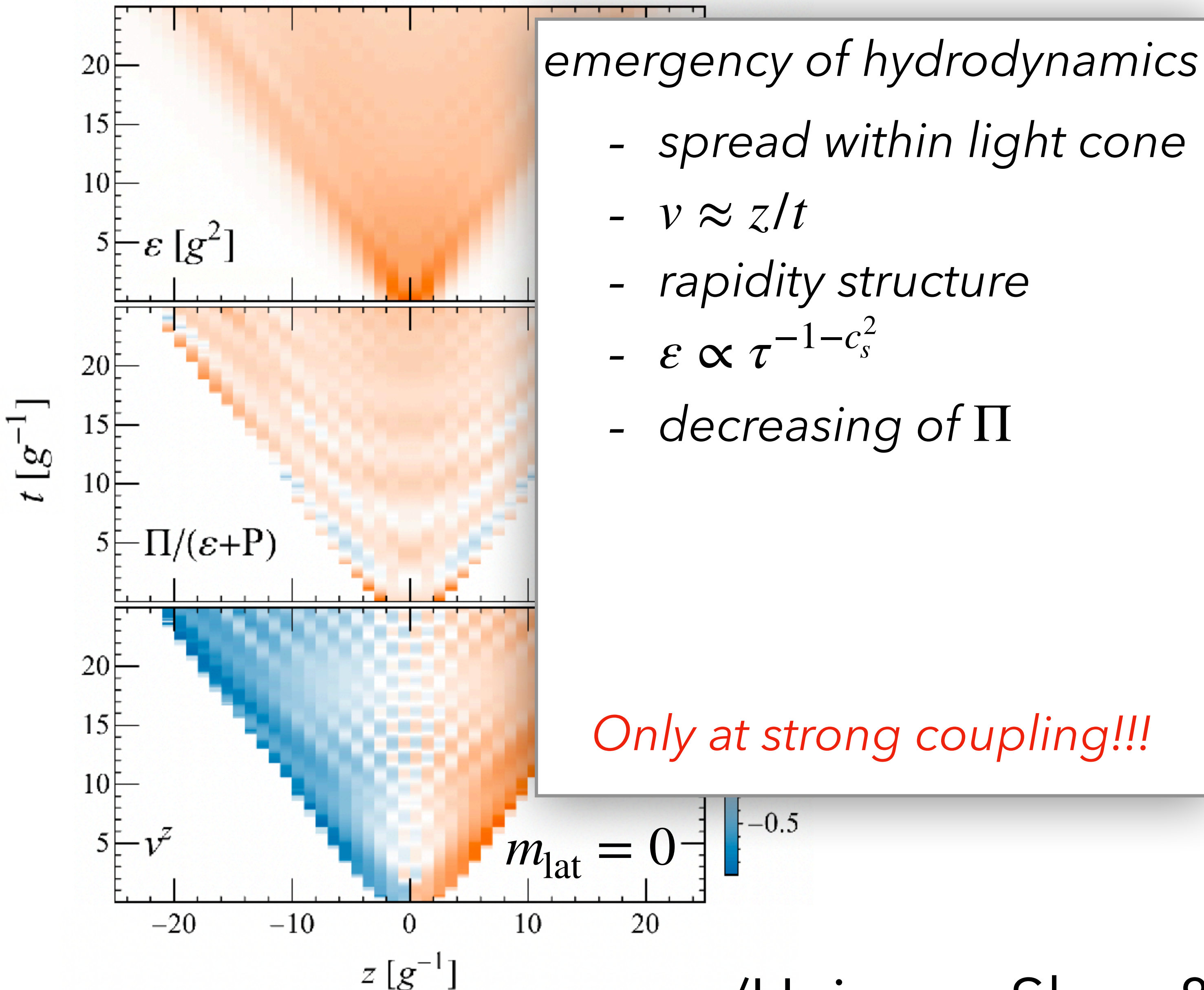












Haiyang Shao(邵海洋)

w/Haiyang Shao & Shile Chen, final preparation

Real-time non-perturbative quantum evolution

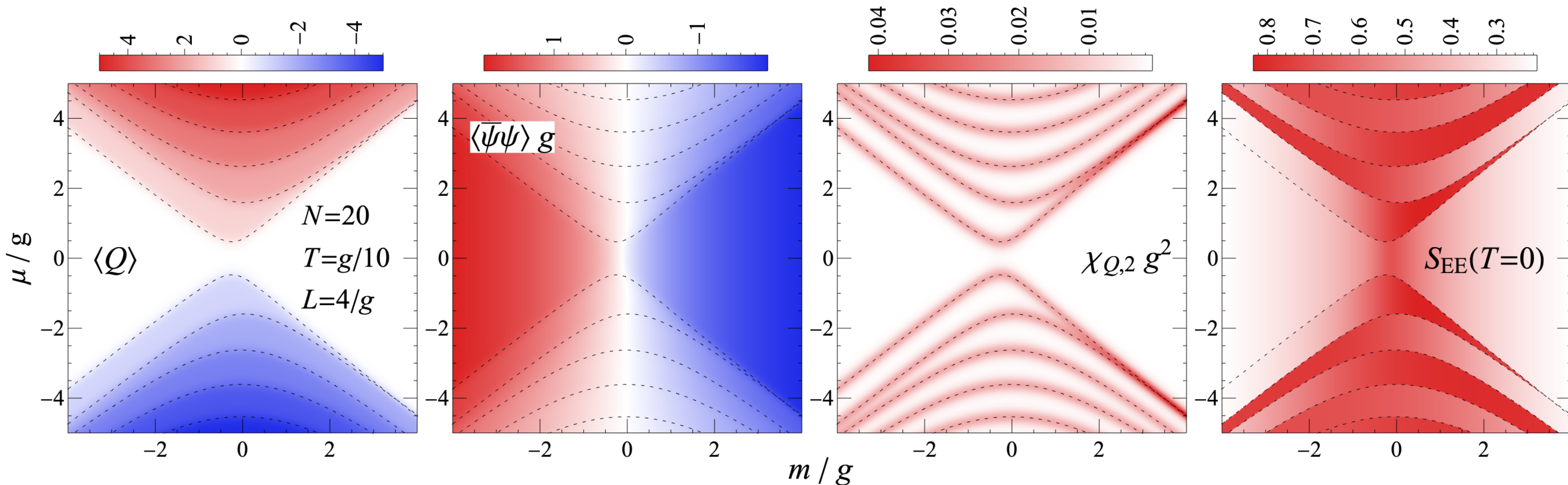
- thermalization of quantum distribution function
 - strong coupling limit: ETH satisfied
 - weak coupling: ETH violated depending on symmetry
- emergence of hydrodynamics

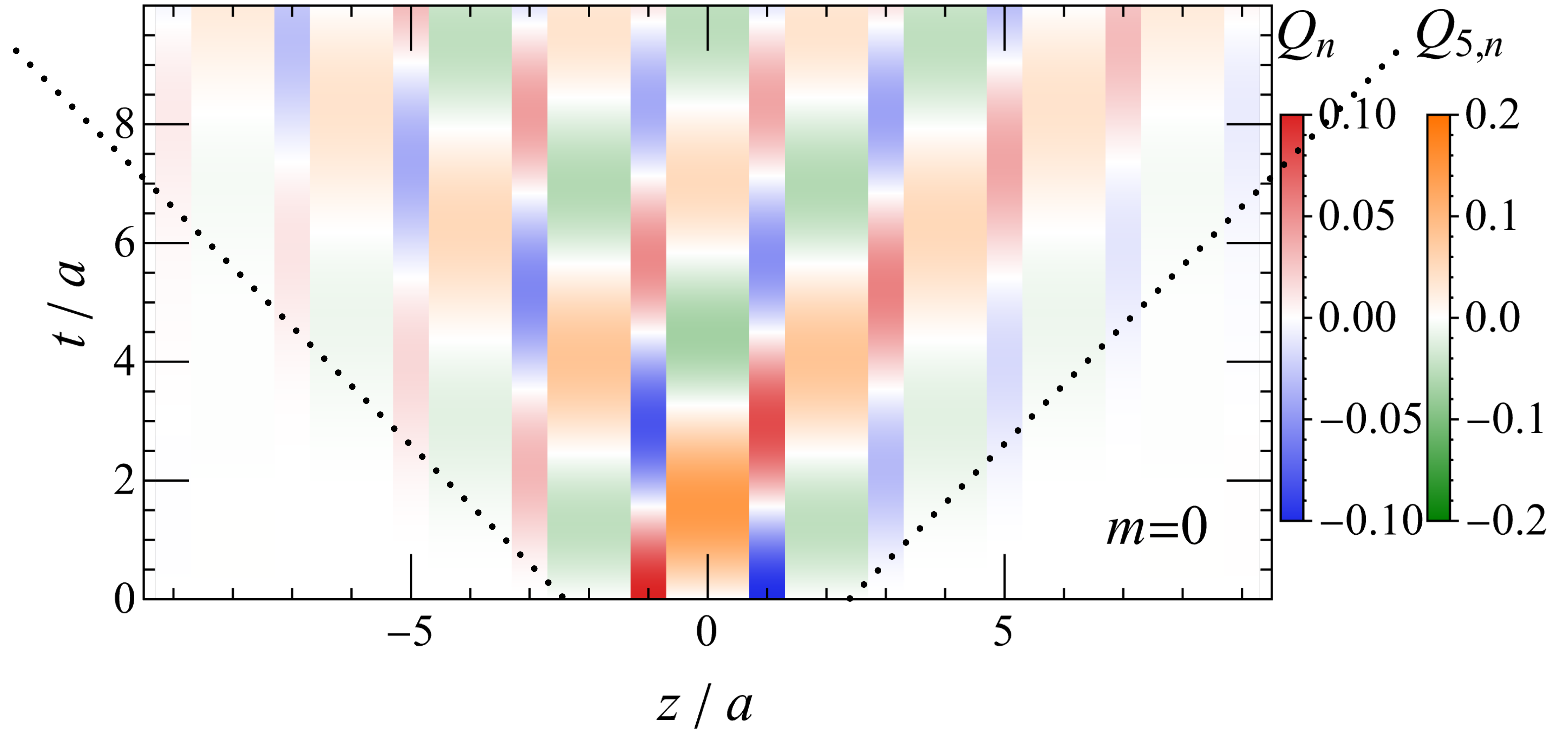
back up slides

- finite temperature, finite chemical potential: $e^{-(H-\mu Q)/T}$

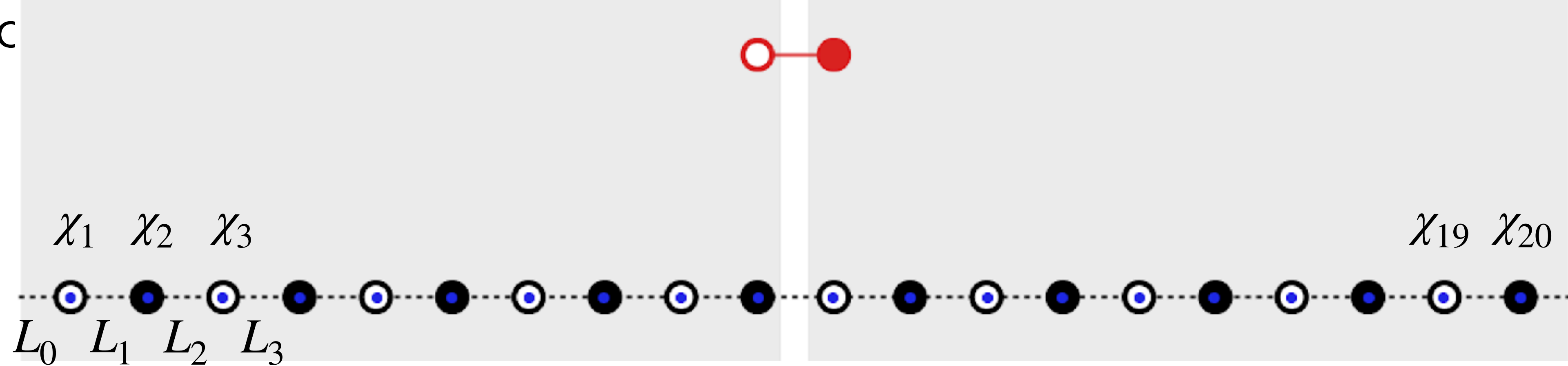
$$\langle O \rangle_{\text{th}} \equiv \text{Tr}(\rho_{\text{th}} O)$$

$$\rho_{\text{th}} \equiv \frac{e^{-(H-\mu Q)/T}}{\text{Tr}(e^{-(H-\mu Q)/T})}$$





1+1D Sc



discretize and matrix(gate) representation:

staggered fermion that satisfied anti-commutation: $\{\psi_a(x), \psi_b^\dagger(y)\} = \delta_{a,b}\delta(x-y)$

gauge field fixed by Gauss' law: $\partial_1 E - g \bar{\psi} \gamma^0 \psi = j_{\text{ext}}^0$

$$E(x = an)$$

$\leftrightarrow$

$$L_n$$

$$L_n - L_{n-1} - \frac{Z_n + (-1)^n}{2} = \frac{1}{g} \int_{(n-1/2)a}^{(n+1/2)a} dx j_{\text{ext}}^0(x, t) ,$$

thermal equilibrium property

$$\hat{\rho}_{\text{th}} \equiv Z^{-1} e^{-(\hat{H} - \mu \hat{Q})/T}$$

$$\langle \hat{O} \rangle_{\text{th}} \equiv \text{Tr}(\hat{\rho}_{\text{th}} \hat{O})$$

real time evolution

$$\partial_t \hat{\rho} = -i[\hat{H}, \hat{\rho}]$$

$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i \hat{H} |\psi(t)\rangle$$

$$q_{n,t} \equiv \langle \psi^\dagger(a n) \psi(a n) \rangle_t = \frac{\langle Z_n \rangle_t + (-1)^n}{2a},$$

$$\nu_{n,t} \equiv \langle \bar{\psi}(a n) \psi(a n) \rangle_t = \frac{(-1)^n \langle Z_n \rangle_t}{2a},$$

$$\Pi_{n,t} \equiv \langle E(a n) \rangle_t = g \langle L_n \rangle_t,$$

⋮

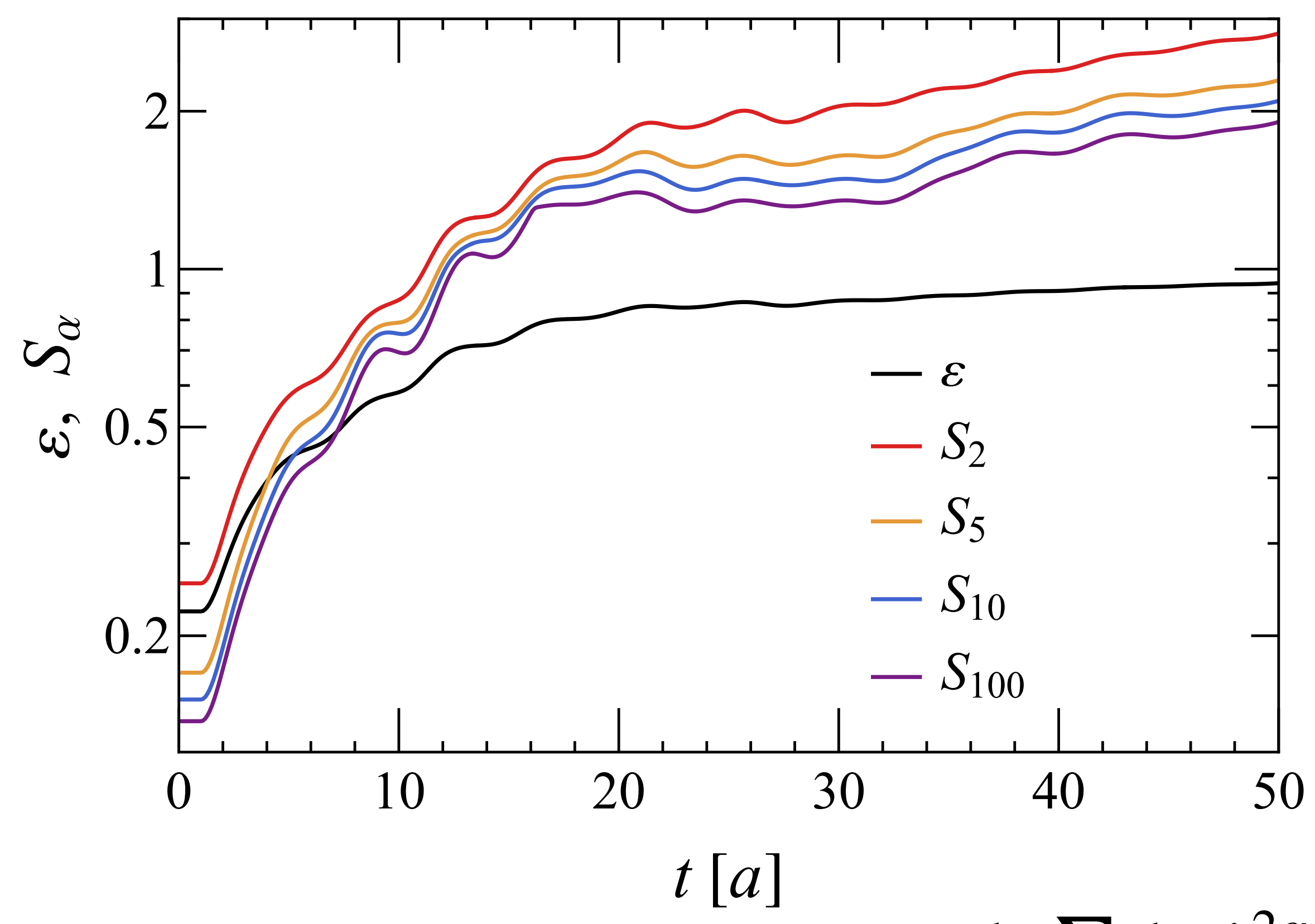
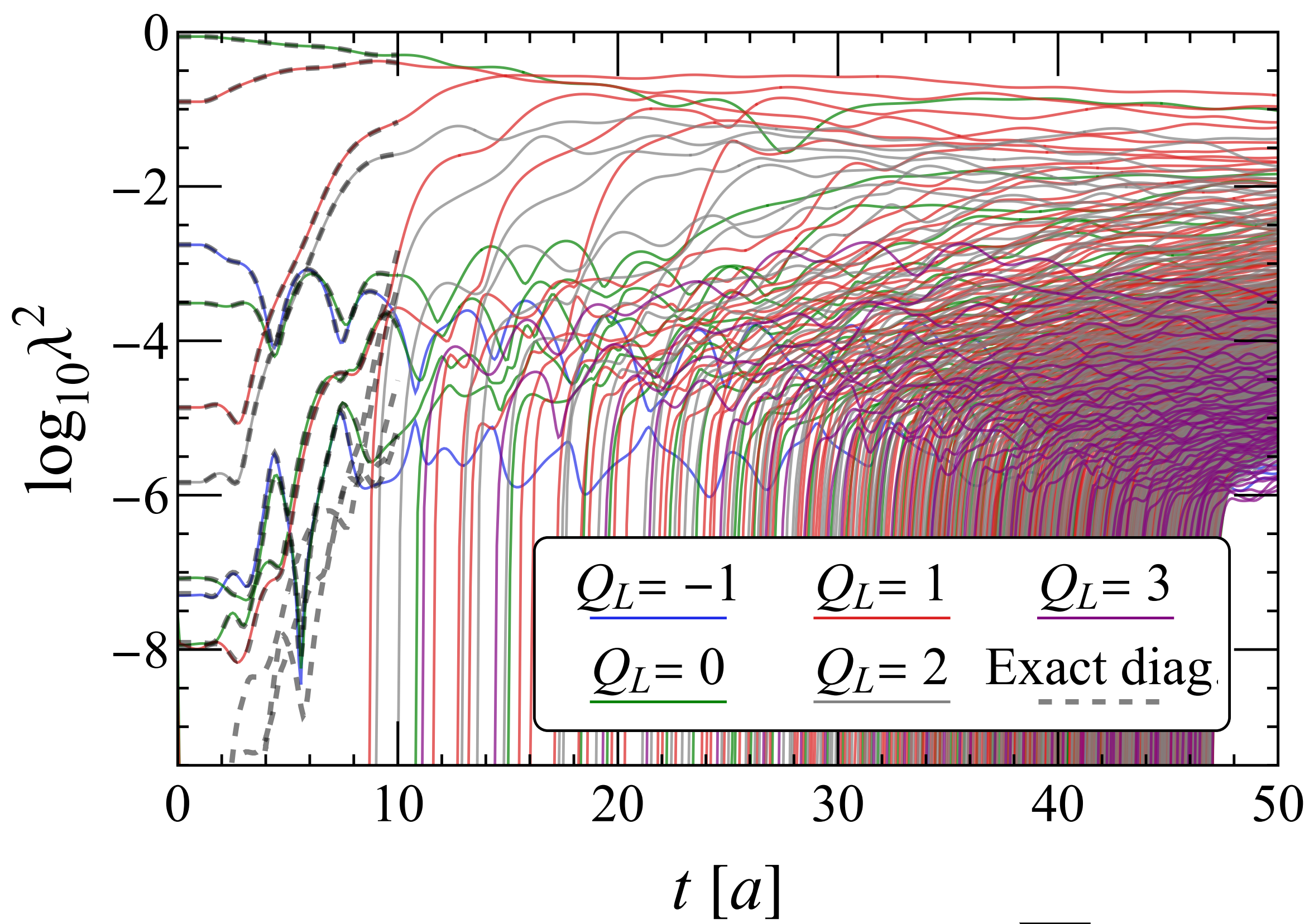
dimension of state vector = 2^N

N : number of lattice sides

dimension of Hamiltonian = ~~$2^N \times 2^N$~~ sparse $\sim 2N \times 2^N$

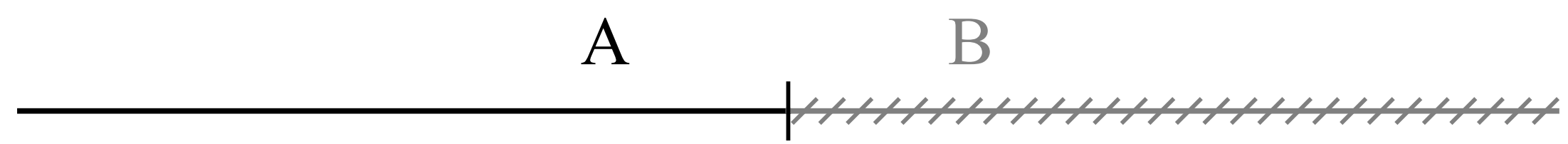
N	dimension	memory of Hamiltonian	# of qubit (N)
8	256	~ 131 kB	8
12	4,096	~ 3.1 MB	12
16	65,536	~ 67 MB	16
20	1,048,576	~ 1.3 GB	20
24	16,777,216	~ 26 classical hardware in this work	
28	268,435,456	~ 481 GB	28

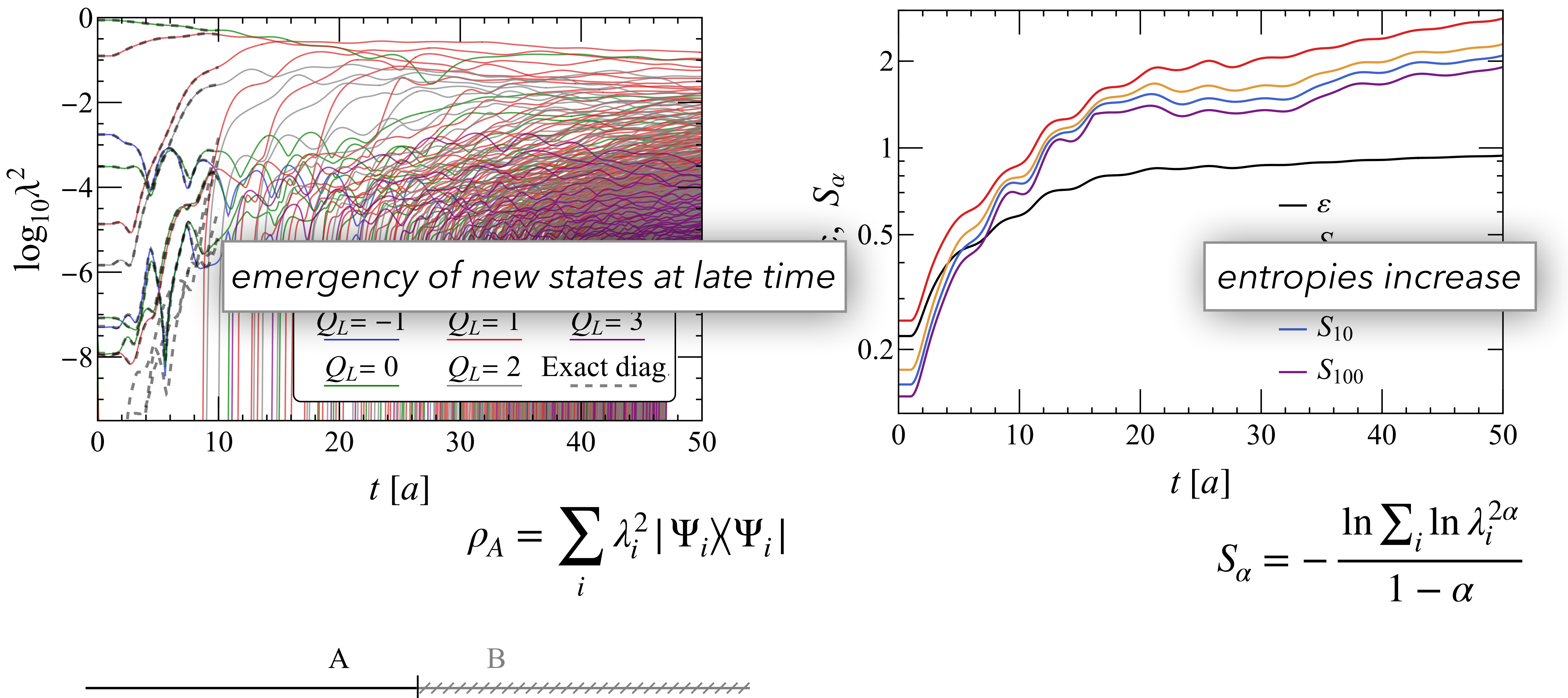
performance not satisfying...

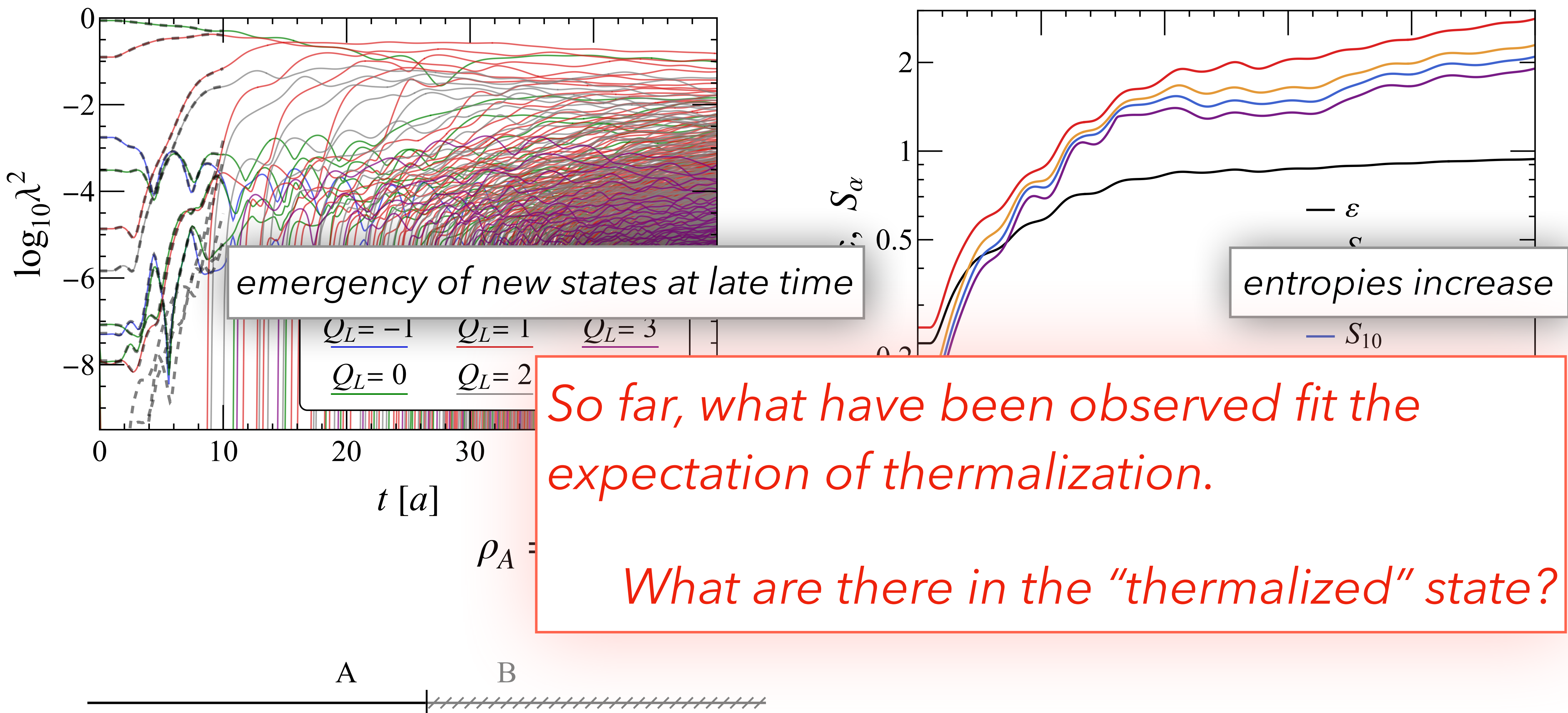


$$\rho_A = \sum_i \lambda_i^2 |\Psi_i\rangle\langle\Psi_i|$$

$$S_\alpha = - \frac{\ln \sum_i \ln \lambda_i^{2\alpha}}{1 - \alpha}$$

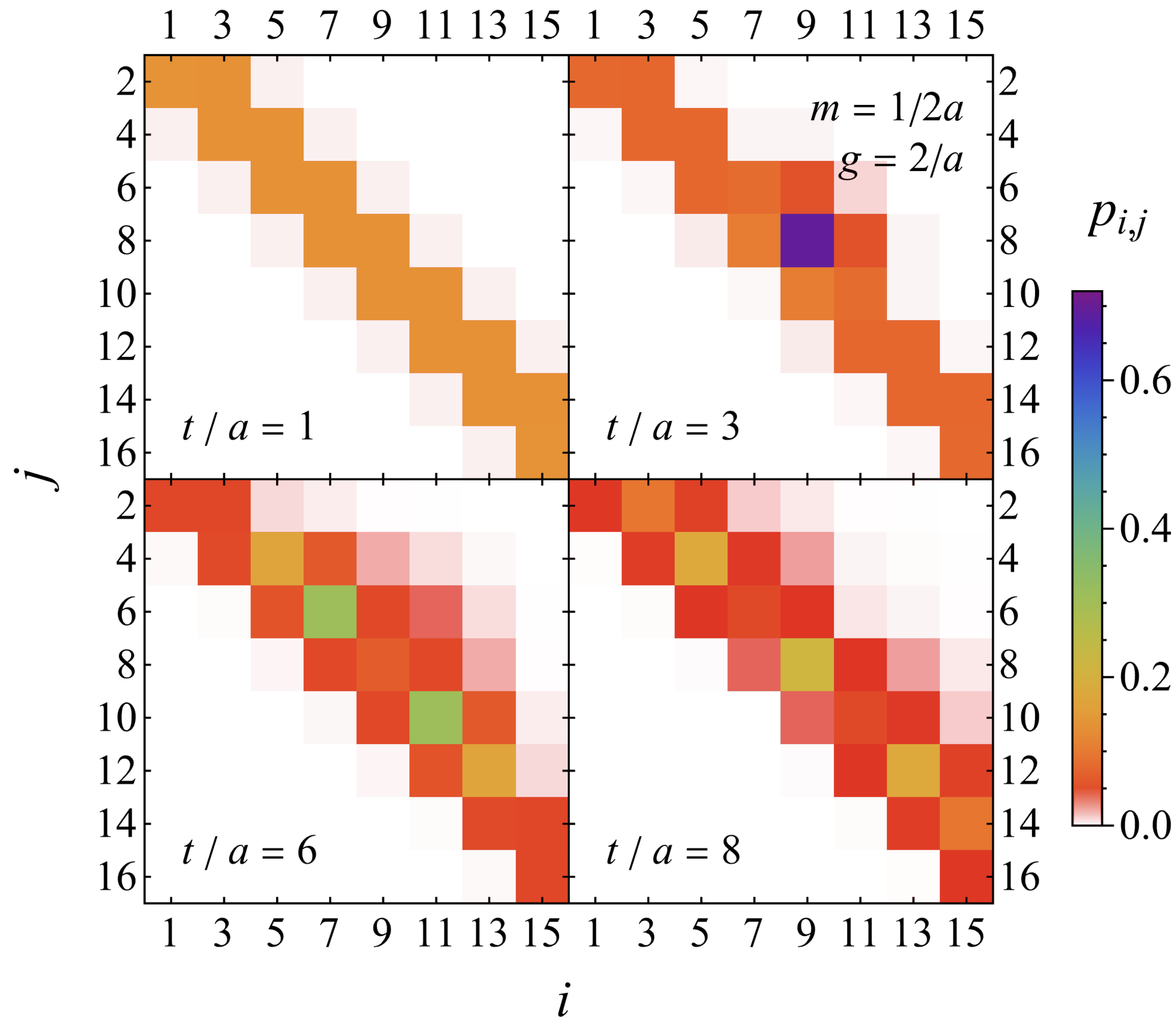






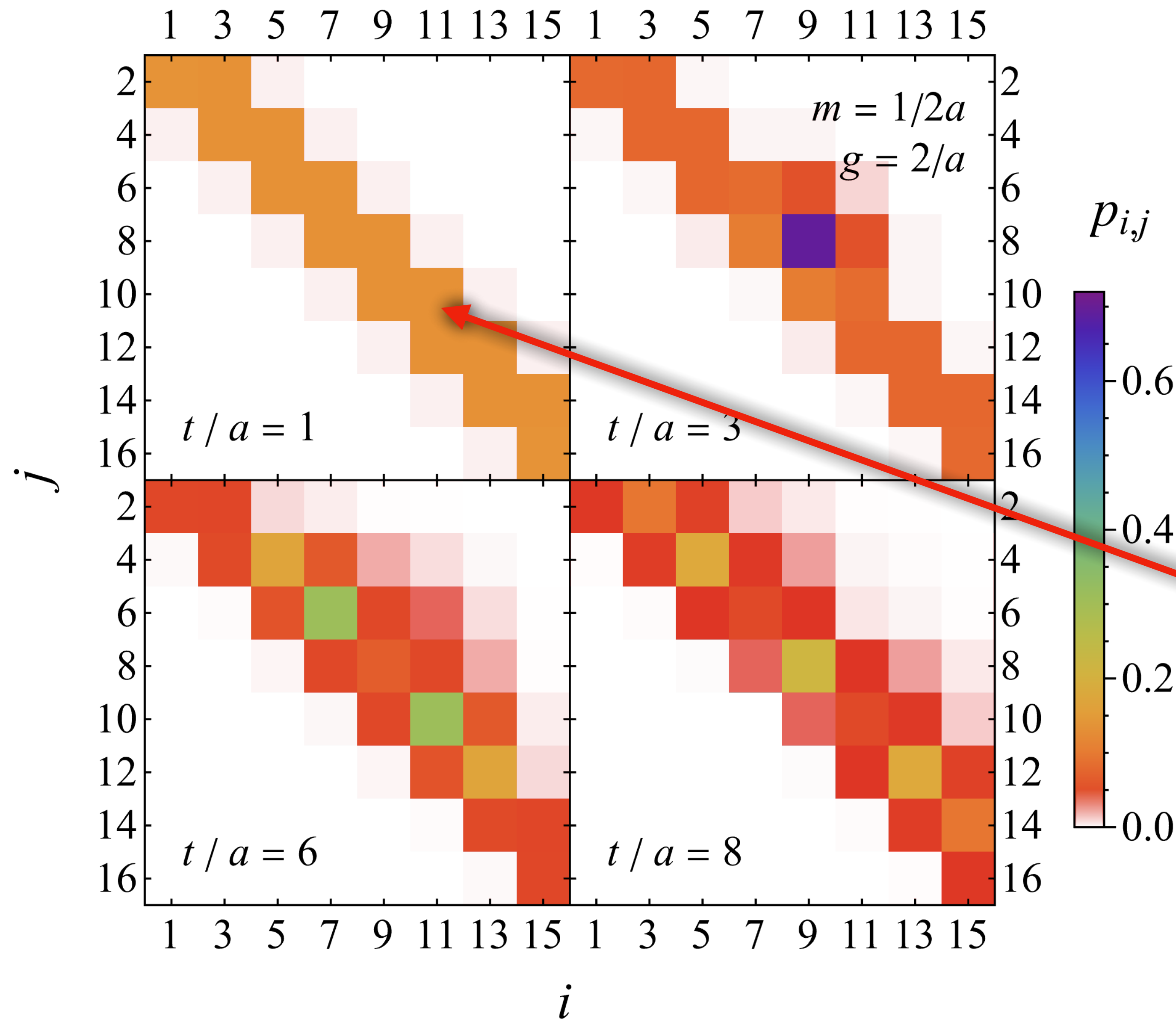
So far, what have been observed fit the expectation of thermalization.

What are there in the "thermalized" state?



$$p_{i,j} = |\langle \Psi_t | \chi_i \chi_j^\dagger | \text{empty} \rangle|^2$$

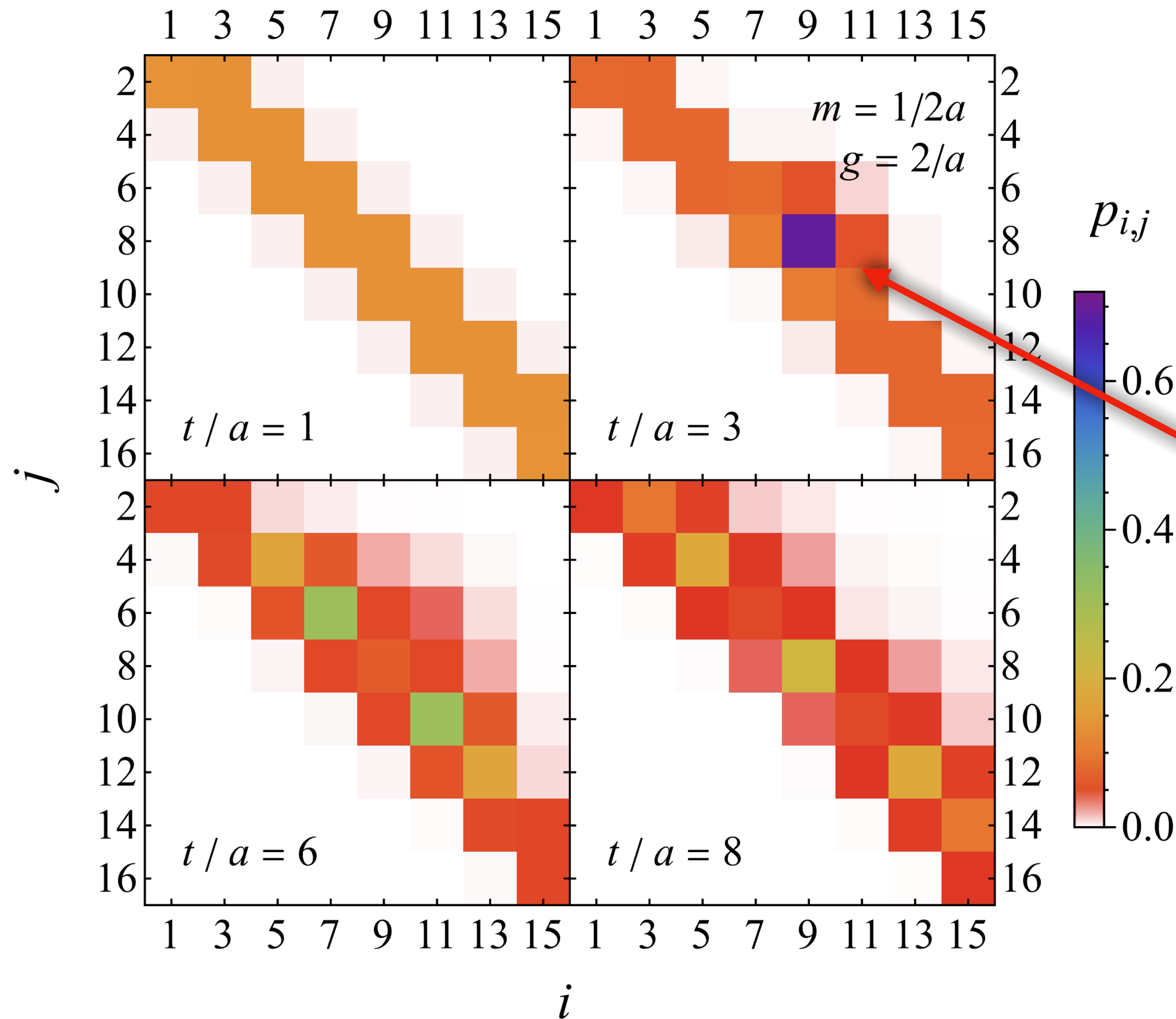
*probability of exciting the
 i^{th} antiquark and j^{th} quark*



$$p_{i,j} = |\langle \Psi_t | \chi_i \chi_j^\dagger | \text{empty} \rangle|^2$$

probability of exciting the
 i^{th} antiquark and j^{th} quark

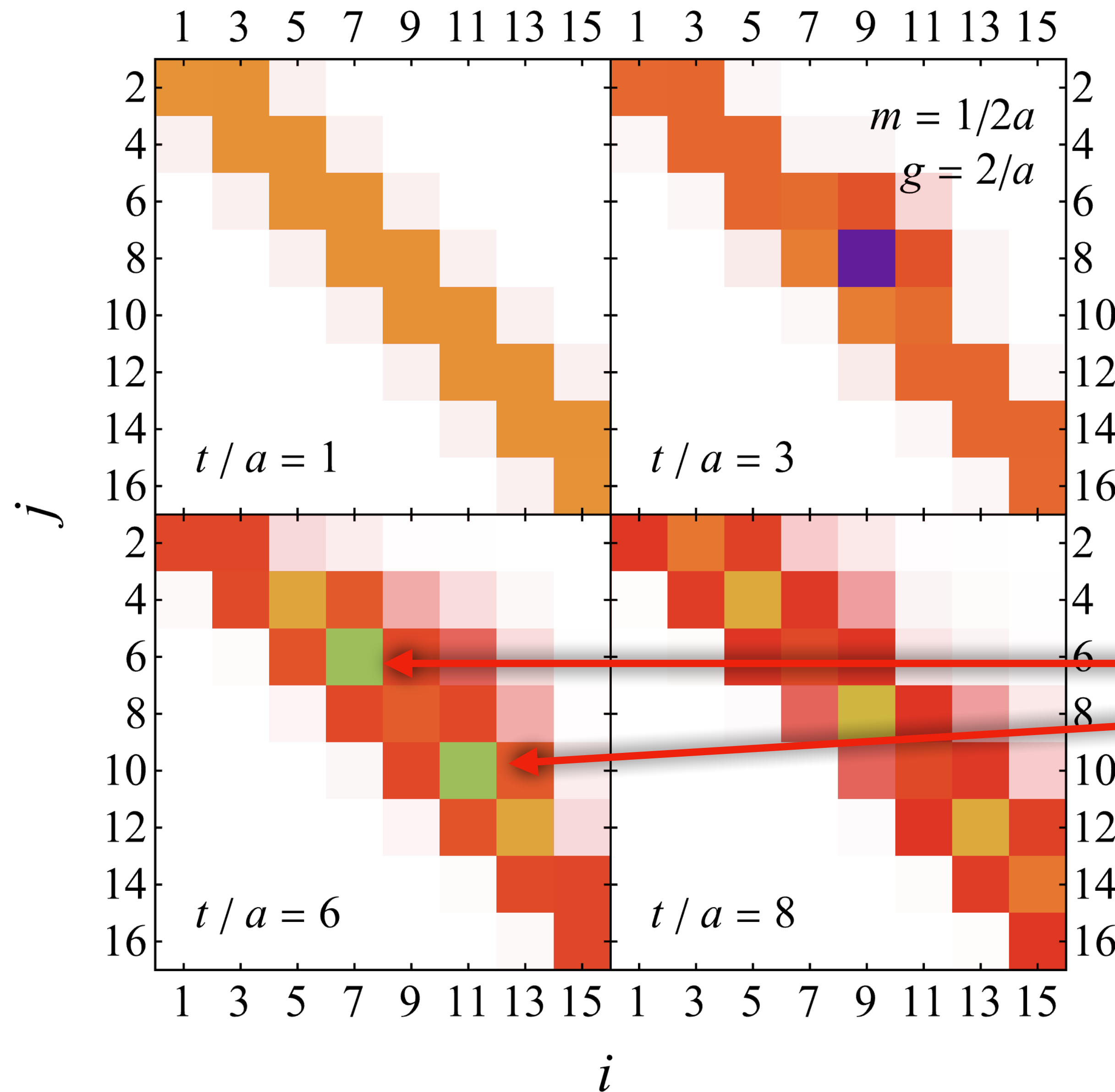
homogeneous vacuum structure



$$p_{i,j} = |\langle \Psi_t | \chi_i \chi_j^\dagger | \text{empty} \rangle|^2$$

probability of exciting the
 i^{th} antiquark and j^{th} quark

creation of meson



$$p_{i,j} = |\langle \Psi_t | \chi_i \chi_j^\dagger | \text{empty} \rangle|^2$$

probability of exciting the
 i^{th} antiquark and j^{th} quark

two mesons

