

# Jet Physics in Heavy-Ion Collisions





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## Probing quark-gluon plasma (QGP) using jets



- Quantify QGP effects on jets by comparing jet spectra between pp and AA collisions
- Nuclear modification factor:

$$R_{\rm AA}(p_{\rm T}) = \frac{dN^{\rm AA}/dp_{\rm T}}{\langle N_{\rm coll}^{\rm AA} \rangle \times dN^{pp}/dp_{\rm T}}$$







## What is a jet depends on how one defines a jet

### R<sub>AA</sub> of single particles





Mueller *et al.*, Ann. Rev. Nucl. Part. Sci. 62, 361 (2012)

### R<sub>AA</sub> of full jets





ATLAS, Phys. Lett. B 790 (2019) 108

## What is a jet depends on how one defines a jet

### R<sub>AA</sub> of single particles





Mueller *et al.*, Ann. Rev. Nucl. Part. Sci. 62, 361 (2012)

### Different clustering algorithms developed for different purposes



*k*<sub>T</sub> algorithm for estimating soft background



### R<sub>AA</sub> of full jets





ATLAS, Phys. Lett. B 790 (2019) 108

## anti-k<sub>T</sub> algorithm

for identifying energetic jets



C/A algorithm for identifying splitting angles



## Jets tagged with heavy quarks





- Produced from initial hard scatterings
- Serve as an ideal probe of the QGP properties
- Provide a unique opportunity for studying the flavor dependence of parton splitting (dead cone effect)

## Searches for the flavor dependence of parton splitting

### Distribution of splitting angles in pp



Clear suppression of splitting at small  $\theta$  in *D*-jets *vs*. inclusive jets

### Hadron R<sub>AA</sub> (parton energy loss)



Phys. Lett. B 782 (2018) 474-496

No clear separation between charged hadrons, *D*, and *B*, except at very low  $p_T$ 

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**Goals:** 

### Hadron *R*<sub>AA</sub> (parton energy loss)



Phys. Lett. B 782 (2018) 474-496

No clear separation between charged hadrons, D, and B, except at very low  $p_{T}$ 

 Consistently understand hadron and full jet observables Consistently understand light and heavy flavor jets Use jets to probe properties of the QGP

## **Theoretical framework of jet quenching**



$$d\sigma_{h} = \sum_{abjd} f_{a/p} \otimes f_{b/p} \otimes d\sigma_{ab \to jd} \otimes D_{h/j} (\text{or } J_{j})$$

- $f_{a/p}, f_{b/p} \rightarrow f_{a/A}, f_{b/B}$ : cold nuclear matter (initial state) effect, e.g., shadowing, Cronin, ..., measured in pA collisions



$$d\tilde{\sigma}_{h} = \sum_{abjd} f_{a|A} \otimes f_{b|B} \otimes d\sigma_{ab \to jd} \otimes \tilde{D}_{h|j} \text{ (or } \tilde{J}_{abjd}$$

•  $D_{h/i} \rightarrow D_{h/i}$ : medium modified fragmentation function, hot nuclear matter (final state) effect • Factorization assumption:  $\tilde{D}_{h/j} = \sum_{i'} P_{j \to j'} \otimes D_{h/j'}$ , nuclear modification of parton j



### Parton transport inside the QGP

### Linear Boltzmann Transport (LBT)

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

### Parton transport inside the QGP

### Linear Boltzmann Transport (LBT)

**Elastic scattering (** $ab \rightarrow cd$ **)** 

$$\mathscr{C}_a^{\text{el}} = \sum_{b,c,d} \int \prod_{i=b,c,d} \frac{d[p_i]}{2E_a} (\gamma_d f_c f_d - \gamma_b)$$

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

 $\int_{a} f_a f_b \cdot (2\pi)^4 \delta^4 (p_a + p_b - p_c - p_d) \left| \mathcal{M}_{ab \to cd} \right|^2$ 

 $2 \rightarrow 2$  scattering matrices

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loss term: scattering rate (for Monte-Carlo simulation)

$$\Gamma_{a}^{\text{el}}(\mathbf{p}_{a}, T) = \sum_{b,c,d} \frac{\gamma_{b}}{2E_{a}} \int \prod_{i=b,c,d} d[p_{i}]f_{b}$$

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

 $\int f_a f_b \cdot (2\pi)^4 \delta^4 (p_a + p_b - p_c - p_d) \left[ \mathcal{M}_{ab \to cd} \right]^2$ 

 $2 \rightarrow 2$  scattering matrices

 $f_b \cdot (2\pi)^4 \delta^{(4)}(p_a + p_b - p_c - p_d) \left| \mathcal{M}_{ab \to cd} \right|^2$ 

### **Inelastic scattering**





• Medium information absorbed in  $\hat{q} \equiv d \langle p_{\perp}^2 \rangle / dt$ 

Majumder PRD 85 (2012); Zhang, Wang and Wang, PRL 93 (2004)

• Higher-twist formalism: collinear expansion (  $\langle k_{\perp}^2 \rangle \ll l_{\perp}^2 \ll Q^2$  )

$$\frac{1}{4}\sin^2\left(\frac{t-t_i}{2\tau_f}\right)$$

## Flavor hierarchy in hadron suppression

Perturbative calculation simultaneously describes the  $R_{AA}$  of light and heavy hadrons at high  $p_T$ [Xing, SC, Qin and Xing, Phys. Lett. B 805 (2020) 135424 ]

pp baseline within the NLO production + fragmentation framework • Crucial contributions from the gluon fragmentation (or  $g \rightarrow Q\bar{Q}$ ) process



**Gluon fragmentation** 

- dominates  $h^{\pm}$  production up to 50 GeV
- contributes to over 40% D up to 100 GeV

## Flavor hierarchy in hadron suppression

### NLO initial production and fragmentation + Boltzmann transport + hydrodynamic medium for QGP

### 1.2 CMS 0-10% ALICE 0-5% 0.8 R 44 0.6 0.4 100 10 p<sub>T</sub> (GeV)

### charged hadron

- g-initiated h & D  $R_{AA} < q$ -initiated h & D  $R_{AA}$  [ $\Delta E_q > \Delta E_{q/c}$ ]



D meson

p<sub>T</sub> (GeV)

•  $R_{AA}(c->D) > R_{AA}(q->h) [\Delta E_q > \Delta E_c], R_{AA}(q->D) < R_{AA}(q->h) [different FFs] => R_{AA}(h) \approx R_{AA}(D)$ • Signature of flavor hierarchy of parton  $\Delta E$  offset by gluon production/fragmentation in hadron  $R_{AA}$ 



## Flavor hierarchy in hadron suppression



- starting from  $p_T \sim 8 \text{ GeV}$
- confirmation from future precision measurement

• A simultaneous description of charged hadron, D meson, B meson, B-decay D meson R<sub>AA</sub>'s

Predict R<sub>AA</sub> separation between B and h / D below 40 GeV, but similar values above – wait for



## From hadrons to full jets



- Jet-medium interactions: medium modification of jets + medium response

 Energy deposition + depletion → jet-induced medium excitation (medium response) Primordial jet partons and medium response cannot be cleanly separated for jet observables

### Jet R<sub>AA</sub> and v<sub>2</sub>

### RAA



- Including medium response reduces jet energy loss and thus increases the jet  $R_{AA}$



• With  $R_{AA}$  fixed, including medium response (coupled to medium flow) increases the jet  $v_2$ 

### Jet substructure



### Transverse (*r*) distribution: jet shape



Tachibana, Chang, Qin, Phys. Rev. C 95 (2017) 044909

### Longitudinal (z) distribution: jet fragmentation function



Chen, SC, Luo, Pang, Wang, Phys. Lett. B 777 (2018) 86-90



## Search for unique signatures of medium response

### Hadron suppression in diffusion wake



- Confirmed by recent CMS data [CMS-PAS-HIN-23-006]

• Hadron suppression predicted in the backward direction of jets at  $1 < p_T^h < 2 \text{ GeV}$ [Chen, SC, Luo, Pang, Wang, PLB 777 (2018) 86, Yang, Luo, Chen, Pang, Wang, PRL 130 (2023) 052301]

## New developments on searching for wake

• Extension from 2D ( $\phi$ -plane) to 3D ( $\eta$ - $\phi$ ) structure of wake [Yang, Luo, Chen, Pang, Wang, Phys. Rev. Lett. 130 (2023) 5, 052301; Yang, Wang, arXiv: 2501.03419]

Example: interference between wave front and energy depletion in di-jet events



Rapidity asymmetry of hadron distribution w.r.t. the leading jet

4 -4

 $\Delta \eta = \eta_h - \eta_{iet}$ 

 $\Delta \eta = \eta_h - \eta_{jet}$ 

## New developments on searching for wake

- From jet-transport-hydrodynamics simulation to  $\phi^4$  field simulation



- Jet energy deposition gradually approaches thermalization •
- Mach cone structure for jet energy deposition
- trajectory)  $\rightarrow$  possible rich pattern for jet-induced QGP polarization





Cartoon from Serenone et. al., PLB 820 (2021) 136500

• Vortex ring structure:  $\vec{\omega}$  parallel to  $\vec{p}_{jet} \times \vec{v}_{med}$  in front, antiparallel in back (close to jet

## Hadron chemistry as a new signature of medium response

### **Baryon enhancement**



Luo, Mao, Qin, Wang, Zhang, PLB 837 (2023) 137638

Luo, SC, Qin, arXiv:2412.19283

- Larger q and s densities in QGP than in vacuum jets
- Baryon and strangeness enhancement around jets in AA
- Stronger enhancement at larger distance from jet axis



## A novel observable: energy-energy correlator (EEC)

![](_page_22_Figure_1.jpeg)

[Komiske et. al., PRL 130 (2023) 051901]

- **EEC**:  $\frac{d\Sigma}{dR_L} = \int d\vec{n}_1 d\vec{n}_2 \frac{\langle \mathscr{E}(\vec{n}_1)\mathscr{E}(\vec{n}_2) \rangle}{Q^2} \delta(\Delta R_{12} - R_L)$
- Proposed by conformal theory,  $EEC(R_{L})$  is scale independent at high energy limit
- Implement a first realistic calculation on light and heavy flavor jet EEC in AA collisions [Xing, SC, Qin, Wang, Phys. Rev. Lett. 134 (2025) 052301]

![](_page_22_Figure_7.jpeg)

[Craft et. al., arXiv:2210.09311]

 $\mathscr{E}$ : energy flow in a given direction,  $\Delta R_{12} = \sqrt{\Delta \phi_{12}^2 + \Delta \eta_{12}^2}$ : relative angle, Q: hard scale

Variation of EEC(R<sub>L</sub>) indicates emergence of new scales: hadronization, QGP, quark mass ...

## Light vs. heavy flavor jet EEC in pp collisions

![](_page_23_Figure_1.jpeg)

• Jet in pp: Pythia 8 simulation

• EEC analysis (*i*, *j* denote jet constituents)

$$\frac{d\Sigma(\theta)}{d\theta} = \frac{1}{\Delta\theta} \sum_{\substack{|\theta_{ij} - \theta| < \Delta\theta/2}} \frac{p_{\mathrm{T},i}(\vec{n}_i) p_{\mathrm{T},j}(\vec{n}_j)}{p_{\mathrm{T},j\text{et}}^2}$$

• Flavor (mass) dependence:

- Overall magnitude: charged jet > D-jet > B-jet
- Typical (peak) angle: charged jet < D-jet < B-jet

Suppression of splitting within  $\theta_0 \sim m_Q/E_Q$  in vacuum

## Light vs. heavy flavor jet EEC in pp collisions

![](_page_24_Figure_1.jpeg)

• Jet in pp: Pythia 8 simulation

• EEC analysis (*i*, *j* denote jet constituents)

$$\frac{d\Sigma(\theta)}{d\theta} = \frac{1}{\Delta\theta} \sum_{\substack{|\theta_{ij} - \theta| < \Delta\theta/2}} \frac{p_{\mathrm{T},i}(\vec{n}_i) p_{\mathrm{T},j}(\vec{n}_j)}{p_{\mathrm{T},j\text{et}}^2}$$

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Jet energy dependence

• Higher  $p_T \rightarrow \Sigma$  peaks at smaller  $\theta$ 

 $p_{\rm T}\theta_{\rm peak}$  ~ transition scale between pert. and non-pert.

## Light vs. heavy flavor jet EEC in central PbPb collisions

![](_page_25_Figure_1.jpeg)

Full simulation — Pythia + LBT in hydro

- Flavor hierarchy of EEC preserves, though being modified, in AA collisions
- enhancement at small  $\theta$  (except for *B*-jet) and large  $\theta$

![](_page_25_Figure_5.jpeg)

• General features of nuclear modification (AA - pp): suppression at intermediate  $\theta$ ,

## **Different contributions to medium modification on EEC**

![](_page_26_Figure_1.jpeg)

- Medium response enhances EEC at large  $\theta$

S: shower partons inherited from Pythia S+R: add medium-induced gluons S+R+M: further add medium response

• Jet energy loss causes suppression over the entire  $\theta$  region - Medium-induced gluon emission enhances EEC at small  $\theta$ 

## **Further studies on EEC**

# Discriminate quark and gluon jet quenching in AA collisions

![](_page_27_Figure_2.jpeg)

[Chen, Shen, Xue, Dai, Zhang, Wang, arXiv:2409.13996]

Can be tested by comparing between inclusive jets and γ-jets

# Measurement of $J/\psi$ energy correlator in *pp* collisions

![](_page_27_Figure_6.jpeg)

[From STAR poster at Quark Matter 2025]

May be used to probe the hadronization process of J/ψ, e.g., c + c̄ → J/ψ + g
[Chen, Liu, Ma, Phy. Rev. Lett. 133 (2024) 191901]

## Constraints on jet transport coefficient inside the QGP

 $\hat{\boldsymbol{q}} \equiv d\langle k_{\perp}^2 \rangle / dt \sim \langle F^{ai+}(0)F_i^{a+}(y^-) \rangle$ 

![](_page_28_Figure_2.jpeg)

 QGP is much more opaque than cold nuclear matter to jet propagation [Xie et al., Phys. Rev. C 108 (2023) 1, L011901]

![](_page_28_Figure_5.jpeg)

• Jet model improvement + Bayesian statistical analysis  $\rightarrow$  temperature and jet energy dependences of  $\hat{q}$ 

## Probing the equation of state of the QGP

[F.-L. Liu, X.-Y. Wu, SC, G-Y. Qin, X.-N. Wang, Phys. Lett. B 848 (2024) 138355]

![](_page_29_Figure_2.jpeg)

![](_page_29_Figure_3.jpeg)

**Strategy:** Fit g from comparing transport model to data Calculate EoS from g

## **EoS of QGP and diffusion coefficient of heavy quarks**

### **Equation of state**

![](_page_30_Figure_2.jpeg)

- Agreement with the lattice data

### **Diffusion coefficient**

![](_page_30_Figure_6.jpeg)

Simultaneous constraint on QGP properties and transport properties of hard probes

## **Probing the specific viscosity of the QGP**

### • Indirect constraint: convert $\hat{q}$ into $\eta/s$

![](_page_31_Figure_2.jpeg)

 $\eta/s \approx 1.25 T^3/\hat{q}$ 

weak coupling calculation

• **Outlook**: constrain specific viscosity and speed of sound using medium response

![](_page_31_Figure_6.jpeg)

![](_page_31_Figure_7.jpeg)

Patterns of jet-induced Mach cone depends on  $\eta/s$  and  $c_s$  (EoS)

Neufeld, Phys. Rev. C 79 (2009) 054909

## Summary

### Probing strongly interacting matter using energetic hadrons and jets

- Perturbative calculation successfully describes hadron suppression at high  $p_{T}$
- Jet-induced medium excitation is crucial for understanding full jet observables
- EEC is an excellent observable for studying flavor dependence of jet quenching
- Jet observables are used to constrain various QGP properties:  $\hat{q}$ , EoS,  $\eta/s$  ...

![](_page_32_Picture_6.jpeg)

Thank. you!

## EEC of partons developed from a single quark

![](_page_34_Figure_1.jpeg)

- Single quark  $\rightarrow$  LBT + static medium  $\rightarrow$  EEC of daughter partons
- Flavor (mass) hierarchy of EEC:

  - Clear strong suppression of  $\Sigma$  below  $\theta_0 \sim m_Q/E_{\rm initial}$
- Contributions form medium response and gluon emission show similar hierarchies

• Magnitude: charged > D > B-jet; peak position: charged < D < B-jet (similar to vacuum jets)

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## Effects of trigger bias on the jet EEC

### $p_{T}$ trigger in both pp and AA

![](_page_35_Figure_2.jpeg)

- AA jets with trigger bias originate from pp jets with higher  $p_T$  and initial virtuality scale → Stronger but narrower vacuum splittings
- Can be tested using  $\gamma$ -jets

 $p_{T}$  trigger only in pp (no trigger bias in AA)

![](_page_35_Figure_6.jpeg)

 $\rightarrow$  Enhances EEC at small  $\theta$ , reduces the suppression/enhancement at intermediate/large  $\theta$