

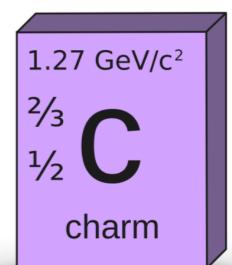


School of Nuclear Science and Technology

Kang-kang Shao

Collaborators: Chun Huang (WashU) , Dong-hao Li (LZU)

Supervisor: Fu-sheng Yu, Qin Qin



第二届武汉高能物理青年论坛
Wu han, China

Content

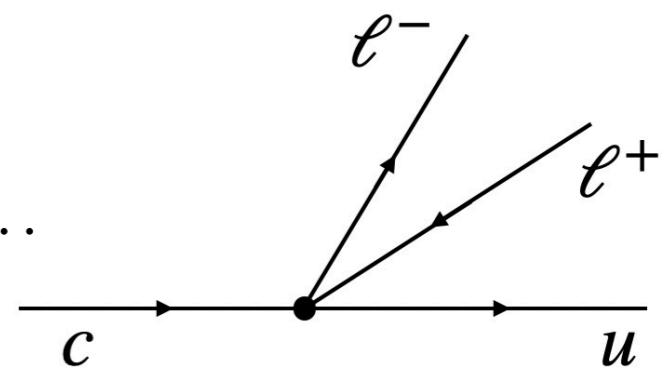
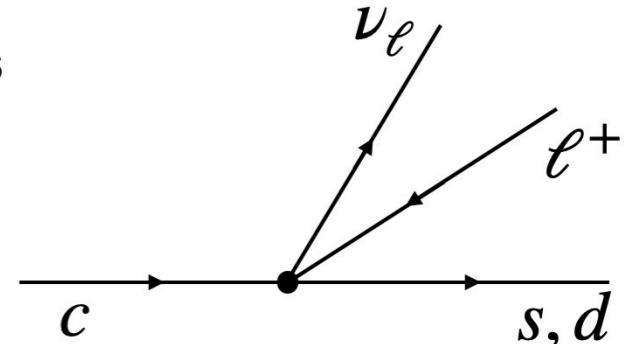
- 1. What and Why Electronic Semi-inclusive Charm decay?**
- 2. Theoretical Framework**
- 3. Experimental Status**
- 4. Phenomenological Analysis (preliminary)**
- 5. Wishlist**

1. What and Why Electronic Semi-inclusive Charm decay?

What and Why Electronic Semi-inclusive Charm decay



- **Experimental detection of partial final state particles**
 - ➔ $D \rightarrow e^+ X$ ($D \rightarrow e^+ \nu_e X$, only e^+ is detected)
- **Sum of a group of exclusive channels**
 - ➔ $D^0 \rightarrow e^+ X_s = D \rightarrow e^+ \nu_e K^-$, $e^+ \nu_e K^- \pi^0$, $e^+ \nu_e \bar{K}^0 \pi^-$, ...
 - ➔ $D^0 \rightarrow e^+ X_d = D \rightarrow e^+ \nu_e \pi^-$, $e^+ \nu_e \pi^- \pi^0$, $e^+ \nu_e \pi^- \pi^+ \pi^-$, ...



- Compared to exclusive decays: **Better** theoretical control
- Compared to beauty decays: **More sensitive to power corrections**



What and Why Electronic Semi-inclusive Charm decay

Charmed hadron lifetimes: theory vs experiment

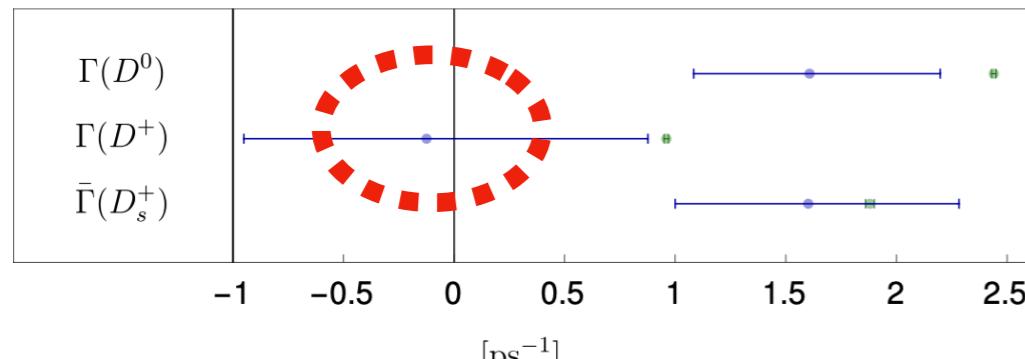


Fig 1

[Lenz et al, '22]

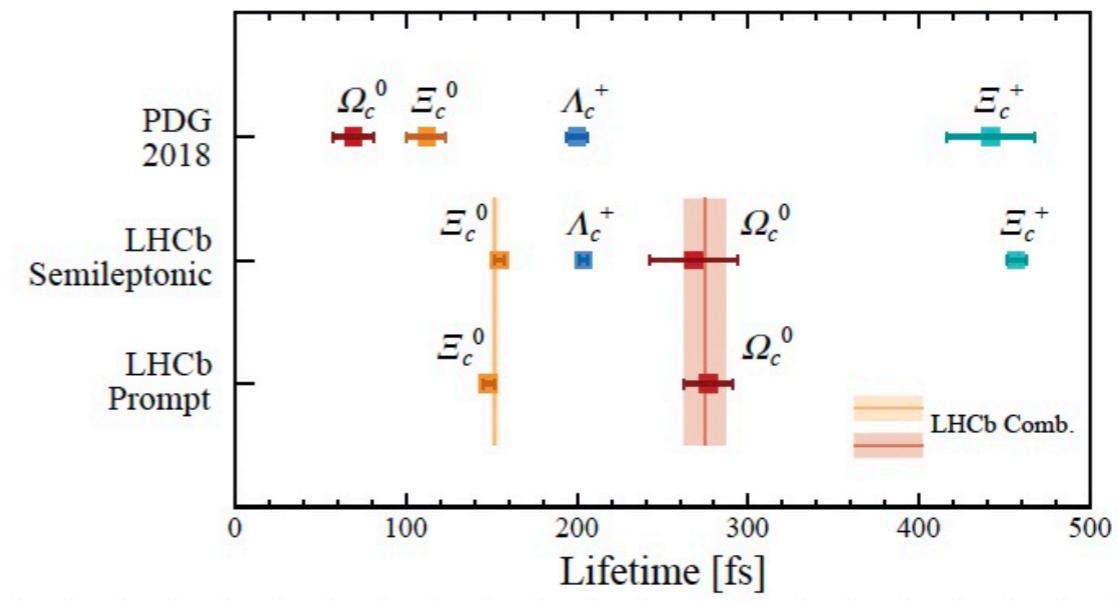


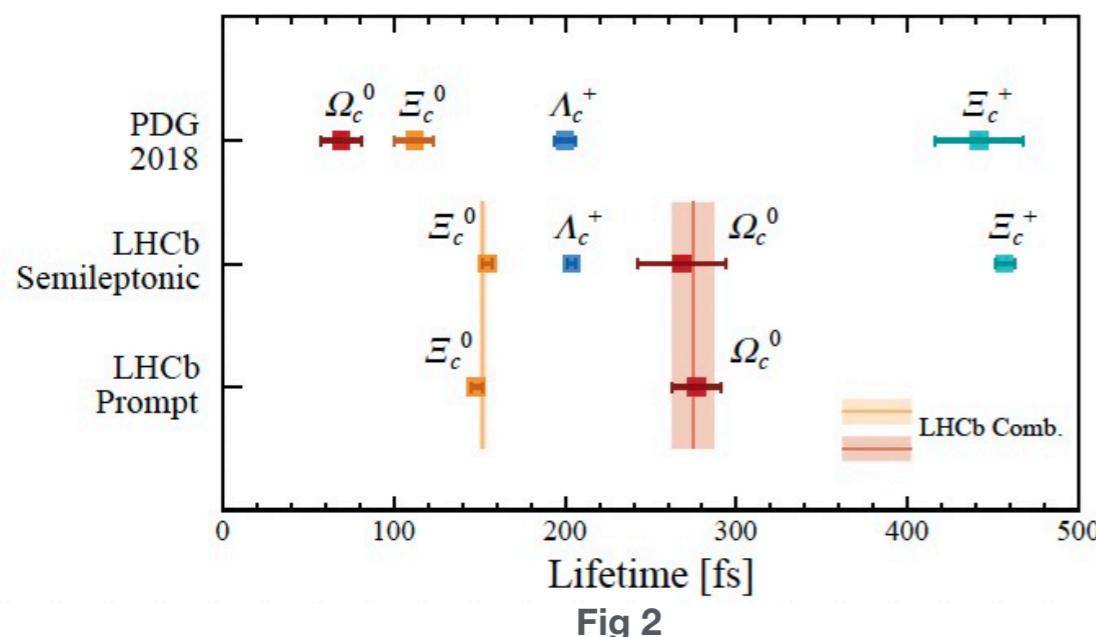
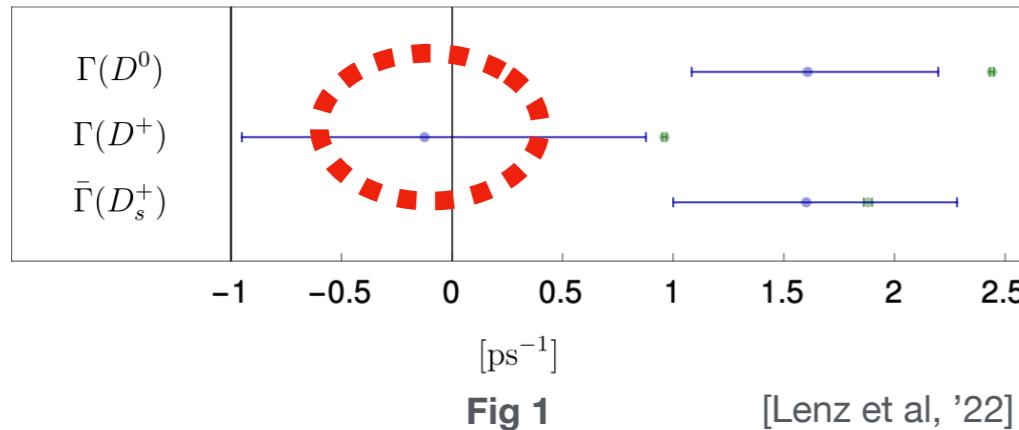
Fig 2

$$\begin{aligned}\mathcal{O}(1/m_c^3) &\Rightarrow \tau(\Xi_c^+) > \tau(\Lambda_c^+) > \tau(\Xi_c^0) > \tau(\Omega_c^0), \\ \mathcal{O}(1/m_c^4) &\Rightarrow \tau(\Omega_c^0) > \tau(\Xi_c^+) > \tau(\Lambda_c^+) > \tau(\Xi_c^0), \\ \mathcal{O}(1/m_c^4) \text{ with } \alpha &\Rightarrow \tau(\Xi_c^+) > \tau(\Omega_c^0) > \tau(\Lambda_c^+) > \tau(\Xi_c^0).\end{aligned}$$

[Cheng, '21]

What and Why Electronic Semi-inclusive Charm decay

Charmed hadron lifetimes: theory vs experiment



$$\mathcal{O}(1/m_c^3) \Rightarrow \tau(\Xi_c^+) > \tau(\Lambda_c^+) > \tau(\Xi_c^0) > \tau(\Omega_c^0),$$

$$\mathcal{O}(1/m_c^4) \Rightarrow \tau(\Omega_c^0) > \tau(\Xi_c^+) > \tau(\Lambda_c^+) > \tau(\Xi_c^0),$$

$$\mathcal{O}(1/m_c^4) \text{ with } \alpha \Rightarrow \tau(\Xi_c^+) > \tau(\Omega_c^0) > \tau(\Lambda_c^+) > \tau(\Xi_c^0).$$

Solutions/hints

- ▶ Dependence on identical hadronic parameters in HQET, $\langle H_c | O_i | H_c \rangle$
- ▶ Extraction in the inclusive decay spectrum and application to lifetime

“ Again a more precise experimental determination of μ_π^2 from fits to semi-leptonic D^+ , D^0 and D_s^+ meson decays — as it was done for the B^+ and B^0 decays — would be very desirable. ”

[Lenz et al, '22]

2. Theoretical Framework

Theoretical Framework

- **Optical theorem**

$$\sum \langle D | H | X \rangle \langle X | H | D \rangle \propto \text{Im} \int d^4x \langle D | T\{H(x)H(0)\} | D \rangle$$

- **Operator product expansion (OPE)**

★ Short distance : $x \sim 1/m_c$

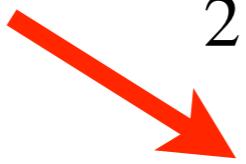
★ Fluctuation in D meson $\sim \Lambda_{\text{QCD}}$

$$T\{H(x)H(0)\} = \sum_n C_n(x) O_n(0) \rightarrow 1 + \frac{\Lambda_{\text{QCD}}}{m_c} + \frac{\Lambda_{\text{QCD}}^2}{m_c^2} + \dots$$

Systematic OPE in HQET.

- **Heavy Quark Effective Theory**

$$h_\nu(x) \equiv e^{-im_c v \cdot x} \frac{1 + \gamma \cdot v}{2} c(x) \quad v = (1, 0, 0, 0) \quad L \ni \bar{h}_\nu i v \cdot D h_\nu$$


Subtract the big intrinsic momentum,
Leave only $\sim \Lambda_{\text{QCD}}$ degrees of freedom.

$$-\bar{h}_\nu \frac{D_\perp^2}{2m_c} h_\nu - a(\mu) g \bar{h}_\nu \frac{\sigma \cdot G}{4m_c} h_\nu + \dots$$

Theoretical Framework

- **OPE**

$$T\{H(x)H(0)\} = \sum_n C_n(x) O_n(0)$$

$$C_n(x) \qquad \qquad \qquad O_n(0)$$

★ LO: $\alpha_s^0(m_c)$

★ Dim-3: $\bar{h}_v h_v$ ($\bar{c}\gamma^\mu c$) → **partonic decay rate.**

★ NLO: $\alpha_s(m_c)$

★ Dim-5: $\bar{h}_v D_\perp^2 h_v$, $g\bar{h}_v \sigma \cdot G h_v$.

★ ...

★ Dim-6: $\bar{h}_v D_\mu(v \cdot D) D^\mu h_v$, $(\bar{h}_v \Gamma_1 q)(\bar{q} \Gamma_2 h_v), \dots$

★ ...

- **Contribute to inclusive decay rate and lifetime**

1. Matrix elements of the **same operators** (SL& NL)

2. Only different short-distance coefficients

$$\lambda_1 \equiv \frac{1}{4m_D} \langle D | \bar{h}_v (iD)^2 h_v | D \rangle = -\mu_\pi^2$$
$$\lambda_2 \equiv \frac{1}{16(s_c \cdot s_q)} \frac{1}{2m_D} \langle D | \bar{h}_v g \sigma \cdot G h_v | D \rangle = \frac{\mu_G^2}{3}$$

Theoretical Framework

- **Structure functions**

$$\frac{d\Gamma}{d\hat{E}_\ell d\hat{q}^2 d\hat{u}} = \frac{G_F^2 m_b^5 |V_{cb}|^2}{16\pi^3} \theta(\hat{u}_+ - \hat{u}) \theta(\hat{E}_\ell) \theta(\hat{q}^2) \times \\ \times \left\{ \hat{q}^2 W_1 - \left[2\hat{E}_\ell^2 - 2\hat{E}_\ell \hat{q}_0 + \frac{\hat{q}^2}{2} \right] W_2 + \hat{q}^2 (2\hat{E}_\ell - \hat{q}_0) W_3 \right\},$$

$$W_i = W_i^{(0)} + \frac{\mu_\pi^2}{2m_c^2} W_i^{(\pi,0)} + \frac{\mu_G^2}{2m_c^2} W_i^{(G,0)} + \frac{\alpha_s}{\pi} \left[C_F W_i^{(1)} + C_F \frac{\mu_\pi^2}{2m_c^2} W_i^{(\pi,1)} + \frac{\mu_G^2}{2m_c^2} W_i^{(G,1)} \right] + \dots$$

$$W_i^{(1)} = w_i^{(0)} \left\{ \mathcal{S}_i \delta(\hat{u}) - \left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ - \left(\frac{7}{4} - 2 \ln w \right) \left[\frac{1}{\hat{u}} \right]_+ + w B(\hat{q}^2, \hat{u}) \theta(\hat{u}) \right\} + \mathcal{R}_i^{(1)} \theta(\hat{u}),$$

$$\mathcal{S}_i = -\frac{5}{4} - \frac{\pi^2}{3} - \text{Li}_2(1-w) - 2 \ln^2 w - \frac{5w-4}{2(1-w)} \ln w + \frac{\ln w}{2(1-w)} \delta_{i2}$$

$$\mathcal{R}_1^{(1)} = \frac{3}{4} + \frac{\hat{u}(12-w-\hat{u})}{2\tilde{\lambda}} + \left(w + \frac{\hat{u}}{2} - \frac{\hat{u}(2\hat{u}+3w)}{\tilde{\lambda}} \right) \mathcal{I}_1$$

$$\mathcal{R}_2^{(1)} = \frac{6\hat{u}(\hat{u}^2 - (3-w)\hat{u} - 12 + 13w)}{\tilde{\lambda}^2} + \frac{\hat{u} - 38 + 21w}{\tilde{\lambda}} \\ - 4 \frac{\frac{w}{2}\hat{u}^3 + (2w^2 - 6)\hat{u}^2 + (7 - 3w + \frac{5}{2}w^2)w\hat{u} + w^3(w-4)}{\tilde{\lambda}^2} \mathcal{I}_1$$

$$\mathcal{R}_3^{(1)} = \frac{3\hat{u} - 8 + 5w}{\tilde{\lambda}} + \frac{\hat{u}^2 - (6-w)\hat{u} + 4w}{\tilde{\lambda}} \mathcal{I}_1$$

$$\mathcal{I}_1 = \frac{1}{\sqrt{\tilde{\lambda}}} \ln \frac{\hat{u} + w + \sqrt{\tilde{\lambda}}}{\hat{u} + w - \sqrt{\tilde{\lambda}}}$$

NLO dim-3
10

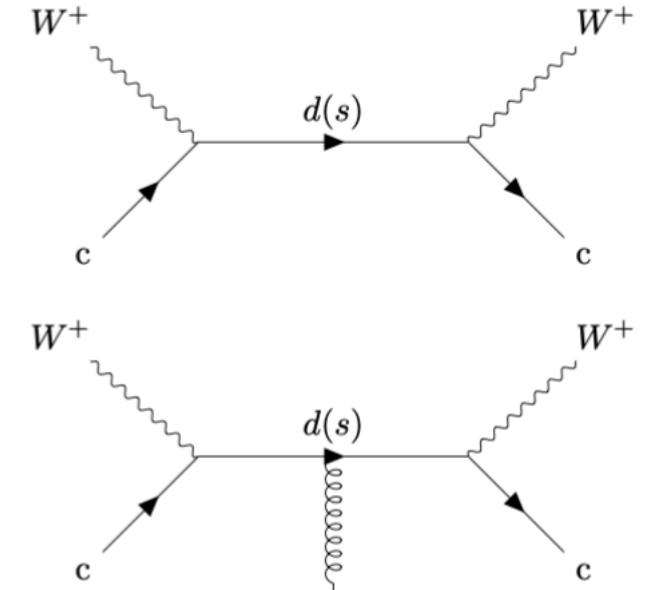


Fig 5. Leading Order

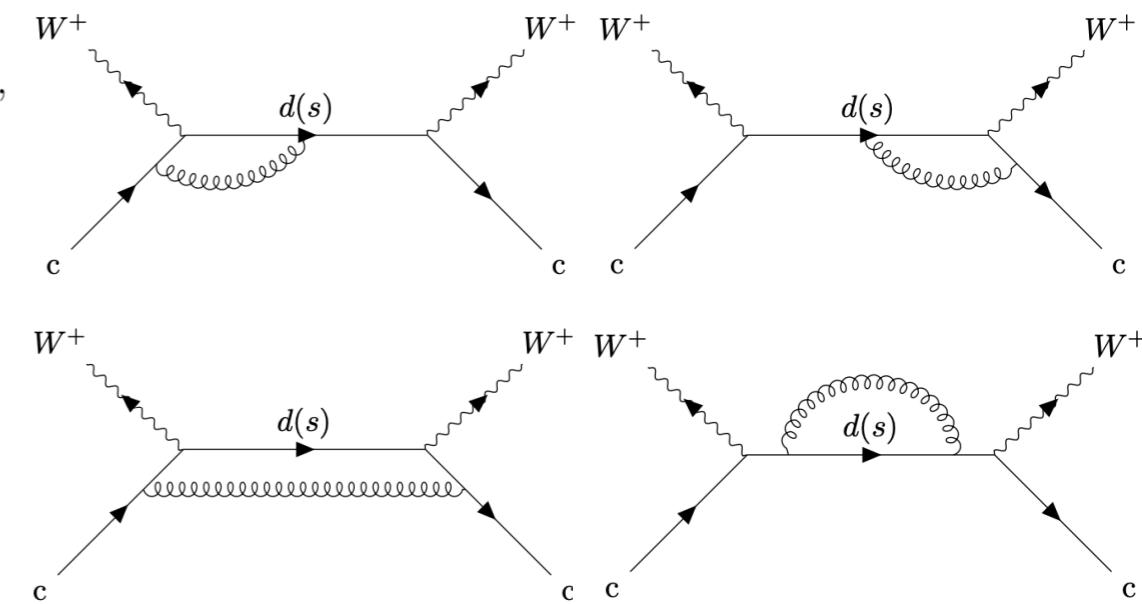


Fig 6. Next to Leading Order

Theoretical Framework

- **Structure functions**

$$W_1^{(\pi,1)} = w \left[B_1 - \frac{C}{2} + \frac{5w-2}{12} \left[\frac{1}{\hat{u}^2} \right]_+ + \left(\frac{16+3w-10w^2}{12} - \frac{8w^3-w^2-14w+8}{6(1-w)} \ln w \right) \delta'(\hat{u}) \right] \\ - \frac{4}{3}(2-w) \left[\left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ + L_w \delta(\hat{u}) \right] + \left(\frac{8}{3}(2-w) \ln w - \frac{4+18w-13w^2}{6w} \right) \left[\frac{1}{\hat{u}} \right]_+ + \mathcal{R}_1^{(\pi)} \theta(\hat{u}) \\ + \left(\frac{13w}{12} - \frac{1}{6} - \frac{1}{3w} - \frac{w^2}{12} + \frac{w^3}{4} + \frac{4+6w-13w^2+3w^3+2w^5}{3w(1-w)} \ln w \right) \delta(\hat{u}) \quad (\text{A.1})$$

$$W_2^{(\pi,1)} = 4B_2 + 6C + \frac{9w-10}{3} \left[\frac{1}{\hat{u}^2} \right]_+ + \left(\frac{4+6w+16w^2}{3} \ln w - \frac{22-21w+10w^2}{3} \right) \delta'(\hat{u}) \\ + \left(w^2 + \frac{116}{3w^2} - 7w - \frac{50}{w} + \frac{88}{3} - 4 \frac{42-34w+17w^2-6w^3+2w^4}{3w^2} \ln w \right) \delta(\hat{u}) \\ + \left(\frac{10}{3} - \frac{68}{3w} + \frac{28}{w^2} \right) \left[\frac{1}{\hat{u}} \right]_+ + \mathcal{R}_2^{(\pi)} \theta(\hat{u}) \quad (\text{A.2})$$

$$W_3^{(\pi,1)} = 2B_3 + C + \left(\frac{7w}{6} - 1 \right) \left[\frac{1}{\hat{u}^2} \right]_+ + \left(\frac{5}{3}(1-w)w + \frac{w(6+3w-8w^2)}{3(1-w)} \ln w \right) \delta'(\hat{u}) \\ + 2 \left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ + \left(\frac{19}{6} - \frac{2}{w} + \frac{4}{w^2} - 4 \ln w \right) \left[\frac{1}{\hat{u}} \right]_+ + \left[2L_w + \frac{w^2}{2} + \frac{14}{3w^2} - \frac{11w}{6} - \frac{20}{3w} \right. \\ \left. + \frac{41}{6} + \left(\frac{7w-6}{1-w} + \frac{4}{3}w - \frac{4}{3}w^2 - \frac{8}{w^2} + \frac{4}{w} \right) \ln w \right] \delta(\hat{u}) + \mathcal{R}_3^{(\pi)} \theta(\hat{u}) \quad (\text{A.3})$$

$$B_i = \frac{w^2}{6} \left(\left[\frac{7}{4} - 2L_w - \frac{2-w}{1-w} \ln w + \delta_{i2} \frac{\ln w}{1-w} \right] \delta''(\hat{u}) - 4 \left[\frac{\ln \hat{u}}{\hat{u}^3} \right]_+ + (8 \ln w - 1) \left[\frac{1}{\hat{u}^3} \right]_+ \right), \\ C = \frac{2(2-w)}{3} \left(- \left[\frac{\ln \hat{u}}{\hat{u}^2} \right]_+ + 2 \ln w \left[\frac{1}{\hat{u}^2} \right]_+ + L_w \delta'(\hat{u}) \right) \\ L_w = \text{Li}_2(1-w) + 2 \ln^2 w + \frac{\pi^2}{3} \quad (\text{A.4})$$

$$\mathcal{R}_1^{(\pi)} = \frac{(4\hat{u}-w)(2-w)\hat{u}+2w^3}{3\hat{u}^3} \ln \frac{\hat{u}}{w^2} + \left[\frac{2w^6}{3\hat{u}^3} + \frac{7w^5}{3\hat{u}^2} - \frac{14-5\hat{u}}{3\hat{u}^2} w^4 - \frac{13\hat{u}+32}{6\hat{u}} w^3 \right. \\ \left. - \frac{23\hat{u}^2-36\hat{u}-48}{6\hat{u}} w^2 - (13\hat{u}^2-58\hat{u}+36) \frac{w}{6} - \frac{\hat{u}}{6} (3\hat{u}^2-26\hat{u}+8) \right] \frac{\mathcal{I}_1}{\lambda} \\ - \frac{4w^2}{3\hat{u}^2} + \frac{2\hat{u}^2+2\hat{u}w-13\hat{u}+17w-28}{3\lambda} + \frac{4w}{3\hat{u}^2} + \frac{2}{3\hat{u}w} - \frac{7\hat{u}+8}{12\hat{u}} \quad (\text{A.5})$$

$$\mathcal{R}_2^{(\pi)} = \frac{12(2-w)\hat{u}+8w^2}{3\hat{u}^3} \ln \frac{\hat{u}}{w^2} + \left[w \left(\frac{8w^4}{3\hat{u}^3} - \frac{40}{3} - \frac{14w}{3} - 2\hat{u} + \frac{(4w-8)w^2}{\hat{u}^2} - 4 \frac{8-8w+w^2}{\hat{u}} \right) \right. \\ \left. + 68+60\hat{u} - \frac{4}{\lambda} (15\hat{u}^3-35\hat{u}^2-76\hat{u}+14w+63w\hat{u}+19w^2) \right] \frac{\mathcal{I}_1}{\lambda} + \frac{16(1+2\hat{u})}{3\hat{u}^2} \\ - \frac{16w}{3\hat{u}^2} - \frac{28}{\hat{u}w^2} + \frac{68}{3\hat{u}w} - \frac{2(9\hat{u}^2+50\hat{u}w-201\hat{u}+86w-78)}{3\hat{u}\lambda} \\ - \frac{4(2\hat{u}^2w+2\hat{u}^3+11\hat{u}^2+49\hat{u}w-81\hat{u}+45w-28)}{\lambda^2} \quad (\text{A.6})$$

$$\mathcal{R}_3^{(\pi)} = - \frac{2(3\hat{u}^2-(2-w)\hat{u}-2w^2)}{3\hat{u}^3} \ln \frac{\hat{u}}{w^2} + \left(8w - \frac{13}{3}w^2 - 4 - \frac{10}{3}\hat{u}(w-2) - \hat{u}^2 \right. \\ \left. + \frac{10\hat{u}(w-2)w^3+4w^5}{3\hat{u}^3} \right) \frac{\mathcal{I}_1}{\lambda} - \frac{2(7\hat{u}^2+11\hat{u}w-19\hat{u}+17w-16)}{3\hat{u}\lambda} \\ - \frac{8w}{3\hat{u}^2} + \frac{8(\hat{u}+1)}{3\hat{u}^2} - \frac{4}{\hat{u}w^2} + \frac{2}{\hat{u}w} \quad (\text{A.7})$$

$$W_1^{(G,1)} = - \frac{2}{3}w \left[G_1 + \left(\frac{C_F}{4} \left(1+8w-5 \frac{w^2 \ln w}{1-w} \right) - \frac{C_A}{4}(1+2w) \right) \delta'(\hat{u}) \right. \\ \left. + C_F \left(5 + \frac{2}{w^2} - \frac{2}{w} \right) \left(\left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ - 2 \ln w \left[\frac{1}{\hat{u}} \right]_+ + L_w \delta(\hat{u}) \right) \right] \\ - \frac{2}{3} \left(\frac{C_A}{4}(8-5w) + C_F \left(\frac{4}{w} - 3 + \frac{5w}{4} \right) \right) \left[\frac{1}{\hat{u}} \right]_+ \\ - \frac{1}{3} \left(C_A \frac{5w^3-34w^2+51w-20}{2(w-1)^2} + C_F \frac{10w^5-21w^4+7w^3-10w^2+28w-16}{(w-1)^2 w} \right) \ln w \delta(\hat{u}) \\ - \frac{1}{3} \left(C_A \frac{2w^4+2w^3-3w^2+5w-4}{2(1-w)w} + C_F \frac{35w^3-25w^2-10w-8}{4(1-w)} \right) \delta(\hat{u}) + \mathcal{R}_1^{(G)} \theta(\hat{u}) \quad (\text{A.9})$$

$$W_2^{(G,1)} = - \frac{8}{3} \left[G_2 + \left(C_F \left(\frac{1}{w} - \frac{11}{4} + 2w - \left(1 - \frac{5w}{4} \right) \ln w \right) + \frac{C_A}{4}(3-2w) \right) \delta'(\hat{u}) \right. \\ \left. + 2C_F \frac{w^2-w-1}{w^3} \left(\left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ - 2 \ln w \left[\frac{1}{\hat{u}} \right]_+ + L_w \delta(\hat{u}) \right) \right] \\ + \left(C_A \left(\frac{8}{w} - \frac{9}{2w^2} - \frac{2}{w^3} - \frac{9}{4} \right) + C_F \left(\frac{7}{w^2} - \frac{6}{w^3} - \frac{5}{2} \right) \right) \left[\frac{1}{\hat{u}} \right]_+ \\ + \left(C_A \frac{9w^3-56w^2+40w+16}{4w^3} + C_F \frac{5w^4-6w^3+3w^2-12w+12}{w^3} \right) \ln w \delta(\hat{u}) \\ - \left(C_A \frac{2w^3+4w^2-23w+16}{4w^2} + C_F \frac{35w^4-98w^3+34w^2-120w+32}{8w^3} \right) \delta(\hat{u}) + \mathcal{R}_2^{(G)} \theta(\hat{u}) \quad (\text{A.10})$$

$$W_3^{(G,1)} = - \frac{4}{3} \left[G_3 + \left(C_F \left(\frac{1}{4} + \frac{5w}{2} - \frac{5w^2 \ln w}{4(1-w)} \right) - \frac{C_A}{4}(1+w) \right) \delta'(\hat{u}) \right] \\ - \frac{2}{3} C_F \frac{5w^2+4w+4}{w^2} \left(\left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ - 2 \ln w \left[\frac{1}{\hat{u}} \right]_+ + L_w \delta(\hat{u}) \right) \\ - \frac{2}{3} \left(C_A \left(\frac{2}{w^2} + \frac{5}{w} - \frac{7}{2} \right) + C_F \left(\frac{4}{w^2} - \frac{5}{4} \right) \right) \left[\frac{1}{\hat{u}} \right]_+ \\ - \frac{1}{3} \left(C_A \frac{7w^4-40w^3+49w^2-6w-8}{(w-1)^2 w^2} + C_F \frac{(20w^5-37w^4-w^3+6w^2+24w-16)}{(w-1)^2 w^2} \right) \ln w \delta(\hat{u}) \\ - \frac{2}{3} \left(C_A \frac{w^2-w+1}{1-w} + C_F \frac{35w^4-85w^3+66w^2-8w-16}{4(1-w)w^2} \right) \delta(\hat{u}) + \mathcal{R}_3^{(G)} \theta(\hat{u}) \quad (\text{A.11})$$

$$G_i = \left(1 + \frac{5}{2}w - 4\delta_{i2} \right) \left[C_F \left(\frac{3-8 \ln w}{4} \left[\frac{1}{\hat{u}^2} \right]_+ + \left[\frac{\ln \hat{u}}{\hat{u}^2} \right]_+ - L_w \delta'(\hat{u}) \right) + \frac{C_A}{2} \ln \frac{\mu}{m_b} \delta'(\hat{u}) \right] \\ + C_A \left[\frac{1+w}{2} \left[\frac{1}{\hat{u}^2} \right]_+ + \ln w \delta'(\hat{u}) \right] - \delta_{i2} \left(\frac{1+2w}{2w} \left[\frac{1}{\hat{u}^2} \right]_+ + \frac{\ln w}{w} \delta'(\hat{u}) \right) - \frac{3C_A}{4} \frac{w_i^{(G,0)}}{w_i^{(0)}} \ln \frac{\mu}{m_b} \delta(\hat{u}) \\ + C_A \left(\frac{1+4w}{2w^2} - \frac{1+2w}{w^3} \delta_{i2} \right) \left[\left[\frac{\ln \hat{u}}{\hat{u}} \right]_+ - 2 \ln w \left[\frac{1}{\hat{u}} \right]_+ + L_w \delta(\hat{u}) \right] \quad (\text{A.12})$$

where

$$w_1^{(G,0)} = - \frac{2}{3}(4-5w), \quad w_2^{(G,0)} = 0, \quad w_3^{(G,0)} = \frac{10}{3}, \quad (\text{A.13})$$

$$\mathcal{R}_1^{(G)} = \frac{C_A}{3} \left[\frac{1}{2} + \frac{\hat{u}+13w-16}{\lambda} + \frac{4w+1}{\hat{u}w} \ln \frac{\hat{u}}{w^2} + \left(\frac{4w+1-6\hat{u}}{\hat{u}} + 2 \frac{3\hat{u}(\hat{u}-3+w)+4w}{\lambda} \right) \mathcal{I}_1 \right] \\ + \frac{C_F}{3} \left[\frac{15\hat{u}-5\hat{u}w-5\hat{u}^2-11w+20}{\lambda} - \frac{4w}{\hat{u}\lambda} - \frac{10w}{\hat{u}} + \frac{8}{\hat{u}w} + \frac{11\hat{u}+24}{4\hat{u}} + \left(\frac{5w^2}{\hat{u}^2} + \frac{2(5\hat{u}+1)w}{\hat{u}^2} + \frac{4-4w}{\hat{u}w} \right) \ln \frac{\hat{u}}{w^2} \right. \\ \left. + \left(\frac{8-3\hat{u}^2-13\hat{u}w+10\hat{u}-12w}{\lambda} + \frac{5w^3}{\hat{u}^2} + \frac{(15\hat{u}+2)w^2}{\hat{u}^2} + \frac{3(5\hat{u}-8)w}{2\hat{u}} + \frac{5\hat{u}}{2} - 2 \right) \mathcal{I}_1 \right] \quad (\text{A.14})$$

$$\mathcal{R}_2^{(G)} = 4C_A \left[\frac{16-13\hat{u}^2-25\hat{u}w+51\hat{u}-29w}{\lambda^2} + \frac{22-15\hat{u}w-9\hat{u}^2+112\hat{u}-32w}{6\lambda\hat{u}} + \frac{w}{\lambda\hat{u}^2} + \frac{16\hat{u}-1}{3\hat{u}^2w} - \frac{3}{\hat{u}w^2} - \frac{4}{3\hat{u}w^3} \right. \\ \left. + \frac{4w^2-3w-2}{3w^3\hat{u}} \ln \frac{\hat{u}}{w^2} + \left(\frac{14\hat{u}^2-26\hat{u}w+58\hat{u}-3w-2}{3\lambda\hat{u}} - \frac{2(3\hat{u}^2w+3\hat{u}^3-5\hat{u}^2+20\hat{u}w-25\hat{u})}{\lambda^2} \right) - \frac{8w}{\lambda^2} + \frac{4}{3\hat{u}} \right) \mathcal{I}_1 \\ + 4C_F \left[\frac{5\hat{u}^2w+42\hat{u}w+5\hat{u}^3-4\hat{u}^2-55\hat{u}+39w-36}{\lambda^2} + \frac{4w}{\lambda^2\hat{u}} + \frac{53\hat{u}w-20\hat{u}^2-155\hat{u}+44w-52}{6\lambda\hat{u}} + \frac{14}{3\hat{u}w^2} - \frac{4}{\hat{u}w^3} - \frac{10}{3\hat{u}} \right. \\ \left. + \frac{4\hat{u}(w^2-w-1)+(w-6)w^3}{3\hat{u}^2w^3} \ln \frac{\hat{u}}{w^2} + \left(\frac{23\hat{u}^2w+13\hat{u}^3-37\hat{u}^2+47\hat{u}w-58\hat{u}+20w-8}{\lambda^2} \right. \right. \\ \left. \left. + \frac{25\hat{u}^2w+15\hat{u}^3-114\hat{u}^2+76\hat{u}w-150\hat{u}+16w-8}{6\lambda\hat{u}} + \frac{5w^2}{3\hat{u}^2} + \frac{(5\hat{u}-6)w}{3\hat{u}^2} - \frac{5\hat{u}+8}{2\hat{u}} \right) \mathcal{I}_1 \right] \quad (\text{A.15})$$

$$\mathcal{R}_3^{(G)} = \frac{4C_A}{3} \left[\frac{15\hat{u}-3\hat{u}^2-3\hat{u}w-5w-2}{2\lambda\hat{u}} + \frac{1}{\hat{u}w^2} + \frac{5}{2\hat{u}^2\hat{u}} \ln \frac{\hat{u}}{w^2} + \frac{w-5\hat{u}-2w\hat{u}+4w^2}{2\lambda\hat{u}} \mathcal{I}_1 \right] \\ + \frac{4C_F}{3} \left[\frac{2\hat{u}^2+7\hat{u}w-9\hat{u}+3w}{\lambda\hat{u}} + \frac{2}{\hat{u}w^2} - \frac{5}{\hat{u}} + \left(\frac{5w}{2\hat{u}^2} + \frac{5\hat{u}+2}{2\hat{u}^2} + \frac{2}{\hat{u}w^2} + \frac{2}{\hat{u}w} \right) \ln \frac{\hat{u}}{w^2} \right. \\ \left. + \left(\frac{5\hat{u}^2+5\hat{u}w-16\hat{u}+12w-12}{2\lambda} + \frac{5w^2}{2\hat{u}^2} + \frac{(5\hat{u}+1)w}{\hat{u}^2} - \frac{5\hat{u}+8}{4\hat{u}} \right) \mathcal{I}_1 \right] \quad (\text{A.16})$$

NLO dim-5: μ_G^2 [Capdevila et al. '22]
shaokk18@lzu.edu.cn

Theoretical Framework

Analytical Result

$$\frac{1}{\Gamma_0} \frac{d\Gamma}{dy} = \left\{ 2(3 - 2y)y^2\theta(1 - y) \right. \\ \left. - \frac{2\lambda_1}{m_b^2} \left[-\frac{5}{3}y^3\theta(1 - y) + \frac{1}{6}\delta(1 - y)\theta(1^+ - y) + \frac{1}{6}\delta'(1 - y)\theta(1^+ - y) \right] \right. \\ \left. - \frac{2\lambda_2}{m_b^2} \left[-y^2(6 + 5y)\theta(1 - y) + \frac{11}{2}\delta(1 - y)\theta(1^+ - y) \right] \right\}$$

LO: dim-3 and dim-5

$$y = \frac{2E_e}{m_c}$$

$$\frac{1}{\Gamma_0} \frac{d\Gamma}{dy} \supset \frac{1}{360} \left(5y \left(-82 + y \left(153 - 86y + 8\pi^2(-3 + 2y) \right) \right) - 24y^3 \left(-50 + y \left(5 + y(-28 + 9y) \right) \right) \text{ArcCoth}[1 + 2\sqrt{y}] + \right. \\ 12y^3 \left(125 - 2(-5 + y)y(-1 + 9y) \right) \text{ArcTanh}[1 - 2\sqrt{y}] - 816 \log[1 - \sqrt{y}] - \\ 600 \log[1 + \sqrt{y}] - 108 \log[(1 + \sqrt{y})^2] + 406 \log[1 - y] + \\ 4(6y^4(1 + y)(5 + 9y) \text{ArcTanh}[1 - 2y] + 6y^3(25 + 9y^2) \log[1 + \sqrt{y}] + 90y \log[1 - y] - \\ 105y^2 \log[1 - y] - 260y^3 \log[1 - y] - 180y^5 \log[1 - y] + 150 \log[1 - y]^2 - 360y \log[1 - y]^2 + \\ 135y^2 \log[1 - y]^2 + 30y^3 \log[1 - y]^2 - 75y^3 \log[y] + 153y^5 \log[y] - 30y^3 \log[y^5]) +$$

... **NLO: dim-3**

$$\frac{1}{\Gamma_0} \frac{d\Gamma}{dy} \supset \frac{1}{810} \left(480\pi^2 + 695y + \frac{7615y^2}{2} - 165\pi^2y^2 - \frac{545y^3}{2} + 540\pi^2y^3 - 6030y^4 + 795\pi^2y^4 + \right. \\ \left. 1800y^5 - 1230\pi^2y^5 + 300\pi^2y^6 + 1440i\pi \log[2] + 2880y \log[2] - 7740y^2 \log[2] + \right.$$

... **NLO: dim-5**

Some tips

- Up to finite power, the obtained differential decay rate is **NOT** the experimental spectrum
- Observables require integration over final states

$$\Gamma = \int \frac{d\Gamma}{dy} dy$$

$$\langle E_\ell \rangle = \frac{1}{\Gamma} \int \frac{d\Gamma}{dy} E_\ell dy$$

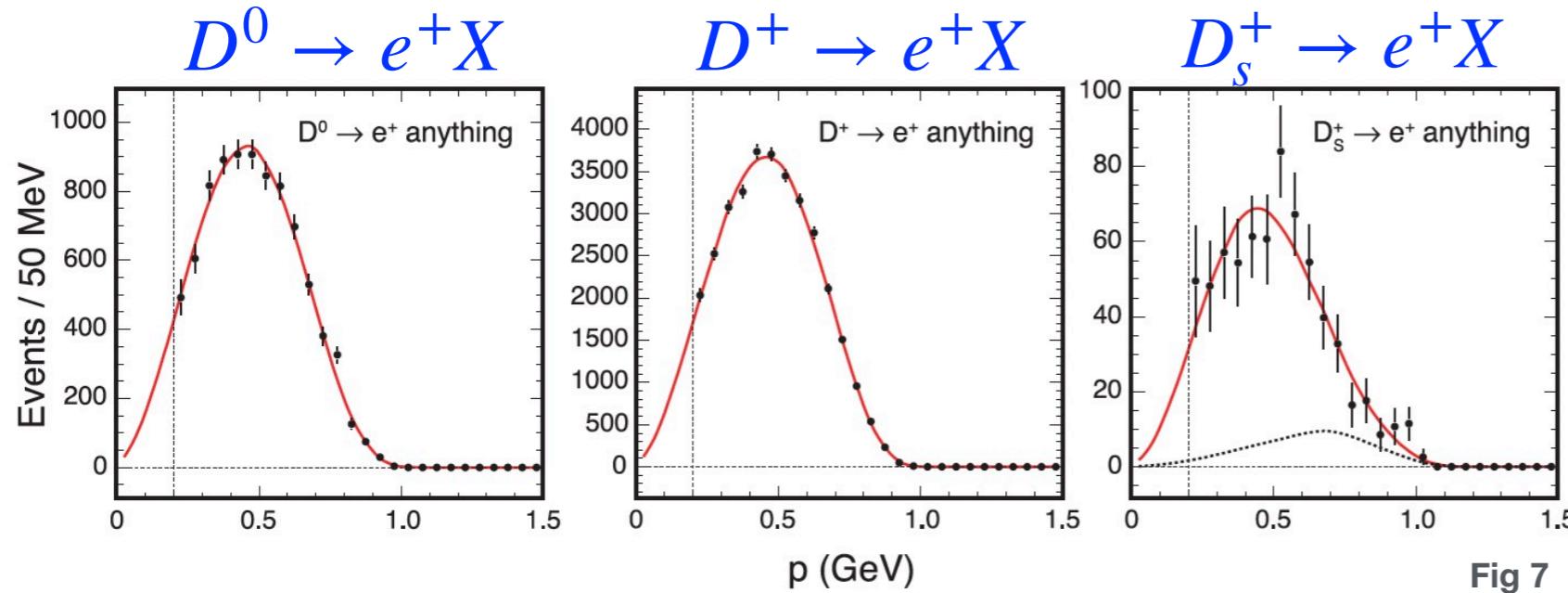
$$\langle E_\ell^2 \rangle = \frac{1}{\Gamma} \int \frac{d\Gamma}{dy} E_\ell^2 dy$$

...

3. Experimental status

Experimental status

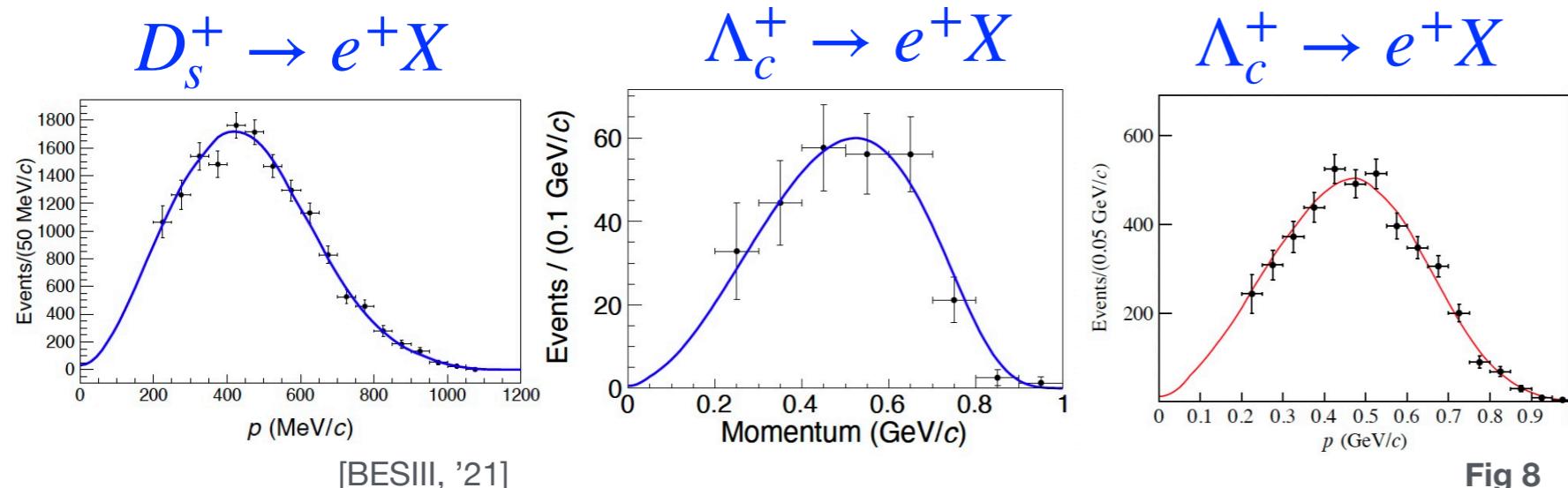
CLEO measurements



$$B(D^0 \rightarrow X e^+ \nu_e) = (6.46 \pm 0.09 \pm 0.11)\%$$
$$B(D^0 \rightarrow X e^+ \nu_e) = (6.46 \pm 0.09 \pm 0.11)\%$$
$$B(D_s^+ \rightarrow X e^+ \nu_e) = (6.52 \pm 0.39 \pm 0.15)\%$$

Fig 7

BESIII measurements



$$B(D_s^+ \rightarrow X e^+ \nu_e) = (6.30 \pm 0.13 \pm 0.10)\%$$

[BESIII, '21]

$$B(\Lambda_c^+ \rightarrow X e^+ \nu_e) = (3.95 \pm 0.34 \pm 0.09)\%$$

[BESIII (567 pb^{-1}), '18]

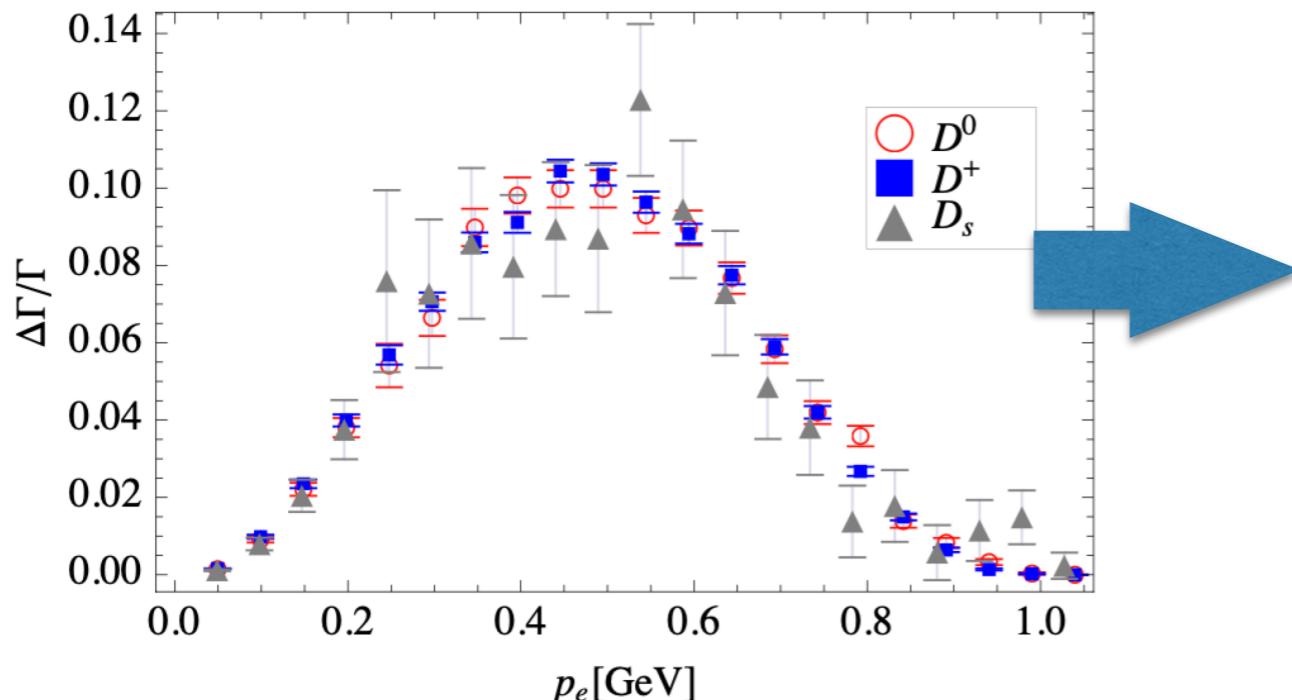
$$B(\Lambda_c^+ \rightarrow X e^+ \nu_e) = (4.06 \pm 0.10_{\text{stat.}} \pm 0.09_{\text{syst.}})\%$$

[BESIII (4.5 fb^{-1}), '23]

Fig 8

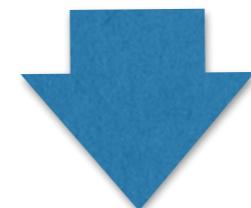
Experimental status

Electronic energy momentum



The laboratory frame of the D meson

$$\begin{aligned}\langle E_e \rangle_{lab}^{D^0} &= 0.465(3) \text{ GeV}, \\ \langle E_e \rangle_{lab}^{D^+} &= 0.459(1) \text{ GeV}, \\ \langle E_e \rangle_{lab}^{D_s} &= 0.466(12) \text{ GeV}, \\ \langle E_e^2 \rangle_{lab}^{D^0} &= 0.248(2) \text{ GeV}^2, \\ \langle E_e^2 \rangle_{lab}^{D^+} &= 0.242(1) \text{ GeV}^2, \\ \langle E_e^2 \rangle_{lab}^{D_s} &= 0.254(13) \text{ GeV}^2.\end{aligned}$$



D mesons rest frame

$$\begin{array}{ll}\langle E_\ell \rangle_{exp}^{D^0} = 0.459(3) \text{ GeV}, & \langle E_\ell^2 \rangle_{exp}^{D^0} = 0.240(2) \text{ GeV}^2, \\ \langle E_\ell \rangle_{exp}^{D^+} = 0.455(1) \text{ GeV}, & \langle E_\ell^2 \rangle_{exp}^{D^+} = 0.236(1) \text{ GeV}^2, \\ \langle E_\ell \rangle_{exp}^{D_s} = 0.456(11) \text{ GeV}, & \langle E_\ell^2 \rangle_{exp}^{D_s} = 0.239(12) \text{ GeV}^2,\end{array}$$

[Gambino,Kamenik, '10]

4. Phenomenological Analysis (preliminary)

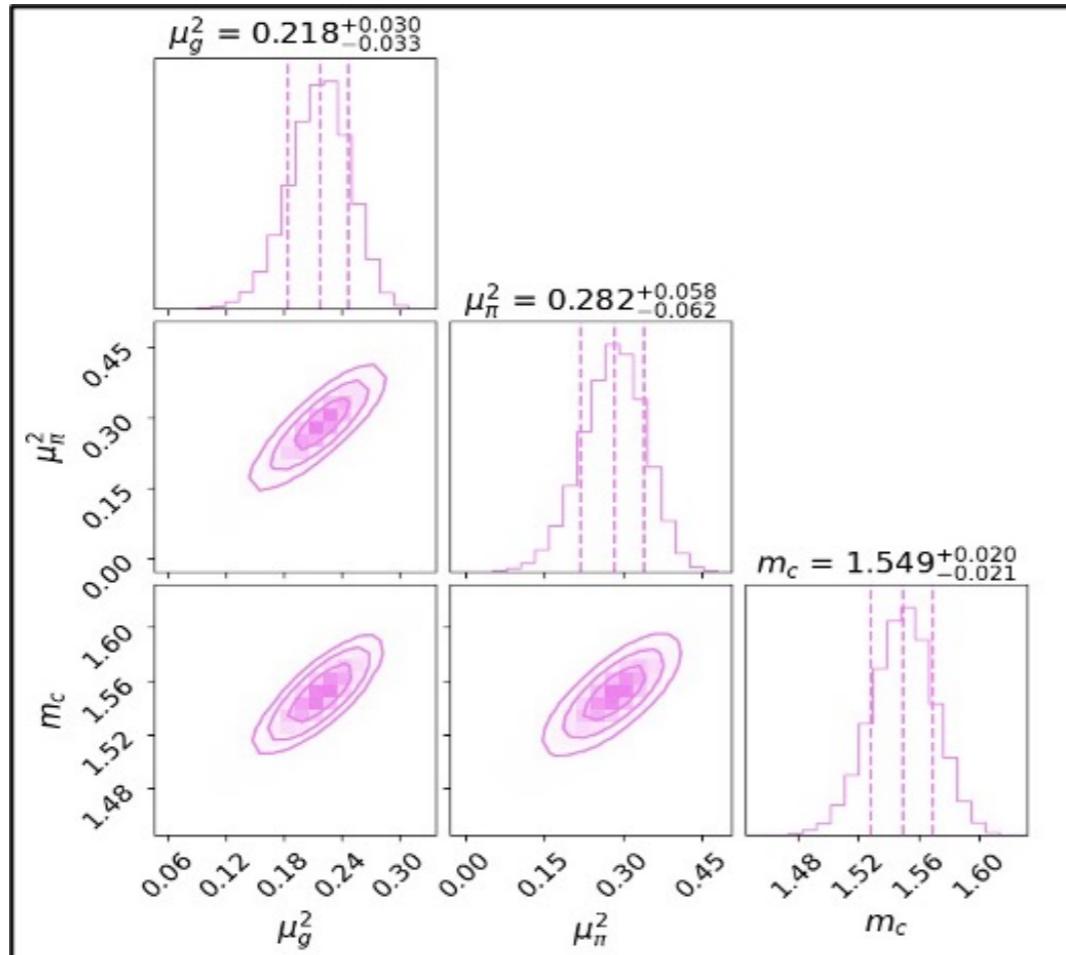
<https://github.com/ChunHuangPhy/CompactObject>

Phenomenological Analysis (preliminary)

CompactObject

[Huang et al, '2024]

★ For the first time, we systematically extract the Model-independent fundamental parameters of HQE from experimental data on semi-leptonic inclusive decays of D mesons.



$$\mu_\pi^2(D) = (0.48 \pm 0.20)\text{GeV}^2$$

$$\mu_G^2(D) = (0.34 \pm 0.10)\text{GeV}^2$$

[Lenz et al, '22]

$$\mu_\pi^2(D) = (0.282^{+0.058}_{-0.062})\text{GeV}^2$$

$$\mu_G^2(D) = (0.218^{+0.030}_{-0.033})\text{GeV}^2$$

$$m_c = (1.549^{+0.020}_{-0.021})\text{GeV}$$

[Our results]

✓ Higher precision

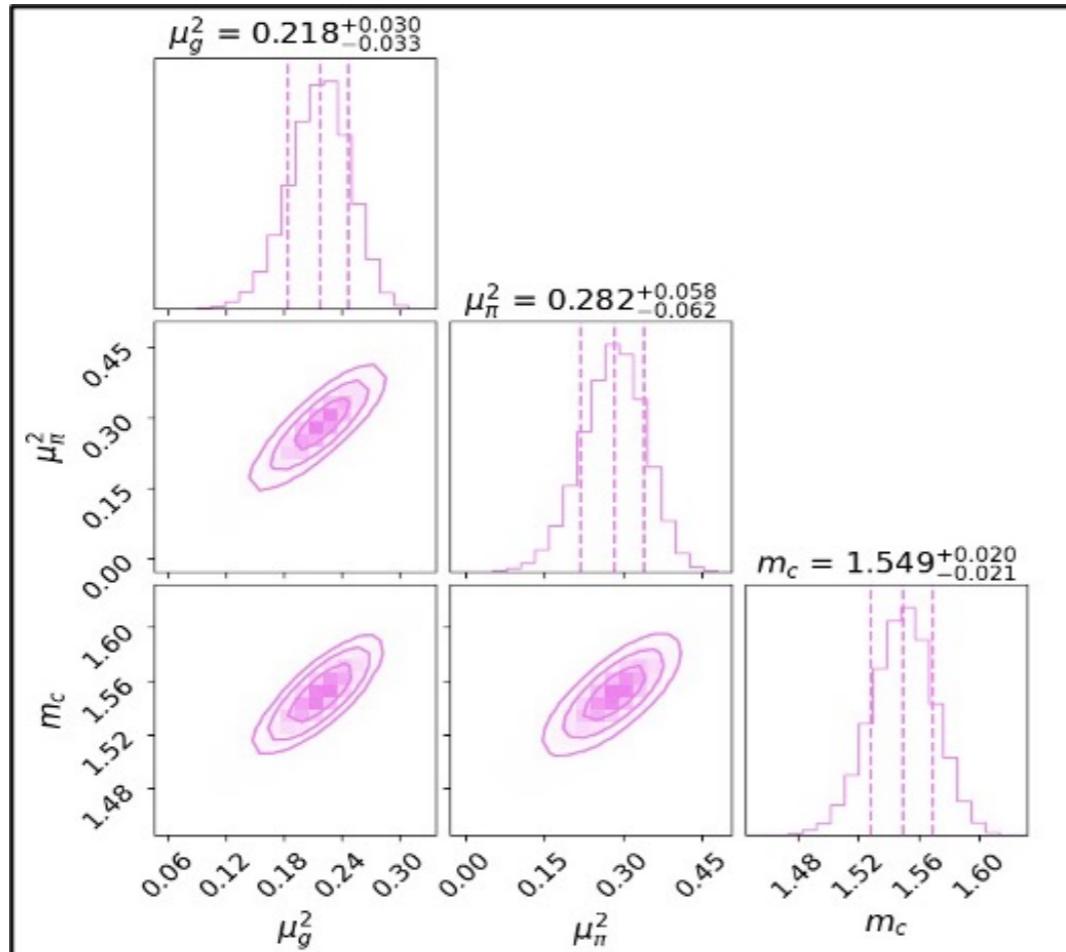
✓ Stronger constraints

Phenomenological Analysis (preliminary)

CompactObject

[Huang et al, '2024]

★ For the first time, we systematically extract the Model-independent fundamental parameters of HQE from experimental data on semi-leptonic inclusive decays of D mesons.



$$\mu_\pi^2(D) = (0.48 \pm 0.20) \text{GeV}^2$$

$$\mu_G^2(D) = (0.34 \pm 0.10) \text{GeV}^2$$

[Lenz et al, '22]

$$\mu_\pi^2(D) = (0.282^{+0.058}_{-0.062}) \text{GeV}^2$$

$$\mu_G^2(D) = (0.218^{+0.030}_{-0.033}) \text{GeV}^2$$

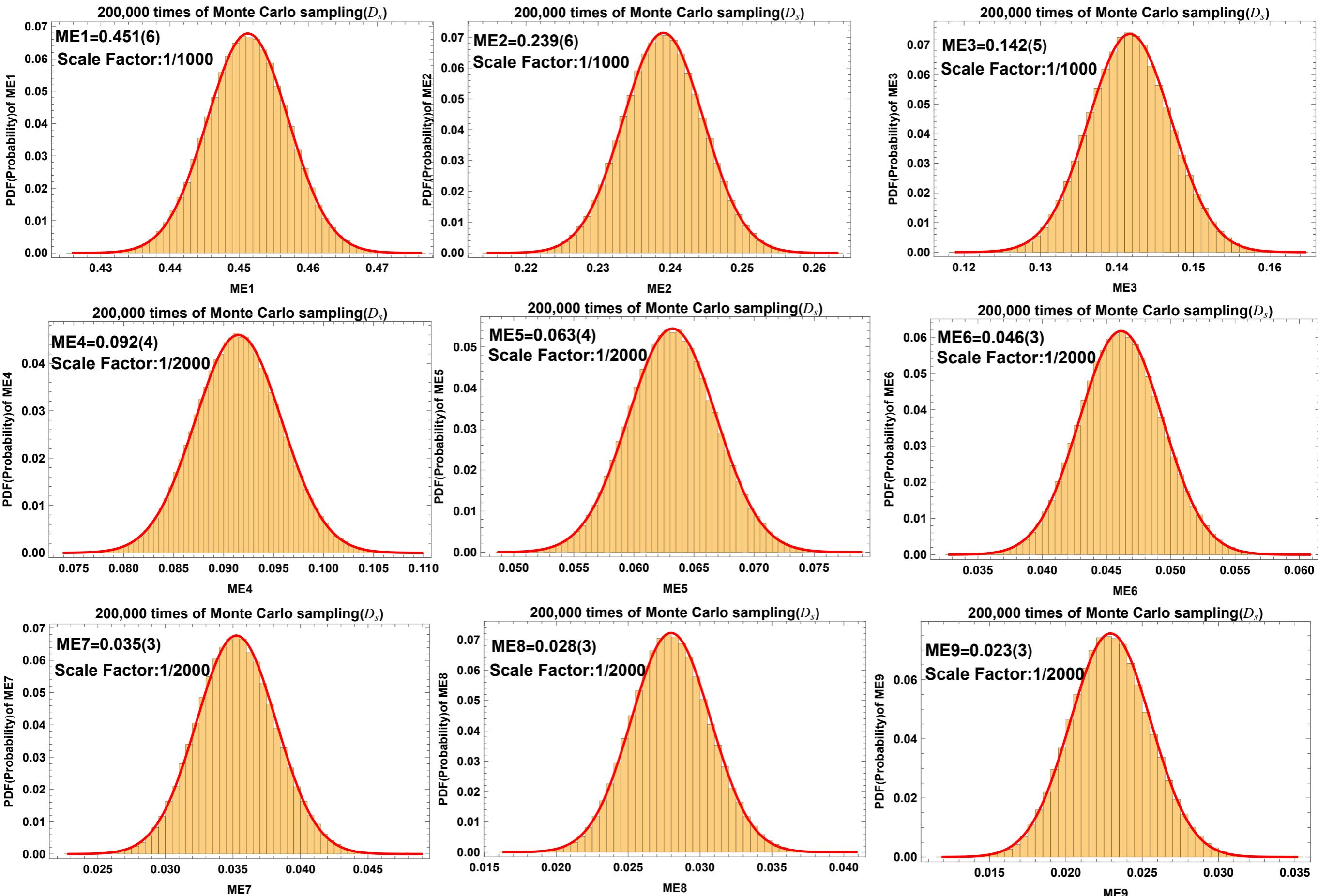
$$m_c = (1.549^{+0.020}_{-0.021}) \text{GeV}$$

✓ Higher precision

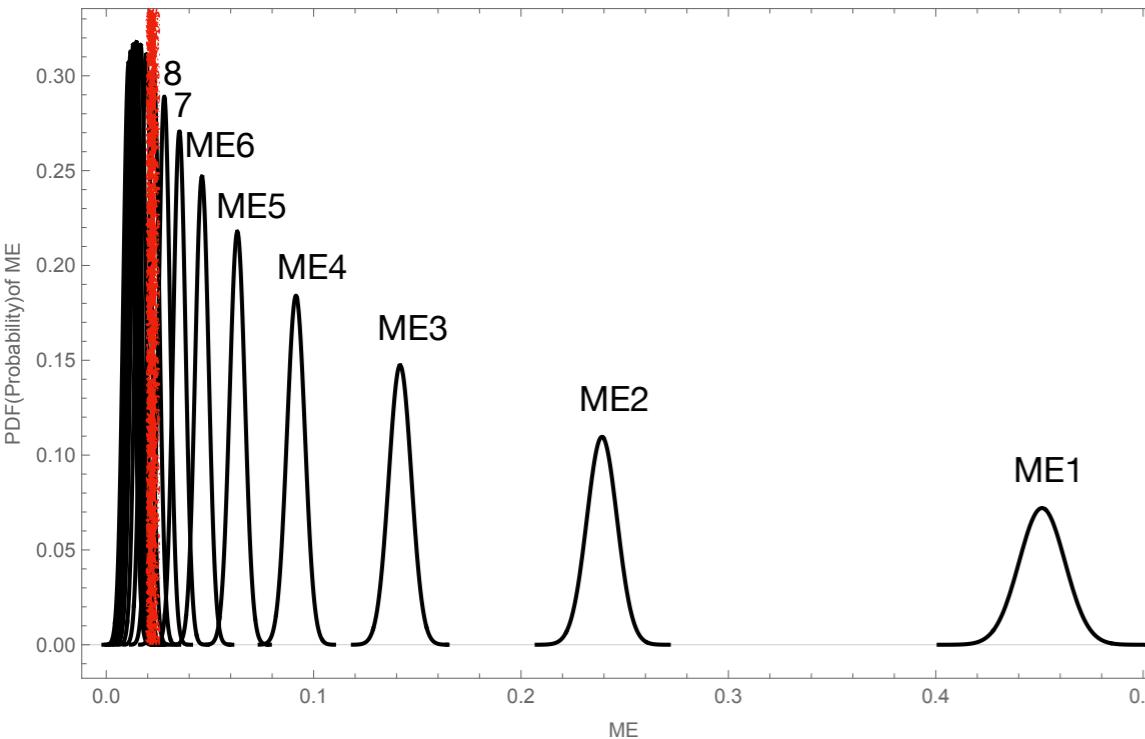
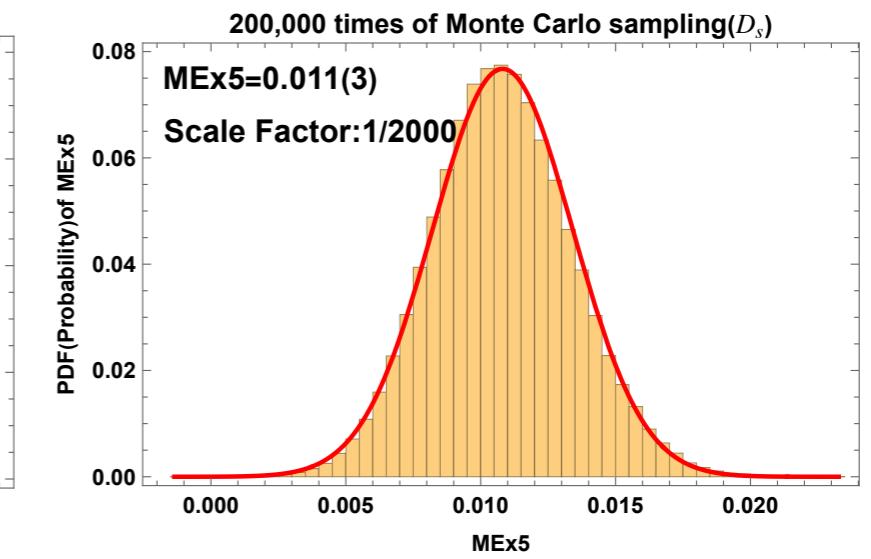
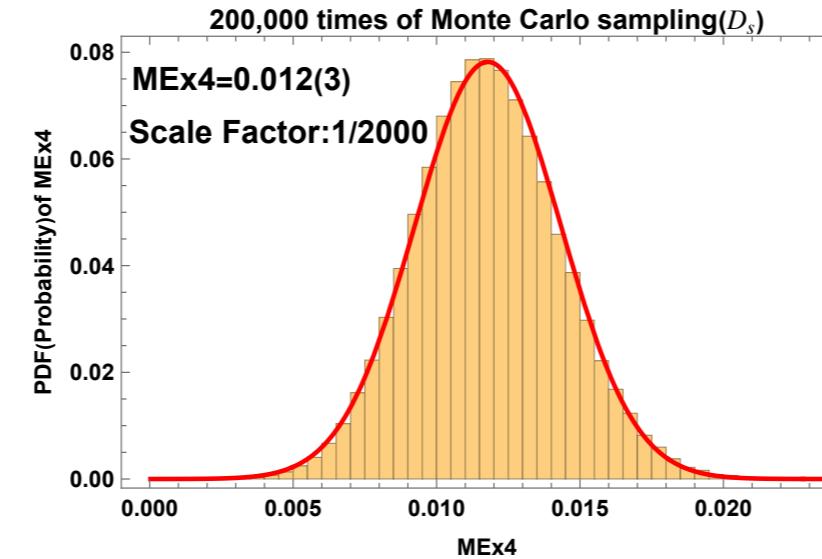
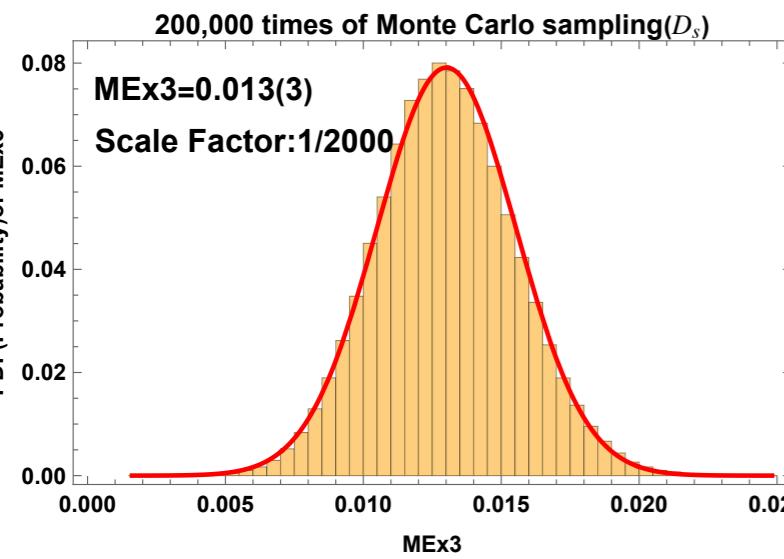
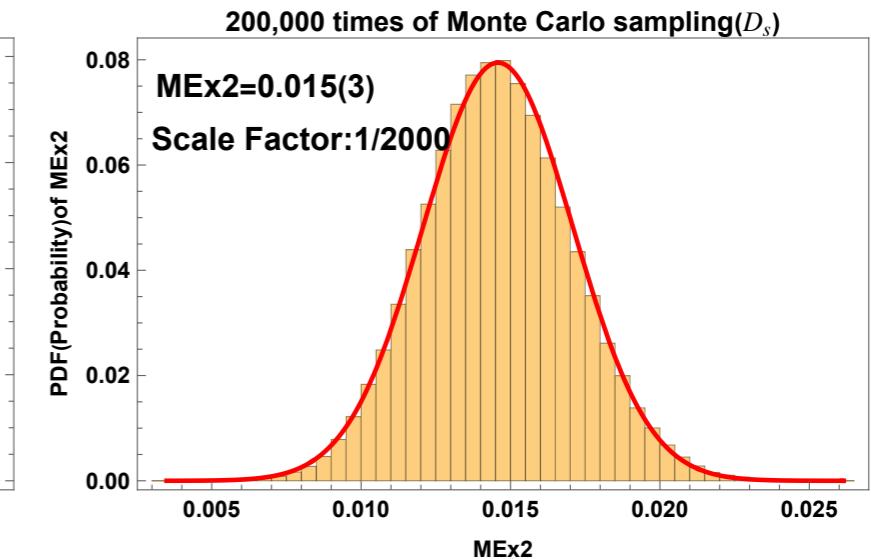
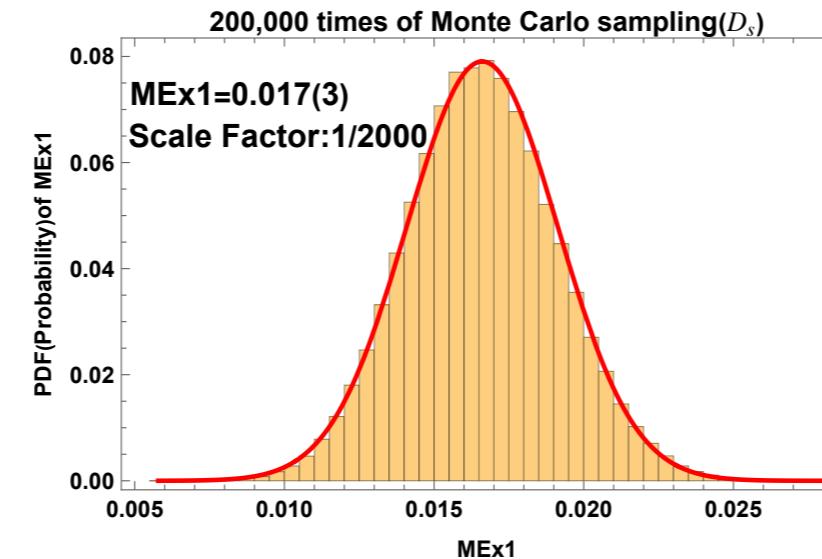
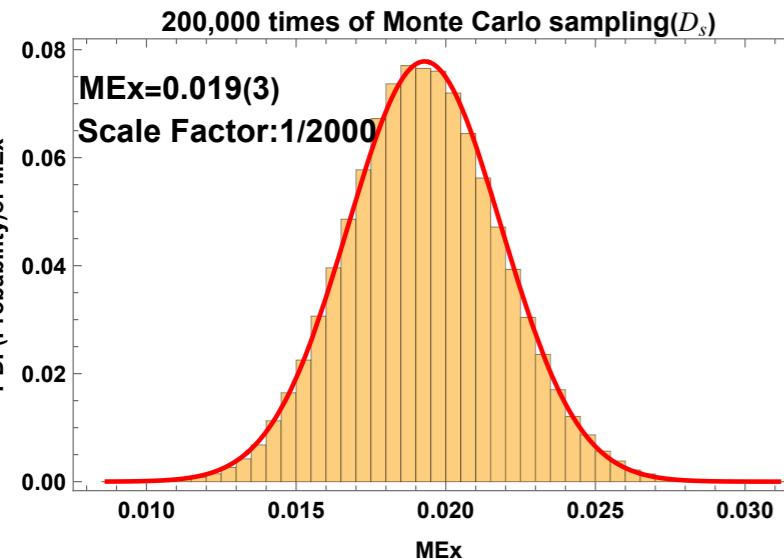
[Our results]

✓ Stronger constraints

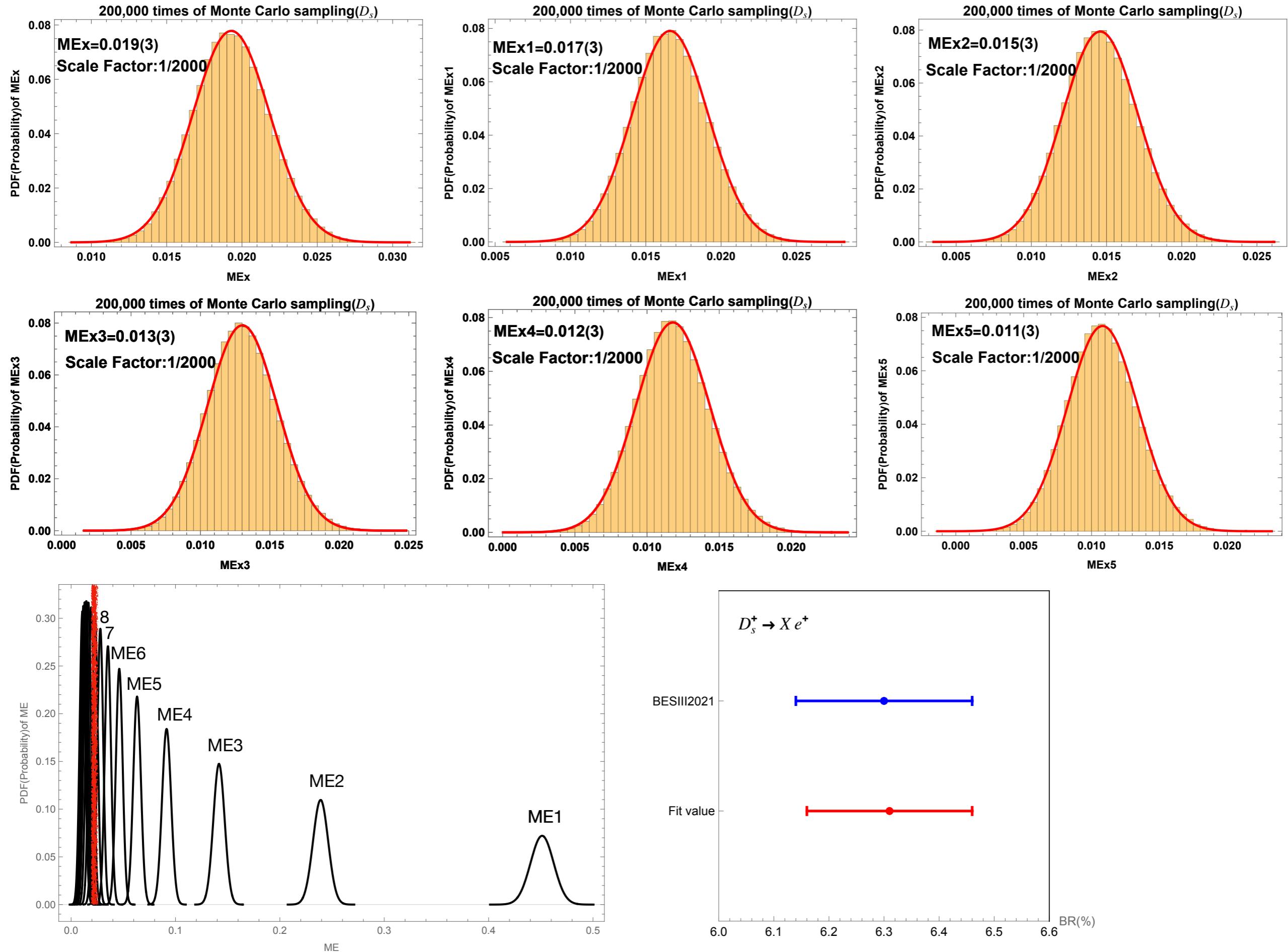
Lab Frame



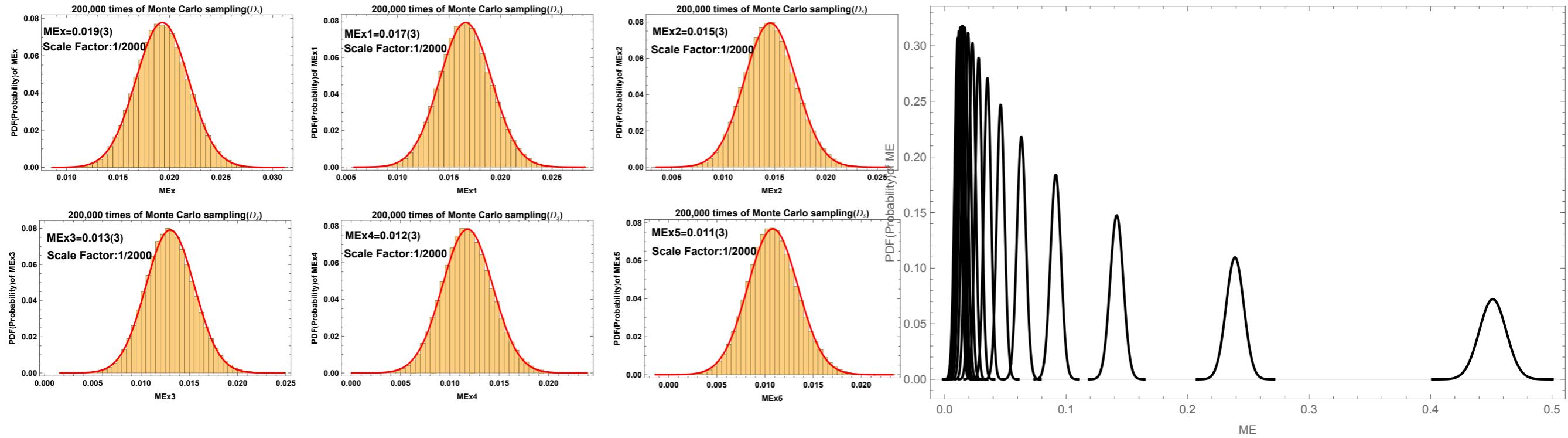
Lab Frame



Lab Frame



Lab Frame



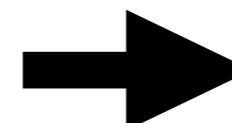
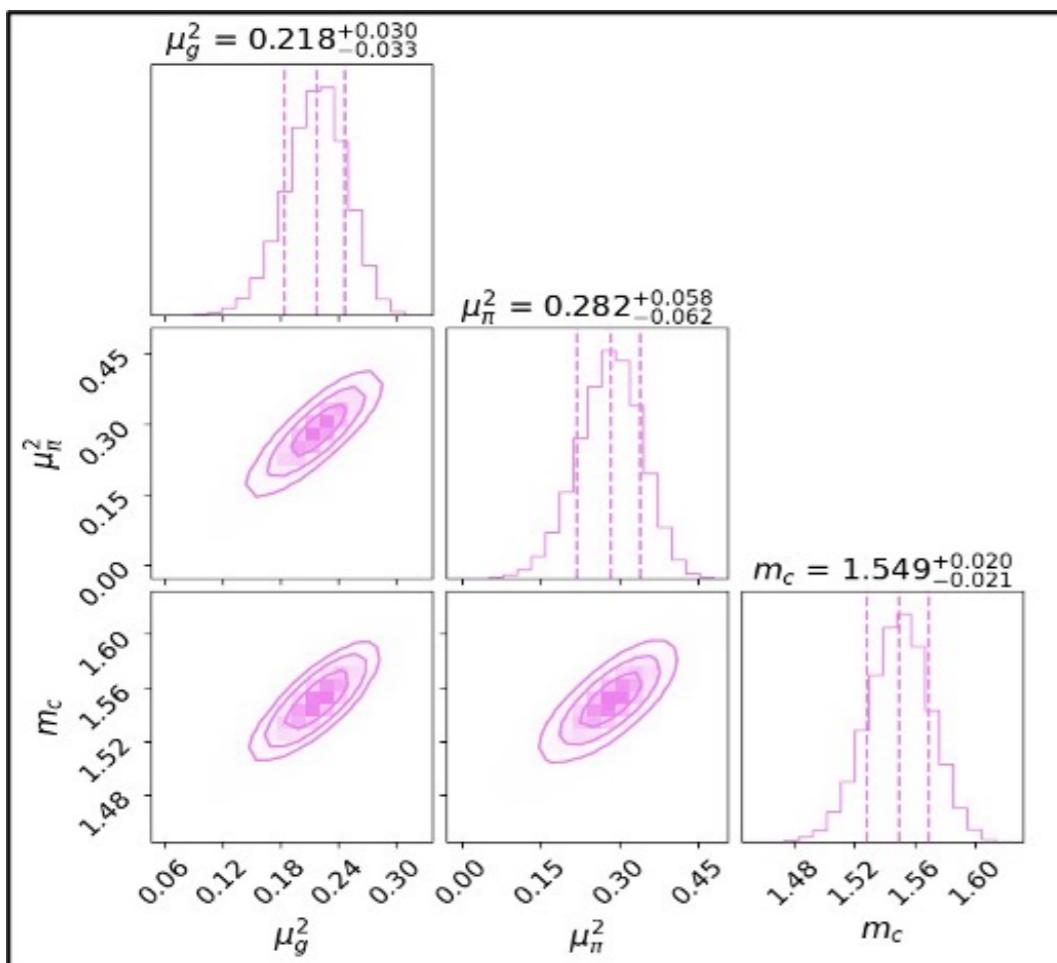
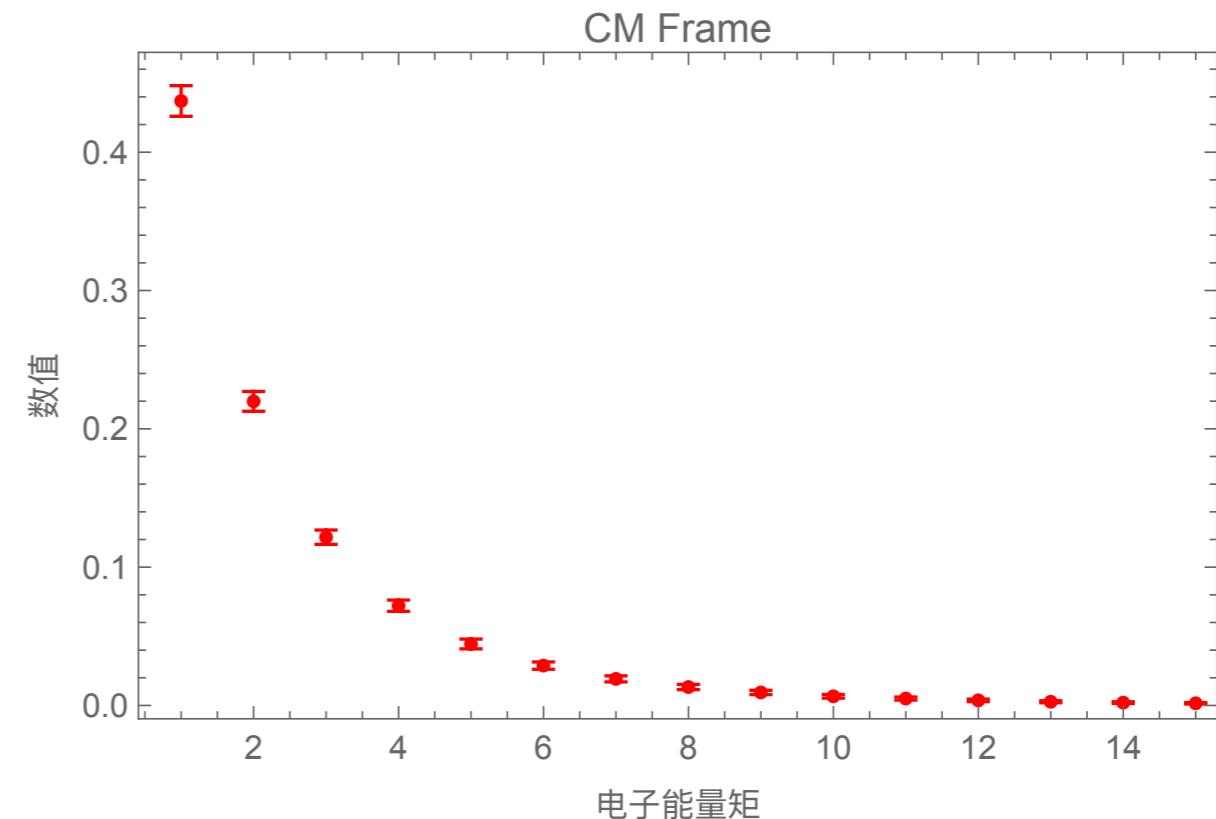
ME1	ME2	ME3	ME4	ME5	ME6	ME7	ME8	ME9	MEx	MEx1	MEx2	MEx3	MEx4	MEx5
<i>//MatrixForm=</i>														
1.	0.936881	0.830061	0.725193	0.631639	0.551242	0.483639	0.427467	0.380973	0.342399	0.310188	0.283049	0.259952	0.240092	0.222844
0.936881	1.	0.967706	0.90011	0.821025	0.742009	0.668842	0.603852	0.547373	0.49875	0.456954	0.420911	0.389647	0.362335	0.338296
0.830061	0.967706	1.	0.978961	0.930231	0.86908	0.804952	0.743163	0.686227	0.634967	0.58931	0.548778	0.512763	0.480654	0.4519
0.725193	0.90011	0.978961	1.	0.984712	0.947972	0.900444	0.849159	0.798258	0.749891	0.704974	0.66374	0.626072	0.5917	0.560306
0.631639	0.821025	0.930231	0.984712	1.	0.988516	0.96047	0.923459	0.882565	0.840891	0.800177	0.761306	0.72466	0.690341	0.658301
0.551242	0.742009	0.86908	0.947972	0.988516	1.	0.991273	0.969692	0.940677	0.907884	0.873633	0.839321	0.805755	0.773375	0.742402
0.483639	0.668842	0.804952	0.900444	0.96047	0.991273	1.	0.993296	0.976489	0.953456	0.926851	0.898428	0.869318	0.840238	0.811635
0.427467	0.603852	0.743163	0.849159	0.923459	0.969692	0.993296	1.	0.994772	0.981464	0.962876	0.940968	0.917089	0.892162	0.866815
0.380973	0.547373	0.686227	0.798258	0.882565	0.940677	0.976489	0.994772	1.	0.995845	0.985106	0.969834	0.951507	0.931183	0.909623
0.342399	0.49875	0.634967	0.749891	0.840891	0.907884	0.953456	0.981464	0.995845	1.	0.996631	0.987795	0.975026	0.959462	0.941949
0.310188	0.456954	0.58931	0.704974	0.800177	0.873633	0.926851	0.962876	0.985106	0.996631	1.	0.997214	0.989811	0.978963	0.965565
0.283049	0.420911	0.548778	0.66374	0.761306	0.839321	0.898428	0.940968	0.969834	0.987795	0.997214	1.	0.997655	0.991352	0.982006
0.259952	0.389647	0.512763	0.626072	0.72466	0.805755	0.869318	0.917089	0.951507	0.975026	0.989811	0.997655	1.	0.997995	0.992554
0.240092	0.362335	0.480654	0.5917	0.690341	0.773375	0.840238	0.892162	0.931183	0.959462	0.978963	0.991352	0.997995	1.	0.998263
0.222844	0.338296	0.4519	0.560306	0.658301	0.742402	0.811635	0.866815	0.909623	0.941949	0.965565	0.982006	0.992554	0.998263	1.

CM Frame

```

= Solve[ME1 == Around[0.451, 0.011], ME1CM] (*CM系下的电子能量一阶矩*)
Solve[ME2 == Around[0.239, 0.007], ME2CM] (*CM系下的电子能量二阶矩*)
Solve[ME3 == Around[0.142, 0.005], ME3CM] (*CM系下的电子能量三阶矩*)
Solve[ME4 == Around[0.092, 0.004], ME4CM] (*CM系下的电子能量四阶矩*)
Solve[ME5 == Around[0.063, 0.004], ME5CM] (*CM系下的电子能量五阶矩*)
Solve[ME6 == Around[0.046, 0.003], ME6CM] (*CM系下的电子能量六阶矩*)
Solve[ME7 == Around[0.035, 0.003], ME7CM] (*CM系下的电子能量七阶矩*)
Solve[ME8 == Around[0.028, 0.003], ME8CM] (*CM系下的电子能量八阶矩*)
Solve[ME9 == Around[0.023, 0.003], ME9CM] (*CM系下的电子能量九阶矩*)
Solve[MEx == Around[0.019, 0.003], MExCM] (*CM系下的电子能量十阶矩*)
Solve[MEx1 == Around[0.017, 0.003], MEx1CM] (*CM系下的电子能量十一阶矩*)
Solve[MEx2 == Around[0.015, 0.003], MEx2CM] (*CM系下的电子能量十二阶矩*)
Solve[MEx3 == Around[0.013, 0.003], MEx3CM] (*CM系下的电子能量十三阶矩*)
Solve[MEx4 == Around[0.012, 0.003], MEx4CM] (*CM系下的电子能量十四阶矩*)
Solve[MEx5 == Around[0.011, 0.003], MEx5CM] (*CM系下的电子能量十五阶矩*)

```



$$\mu_\pi^2(D) = (0.48 \pm 0.20)\text{GeV}^2$$

$$\mu_G^2(D) = (0.34 \pm 0.10)\text{GeV}^2$$

[Lenz et al, '22]

好期待.....

$$\mu_\pi^2(D) =$$

$$\mu_G^2(D) =$$

$$m_c =$$

...



What did we do?

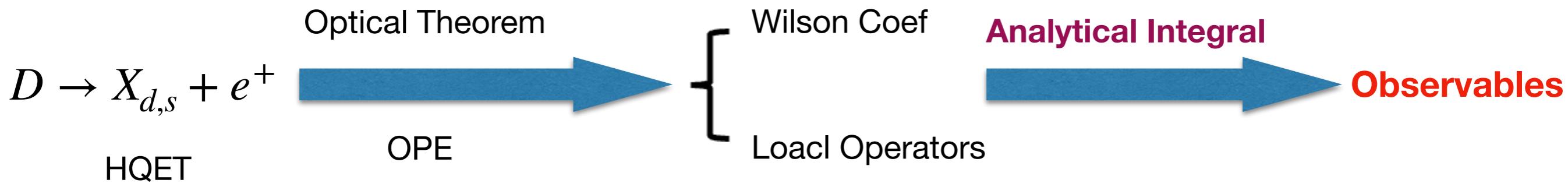
Theory:

Experiment:

Phenomenological Analysis:

What did we do?

Theory:

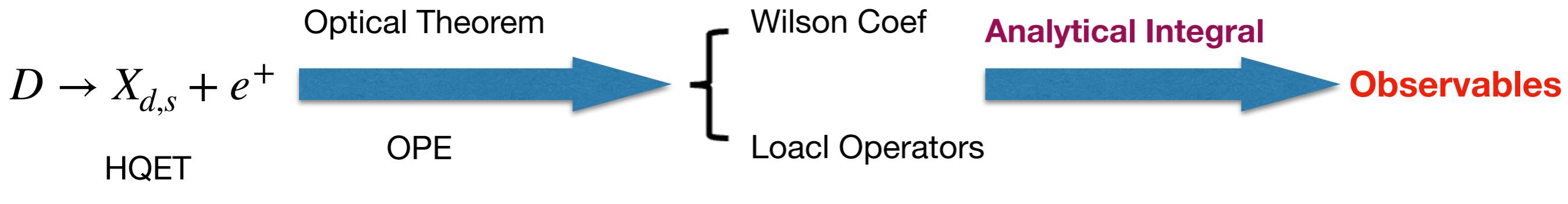


Experiment:

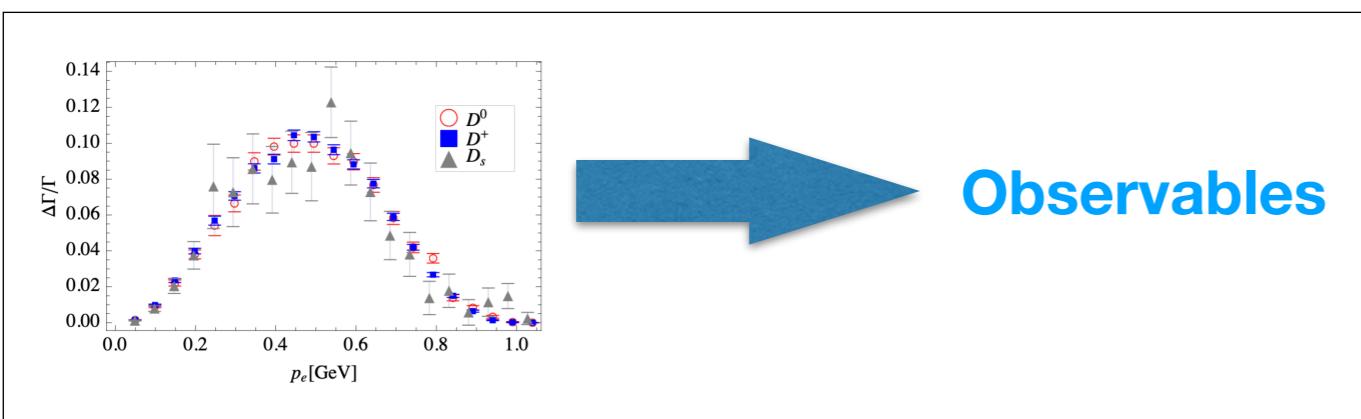
Phenomenological Analysis:

What did we do?

Theory:



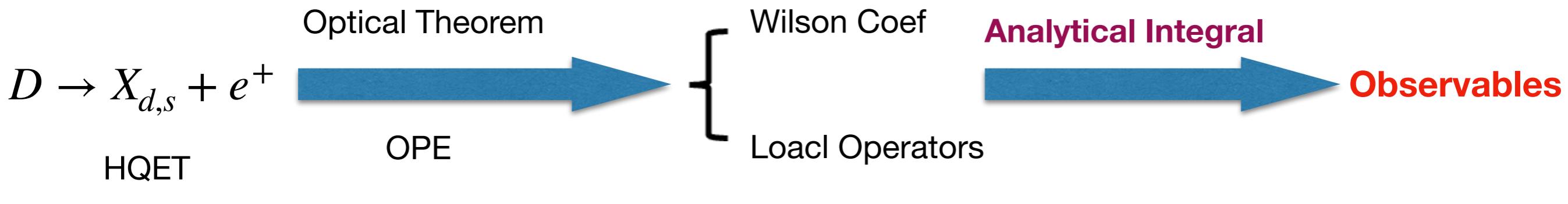
Experiment:



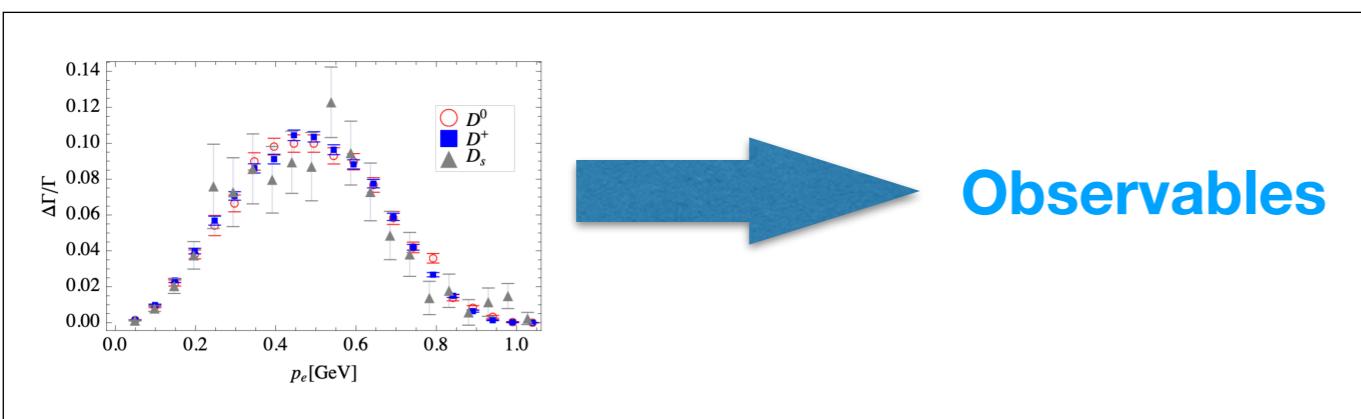
Phenomenological Analysis:

What did we do?

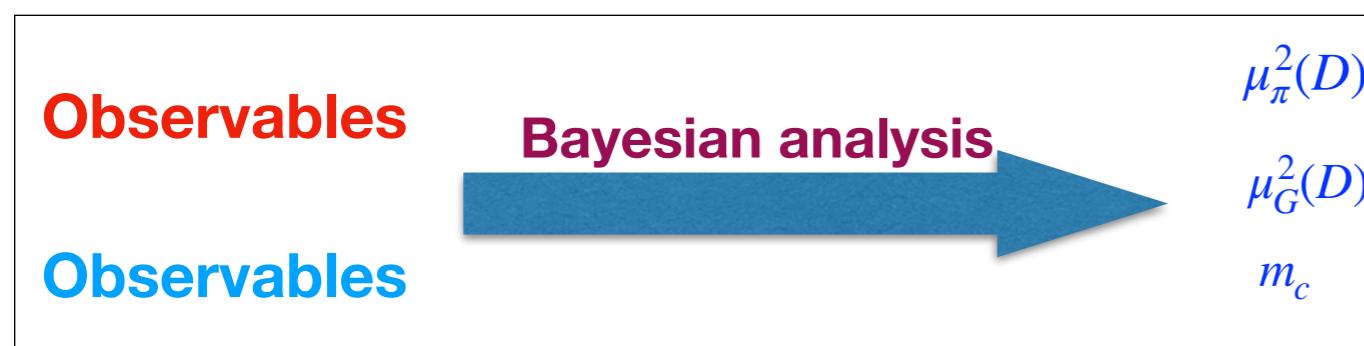
Theory:



Experiment:

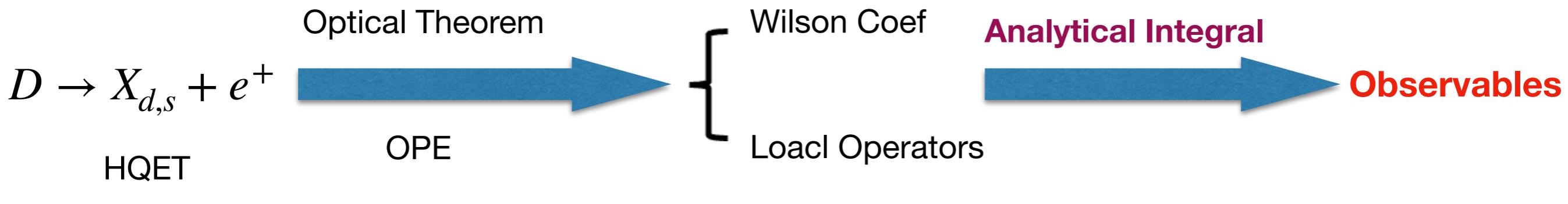


Phenomenological Analysis:

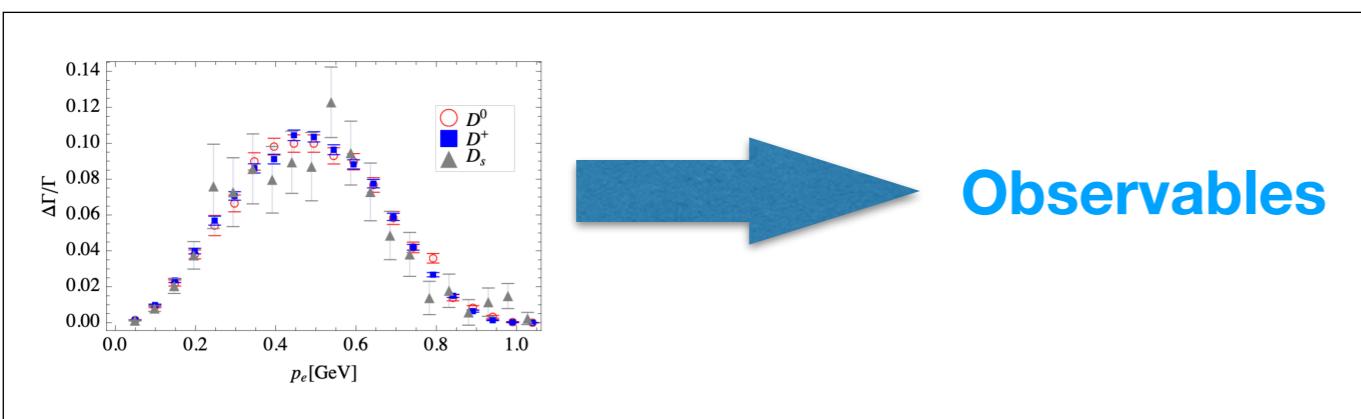


What did we do?

Theory:



Experiment:



Observables



Phenomenological Analysis:



Wishlist

- Precision measurements of leptonic energy spectrum **in the rest frame** of charmed hadrons
- q^2 spectrum, good for **higher-dimensional operators**
- Separate X_d , X_s , to give **first measurements** of V_{cd} , V_{cs}
- Rare decays: $D \rightarrow X_u \ell \ell$

Thank you!

Wen-jie Song (宋雯捷), Dong-Xiao (肖栋) , Ying-ao Tang (唐迎澳), Ji-xin Yu(余纪新),
Yong-Zheng(郑勇), Bo-nan zhang(张博楠), Yin-fa Shen (沈胤发) et, al

Appendix

Summary: observable

$$\Gamma = \frac{G_F^2 m_c^5 |V_{CKM}|}{192\pi^3}$$

Cons

$$\langle E_e \rangle = \frac{G_F^2 m_c^5 |V_{CKM}|}{192\pi^3 \Gamma}$$

$$\langle E_e^2 \rangle = \frac{G_F^2 m_c^5 |V_{CKM}|}{192\pi^3 \Gamma}$$

M1=M1

M2=M1+M2

M3=M1+M2+M3

M4=M1+M2+M3+M4

$$\text{Cons} = \frac{=}{2.28252 \times 10^{-14} m_c^5}$$

Prior Distribution for Free Parameters

$$m_c \in [1, 2] \text{ GeV}$$

$$\begin{array}{|c|c|} \hline & \mu_\pi^2 \in [-0.28, 1.08] \text{ GeV}^2 \\ \hline & \mu_G^2 \in [0.04, 0.64] \text{ GeV}^2 \\ \hline \end{array}$$

A: Daniel King,
Alexander Lenz et al.
“Revisiting inclusive
decay widths of
charmed mesons”,
[JHEP08\(2022\)241](#).

$$\text{Fixed } \alpha_s(\bar{m}_c = 1.273 \text{ GeV}) = 0.378387$$

Prof Qin's code: running from
 $\alpha_s(m_z = 91.1880 \text{ GeV}) = 0.1179$ at four loop level.

Note!

M1=M1

Tree Level dim-3 operator

M2=M1+M2

Tree Level : dim-3 operator + μ_π^2

M3=M1+M2+M3

Tree Level : dim-3 operator + $\mu_\pi^2 + \mu_G^2$

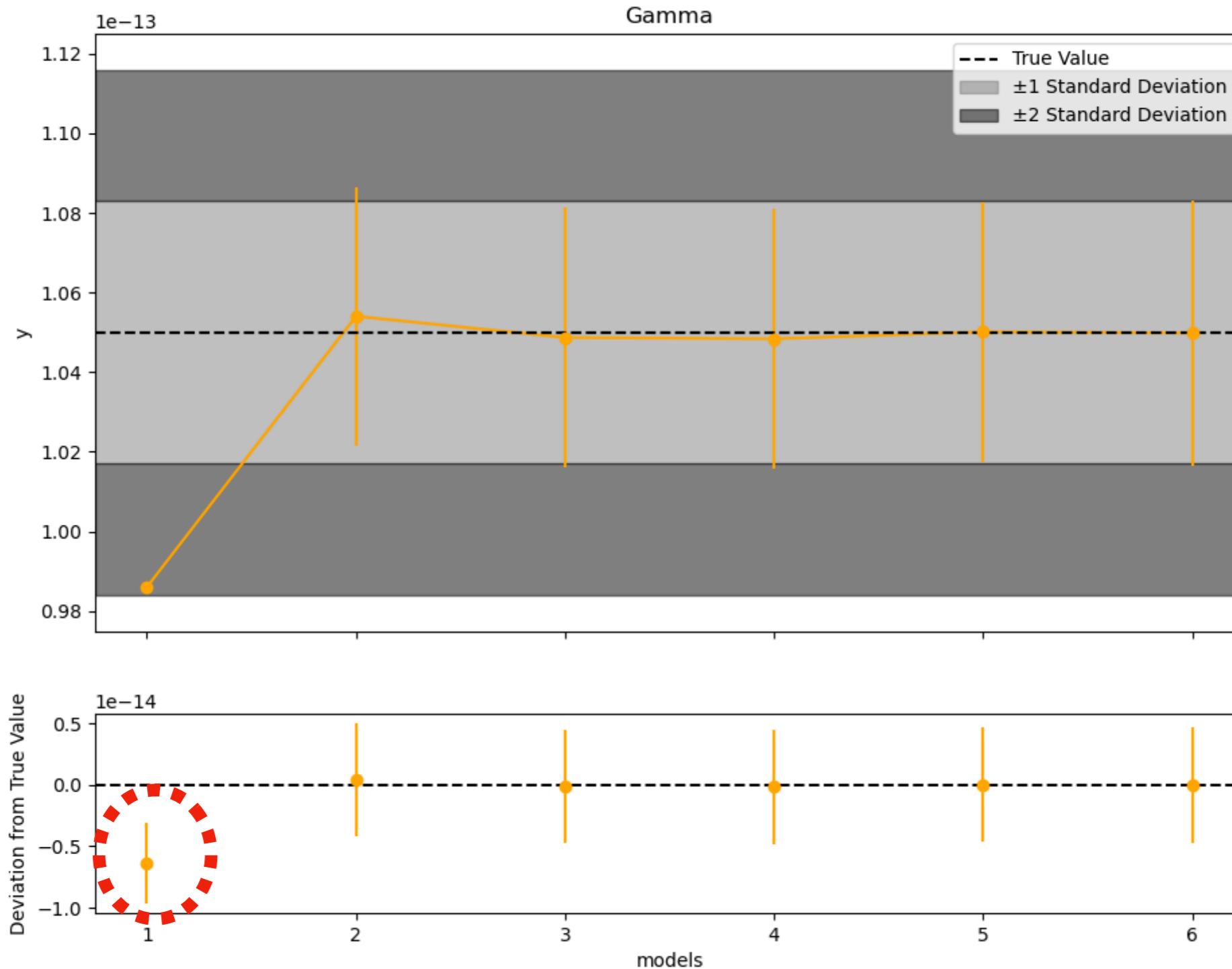
M4=M1+M2+M3+M4

Tree Level : dim-3 operator + $\mu_\pi^2 + \mu_G^2$ + dim-3 (NLO)

Phenomenological Analysis (preliminary)

CompactObject

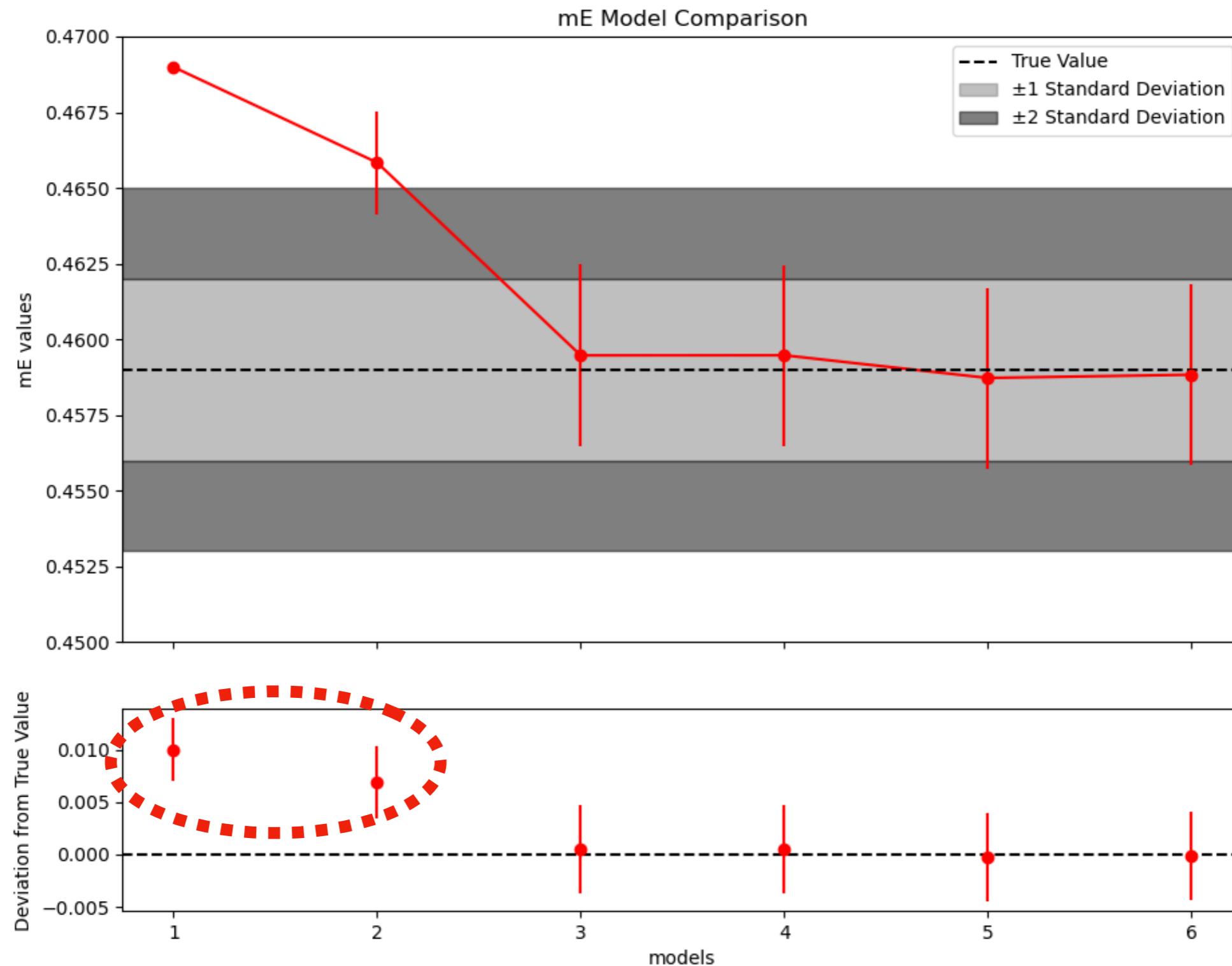
[Huang et al, '2024]



Phenomenological Analysis (preliminary)

CompactObject

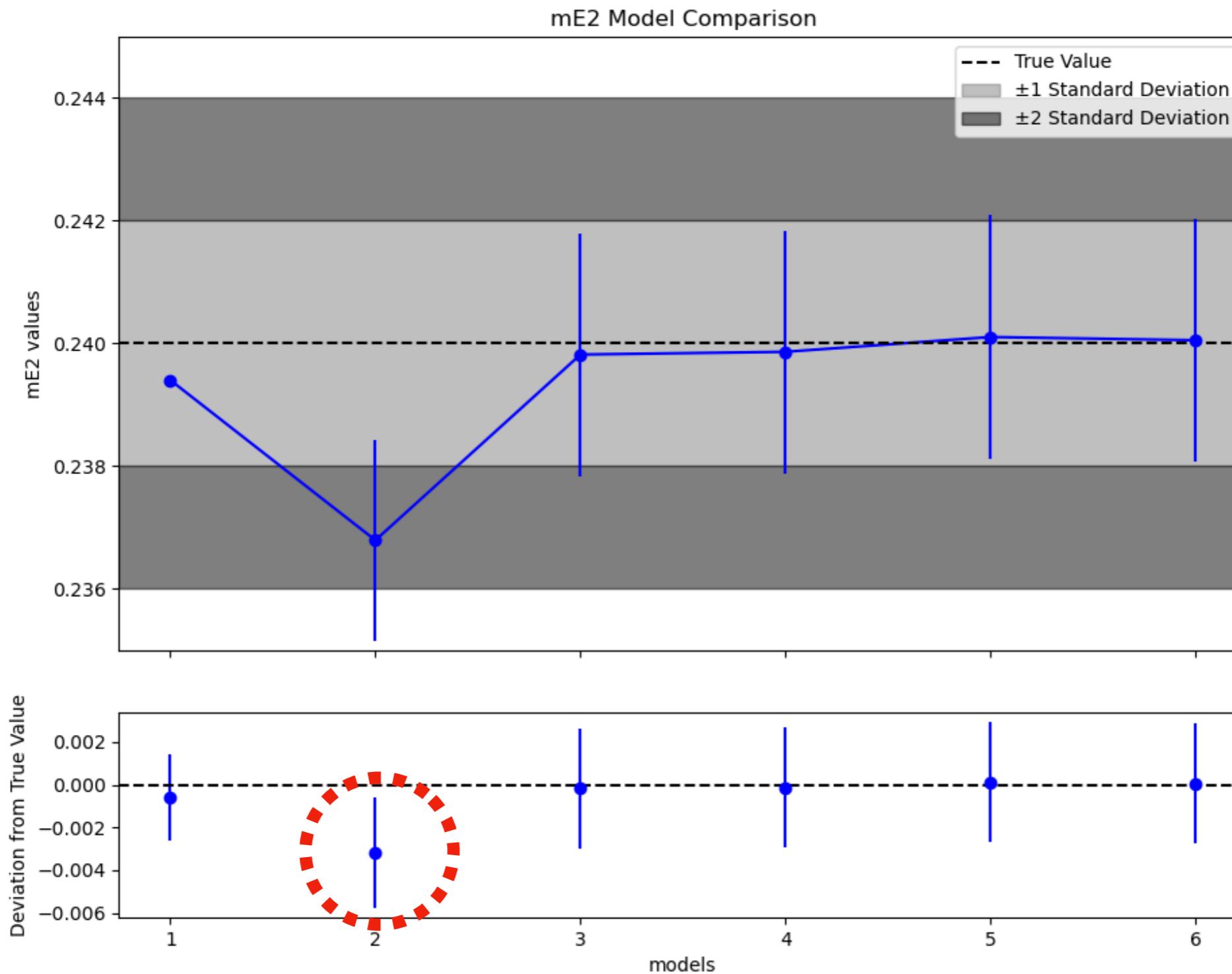
[Huang et al, '2024]



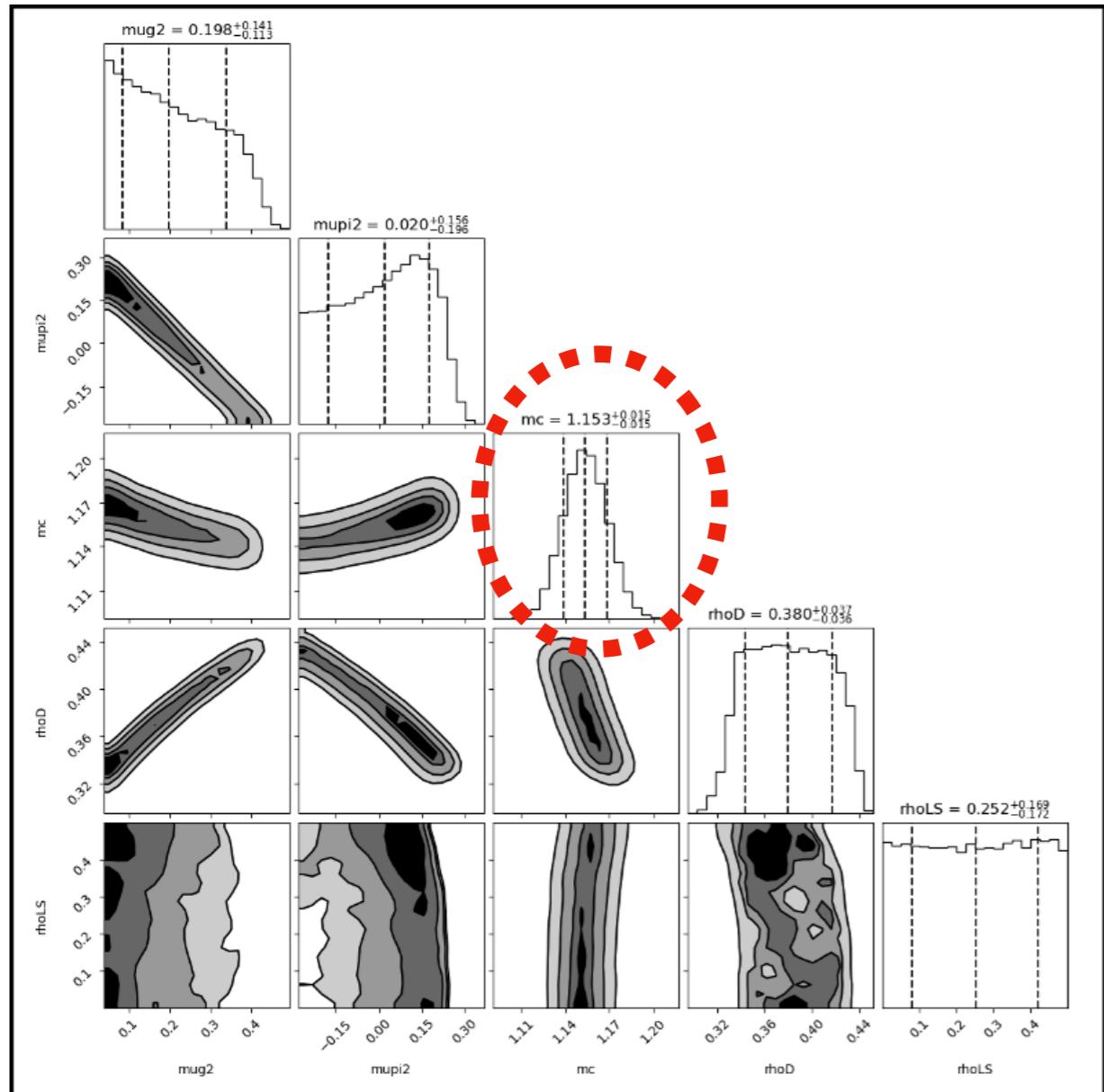
Phenomenological Analysis (preliminary)

CompactObject

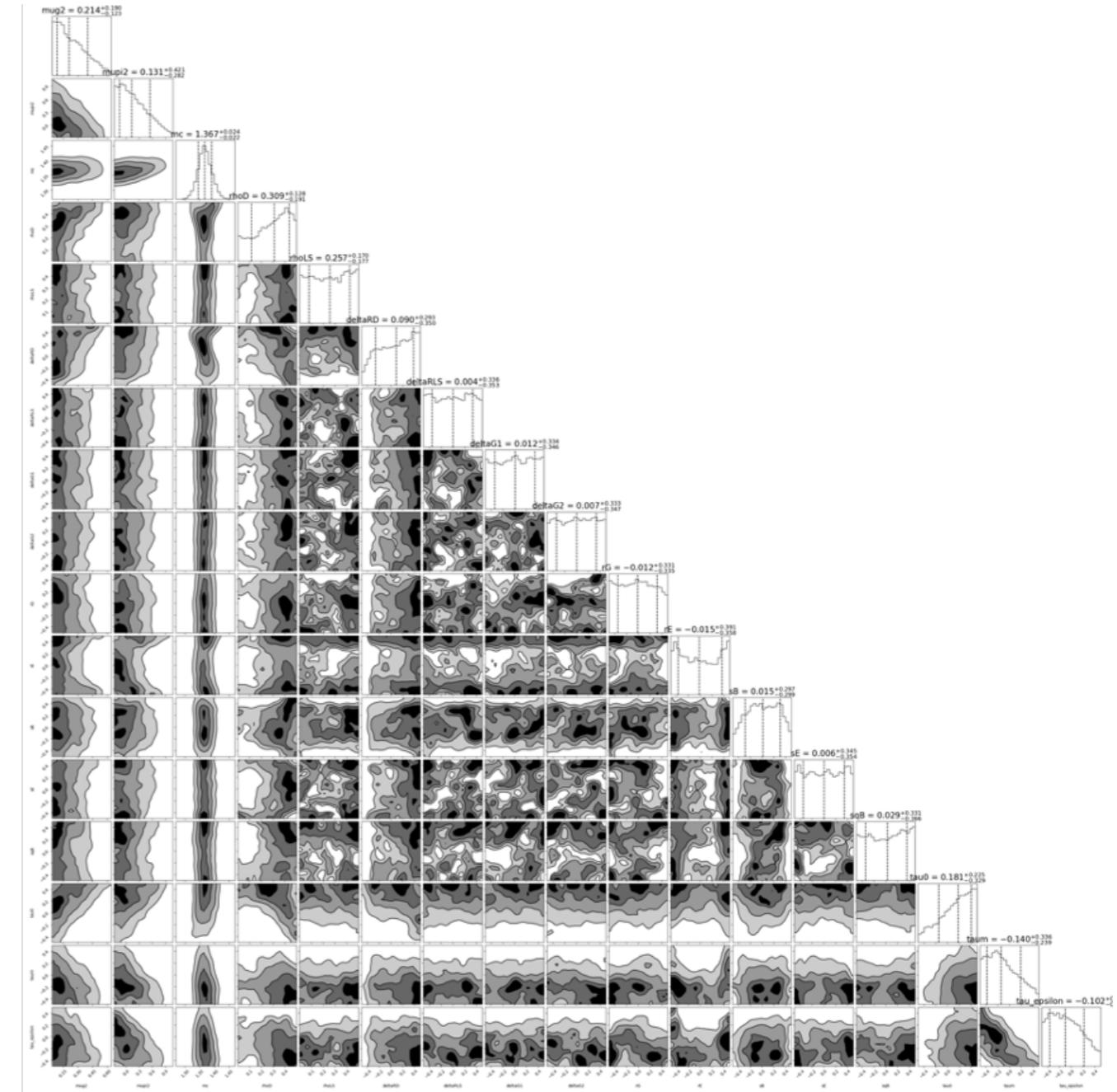
[Huang et al, '2024]



Appendix



[M5]



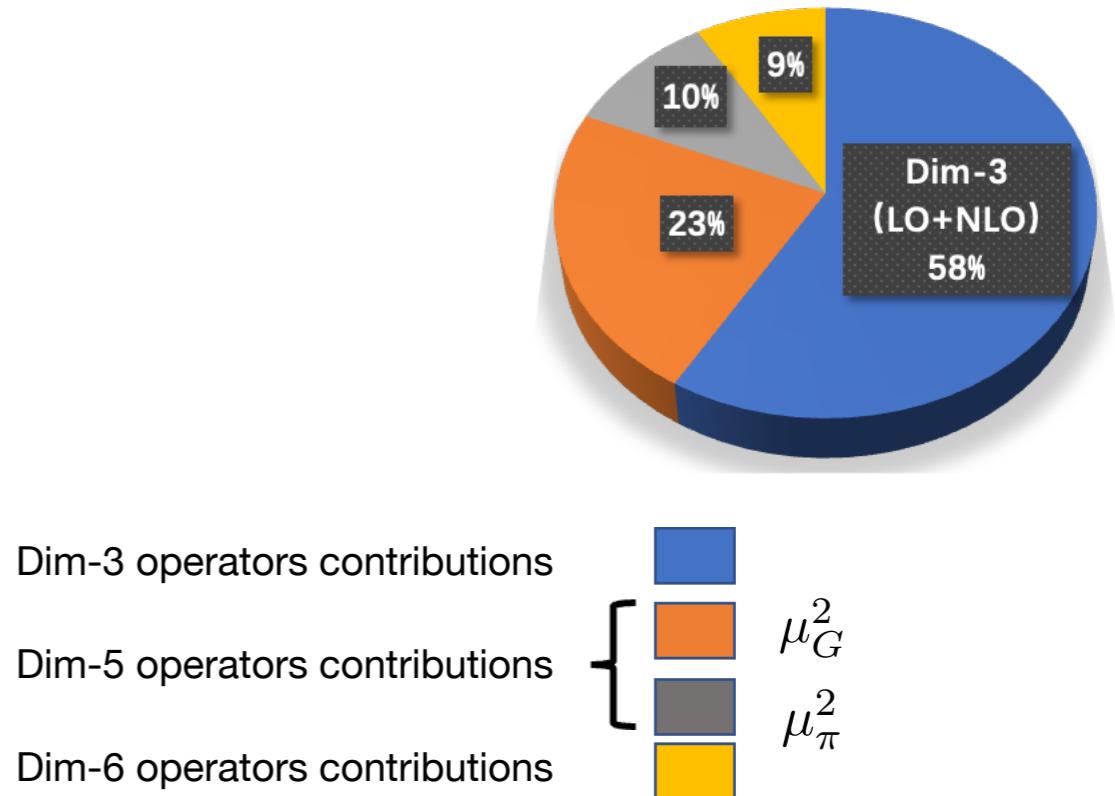
[M6]

Appendix

In the framework of the HQET,

SL:

$$\begin{aligned}
 \Gamma_{sl}^{D^+} &= \Gamma_0 \left[\underbrace{1.02}_{c_3^{\text{LO}}} + \underbrace{0.16}_{\Delta c_3^{\text{NLO}}} - 0.27 \frac{\mu_\pi^2(D)}{\text{GeV}^2} - 0.84 \frac{\mu_G^2(D)}{\text{GeV}^2} + 2.48 \frac{\rho_D^3(D)}{\text{GeV}^3} + \underbrace{0.00}_{\text{dim-7,VIA}} \right. \\
 &\quad \left. - 0.28 \tilde{B}_1^q + 0.28 \tilde{B}_2^q - 0.09 \tilde{\epsilon}_1^q + 0.09 \tilde{\epsilon}_2^q - 5.24 \tilde{\delta}_1^{sq} + 5.24 \tilde{\delta}_2^{sq} \right] \\
 &= 1.02 \Gamma_0 \left[1 + 0.16 - 0.13 \frac{\mu_\pi^2(D)}{0.48 \text{GeV}^2} - 0.28 \frac{\mu_G^2(D)}{0.34 \text{GeV}^2} + 0.20 \frac{\rho_D^3(D)}{0.082 \text{GeV}^3} \right. \\
 &\quad \left. - \underbrace{0.00}_{\text{dim-6,7,VIA}} - 0.005 \frac{\delta \tilde{B}_1^q}{0.02} + 0.005 \frac{\delta \tilde{B}_2^q}{0.02} + 0.004 \frac{\tilde{\epsilon}_1^q}{-0.04} - 0.004 \frac{\tilde{\epsilon}_2^q}{-0.04} \right. \\
 &\quad \left. - 0.0118 r_1^{sq} - 0.0088 r_2^{sq} \right], \tag{4.13}
 \end{aligned}$$

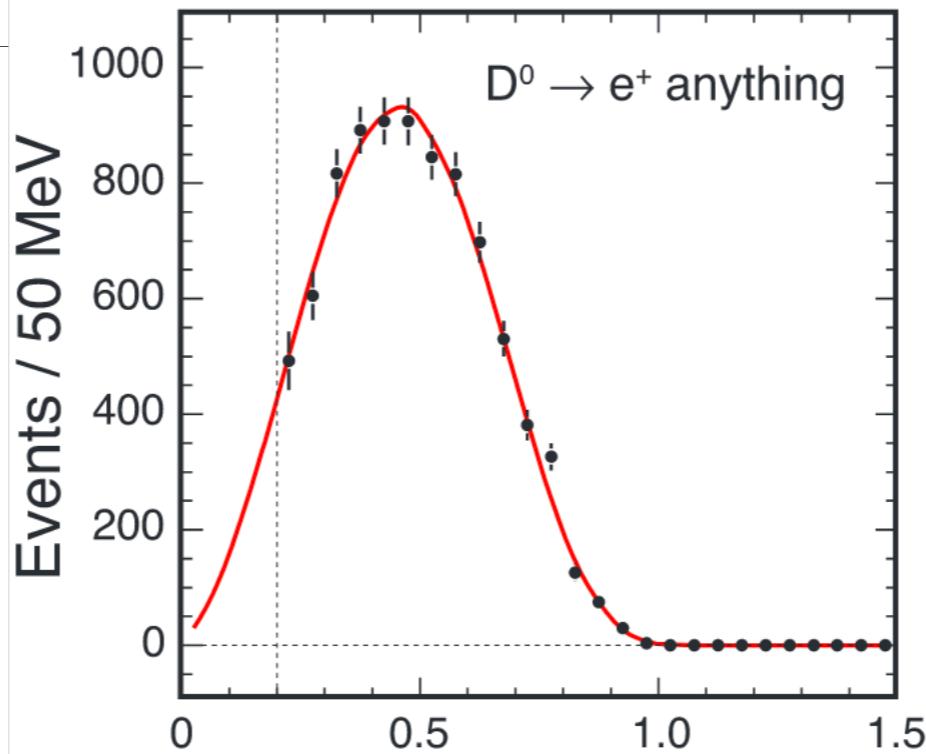


SL+NL:

$$\begin{aligned}
 \Gamma(D^+) &= \Gamma_0 \left[\underbrace{6.15}_{c_3^{\text{LO}}} + \underbrace{2.95}_{\Delta c_3^{\text{NLO}}} - 1.66 \frac{\mu_\pi^2(D)}{\text{GeV}^2} + 0.13 \frac{\mu_G^2(D)}{\text{GeV}^2} + 23.6 \frac{\rho_D^3(D)}{\text{GeV}^3} \right. \\
 &\quad \left. - 16.9 \tilde{B}_1^q + 0.56 \tilde{B}_2^q + 84.0 \tilde{\epsilon}_1^q - 1.34 \tilde{\epsilon}_2^q + \underbrace{6.76}_{\text{dim-7}} \right. \\
 &\quad \left. - 0.06 \tilde{\delta}_1^{qq} + 0.06 \tilde{\delta}_2^{qq} - 16.8 \tilde{\delta}_3^{qq} + 16.9 \tilde{\delta}_4^{qq} - 29.3 \tilde{\delta}_1^{sq} + 28.8 \tilde{\delta}_2^{sq} + 0.56 \tilde{\delta}_3^{sq} + 2.36 \tilde{\delta}_4^{sq} \right] \\
 &= 6.15 \Gamma_0 \left[1 + 0.48 - 0.13 \frac{\mu_\pi^2(D)}{0.48 \text{GeV}^2} + 0.01 \frac{\mu_G^2(D)}{0.34 \text{GeV}^2} + 0.31 \frac{\rho_D^3(D)}{0.082 \text{GeV}^3} \right. \\
 &\quad \left. - \underbrace{2.66}_{\text{dim-6,VIA}} - 0.055 \frac{\delta \tilde{B}_1^q}{0.02} + 0.002 \frac{\delta \tilde{B}_2^q}{0.02} - 0.546 \frac{\tilde{\epsilon}_1^q}{-0.04} + 0.009 \frac{\tilde{\epsilon}_2^q}{-0.04} + \underbrace{1.10}_{\text{dim-7,VIA}} \right. \\
 &\quad \left. - 0.0000 r_1^{qq} - 0.0000 r_2^{qq} + 0.0011 r_3^{qq} + 0.0008 r_4^{qq} \right. \\
 &\quad \left. - 0.0109 r_1^{sq} - 0.0080 r_2^{sq} - 0.0000 r_3^{sq} + 0.0001 r_4^{sq} \right], \tag{4.6}
 \end{aligned}$$

Discussion about Energy Spectrum

$p(\text{GeV})$	$\Delta B(D^0 \rightarrow X e^+ \nu_e)(\%)$
0.200 – 0.250	0.347 ± 0.036
0.250 – 0.300	0.426 ± 0.030
0.300 – 0.350	0.576 ± 0.031
0.350 – 0.400	0.629 ± 0.030
0.400 – 0.450	0.640 ± 0.031
0.450 – 0.500	0.640 ± 0.031
0.500 – 0.550	0.596 ± 0.029
0.550 – 0.600	0.575 ± 0.029
0.600 – 0.650	0.492 ± 0.026
0.650 – 0.700	0.374 ± 0.023
0.700 – 0.750	0.269 ± 0.019
0.750 – 0.800	0.230 ± 0.017
0.800 – 0.850	0.089 ± 0.011
0.850 – 0.900	0.053 ± 0.008
0.900 – 0.950	0.021 ± 0.005
0.950 – 1.000	0.002 ± 0.002
1.000 – 1.050	...



$$\mathcal{B}(D^0 \rightarrow X e^+ \nu_e) = (6.46 \pm 0.09 \pm 0.11)\%$$

Channel	$\mathcal{B}(\%)$	Form factor	Comment
$D^0 \rightarrow K^{*-} e^+ \nu_e$	2.16(17)[1]	SPOLE	$r_V = 1.62(8)$ and $r_2 = 0.83(5)$ [17]
$D^0 \rightarrow K^- e^+ \nu_e$	3.50(5)[5]	BK	$\alpha_{\text{BK}} = 0.30(3)$ [5]
$D^0 \rightarrow K_1^- e^+ \nu_e$	0.11(11)	ISGW2	\mathcal{B} from Ref. [10] scaled by Ref. [5]
$D^0 \rightarrow K_2^{*-} e^+ \nu_e$	0.11(11)	ISGW2	\mathcal{B} set to same as $D^0 \rightarrow K_1^- e^+ \nu_e$
$D^0 \rightarrow \bar{K} \pi e^+ \nu_e$	0.12(3)[17, 29]	PHSP	Nonresonant
$D^0 \rightarrow \pi^- e^+ \nu_e$	0.288(9)[5]	BK	$\alpha_{\text{BK}} = 0.21(7)$ [5]
$D^0 \rightarrow \rho^- e^+ \nu_e$	0.16(2)[2]	SPOLE	$r_V = 1.4(3)$ and $r_2 = 0.6(2)$ [2]

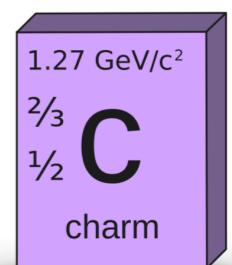


School of Nuclear Science and Technology

Kang-kang Shao

Collaborators: Chun Huang (WashU) , Dong-hao Li (LZU)

Supervisor: Fu-sheng Yu, Qin Qin



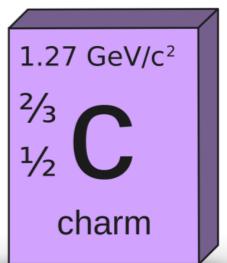
HFCPV2024
Heng Yang, China



Kang-kang Shao

Collaborators: Chun Huang (WashU) , Dong-hao Li (LZU)

Supervisor: Fu-sheng Yu, Qin Qin



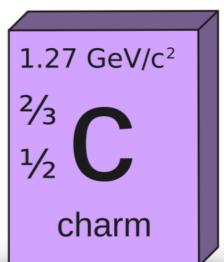
**HFCPV2024
Heng Yang, China**



Kang-kang Shao

Collaborators: Chun Huang (WashU) , Dong-hao Li (LZU)

Supervisor: Fu-sheng Yu, Qin Qin



**HFCPV2024
Heng Yang, China**