

#### CENTRAL CHINA NORMAL UNIVERSITY

## Study of the $B^- \rightarrow K^- \eta \eta_c$ decay due to the $D\overline{D}$ bound state

Le-Le Wei

Central China Normal University

llwei@mails.ccnu.edu.cn

第二届武汉高能物理青年论坛

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Formalism

## Results & Discussion





### **Exotic state**



- Hadron: strong interaction, composites of quarks and gluons
- Traditional quark model:

Hadron  $\begin{cases} \text{Meson } (q\bar{q}) \\ \text{Baryon } (qqq, \bar{q}\bar{q}\bar{q}) \\ \text{Phys. Lett. 8, 214 (1964)} \\ \text{CERN-TH-401, CERN-TH-412} \end{cases}$ 

Quantum Chromodynamics:

. . .

ExoticMolecular  $(M\overline{M}, B\overline{B}...)$ ExoticTetraquark  $(qq\overline{q}\overline{q}...)$ Pentaquark  $(qqqq\overline{q}...)$ StateHybrid  $(q\overline{q}g, qqqg...)$ Glueball (gg, ggg...)



#### **Exotic state**



- 2003, Belle Collaboration, X(3872) in the J/ψπ<sup>+</sup>π<sup>-</sup> invariant mass distribution
   [Belle] Phys. Rev. Lett. 91, 262001 (2003)
   LHCb, Belle, BABAR, BESIII... many exotic states have been observed experimentally
- ➢ Exotic states near the thresholds of a pair of hadrons:
  X(3872) ↔ DD̄\* [Belle] Phys. Rev. Lett. 91, 262001 (2003)
  Z<sub>c</sub>(3900) ↔ DD̄\* [BESIII] Phys. Rev. Lett. 110, 252001 (2013)
  Z<sub>cs</sub>(3985) ↔ DD̄<sub>s</sub>\* [BESIII] Phys. Rev. Lett. 126, 102001 (2021)
  P<sub>cs</sub>(4459) ↔ Ξ<sub>c</sub>D̄\* [LHCb] Sci. Bull. 66, 1278 (2021)
  T<sub>cc</sub><sup>+</sup> ↔ DD\* [LHCb] Nature Phys. 18, 751 (2022) ...
  Explanations: compact tetraquark, pentaquark, or hexaquark, hadronic molecule, kinematic effect (threshold cusp, triangle singularity) ...

#### Hadronic molecule:

Two or more hadrons bounded through the strong interactions Mass around the threshold of a pair of hadrons Quantum number  $J^{PC} \leftrightarrow S$ -wave of the composites Rev. Mod. Phys. 90, 015003 (2018)

## $D\overline{D}$ bound state



- > 2007,  $\pi^+\pi^-$ ,  $\pi^0\pi^0$ ,  $K^+K^-$ ,  $K^0\overline{K}^0$ ,  $D^+D^-$ ,  $D^0\overline{D}^0$ ,  $D_s^+D_s^-$ ,  $\eta\eta$ ,  $\eta\eta_c$  coupled channels, unitary coupled-channel approach, Phys. Rev. D 76, 074016 (2007)  $D\overline{D}$  bound state with  $I(J^{PC}) = 0(0^{++})$ ,  $M \approx 3720$  MeV, denote as X(3700)
- ➤ Unitary coupled-channel approach:  $\sqrt{s} = (3722 i18)$  MeV Eur. Phys. J. A 41, 85 (2009) Lattice QCD: binding nergy  $B = (4.0^{+5.0}_{-3.7})$  MeV J. High Energy Phys. 06, 035 (2021) QCD sum rule:  $M_{D\overline{D}} = (3.73 \pm 0.08)$  GeV Phys. Rev. D 105, 094003 (2022) Nonrelativistic EFT:  $M = (3739.3 \pm 0.1)$  MeV Phys. Rev. D 105, 034024 (2022)
- Some theoretical studies of the experimental measured process support the existence of such a DD̄ bound state, e.g., e<sup>+</sup>e<sup>-</sup> → J/ψDD̄<sup>(\*)</sup>, B → DD̄K, γγ → DD̄.
   Eur. Phys. J. A 36, 189 (2008), Eur. Phys. J. C 76, 121 (2016), Phys. Lett. B 827, 136982 (2022), Eur. Phys. J. A 57, 38 (2021), Phys. Rev. D 103, 054008 (2021)
- Some process suggested to search for the  $D\overline{D}$  bound state:  $\psi(3770)/\psi(4040) \rightarrow \gamma X(3700) \rightarrow \gamma \eta \eta', e^+e^- \rightarrow J/\psi X(3700) \rightarrow J/\psi \eta \eta',$   $\psi(3770) \rightarrow \gamma D\overline{D}, \Lambda_b \rightarrow \Lambda D\overline{D}, B^+ \rightarrow K^+\eta \eta$ Eur. Phys. J. A 41, 85 (2009), Eur. Phys. J. A 49, 52 (2013), Eur. Phys. J. C 80, 510 (2020), Phys. Rev. D 103, 114013 (2021), Phys. Rev. D 108, 054004 (2023)

## $\Lambda_b\to\Lambda D\overline{D}$









The intermediate state  $\psi(3770)$ :

$$t^{P-\text{wave}} = \frac{\beta \times V_p \times M_{\psi(3770)} \tilde{p}_D}{M_{D\bar{D}}^2 - M_{\psi(3770)}^2 + iM_{\psi(3770)} \tilde{\Gamma}_{\psi(3770)}}$$

$$S^{-\text{wave}}_{\Lambda_{b}\to\Lambda D^{0}\overline{D}^{0}} = V_{p}(1 + G_{D^{+}D^{-}t_{D^{+}D^{-}\to D^{0}\overline{D}^{0}} + G_{D^{0}\overline{D}^{0}}t_{D^{+}D^{-}\to D^{0}\overline{D}^{0}} + (1 + C)G_{D_{s}^{+}D_{s}^{-}}t_{D_{s}^{+}D_{s}^{-}\to D^{0}\overline{D}^{0}})$$

$$S^{-\text{wave}}_{\Lambda_{b}\to\Lambda D^{+}D^{-}} = V_{p}(1 + G_{D^{+}D^{-}t_{D^{+}D^{-}\to D^{+}D^{-}} + G_{D^{0}\overline{D}^{0}}t_{D^{0}\overline{D}^{0}\to D^{+}D^{-}} + (1 + C)G_{D_{s}^{+}D_{s}^{-}}t_{D_{s}^{+}D_{s}^{-}\to D^{+}D^{-}} + (1 + C)G_{D_{s}^{+}D_{s}^{-}}t_{D_{s}^{+}D_{s}^{-}\to D^{+}D^{-}})$$
BS equation:  $t_{i+i+i} = [1 - VG]^{-1}V$ 

Phys. Rev. D 103, 114013 (2021)

The  $D^0\overline{D}^0$  (left) and  $D^+D^-$  (right) invariant mass distributions: J. High Energy Phys. 07,140 (2024)

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- An enhancement near  $D\overline{D}$  threshold
- Partial wave analysis

Sizeable coupling between X(3700) and  $\eta \eta_c$ ,  $M_{X(3700)} - (M_{\eta} + M_{\eta_c}) \approx 200 \text{ MeV}$  $\rightarrow \text{A clear peak in the } \eta \eta_c \text{ invariant mass distribution} \rightarrow B^- \rightarrow K^- \eta \eta_c \text{ decay}$ 





Formalism

## Results & Discussion



Bethe-Salpeter equation:



Scattering amplitude:  $T = V + VGT \rightarrow T = [1 - VG]^{-1}V$ 

*V*: transition potential between the coupled channels

G: loop function for the two-meson propagation in the lth coupled channel

$$G_{l}(s) = \frac{1}{16\pi^{2}} \left[ \alpha_{l} + \ln \frac{m_{1}^{2}}{\mu^{2}} + \frac{m_{2}^{2} - m_{1}^{2} + s}{2s} \ln \frac{m_{2}^{2}}{m_{1}^{2}} + \frac{m_{2}^{2} - m_{1}^{2} + s}{\sqrt{s}} \ln \frac{m_{2}^{2}}{m_{1}^{2}} + \frac{p_{1}}{\sqrt{s}} \times \left( \ln \frac{s - m_{2}^{2} + m_{1}^{2} + 2p\sqrt{s}}{-s + m_{2}^{2} - m_{1}^{2} + 2p\sqrt{s}} + \ln \frac{s + m_{2}^{2} - m_{1}^{2} + 2p\sqrt{s}}{-s - m_{2}^{2} + m_{1}^{2} + 2p\sqrt{s}} \right) \right]$$
Prog. Part. Nucl. Phys. 45, 157 (2000)

➢ Pole in the complex energy plane (T) ↔ bound state of a pair of hadrons  $\sqrt{s}_{p} = (M_{R} - i\Gamma_{R}/2) \text{ MeV} \qquad M_{R} ↔ \text{mass}; \ \Gamma_{R} ↔ \text{width}$ 

## $B^- \rightarrow K^- \eta \eta_c$





The intermediate state  $K^*(1430)$ :

$$T_{K^*} = \frac{V_p \times \beta \times M_{K_0^*(1430)}^2}{M_{K^-\eta}^2 - M_{K_0^*(1430)}^2 + iM_{K_0^*(1430)}\Gamma_{K_0^*(1430)}}$$



 $K_0^*(1430)$ 

Doubly differential decay width:  $\frac{\mathrm{d}^{2}\Gamma}{\mathrm{d}M_{nn_{c}}\mathrm{d}M_{K^{-}n}} = \frac{1}{(2\pi)^{3}} \frac{M_{K^{-}\eta}M_{\eta\eta_{c}}}{8M_{B^{-}}^{3}} \left|T_{X(3700)} + e^{i\varphi}T_{K^{*}}\right|^{2}$ Differential decay width:  $\frac{d\Gamma}{dM_{12}} = \int \frac{d^2\Gamma}{dM_{12}dM_{22}} dM_{23}$ 







#### The $\eta \eta_c$ (a) and $K^- \eta$ (b) invariant mass distributions (fixed parameter):



 $\blacktriangleright$  a clear peak around 3720 MeV in the  $\eta \eta_c$  IMD associated with the  $D\overline{D}$  bound state

The  $\eta \eta_c$  (a) and  $K^- \eta$  (b) invariant mass distributions with  $\beta = 0.009, 0.012, 0.015, 0.018$ :



## **Results and discussion**



The  $\eta \eta_c$  (a) and  $K^- \eta$  (b) invariant mass distributions with  $\varphi = 0, \frac{\pi}{3}, \frac{2\pi}{3}, \pi, \frac{4\pi}{3}, \frac{5\pi}{3}$ .



The  $\eta \eta_c$  (a) and  $K^- \eta$  (b) invariant mass distributions with C = 3.0, 2.5, 2.0:



> Within the variation ranges of parameters, always a clear peak in the  $\eta \eta_c$  IMD





Formalism

## Results & Discussion

## Summary



- Sizeable coupling between X(3700) and  $\eta\eta_c$ ,  $M_{X(3700)} (M_\eta + M_{\eta_c}) \approx 200 \text{ MeV}$ Many exotic states discovered in the *B*-meson three-body decay  $\rightarrow$  Search for the  $D\overline{D}$  bound state in the  $B^- \rightarrow K^- \eta\eta_c$  process
- ➢ Unitary coupled-channel approach, calculate the  $\eta\eta_c$  invariant mass distributions X(3700): S-wave pseudoscalar meson-pseudoscalar meson interactions Intermediate resonance K<sup>\*</sup><sub>0</sub>(1430)
- >  $\eta \eta_c$  invariant mass distribution: A clear peak appears around 3720 MeV associated with the  $D\overline{D}$  bound state X(3700) Intermediate resonance  $K_0^*(1430)$  gives a smooth contribution
- A more precise measurement of the  $B^- \to K^- \eta \eta_c$  decay at the Belle II and LHCb experiments in the future, to confirm the existence of such a predicted  $D\overline{D}$  bound state and measure its mass and width

# Thank you for your attention!



Doubly differential decay width of the  $B^- \to K^- \eta \eta_c$  decay in the  $(M^2_{\eta \eta_c}, M^2_{K^- \eta})$  plane



> The X(3700) and  $K_0^*(1430)$  resonances can be clearly seen



 $B^{\pm} \to K^{\pm} \eta_c \eta$  Belle Collaboration: 772 × 10<sup>6</sup>  $B\bar{B}$  pairs collected at  $\Upsilon(4S)$ 

No obvious X(3730) signal, the upper limits of the branching ratios:

 $\begin{aligned} \mathcal{B}(B^{\pm} \to K^{\pm} \eta_c \eta) &< 2.2 \times 10^{-4} \end{aligned} \qquad \text{[Belle] J. High Energy Phys. 06, 132 (2015)} \\ \mathcal{B}(B^{\pm} \to K^{\pm} X(3730)) \times \mathcal{B}(X(3730) \to \eta_c \eta) &< 4.6 \times 10^{-5} \end{aligned}$ 





#### Other quark level diagrams:



Components of the meson systems after hadronization:

$$\begin{split} |H\rangle^{3a} &= V_p V_{ub} V_{us} \left( \frac{1}{3} \eta \eta + \frac{1}{2} \pi^0 \pi^0 + \frac{1}{6} \eta' \eta' + \frac{1}{\sqrt{18}} \eta \eta' + \pi^+ \pi^- + K^+ K^- + D^0 \overline{D}^0 \right) K^- \\ |H\rangle^{3b} &= C \times V_p V_{ub} V_{us} \left( \frac{1}{3} \eta \eta + \frac{1}{2} \pi^0 \pi^0 + \frac{1}{6} \eta' \eta' + \frac{1}{\sqrt{18}} \eta \eta' + \pi^+ \pi^- + K^+ K^- + D^0 \overline{D}^0 \right) K^- \\ |V_{ub}| &= 0.00369 \pm 0.00011 \qquad |V_{us}| = 0.22500 \pm 0.00067 \\ |V_{cb}| &= 0.04182^{+0.00085}_{-0.00074} \qquad |V_{cs}| = 0.97349 \pm 0.00016 \end{split}$$

Prog. Theor. Exp. Phys. 2022, 083C01 (2022)



#### Hadronization:

L. R Dai, et al. Eur. Phys. J. C 76, 121 (2016)



 $q\bar{q}$  matrix expressed in terms of the physical pseudoscalar mesons:

$$M \Rightarrow \begin{pmatrix} \frac{\eta}{\sqrt{3}} + \frac{\pi^{0}}{\sqrt{2}} + \frac{\eta'}{\sqrt{6}} & \pi^{+} & K^{+} & \overline{D}^{0} \\ \\ \pi^{-} & \frac{\eta}{\sqrt{3}} - \frac{\pi^{0}}{\sqrt{2}} + \frac{\eta'}{\sqrt{6}} & K^{0} & D^{-} \\ \\ \\ K^{-} & \overline{K}^{0} & \frac{\sqrt{2}}{\sqrt{3}}\eta' - \frac{\eta}{\sqrt{3}} & D_{s}^{-} \\ \\ D^{0} & D^{+} & D_{s}^{+} & \eta_{c} \end{pmatrix}$$