









Tensions in Cosmology: Are we Approaching New Physics?

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IHEP, Beijing July 2025



- The history of Science is a history of tensions between theoretical predictions and observations
 - Astrophysics and Cosmology have become precision sciences with an incredibly huge amount of data
 - New Tensions appear.
 Are we approaching New Physics?

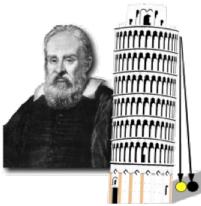
History of Science - History of Tensions



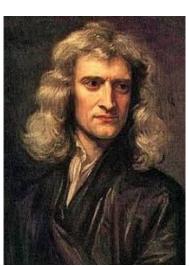








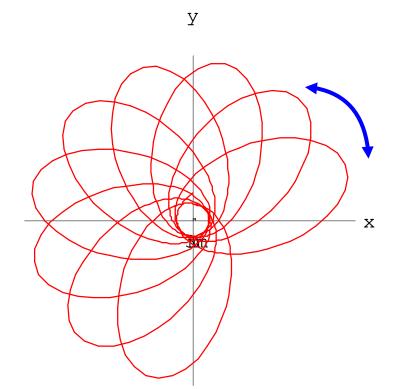






Mercury periliheimum - 1859

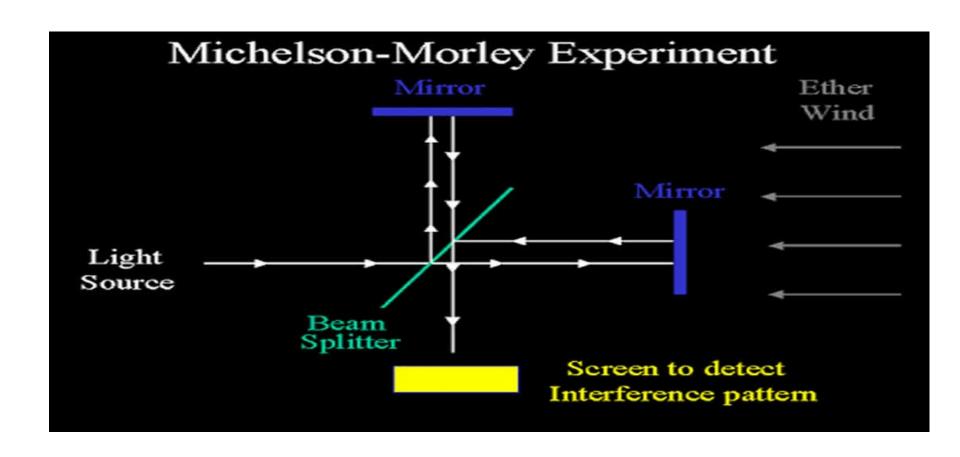
• The true orbits of planets, even if seen from the SUN are not ellipses. They are rather curves of this type:



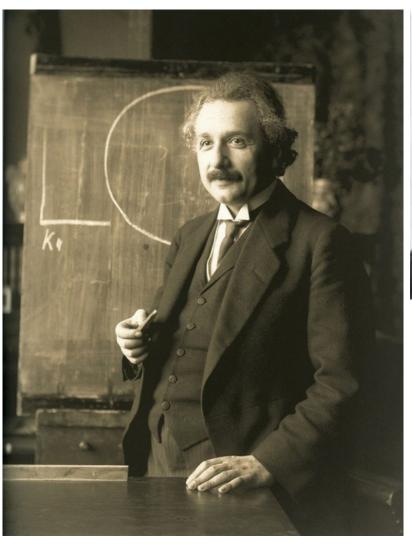
For the planet Mercury it is

 $\Delta \varphi = 43$ " of arc per century

Michelson-Morley experiment - 1887



The Theory of Relativity









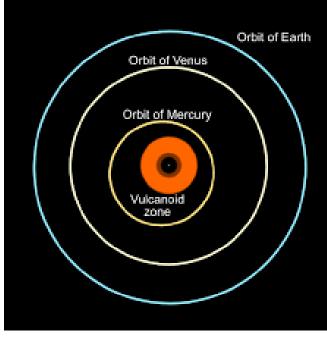
$$S = \frac{1}{16 \pi G} \int d^4 x \sqrt{-g} [R - 2\Lambda] + \int d^4 x L_m (g_{\mu\nu}, \psi)$$

$$\Rightarrow R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + g_{\mu\nu} \Lambda = 8 \pi G T_{\mu\nu}$$

$$T^{\mu\nu} \equiv \frac{2}{\sqrt{-g}} \frac{\delta L_m}{\delta g_{\mu\nu}}$$

60 years of tension





Modified Gravity before General Relativity

- Modifications to Newton's Law
- Inverse Cube Law.
- Extended Inverse-Square Law (Simon Newcomb -1880's)
- Lord Kelvin theory of everything (end of 19th century)
- Hendrik Lorentz: gravity on the basis of his ether theory and Maxwell's equations. (1900)
- Nordström's theory of gravitation (1912 and 1913)
- Einstein's scalar theory of gravity (1913)

History of Major Shifts of Cosmological Models

Aristarchus of Samos (c. 310-230 BCE)

- ► Geocentric (Ptolemy) to Heliocentric (Copernicus, 16th-17th c.):
 - Retrograde motion (Copernicus)
 - ▶ Phases of Venus, Moons of Jupiter (Galileo, 1610)
- ► Heliocentric to Infinite Universe (18th-19th c.):
 - ▶ Improved telescopes, Uranus discovery (Herschel, 1781)
 - Stellar parallax (Bessel, 1838)
- ▶ Infinite to Static Universe (Einstein, early 20th c.):
 - Nebulae spectroscopy (Slipher)
 - Stellar distances (Leavitt, Hertzsprung)
- Static to Expanding Universe (Lemaître, Hubble, 1920s-30s):
 - ► Galactic redshift (Slipher, 1912-14)
 - Hubble's law (Hubble, 1929)
- Expanding Universe to Inflationary Big Bang (1960s-80s):
 - CMB (Penzias & Wilson, 1964)
 - Light element abundance (Alpher, Herman)
 - Inflation theory (Guth, Linde, 1980s)

- ▶ Introduction of Dark Matter (1970s-80s):
 - ► Galaxy rotation curves (Rubin, 1970s)
 - ► Gravitational lensing (Walsh et al., 1979)
 - Galaxy clusters (Zwicky, 1930s)
- ► Lambda-CDM (late 1990s-present):
 - Supernova observations (Perlmutter, Schmidt, Riess, 1998-99)
 - ► CMB (WMAP, Planck), BAO (SDSS, 2005)
- ► Potential Future Shift (2020s-?):
 - ► Hubble tension (Riess et al. vs Planck Collaboration)
 - ▶ S8 (growth rate) tension (KiDS, DES, Planck collaborations)
 - Cosmic dipoles tension (Various teams)
 - CMB anomalies (Planck Collaboration)
 - ► ISW (Integrated Sachs-Wolfe) tension
 - Lithium problem (Primordial Nucleosynthesis)



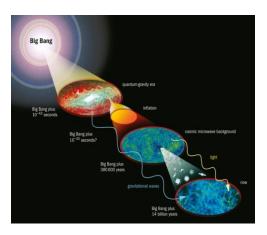
New Astronomy Reviews

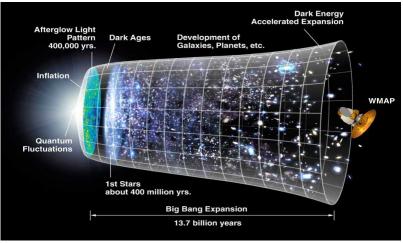
Volume 95, December 2022, 101659

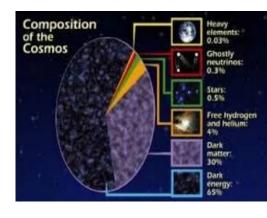


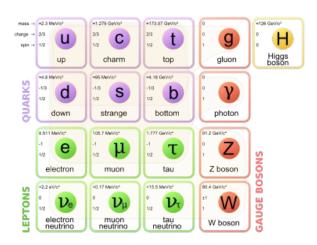
Challenges for Λ CDM: An update

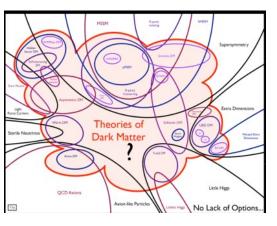
Summary of 20th century Observations

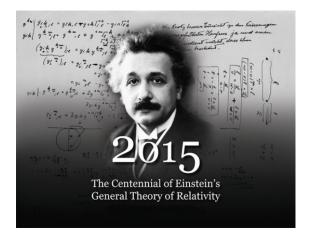








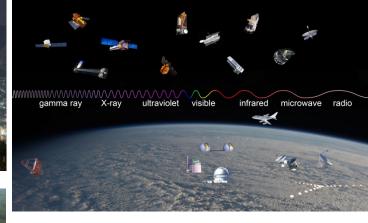


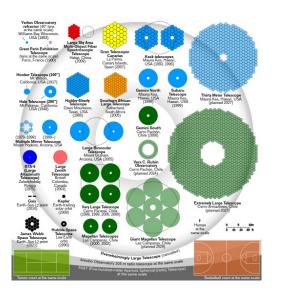


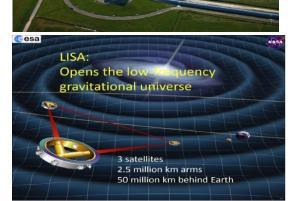
Astrophysics and Cosmology in the 21st century

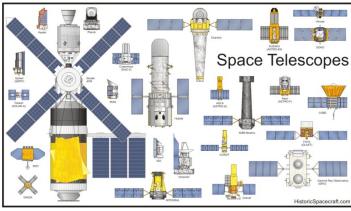








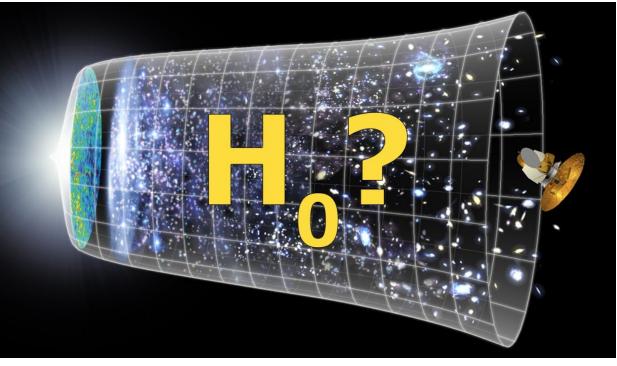


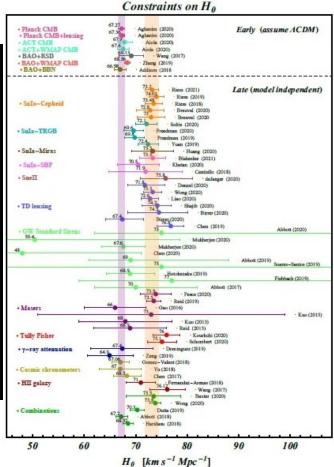


So what do these new observations tell us?

H0 tension

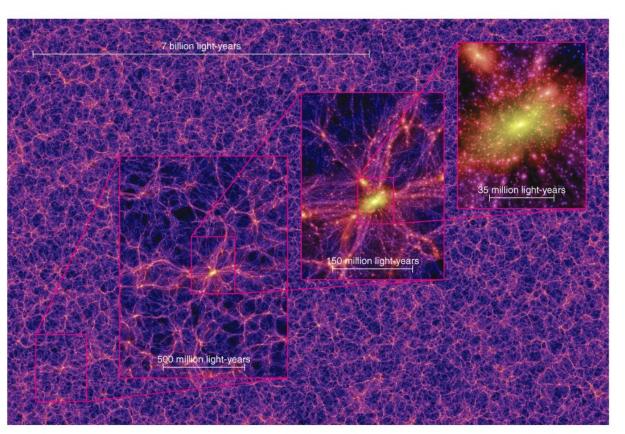
The Universe expands faster than expected!



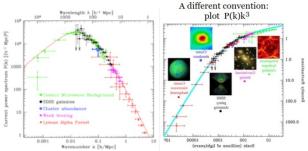


S8 Tension

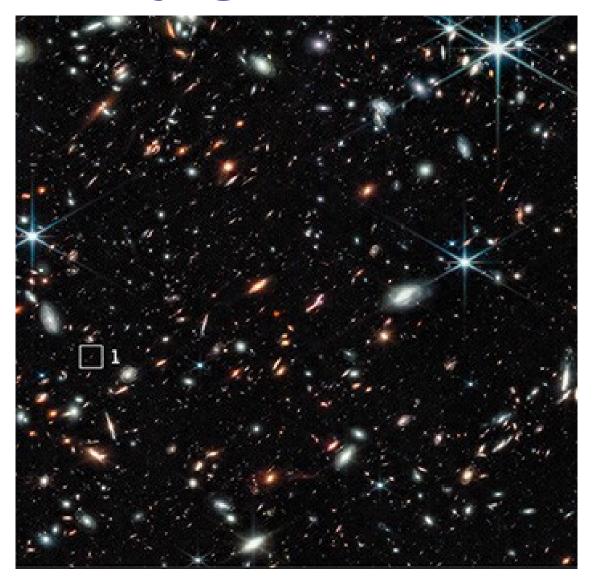
"Less" Matter clustering than expected!



Matter Density Fluctuation Power Spectrum

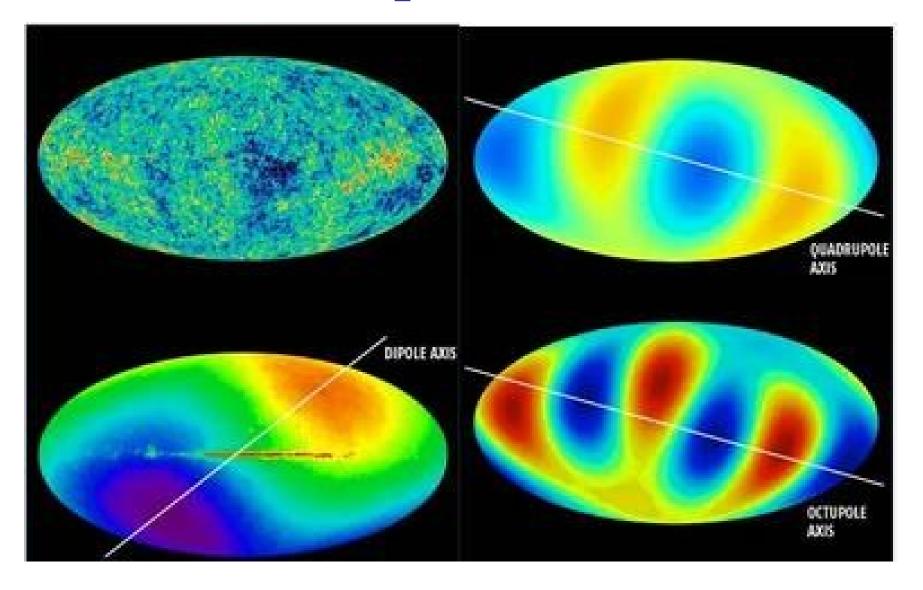


Too many galaxies too early!

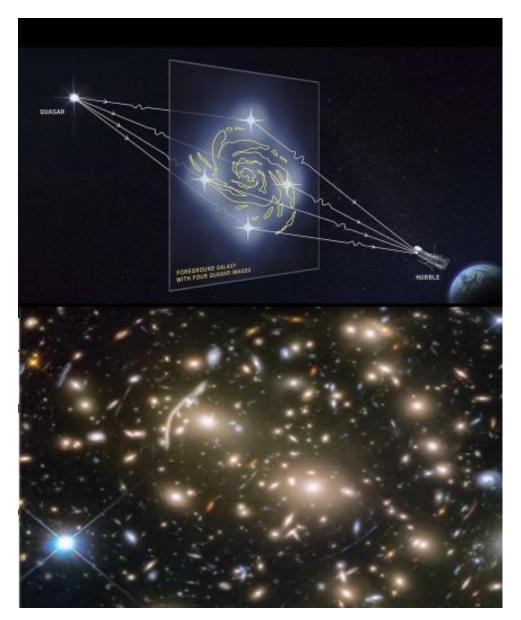


James Webb space telescope

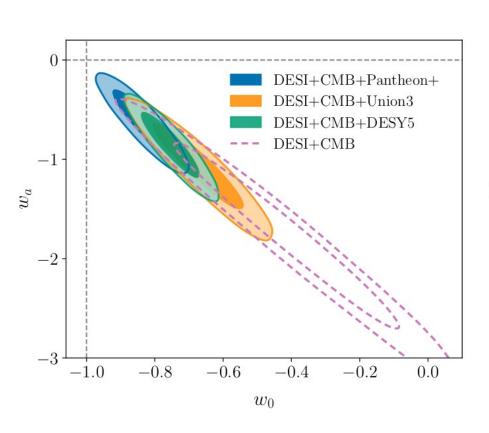
Cosmic dipole tension!

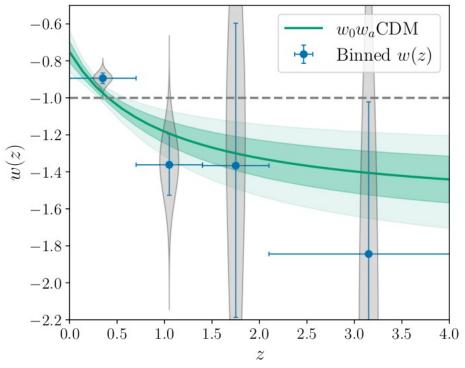


The lensing anomaly



DESI2 DATA and QUINTOM BEHAVIOR





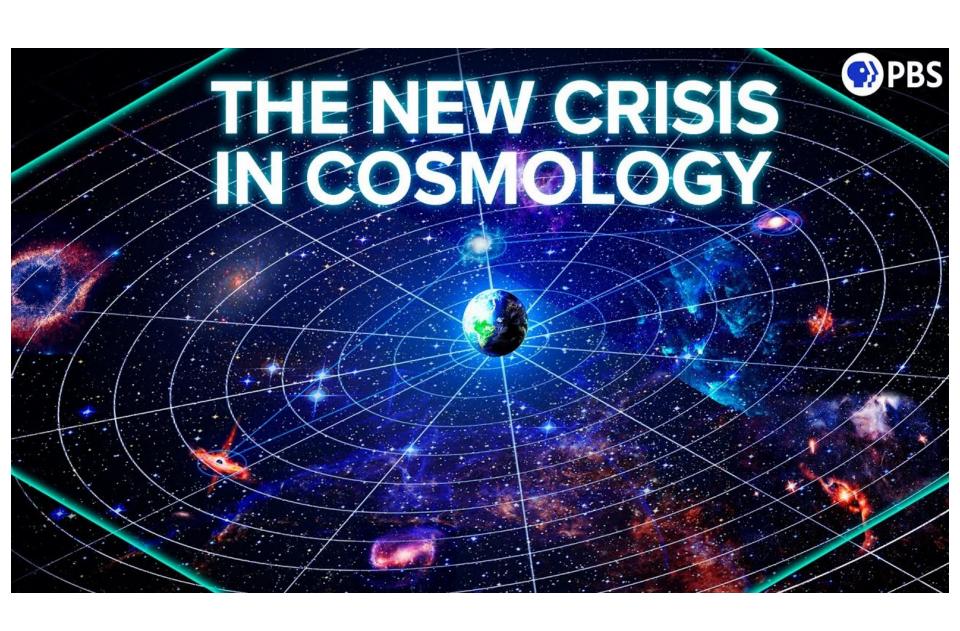
Submitted to the Proceedings of the US Community Study on the Future of Particle Physics (Snowmass 2021)

Cosmology Intertwined: A Review of the Particle Physics, Astrophysics, and Cosmology Associated with the Cosmological Tensions and Anomalies

Elcio Abdalla, Guillermo Franco Abellán, Amin Aboubrahim, Adriano Agnello, Özgür Akarsu, Yashar Akrami ^{6,7,8,9} George Alestas, ¹⁰ Daniel Aloni, ¹¹ Luca Amendola, ¹² Luis A. Anchordoqui, ^{13,14,15} Richard I. Anderson, ¹⁶ Nikki Arendse, ¹⁷ Marika Asgari, ^{18,19} Mario Ballardini, ^{20,21,22,23} Vernon Barger, ²⁴ Spyros Basilakos, ^{25,26} Ronaldo C. Batista, ²⁷ Elia S. Battistelli, ^{28,29} Richard Battye, ³⁰ Micol Benetti, ^{31,32} David Benisty, ³³, ³⁴, ³⁵ Asher Berlin, ³⁶ Paolo de Bernardis, ²⁸, ²⁹ Emanuele Berti, ³⁷ Bohdan Bidenko, ³⁸, ³⁹ Simon Birrer, ⁴⁰ John P. Blakeslee, 41 Kimberly K. Boddy, 42 Clecio R. Bom, 43, 44 Alexander Bonilla, 45 Nicola Borghi, 46, 47 Francois R. Bouchet, 48 Matteo Braglia, 49, 21 Thomas Buchert, 50 Elizabeth Buckley-Geer, 51, 52 Erminia Calabrese, 53 Robert R. Caldwell, ⁵⁴ David Camarena, ⁵⁵ Salvatore Capozziello, ^{56,31} Stefano Casertano, ⁵⁷ Angela Chen, ^{58,59} Geoff C.-F. Chen, 60 Hsin-Yu Chen, 61 Jens Chluba, 30 Anton Chudaykin, 62 Michele Cicoli, 20, 22 Craig J. Copi, 6 Fred Courbin, 16 Francis-Yan Cyr-Racine 63 Bozena Czerny 64 Maria Dainotti 65, 66, 67 Guido D'Amico 68, 69 Anne-Christine Davis, 33, 34 Javier de Cruz Pérez, 70 Jaume de Haro, 71 Jacques Delabrouille, 72, 73, 74, 75 Peter B. Denton, 76 Suhail Dhawan, 77 Keith R. Dienes, 78, 79 Eleonora Di Valentino, 80, 8 Pu Du, 81 Dominique Eckert, 82 Celia Escamilla-Rivera, 83 Agnès Ferté, 84 Fabio Finelli, 85, 22 Pablo Fosalba, 86, 87 Wendy L. Freedman, 52 Noemi Frusciante. 88 Enrique Gaztañaga. 86, 87 William Giarè. 89, 29 Elena Giusarma. 90 Adrià Gómez-Valent. 91 Will Handley, 92, 93 Ian Harrison, 94 Luke Hart, 30 Dhiraj Kumar Hazra, 95 Alan Heavens, 7 Asta Heinesen, 50 Hendrik Hildebrandt, 96 J. Colin Hill, 97,98 Natalie B. Hogg, 99 Daniel E. Holz, 52,100,101 Deanna C. Hooper, 102 Nikoo Hosseinineiad, 103 Dragan Huterer, 104, 105 Mustapha Ishak, 106 Mikhail M. Ivanov, 107 Andrew H. Jaffe, 7 In Sung Jang, 52 Karsten Jedamzik, 108 Raul Jimenez, 109, 110 Melissa Joseph, 11 Shahab Joudaki, 111, 112 Marc Kamionkowski, 37 Tanvi Karwal, 113 Lavrentios Kazantzidis, 10 Ryan E. Keeley, 114 Michael Klasen, 3 Eiichiro Komatsu, 115, 116 Léon V.E. Advinta Wat, Lavelinos Kazalitzus, 1941 E. Recety, Sindar Kassar, 1951 Chung-Chi Lee, 120 Julien Lesgourgus, 121 Jackson Levi Said, 122, 123 Tiffany R. Lewis, 124 Benjamin L'Huillier, 125 Matteo Lucca, 126 Roy Maartens, 23, 127, 128 Jackson Levi Sand, Denny Marfatia, 139 Valerio Marra, 131, 132, 133 Carlos J. A. P. Martins, 134, 135 Silvia Masi, 28, 29 Sabino Matarrese, 136, 137, 138, 139 Arindam Mazumdar, 140 Alessandro Melchiorri, 28, 29 Olga Mena, 141 Laura Mersini-Houghton, 142 James Mertens, 143 Dinko Milaković, 133, 132, 144 Yuto Minami, 145 Vivian Miranda, 146 Cristian Moreno-Pulido. 147 Michele Moresco. 46, 47 David F. Mota. 148 Emil Mottola, 63 Simone Mozzon. 149 Jessica Muir. 150 Ankan Mukherjee, ¹⁵⁴ Suvodip Mukherjee, ¹⁵⁰ Pavel Naselsky, ¹⁵² Pran Nath, ¹⁵³ Savvas Nesseris, ⁹⁹ Florian Niedermann, ¹⁵⁴ Alessio Notari, ¹⁵⁵ Rafael C. Nunes, ¹⁵⁶ Eoin Ó Colgáin, ¹⁵⁷, ¹⁵⁸ Kayla A. Owens, ⁵² Emre Özülker, ⁵ Francesco Pace, ¹⁵⁹, ¹⁶⁰ Andronikos Paliathanasis, ¹⁶¹, ¹⁶² Antonella Palmese, ¹⁶³ Supriya Pan, ¹⁶⁴ Daniela Paoletti, ⁸⁵, ²² Santiago E, Perez Bergliaffa, 165 Leandros Perivolaropoulos, 10 Dominic W, Pesce, 166, 167 Valeria Pettorino, 168 Oliver H. E. Philcox, 169, 107 Levon Pogosian, 170 Vivian Poulin, Gaspard Poulot, 80 Marco Raveri, 171 Mark J. Reid, 172 Fabrizio Renzi, 173 Adam G. Riess, 37 Vivian I. Sabla, 54 Paolo Salucci, 174, 175 Vincenzo Salzano, 176 Emmanuel N. Saridakis, 26,75,177 Bangalore S. Sathyaprakash, 178,179,94 Martin Schmaltz, 11 Nils Schöneberg, 180 Dan Scolnic, 181 Anjan A. Sen. 182, 183 Neelima Sehgal, 184 Arman Shafieloo, 185 M.M. Sheikh-Jabbari, 186 Joseph Silk, 187 Alessandra Silvestri, ¹⁷³ Foteini Skara, ¹⁰ Martin S. Sloth, ¹⁸⁸ Marcelle Soares-Santos, ⁵⁸ Joan Solà Peracaula, ¹⁴⁷ Yu-Yang Songsheng ⁸¹ Jorge F. Soriano ^{13,14} Denitsa Staicova ¹⁸⁹ Glenn D. Starkman ^{6,7} István Szapudi ¹⁹⁰ Elsa M. Teixeira, 80 Brooks Thomas, 191 Tommaso Treu, 60 Emery Trott, 58 Carsten van de Bruck, 80 J. Alberto Vazquez, 192 Licia Verde, 193, 194 Luca Visinelli, 195 Deng Wang, 196 Jian-Min Wang, 81 Shao-Jiang Wang, 197 Richard Watkins, 198 Scott Watson, 199 John K. Webb, 120 Neal Weiner, 200 Amanda Weltman, 201 Samuel J. Witte, 202 Radoslaw Wojtak, Anil Kumar Yadav, Weiqiang Yang, 204 Gong-Bo Zhao, 205, 206 and Miguel Zumalacárregui 207 ¹Instituto de Física, Universidade de São Paulo - C.P. 66318, CEP: 05315-970, São Paulo, Brazil ²Laboratoire Univers & Particules de Montpellier (LUPM), Université de Montpellier (UMR-5299)

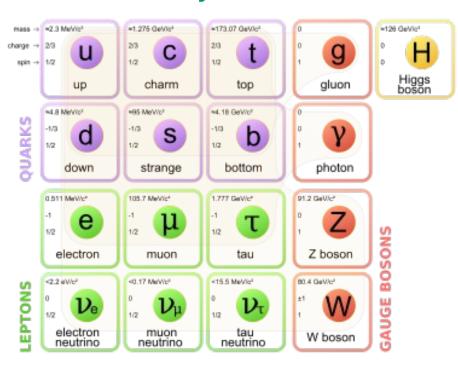
Challenges for Λ CDM Beyond H_0 and S_8

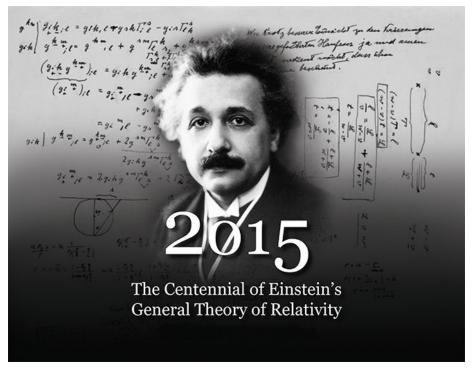
- A. The A_{lens} Anomaly in the CMB Angular Power Spectrum
- B. Hints for a Closed Universe from Planck Data
- C. Large-Angular-Scale Anomalies in the CMB Temperature and Polarization
 - 1. The Lack of Large-Angle CMB Temperature Correlations
 - 2. Hemispherical Power Asymmetry
 - 3. Quadrupole and Octopole Anomalies
 - 4. Point-Parity Anomaly
 - 5. Variation in Cosmological Parameters Over the Sky
 - 6. The Cold Spot
 - 7. Explaining the Large-Angle Anomalies
 - 8. Predictions and Future Testability
 - 9. Summary
- D. Abnormal Oscillations of Best Fit Parameter Values
- E. Anomalously Strong ISW Effect
- F. Cosmic Dipoles
 - 1. The α Dipole
 - 2. Galaxy Cluster Anisotropies and Anomalous Bulk Flows
 - 3. Radio Galaxy Cosmic Dipole
 - 4. QSO Cosmic Dipole and Polarisation Alignments
 - 5. Dipole in SNIa
 - 6. Emergent Dipole in H_0
 - 7. CMB Dipole: Intrinsic Versus Kinematic?
- G. The Ly- α Forest BAO and CMB Anomalies
 - 1. The Ly- α Forest BAO Anomaly
 - 2. Ly- α -Planck 2018 Tension in n_s - Ω_m
- H. Parity Violating Rotation of CMB Linear Polarization
- I. The Lithium Problem
- J. Quasars Hubble Diagram Tension with Planck- Λ CDM
- K. Oscillating Force Signals in Short Range Gravity Experiments
- L. ACDM and the Dark Matter Phenomenon at Galactic Scales



Are we approaching a Revolution in Physics?

Knowledge of Physics: Standard Model + General Relativity

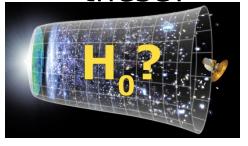


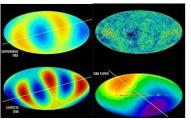


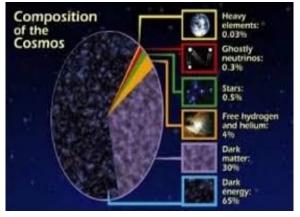
Are we approaching a Revolution in Physics?

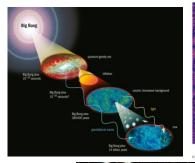
So can our knowledge of Physics describes all

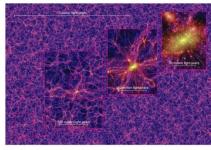
these?













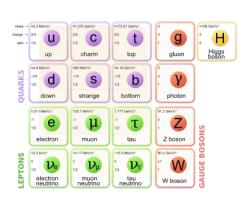


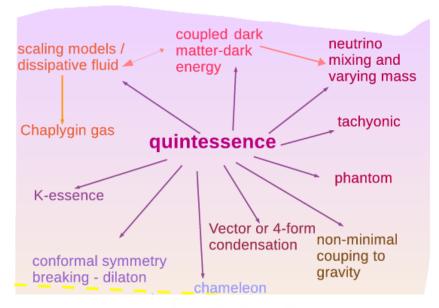
NO! We need new physics!

We need to modify something:

The universe content or The theory of Gravity

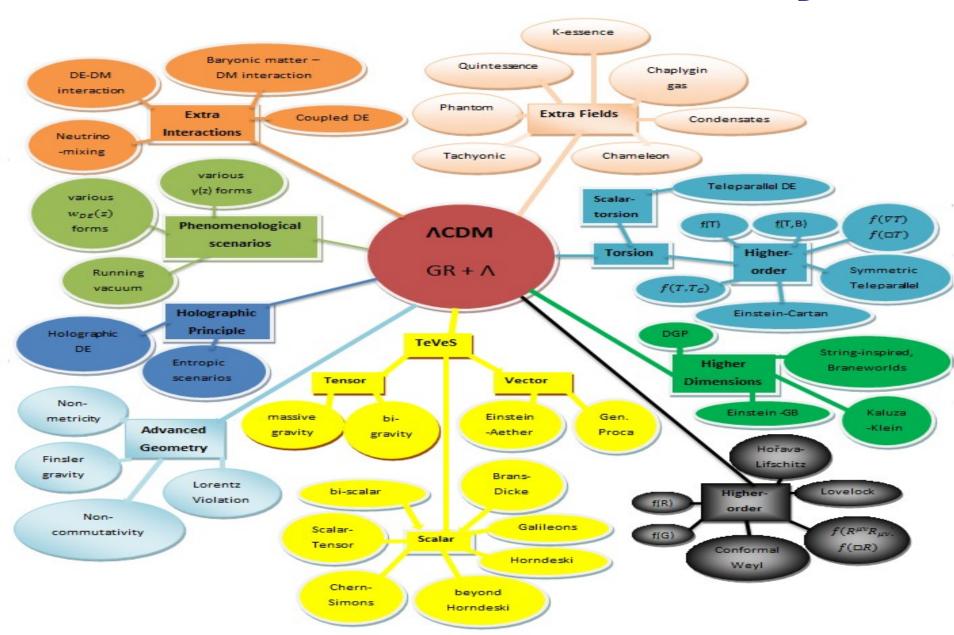
New particles/interactions







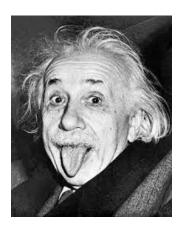
Modified Gravity



Standard Model vs General Relativity

$$\begin{array}{c} -\frac{1}{2}\partial_{\nu}g_{\mu}^{a}\partial_{\nu}g_{\mu}^{a} - g_{\nu}f^{abc}\partial_{\mu}g_{\nu}^{a}g_{\nu}^{b}g_{\nu}^{c} - \frac{1}{4}g_{\nu}^{2}f^{abc}f^{adg}g_{\mu}^{b}g_{\nu}^{c}g_{\mu}^{d}g_{\nu}^{c} + \\ -\frac{1}{2}ig_{\nu}^{2}(g_{\nu}^{a})^{\mu}g_{\mu}^{c}+G^{a}\partial^{2}G^{a} + g_{\nu}f^{abc}\partial_{\mu}G^{a}G^{b}g_{\mu}^{c}] - \partial_{\nu}W_{\mu}^{+}\partial_{\nu}W_{\mu}^{-} - \\ 2 & M^{2}W_{\mu}^{+}W_{\mu}^{-} - \frac{1}{2}\partial_{\nu}Z_{\mu}^{0}\partial_{\nu}Z_{\mu}^{0} - \frac{1}{2c_{\nu}^{2}}M^{2}Z_{\mu}^{0}Z_{\mu}^{0} - \frac{1}{2}\partial_{\mu}A^{0}\partial_{\mu}A_{\nu} - \frac{1}{2}\partial_{\mu}H\partial_{\mu}H - \\ \frac{1}{2}m_{h}^{2}H^{2} - \partial_{\mu}\phi^{+}\partial_{\mu}\phi^{-} - M^{2}\phi^{+}\phi^{-} - \frac{1}{2}\partial_{\mu}\phi^{0}\partial_{\mu}\phi^{0} - \frac{1}{2c_{\nu}^{2}}M\phi^{0}\phi^{0} - \beta_{h}[\frac{2M^{2}}{2g^{2}} + \frac{2M}{2}M + \frac{1}{2}(H^{2} + \phi^{0}\phi^{0} + 2\phi^{+}\phi^{-})] + \frac{2M^{2}}{2}\alpha_{h} - igc_{w}(\partial_{\nu}Z_{\mu}^{0}(W_{\mu}^{+}W_{\nu}^{-} - W_{\nu}^{-}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}W_{\nu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu}^{-}) - A_{\nu}(W_{\nu}^{+}\partial_{\nu$$

$$S = -\frac{1}{16\pi G} \int \sqrt{-g} (R(g) + 2\Lambda) d^4x$$



General Relativity Assumptions and Considerations

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left[R - 2\Lambda \right] + \int d^4x \left[L_m \left(g_{\mu\nu}, \psi \right) \right]$$

- Diffeomorphism invariance
- Spacetime dimensionality=4
- Geometry=Curvature (connection=Levi Civita)
- Linear in Ricci scalar
- Metric compatibility (zero non-metricity)
- Minimal matter coupling
- Equivalence principle
- Lorentz invariance
- Locality

"Those that do not know geometry are not allowed to enter".

Front Door of Plato's Academy



Descriptions of Gravity

 Einstein 1916: General Relativity: energy-momentum source of spacetime Curvature

Levi-Civita connection: Zero Torsion

Einstein 1928: Teleparallel Equivalent of GR: Weitzenbock connection: Zero Curvature

[Cai, Capozziello, De Laurentis, Saridakis, Rept.Prog.Phys. 79]

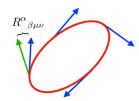
$$\begin{Bmatrix} {}^{\alpha}_{\mu\nu} \end{Bmatrix} = \frac{1}{2} g^{\alpha\lambda} (g_{\lambda\nu,\mu} + g_{\mu\lambda,\nu} - g_{\mu\nu,\lambda}). \tag{1.3}$$

The corresponding covariant derivative will be denoted by \mathcal{D} so that we will have $\mathcal{D}_{\alpha}g_{\mu\nu} = 0$. A general connection $\Gamma^{\alpha}_{\mu\nu}$ then admits the following convenient decomposition:

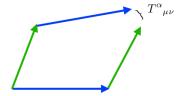
$$\Gamma^{\alpha}_{\ \mu\nu} = \begin{Bmatrix} {}^{\alpha}_{\mu\nu} \end{Bmatrix} + K^{\alpha}_{\ \mu\nu} + L^{\alpha}_{\ \mu\nu} \tag{1.4}$$

with

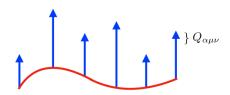
$$K^{\alpha}_{\ \mu\nu} = \frac{1}{2} T^{\alpha}_{\ \mu\nu} + T^{\alpha}_{(\mu\ \nu)}, \quad L^{\alpha}_{\ \mu\nu} = \frac{1}{2} Q^{\alpha}_{\ \mu\nu} - Q^{\alpha}_{(\mu\ \nu)}$$
 (1.5)



The rotation of a vector transported along a closed curve is given by the curvature: General Relativity.



The non-closure of parallelograms formed when two vectors are transported along each other is given by the torsion: Teleparallel Equivalent of General Relativity.



The variation of the length of a vector as it is transported is given by the non-metricity:

Symmetric Teleparallel Equivalent of General Relativity.

Metric-Affine Modified Gravity

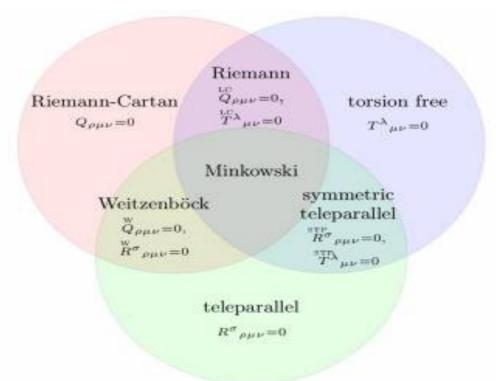


FIG. 1. Subclasses of metric-affine geometry, depending on the properties of connection.

$$\begin{split} S_{\rm GR} &= \frac{1}{2\kappa^2} \int \Big\{ g^{\mu\nu} \hat{R}_{\mu\nu} + \lambda^{\mu\nu\lambda}_{(1)} T_{\mu\nu\lambda} + \lambda^{\mu\nu\lambda}_{(2)} Q_{\mu\nu\lambda} \Big\} \sqrt{-g} \, d^4x \,, \\ S_{\rm total} &= S_{\rm GR} + S_{\rm matter} \,, \end{split}$$

Curvature and Torsion

- Vierbeins e_A^μ : four linearly independent fields in the tangent space $g_{\mu\nu}(x) = \eta_{AB} \ e_\mu^A(x) \ e_\nu^B(x)$
- Connection: ω_{ABC}
- Curvature tensor: $R_{B\mu\nu}^A = \omega_{B\nu,\mu}^A \omega_{B\mu,\nu}^A + \omega_{C\mu}^A \omega_{B\nu}^C \omega_{C\nu}^A \omega_{B\mu}^C$
- Torsion tensor: $T_{\mu\nu}^{A} = e_{\nu,\mu}^{A} e_{\mu,\nu}^{A} + \omega_{B\mu}^{A} e_{\nu}^{B} \omega_{B\nu}^{A} e_{\mu}^{B}$
- Levi-Civita connection and Contortion tensor: $\omega_{ABC} = \Gamma_{ABC} + K_{ABC}$ $K_{ABC} = \frac{1}{2} (T_{CAB} T_{BCA} T_{ABC}) = -K_{BAC}$
- Curvature and Torsion Scalars:

$$R = g^{\mu\nu}R_{\mu\nu} = g^{\mu\nu}R^{\rho}_{\mu\rho\nu}$$

$$R = \overline{R} + T - 2(T_{\nu}^{\nu\mu})_{;\mu}$$

$$T = \frac{1}{4} T^{\rho\mu\nu} T_{\rho\mu\nu} + \frac{1}{2} T^{\rho\mu\nu} T_{\nu\mu\rho} - T^{\rho}_{\rho\mu} T^{\nu\mu}_{\nu}$$

f(T) Gravity and f(T) Cosmology

- f(T) Gravity: Simplest torsion-based modified gravity
- Generalize T to f(T) (inspired by f(R))

$$S = \frac{1}{16\pi G} \int d^4x \ e \left[T + f(T) \right] + S_m$$

Equations of motion:

$$e^{-1}\partial_{\mu}\Big(ee^{\rho}_{\scriptscriptstyle{A}}S^{\mu\nu}_{\rho}\Big)\Big(1+f_{\scriptscriptstyle{T}}\Big)-e^{\lambda}_{\scriptscriptstyle{A}}T^{\rho}_{\scriptscriptstyle{\mu\lambda}}S^{\nu\mu}_{\rho}+e^{\rho}_{\scriptscriptstyle{A}}S^{\mu\nu}_{\rho}\partial_{\mu}(T)f_{\scriptscriptstyle{TT}}-\frac{1}{4}e^{\nu}_{\scriptscriptstyle{A}}[T+f(T)]=4\pi Ge^{\rho}_{\scriptscriptstyle{A}}T^{\nu\{\rm EM\}}_{\rho}$$

f(T) Cosmology: Apply in FRW geometry:

$$e^A_\mu = diag(1, a, a, a) \implies ds^2 = dt^2 - a^2(t)\delta_{ij}dx^idx^j$$
 (not unique choice)

Friedmann equations:

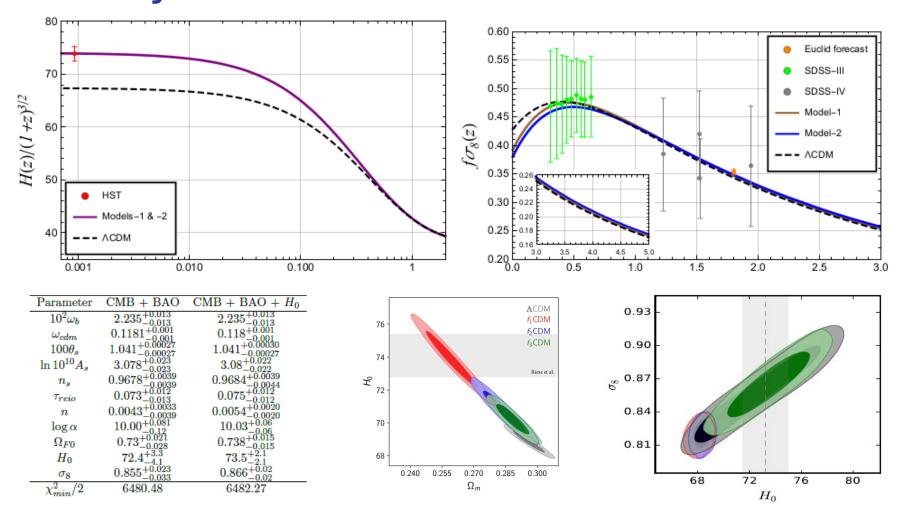
$$H^{2} = \frac{8\pi G}{3} \rho_{m} - \frac{f(T)}{6} - 2f_{T}H^{2}$$

$$\dot{H} = -\frac{4\pi G(\rho_m + p_m)}{1 + f_T - 12H^2 f_{TT}}$$

Find easily

$$T = -6H^2$$

Solving H0 and S8 tensions in f(T) Gravity

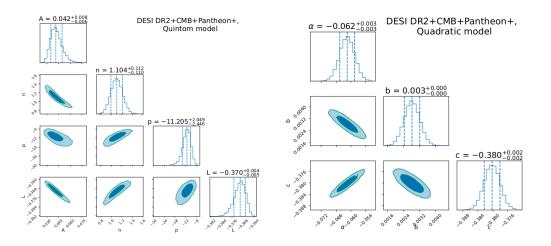


[S-F Yan, P. Zhang, J-W Chen, X_Z Zhang, Y-F Cai, E.N. Saridakis, PRD 101] [J-W Chen, W. Luo, Y-F Cai, E.N. Saridakis, PRD 102]

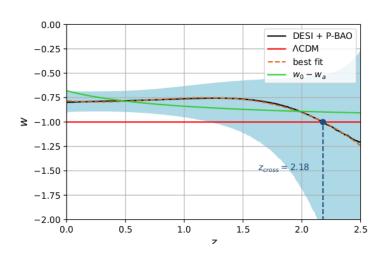
Modified gravity realizations of Quintom after DESI DR2

[Z-K Guo, Y-S Piao, X-M Zhang, Y-Z Zhang, Phys.Lett.B 608 (2005)]

$$f(T) = T + \alpha (-T)^n \left[1 - e^{pT_0/T} \right] - 2\Lambda$$



$$\frac{f(T)}{T_0} = \frac{T}{T_0} + \alpha \frac{T}{T_0} + b \frac{T^2}{T_0^2} - 2c$$



DESI DR2+CMB+Pantheon+ datasets		
Criteria	Quintom model	Quadratic model
AIC	97.81	95.146
BIC	117.824	117.1

[Y. Yang, X. Ren, Q. Wang, Z.Lu, D. zhang, YF. Cai, E. N. Saridakis. Sci.Bull. 69 (2024)] [Y. Yang, Q. Wang, X. Ren, E. N. Saridakis. YF. Cai, 2504.06784 [astro-ph.CO]]

f(Q) gravity

$$T^{\lambda}{}_{\mu\nu} \equiv \Gamma^{\lambda}{}_{\mu\nu} - \Gamma^{\lambda}{}_{\nu\mu}$$

$$R^{\sigma}{}_{\rho\mu\nu} \equiv \partial_{\mu}\Gamma^{\sigma}{}_{\nu\rho} - \partial_{\nu}\Gamma^{\sigma}{}_{\mu\rho} + \Gamma^{\alpha}{}_{\nu\rho}\Gamma^{\sigma}{}_{\mu\alpha} - \Gamma^{\alpha}{}_{\mu\rho}\Gamma^{\sigma}{}_{\nu}(5)$$

while the nonmetricity can be expressed as

$$Q_{\rho\mu\nu} \equiv \nabla_{\rho}g_{\mu\nu} = \partial_{\rho}g_{\mu\nu} - \Gamma^{\beta}{}_{\rho\mu}g_{\beta\nu} - \Gamma^{\beta}{}_{\rho\nu}g_{\mu\beta} \,. \tag{6}$$

$$R^{\sigma}{}_{\rho\mu\nu} \equiv \partial_{\mu}\Gamma^{\sigma}{}_{\nu\rho} - \partial_{\nu}\Gamma^{\sigma}{}_{\mu\rho} + \Gamma^{\alpha}{}_{\nu\rho}\Gamma^{\sigma}{}_{\mu\alpha} - \Gamma^{\alpha}{}_{\mu\rho}\Gamma^{\sigma}{}_{\nu}(5) \qquad Q = -\frac{1}{4}Q_{\alpha\beta\gamma}Q^{\alpha\beta\gamma} + \frac{1}{2}Q_{\alpha\beta\gamma}Q^{\gamma\beta\alpha} + \frac{1}{4}Q_{\alpha}Q^{\alpha} - \frac{1}{2}Q_{\alpha}\tilde{Q}^{\alpha},$$
while the nonmetricity can be expressed as

where $Q_{\alpha} \equiv Q_{\alpha \mu}^{\ \mu}$, and $\tilde{Q}^{\alpha} \equiv Q_{\mu}^{\ \mu\alpha}$.

$$S = -\frac{1}{2} \int d^4x \sqrt{-g} f(Q). \tag{8}$$

$$\frac{2}{\sqrt{-g}} \nabla_{\alpha} \left\{ \sqrt{-g} g_{\beta\nu} f_Q \left[-\frac{1}{2} L^{\alpha\mu\beta} + \frac{1}{4} g^{\mu\beta} \left(Q^{\alpha} - \tilde{Q}^{\alpha} \right) \right. \right. \\
\left. -\frac{1}{8} \left(g^{\alpha\mu} Q^{\beta} + g^{\alpha\beta} Q^{\mu} \right) \right] \right\} \\
+ f_Q \left[-\frac{1}{2} L^{\mu\alpha\beta} - \frac{1}{8} \left(g^{\mu\alpha} Q^{\beta} + g^{\mu\beta} Q^{\alpha} \right) \right. \\
\left. + \frac{1}{4} g^{\alpha\beta} \left(Q^{\mu} - \tilde{Q}^{\mu} \right) \right] Q_{\nu\alpha\beta} + \frac{1}{2} \delta^{\mu}_{\nu} f = T^{\mu}_{\nu}, \quad (9)$$
with $f_Q = \partial f / \partial Q$.

f(Q) cosmology

(20)

$$6f_Q H^2 - \frac{1}{2}f = \rho_m,$$

$$(12H^2 f_{QQ} + f_Q)\dot{H} = -\frac{1}{2}(\rho_m + p_m). \tag{12}$$

$$-a^{2}\delta\rho = 6\left(f_{Q} + 12a^{-2}\mathcal{H}^{2}f_{QQ}\right)\mathcal{H}(\mathcal{H}\phi + \varphi') + 2f_{Q}k^{2}\psi$$

$$-2\left[f_{Q} + 3a^{-2}f_{QQ}\left(\mathcal{H}' + \mathcal{H}^{2}\right)\right]\mathcal{H}k^{2}B. \quad (19)$$

$$\frac{1}{2}a^{2}\delta\rho = \left(f_{Q} + 12a^{-2}f_{QQ}\mathcal{H}^{2}\right)\left(\mathcal{H}\phi' + \varphi''\right) + \left[f_{Q}\left(\mathcal{H}' + 2\mathcal{H}^{2} - \frac{1}{3}k^{2}\right) + 12a^{-2}f_{QQ}\mathcal{H}^{2}\left(4\mathcal{H}' - \mathcal{H}^{2}\right) + 12a^{-2}\frac{df_{QQ}}{d\tau}\mathcal{H}^{3}\right]\phi$$

$$+ 2\left[f_{Q} + 6a^{-2}f_{QQ}\left(3\mathcal{H}' - \mathcal{H}^{2}\right) + 6a^{-2}\frac{df_{QQ}}{d\tau}\mathcal{H}\right]\mathcal{H}\varphi' + \frac{1}{3}f_{Q}k^{2}\psi$$

$$- \frac{1}{3}\left(f_{Q} + 6a^{-2}f_{QQ}\mathcal{H}^{2}\right)k^{2}B' - \frac{1}{3}\left[2f_{Q} + 3a^{-2}f_{QQ}\left(5\mathcal{H} - \mathcal{H}^{2}\right) + 6a^{-2}\frac{df_{QQ}}{d\tau}\mathcal{H}\right]\mathcal{H}k^{2}B, \quad (21)$$

$$+ 6a^{-2}f_{QQ}\mathcal{H}^{2}\varphi' - 9a^{-2}f_{QQ}\left(\mathcal{H}' - \mathcal{H}^{2}\right)\mathcal{H}\varphi$$

$$+ f_{Q}\psi' - a^{-2}f_{QQ}\mathcal{H}^{2}k^{2}B, \quad (20)$$

$$\delta' = (1+w)\left(-k^2v - k^2B + 3\varphi'\right) + 3\mathcal{H}\left(w\rho - \frac{\delta p}{\rho}\right), (22)$$

$$v' = -\mathcal{H}\left(1 - c_s^2\right)v + \frac{\delta p}{\rho + p} + \phi. \quad (23)$$

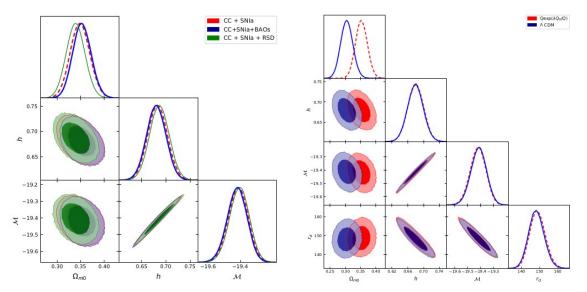
$$- f_{QQ}\mathcal{H}\left[2\mathcal{H}\varphi' + \left(\mathcal{H}' + \mathcal{H}^2\right)\phi + \left(\mathcal{H}' - \mathcal{H}^2\right)\left(\psi - B'\right)\right]$$

$$- \left[f_{QQ}\left(\mathcal{H}'^2 + \mathcal{H}\mathcal{H}'' - 3\mathcal{H}^2\mathcal{H}' - \frac{1}{3}\mathcal{H}^2k^2\right)\right]$$

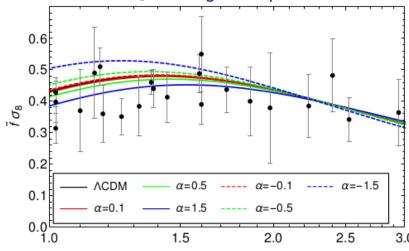
$$+ \frac{df_{QQ}}{d\tau}(\mathcal{H}' - \mathcal{H}^2)\mathcal{H}B = 0, \quad (24)$$

$$\delta'' + \mathcal{H}\delta' = \frac{4\pi G\rho}{f_Q}\delta, \qquad (30) \qquad G_{eff} \equiv \frac{G}{f_Q}, \qquad (31)$$

Solving the tensions in f(Q) gravity



[F. Anagnostopoulos, S. Basilakos, E.N.Saridakis, JCAP 2019]

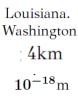


z+1 [S. Peirone, G. Benevento, N. Frusciante, S. Tsujikawa, PRD 100]

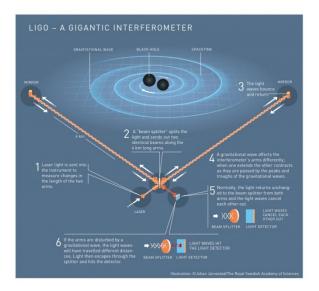
Testing GR and Modified Gravity

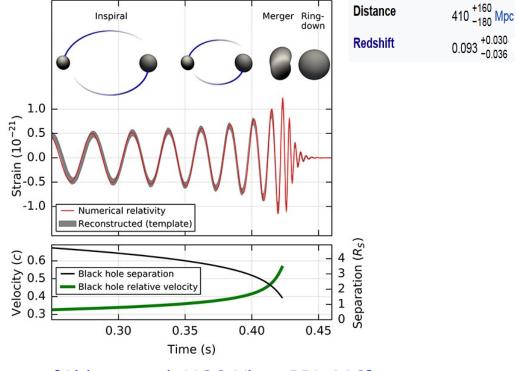
- Solar-System data
- Galaxy data
- Galaxy-cluster data
- Cosmological data (SNIa, BAO,CMB,CC, LSS)
- Early Universe (Inflation, Baryogenesis, BBN)
- Black-Hole-shadow data
- Gravitational-wave data (multi-messenger astronomy)

■ GW150914: Two black holes w36 ⁺⁵₋₄ M☉ and 29 ⁺⁴₋₄ M☉, resulting in a 62 ⁺⁴₋₄ M☉ black hole









[Abbott et al, LIGO Virgo PRL 116]
2017 Nobel Price in Physics

 In case of GWs from black hole mergers we know their properties at the moment of detection, and their direction (in case of three detectors).
 Assuming GR and ACDM we can extract their speed, distance, and properties at the moment of emission.

In case of GWs from neutron star mergers, and their E/M counterpart, we know their properties at the moment of detection and their direction, but using the implied physics from the E/M information we can extract their speed, distance and properties at the moment of emission, independently of the underlying gravitational theory and cosmological scenario.

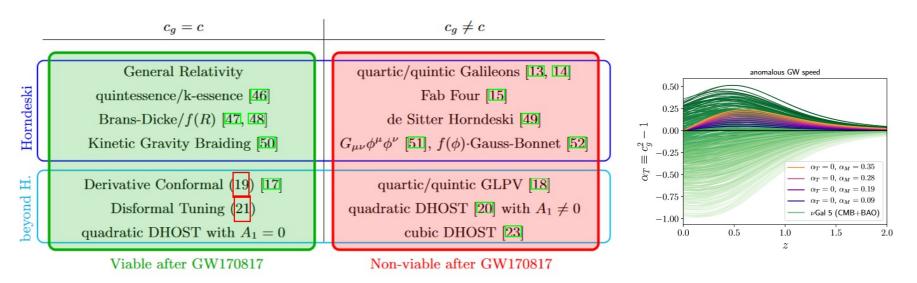
Great tool for testing General Relativity and cosmological scenarios!

• An immediate result: The speed of GWs is equal to the speed of light!

GW170817 time delay $1.74 \pm 0.05s$ constrains:

 $-3 \cdot 10^{-15} \le c_g/c - 1 \le 7 \cdot 10^{-16}$

Excludes a large number of theories that were consistent with other data!



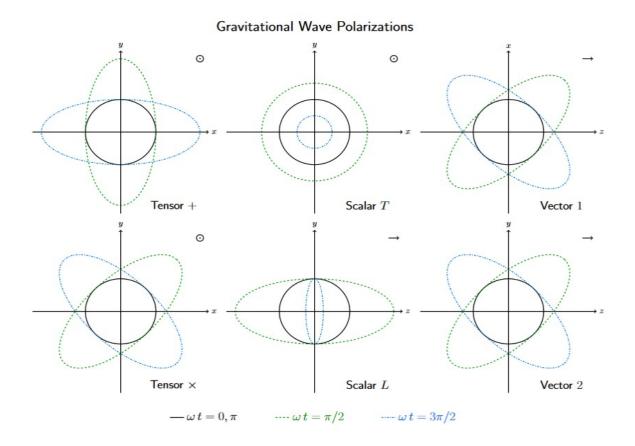
For tensor perturbatior $g_{00} = -1$, $g_{0i} = 0$, $g_{ij} = a^2 \left(\delta_{ij} + h_{ij} + rac{1}{2} h_{ik} h_{kj}
ight)$

$$\ddot{h}_{ij} + (3 + \alpha_M)\dot{h}_{ij} + (1 + \alpha_T)\frac{k^2}{a^2}h_{ij} = 0$$

$$\alpha_M = \frac{d\log(M_*^2)}{d\log a} \qquad c_g^2 = (1 + \alpha_T)$$

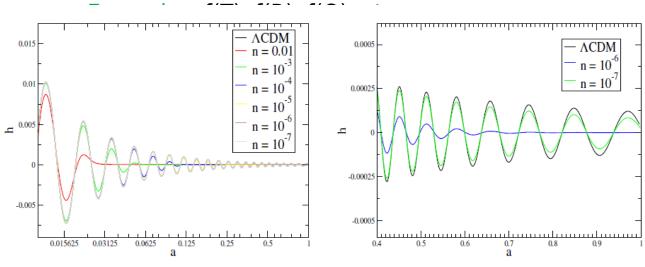
$$h_{\rm GW} \sim h_{\rm GR} \quad \underbrace{e^{-\frac{1}{2}\int \nu \mathcal{H} d\eta}}_{\text{Affects amplitude}} \underbrace{e^{ik\int (\alpha_T + a^2m^2/k^2)^{1/2} d\eta}}_{\text{Affects phase}}$$

Polarizations:



Gravitational waves in modified gravity

- Testing General Relativity, modified gravities, and various cosmological scenarios.
- The GWs properties at emission and detection are determined by them.



$$h_{\mu\nu}^{(1)} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 2\gamma_1^{(1)1} & B_1^2 \exp(ip_\mu x^\mu) & 0 \\ 0 & B_1^2 \exp(ip_\mu x^\mu) & -2\gamma_1^{(1)1} & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

[Cai, Li, Saridakis, Xue PRD97]

[Li, Cai, Cai, Saridakis, JCAP 1810]

[Farrugia, Said, Gakis, Saridakis, PRD97]

Gravitational Waves in Modified Teleparallel Theories

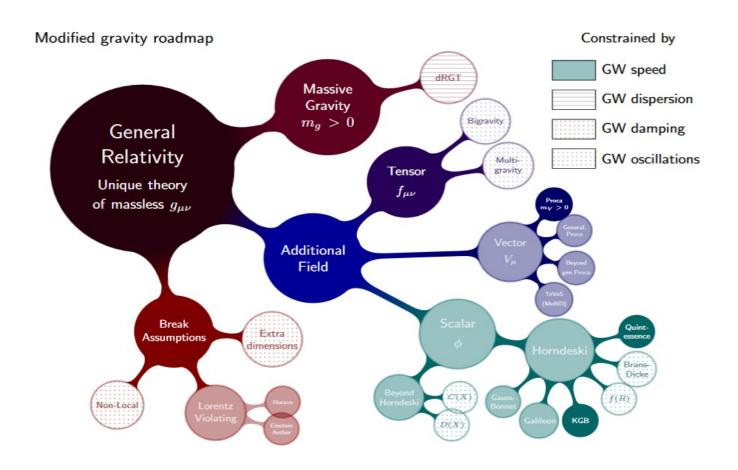
$$S = \frac{1}{16\pi G} \int d^4x \, e \, f(T,B) + \int d^4x \, e \, \mathcal{L}_m$$
 $R = -T - 2\nabla^{\mu}T^{\nu}_{\ \mu\nu}$

$$-f_T G_{\mu\nu} + (g_{\mu\nu} \Box - \nabla_{\mu} \nabla_{\nu}) f_B$$
$$+ \frac{1}{2} g_{\mu\nu} (f_B B + f_T T - f)$$
$$+ 2S_{\nu \mu}^{\alpha} \partial_{\alpha} (f_T + f_B) = 8\pi G \Theta_{\mu\nu}$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}^{(1)} + \mathcal{O}\left(h_{\mu\nu}^{(2)}\right)$$

$$h_{\mu\nu}^{(1)} = \begin{pmatrix} -2A \exp(ik_{\mu}x^{\mu}) - \frac{f_{BB}^{(0)}B^{(1)}}{f_{T}^{(0)}} & B_{1} \exp(ik_{\mu}x^{\mu}) & B_{2} \exp(ik_{\mu}x^{\mu}) & -2A \exp(ik_{\mu}x^{\mu}) \\ B_{1} \exp(ik_{\mu}x^{\mu}) & h_{+} + \frac{f_{BB}^{(0)}B^{(1)}}{f_{T}^{(0)}} & h_{\times} & B_{1} \exp(ik_{\mu}x^{\mu}) \\ B_{2} \exp(ik_{\mu}x^{\mu}) & h_{\times} & -h_{+} + \frac{f_{BB}^{(0)}B^{(1)}}{f_{T}^{(0)}} & B_{2} \exp(ik_{\mu}x^{\mu}) \\ -2A \exp(ik_{\mu}x^{\mu}) & B_{1} \exp(ik_{\mu}x^{\mu}) & B_{2} \exp(ik_{\mu}x^{\mu}) & -2A \exp(ik_{\mu}x^{\mu}) + \frac{f_{BB}^{(0)}B^{(1)}}{f_{T}^{(0)}} \end{pmatrix}$$

Gravitational waves and Modified Gravity



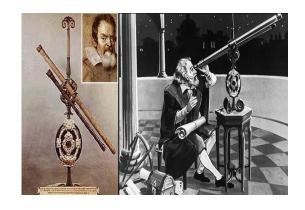
5000 years of observations 500 years of organized observations



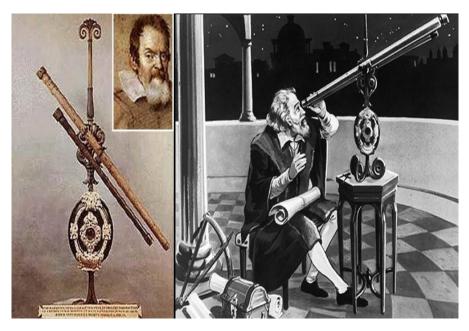








We are approaching a Revolution!

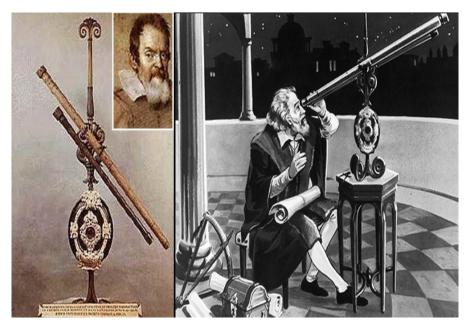




EM observations: 400 years

GW observations: 10 years

We are approaching a Revolution!





EM observations: 400 years

GW observations: 10 years

Thank you!