

Spin polarization in off-equilibrium medium at weak and strong coupling



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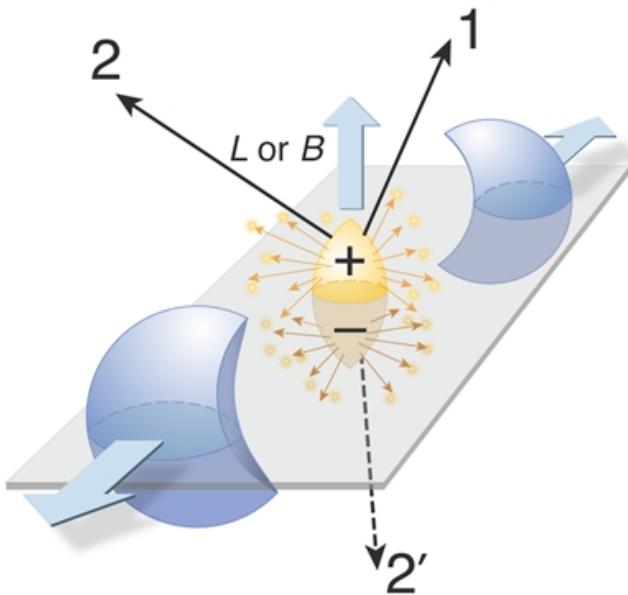
Holographic applications: from Quantum Realms to the Big Bang,
Beijing, July 11-19, 2025

Li, SL, 2505.21883
SL, Tian, 2410.22935

Outline

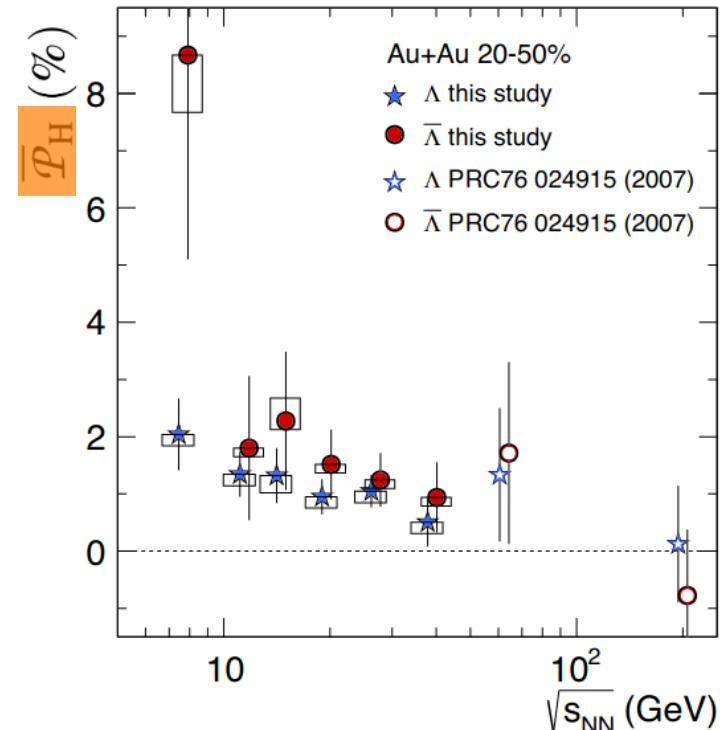
- ◆ Rich spin phenomena in heavy ion collisions
- ◆ What insights we wish to gain from holography
- ◆ Spectral function of particle in off-equilibrium state
- ◆ Polarized spectral function at weak coupling
- ◆ Polarized spectral function at strong coupling
- ◆ Summary

Global polarization of Λ ($S=1/2$)



$$L_{ini} \sim 10^5 \hbar \rightarrow S_{final}$$

Liang, Wang, PRL 2005, PLB 2005

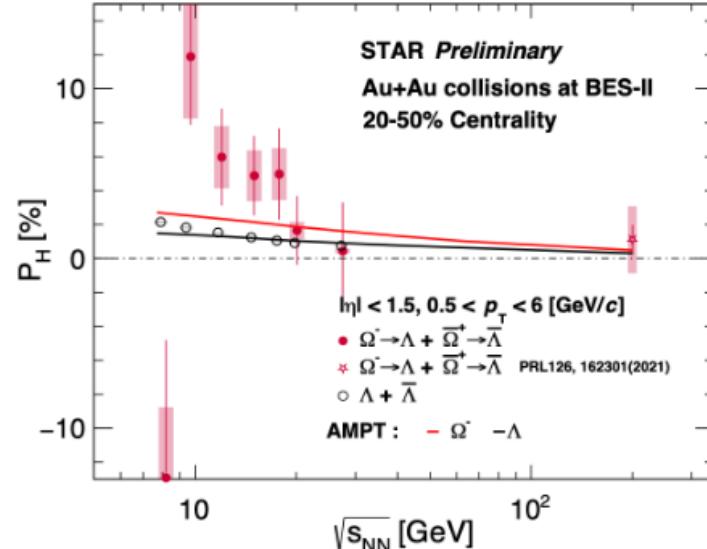
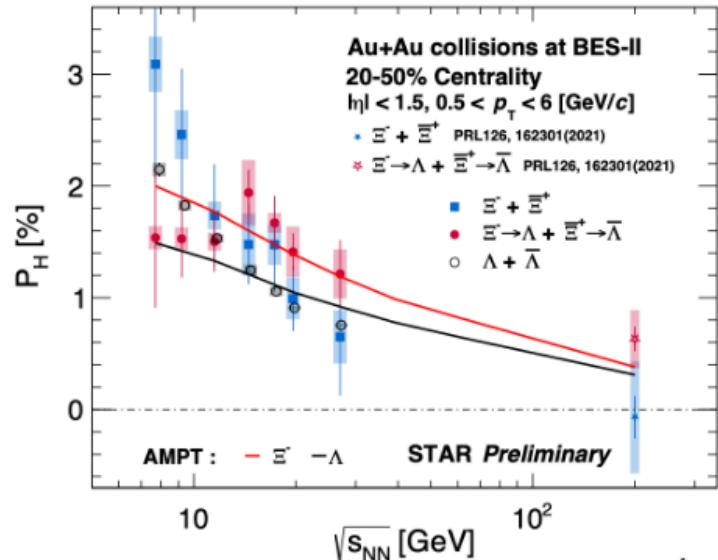


STAR collaboration,
Nature 2017

$$e^{-\beta(H_0 - \mathbf{S} \cdot \boldsymbol{\omega})}$$

Becattini et al, PRC 2017

Global polarization of Ξ , Ω ($S=1/2$)



point particle

$$P_H \simeq \frac{(s+1)}{3} \frac{\omega}{T}$$

Becattini et al, PRC 2017

quark model

$$P_\Lambda = P_s \quad P_\Xi \simeq \frac{4 P_s - P_d}{3}$$

$$P_\Omega \simeq \frac{5 P_s}{3}$$

Liang, Wang, PRL 2005

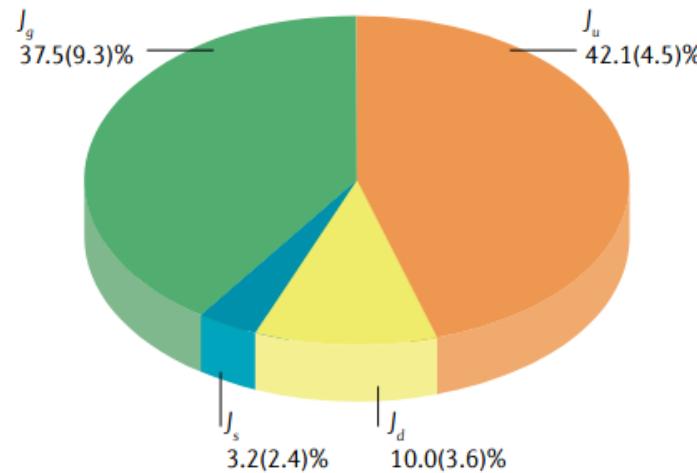
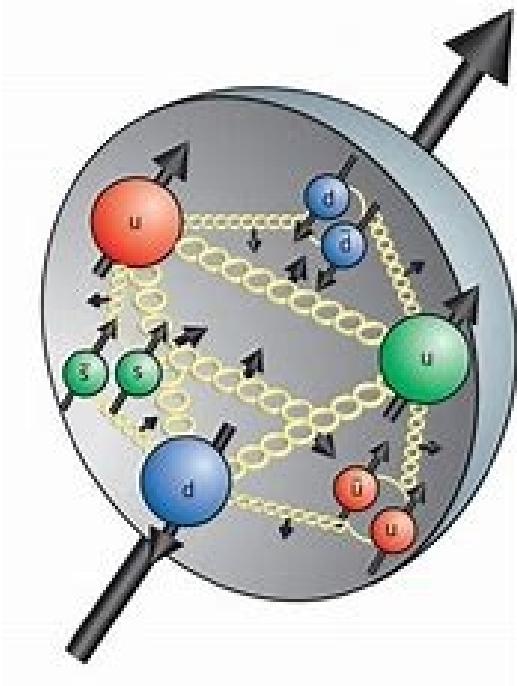
→ $P_\Omega \simeq \frac{5 P_\Lambda}{3}$

$$P_\Xi \simeq P_\Lambda$$

same prediction from non-interacting nature of models

disfavored by data?

Interaction effect significant

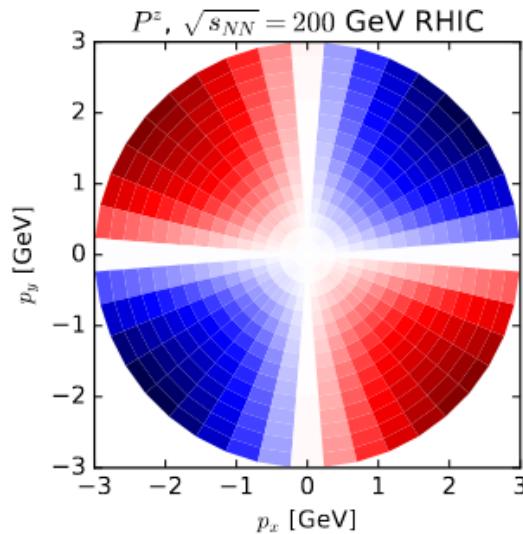
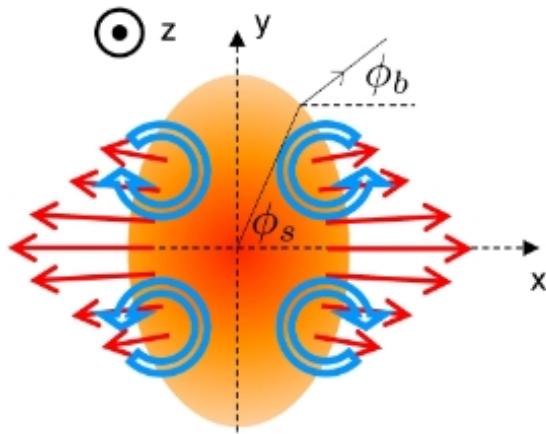


$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + L_q + \Delta G + L_g$$

~40%

Alexandrou et al, PRD 2020

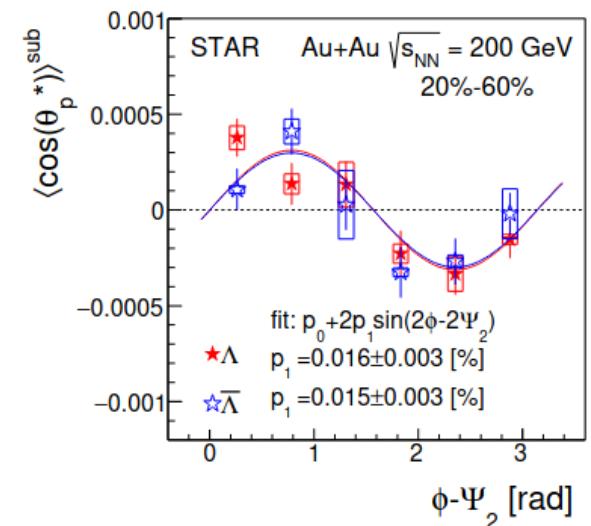
Local polarization of Λ : vorticity only



$$S^i \sim \omega^i$$

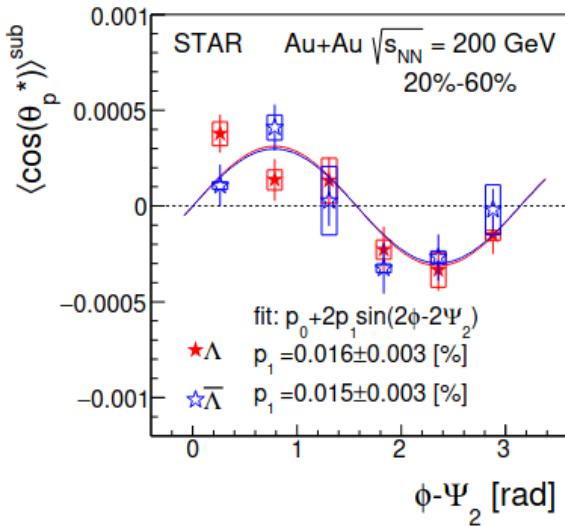
Becattini, Karpenko, PRL 2018
Wei, Deng, Huang, PRC 2019
Wu, Pang, Huang, Wang, PRR 2019
Fu, Xu, Huang, Song, PRC 2021

wrong sign

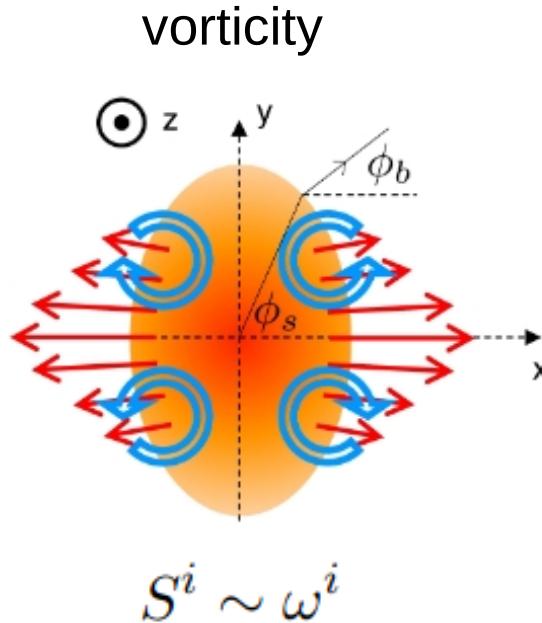


STAR collaboration, PRL 2019

Local polarization of Λ : vorticity + shear



STAR collaboration, PRL
2019



wrong sign

shear

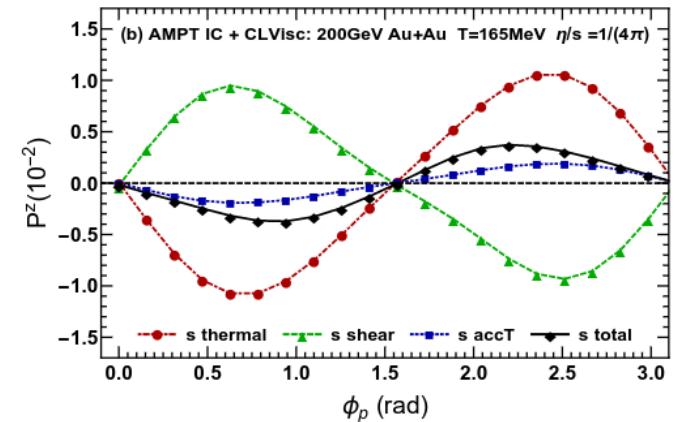
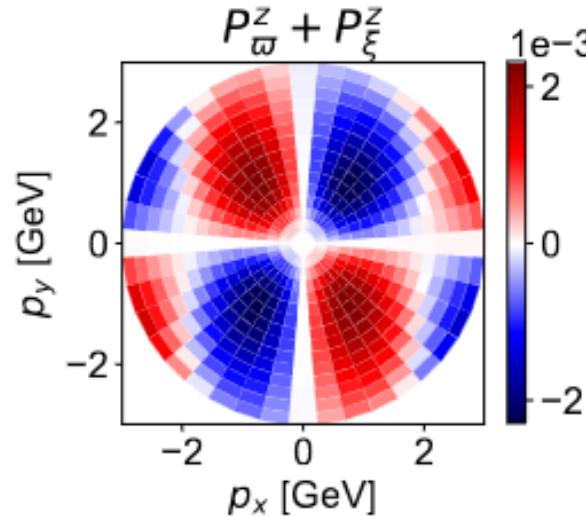
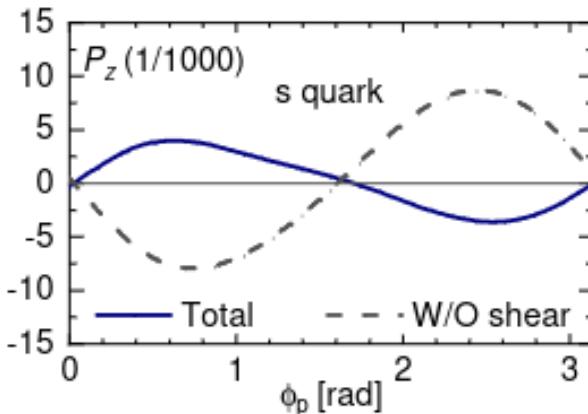
$$S^i \sim \epsilon^{ijk} \hat{p}_k \hat{p}_l \sigma_{jl}$$

$$S^z \sim (\langle p_y^2 \rangle - \langle p_x^2 \rangle) \partial_y u_x$$

right sign

Hidaka, Pu, Yang, PRD 2018
Liu, Yin, JHEP 2021
Becattini, et al, PLB 2021

Sensitivity on interaction: shear



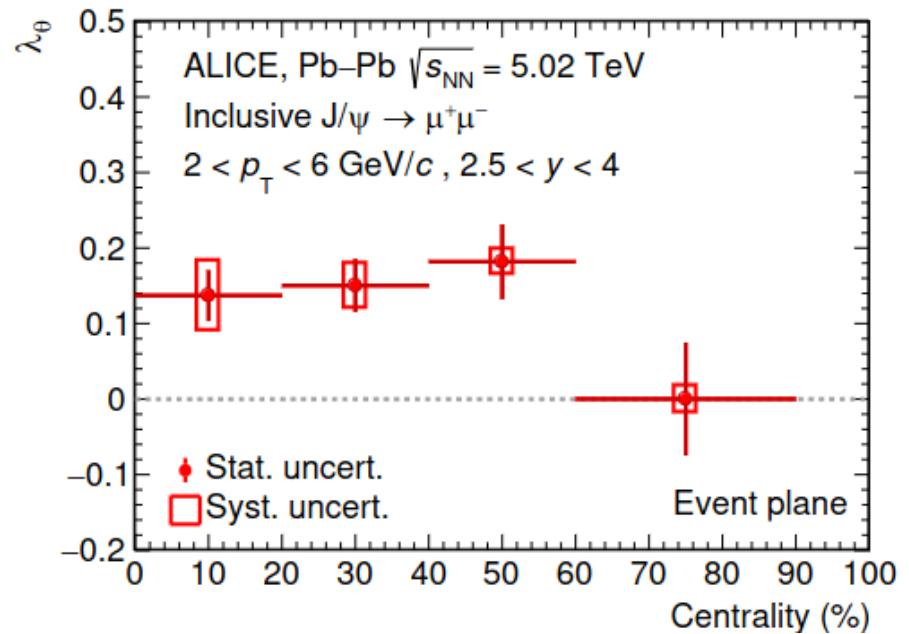
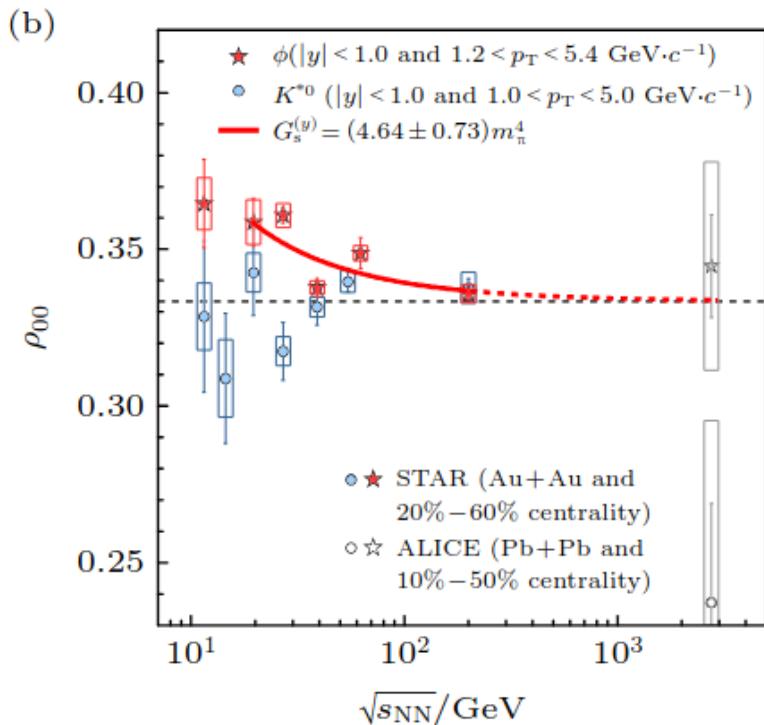
two scenarios
 Λ : point particle
 Λ : quark model

$$S^\mu(p^\alpha) = \frac{1}{4m} \frac{\int d\Sigma \cdot p n_0 (1 - n_0) \mathcal{A}^\mu}{\int d\Sigma \cdot p n_0},$$

Interaction effect?

Fu, Liu, Pang, Song, Yin, PRL 2021
 Becattini, et al, PRL 2021
 Yi, Pu, Yang, PRC 2021

Spin alignment of ϕ , J/ψ ($S=1$)



splitting of longitudinal/transverse polarizations

Holographic studies: Hou et al PRD 2024, JHEP 2024, Huang et al PRD 2025

Interaction effect in medium: weak coupling picture

free theory $S^i \sim \left(1 \omega^i + 1 \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} + 1 \epsilon^{ijk} \hat{p}_j \frac{\partial_k \beta}{\beta} \right) \delta(p^2 - m^2)$

- distribution function correction

$$(\partial_t + \hat{p} \cdot \nabla_x) f(t, x, p) = -C[f] \quad \delta f \propto \frac{\nabla_x f}{g^4}$$

$$S^i \sim \left((1 + O(1)) \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} \right) \delta(P^2 - m^2)$$

deviation of steady state
from local equilibrium

SL, Wang, JHEP 2022, PRD 2025
Fang, Pu, Yang, PRD 2024
Fang, Pu, PRD 2025

- spectral function correction in off-equilibrium state **focus of this talk**

SL, Tian, 2410.22935
Fang, Pu, Yang, 2503.13320

Spectral function of probe “particle” in medium

$$\rho_{\alpha\beta}(\omega, \vec{p}) = \int d^4x e^{i\vec{p}\cdot\vec{x}} \left\langle \Psi_\alpha(x) \Psi_\beta^\dagger(0) + \Psi_\beta^\dagger(0) \Psi_\alpha(x) \right\rangle$$

$\omega, p \gg \partial_X$ particle localized in fluid element

$\langle \dots \rangle = \text{Tr}[D \dots]$ D: **density matrix**

probe particle in off-equilibrium fluid $D = D_{\text{probe}} \otimes D_{\text{fluid}}$.

$$\begin{array}{ccc} D_{\text{probe}}^{(0)} \otimes D_{\text{fluid}}^{(1)} & & D_{\text{probe}}^{(1)} \otimes D_{\text{fluid}}^{(0)} \\ \rho = 2 \text{Im } G^R \quad \longrightarrow \quad S^i \sim \text{tr}[\rho \sigma^i] & & \end{array}$$

Off-equilibrium spectral function at weak coupling

- Probe quark

Spectral function from self-energy

$$\frac{i}{2} \not{\partial} S_R(X, P) + \not{P} S_R(X, P) - \left(\Sigma_R(X, P) S_R(X, P) + \frac{i}{2} \{ \Sigma_R(X, P), S_R(X, P) \}_{\text{PB}} \right) = -1$$

$$S_R = S_R^{(0)} + S_R^{(1)} + \dots , \quad \{A, B\}_{\text{PB}} = \partial_P A \cdot \partial_X B - \partial_X A \cdot \partial_P B$$

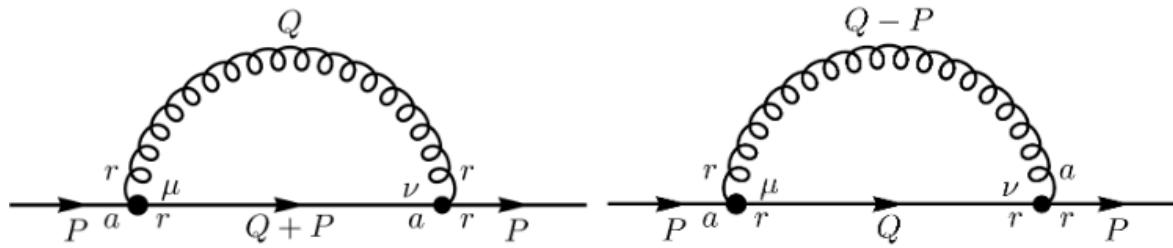
$$S_R^{(0)} = -\frac{1}{\not{P}} - \frac{1}{\not{P}} \Sigma_R \frac{1}{\not{P}}$$

$$S_R^{(1)} = -\frac{1}{\not{P}} \delta \Sigma_R \frac{1}{\not{P}} + \gamma^5 \gamma^\beta P^\nu T^{\mu\lambda} \epsilon_{\beta\lambda\mu\nu} \frac{-1}{(P^2)^2} \quad T_{\mu\lambda} = \partial_{[\mu} \Sigma_{\lambda]}^R$$

off-equilibrium self-energy

gradient of equilibrium self-energy

Equilibrium self-energy



$P \gg T$, $P^2 \ll p/\beta$,
energetic quark
nearly on-shell

$$\frac{\Sigma_{ar}}{g^2 C_F} = 2i\cancel{P}(A + B) + 4ip_0\gamma^0 A.$$

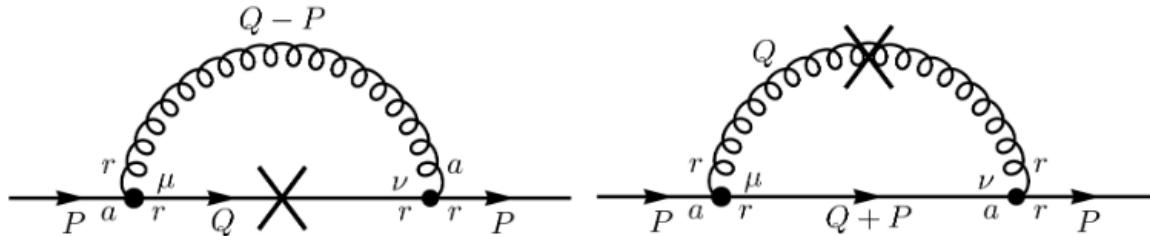
→ $p_0 = p(1 + 8A)$

$$A = \frac{1}{2(2\pi)^2} \frac{-i\pi}{2p\beta}.$$

modified dispersion:
finite damping

$$S_R^{(1)} = \gamma^5 \gamma_\beta P_\nu \partial_\mu \Sigma_\lambda^R \frac{-1}{(P^2)^2} \epsilon^{\beta\lambda\mu\nu} \sim O(T\partial) \quad D^{(0)}_{\text{probe}} \otimes D^{(1)}_{\text{fluid}}$$

Off-equilibrium self-energy



off-equilibrium propagators

Hidaka, Pu, Yang 2017

Huang et al 2020

Hattori et al 2020

$$\frac{\delta \Sigma_{ar}}{g^2 C_F} = \gamma^5 \gamma^\mu \mathcal{A}_\mu,$$

$$\mathcal{A}^0 = i\omega^i p_i \beta (-4\delta A - 2\delta B),$$

$$\mathcal{A}^k = i\omega_\parallel^k \beta p (-4\delta A - 2\delta B) + \omega_\perp^k \beta p (-4\delta A - \delta B) + \epsilon^{ijk} \hat{p}_i \hat{p}_l \sigma_{jl} \beta p (-\delta B) + \epsilon^{ijk} p_i \partial_j \beta (-\delta C)$$

$$\delta S_R^{(0)} = -\frac{1}{P} \delta \Sigma_R \frac{1}{P} \sim O(p\partial) \quad \text{dominant contribution} \quad P \gg T.$$

$$D^{(1)}_{\text{probe}} \otimes D^{(0)}_{\text{fluid}}$$

Weak coupling summary

free

$$\gamma^5 \gamma_i \frac{2\pi}{2} \left(\omega^i + \epsilon^{ijk} \hat{p}_k \frac{\partial_j \beta}{\beta} + \epsilon^{ijk} \hat{p}_k \hat{p}_l \sigma_{lj} \right) \beta \tilde{f}(p),$$

interaction

$$\int dp_0 \delta\rho(P) \tilde{f}(p_0)$$

$$O(\lambda) = \frac{g^2 C_F}{2(2\pi)^2} \frac{\pi}{2p} \gamma^5 \gamma_i \left[0.95 \omega_{||}^i + 1.48 \omega_{\perp}^i - 0.52 \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} - 0.02 \epsilon^{ijk} \hat{p}_j \frac{\partial_k \beta}{\beta} \right] \tilde{f}(p)$$

Polarized spectral function

Off-equilibrium spectral function at strong coupling

- Probe baryon

Holographic model for baryon

$$S = i \int d^{D+1}x \sqrt{-g} \bar{\psi} (\Gamma^M \nabla_M - m) \psi,$$

Iqbal, Liu 2009

5D Dirac fermion  4D Weyl fermion

$$\psi = \begin{pmatrix} \psi_R \\ \psi_L \end{pmatrix}$$

“lump of quark/gluon”

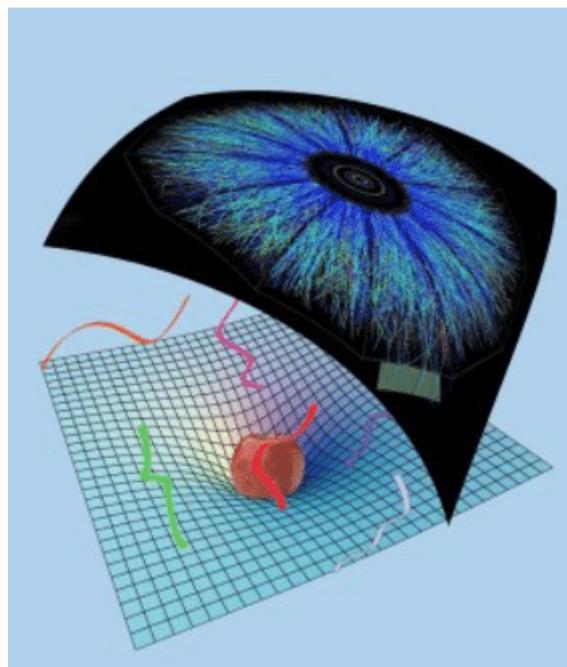
One baryon = R + L Weyl fermions

focus on spectral function in off-equilibrium QGP

Holographic model for QGP

$$ds^2 = -2u_\mu(x)dx^\mu dr - r^2 f(b(x)r)u_\mu u_\nu dx^\mu dx^\nu + r^2 P_{\mu\nu} dx^\mu dx^\nu \quad \text{local equilibrium}$$

$$+ 2r^2 b F(br) \sigma_{\mu\nu} dx^\mu dx^\nu + \frac{2}{3} r u_\mu u_\nu \partial_\lambda u^\lambda dx^\mu dx^\nu - r u^\lambda \partial_\lambda (u_\mu u_\nu) dx^\mu dx^\nu \quad \text{steady state}$$



Bhattacharyya et al,
JHEP 2008

$$O(\partial^0)$$

$$O(\partial)$$

$$T^{\mu\nu} = (\pi T)^4 (\eta^{\mu\nu} + 4 u^\mu u^\nu) - 2 (\pi T)^3 \sigma^{\mu\nu}$$

local equilibrium

η
steady state

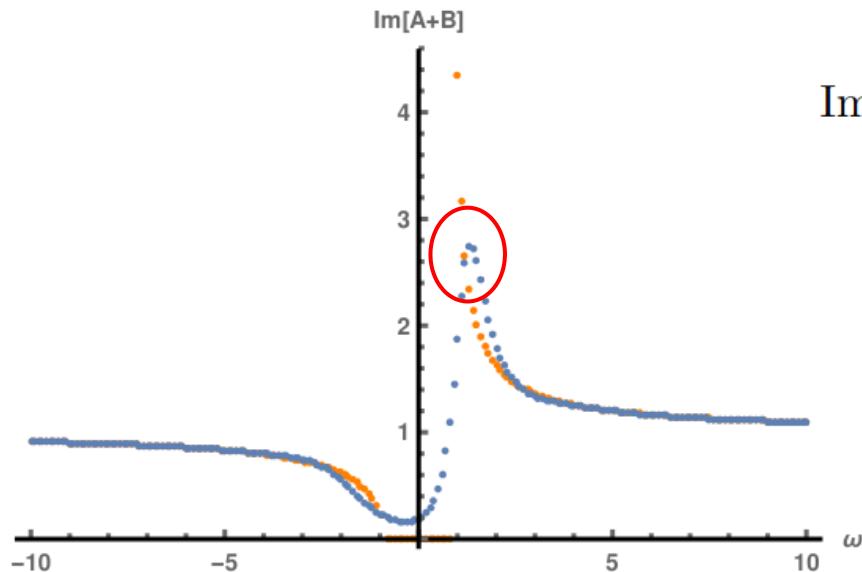
$$\sigma_{ij} - \frac{\partial_0 b}{3} \partial_i u_i - \partial_i b = \partial_0 u_i$$

T-grad =
acceleration

Equilibrium spectral function

$$G_R = A + B\hat{p} \cdot \vec{\sigma}$$

$$\rho = 2\text{Im}G_R = \text{Im}[A + B](1 - \hat{p} \cdot \vec{\sigma}) + \text{Im}[A - B](1 + \hat{p} \cdot \vec{\sigma})$$



orange: vacuum

Iqbal, Liu 2009

$$\text{Im}[A + B] = \text{sgn}(\omega)\theta(\omega^2 - p^2) \left(\frac{\omega + p}{\omega - p} \right)^{1/2}$$

no spacelike spectral

blue: equilibrium QGP
soften the singularity
develop spacelike spectral

focus on timelike spectral for baryon

Gradient corrections

$$(\Gamma^M \nabla_M - m) \psi = 0$$

gradient correction to Dirac field  gradient correction to baryon

$$D_{\text{fluid}}^{(0)} \otimes D_{\text{baryon}}^{(1)}$$

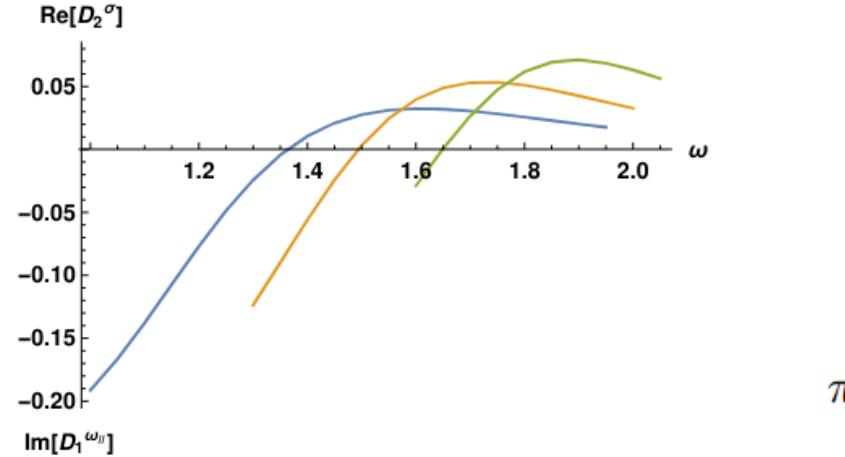
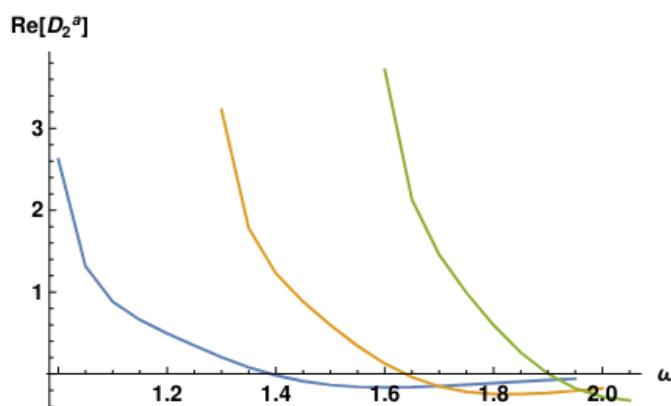
$$(\Gamma^M \nabla_M - m) \psi = 0$$

gradient correction to Dirac operator  gradient correction to medium

$$D_{\text{fluid}}^{(1)\text{le}} \otimes D_{\text{baryon}}^{(0)} \\ D_{\text{fluid}}^{(1)\text{ss}}$$

$$\frac{1}{2} \text{tr} [(\delta\rho^L + \delta\rho^R)\sigma_k] = 4 \left(-\text{Re}[D_2^a]\epsilon^{ijk}\hat{p}_j\partial_0 u_i - \text{Re}[D_2^\sigma]\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il} + \text{Im}[D_1^\omega]\omega_k + \text{Im}[D_1^{\omega\parallel}]\omega_k^\parallel \right)$$

$$D_{\text{fluid}}^{(0)} \otimes D_{\text{baryon}}^{(1)}$$

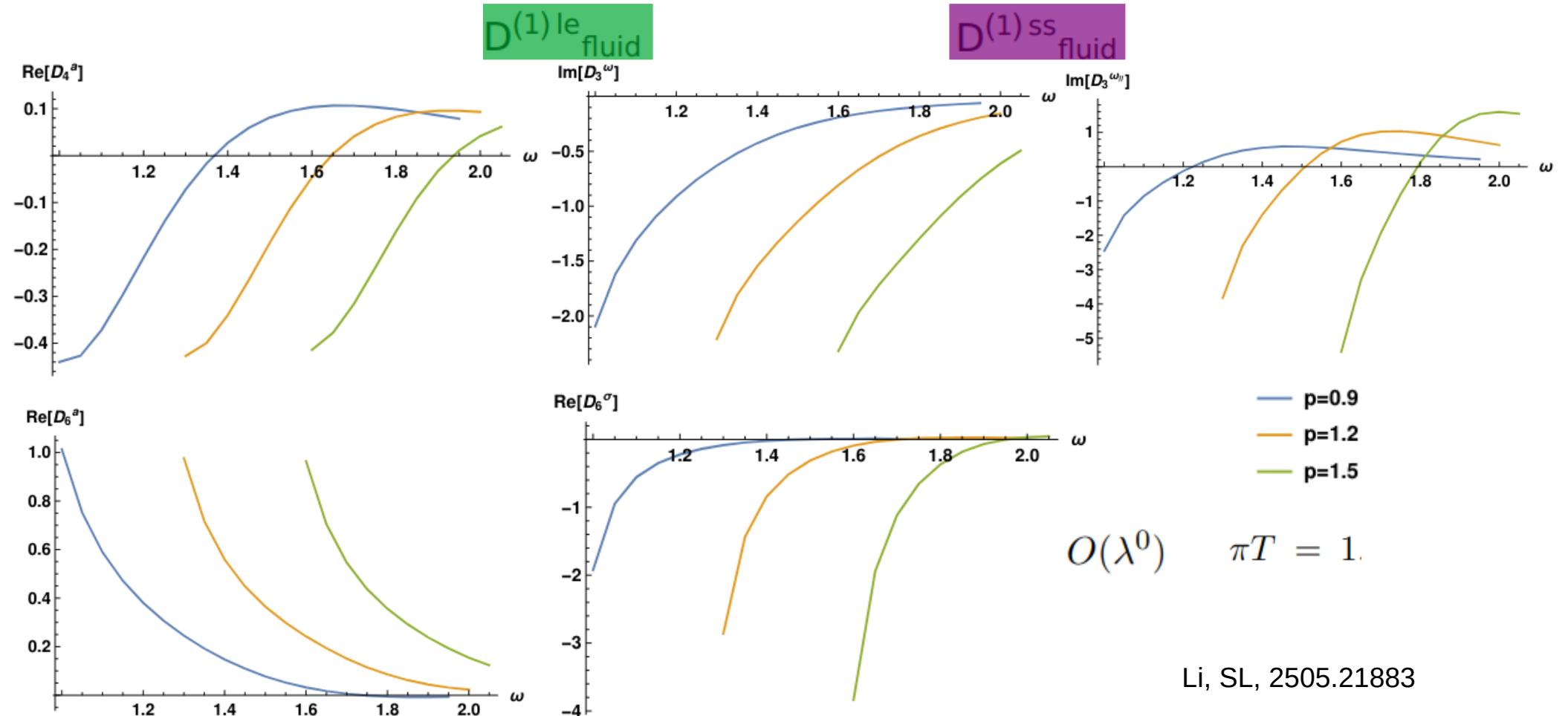


$$\pi T = 1$$

$$O(\lambda^0)$$

Li, SL, 2505.21883

$$\frac{1}{2} \text{tr}[(\delta\rho^L + \delta\rho^R)\sigma_k] = 4 \left[-\text{Re}[D_4^a + D_6^a]\epsilon^{ijk}\hat{p}_j\partial_0 u_i - \text{Re}[D_6^\sigma]\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il} + \text{Im}[D_3^\omega]\omega_k + \text{Im}[D_3^\omega]\omega_k^\parallel \right]$$



Strong coupling summary

$\text{tr}[(\delta\rho^L + \delta\rho^R)\sigma_k]$	$\epsilon^{ijk}\hat{p}_j\partial_0 u_i$	$\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il}$	ω_k^\perp	ω_k^\parallel
$D_{\text{probe}}^{(1)}$	$-4\text{Re}[D_2^a]$	$-4\text{Re}[D_2^\sigma]$	$4\text{Im}[D_1^\omega]$	$4\text{Im}[D_1^\omega + D_1^{\omega\parallel}]$
$D_{\text{fluid}}^{(1)\text{le}}$	$-4\text{Re}[D_4^a]$		$4\text{Im}[D_3^\omega]$	$4\text{Im}[D_3^\omega + D_3^{\omega\parallel}]$
$D_{\text{fluid}}^{(1)\text{ss}}$	$-4\text{Re}[D_6^a]$	$-4\text{Re}[D_6^\sigma]$		

Polarized spectral function

Weak (quark) vs Strong (baryon)

Common

- ◆ Polarized spectral function in off-equilibrium state
 $D_{\text{probe}}^{(1)}$ $D_{\text{fluid}}^{(1)\text{le}}$ $D_{\text{fluid}}^{(1)\text{ss}}$
- ◆ Correction suppressed for energetic particles

Difference

- ◆ Weak $\mathcal{O}(\lambda)$ vs strong $\mathcal{O}(\lambda^0)$

Summary

- ◆ Sensitiveness of polarization to interaction
- ◆ Spectral function in off-equilibrium state
- ◆ vorticity, shear, T-grad generically lead to (interaction dependent) polarized spectral function

Outlook

- ◆ Mass dependence of polarized spectral function
- ◆ Direct access to polarization: Schwinger-Keldysh extended holography