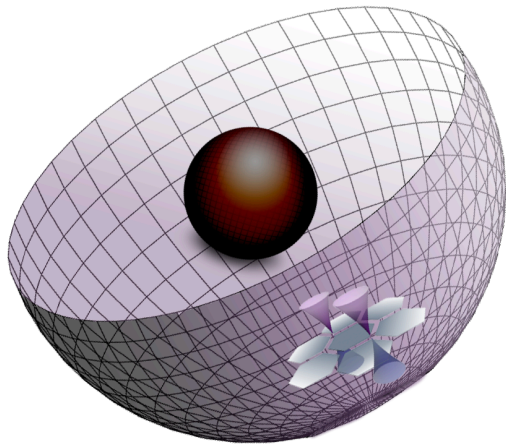


ER=EPR, and Strange metal from Entanglement

A parallelism between Y-SYK and quantum gravity



Sang-Jin Sin (Hanyang U.)
2025.07@Beijing

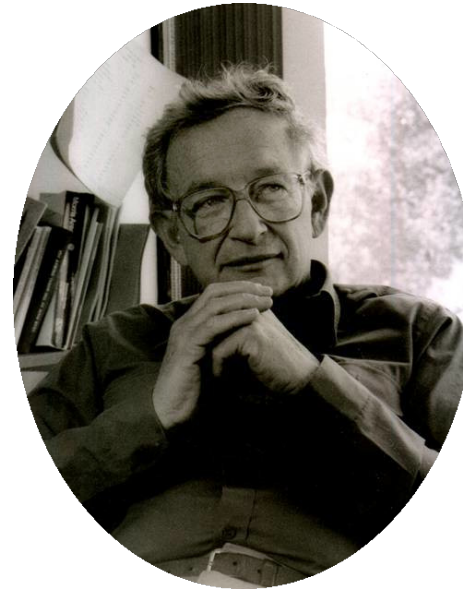
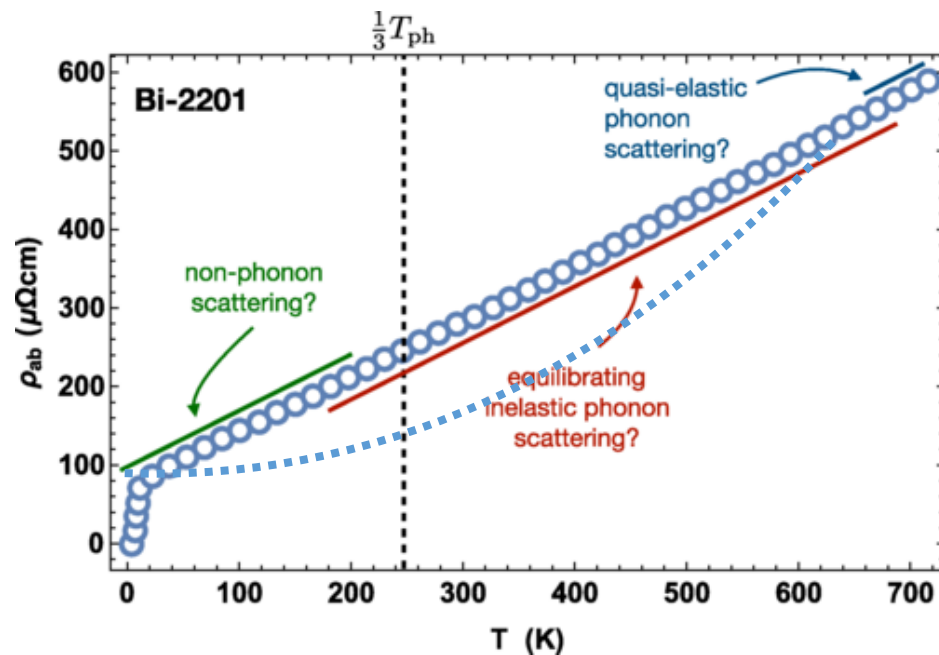
Strange Metal.

1

$$\rho \sim T$$

2

$$\sigma_{xx}/\sigma_{xy} \sim T^2$$



Essence of High T_c is in the Normal state!

References

- P.W. Anderson, Phys. Rev. Lett. 67 (1991) 2092.
- P.W. Anderson, Physics Today 66 (2013) 9.

Problem

- Exp: Universal behavior of strongly correlated metal,
Theory: Difficult to write a (solvable QFT) model for 35 years.
- Even the lowest dim. H_{int} produces $\rho \sim T^\alpha$, $\alpha > 1$,
==> No microscopic understanding for mechanism
- Phenomenology: momentum is dissipated Strangely Fast : $1/\tau \sim T$
You will see why!

A model for Linear Resistivity

2023.Aug. Patel and Sachdev et.al, wrote 2+1 model based on SYK model.

The screenshot shows the top portion of a Science journal article page. At the top left is the 'Science' logo in red. To its right are navigation links: 'Current Issue', 'First release papers', 'Archive', and 'About' with a dropdown arrow. A 'Submit mail' button is on the far right. Below this is a breadcrumb trail: 'HOME > SCIENCE > VOL. 381, NO. 6659 > UNIVERSAL THEORY OF STRANGE METALS FROM SPATIALLY RANDOM INTERACTIONS'. Underneath the breadcrumb is a row with a lock icon, the text 'RESEARCH ARTICLE | PHYSICS', and a series of social media icons (Facebook, X, LinkedIn, etc.). The main title 'Universal theory of strange metals from spatially random interactions' is prominently displayed in large, bold black font. Below the title, the authors are listed: 'AAVISHKAR A. PATEL', 'HAOYU GUO', 'ILYA ESTERLIS', and 'SUBIR SACHDEV', each followed by an ORCID icon. A link for 'Authors Info & Affiliations' is to the right. Below the authors, the publication details are given: 'SCIENCE • 17 Aug 2023 • Vol 381, Issue 6659 • pp. 790-793 • DOI: 10.1126/science.abq6011'. At the bottom of this section, there is a download icon followed by the number '7,449', and on the right, icons for notifications, bookmarks, and a red circular icon with a white document symbol.

Science

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HOME > SCIENCE > VOL. 381, NO. 6659 > UNIVERSAL THEORY OF STRANGE METALS FROM SPATIALLY RANDOM INTERACTIONS

RESEARCH ARTICLE | PHYSICS

Universal theory of strange metals from spatially random interactions

AAVISHKAR A. PATEL, HAOYU GUO, ILYA ESTERLIS, AND SUBIR SACHDEV Authors Info & Affiliations

SCIENCE • 17 Aug 2023 • Vol 381, Issue 6659 • pp. 790-793 • DOI: 10.1126/science.abq6011

7,449

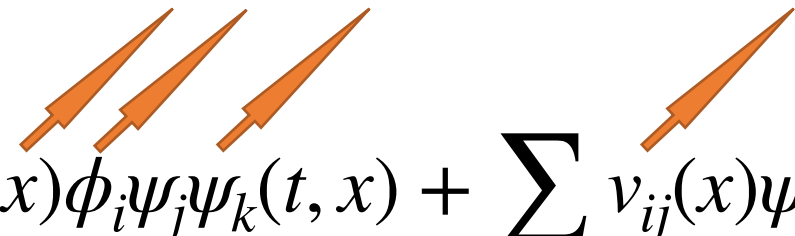
Based on SYK : all to all with Random coupling

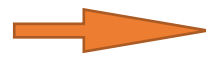
$$L = \frac{1}{2} \sum_{i=1}^N \chi_i \partial_\tau \chi_i - \frac{1}{4!} \sum_{i,j,k,l=1}^N J_{ijkl} \chi_i \chi_j \chi_k \chi_l,$$

$$\langle J_{ijkl} J_{i'j'k'l'} \rangle = J^2 \delta_{ii'} \delta_{jj'} \delta_{kk'} \delta_{ll'}$$

1. Solvable.
2. $\Gamma \sim T$
3. Maximal Chaos
- ****
4. 0-dim => No transport!

Yukawa-SYK in 2+1 dim.

$$\mathcal{H} = \mathcal{H}_\psi + \mathcal{H}_\phi + \sum_{ijk} g_{ijk}(x) \phi_i \psi_j \psi_k(t, x) + \sum_{ij} v_{ij}(x) \psi_i \psi_j$$



$$\left\{ \begin{array}{l} \langle g_{ijl}(\mathbf{r}) \rangle = 0 \\ \langle g_{ijl}^*(\mathbf{r}) g_{i'j'l'}(\mathbf{r}') \rangle = g^2 \delta_{ii'} \delta_{jj'} \delta_{ll'} \delta(\mathbf{r} - \mathbf{r}') \end{array} \right.$$

$$\rho = \rho_0 + AT$$

Resistivity linear in T at Low tem.

References

- A.A. Patel, H. Guo, I. Esterlis and S. Sachdev, Science 381 (2023) abq6011
- I. Esterlis, H. Guo, A.A. Patel and S. Sachdev, Phys. Rev. B 103 (2021) 235129
- H. Guo, A.A. Patel, I. Esterlis and S. Sachdev, Phys. Rev. B 106 (2022) 115151



Yili Wang, XianHui Ge, YK Han + SJS

1. Yili Wang, Xian-Hui Ge, and S. Sin, arXiv: 2406.11170
2. Yili Wang, Young-Kwon Han, Xian-Hui Ge, and S. Sin, arXiv: 2501.07792
3. ER=EPR and strange metal from quantum entanglement.
4. Most general strange metal model. TBA

1. How universal is the (spatially random disorder) mechanism?

Yukawa-SYK, & Vector-SYK, Model. $\phi\psi\psi$, $A_{\mu}^{ext}\psi\partial_{\mu}\psi$

2. Does it work in 3+1 ? No.

=> layered structure is important. (CuO_2 plane)

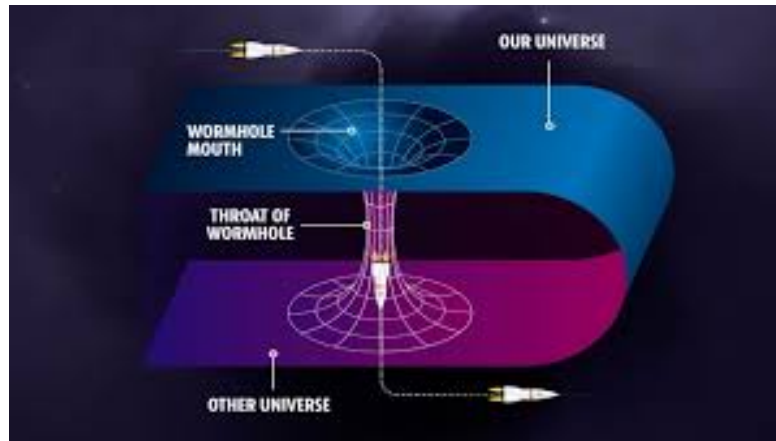
3. T^2 inverse Hall angle behavior? No!

4. The most general QFT model?

$\phi\psi\psi$, $A_{\mu}^{ext}\psi\partial_{\mu}\psi$, $\dim 5/2$ => all others are irrelevant

Our Q: why it works as a Mechanism for the Planckian dissipation

ANS: randomness brings us Wormhole, a short cut in space



Claim:

$$\langle \longleftrightarrow \rangle \quad \langle g_{ijl}^*(\mathbf{r}) g_{i'j'l'}(\mathbf{r}') \rangle = g^2 \delta_{ii'} \delta_{jj'} \delta_{ll'} \delta(\mathbf{r} - \mathbf{r}')$$

- Gravitational Wormhole
- Field theory wormhole

Quenched vs Annealed

$$W_q = \int D[g] P[g] \left[\ln \left(\int D[\Psi] e^{-S_{\text{tot}}[g, \psi, \phi]} \right) \right] = \langle \log Z \rangle_{dis}$$

$$W_a = \ln \left(\int D[\Psi] \int D[g] P[g] e^{-S_{\text{tot}}[g, \psi, \phi]} \right) = \log \langle Z \rangle_{dis}$$

Quenched case: replica trick.
Replica off diagonal case is
1/N suppressed.

For annealed case, $H_{int}(x) = g(x)^2 + g(x) \int dt \mathcal{O}(x, t) = [g(x) + \int dt \mathcal{O}(x, t)]^2 - \boxed{\iint dt_1 dt_2 \mathcal{O}(x, t_1) \mathcal{O}(x, t_2)}$

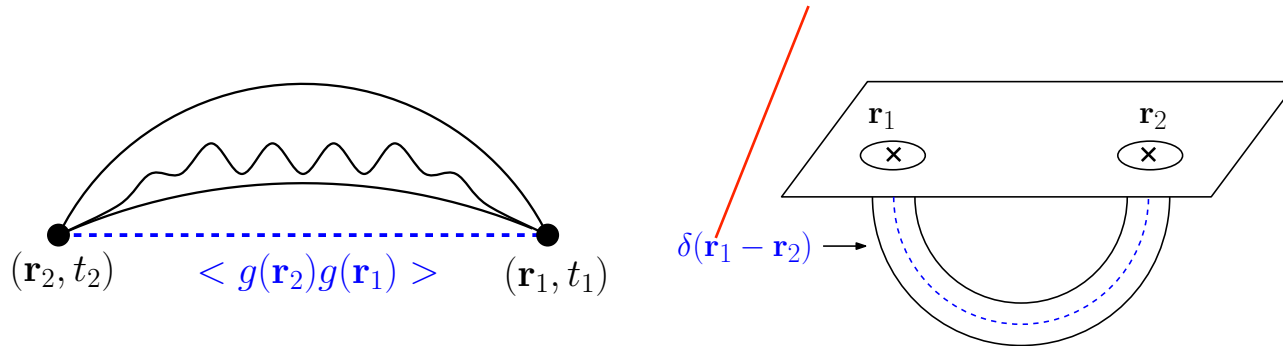
$$S_{\text{int}} \equiv -\frac{g^2}{2N^2} \int d\tau_1 d\tau_2 d^2 \mathbf{r} \sum_{a,b,c=1}^N \psi_a^\dagger \psi_b(\mathbf{r}, \tau_1) \phi_c(\mathbf{r}, \tau_1) \\ \times \psi_b^\dagger \psi_a(\mathbf{r}, \tau_2) \phi_c(\mathbf{r}, \tau_2). \quad (8)$$

Quenched => wormhole

$$-F_q = \langle \log Z \rangle_{dis} = \sum_{n=0}^{\infty} \int_{\mathbf{x}_1 \dots \mathbf{x}_n} \left\langle \left\langle \prod_{a=1}^n g_{ijk}(\mathbf{x}_a) H_{ijk}(\mathbf{x}_a) \right\rangle_{\Phi, c} \right\rangle_{dis},$$

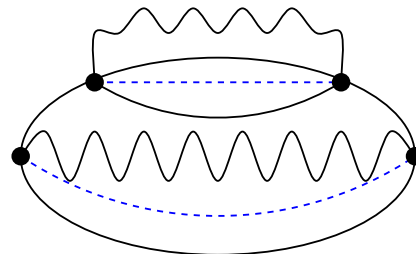
$$H_{ijk} = \phi_i \psi_j \psi_k(t, x)$$

$$n=2 : \langle g(x)g(y) \rangle \langle \phi \psi \psi(x) \phi \psi \psi(y) \rangle$$



Wormhole

n=4

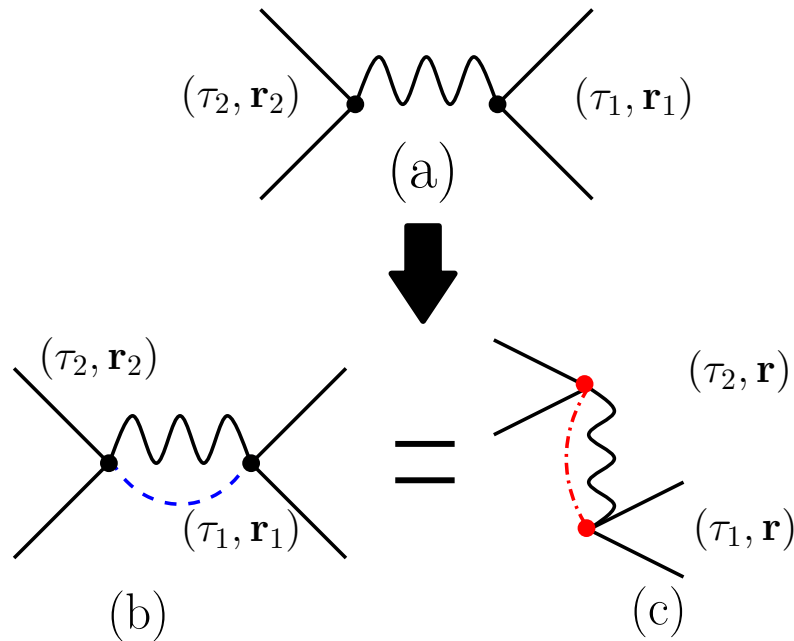


Many Wormholes

The role of the QFT-WH:

Without it, distance dependent suppression

With \Rightarrow no such suppression



$$\frac{1}{r} e^{-mx}$$

Causal
suppression!



1

In many body:
 $m \rightarrow 1/L$, $L \sim \text{MFP}$

Is SM a **weak coupling phenomena?**

Even with small coupling, achieve strongly correlated system.

Disorder => wormhole: decoherence?

1. Use equivalence of Quenched = Annealed In large N limit.

We calculated $\rho(T)$ and show it is the same as that of Patel. et.al

2. Ask what we mean by spatial randomness!

Annealed picture: vertex gives EPR

$$H_{\text{int}}(x) = g(x)^2 + g(x) \int dt \mathcal{O}(x, t) = \left[g(x) + \int dt \mathcal{O}(x, t) \right]^2 - \iint dt_1 dt_2 \mathcal{O}(x, t_1) \mathcal{O}(x, t_2)$$

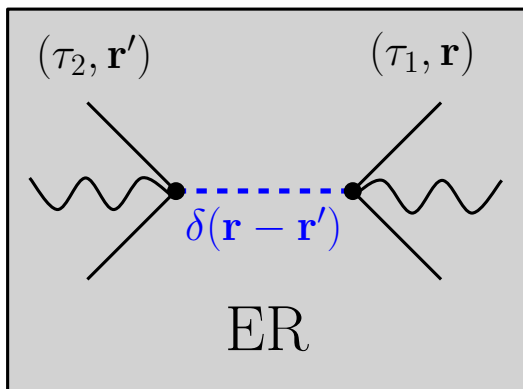
$$S_{\text{int}} = -\frac{g^2}{2N^2} \int d\tau_1 d\tau_2 d^2 \mathbf{r} \sum_{a,b,c=1}^N \psi_a^\dagger \psi_b(\mathbf{r}, \tau_1) \phi_c(\mathbf{r}, \tau_1) \times \psi_b^\dagger \psi_a(\mathbf{r}, \tau_2) \phi_c(\mathbf{r}, \tau_2). \quad (8)$$

This vertex gives Entanglement!

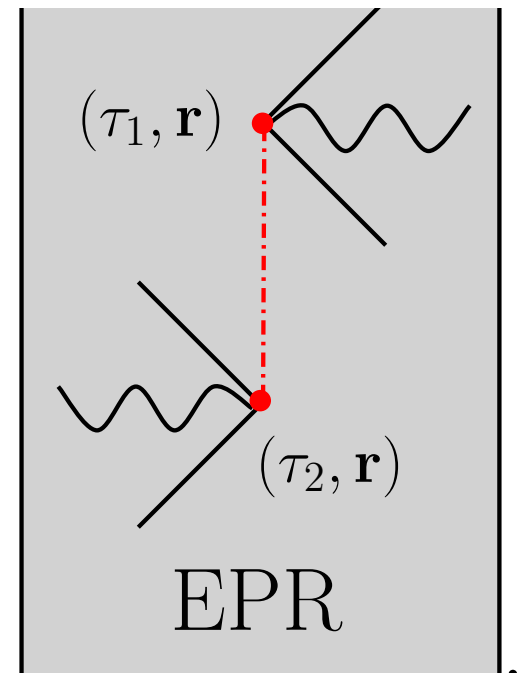
Only one p conservation for two vertex->

Long range momentum correlation

$$|\Psi\rangle = \int d\mathbf{p} \sqrt{P(\mathbf{p})} |\mathbf{p}, -\mathbf{p}\rangle, \text{ with } \int d\mathbf{p} P(\mathbf{p}) = 1.$$



In quenched ==> In annealed
ER ==> EPR



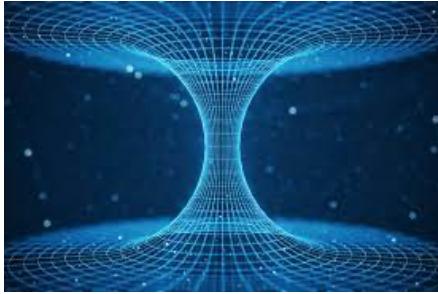
A Parallelism in information loss prob.

Quantum gravity vs Yukawa-SYK

Wormhole in QG and information loss

Hawking (1987)

deCoherence by the Wormhole :



Info. can leak away by WormHole

VS

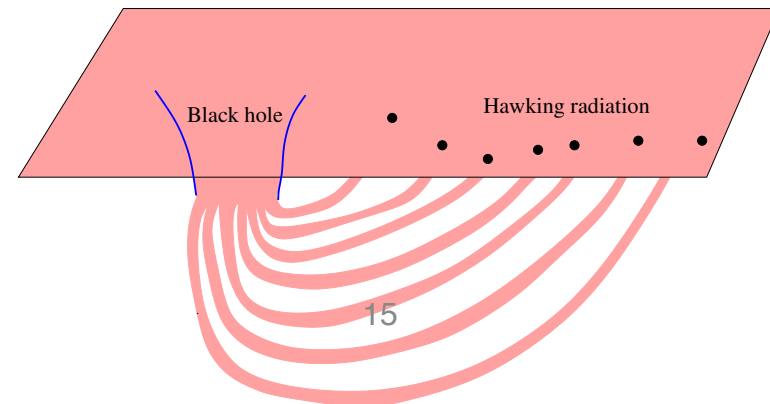
- Maldacena+Susskind :
arXiv:1306.0533

Any entangled pair is connected by a microscopic wormhole

ER=EPR



- Hawking Radiation is entangled by WH
- Not thermal, => info not lost.



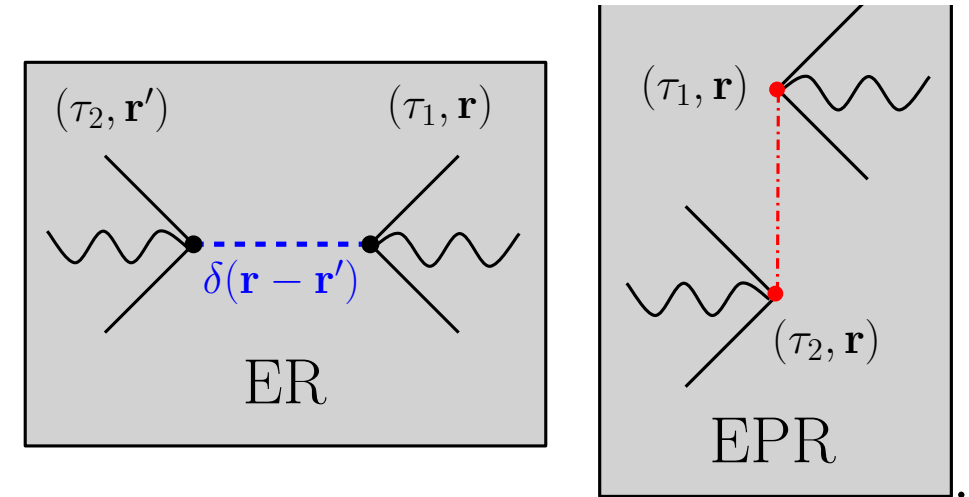
Parallelism in Yukawa-SYK and information loss

Usual SYK is **disorder averaged** hence describe a **Mixed state** !

the coupling correlation is the way to deliver the information from one vertex to the other.

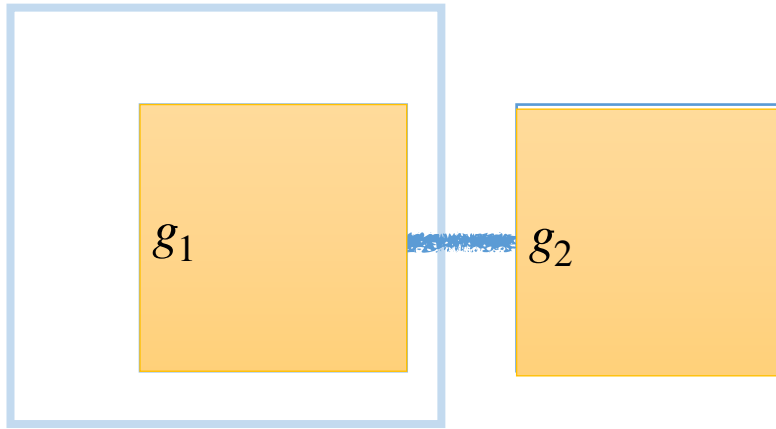
Consistent with Hawking's view :

=> SM is mixed state property not really a qm property.



- 1. Quenched \Leftrightarrow Annealed
 \Leftrightarrow **ER=EPR**
- 2. This is the **origin of the Planckian Dissipation.**
(Because this is mechanism of **long range p delivery.**)

wormhole and information loss: g vs $g(x)$?

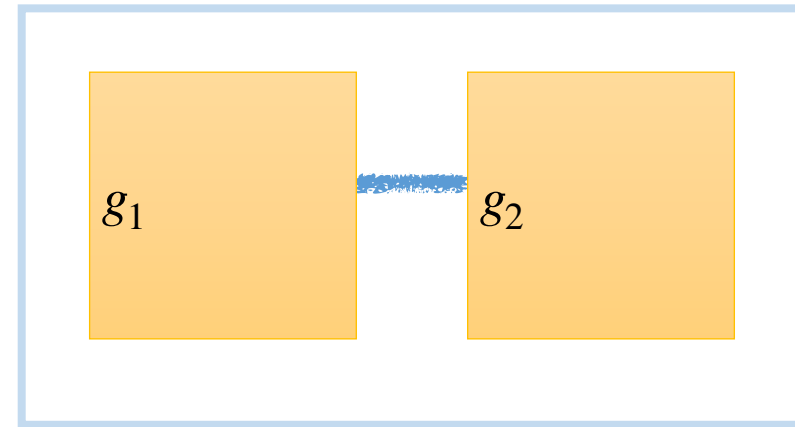


System

Open system: information leaks

sample = Each box ,

Mixed: sum of independent many ,



System include both

Closed system: No information loss

What do we mean by spatially Random $g(x)$

- 1. Spatial inhomogeneity is important.

$$Sys = \sum_i (subSys)_i, \quad \text{NOT} \quad Sys = \langle Sys_i \rangle$$

- \Leftrightarrow system is sum of many small pieces of subregions



Glue all to all fashion

Interacting \Leftrightarrow Entangled

1. 2. sample= sum of the boxes
qm: the whole is 1 system by connection.

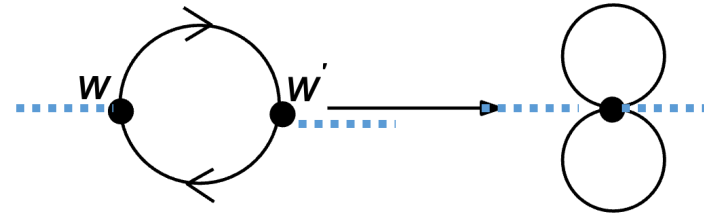
Parallel to the Raamsdonk's idea! Entanglement is a glue of spacetime.

Strange metallicity is more than Entanglement

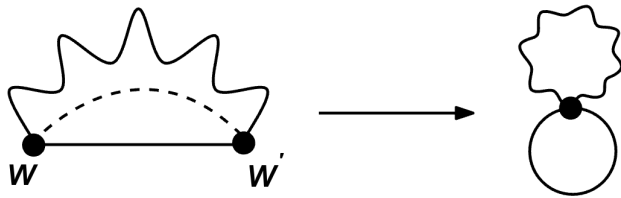
Need all 3 conditions

1. 2+1 dim!
2. Spatial dependence of g !
3. Inclusion of minimal dim operator.

$$\langle g_{ijl}^*(\mathbf{r}) g_{i'j'l'}(\mathbf{r}') \rangle = g^2 \delta_{ii'} \delta_{jj'} \delta_{ll'} \delta(\mathbf{r} - \mathbf{r}')$$



Boson self-energy



Electron self-energy

$$\Sigma_v(i\omega_n) \equiv -i \frac{\Gamma}{2} \text{sgn}(\omega_n)$$

$$\Sigma_K(i\omega_n) \equiv -i c_K \omega_n \ln \left(\frac{c_0}{\omega_n} \right)$$

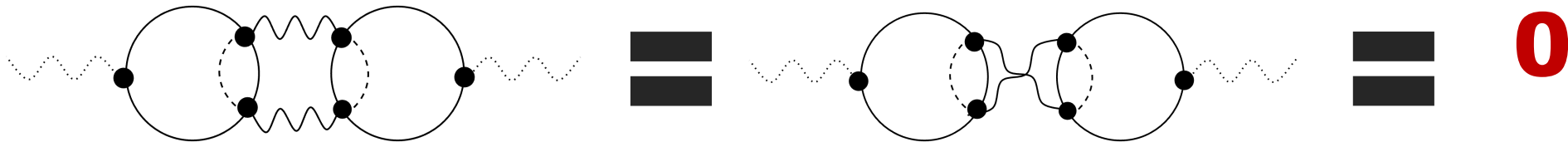
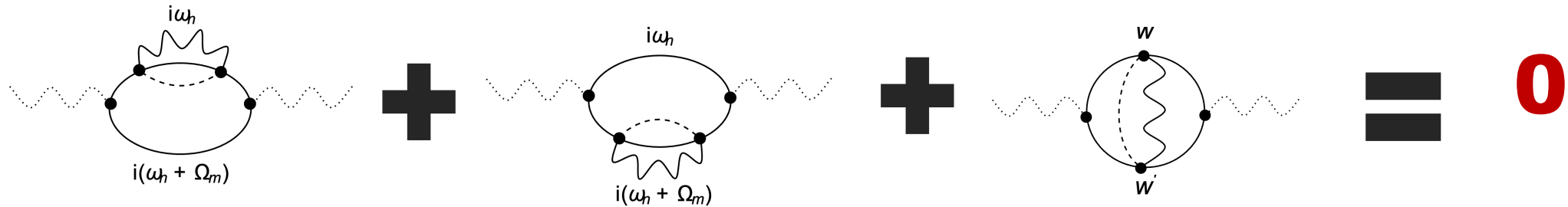
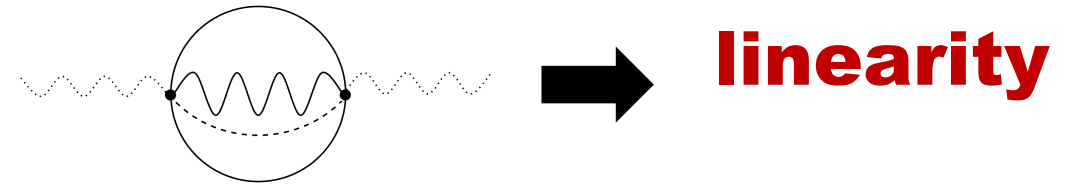
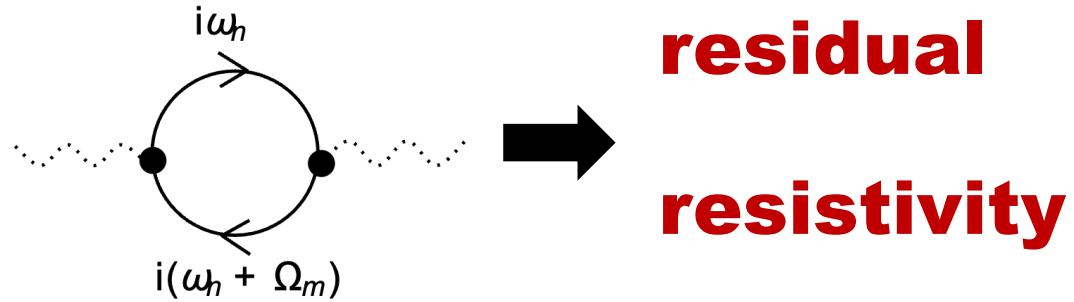
Potential disorder

Electron-boson coupling

Momentum conservation at each vertex is lost!
together with two vertex=> conserved!

$$\rho = \rho_0 + T^2 \rightarrow \rho_0 + T^1$$

spatial randomness: changing the scaling law



Summary and future work

- Spatial random coupling \Rightarrow Wormhole
 \Rightarrow EPR+ Strange transport (long range momentum delivery).
- (quenched \longleftrightarrow annealed) \Rightarrow (ER = EPR)
- Remaining Q: In what system Y-SYK is realized and How Non-disordered system can be SM?

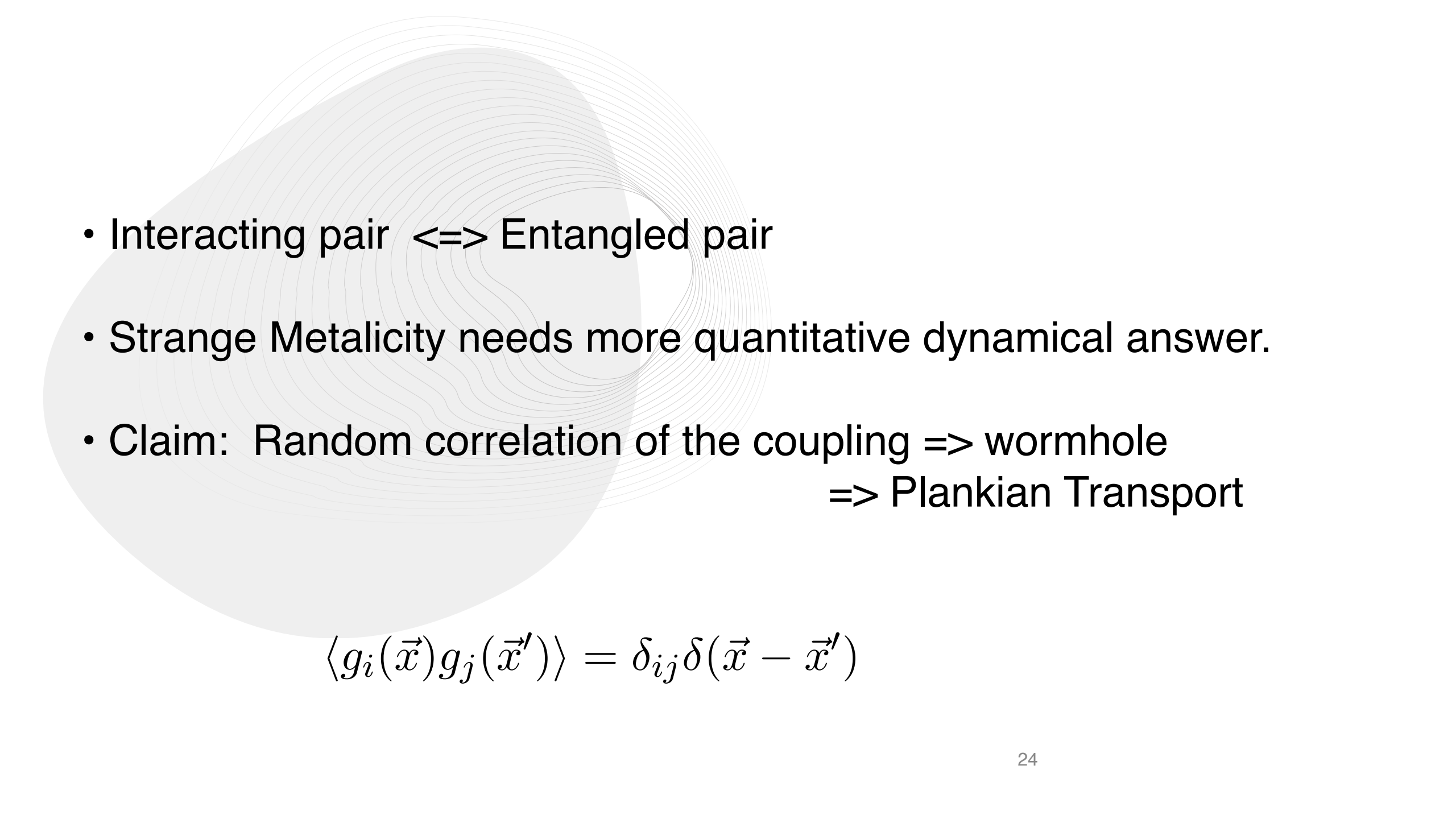
Epilogue from the movie “Dr. Strange”

- Dr. Strange came through the wormhole and shouted,
운동량 배달 왔어요!



- Momentum delivery! 运动量快递

Thank you

- 
- Interacting pair \Leftrightarrow Entangled pair
 - Strange Metalicity needs more quantitative dynamical answer.
 - Claim: Random correlation of the coupling \Rightarrow wormhole
 \Rightarrow Plankian Transport

$$\langle g_i(\vec{x}) g_j(\vec{x}') \rangle = \delta_{ij} \delta(\vec{x} - \vec{x}')$$

III. Equivalence of the Quenched and Annealed disorder.

$$W_q = \int D[g] P[g] \left[\ln \left(\int D[\Psi] e^{-S_{\text{tot}}[g, \psi, \phi]} \right) \right],$$

$$W_a = \ln \left(\int D[\Psi] \int D[g] P[g] e^{-S_{\text{tot}}[g, \psi, \phi]} \right)$$

$$W_q = W_a.$$

$$g^2 + g(x)\mathcal{O}(x) = (g + \cancel{\mathcal{O}(x)})^2 - \mathcal{O}(x)\mathcal{O}(x)$$

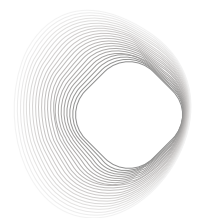
For annealed case,

one can integrate out the disorder variable $g_{ijk}(\mathbf{r})$ to get an effective theory, which can be described by replacing the interaction (1) with

Quenched case: replica trick.
Replica off diagonal case is 1/N suppressed.

$$S_{\text{int}} \equiv -\frac{g^2}{2N^2} \int d\tau_1 d\tau_2 d^2\mathbf{r} \sum_{a,b,c=1}^N \psi_a^\dagger \psi_b(\mathbf{r}, \tau_1) \phi_c(\mathbf{r}, \tau_1) \\ \times \psi_b^\dagger \psi_a(\mathbf{r}, \tau_2) \phi_c(\mathbf{r}, \tau_2). \quad (8)$$

- For $g_{ijk}(x)\phi_i\psi_j\psi_k(t, x)$, $i=1,\dots, N$, with large N , the replica index a from $g_{ijk}(x)\phi_i^a\psi_j^a\psi_k^a(t, x)$ can be dropped after $R \rightarrow 0$.
- The equivalence can be shown case by case.
We have done at the level of conductivity to discuss SM.



SYK-rised Vector Model

$$\begin{aligned}
 S = \int d\tau d^2r & \left[\sum_{i=1}^N \psi_i^\dagger(r, \tau) \left(\partial_\tau - \frac{\nabla^2}{2m} - \mu \right) \psi_i(r, \tau) \right. \\
 & + \sum_{i,j=1}^N \boxed{v_{ij}(r)} \psi_i^\dagger(r, \tau) \psi_j(r, \tau) - \frac{1}{2K^2} \sum_{l=1}^N g_{ab} a_l^a (-\partial_\tau^2 + \mathbf{q}^2) \boxed{a_l^b} \\
 & \left. + \sum_{i,j,l=1}^N \frac{\boxed{K_{ijl}(r)}}{KN} \frac{i}{m} \psi_i^\dagger \nabla_a \psi_j \mathbf{a}_l^a + \frac{1}{K^2 N^{3/2}} \sum_{i,j,s,t=1}^N \frac{\boxed{\tilde{K}_{ijst}(r)}}{2m} \mathbf{a}_s \cdot \mathbf{a}_t \psi_i^\dagger \psi_j \right]
 \end{aligned}$$

potential disorder
vector field

SYK-rised random coupling

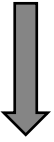
$$\begin{aligned}
 \langle K_{ijl}(r) \rangle &= 0, & \langle K_{ijl}^*(r) K_{i'j'l'}(r') \rangle &= K^2 \delta_{ii'} \delta_{jj'} \delta_{ll'} \delta(r - r') \\
 \langle \tilde{K}_{ijl}(r) \rangle &= 0, & \langle \tilde{K}_{ijl}^*(r) \tilde{K}_{i'j'l'}(r') \rangle &= \tilde{K}^2 \delta_{ii'} \delta_{jj'} \delta_{ll'} \delta(r - r') \\
 \langle v_{ij}(r) \rangle &= 0, & \langle v_{ij}(r) v_{i'j'}(r') \rangle &= \delta_{ii'} \delta_{jj'} \delta(r - r')
 \end{aligned}$$

SYK-rised Vector Model

$$G(x_1, x_2) \equiv -\frac{1}{N} \sum_i \langle \mathcal{T} \left(\psi_i(x_1) \psi_i^\dagger(x_2) \right) \rangle$$

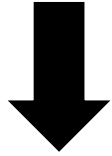
$$D^{\mu\nu}(x_1, x_2) \equiv \frac{1}{N} \sum_l \langle \mathcal{T} \left(a_l^\mu(x_1) a_l^\nu(x_2) \right) \rangle$$

Vector
model



G-Σ
action

Saddle Point Equation



Dyson's Equations

$$\begin{aligned} 0 &= \frac{\delta S}{N} \\ &= \text{Tr}(\delta \Sigma (G_*[\Sigma] - G) + \delta G (\Sigma_*[G] - \Sigma)) \\ &\quad + \frac{1}{2} \delta \Pi_{ab} (D^{ab} - D_*^{ab}[\Pi^{ab}]) + \delta D_{ab} (\Pi^{ab} - \Pi_*^{ab}[D^{ab}]) \end{aligned}$$

$$G = G_* = (-\partial_\tau - \varepsilon_k + \mu - \Sigma)^{-1}$$

$$\Sigma = \Sigma_* = v^2 G(\tau, r=0) \delta^3(r) + \frac{(k_1 + k_2)^a (k_1 + k_2)^b}{4m^2} D_{ab} G \bar{\delta} + \frac{1}{4m^2} D_{ab} D^{ab} G \bar{\delta}$$

$$D_{ab} = D_{*ab} = K^2 (-g^{ab} (-\partial_\tau^2 + \mathbf{q}^2) - K^2 \Pi^{ab})^{-1}$$

$$\Pi_{ab} = \Pi_{*ab} = -\frac{(k_1 + k_2)_a (k_1 + k_2)_b}{4m^2} G \bar{\delta} \cdot G - \frac{1}{4m^2} G D_{ab} G \bar{\delta}$$

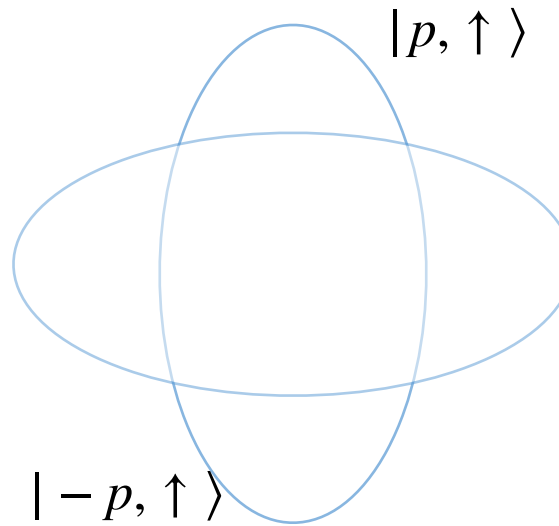
III. Equivalence of the Quenched and Annealed disorder.

- Quenched: $H_{int} = J\mathcal{O}$, should use replica trick.

$$\langle O \rangle_Q = -\frac{\delta}{\delta J} \langle \ln Z[J] \rangle_{dis} = -\lim_{R \rightarrow 0} \frac{1}{R} \frac{\delta Z^R[J]}{\delta J}$$

- Annealed :

$$\langle O \rangle_A = -\frac{\delta}{\delta J} \ln \langle Z[J] \rangle_{dis}$$



So the pair is coming from $|p, \uparrow\rangle \otimes |-p, \uparrow\rangle$ which is a triplet. So triplet superconductor is possible. While the singlet is possible only for the discrete 4 points. It is similar to the superconductor out of spineless fermions. The point of the chiral superconductor.