

# Entanglement and Symmetries in Conformal Field Theories and Holography

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# People behind the research



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## Charged Moments in $W_3$ Higher Spin Holography

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## Symmetry-resolved entanglement in $AdS_3/CFT_2$ coupled to $U(1)$ Chern-Simons theory



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## Symmetry-resolved entanglement for excited states and two entangling intervals in $AdS_3/CFT_2$

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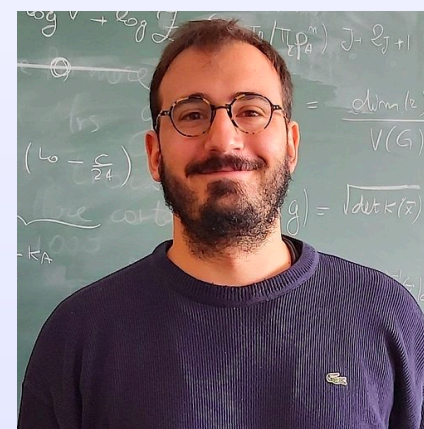
Suting Zhao,<sup>a,1</sup>, Christian Northe,<sup>a,2</sup>, Konstantin Weisenberger,<sup>a,3</sup> René Meyer,<sup>a,4</sup>

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## On the boundary conformal field theory approach to symmetry-resolved entanglement

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# Quantum Information and AdS/CFT

## Building up spacetime with quantum entanglement

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### Abstract

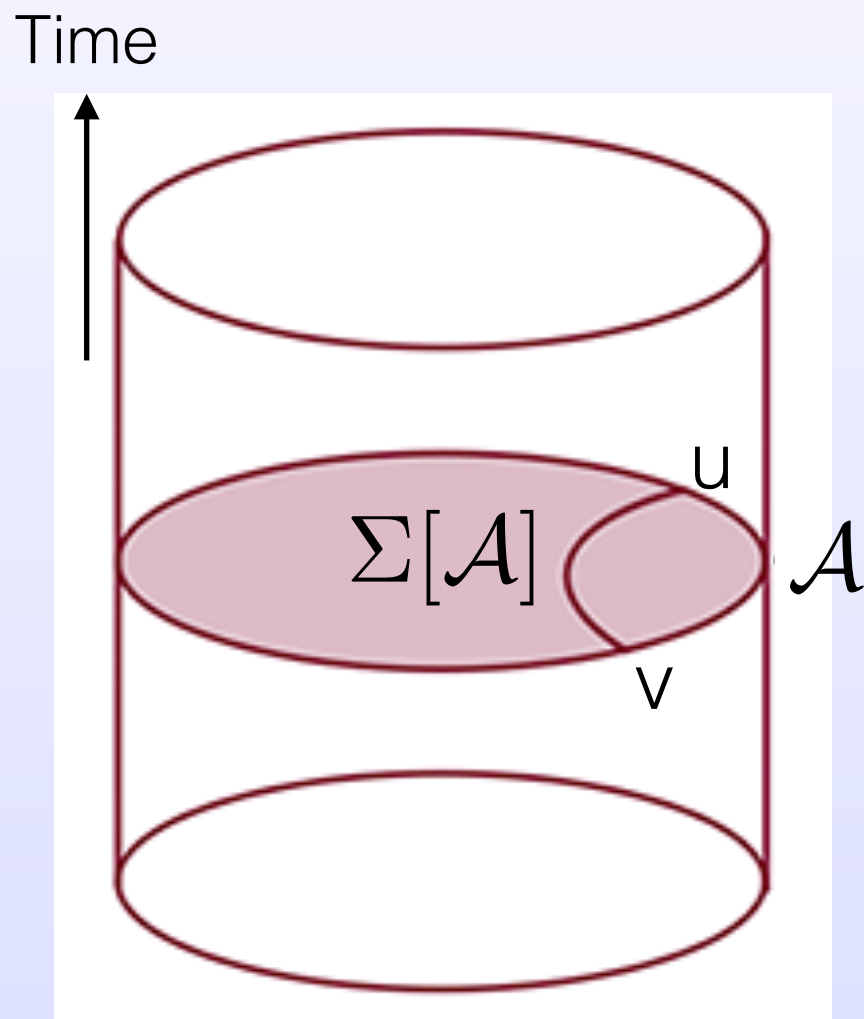
In this essay, we argue that the emergence of classically connected spacetimes is intimately related to the quantum entanglement of degrees of freedom in a non-perturbative description of quantum gravity. Disentangling the degrees of freedom associated with two regions of spacetime results in these regions pulling apart and pinching off from each other in a way that can be quantified by standard measures of entanglement.

[arXiv:1005.3035 \[hep-th\]](https://arxiv.org/abs/1005.3035)

Fine-grained information on bulk (quantum) gravity from symmetry resolved entanglement?

# Entanglement Entropy in AdS3/CFT2

Minimal length curve (geodesic)  
anchored at the ends of the entangling interval



$$S(\mathcal{A}) = \frac{\text{Length}(\Sigma[\mathcal{A}])}{4G_3}$$

$$c = \frac{3L}{2G_3} \gg 1$$

L... Curvature Radius of AdS3 space-time

For ground state of 2D CFTs:

$$S(\mathcal{A}) = \frac{c}{3} \log \frac{|v-u|}{\epsilon}$$

$\epsilon$  ... Short Distance Cutoff

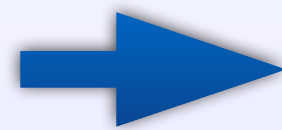
# Symmetry Resolved Entanglement

Entanglement entropy in each charge sector, e.g. U(1)

charge operator  $Q$

$$Q = Q_A \oplus Q_B.$$

eigenstate of  $Q$ :  $[\rho, Q] = 0.$



$$[\rho_A, Q_A] = 0.$$

Block decomposition:

$$\rho_A = \bigoplus_q \rho_A(q)$$

Symmetry Resolved Renyi and Entanglement Entropy:

$$S_n(q) = \frac{1}{1-n} \log \text{Tr} \left( \frac{\rho_A(q)}{P_A(q)} \right)^n$$

$$P_A(q) = \frac{\text{Tr} \rho_A(q)}{\text{Tr} \rho_A} = \text{Tr} \rho_A(q)$$

$$S_1(q) = \lim_{n \rightarrow 1} S_n(q) = -\text{Tr} \left( \frac{\rho_A(q)}{P_A(q)} \log \frac{\rho_A(q)}{P_A(q)} \right)$$

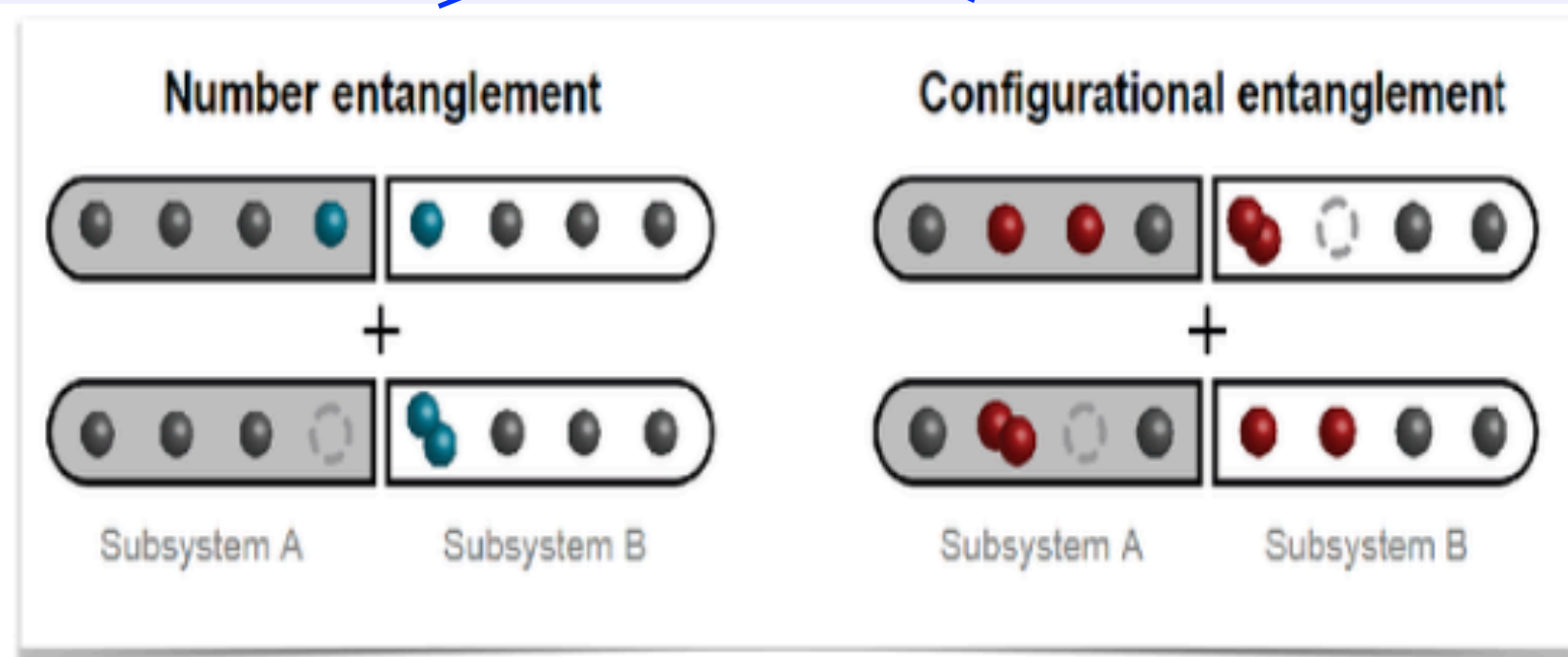


# Symmetry Resolved Entanglement

Entanglement entropy in each charge sector

$$S_1 = \sum_q P_{\mathcal{A}}(q) S_1(q) - \sum_q P_{\mathcal{A}}(q) \log P_{\mathcal{A}}(q)$$

$$P_{\mathcal{A}}(q) = \frac{\text{Tr} \rho_{\mathcal{A}}(q)}{\text{Tr} \rho_{\mathcal{A}}} = \text{Tr} \rho_{\mathcal{A}}(q)$$



A. Lukin, M. Rispoli, R. Schittko, M. E. Tai, A. M. Kaufman, S. Choi, V. Khemani, J. Leonard, and M. Greiner, Probing entanglement in a many-body localized system, *Science* 364, 6437 (2019).

# Example: 2-Qubit system

$$|EPR\rangle = \frac{1}{\sqrt{2}} (|0\rangle_A \times |1\rangle_B + |1\rangle_A \times |0\rangle_B)$$

$$\rho_A = \frac{1}{2} (|0\rangle_A \langle 0|_A + |1\rangle_A \langle 1|_A)$$

Conserved Charge: Occupation number (0 or 1)

$$\rho_A(0) = |0\rangle_A \langle 0|_A \quad \rho_A(1) = |1\rangle_A \langle 1|_A$$

$$P_A(0) = P_A(1) = \frac{1}{2}$$

$$S_1(0) = S_1(1) = 0$$

Only number entanglement (classical charge correlations)  
No configurational entanglement (quantum correlations)

# U(1) Kac-Moody CFTs

Holographic CFT with U(1) conserved current:  $c = \frac{3L}{2G_3} \gg 1$

$$J(z) = \sum_{n=-\infty}^{\infty} \frac{J_n}{z^{n+1}}$$

Conserved current

$$T(z) = \sum_{n=-\infty}^{\infty} \frac{L_n}{z^{n+2}}$$

Energy-Momentum-Tensor

U(1)<sub>k</sub> Kac-Moody algebra at level k

$$[J_n, J_m] = \frac{1}{2}nk\delta_{m+n}$$

$$[L_n, J_m] = -mJ_{n+m}$$

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2012.11274)

K. Weisenberger, S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2108.09210)



# AdS<sub>3</sub> dual to U(1)<sub>k</sub> Kac-Moody CFT

3D Einstein-Hilbert gravity

$$S_g = \frac{1}{16\pi G_3} \int d^3x \sqrt{g} \left( R + \frac{2}{L^2} \right)$$

L... Curvature Radius of AdS<sub>3</sub> space-time

To suppress quantum gravity effects:  $c = \frac{3L}{2G_3} \gg 1$

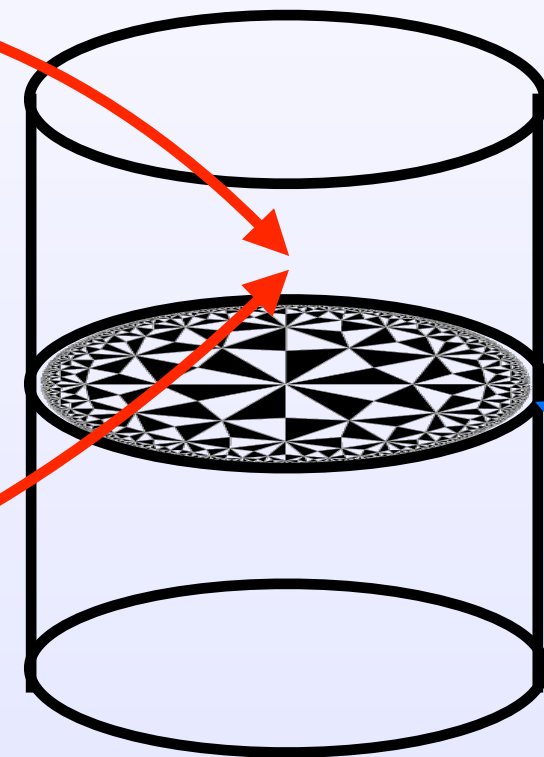
U(1)<sub>k</sub> Chern-Simons theory

$$S_{CS} = \frac{ik}{8\pi} \int A \wedge dA$$

k... Chern-Simons level

Asymptotic symmetry analysis:

Boundary theory has U(1)<sub>k</sub> Kac-Moody symmetry

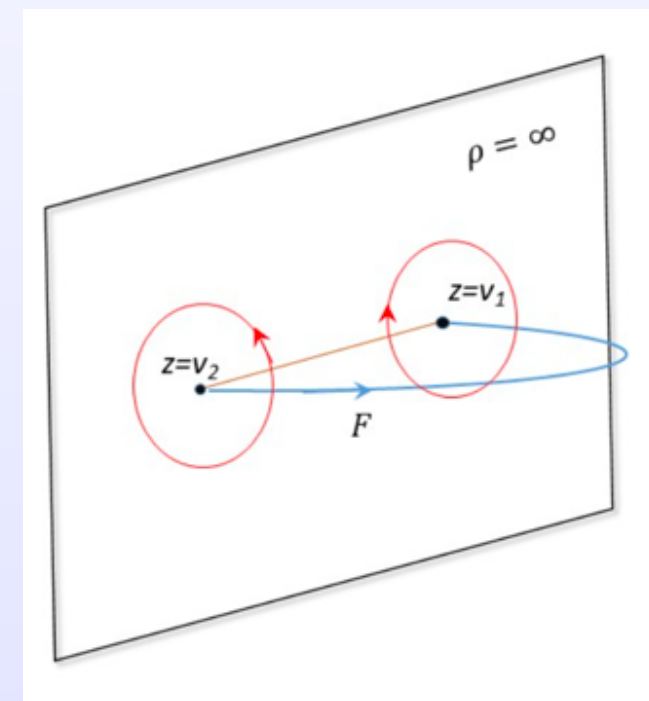
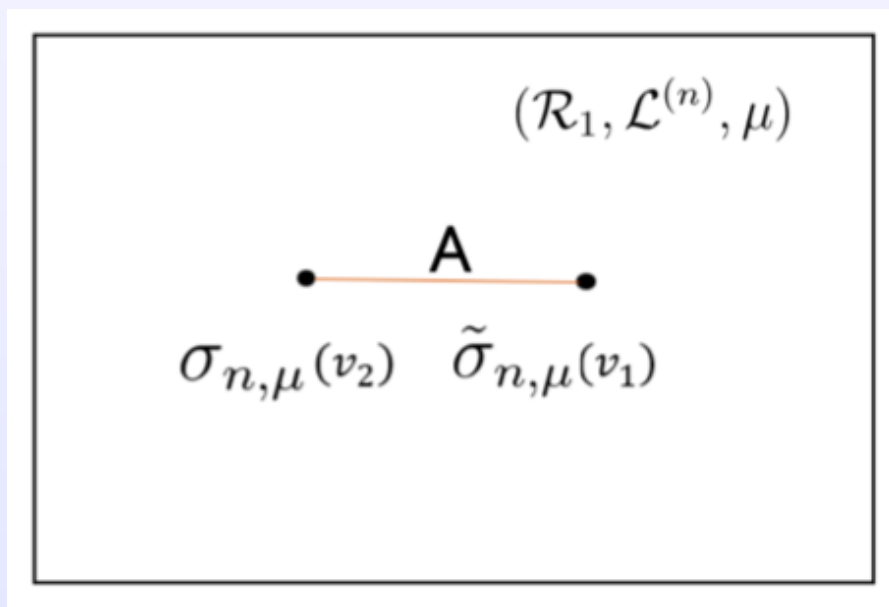


# Single Interval SREE in AdS<sub>3</sub>/CFT<sub>2</sub>

Charged twist operator induces flux:

Sela, Goldstein PRL 2018

Wilson line following the Ryu-Takayanagi geodesic



CFT<sub>2</sub> result matches AdS<sub>3</sub> result for  $c \gg 1$

$$S_1(q) = \frac{c}{6} \ell - \frac{1}{2} \log \left( \frac{k\ell}{2\pi} \right) \quad \text{with} \quad \ell = 2 \log \frac{|v_1 - v_2|}{\epsilon}$$

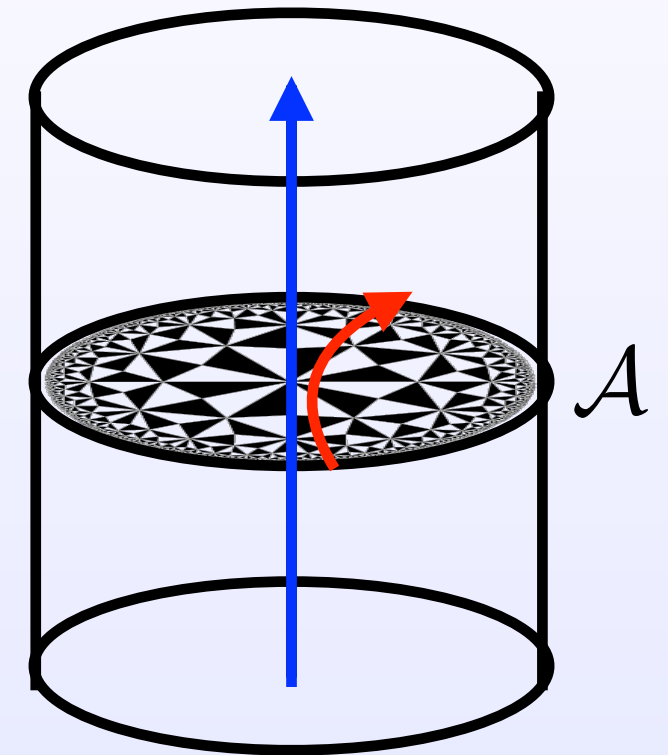
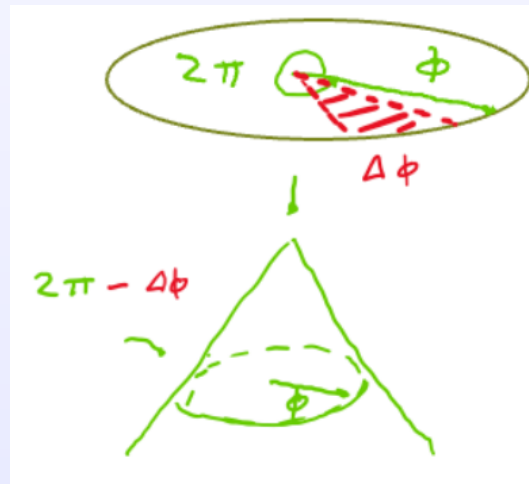
Equipartition of entanglement!

Suting Zhao, Christian Northe, RM, JHEP 2021 (arXiv: 2012.11274),

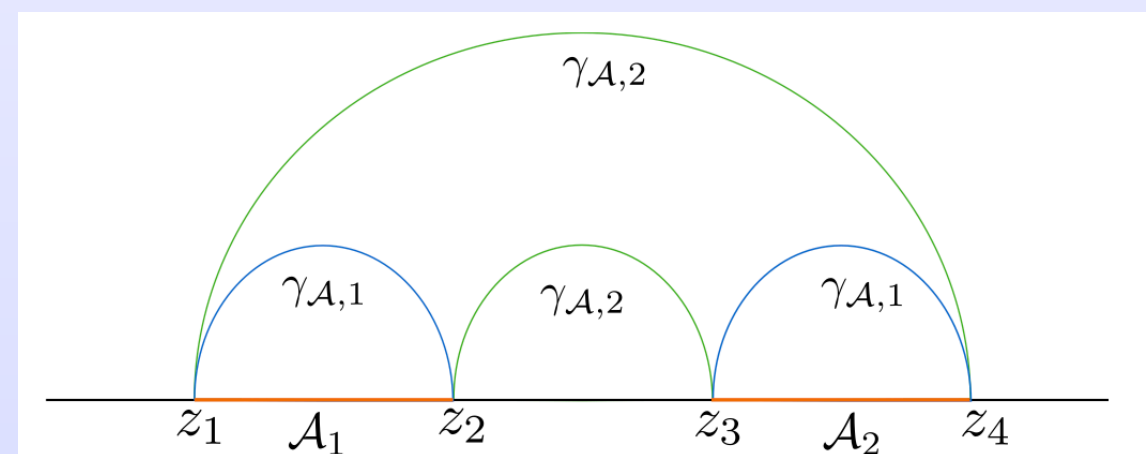
Belin et.al. 2013, Xavier et.al.. 2019

# Further Checks

- Single interval with uncharged/charged heavy primary insertions



- Two intervals in the ground state



S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2012.11274)

K. Weisenberger, S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2108.09210)

# Breakdown of Equipartition

SL(3,R) Higher Spin Gravity in 3D:  $W_3$  symmetric CFT  
Energy-momentum tensor plus Spin 3 current

$$T(z)W(w) = \frac{3W(w)}{(z-w)^2} + \frac{\partial W(w)}{z-w} + \dots$$

Charged moments for a single interval:

Topological black hole grand canonical partition function

Perturbative result to quartic order in  $\mu$

$$\log \text{Tr} \left( e^{-2\pi n \mathcal{H} + 2\pi i \mu Q_{\mathcal{A}}} \right) = \frac{c\ell}{6n} \left( -\frac{1}{3} \frac{\mu^2}{n^4} + \frac{10}{27} \frac{\mu^4}{n^8} + \dots \right)$$

Fourier transformation and taking the replica limit  
yields breakdown of equipartition in SREE at large  $c$ .

# Boundary Conformal Field Theory

## Subleading corrections to (SR)EE: Boundary Conditions

- Definition of entanglement entropy requires bipartition

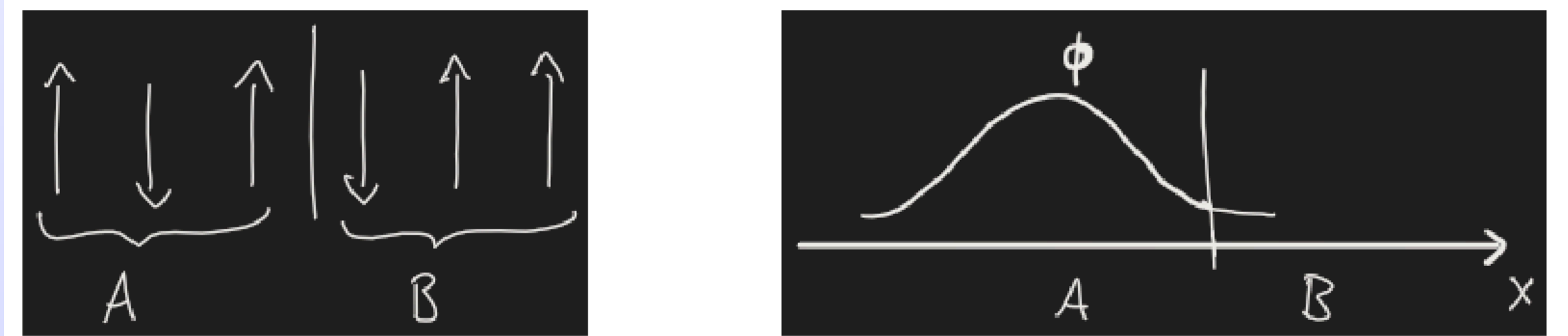
(Ohmori, Tachikawa, 2015)

$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$$

Quantum fields are distributions

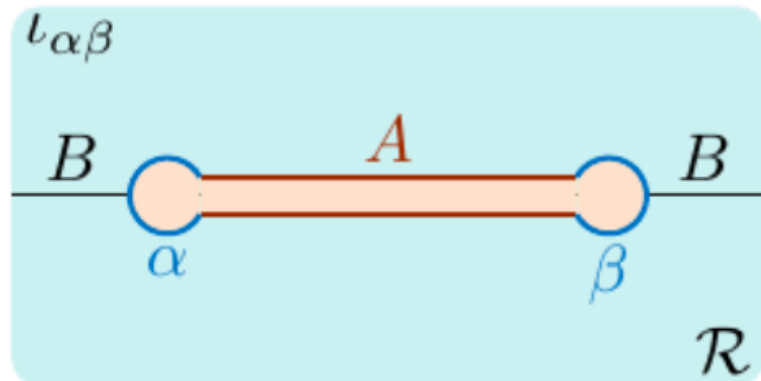
Must be smeared against test functions

Naive bipartition is insufficient



# Boundary Conformal Field Theory

## Solution: Boundary Conditions



Choose boundary conditions which preserve conformal symmetry and  $U(1)$

$$T = \bar{T}|_{bdy} \quad J = \pm \bar{J}|_{bdy}$$

Provide local set of conserved charges: Chiral algebra

$$K_A = \int_A \frac{(2A)^2 - x^2}{4A} (T(x) + \bar{T}(x)) \quad Q_A = \int_A (J(x) \mp \bar{J}(x))$$

Their eigenstates furnish the Hilbert space  $\mathcal{H}_{A,\alpha\beta}$ , i.e. the **entanglement spectrum** associated with the region  $A$ .

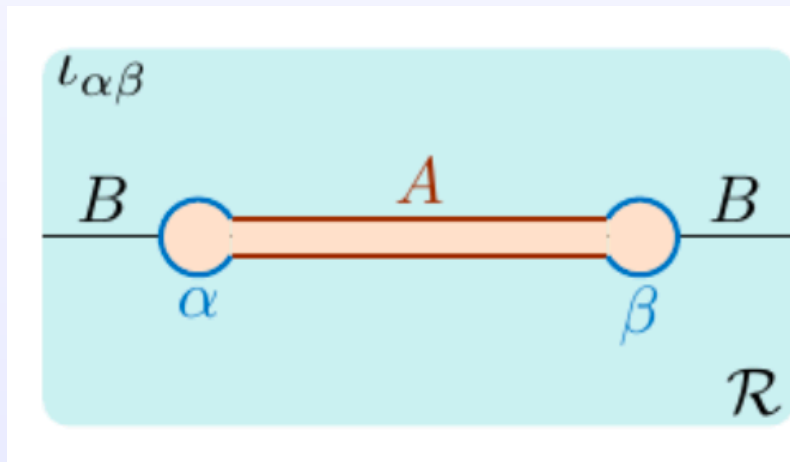
Mathematically, this is encoded in a **factorization map**

$$\iota_{\alpha\beta} : \mathcal{H} \rightarrow \mathcal{H}_{A,\alpha\beta} \otimes \mathcal{H}_{B,\alpha\beta}$$



# Example: Compact Free Boson

Two conserved currents:  $J^\mu = \partial^\mu \varphi$        $\tilde{J}^\mu = \epsilon^{\mu\nu} \partial_\nu \varphi$



Choose boundary conditions which preserve conformal symmetry and  $U(1)$

$$T = \bar{T}|_{bdy}$$

$$J = \pm \bar{J}|_{bdy}$$

Dirichlet

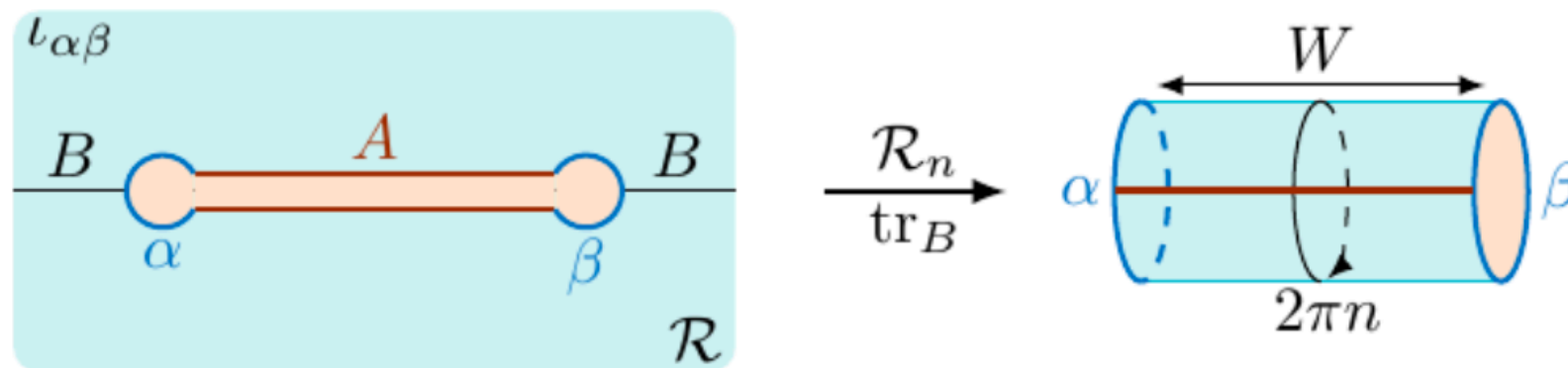
Neumann

- After cutting: NN and DD boundary conditions preserve conservation law for either  $w$  or  $m$
- ND and DN boundary conditions break both conservation laws, remaining  $\mathbb{Z}_2$  symmetry

# Boundary Conformal Field Theory

## Solution: Boundary Conditions

- Replicate geometry  $n$  times  $\rightarrow \mathcal{R}_n$
- trace over  $B$
- map onto **annulus** of width  $W = 2 \log A/\epsilon$



Focus on vacuum of CFT,  $\rho = |0\rangle\langle 0|$ . **Reduced density matrix**

$$\rho_A = \frac{q^{L_0 - c/24}}{Z_{\alpha\beta}(q)},$$

where  $Z_{\alpha\beta}(q) = \text{tr}_A q^{L_0 - c/24}$  and  $q = e^{-2\pi^2/W}$

# U(1) Resolution and BCFT

## Direct calculation of charged partition functions

Boundary conditions are  $J = \pm \bar{J}|_{bdy}$  and  $Z_{\pm}(q) = \sum_{Q \in \sigma_{\pm}} \chi_Q(q)$

$$\chi_Q(q) = \frac{q^{Q^2}}{\eta(q)} \quad \eta(q) = q^{1/24} \prod_{k=1}^{\infty} (1 - q^k)$$

These correspond to Fock spaces of  $U(1)$  charge  $Q$

$$a_{-n_1} \dots a_{-n_2} a_{-n_1} |Q\rangle$$

Charge-projected partition functions select one representation in  $Z_{\pm}$

$$\mathcal{Z}_n(Q) = \text{tr}_A[\Pi_Q \rho_A^n] = \frac{\chi_Q(q^n)}{(Z_{\pm}(q))^n} \quad p_Q = \frac{\chi_Q(q)}{Z_{\pm}(q)}$$

It vanishes when  $Q$  does not appear in  $Z_{\pm}$

# U(1) Resolution and BCFT

Equipartition to all orders in the UV cutoff expansion

Charge dependence cancels out to all orders in symmetry-resolved entropies

$$S_n(Q) = \frac{1}{1-n} \log \left[ \frac{q^{nQ^2} \eta^n(q)}{(q^{Q^2})^n \eta(q^n)} \right]$$

All **higher orders** are easily found ( $W = 2 \log A/\epsilon$ )

$$S_n(Q) = \underbrace{\frac{W}{12} \frac{n+1}{n} - \frac{1}{2} \log \left[ \frac{W}{\pi} \right] + \frac{1}{2} \frac{\log n}{1-n}}_{\text{Known terms}} + \frac{1}{1-n} \sum_{k=1}^{\infty} \log \left[ \frac{(1 - e^{-2Wk})^n}{1 - e^{-2Wk/n}} \right]$$

These expressions are reasonable only when  $Q$  appears in  $Z_{\pm}$ , i.e. in  $\mathcal{H}_A \rightsquigarrow$

**Determined by boundary conditions**

# Conclusions & Outlook

- New insight into AdS/CFT from quantum information, e.g. Entanglement entropy, **Symmetry resolved entanglement**
- **New tests of AdS<sub>3</sub>/CFT<sub>2</sub> plus Chern-Simons**
- **First example of breakdown of equipartition at large c**
- **BCFT approach yields exact results due to symmetries, U(1) equipartition to all orders for free bosons**
  
- SREE Analysis of further instances of AdS/CFT **WIP**
- Symmetry resolving other quantum information measures **WIP**
- Higher dimensions? **P. Bueno, P.A. Cano, A. Murcia, A.R. Sanchez, PRL 2022**
- Other symmetry groups? **Y. Kusuki et. al., arXiv:2309.03287**
- Implications on bulk entanglement? **WIP w. T. Kögel, J. Xu, W. Weng, G. Di Giulio**
- Timelike (Pseudo) SREE **WIP w. T. Kögel, C. Englert**
- Entanglement in non-hermitean systems **WIP w. Z. Chen, Z.-Y. Xian**
- Field theory and CondMat Applications?  
**C. Northe, Phys. Rev. Lett., arXiv:2303.07724**

