

# Gravitational form factors in holographic QCD

Yang Li

University of Science & Technology of China, Hefei

In collaboration with:

X. Cao, B. Gurjar, C. Mondal, Q. Wang and J.P. Vary

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# Gravitational form factors

## ■ Big puzzles of the strong force within hadrons

- Origin of confinement
- Origin of >99% nucleon mass
- Origin of the nucleon spin
- .....

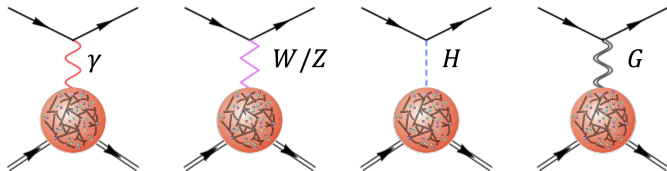


Gross and Klempt (eds.), *50 Years of quantum chromodynamics*

- ## ■ Hadronic energy-momentum tensor encodes the energy, spin and stress distributions within hadrons

[Kobzarev:1962wt, Pagels:1966zza]

$$\langle p', s' | T^{\mu\nu}(0) | p, s \rangle = \frac{1}{M} \bar{u}_{s'}(p') \left[ P^\mu P^\nu A(q^2) + \frac{1}{2} i P^{\{\mu} \sigma^{\nu\} \rho} q_\rho J(q^2) + \frac{1}{4} (q^\mu q^\nu - g^{\mu\nu} q^2) D(q^2) \right] u_s(p)$$



$$\langle \Psi | T^{\alpha\beta}(x) | \Psi \rangle = \langle \mathcal{E} \mathcal{U}^\alpha \mathcal{U}^\beta - \mathcal{P} \Delta^{\alpha\beta} + \frac{1}{2} \partial_\rho (\mathcal{U}^{\{\alpha} \mathcal{S}^{\beta\}} \rho) + \Pi^{\alpha\beta} - g^{\alpha\beta} \Lambda \rangle_\Psi$$

where,  $\langle \mathcal{O}(x) \rangle_\Psi = \int d^3z \bar{\Psi}(z) \mathcal{O}(x-z) \Psi(z)$  is the convolution with the hadronic wavepacket  $\Psi(x)$ .

- Gravitational form factors are F.T. of physical densities

energy density:  $\mathcal{E}(x) = M \int \frac{d^3q}{(2\pi)^3} e^{iq \cdot x} \left\{ \left( 1 - \frac{q^2}{4M^2} \right) A(q^2) + \frac{q^2}{4M^2} [2J(q^2) - D(q^2)] \right\},$

spin density:  $\mathcal{S}^{\alpha\beta}(x) = \int \frac{d^3q}{(2\pi)^3} e^{iq \cdot x} \left\{ i\sigma^{\alpha\beta} \sqrt{1 - \frac{q^2}{4M^2}} - \frac{U^{[\alpha} q^{\beta]}}{2M} \right\} J(q^2),$

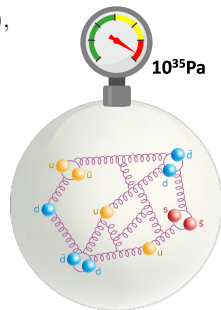
pressure:  $\mathcal{P}(x) = \frac{1}{6M} \int \frac{d^3q}{(2\pi)^3} e^{iq \cdot x} q^2 D(q^2),$

shear:  $\Pi^{\alpha\beta}(x) = \frac{1}{4M} \int \frac{d^3q}{(2\pi)^3} e^{iq \cdot x} \left( q^\alpha q^\beta - \frac{q^2}{3} \Delta^{\alpha\beta} \right) D(q^2),$

cosmological constant:  $\Lambda = -M^2 \int \frac{d^3q}{(2\pi)^3} e^{iq \cdot x} \bar{c}(q^2)$

- Factorization of the dependence on hadron momentum:

$\vec{P} = 0$  - Sachs distribution,  $P_z \rightarrow \infty$  - light-front distribution



# Mechanical properties of hadrons

- Mechanical equilibrium  $\rightarrow$  mechanical stability

[Polyakov:2018zvc]

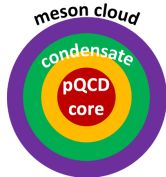
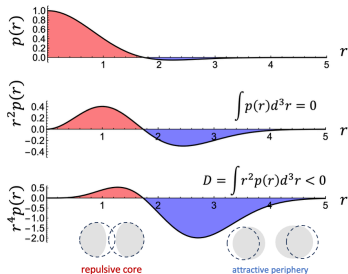
von Laue condition:  $\int d^3r \mathcal{P}(r) = 0, \quad \rightarrow \quad D = \int d^3r r^2 \mathcal{P}(r) \stackrel{???}{<} 0$

- Trace anomaly:  $\Theta \equiv T^\mu_\mu = \frac{\beta(g_s)}{2g_s} G^{\mu\nu a} G_{\mu\nu}^a + O(m_q)$

$$\langle p', s' | T^\mu_\mu | p, s \rangle = \bar{u}_{s'}(p') \left\{ A(Q^2) + \frac{Q^2}{4M^2} [A(Q^2) - 2J(Q^2) + 3D(Q^2)] \right\}$$

- Analogy to superconductor, and dark energy

[Teryaev:2016edw, Ji:2021mtz, Liu:2023cse]



$-6A'(0)$   
valence

$-6E'(0) = r_A^2 - \frac{3}{2} \lambda_c^2 D$   
condensate

$-6\Theta'(0) = r_A^2 - \frac{9}{2} \lambda_c^2 D$   
glueball,  
meson cloud



PARTICLE AND NUCLEAR | RESEARCH UPDATE

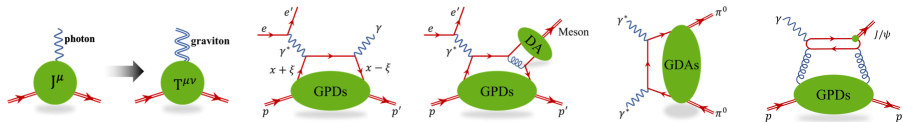
Charmonium's onion-like structure is revealed by new calculations

05 Jul 2024



# The last global unknown

[Reviews: Polyakov:2018zvc, Burkert:2023wzr]



- Ji's sum rules: one graviton  $\approx$  two photons  $\rightarrow$  generalized parton distributions (GPDs)

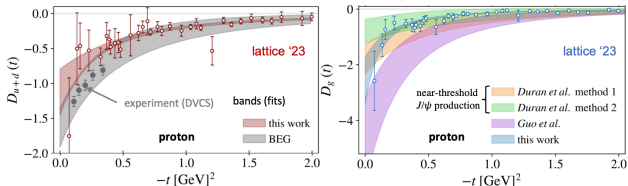
$$\int_{-1}^1 dx x H^{q,g}(x, \xi, t) = A^{q,g}(t) + \xi^2 D^{q,g}(t)$$

[Ji:1996nm, Polyakov:2002yz]

- Deeply virtual Compton scattering & deeply vector meson production [Burkert:2018bqq, Burkert:2021ith]
- Di-photon pair production [Kumano:2017lhr]
- Near threshold vector meson production [Kharzeev:2021qkd, Duran:2022xag]
- Large uncertainties from both the theory and experiments  $\rightarrow$  Electron-Ion Colliders

electromagnetic	$G_E(0) = Q = 1.602176487(40) \times 10^{-19} \text{C}$ $G_M(0) = \mu = 2.792847356(23) \mu_N$
weak	$G_A(0) = g_A = 1.2694(28)$ $G_P(0) = g_p = 8.06(55)$
gravitational	$m_A(0) = m = 938.272013(23) \text{ MeV}/c^2$ $J(0) = J = \frac{1}{2}$ $D(0) = D = ?$

[Lattice '23: Hackett:2023nkr; DR: Broniowski:2024mpw, Cao:2024zlf, Cao:2025dkv]

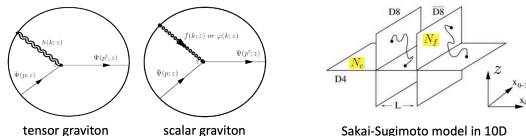


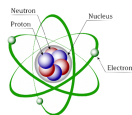


The Answer to the Great Question Of Life, the Universe and Everything is 42!



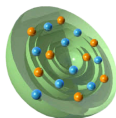
- AdS/QCD: semiclassical field theory in 5D anti-de Sitter space w./ input from QCD
  - Introduce a dilaton field  $\Phi(z)$  to break the conformal symmetry in IR: hard-wall, soft-wall models
  - Good agreement with QCD within 10%-15% level [Erlach:2005qh, Karch:2006pv]
- GFFs in AdS/QCD: graviton-hadron scattering in 5D [Abidin:2008hn, Abidin:2008ku, Abidin:2009hr]
  - Problem: GFF  $D(Q^2)$  inaccessible, since graviton only couples to the traceless part of  $T^{\mu\nu}$
  - Solutions beyond AdS/QCD: scalar graviton, Sakai-Sugimoto model [Mamo:2019mka, Fujita:2022jus]





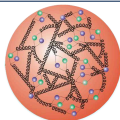
Bohr Model  

$$\left(\frac{\vec{p}^2}{2m} - \frac{Z\alpha}{r}\right)\psi = E\psi$$

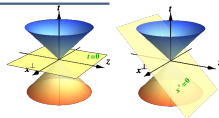


Shell Model  

$$\left(\frac{\vec{p}^2}{2m} + \frac{m\omega^2}{2}r^2\right)\psi = E\psi$$



?



$$x^\pm = t \pm z/c$$

$$H_{\text{QCD}}|\psi\rangle = M^2|\psi\rangle \quad \text{semiclassical appr.} \Rightarrow \left[ -\nabla_{\zeta_\perp}^2 + V_{q\bar{q}}^{\text{eff}}(\vec{\zeta}_\perp) \right] \varphi_h(\vec{\zeta}_\perp) = M_h^2 \varphi_h(\vec{\zeta}_\perp)$$

where  $\vec{\zeta}_\perp = \sqrt{x(1-x)}\vec{r}_\perp$



semiclassical LFQCD

↔

semiclassical field theory in AdS<sub>5</sub>

$$\zeta_\perp = \sqrt{x(1-x)}r_\perp$$

↔

fifth coordinate  $z$ ,

LF amplitude

↔

string amplitude

confining potential  $V_{q\bar{q}}^{\text{eff}}$

↔

dilation field  $\Phi$

$$L^2 - (J-2)^2$$

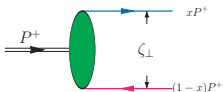
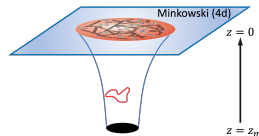
↔

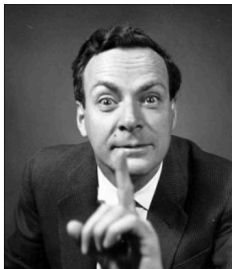
$$(\mu R)^2$$

form factors

↔

form factors





If you can't solve a problem in three different ways, you don't really understand it.

— R. P. Feynman

# Pion charge form factor $F_\pi$ in holographic QCD

- Drell-Yan-West formula in LFQCD:

[Brodsky:2007hb]

$$F_\pi(Q^2) = \int \zeta_\perp d\zeta_\perp |\varphi_\pi(\zeta_\perp)|^2 \zeta_\perp Q K_1(\zeta_\perp Q) + \text{higher Fock sector contributions}$$

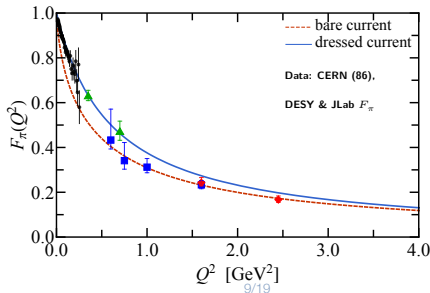
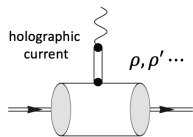
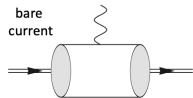
- Electromagnetic coupling in AdS<sub>5</sub>:

[Grigoryan:2007wn, Abidin:2009hr]

$$S_{\text{int}} = e_5 \int d^5x \sqrt{g} g^{NM} \Phi^*(x) i \overleftrightarrow{\nabla}_N \Phi(x) A_M(x) \Rightarrow F_\pi(Q^2) = \int z dz |\varphi_\pi(z)|^2 V(Q^2, z)$$

where,  $V(Q^2, z)$  is the bulk-to-boundary propagator of the 5D EM field  $A_N(x)$ . In soft-wall model,

$$V(Q^2, z) = \Gamma(1 + \frac{Q^2}{4\kappa^2}) U(\frac{Q^2}{4\kappa^2}, 0; \kappa^2 z^2) \stackrel{Q^2 \rightarrow \infty}{=} z Q K_1(zQ) [1 + O(Q^{-1})]$$



Parameters:

soft-wall:  $\kappa = \frac{1}{2} m_\rho = 0.388$  GeV;  
 hard-wall:  $\Lambda_{\text{QCD}} = z_m^{-1} = 0.322$  GeV

Soft-wall:  $r_C = 0.62$  fm  
 Hard-wall:  $r_C = 0.57$  fm  
 Experiment:  $r_C = 0.67$  fm

# Pion gravitational form factor $A_\pi$ in holographic QCD

- Brodsky-Hwang-Ma-Schmidt formula in LFQCD:

[Brodsky:2008pf]

$$A_\pi(Q^2) = \int \zeta_\perp d\zeta_\perp |\varphi_\pi(\zeta_\perp)|^2 \frac{1}{2} \zeta_\perp^2 Q^2 K_2(\zeta_\perp Q) + \text{higher Fock sector contributions}$$

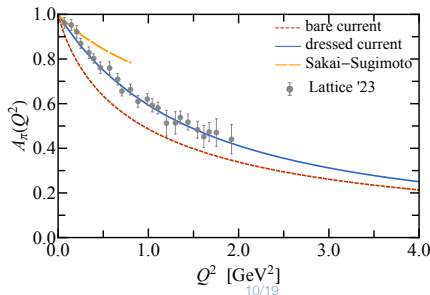
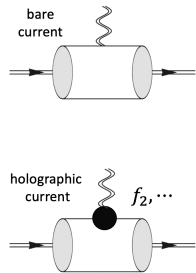
- Gravitational coupling in AdS<sub>5</sub>:  $g_{NM} \rightarrow g_{NM} + \delta g_{NM}$

[Abidin:2008hn, Katz:2005ir]

$$A_\pi(Q^2) = \int z dz |\varphi_\pi(z)|^2 H(Q^2, z)$$

where,  $H(Q^2, z)$  is the bulk-to-boundary propagator of the 5D gravitational field. In soft-wall model,

$$H(Q^2, z) = \Gamma(2 + \frac{Q^2}{8\kappa^2}) U(\frac{Q^2}{8\kappa^2}, -1; 2\kappa^2 z^2) \stackrel{Q^2 \rightarrow \infty}{\sim} \frac{1}{2} z^2 Q^2 K_2(zQ) [1 + O(Q^{-1})]$$



Parameters:

soft-wall:  $\kappa = \frac{1}{2} m_\rho = 0.388 \text{ GeV}$ ;  
hard-wall:  $\Lambda_{\text{QCD}} = z_m^{-1} = 0.322 \text{ GeV}$

Soft-wall:  $r_A = 0.39 \text{ fm}$   
Hard-wall:  $r_A = 0.32 \text{ fm}$   
Lattice '23:  $r_A = 0.41(1) \text{ fm}$

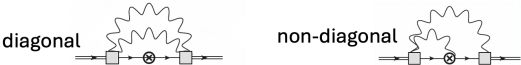
[Hackett:2023nkr]

## Forces inside hadrons: Pressure, surface tension, mechanical radius, and all that

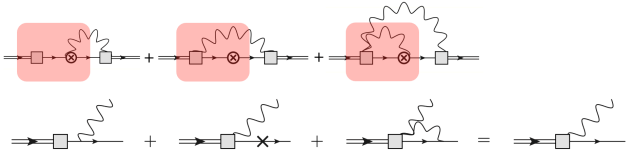
Maxim V. Polyakov and Peter Schweitzer ✉

<https://doi.org/10.1142/S0217751X18300259> | Cited by: 212 (Source: Crossref)

$\hat{T}_{++}$  of the EMT. Being related to the stress tensor  $\hat{T}_{ij}$  the form factor  $D(t)$  naturally “mixes” good and bad light-front components and is described in terms of transitions between different Fock state components in overlap representation. As a quantity intrinsically nondiagonal in a Fock space, it is difficult to study the  $D$ -term in approaches based on light-front wave functions. This is due to the rela-



- In an explicit model calculation (scalar model in 3+1D), we showed that all non-diagonal contributions add up to a diagonal contribution [Cao:2023ohj]



- How could this miracle happen?



# Light-front wave function representation of $D$ -term

In light-front dynamics, operators can be **exactly localized** in the transverse direction and then form factors are expressed as F.T. of one-body densities,

- Localizing the charge: Drell-Yan-West formula

[Drell:1969km, West:1970av]

$$Q = \int d^3x J^+ \rightarrow J_{\text{LF}}^+(r_\perp) = \sum_i e_i \delta^2(r_\perp - r_{i\perp}) \rightarrow \mathcal{F}(r_\perp) = \left\langle \sum_i e_i \delta^2(r_\perp - r_{i\perp}) \right\rangle$$

- Localizing the momentum: Brodsky-Hwang-Ma-Schmidt formula

[Brodsky:2008pf]

$$P^{+,\perp} = \int d^3x T^{++,\perp} \rightarrow \mathcal{A}(r_\perp) = \left\langle \sum_i x_i \delta^2(r_\perp - r_{i\perp}) \right\rangle$$

- Localizing the light-front Hamiltonian: **Cao-Li-Vary formula**

[Cao:2023ohj, Xu:2024hfx]

$$P^- = \int d^3x T^{+-} \rightarrow \mathcal{T}^{+-}(r_\perp) = \left\langle \sum_i \left\{ \frac{-\frac{1}{4}\vec{\nabla}_{i\perp}^2 + m_i^2 + \frac{1}{4}\nabla_\perp^2}{x_i} + V_i \right\} \delta^2(r_\perp - r_{i\perp}) \right\rangle$$

where,  $V = M^2 - \sum_i \frac{-\vec{\nabla}_{i\perp}^2 + m_i^2}{x_i}$  is the potential energy,  $\langle O \rangle \equiv \sum_n \int [dx_i d^2r_{i\perp}]_n \tilde{\psi}_n^*(\{x_i, \vec{r}_{i\perp}\}) O_n \tilde{\psi}_n(\{x_i, \vec{r}_{i\perp}\})$ .

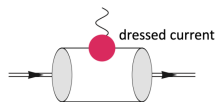
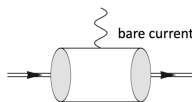
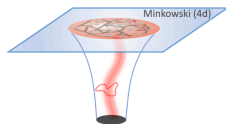
$$T^{+-}(Q^2) = (M_\pi^2 + \frac{1}{4}Q^2)A_\pi(Q^2) + \frac{1}{2}Q^2 D_\pi(Q^2)$$

$$\Rightarrow D_\pi(Q^2) = \int z dz |\varphi_\pi(z)|^2 \left\{ -\frac{1}{4}z^2 Q^2 K_2(zQ) - 2K_0(zQ) - \frac{2U(z)}{Q^2} \left[ zQ K_1(zQ) - \frac{z^2 Q^2}{2} K_2(zQ) \right] \right\}$$

$\Downarrow$

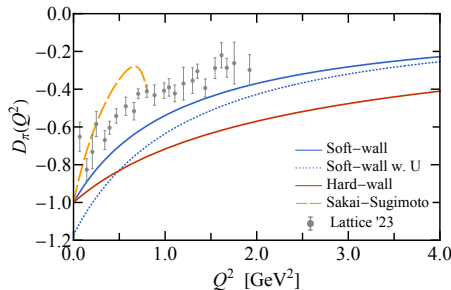
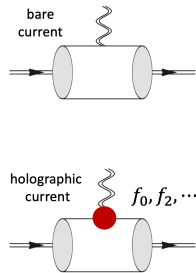
$$D_\pi^{\text{dress}}(Q^2) = \int z dz |\varphi_\pi(z)|^2 \left\{ -\frac{1}{2}H(Q^2, z) - \frac{1}{2}S(Q^2, z) - \frac{2U(z)}{Q^2} [V(Q^2, z) - H(Q^2, z)] \right\}$$

- Large  $Q^2$  scaling consistent with pQCD prediction  $D_\pi(Q^2) \sim 1/Q^2$  [Tong:2021ctu, Tong:2022zax]
- The dressed scalar current  $S(Q^2, z) = \Gamma(2 + \frac{Q^2}{4\kappa^2})U(\frac{Q^2}{4\kappa^2}, -1, \kappa^2 z^2)$  is the bulk-to-boundary propagator of a scalar field  $\text{tr}G^2$  ( $\Delta = 4$ ), i.e. **scalar glueball** [Colangelo:2007pt, Forkel:2007ru, Chen:2015zhx]



$$D_\pi(Q^2) = \int z dz |\varphi_\pi(z)|^2 \left\{ -\frac{1}{2} H(Q^2, z) - \frac{1}{2} S(Q^2, z) - \frac{2U(z)}{Q^2} \left[ V(Q^2, z) - H(Q^2, z) \right] \right\}$$

- In hard-wall model, the  $U$ -term vanishes and  $D_\pi^{\text{HW}}(0) = -1$ , consistent with  $\chi\text{PT}$
- In soft-wall model, the  $U$ -term leads to a small deviation from  $-1$ :  $D_\pi^{\text{SW}}(0) = -1.17$  due to the imperfect  $\chi\text{SB}$ , and should be neglected
- Glueball dominance: scalar glueball  $D_\pi^S(0) = -0.5$ , tensor meson  $D_\pi^T(0) = -0.5$



Parameters:  
 soft-wall:  $\kappa = \frac{1}{2} m_\rho = 0.388 \text{ GeV}$ ;  
 hard-wall:  $\Lambda_{\text{QCD}} = z_m^{-1} = 0.322 \text{ GeV}$

Soft-wall:  $r_D = 0.43 \text{ fm}$   
 Hard-wall:  $r_D = 0.32 \text{ fm}$   
 Lattice '23:  $r_D = 0.61(7) \text{ fm}$   
 [Hackett:2023nkr, ( $M_\pi = 170 \text{ MeV}$ )]

# Nucleon gravitational form factor $A_N$

- Brodsky-Hwang-Ma-Schmidt formula:

[Brodsky:2008pf]

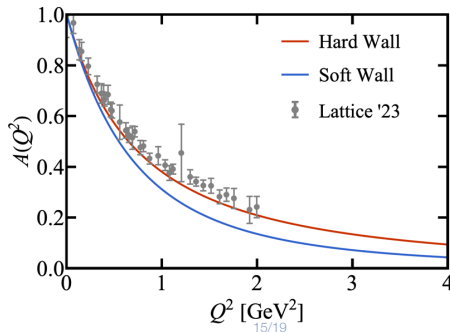
$$A_N(Q^2) = \int \zeta_\perp d\zeta_\perp \frac{1}{2} \left\{ |\varphi_+(\zeta_\perp)|^2 + |\varphi_-(\zeta_\perp)|^2 \right\} \frac{1}{2} \zeta_\perp^2 Q^2 K_2(\zeta_\perp Q) + \text{higher Fock sector contributions}$$

- Gravitational coupling in AdS<sub>5</sub>:  $g_{NM} \rightarrow g_{NM} + \delta g_{NM}$

[Abidin:2008hn]

$$A_N(Q^2) = \int z dz \frac{1}{2} \left\{ |\varphi_+(z)|^2 + |\varphi_-(z)|^2 \right\} H(Q^2, z)$$

here, recall  $H(Q^2, z)$  is the bulk-to-boundary propagator of the 5D gravitational field.



Parameters (simultaneously fixed to the proton mass  $M_p = 0.98$  GeV and rho mass

$$m_\rho = 0.70 \text{ GeV):}$$

soft-wall:  $\kappa = 0.35$  GeV;

hard-wall:  $\Lambda_{\text{QCD}} = z_m^{-1} = 0.245$  GeV;

soft-wall:  $r_A = 0.562$  fm; hard-wall:  $r_A = 0.562$  fm

Lattice:  $r_A = 0.51(1)$  fm

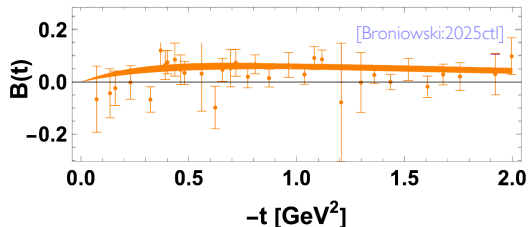
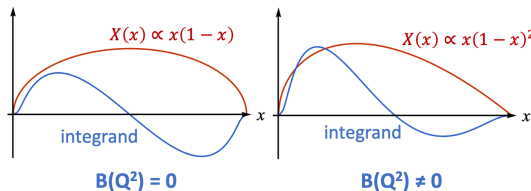
[Broniowski:2025ctl, Cao:2025dkv]

# Nucleon gravitational form factor $B_N$

- $B(0) = 0$  regardless of the choice of  $\varphi_{\pm}$ , consistent with the Equivalence Principle [Teryaev:1999su]
- In  $\text{AdS}_5$ ,  $B(Q^2) = 0$  for all  $Q^2$  due to the absence of the tensor coupling to a bulk Dirac fermion
- In LFQCD, Brodsky-Hwang-Ma-Schmidt formula: [Brodsky:2008pf]

$$\frac{q^1 + iq^2}{2M_N} B(Q^2) = \int dx \int d^2\zeta_{\perp} \varphi_{-}^{*}(\zeta_{\perp}) \varphi_{+}(\zeta_{\perp}) \left\{ x e^{-i\sqrt{\frac{1-x}{x}} \vec{q}_{\perp} \cdot \vec{\zeta}_{\perp}} + (1-x) e^{+i\sqrt{\frac{x}{1-x}} \vec{q}_{\perp} \cdot \vec{\zeta}_{\perp}} \right\} + \dots$$

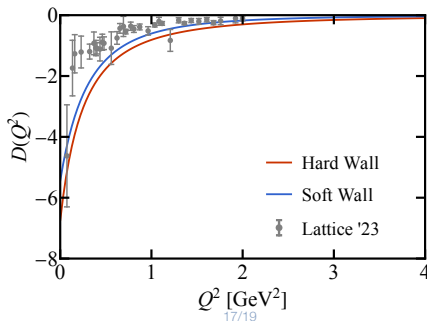
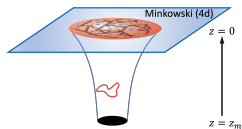
- $B(Q^2) = 0$  also for all  $Q^2 \neq 0$  due to a symmetry of  $\varphi_{\pm}(\zeta_{\perp})$  under swapping  $x \leftrightarrow 1-x$
- $\Rightarrow$  For a general longitudinal profile  $\psi(x, r_{\perp}) = X(x)\varphi(\zeta_{\perp})$ ,  $B(Q^2)$  is non-vanishing, but  $B(Q^2)$  is still **suppressed for longitudinal S-waves**



# Nucleon gravitational form factor $D_N$

$$D_N(Q^2) = \int z dz \left[ |\varphi_+(z)|^2 + |\varphi_-(z)|^2 \right] \left\{ -\frac{1}{4} S(Q^2, z) + \frac{4M^2}{Q^2} [V(Q^2, z) - H(Q^2, z)] \right\} \\ - \int d^2 z \left[ U_+(z) |\varphi_+(z)|^2 + U_-(z) |\varphi_-(z)|^2 \right] \frac{2}{Q^2} \{ V(Q^2, z) - H(Q^2, z) \}$$

- Large  $Q^2$  scaling  $D_N(Q^2) \sim 1/Q^4$  in contrast to the pQCD prediction
- Predictions to the  $D$ -term:  $D_N^{\text{SW}}(0) = -5.4$ ,  $D_N^{\text{HW}}(0) = -6.7$ , compared to the MIT Lattice result  $D = -3.87(97)$



Parameters (simultaneously fixed to the proton mass  $M_p = 0.98$  GeV and rho mass

$m_\rho = 0.70$  GeV):

soft-wall:  $\kappa = 0.35$  GeV;

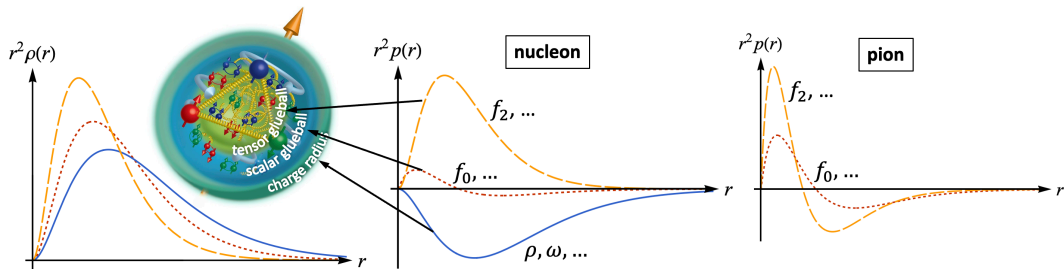
hard-wall:  $\Lambda_{\text{QCD}} = z_m^{-1} = 0.245$  GeV;

soft-wall:  $r_D = 0.882$  fm; hard-wall:  $r_D = 0.953$  fm

# Physical contributions at low $Q^2$

- Vector current, e.g.  $\rho, \omega$ , is attractive and long-ranged
- Tensor current, e.g.  $f_2$ , is repulsive and short-ranged
- Scalar current, e.g.  $f_0$ , is in mechanical equilibrium, and mid-ranged
- For the pion, both the scalar glueball and tensor glueball are self-balanced
- Large- $Q^2$  scaling is related to the parton distributions through Brodsky-Farrar parton counting rule

[Liu:2019vsn, deTeramond:2021lxc]

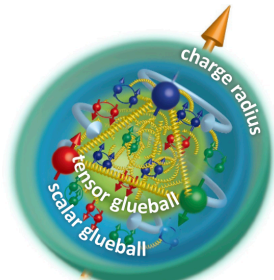


# Summary

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- The gravitational form factors emerge as a vital observable to unravel the internal structure of hadrons
- Holographic QCD provides a unique non-perturbative access to these observables with convincing physical pictures
- By exploiting the remarkable correspondence between light-front QCD and AdS/QCD, we computed the GFFs of the pion and the nucleon, and compared the results with recent Lattice simulations

Thank you!





A scene featuring three Stormtroopers in white armor. The trooper on the left is aiming a blaster upwards, with a red laser beam visible. The trooper in the center is crouching and aiming a blaster. The trooper on the right is aiming a blaster forward, with a bright red light effect at the muzzle. The background is a dark, futuristic interior with blue and red lighting.

I think we're gonna need

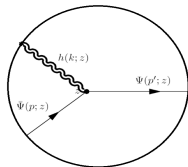
**Backup**

# GFFs in AdS/QCD in 5D

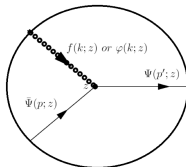
- Access to the tensor GFF  $A(Q^2)$ : graviton-hadron scattering in  $\text{AdS}_5$  [Abidin:2008hn, Abidin:2008ku, Abidin:2009hr]
- NO access to GFF  $D(Q^2)$ , since graviton only couples to the traceless part of  $T^{\mu\nu}$
- Adding the scalar graviton allows the access to GFF  $D(q^2)$ ; however, GFF  $D(q^2)$  vanishes at large  $N_c$  due to the degeneracy between the scalar and tensor glueballs [Mamo:2019mka, Mamo:2021krl, Mamo:2022eui]

$$D_N(q^2) = \frac{4M_N^2}{3q^2} \left[ A(q^2) - \Theta(q^2) \right]$$

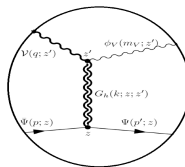
- Fit scalar and tensor glueballs separately and obtained  $A_N(0) = 0.43$  and  $D_N(0) = -1.275$
- GFFs interpreted as the gluon contributions, more relevant for the near-threshold VM production



tensor graviton coupling



scalar graviton coupling



Near threshold VM photo-production

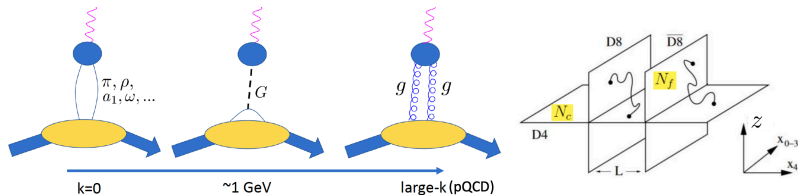
# GFFs in top-down holographic QCD

## Sakai-Sugimoto (SS) model in 10D

[Fujita:2022jus, Fujii:2024rqd]

- Key observation: the 4D part of the traceless  $T^{\mu\nu}$  becomes traceful
- Obtained  $D_\pi(0) = -1$  consistent with  $\chi$ EFT
- Large  $q^2$  scaling different from pQCD prediction
- $D_N(0) = -0.140(22)$  resulting from cancellation between  $U(1)$  and  $SU(2)$  fields
- New extraction using non-perturbative soliton solution gives  $D_N(0) = -2.05$

[Sugimoto:2025btn]



# Holographic light-front QCD

- The  $D$ -term involves non-minimal coupling terms in gravitational EFT,

[Donoghue:1994dn]

$$S_D = -\frac{D}{4} \int d^d x \sqrt{-g} R \phi^2$$

- Need constraints from both the QCD side and the gravity side
- Light-front holography: correspondence between semi-classical LFQCD and AdS/QCD in 5D
  - HLFQCD allows us to impose constraints from both the QCD and the gravity sides
  - Further insights: super-conformal algebra, Veneziano amplitudes, parton counting rules and GPD sum rules

[Review: Brodsky:2014yha]

