

粒子与天体物理学多前沿奇异现象研讨会 (MEPA 2025)

Probing light dark matters using gravitational waves, superfluid and Transmon Qubits

晁伟 (WEI CHAO)

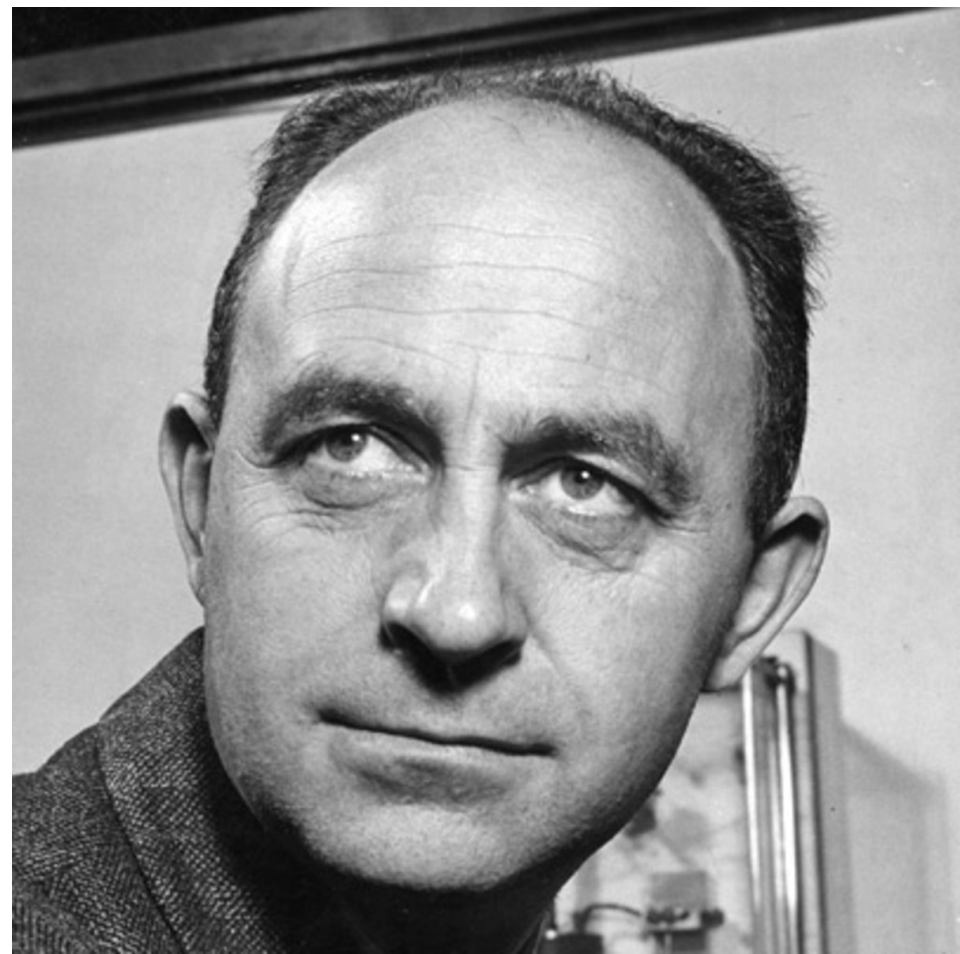
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2025.04.13@NANJING



Possible dark matter candidate

The WIMP Miracle!



Assuming DM is in thermal equilibrium in the early universe

Fermi's constant introduced in 1930s to describe beta decay

$G_F \approx 1.1 \times 10^{-5} \text{ GeV}^{-2}$

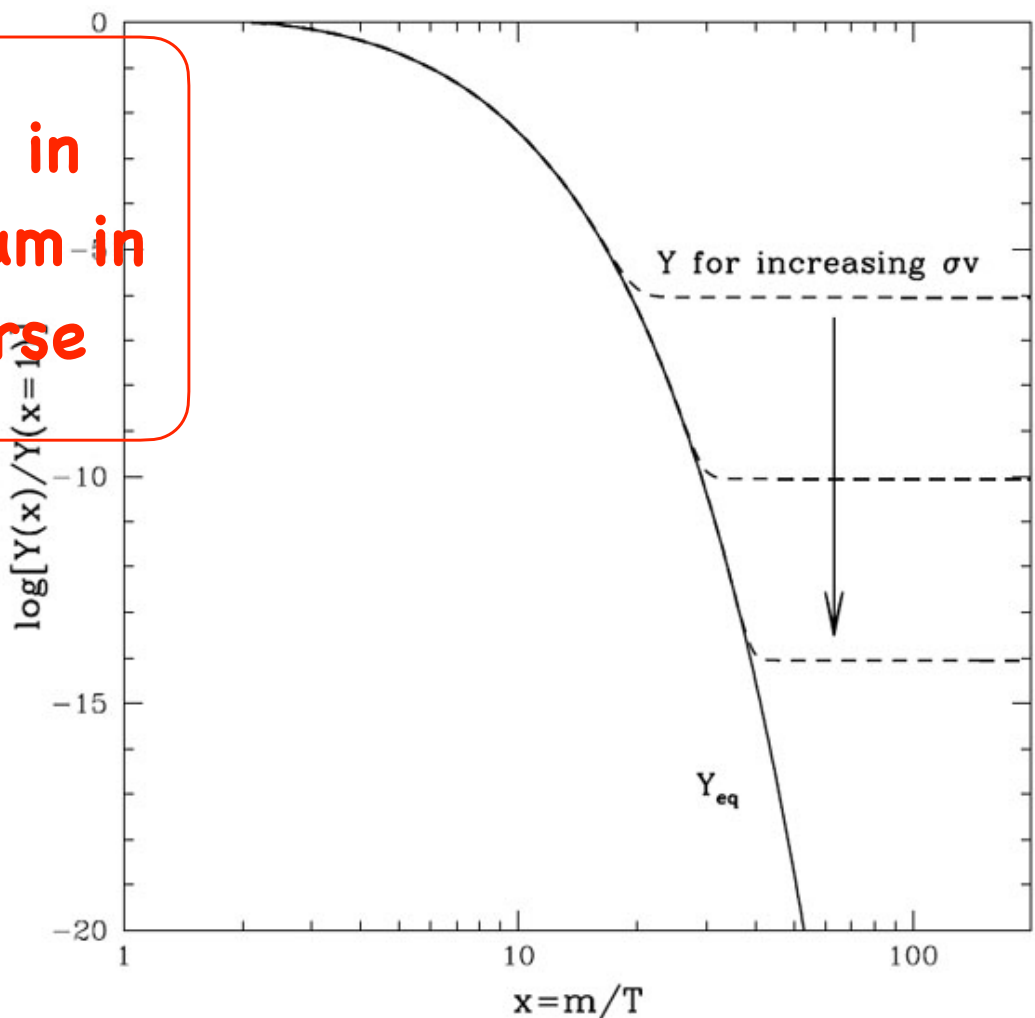
New mass scale: 100 GeV

The weak coupling

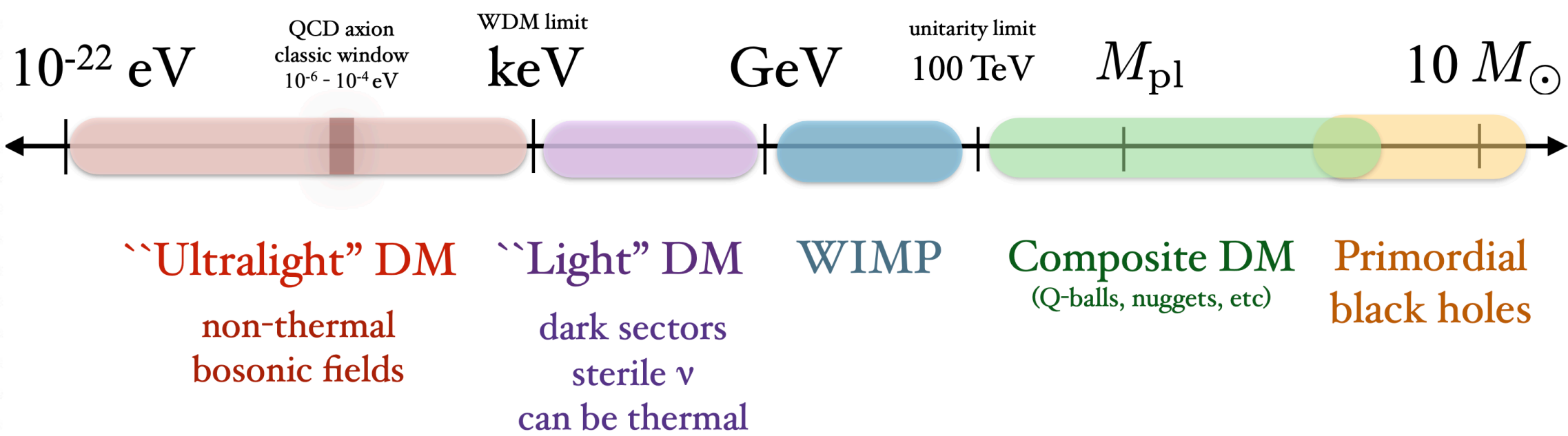
The relic density:

$$\Omega_\chi \sim \frac{1}{\langle \sigma v \rangle} \sim \frac{m_\chi^2}{g_\chi^4}$$

 $m_\chi \sim 100 \text{ GeV}$
 $g_\chi \sim 0.6$
 $\Omega_\chi \sim 0.1$



Dark matter candidate



Brief History:

WIMP

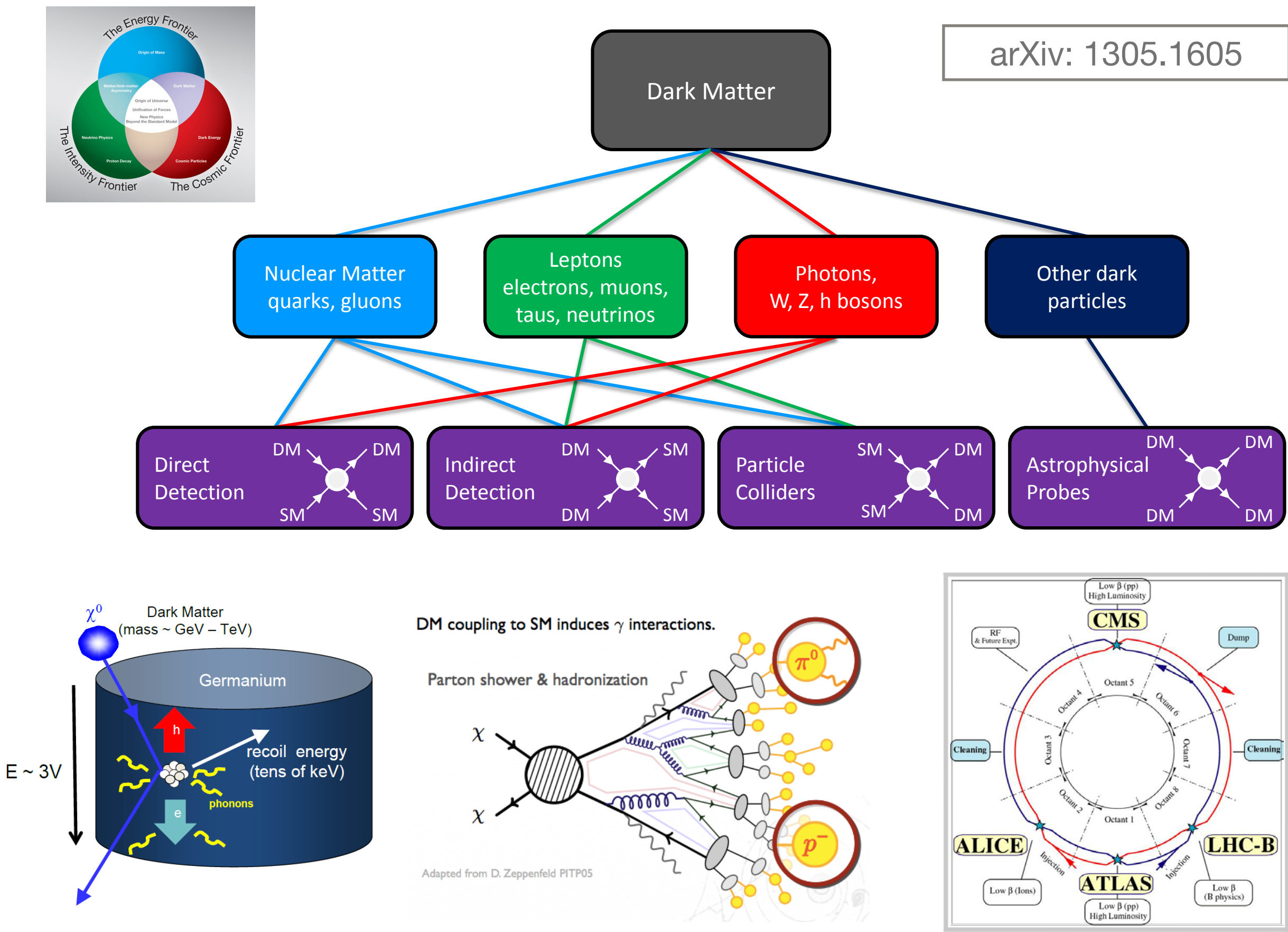
- Axion
- Sub-GeV
- Primordial black hole

- Super heavy DM
- Super light DM

...

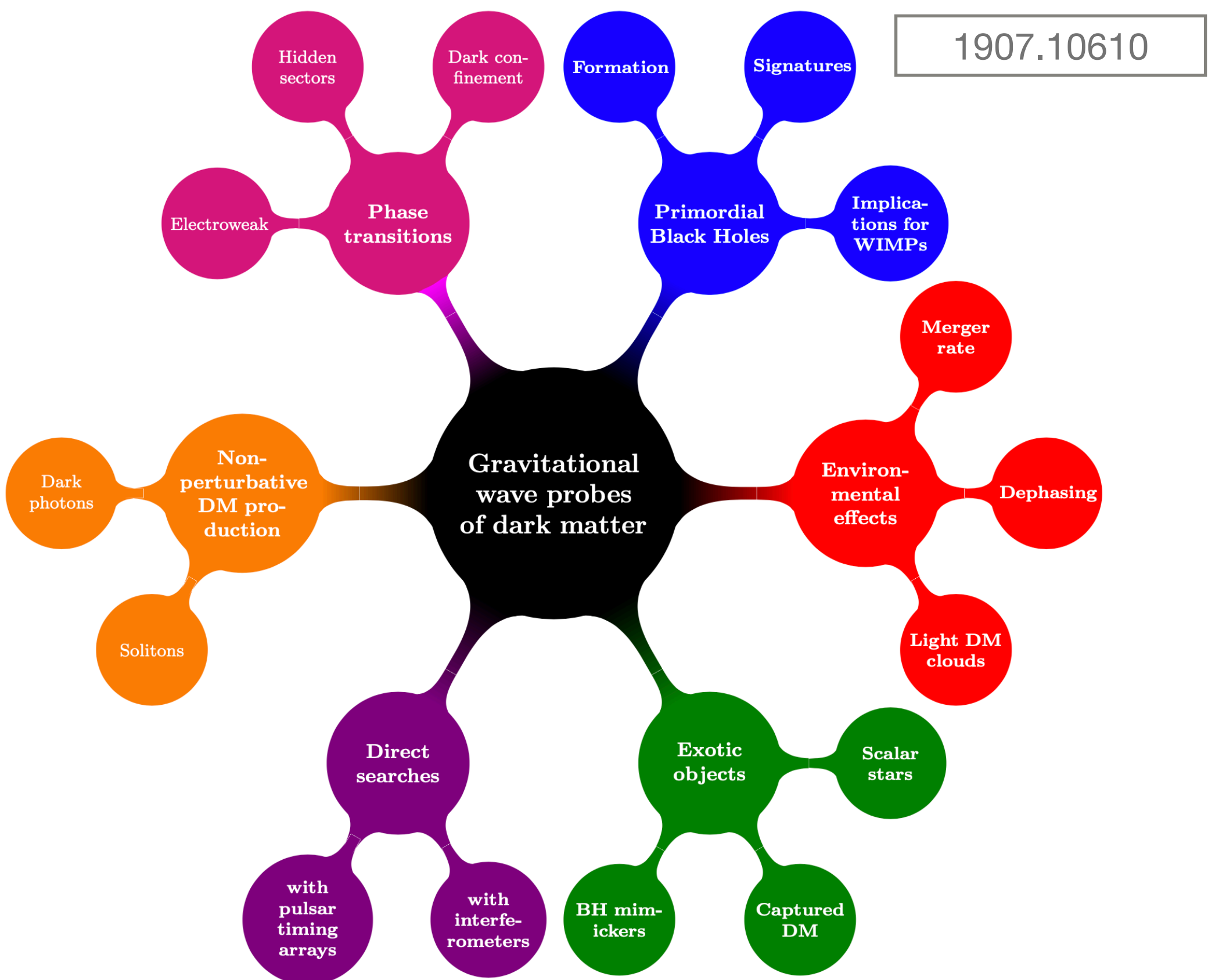
DM Dectections

Traditional methods



arXiv: 1305.1605

Modern strategies



Outline

- Gravitational wave signals of dark photon dark matter.

Oscillations of Ultralight Dark Photon into Gravitational Waves

Chao, Jingjing Feng, Huai-ke Gu and , Tong Li, Nucl.Phys.B(2024)

- Constraining exotic light DM with superfluid.

Axion and Dark Fermion Electromagnetic Form Factors in Superfluid He-4,

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- Constraining bosonic DM using transmon qubit

Probing light bosonic DM with transmon qubits,

Chao, Yu Gao, Mingjie Jin, Xiaosheng Liu and Xilei Sun, e-Print: 2412.20850 [hep-ph]

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Gravitational waves from DP oscillations

Where to start?: Photon-Graviton oscillation

G. Raffelt and L. Stodolsky, *Phys. Rev. D* 37, 1237 (1988).

Action

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{2} M_{pl}^2 R - \frac{1}{4} \hat{F}_{\mu\nu} g^{\alpha\mu} \hat{F}_{\alpha\beta} g^{\beta\nu} - \frac{\varepsilon}{2} F_{\mu\nu} g^{\alpha\mu} \hat{F}_{\alpha\beta} g^{\beta\nu} \right)$$

EOM

$$i \frac{\partial}{\partial z} \begin{pmatrix} \gamma_{\parallel} \\ \gamma'_{\parallel} \\ G_{\times} \end{pmatrix} = \left[\omega + \begin{pmatrix} \Delta_{\gamma\gamma} & \Delta_{\gamma\gamma'} & \Delta_{\gamma h} \\ \Delta_{\gamma\gamma'} & \Delta_{\gamma'\gamma'} & \Delta_{\gamma'h} \\ \Delta_{\gamma h} & \Delta_{\gamma'h} & 0 \end{pmatrix} \right] \begin{pmatrix} \gamma_{\parallel} \\ \gamma'_{\parallel} \\ G_{\times} \end{pmatrix}$$

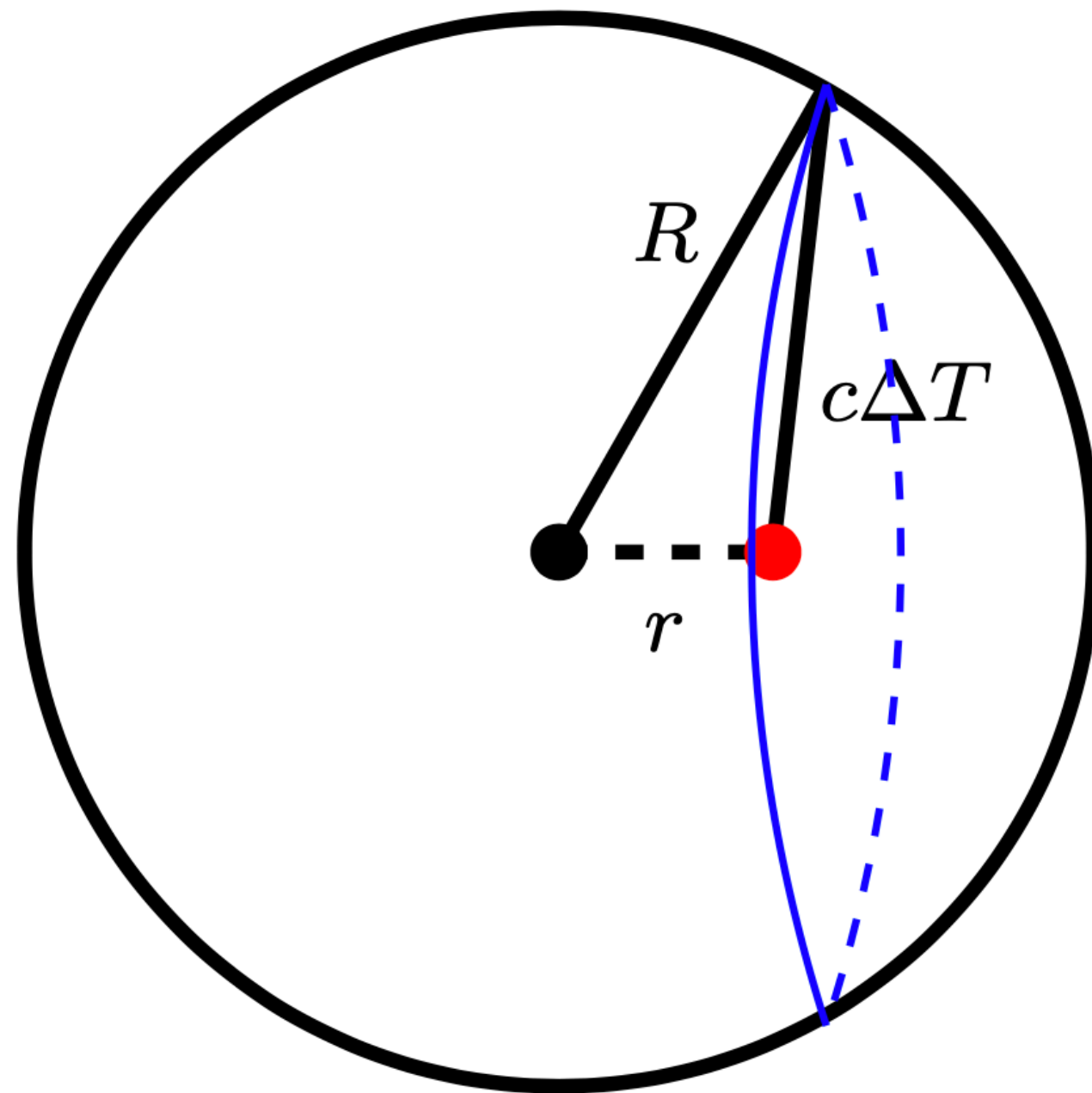
Oscillation probability

$$P(i \rightarrow j) = \left| \sum_{\alpha} \mathcal{U}_{i\alpha}^* e^{-i\rho_{\alpha} z} \mathcal{U}_{j\alpha} \right|^2, \quad i, j \in \gamma, \gamma', G_{\times}.$$

Two flavors

$$P(\gamma' \rightarrow G_{\times}) = \frac{4\Delta_{\gamma'h}^2}{4\Delta_{\gamma'h}^2 + \Delta_{\gamma'\gamma'}^2} \sin^2 \left(\frac{1}{2} \sqrt{4\Delta_{\gamma'h}^2 + \Delta_{\gamma'\gamma'}^2} z \right).$$

Local relic density of GW



$$f(r, \Delta T) = \begin{cases} 1, & c\Delta T \leq R - r, \\ \frac{1}{2} \left(1 - \frac{r^2 + c^2 \Delta T^2 - R^2}{2rc\Delta T} \right), & c\Delta T > R - r, \\ 0, & c\Delta T \geq R + r, \end{cases}$$

Energy density

$$\rho_{\text{LGW}}(t_0) = \frac{1}{V_{\text{GC}}} \int_{t_0-2T}^{t_0} dt \int_0^R dr \mathcal{L}(r) \cdot f(r, t_0 - t) 4\pi r^2.$$

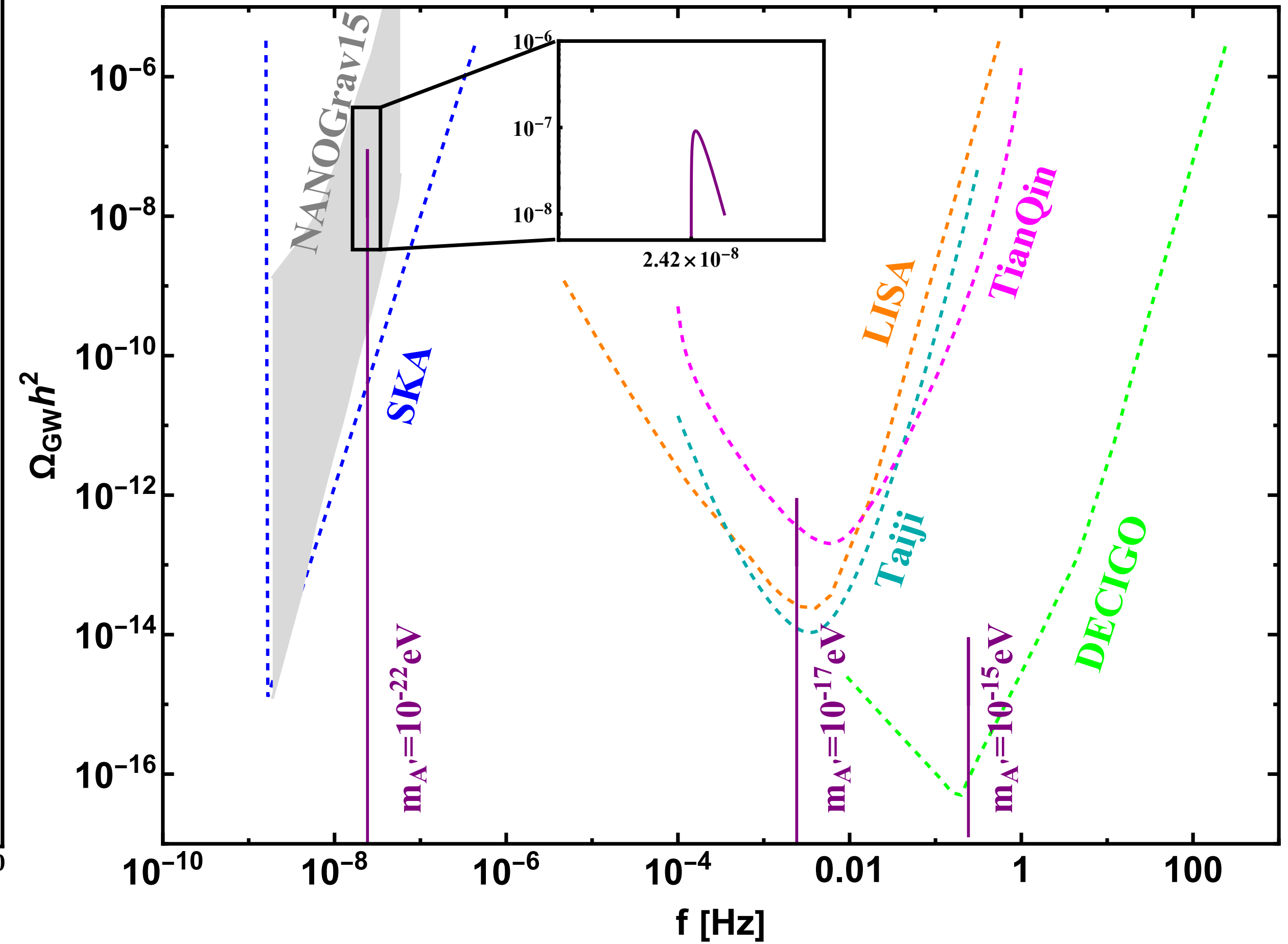
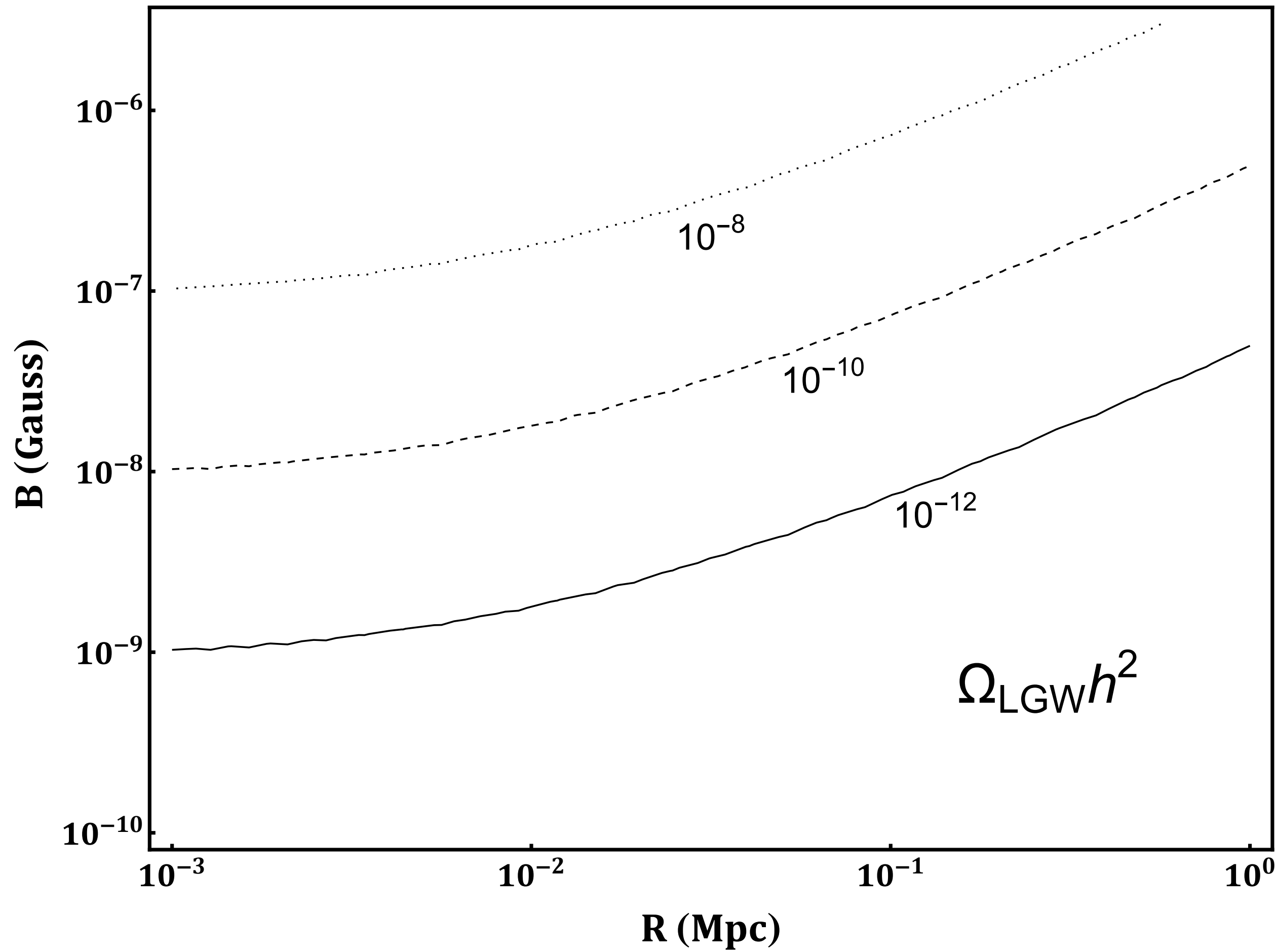
$\mathcal{L}$: luminosity of the oscillation

Flux of GW

$$\mathcal{D}_{\text{GW}} = \int \frac{dn_{\text{DM}}(r)}{dE} \frac{\mathcal{R} d^3 r}{4\pi |\vec{r} - \vec{d}|^2} \quad \mathcal{R} = \begin{cases} \frac{2\Delta_{\gamma'h}^2}{4\Delta_{\gamma'h}^2 + \Delta_{\gamma'\gamma'}^2} \frac{1}{T}, & L \gg vT, \\ P(\gamma' \rightarrow G_\times) \frac{1\text{kpc}}{L} \times 1.02 \times 10^{-14}, & L \ll vT, \end{cases}$$

$$\sim 1.72 \times 10^{-7} / m_{\text{ULDP}}^2 \text{ cm}^{-2} \cdot \text{s}^{-1} \cdot \text{eV}^{-1}$$

Local relic density of GW



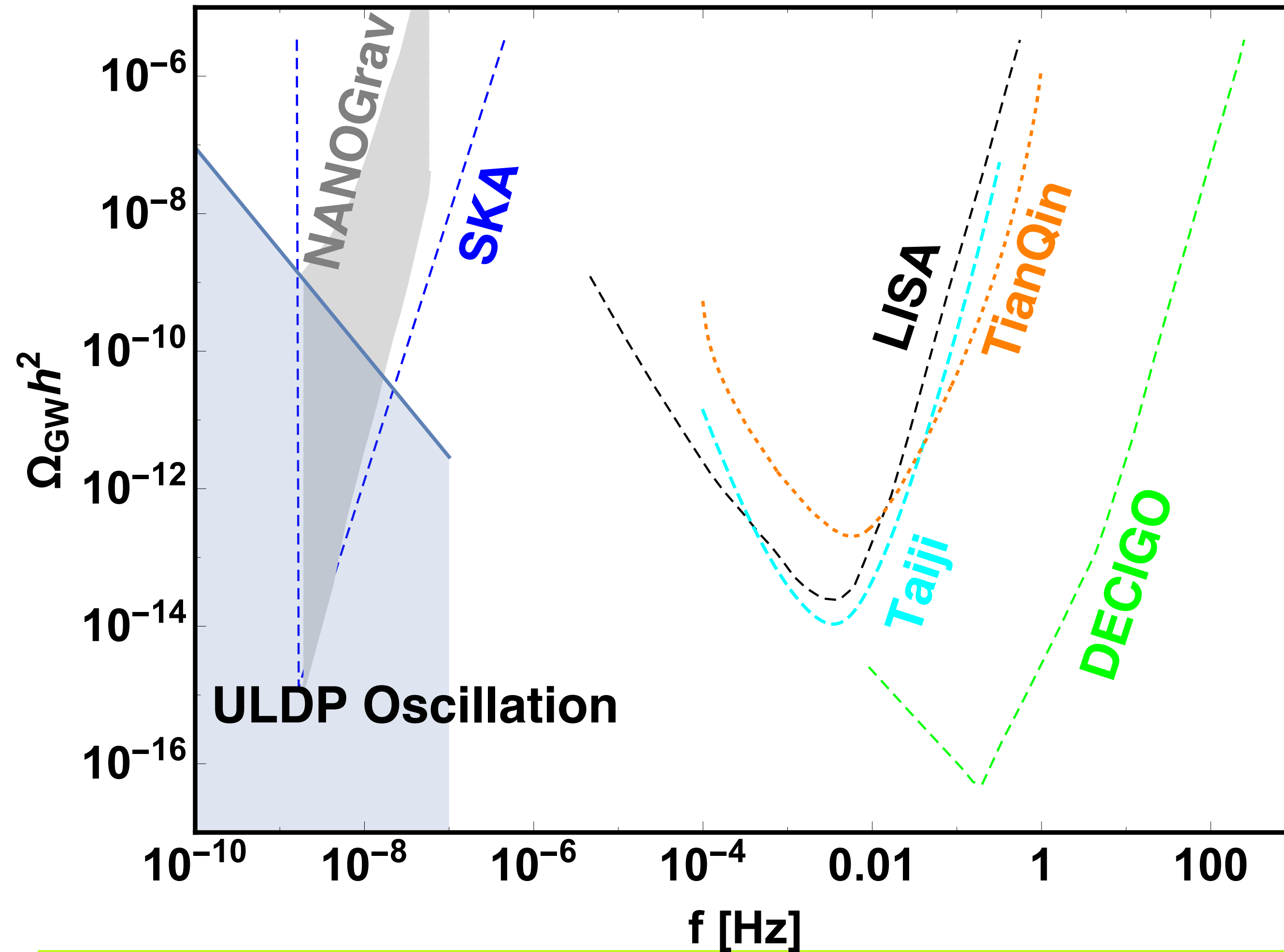
Dark photon oscillations in the early universe

Mechanisms for the production of primordial dark magnetic field

- PDMF can be generated during the pseudo-scalar inflation, ϕ , with the Chern-Simons interaction, $\frac{\phi}{4\pi f_a} F \widetilde{F}$.
- PDMF can be produced from the axion oscillations.
- PDMF can be generated from the oscillation of photon in the axion or some other exotic background.

Redshifted as $(1+z)^{-2}$

Dark photon oscillations in the early universe



$$\frac{d\rho_{\text{GW}}}{dt} = E_{\text{ULDP}} n_{\text{ULDP}} \mathcal{R}(\gamma' \rightarrow \text{SGW}) ,$$

$$\rho_{\text{SGW}} = \int_{z_0}^{z_1} n_{\text{ULDP}}(z_0) \mathcal{R}(z) m_{\text{ULDP}} \frac{dz}{(1+z)^2 H(z)}$$

$$\Omega_{\text{GW}}(f) = \frac{n_{\text{ULDP}} \mathcal{R}(f)}{H(f)} \frac{f}{\rho_C} ,$$

$$B(z_1) \sim 1 \text{ Gauss}$$

$$m_{\text{DP}} \sim 10^{-20} \text{ eV}$$

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Direct detections of light DM

Problem: Kinetic energy of light DM is too small

$$v_{\text{DM}} \sim 10^{-3}$$

$$E_R \sim \frac{(p_{\text{DM}} - p'_{\text{DM}})^2}{2m_T} \ll E_{\text{Threshold}}$$

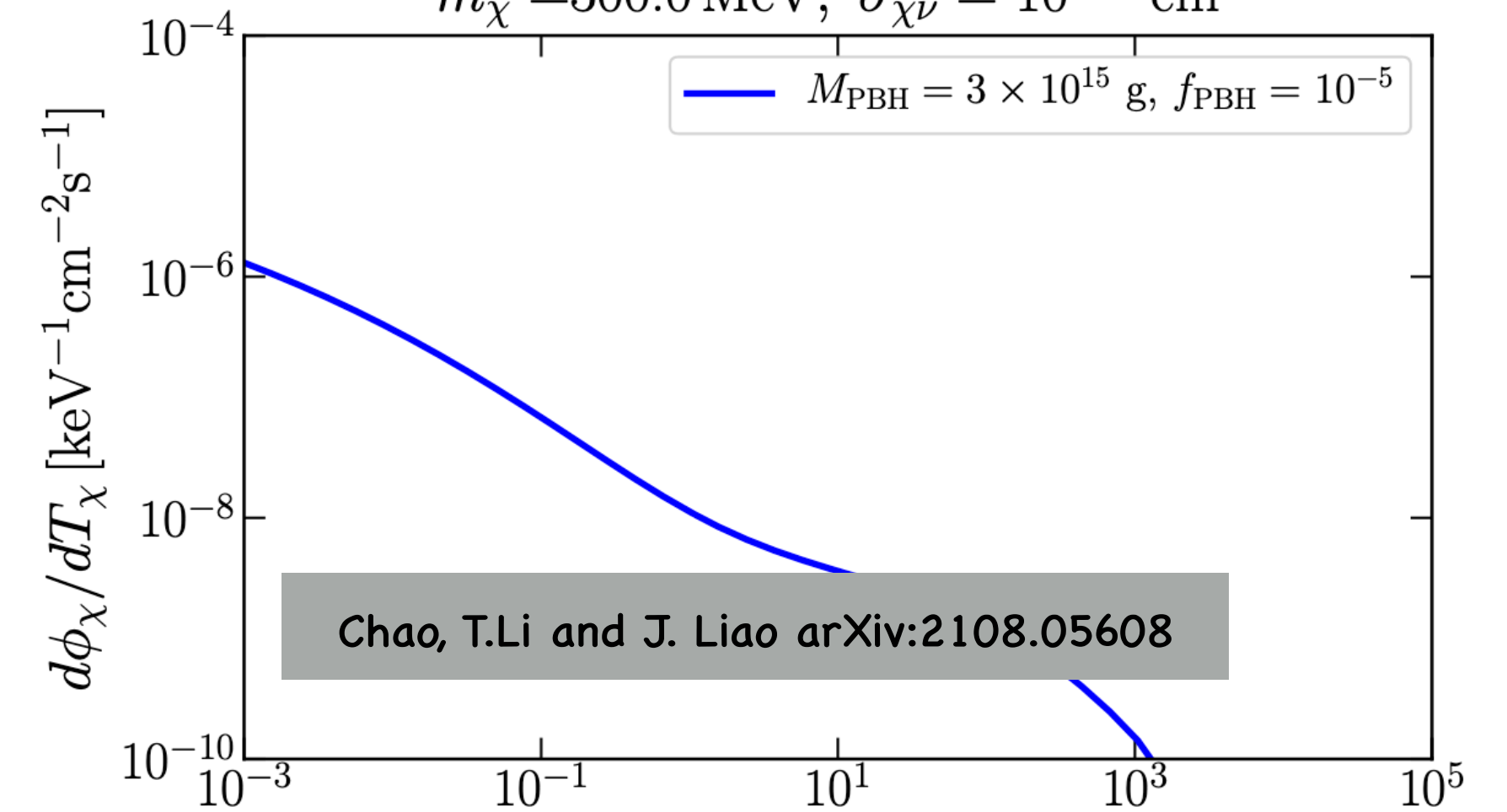
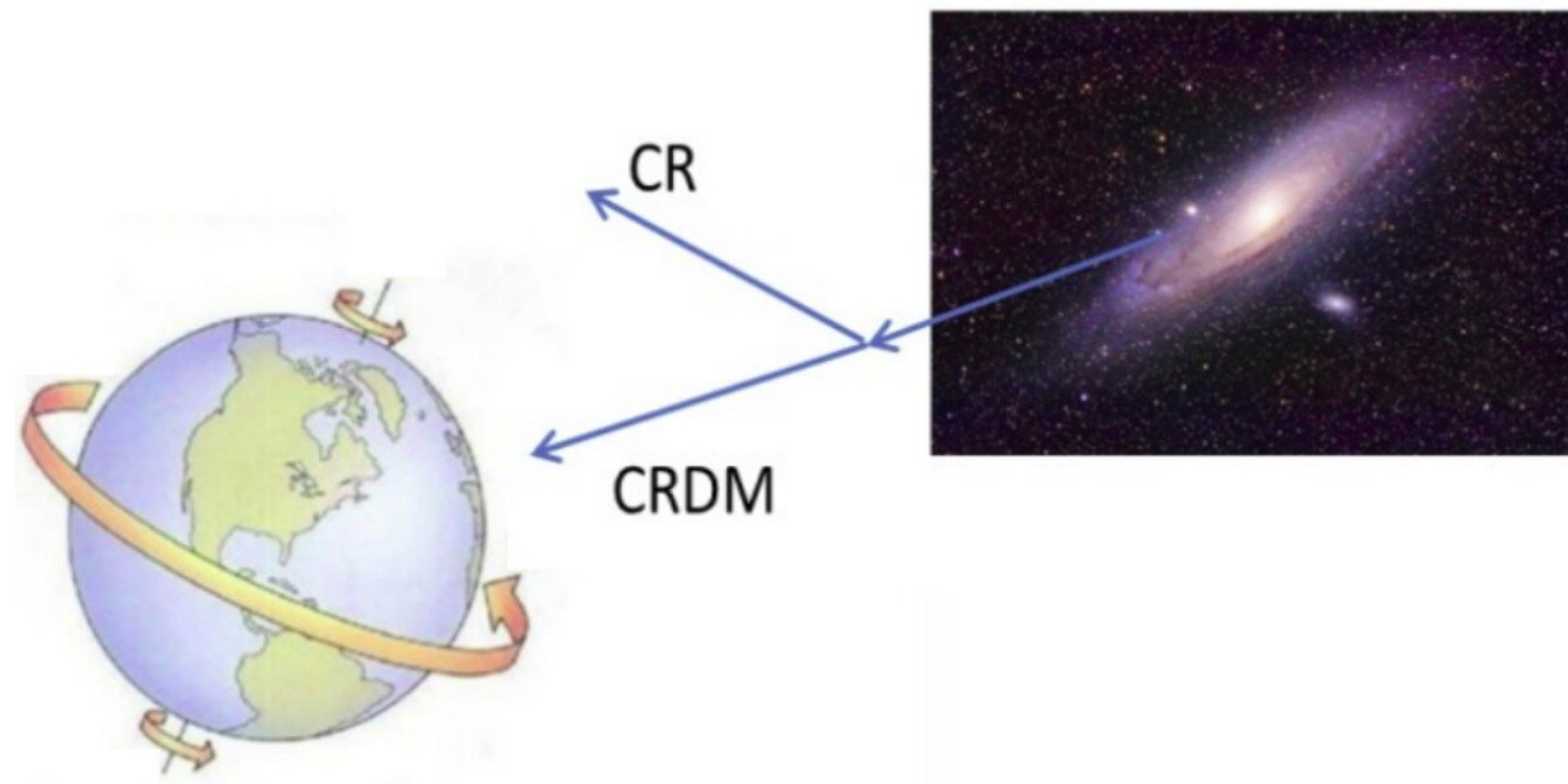
Strategy-1: Boosted DM

$$\frac{d\phi_\chi}{dE_\nu} = \int \frac{d\Omega}{4\pi} \int d\ell \frac{\rho_\chi}{m_\chi} \Phi_\nu(E_\nu) \sigma_{\chi\nu} = D_{\text{halo}} \Phi_\nu(E_\nu) \frac{\sigma_{\chi\nu}}{m_\chi}$$

$$D_{\text{halo}} = 2.02 \times 10^{25} \text{MeV} \cdot \text{cm}^{-2}$$

$$\frac{d\phi_\chi}{dT_\chi} = \int dE_\nu \frac{d\phi_\chi}{dE_\nu} \frac{1}{T_\chi^{\text{max}}(E_\nu)} \Theta(T_\chi^{\text{max}} - T_\chi)$$

$$m_\chi = 300.0 \text{ MeV}, \sigma_{\chi\nu} = 10^{-28} \text{ cm}^2$$



Direct detections of light DM

Strategy-II: Searching for DM using Condense Matter system!

DM mass	DM energy or momentum	CM scale
50 MeV	$p_\chi \sim 50 \text{ keV}$	zero-point ion momentum in lattice
20 MeV	$E_\chi \sim 10 \text{ eV}$	atomic ionization energy
2 MeV	$E_\chi \sim 1 \text{ eV}$	semiconductor band gap
100 keV	$E_\chi \sim 50 \text{ meV}$	optical phonon energy

Fermi's Golden Rule

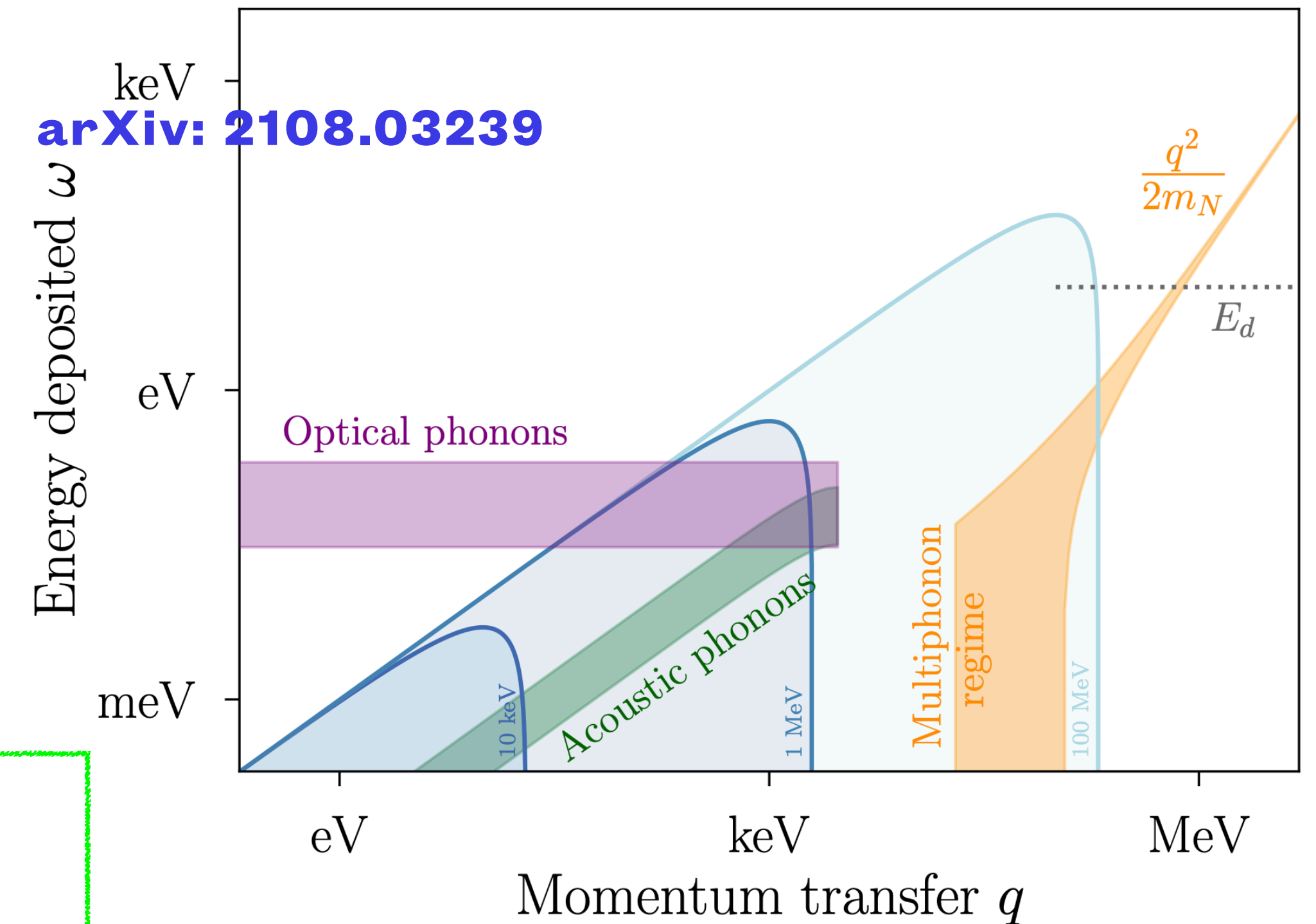
$$\begin{aligned}
 R_\chi &= \frac{1}{\rho_T} \frac{\rho_\chi}{m_\chi} \int d^3v f_\chi(v) \frac{V d^3p'_\chi}{(2\pi)^3} \sum_f |\langle f, p'_\chi | \Delta H_\chi | i, p_\chi \rangle|^2 2\pi \delta(E_f - E_i - \omega) \\
 &= \frac{1}{\rho_T} \frac{\rho_\chi}{m_\chi} \int \frac{d^3q}{(2\pi)^3} d\omega g(q, \omega) \tilde{V}(q) S(q, \omega) \\
 &= \frac{1}{\rho_T} \frac{\rho_\chi}{m_\chi} \frac{\pi \bar{\sigma}(q)}{\mu_\chi^2} \int \frac{q dq}{(2\pi)^2} d\omega \eta(v_{\min}(q, \omega)) \times S(q, \omega)
 \end{aligned}$$

■ $S(q, \omega)$: dynamical structure factor from correlation function

■ $\tilde{V}(q)$: potential

- $S(q, \omega)$: dynamical structure factor from condensed matter

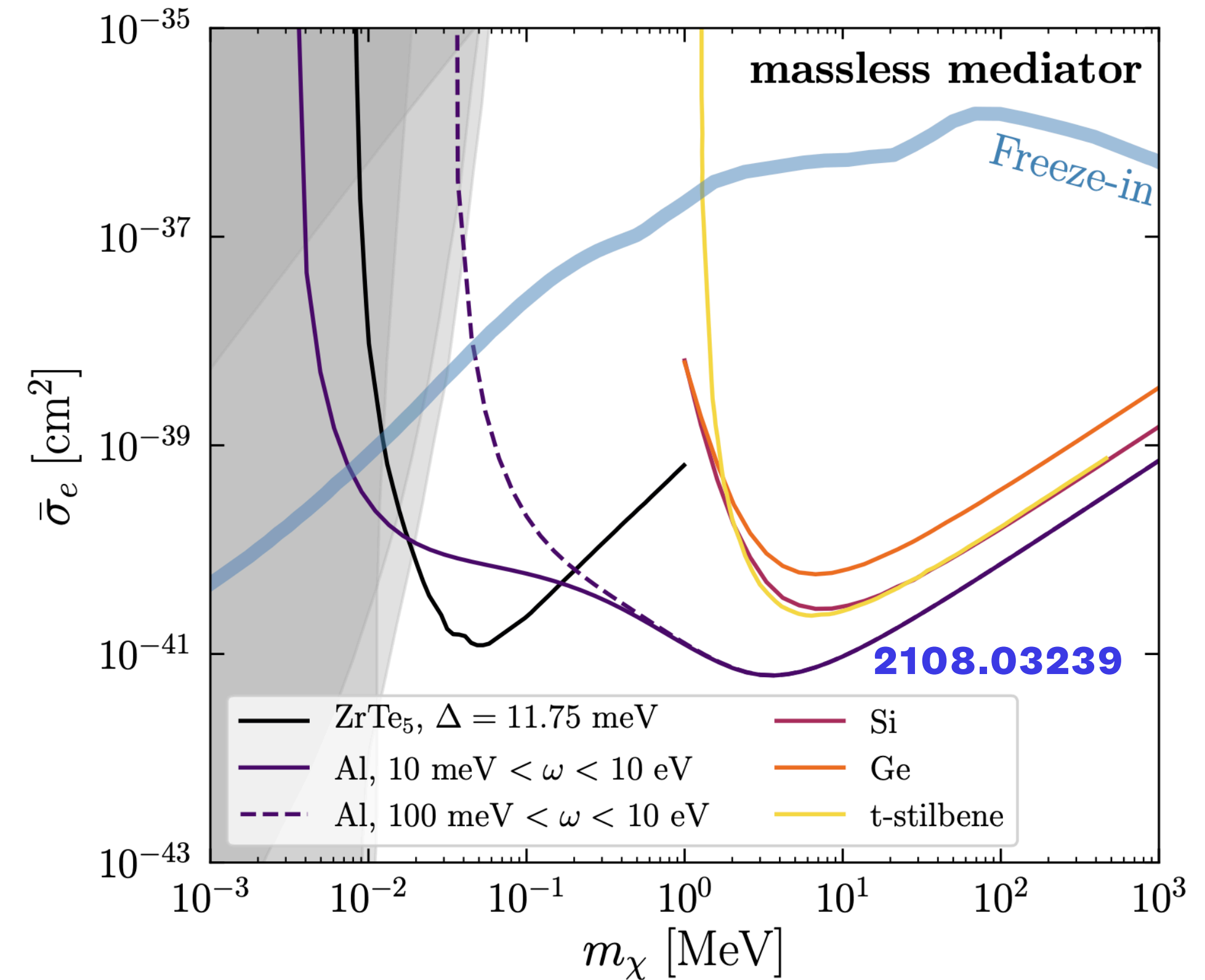
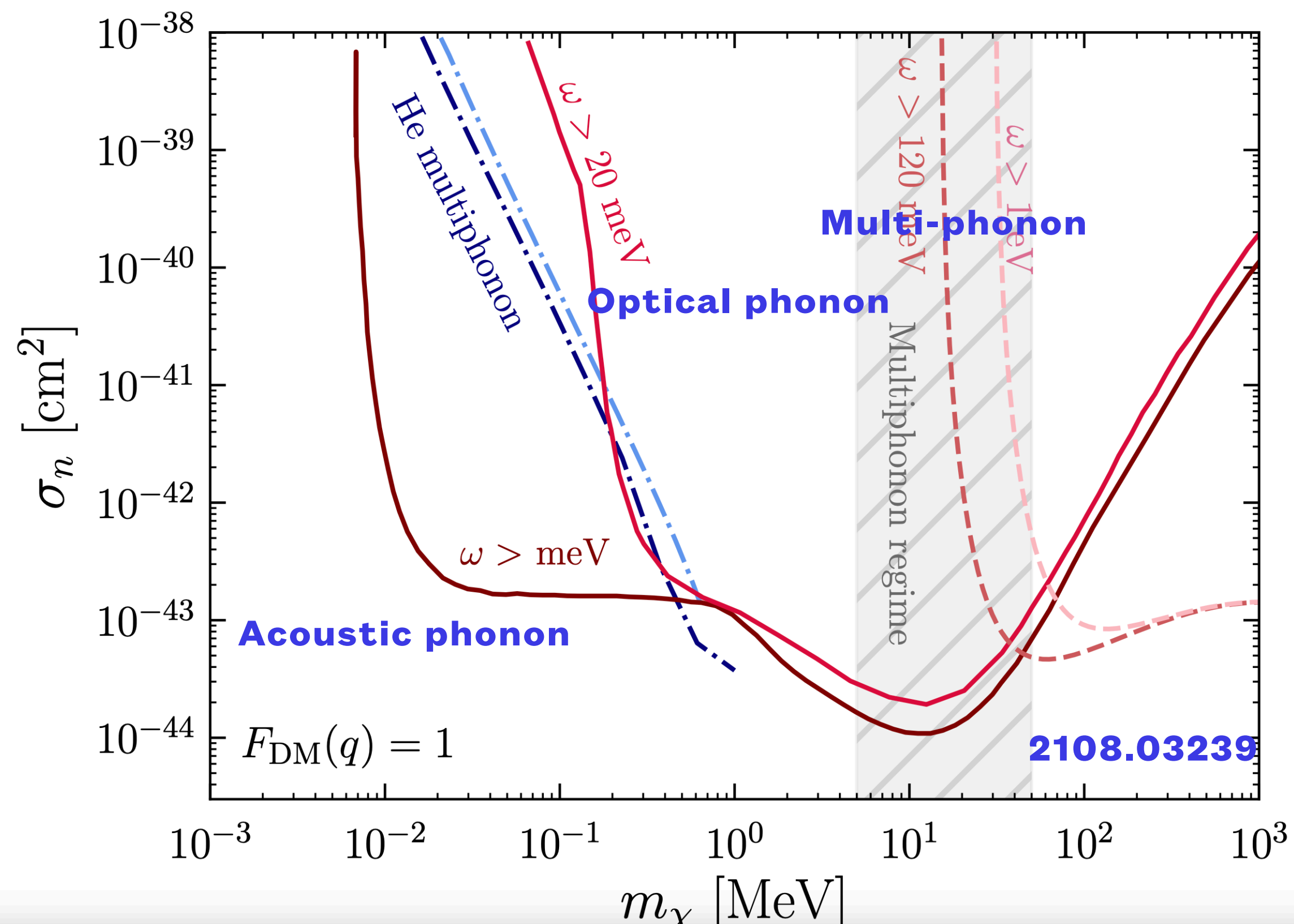
■ $\tilde{V}(q)$: potential felled by DM



Direct detections of light DM

$$R = \frac{1}{\rho_T} \frac{\rho_\chi}{m_\chi} \frac{\pi \bar{\sigma}(q)}{\mu_\chi^2} \int \frac{qdq}{(2\pi)^2} d\omega \eta(v_{\min}(q, \omega)) \times S(q, \omega)$$

$$\frac{dR}{d \ln E_R} = N_T \frac{\rho_\chi}{m_\chi} \frac{\bar{\sigma}_e}{8\mu_{\chi e}^2} \int dq q |F_{\text{DM}}(q)|^2 |f_{\text{ion}}(k, q)|^2 \eta(v_{\min})$$



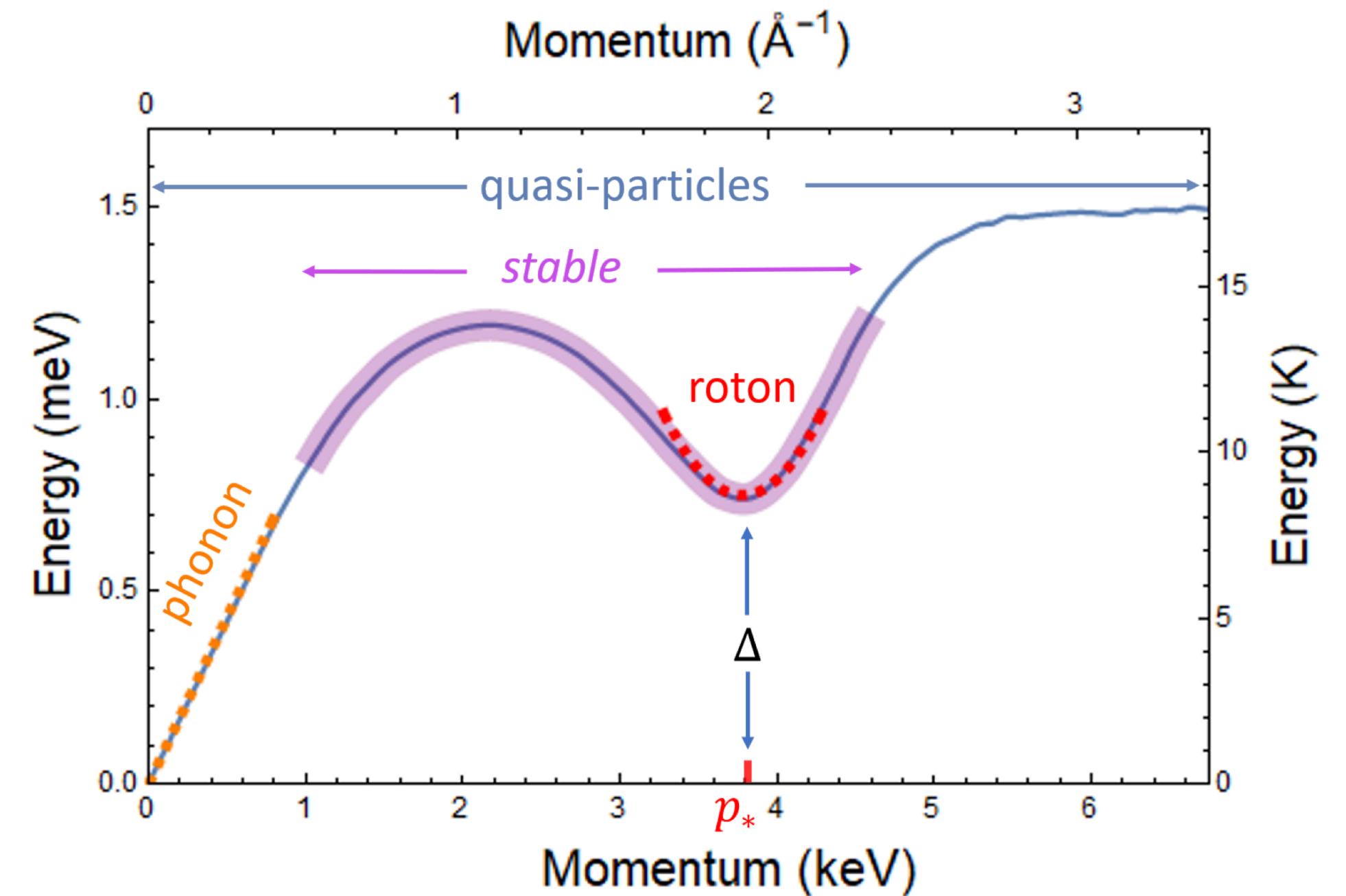
Detecting light DM using superfluid

Phonon (quasiparticle) in Superfluid

Helium-4: A Goldstone-like particle from the spontaneous breaking of the U(1) symmetry as well as the breaking of the boosts and the time translations in the superfluid He-4.

Action of phonon field

$$S_{int} \sim \int d^4x \sqrt{\frac{\mu}{\bar{n}}} c_s \left[\left(\frac{\mu^2}{2} \frac{db}{d\mu} + \mu b \right) \dot{\pi} F^{0\rho} F^0_{\rho} - \mu b \partial_j \pi F^{ij} F_{0i} + \frac{b}{2} \sqrt{\frac{\mu}{\bar{n}}} c_s \partial_\mu \pi \partial_\nu \pi F^{\mu\rho} F^\nu_{\rho} \right]$$



Phonon-photon conversion in external electric field

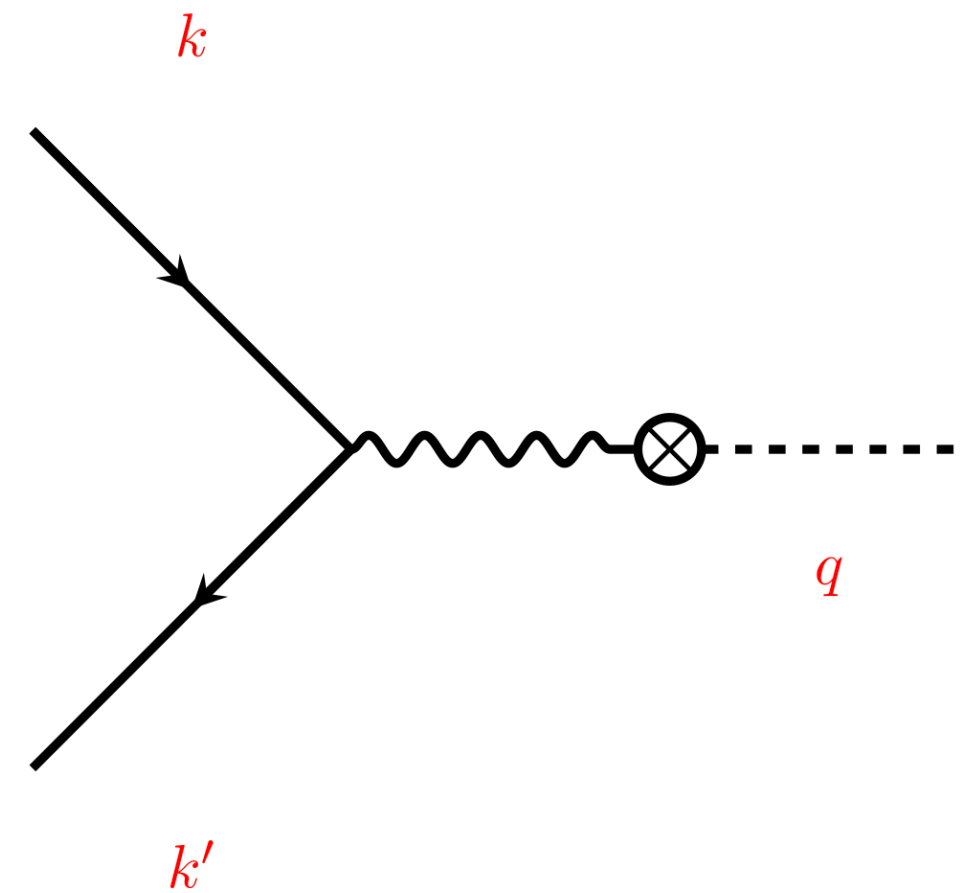
- A. Caputo, A. Esposito, E. Geoffray, A. D. Polosa, and S. Sun, Phys. Lett. B 802, 135258 (2020), arXiv:1911.04511 [hep-ph].

DM electromagnetic form factors in superfluid

Lagrangian:
$$S_{\text{eff}} = \int d^4x \left[\bar{\chi}(i\not{\partial} - m_\chi)\chi - \alpha\bar{\chi}\gamma^\mu\chi\partial^\nu F_{\mu\nu} - \beta\bar{\chi}\gamma^\mu\gamma^5\chi\partial^\nu F_{\mu\nu} - \frac{\lambda}{2}\bar{\chi}\sigma^{\mu\nu}\chi F_{\mu\nu} - \frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu} - \frac{1}{2}g_{a\gamma\gamma'}aF_{\mu\nu}\tilde{F}'^{\mu\nu} + \text{H.c.} \right. \\ \left. + \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}a(X)F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}b(X)F^{\mu\rho}F^\nu_\rho\partial_\mu\phi\partial_\nu\phi \right],$$

α : charge radius; β : Anapole moment; λ : magnetic moment.

Feynman diagram for DM-Helium scattering:



$$= i \frac{\alpha_E \bar{n}}{m_{He} c_s} \sqrt{\frac{m_{He}}{\bar{n}}} E_i \omega_q (\omega_q \delta_i^\mu + q_i \delta_0^\mu),$$

DM electromagnetic form factors in superfluid

Event rate:

$$R = \int dv_\chi f_{\text{MB}}(v_\chi) \frac{\rho_\chi}{m_{\text{He}} \bar{n} m_\chi} \int_{\omega_{\text{min}}}^{\omega_{\text{max}}} d\omega \frac{d\Gamma}{d\omega},$$

$$\frac{d\Gamma}{d\omega} = \frac{\bar{n} \alpha_E^2 \alpha^2}{4\pi c_s^6 m_{\text{He}} v_\chi} |\mathbf{E}|^2 \omega_q^4 \cos^2 \theta_E,$$

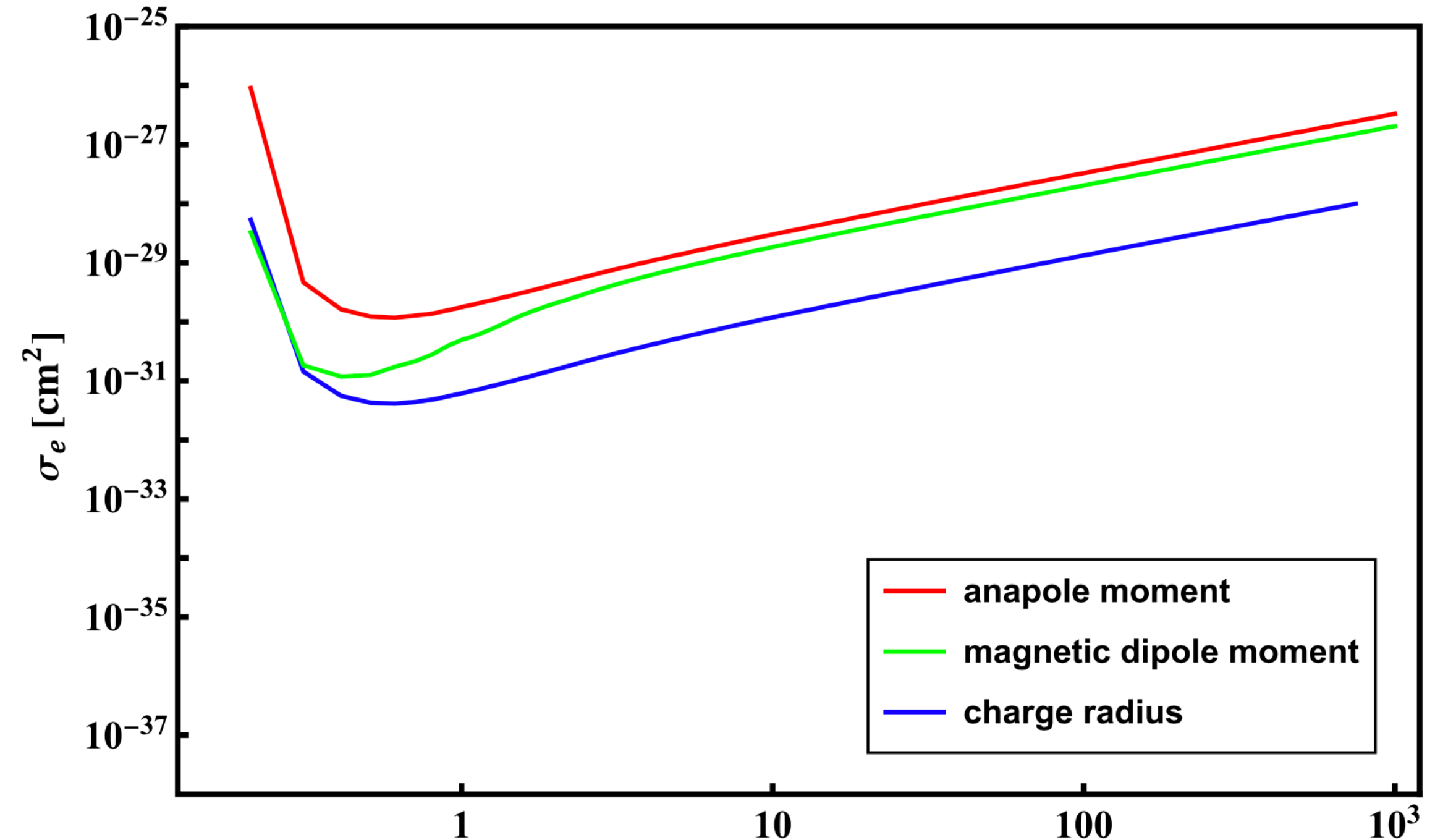
Projected constraints:

$$\sigma_e \equiv \frac{1}{16\pi(m_e + m_\chi)^2} \overline{|\mathcal{M}_{\chi e}(q)|^2}_{q^2=\alpha_{em}^2 m_e^2},$$

$$\text{charge radius: } \overline{|\mathcal{M}_{\chi e}(q)|^2}_{q^2=\alpha_{em}^2 m_e^2} = 64\pi\alpha^2 \alpha_{em} m_e^2 m_\chi^2, \quad (20a)$$

$$\text{anapole moment: } \overline{|\mathcal{M}_{\chi e}(q)|^2}_{q^2=\alpha_{em}^2 m_e^2} = 16\pi\beta^2 \alpha_{em} m_\chi^2 q^2, \quad (20b)$$

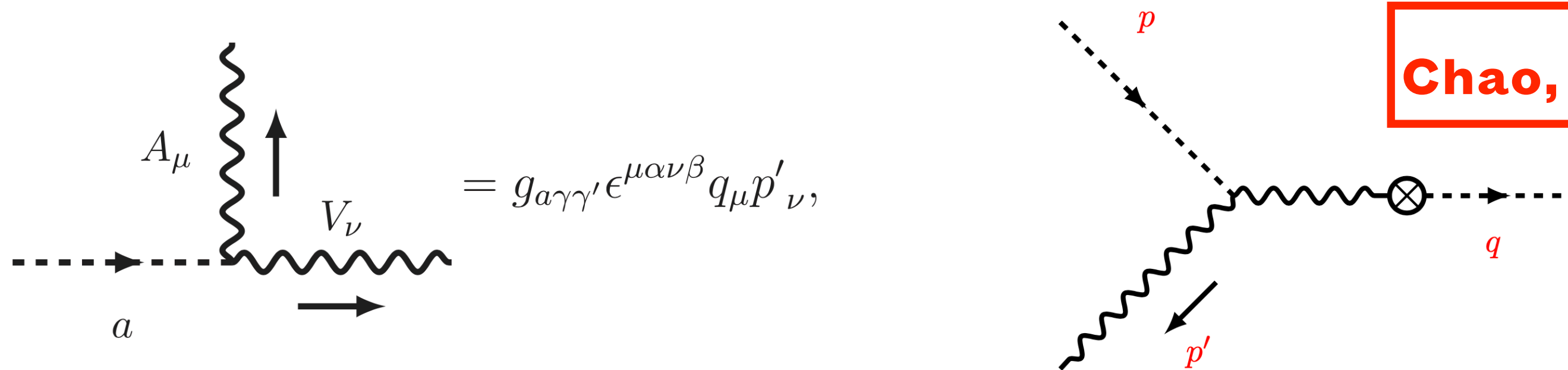
$$\text{magnetic moment: } \overline{|\mathcal{M}_{\chi e}(q)|^2}_{q^2=\alpha_{em}^2 m_e^2} = 16\pi\lambda^2 \alpha_{em} m_\chi^2$$



Chao, Sichun Sun, Xin Wang, and Chenhui Xie, PRD 2024

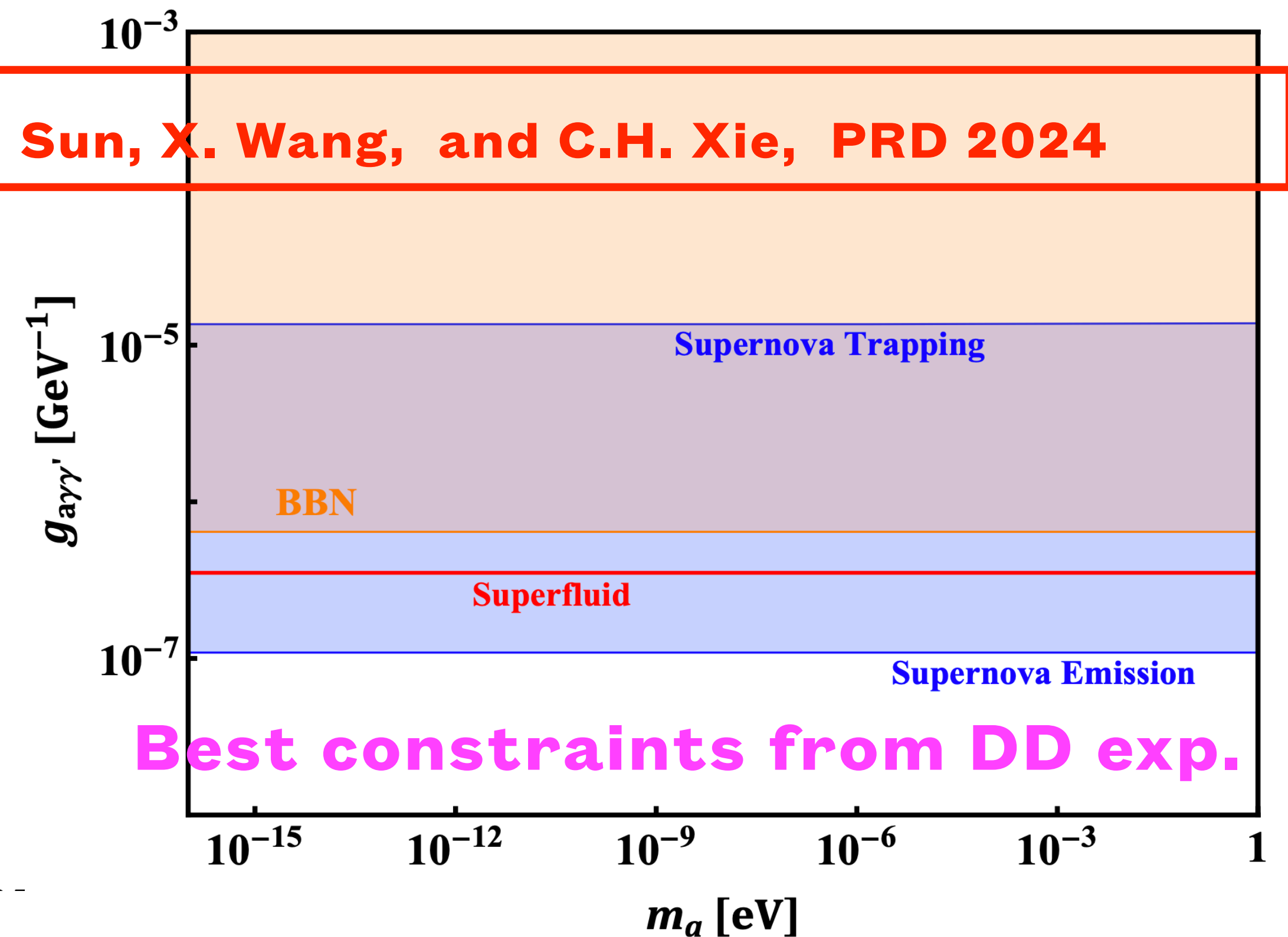
Detecting ALP using superfluid

Why? (1) No DD constraint on this coupling ; (2) There are already strong constraint on $g_{a\gamma\gamma}$.



Chao, S. Sun, X. Wang, and C.H. Xie, PRD 2024

$$\frac{d\Gamma}{d\omega} = \frac{g_{a\gamma\gamma'}^2 \bar{n} \alpha_E^2}{16\pi m_a \omega^2 m_{He} E v_a} \frac{|\mathbf{E}|^2 \omega^2}{c_s^2} \left\{ (\cos^2 \theta_E - c_s^2) \omega^2 \left(1 - \frac{1}{c_s^2} \right) \left[E^2 (1 - v_a^2) + 2E\omega \left(\frac{v_a}{c_s} \cos \theta - 1 \right) + \omega^2 \left(1 - \frac{1}{c_s^2} \right) \right] - \omega^2 \left(1 - \frac{1}{c_s^2} \right) (E \cos \theta_E - |\mathbf{p}|^2 c_s \cos \theta_a)^2 - (\cos^2 \theta_E - c_s^2) \left[\omega \left(E - \frac{|\mathbf{p}|}{c_s} \cos \theta \right) - \omega^2 \left(1 - \frac{1}{c_s^2} \right) \right]^2 \right\}.$$



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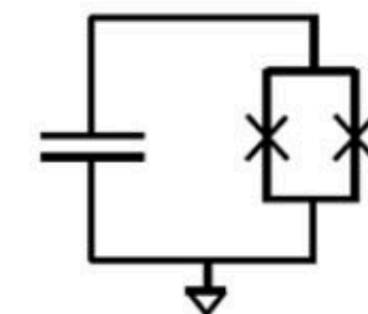
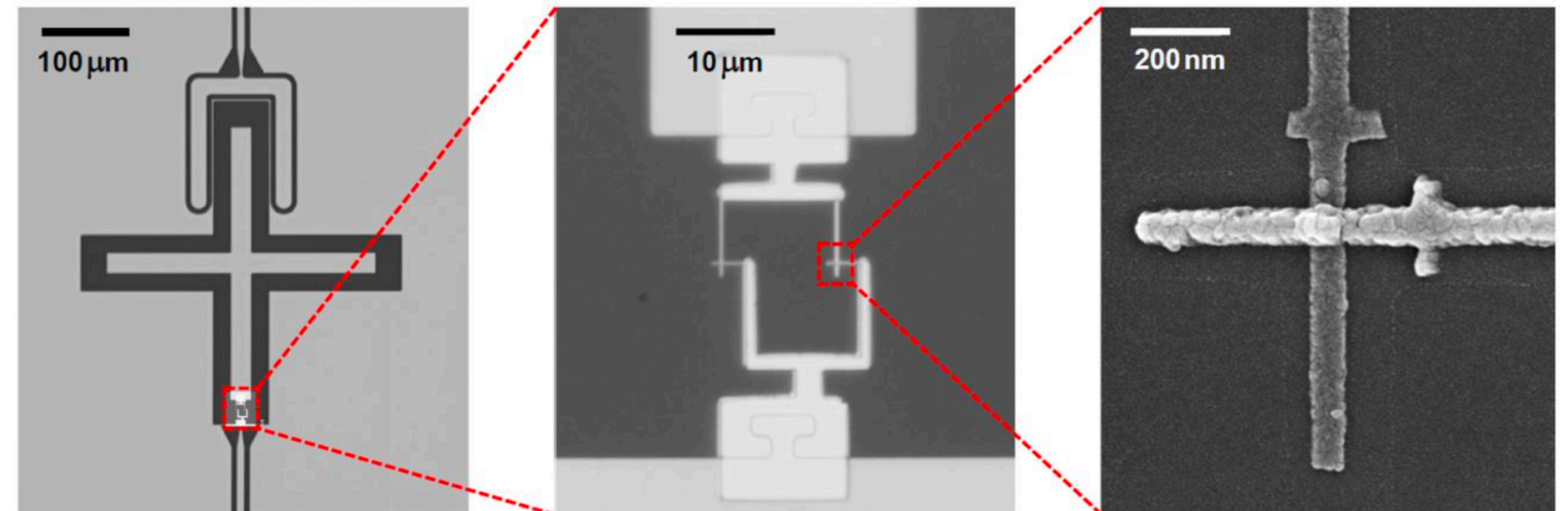
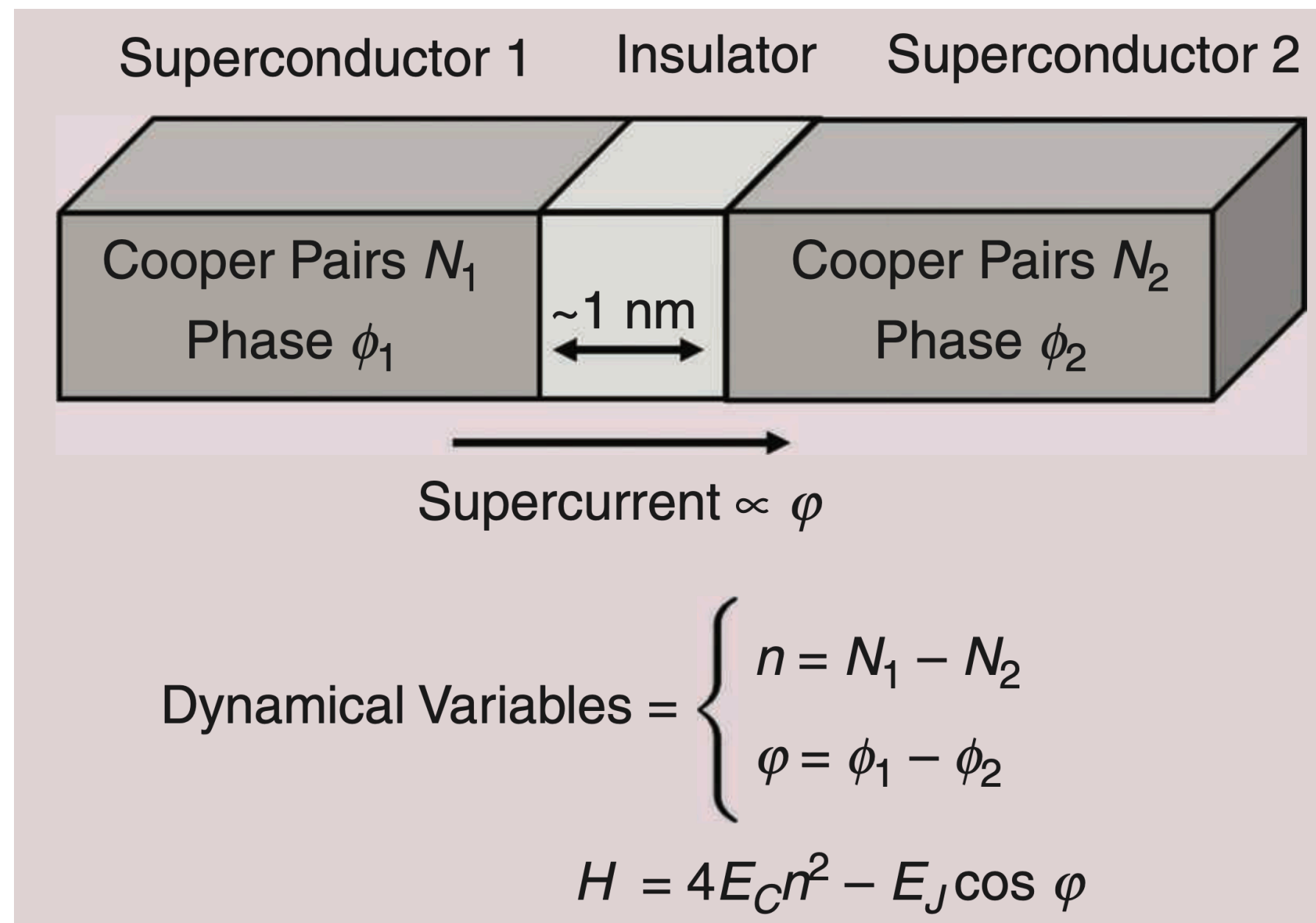
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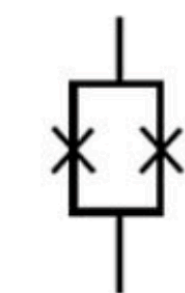
Chao, Yu Gao, Mingjie Jin, Xiaosheng Liu and Xilei Sun, e-Print: 2412.20850 [hep-ph]



Transmon qubit



Transmon qubit
(a)

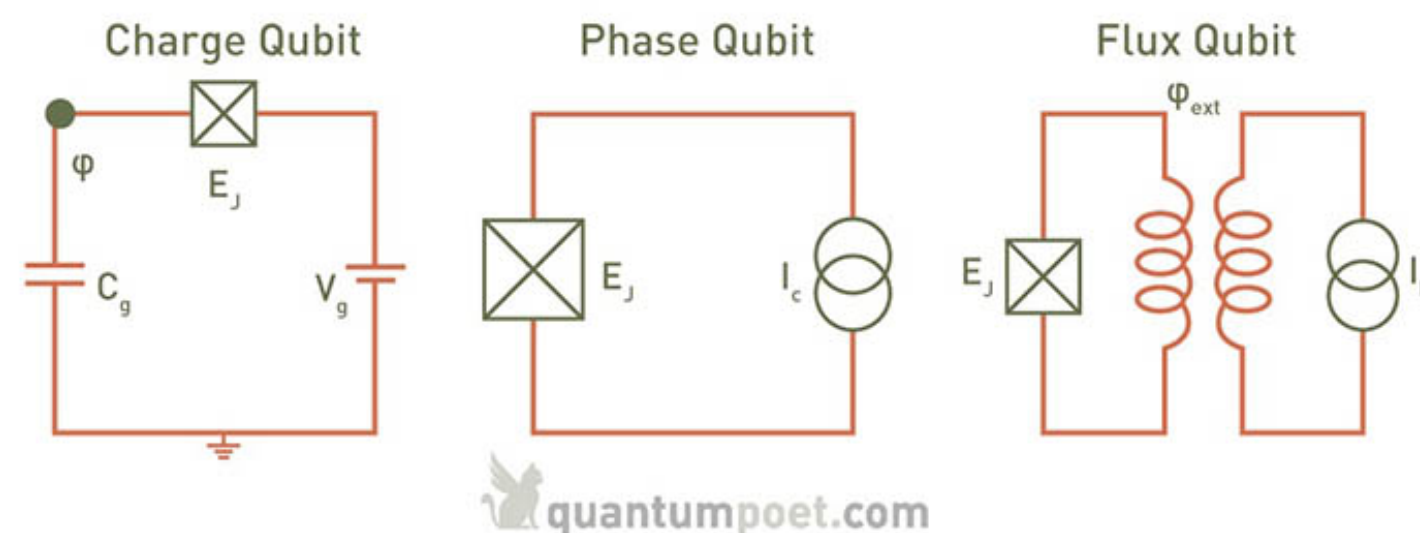


SQUID
(b)



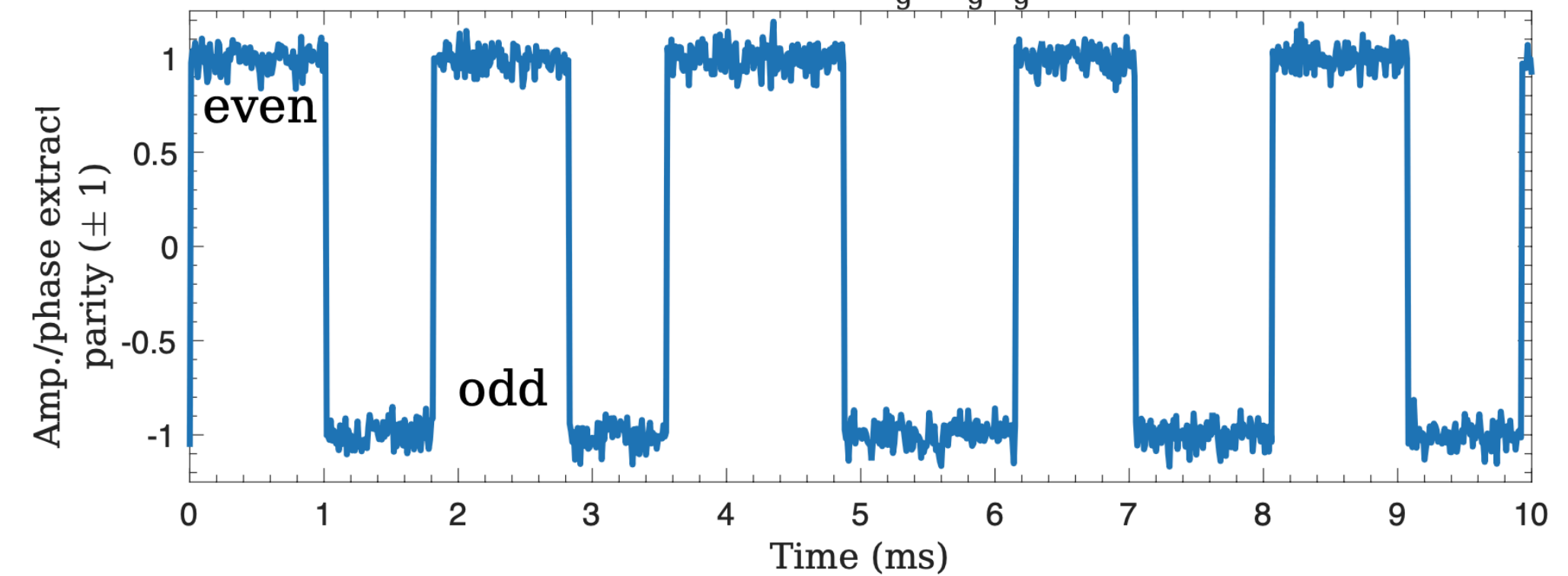
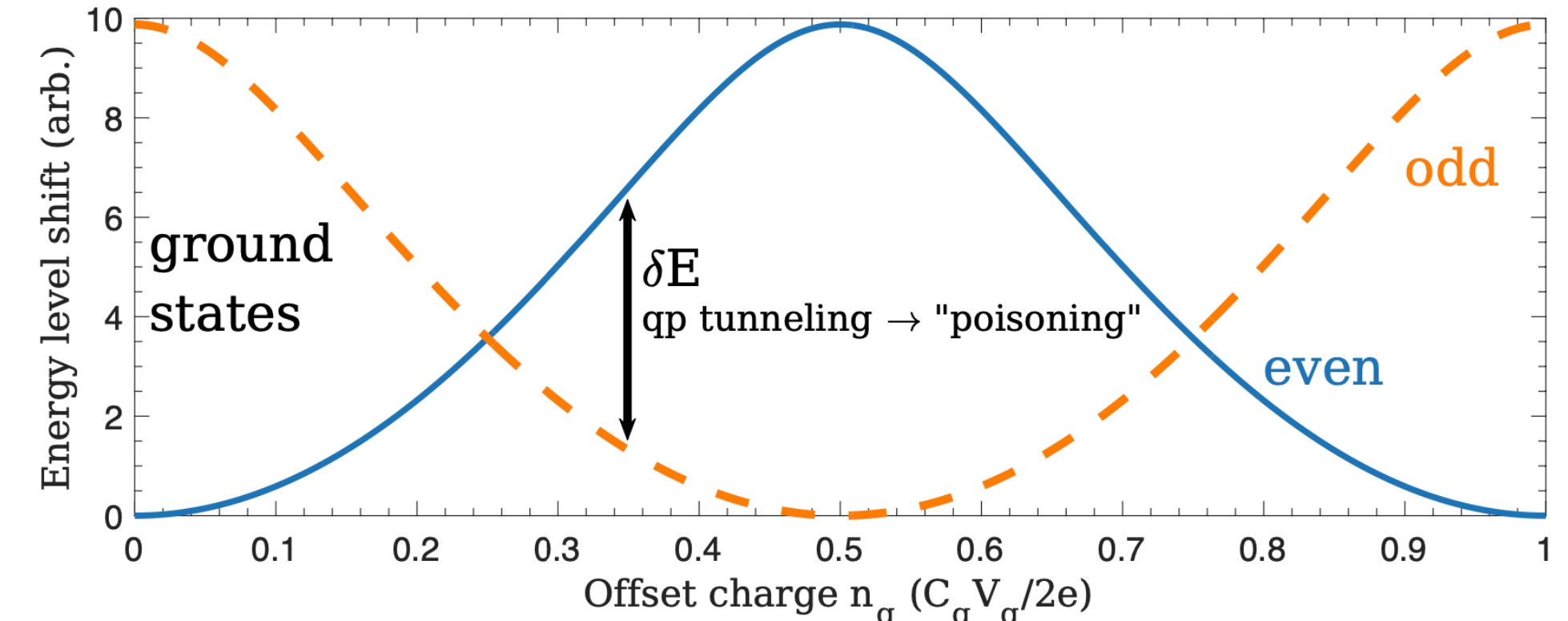
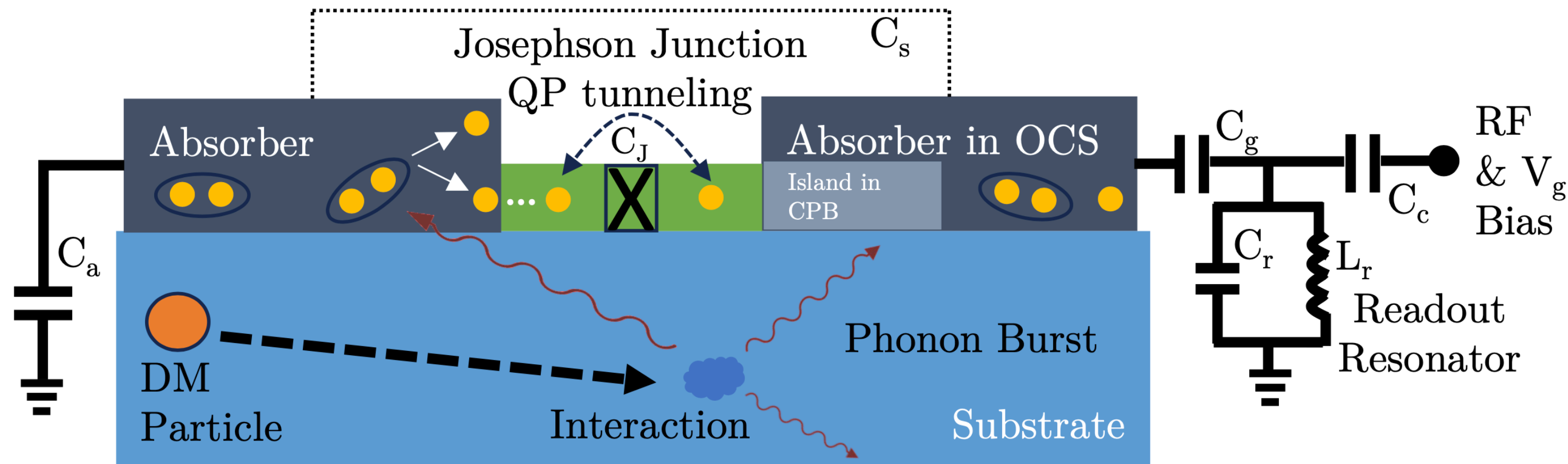
Josephson junction
(c)

The 3 Main Types of Superconducting Qubits



Equivalent to simple mechanical pendulum.

Superconducting qubit as a light DM detector



$$H_{CPB} = 4E_C \left(\hat{n} - n_g + \frac{P-1}{4} \right)^2 - E_j \cos \hat{\phi}$$

E_j : Josephson energy

P : parity dependent shift

$\hat{n}$: number of Cooper pairs on island

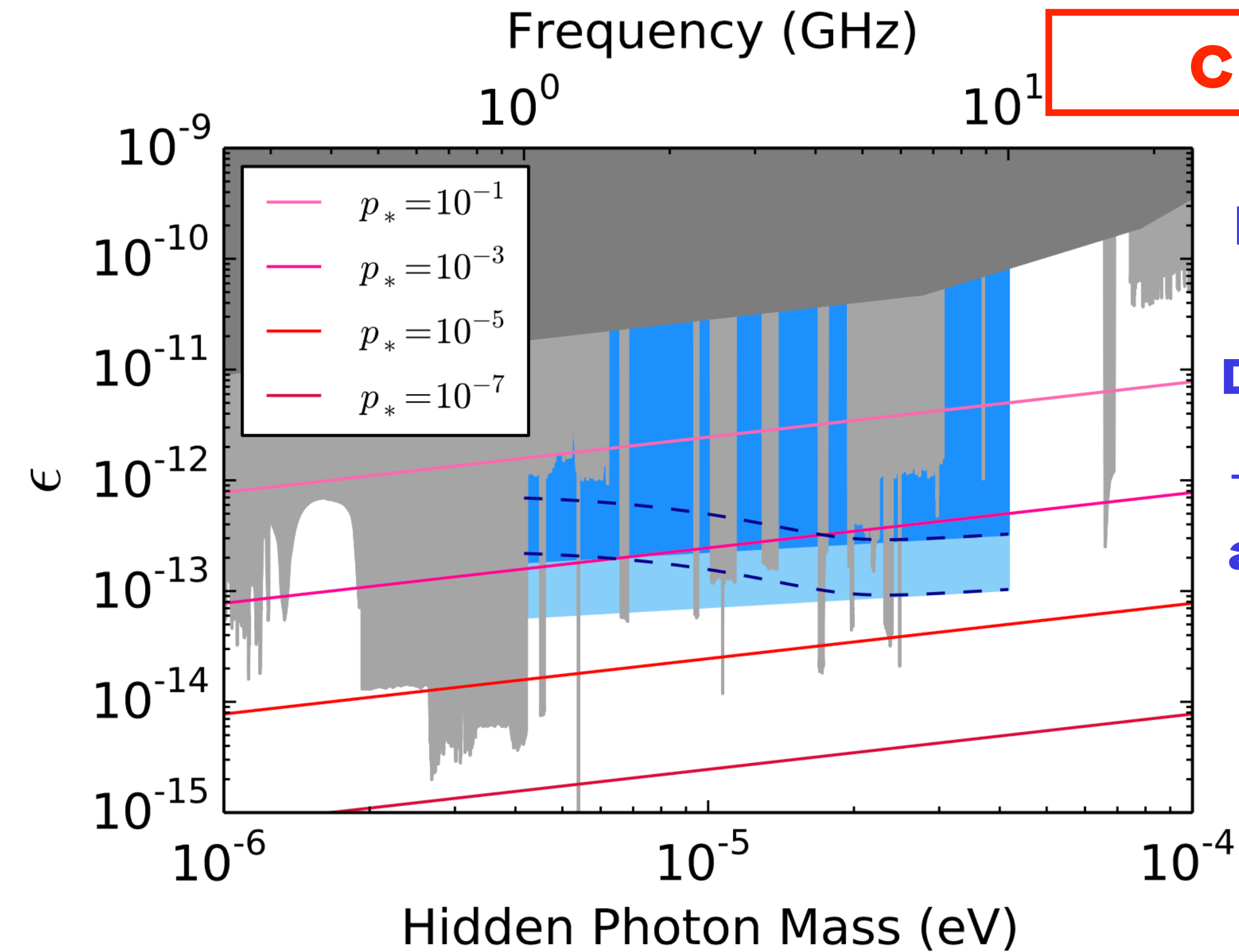
$\hat{\phi}$: phase difference

n_g : offset charge

Plots taken from arXiv: 2405.17192

Transmon qubit as a dark photon detector

Chen, et al., PRL 131, 211001 (2023)



Hamiltonian

$$H_0 = \frac{1}{2}CV^2 - J \cos \theta$$

DP induced electric field $\vec{E}^X = \epsilon \sqrt{2\rho_{\text{DM}}} \vec{n}_X \sin m_x t$

→ Oscillations between the ground state and the excited state!

$$p_{ge}(\tau) \simeq 0.12 \times \kappa^2 \cos^2 \Theta \left(\frac{\epsilon}{10^{-11}} \right)^2 \left(\frac{f}{1 \text{ GHz}} \right) \\ \times \left(\frac{\tau}{100 \text{ } \mu\text{s}} \right)^2 \left(\frac{C}{0.1 \text{ pF}} \right) \left(\frac{d}{100 \text{ } \mu\text{m}} \right)^2 \\ \times \left(\frac{\rho_{\text{DM}}}{0.45 \text{ GeV/cm}^3} \right),$$

Transmon qubit constraint on bosonic DM

Quasiparticle excitation:

$$\frac{dn_{qp}}{dt} = -\Gamma_R - \Gamma_T + \Gamma_G \quad \bullet \text{ Phys. Rev. Lett. 117, 117002 (2016).}$$

Γ_R : **Recombination rate**

$$\Gamma_R = \bar{\Gamma} n_{PQ}^2$$

Γ_T : **Trapping rate**

$$\Gamma_T = \bar{\Gamma}' n_{PQ}$$

Γ_G : **Generating rate induced by DM absorption or scattering!**

In equilibrium:

$$\Gamma_G = \bar{\Gamma} n_{QP}^2 + \Gamma_T n_{QP}$$

Transmon qubit constraint on bosonic DM

The number of phonon per unit time per unit mass:

$$R = \frac{\pi \sigma_{\chi} n_{\chi}}{\rho_T \mu^2} \int d^3 v f_{\chi}(v) \int \frac{d^3 q}{(2\pi)^3} F_{\text{med}}^2(q) S(q, \omega_q)$$

Full phonon energy:

$$E = \int d\omega \omega \times \frac{dR}{d\omega} \times M \times T$$

Number of quasiparticles produced:

$$N = \frac{\varepsilon}{2\Delta} \omega R M T$$

Produced quasi particle density per unit time:

$$n'_{qp} = N/V/T = \frac{\varepsilon}{2\Delta} \omega R \rho_T = \frac{\varepsilon}{2\Delta} \int d\omega \omega \frac{dR}{d\omega} \times \rho_T$$

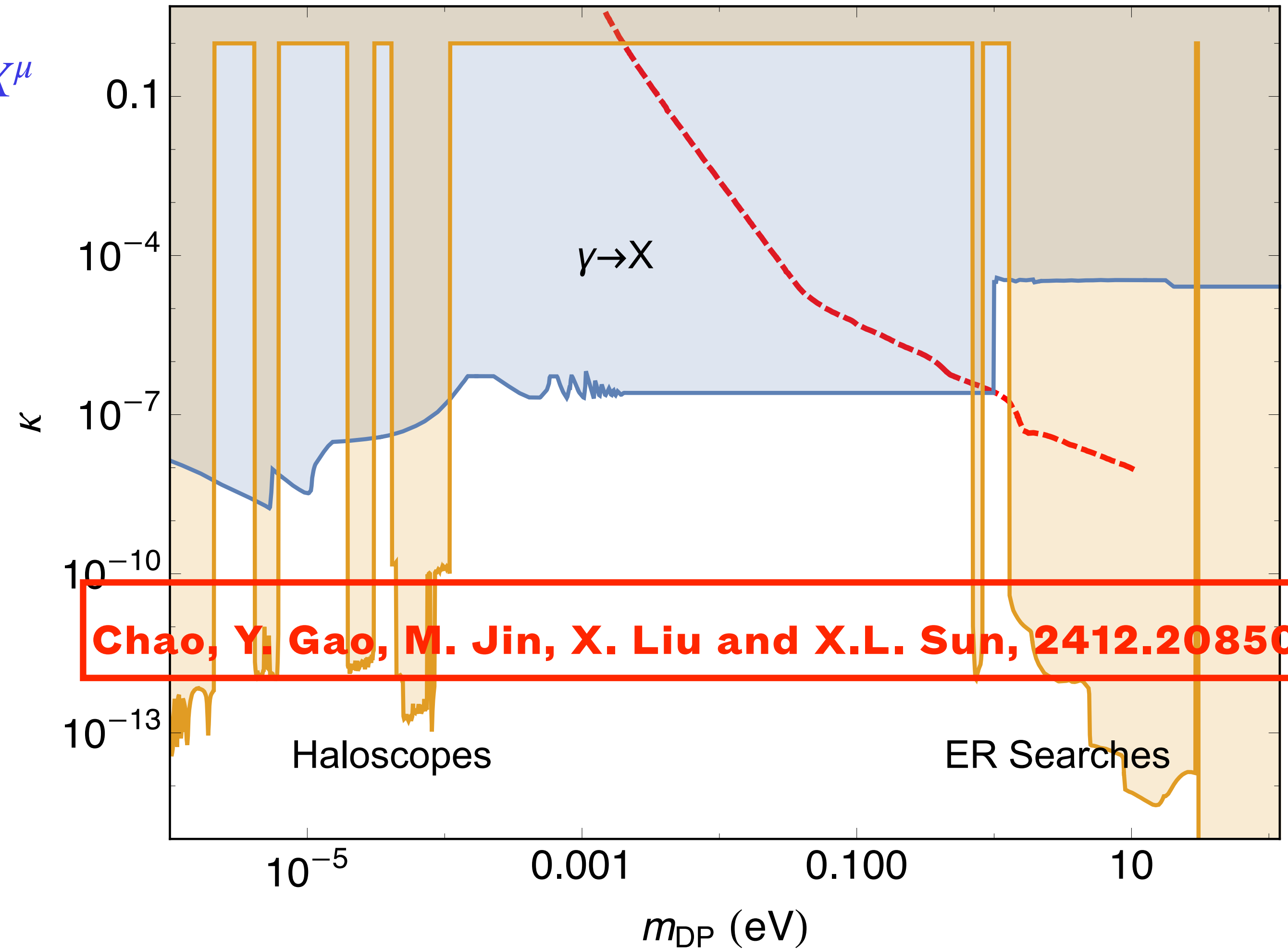
Transmon qubit constraint on the DP

$$\mathcal{L}_{\text{DP}} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\varepsilon}{2}F_{\mu\nu}X^{\mu\nu} + \frac{1}{2}m_X^2 X_\mu X^\mu$$

$$R_{\text{DP}} = \frac{\rho_{\text{DP}}}{\rho_T} \kappa^2 \text{Im} \left[-\frac{1}{\varepsilon(\omega)} \right]$$

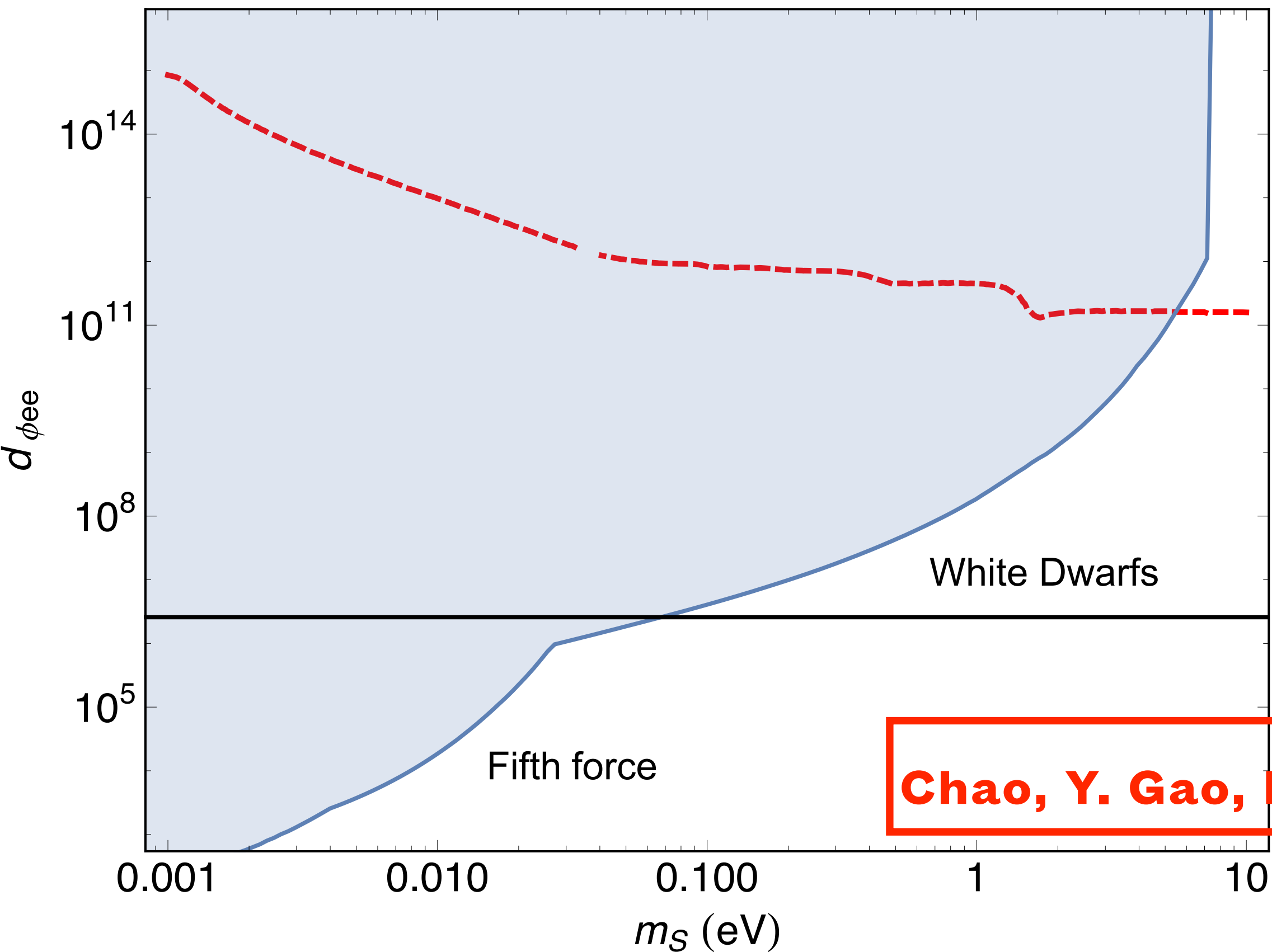
$$n_{qp} < 0.04 \mu\text{m}^{-3}$$

D. Riste et al., Nature Communications 4 (2013), 10.1038/ncomms2936

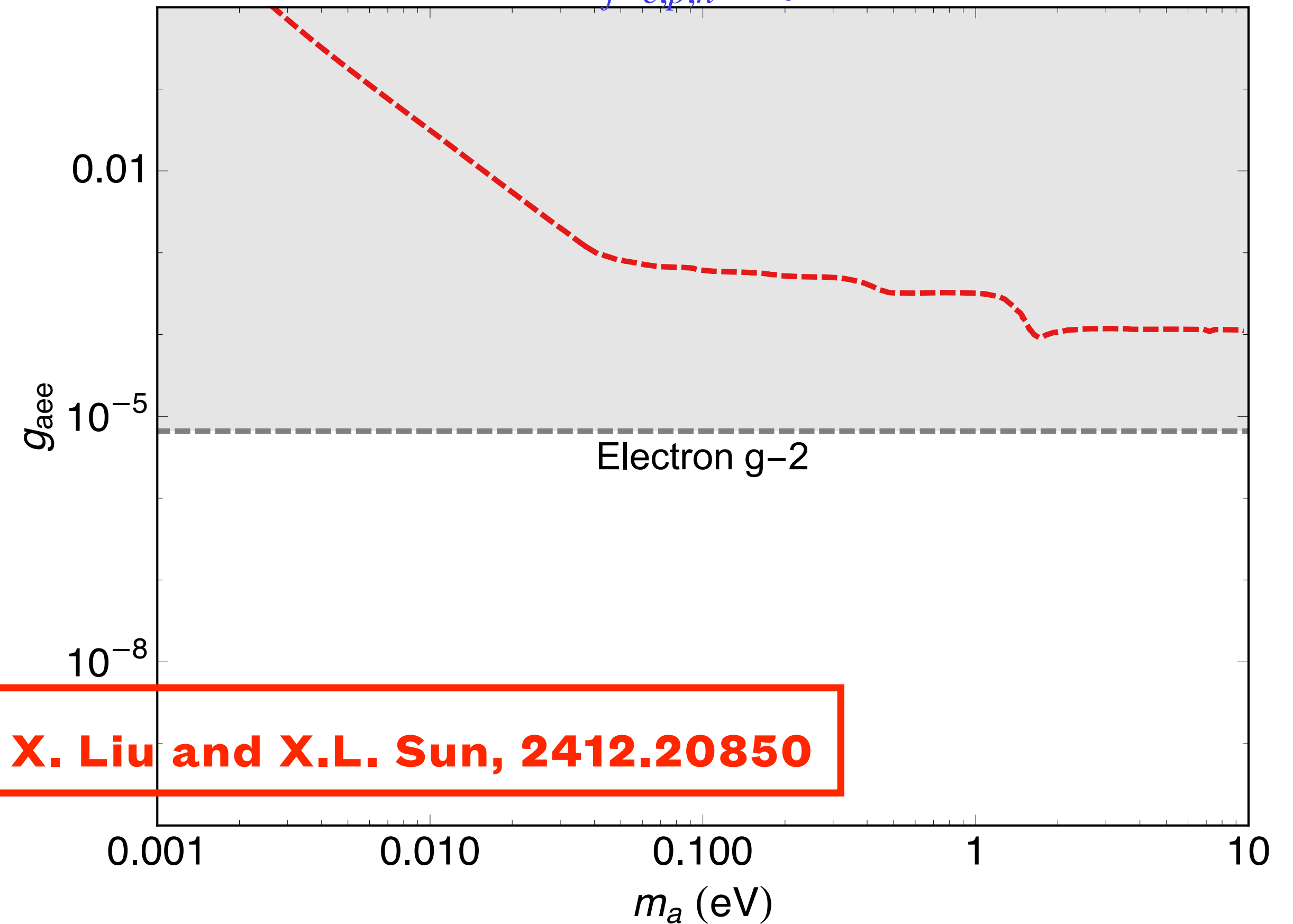


Transmon qubit constraint on Scalar and ALP DM

$$\mathcal{L} \supset \sum_{f=e,p,n} d_{\phi ff} \phi \bar{f} f$$



$$\mathcal{L} \supset \sum_{f=e,p,n} \frac{g_{aff}}{2m_f} \partial_\mu a \bar{f} \gamma^\mu \gamma^5 f$$



Chao, Y. Gao, M. Jin, X. Liu and X.L. Sun, 2412.20850

Summary

New strategy of detecting light dark matters, is a topic deserving deep investigation.

A. Possible gravitational signal from dark photon oscillation is investigated.

B. Detecting light dark matters using superfluid is investigated.

C. New strategy for the DD of bosonic DM using superconducting qubit is investigated.

Thank you for your attention!