

An Efficient Learning Method to Connect Observables

Hang Yu
University of Tsukuba
3/3

Frontiers in NLEFT, Beihang



Collaborator: Takayuki Miyagi

P. Cook et.al. arxiv: 2401.11694 [cs.LG] (2024)

Parametric Matrix Model

Data Driven

Gaussian Process
Neural Network

...

Our method

Model Driven

Eigenvector Continuation

D. Frame et.al. PRL 121 032501 (2018)

Reduce order models

it/6231137 [nucl-th] 2 Mar 2025

HY and T. Miyagi [soonTM(for real)]

An Efficient Learning Method to Connect Observables

Hang Yu^{1,*} and Takayuki Miyagi^{1,†}

¹Center for Computational Sciences, University of Tsukuba, Tsukuba, Ibaraki 305-8577, Japan

Constructing fast and accurate surrogate models is a key ingredient for making robust predictions in many topics. We introduce a new model, the Multiparameter Eigenvalue Problem (MEP) emulator. The new method connects emulators and can make predictions directly from observables to observables. We present that the MEP emulator can be trained with data from Eigenvector Continuation (EC) and Parametric Matrix Model (PMM) emulators. A simple simulation on a one-dimensional lattice confirms the performance of the MEP emulator. Using ^{28}O as an example, we also demonstrate that the predictive probability distribution of the target observables can be easily obtained through the new emulator.

A major theme in physics is to discover and understand phenomena. The gold standard for theoretical work is to explain experimental measurements and observations while also predicting unseen results. However, complexity in theory sometimes prevents all essential parameters from being uniquely determined. These parameters often appear in important constituents, such as a coupling strength in a Hamiltonian. Due to the variations of the parameters, calibrating the parameters and making a prediction often need to be done separately. These separate procedures work well when the underlying theory is simple enough and has only almost uniquely determined parameters. With the advance of physics, theoretical models tend to become complicated, ranging from cosmology models that have many parameters to be optimized [1,2] to *almost* parameter-free theory of the underlying strong forces that is difficult to solve dynamically.

Low-energy nuclear physics lies in the intersection of computationally demanding and multiparametric. Delicate interplays of two- and three-body interactions have been one of the barriers to our theoretical progress, with many parameters called low-energy constants (LECs) appearing in the same order of the underlying effective field

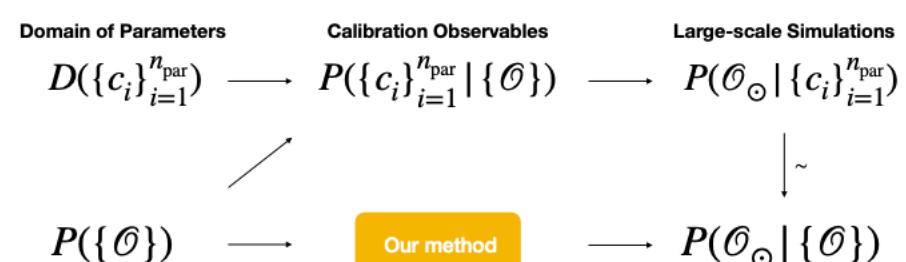


Figure 1. Workflow of our method compared with existing statistical procedures

replace complicated statistical workflow found, for example, in Ref. [1, 2, 11–13]. We can directly obtain the conditional probability $P(\mathcal{O}_0 | \{\mathcal{O}\})$ for the target \mathcal{O}_0 under the calibration data $\{\mathcal{O}\}$, once the probability density $P(\{\mathcal{O}\})$ of calibration observables is given, without going through the complicated workflow that is also dependent on the domain of parameters [13]. Our emulator has its root in the Ritz approximation method, enabling possible generalization to all problems that need multiple eigenvalues/parameters.

We first briefly discuss our motivation. In many sit-

A Short Review of EC

When Perturbation Theory met Variational Principle...

$$H(c) = H_0 + cV$$

$$\psi(c) = \psi_0 + c\psi_1 + c^2\psi_2 + \dots$$

$$E_{\text{gs}}(c) \leq \frac{\langle \psi | H(c) | \psi \rangle}{\langle \psi | \psi \rangle}$$

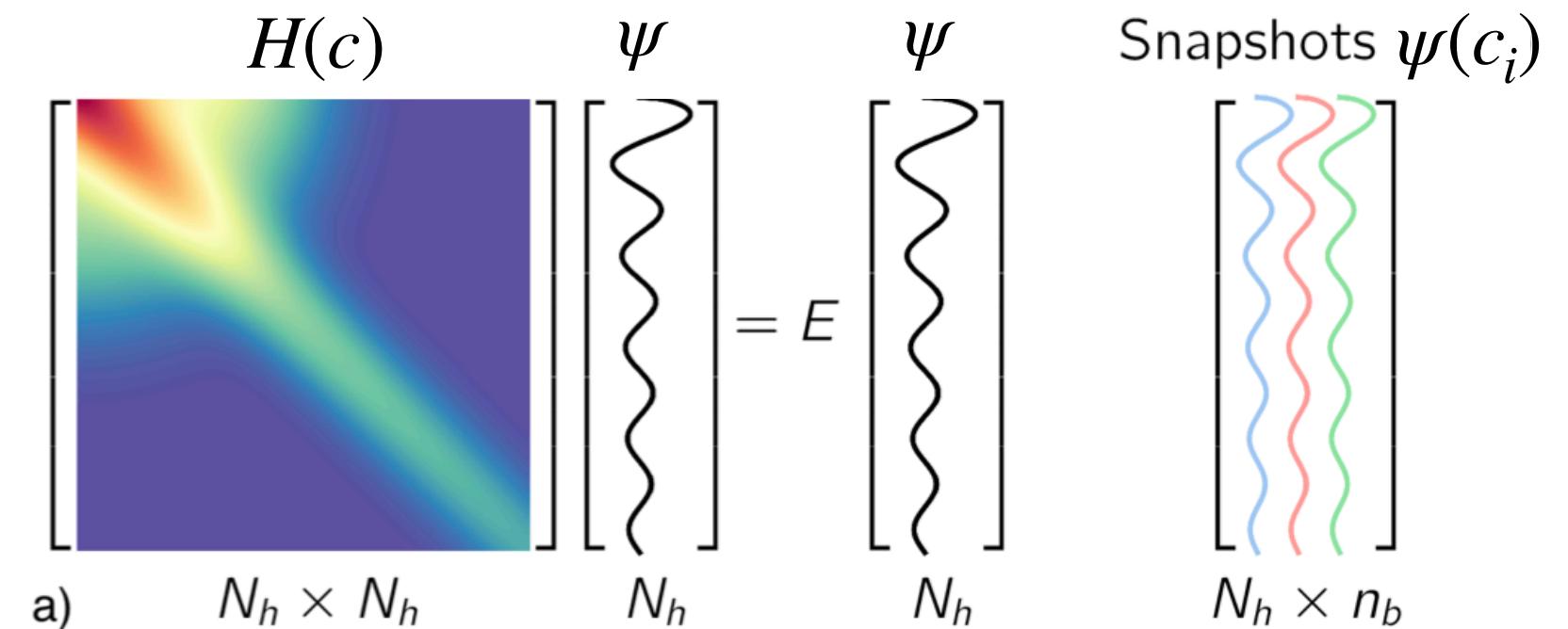
$$S(\psi) = \text{span}(\{\psi_0, \psi_1, \psi_2, \dots\})$$

Eigenvector Continuation

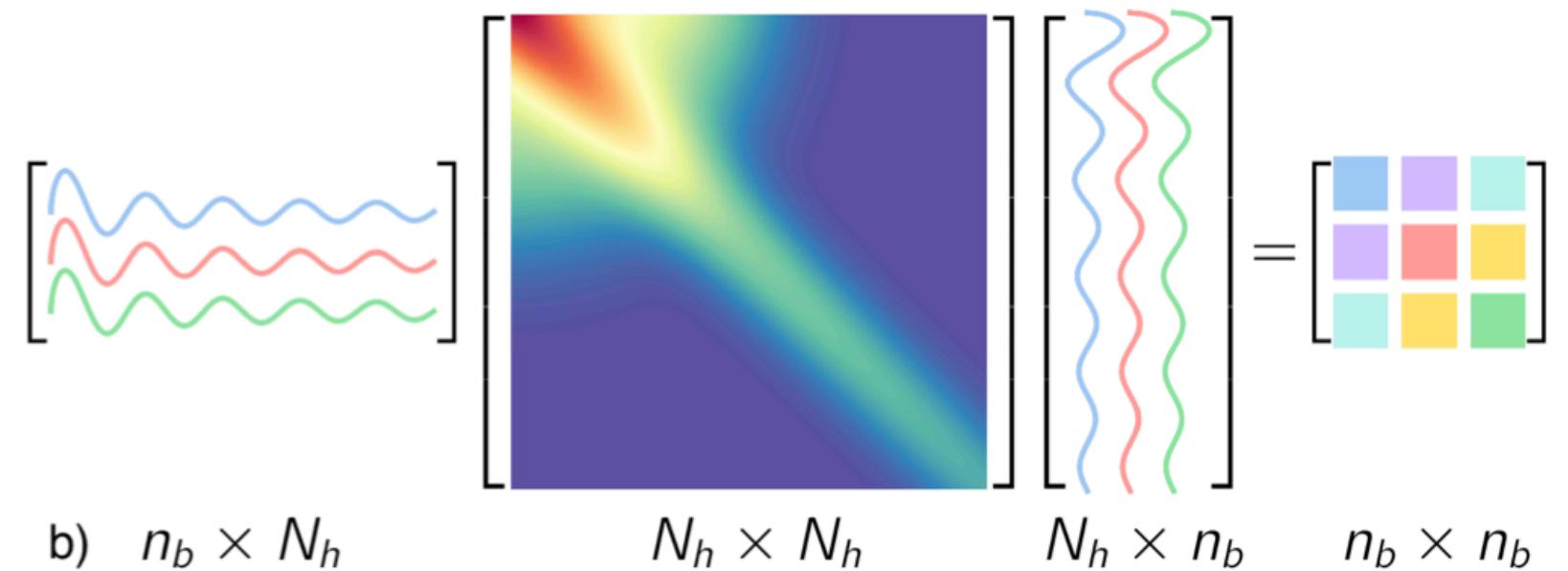
$$(\mathcal{H}(c) - \tilde{E}\mathcal{N})\tilde{\psi}(c) = 0$$

$$\mathcal{H}(c)_{ij} = \langle \psi(c_i) | H(c) | \psi(c_j) \rangle$$

$$\mathcal{N}_{ij} = \langle \psi(c_i) | \psi(c_j) \rangle$$



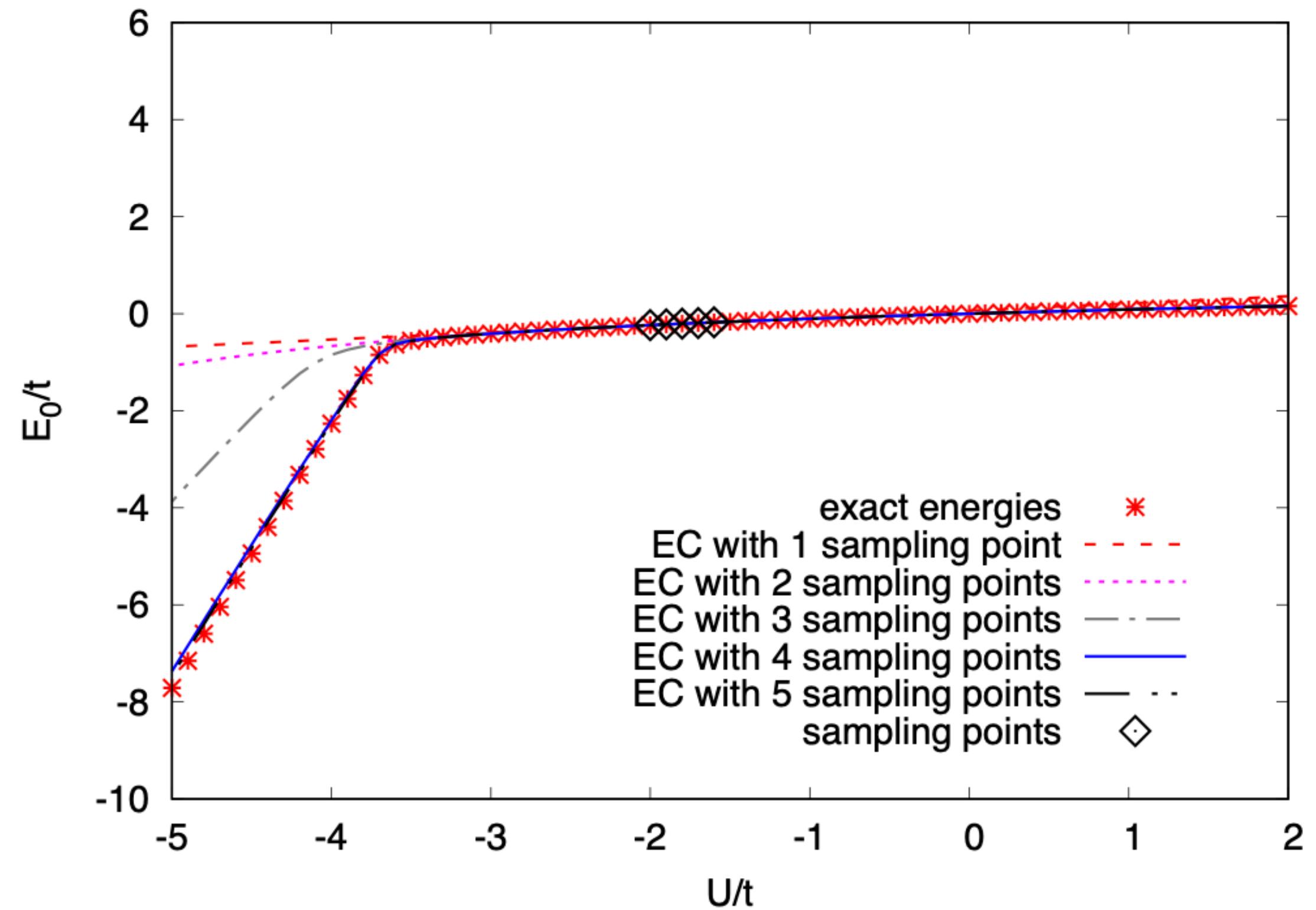
Projection (after orthonormalizing snapshots) $\mathcal{N} = I$



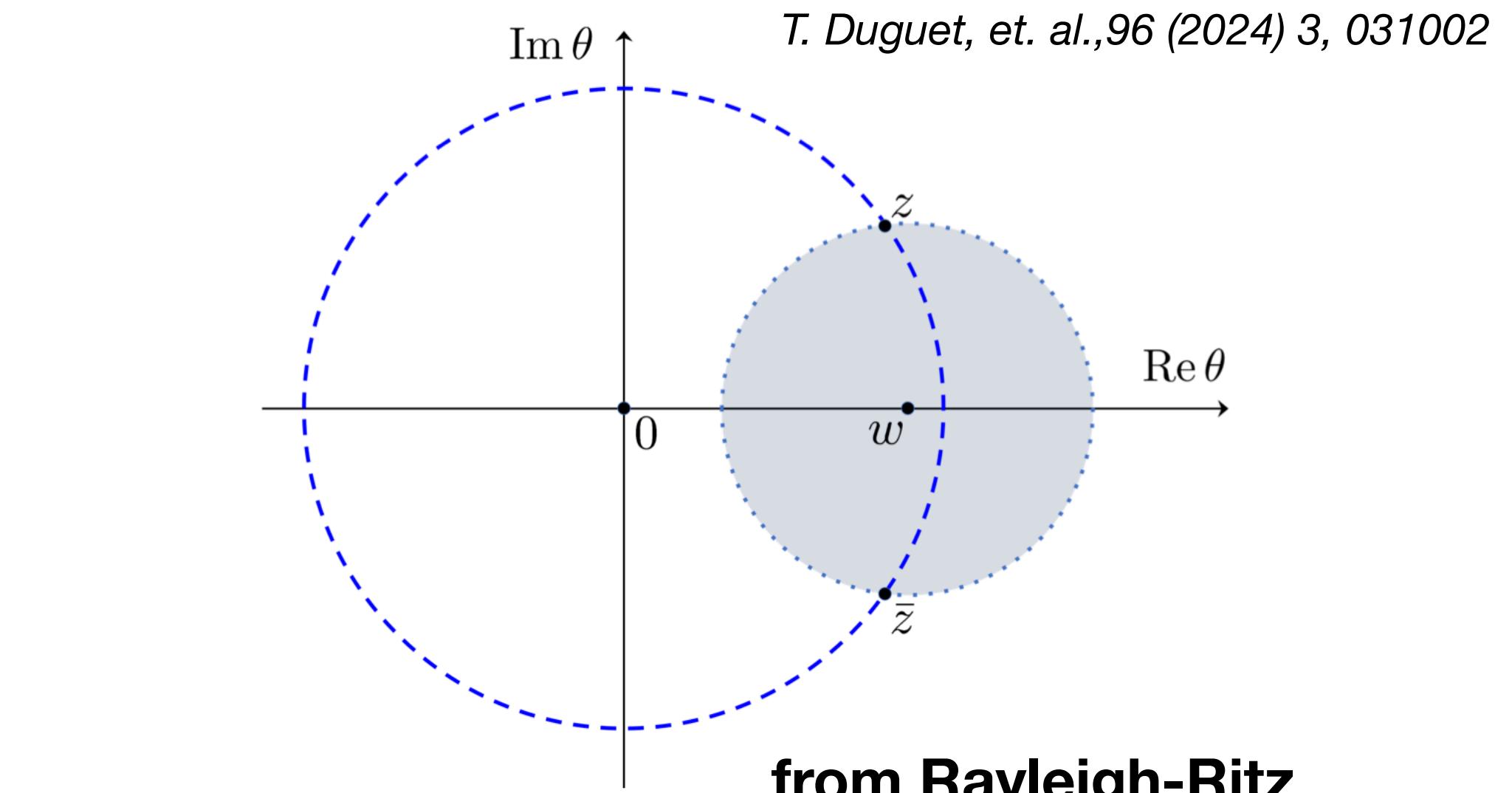
Emulation ($E \approx \tilde{E}$)

$$\mathcal{H}(c)\tilde{\psi}(c) = \tilde{E}\tilde{\psi}(c)$$

$$\begin{bmatrix} \text{color squares} \end{bmatrix} \begin{bmatrix} \bullet \\ \diamond \\ \triangledown \end{bmatrix} = \tilde{E} \begin{bmatrix} \bullet \\ \diamond \\ \triangledown \end{bmatrix}$$

b

D. Frame et.al. PRL 121 032501 (2018)

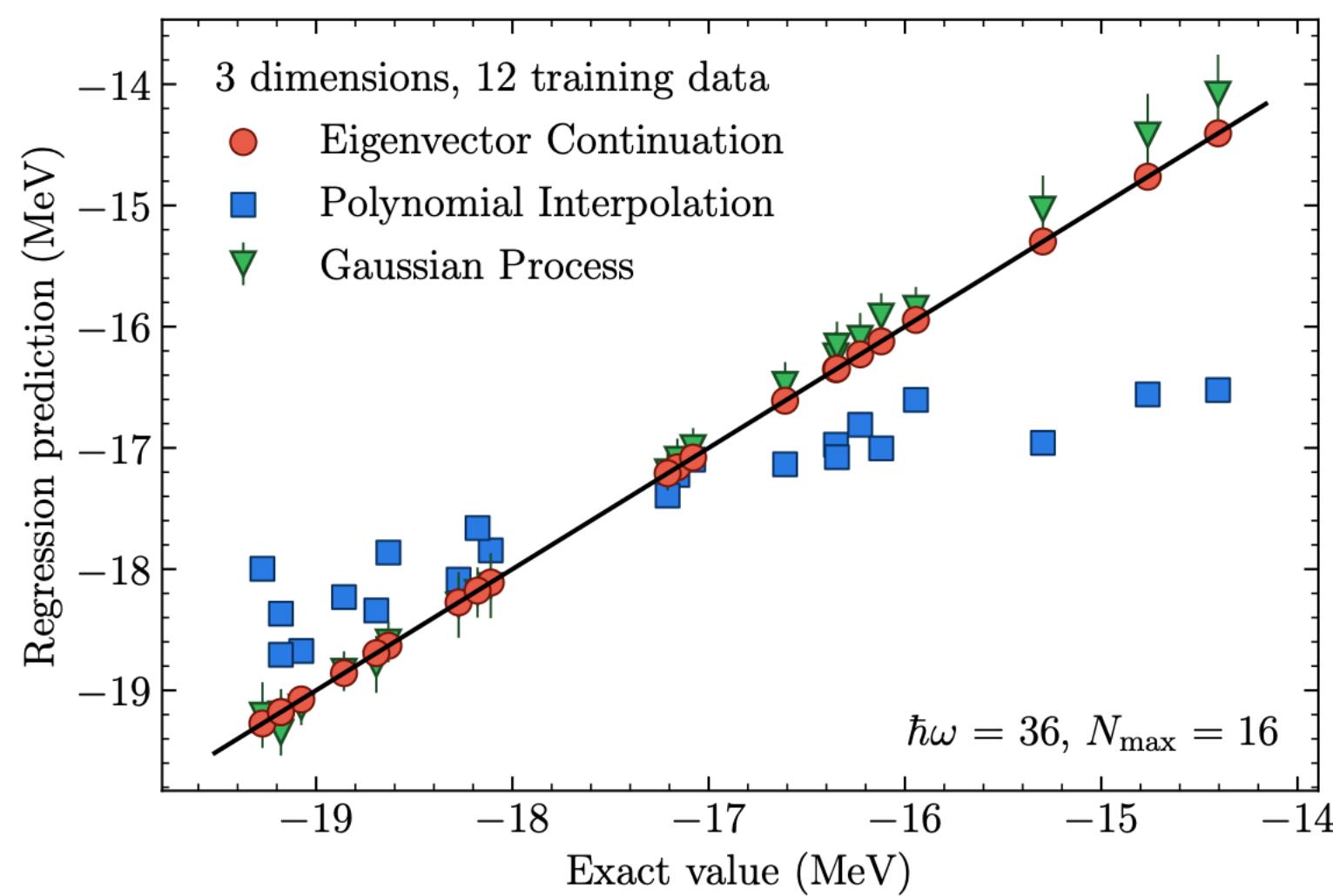
Bypass convergence radii

and requires very few snapshots

$$S(\psi) = \text{span}(\{\psi_0, \psi_1, \psi_2, \dots\})$$

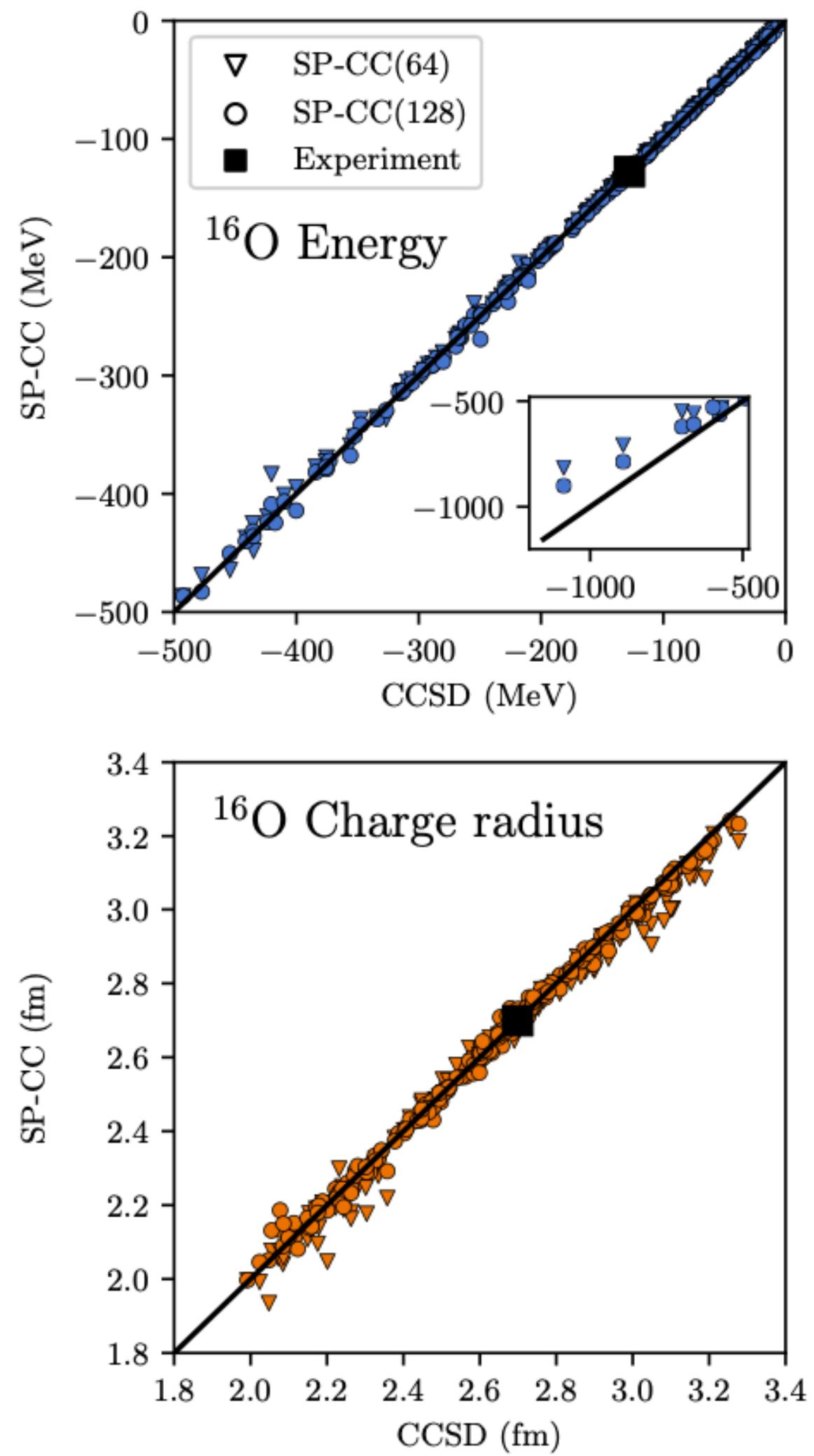
from Perturbation Theory

Examples



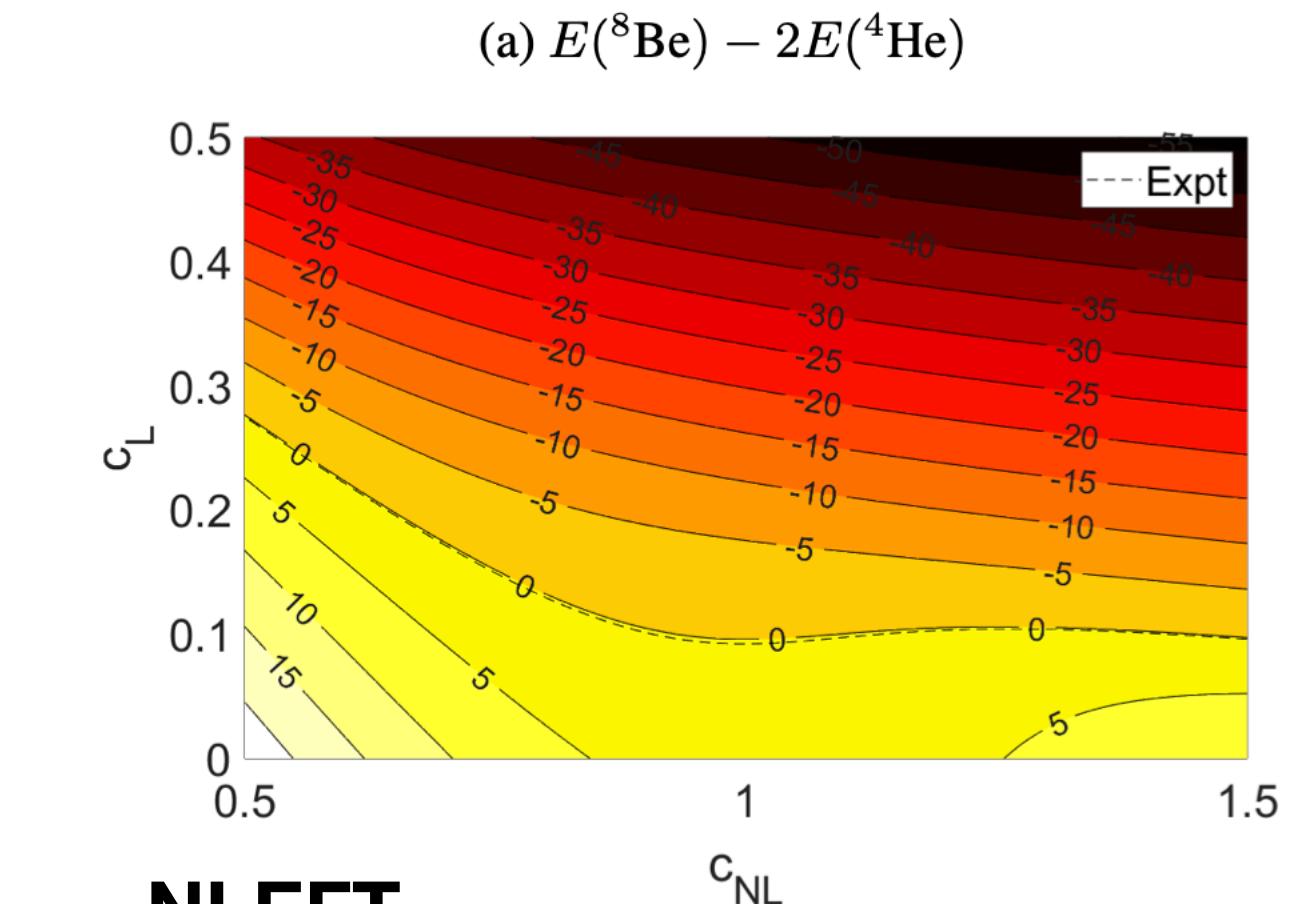
NCSM

S.Koenig et.al. PLB 810 135814 (2019)



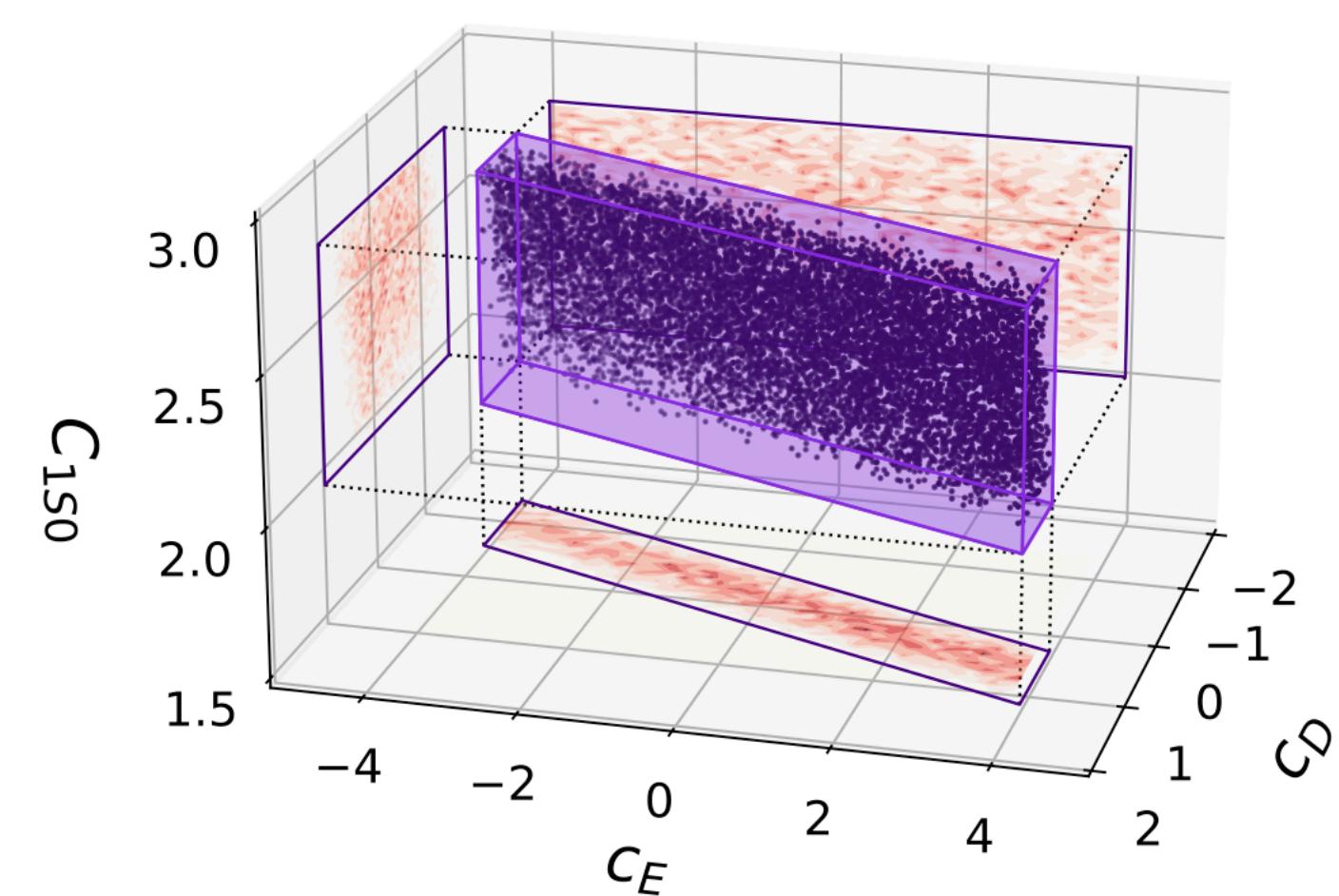
Coupled-Cluster

A. Ekstrom, et.al., PRL 123 252501 (2018)



NLEFT

A. Sarkar, et.al., PRL 131 242503 (2023)



Uncertainty Quantification

W.G. Jiang et.al., PRC 109 064314 (2024)

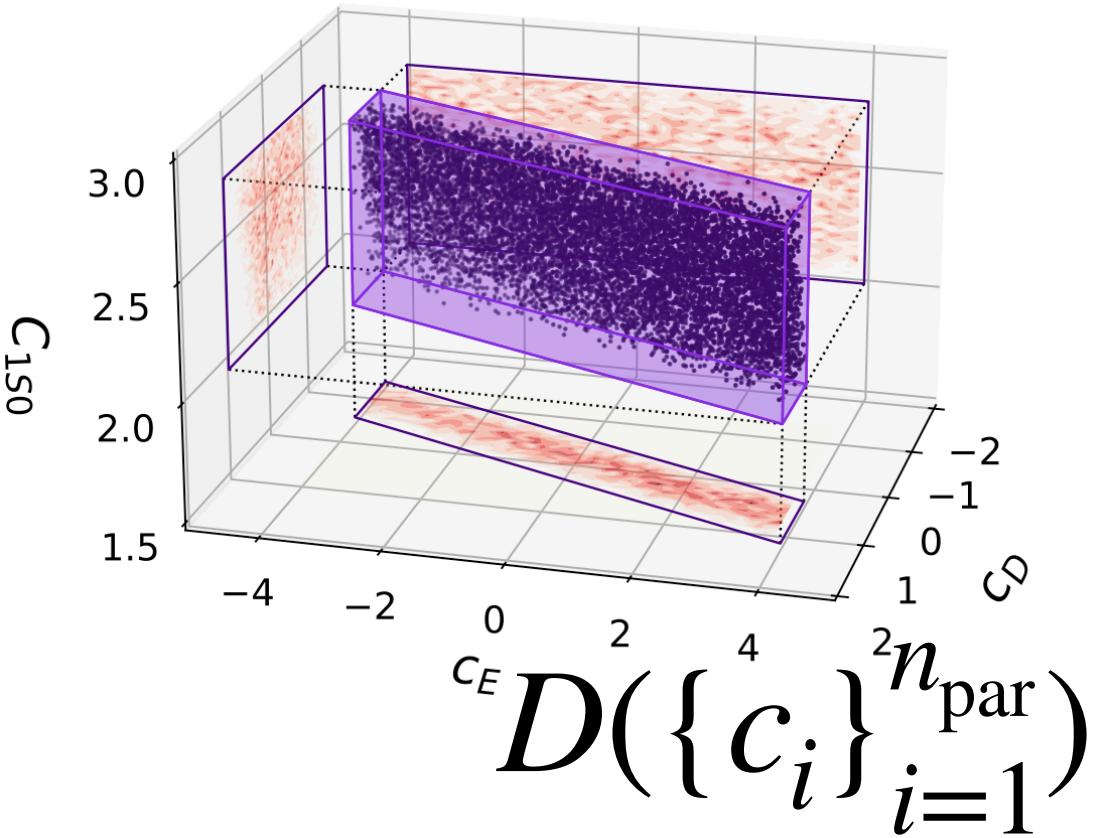
8188 ~ 8218 samples of LEC domain

Motivation

Coupling constants (LECs) are not observables...

$$P(\mathcal{O}_i) \longrightarrow P(c_j)$$

Spin polarized n-d scattering



To search for amplified effects
of N(2-4)LO χ – EFT...

Obs

EFT

Emulators

Uncertainty
Quanti.



:

See talk by H. Krebs



+ 10 – ish Threebody LECs

Nuclear fusion & fission

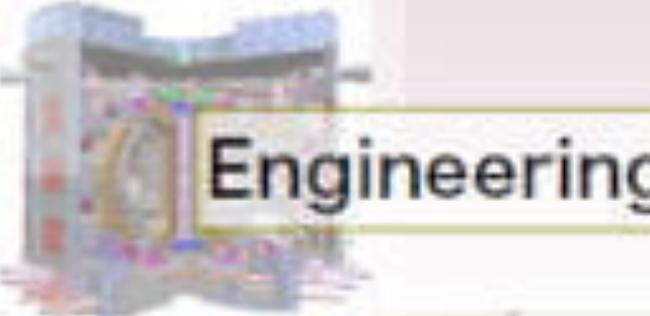
Nucleosynthesis

Neutron star



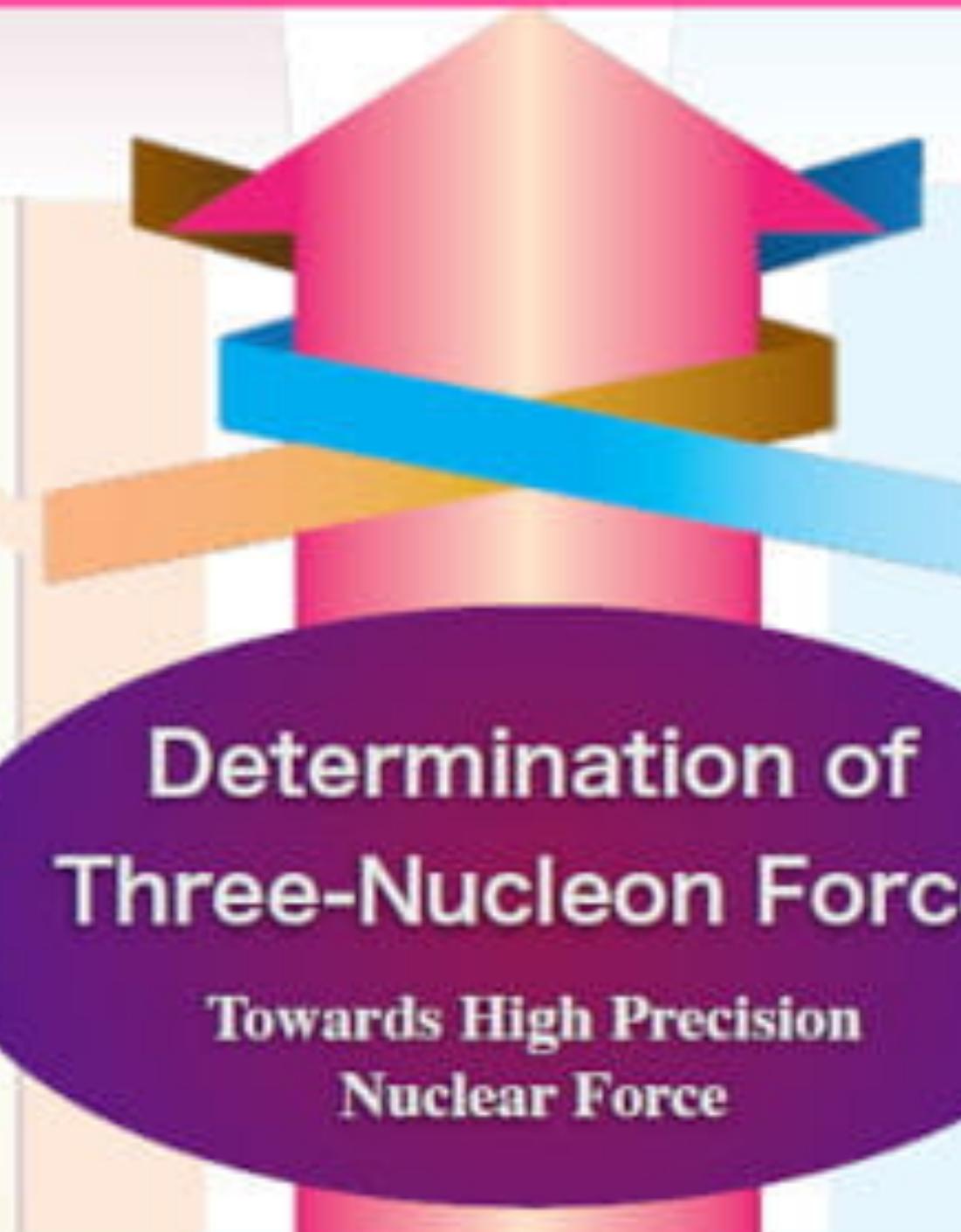
Nuclear Medicine

RI production



Engineering

Applied Science Evolution of Nuclear Data



Determination of
Three-Nucleon Force
Towards High Precision
Nuclear Force

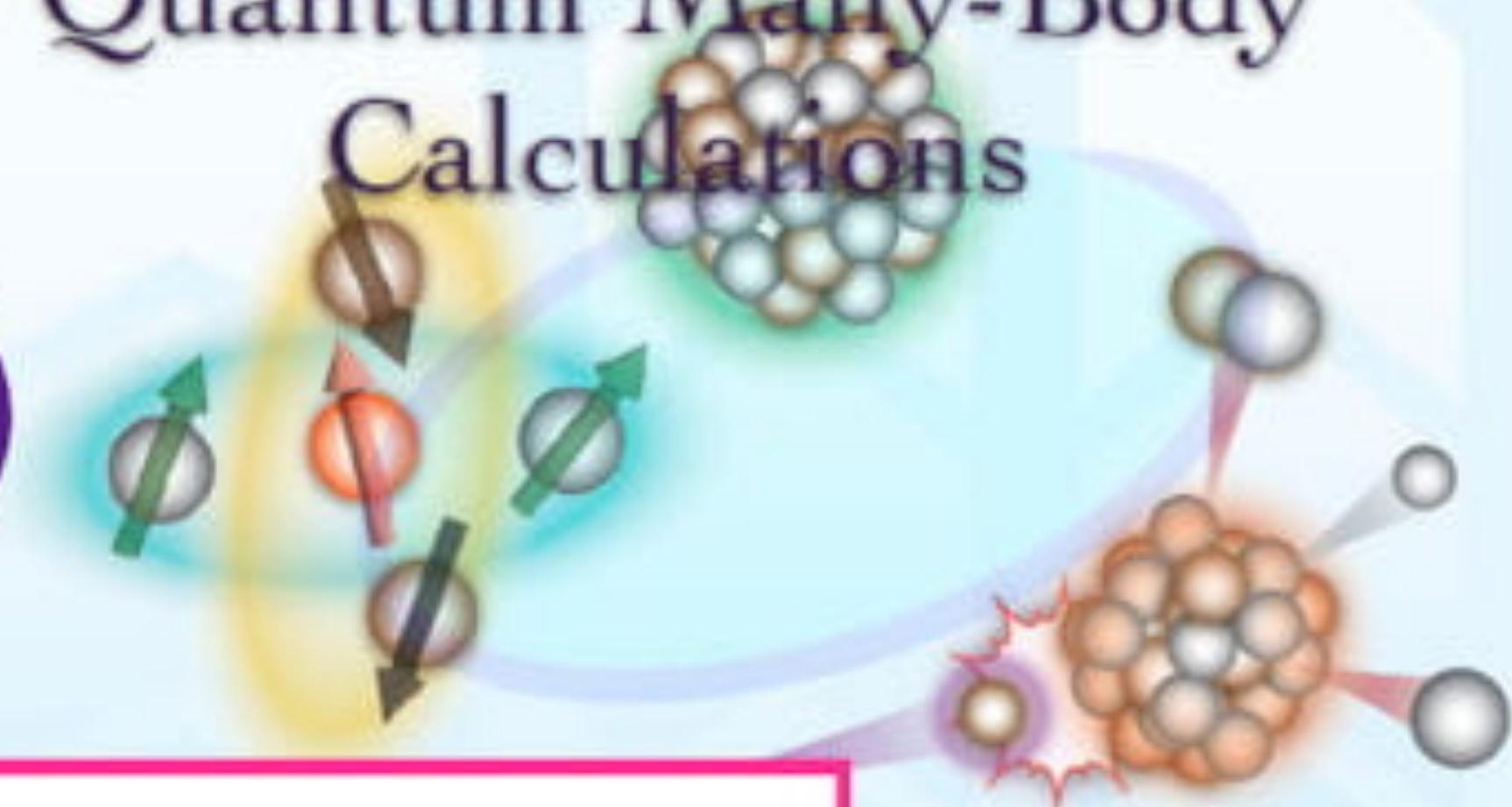
Polarization Experiment - Few-Nucleon Systems-



Ultra Cold Atom Experiment

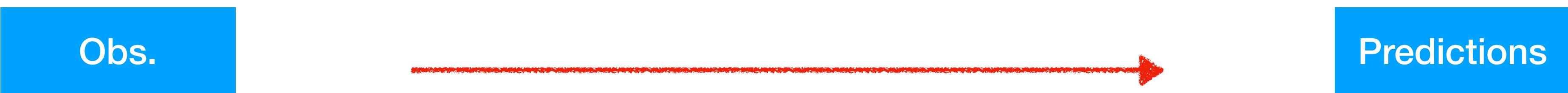
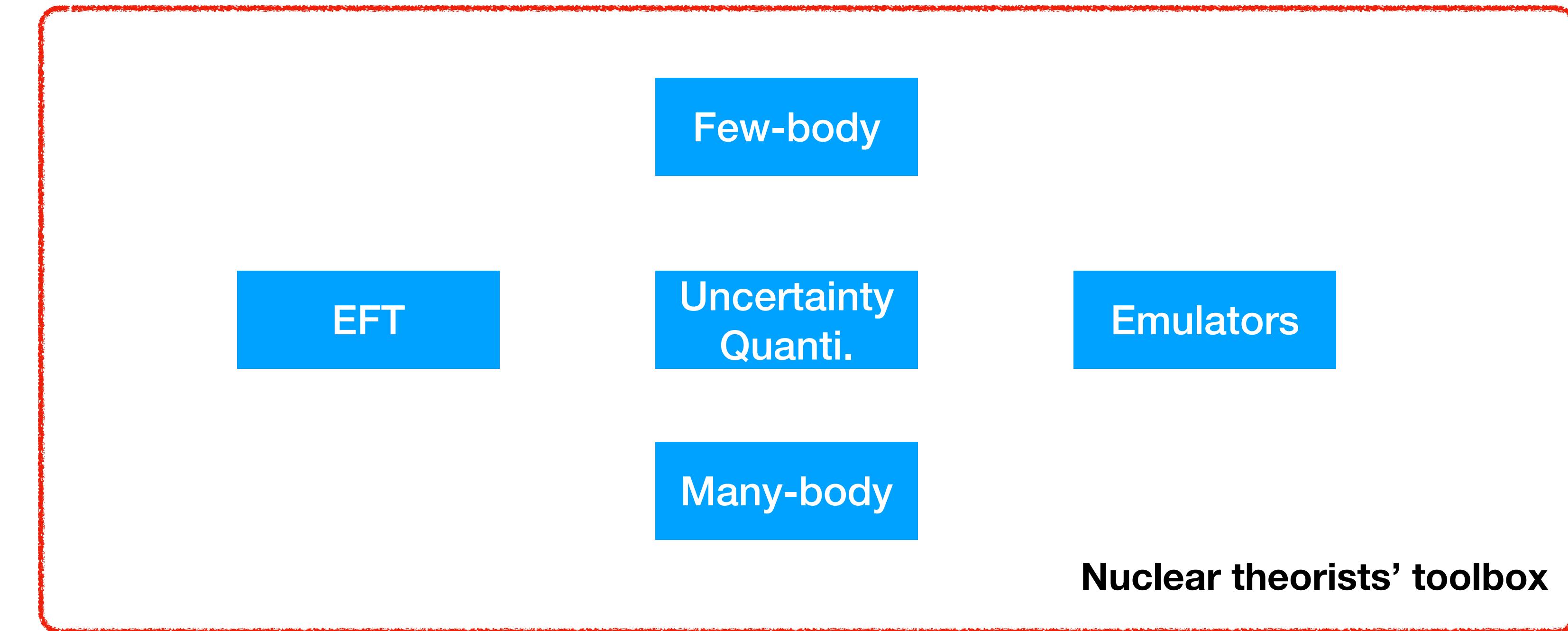


High-Accuracy Quantum Many-Body Calculations



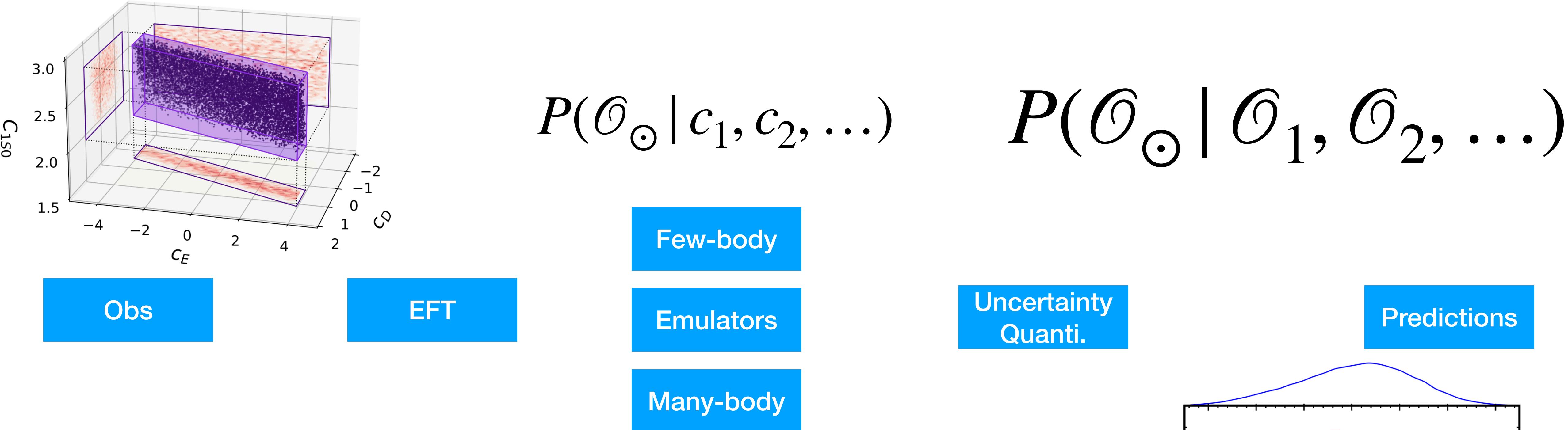
Fundamental Science Descriptions of Nuclei from First Principles

Establishment of Quantum Many-Body Simulation Tool of Nuclear Phenomena
with High-predictive Power



$$P(\mathcal{O}_\odot | \mathcal{O}_1, \mathcal{O}_2, \dots)$$

statistically connecting observables...

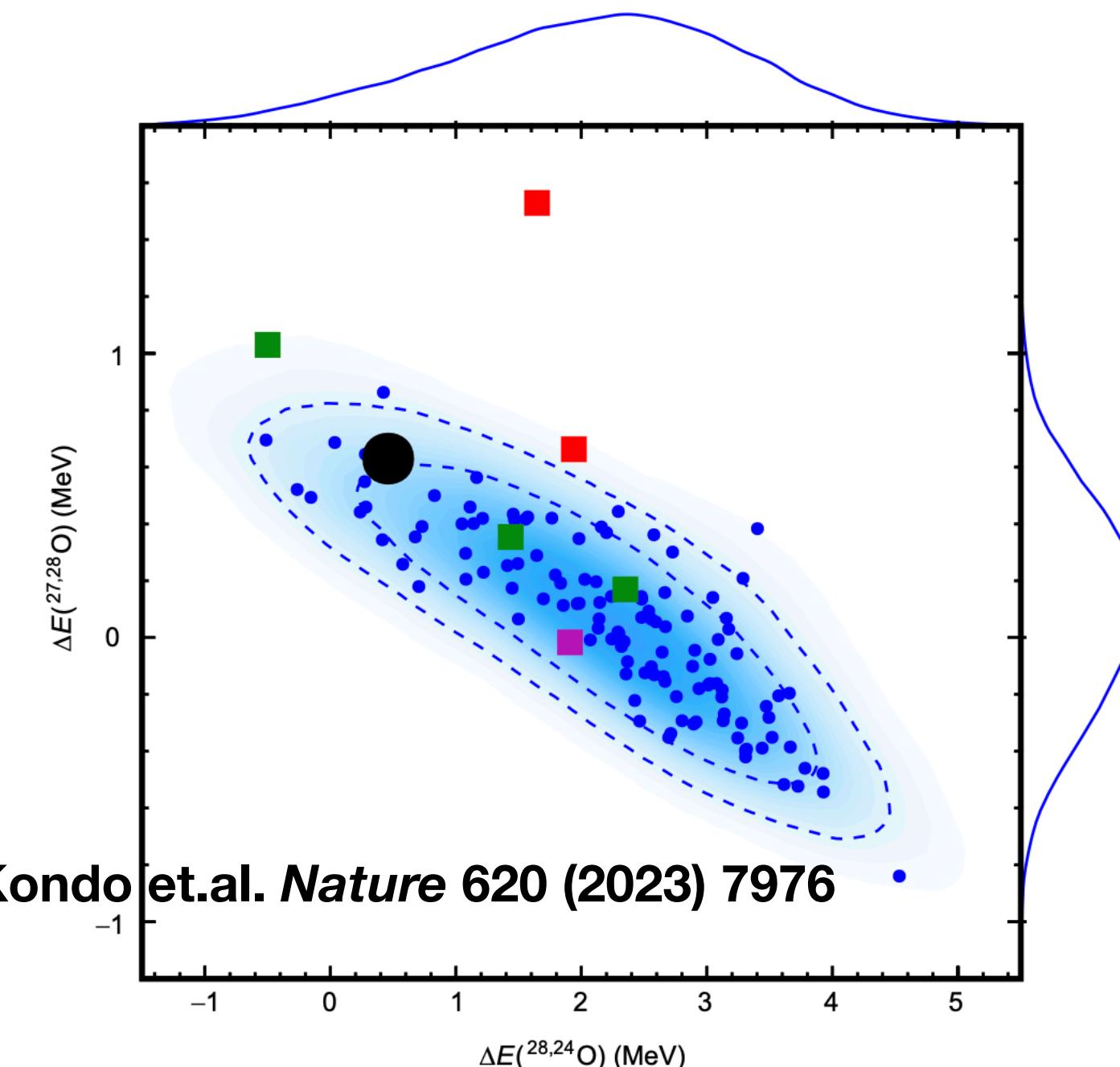


$$P(\mathcal{O}_1, \mathcal{O}_2, \dots)$$

$$P(c_1, c_2, \dots | \mathcal{O}_1, \mathcal{O}_2, \dots)$$

Too complicated

Also could fail where NN..LO contribution is large



$$\text{Domain of Parameters} \quad D(\{c_i\}_{i=1}^{n_{\text{par}}}) \longrightarrow \text{Calibration Observables} \quad P(\{c_i\}_{i=1}^{n_{\text{par}}} | \{\mathcal{O}\}) \longrightarrow \text{Large-scale Simulations} \quad P(\mathcal{O}_\odot | \{c_i\}_{i=1}^{n_{\text{par}}})$$

$$P(\{\mathcal{O}\}) \longrightarrow \boxed{\text{Our method}} \longrightarrow P(\mathcal{O}_\odot | \{\mathcal{O}\})$$

Obs

Predictions



Lets look at the exact math models ...

$$\left(H_0^{[a]} + c_1 H_1^{[a]} + c_2 H_2^{[a]} + \dots + c_{n_{\text{par}}} H_{n_{\text{par}}}^{[a]} - E^{[a]} N^{[a]} \right) \mathbf{y}^{[a]} = 0$$

Bound
 $E^{[a]}$ are eigenvalues

$$\left(H_0^{[b]} + c_1 H_1^{[b]} + c_2 H_2^{[b]} + \dots + c_{n_{\text{par}}} H_{n_{\text{par}}}^{[b]} - E^{[b]} N^{[b]} \right) \mathbf{y}^{[b]} = 0$$

Continuum
 $E^{[b]}$ are parameters

$$\left(H_0^{[c]} + c_1 H_1^{[c]} + c_2 H_2^{[c]} + \dots + c_{n_{\text{par}}} H_{n_{\text{par}}}^{[c]} - E^{[c]} N^{[c]} \right) \mathbf{y}^{[c]} = 0$$

?

$E^{[b]}$ are parameters
 $E^{[c]}, c_i$ are eigenvalues

•
•
•

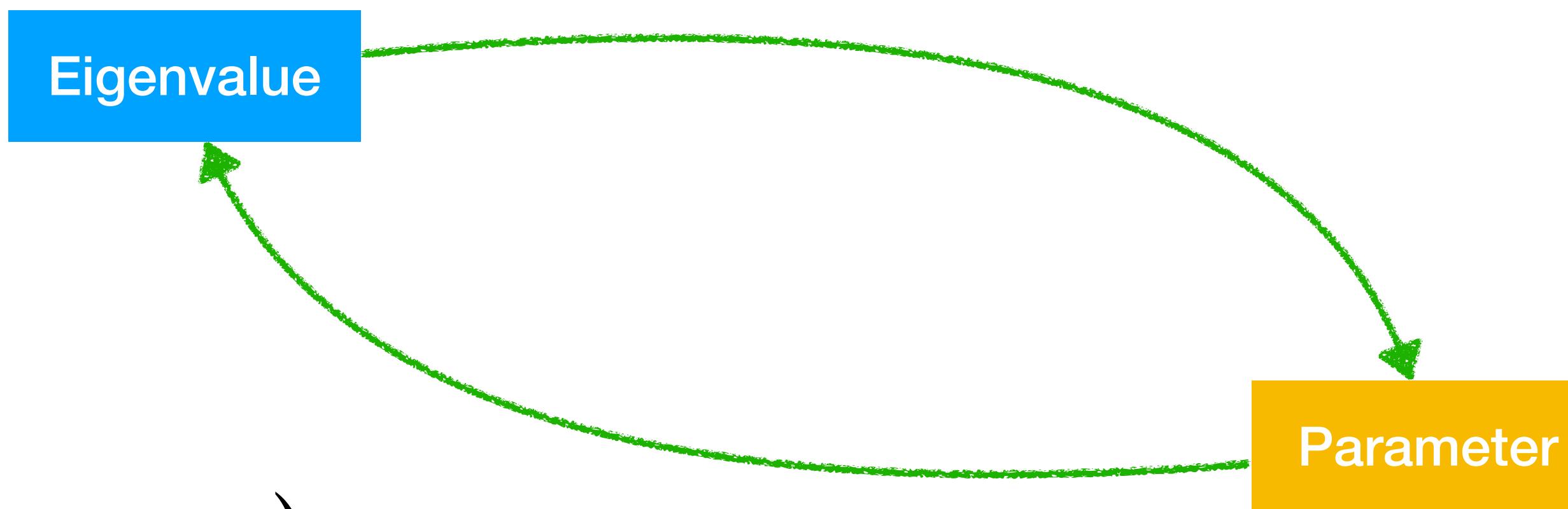
N equations to cancel N non-measurable!

Our Motivation

Our Motivation

N = 1 case

$$\left(H_0 + \sum_{i=1}^{n_{\text{par}}} c_i H_i - EN \right) \mathbf{y} = 0$$



$$H_{n_{\text{par}}}^{-1} \left(H_0 + \sum_{i=1}^{n_{\text{par}}-1} c_i H_i - EN \right) \mathbf{y} + c_{n_{\text{par}}} \mathbf{y} = 0$$

General N>1?

$$\left(H'_0 + \sum_{i=1}^{n_{\text{par}}} c_i H'_i - E^{\text{new}} N' \right) \mathbf{y}' = 0$$

Multiparameter Eigenvalue Problem

$$\left(H_0^{[j]} + \sum_{i=1}^{n_{\text{par}}} c_i H_i^{[j]} - E^{[j]} N^{[j]} \right) \mathbf{y}^{[j]} = 0 \quad 1 \leq j \leq n_{\text{par}}$$



$$\left(H_0^{[j]} + \sum_{i=1}^{n_{\text{par}}} c_i H_i^{[j]} - E^{[j]} N^{[j]} \right) \mathbf{y}^{[j]} = 0 \quad j > n_{\text{par}}$$

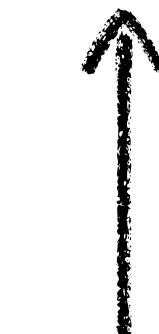
Method (Model Driven)

The *Multiparameter* in MEP
means

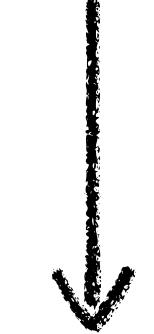
$$\left(H_0^{[j]} + \sum_{i=1}^{n_{\text{par}}} c_i H_i^{[j]} - E^{[j]} N^{[j]} \right) \mathbf{y}^{[j]} = 0$$

Multiparametric Eigenvalues

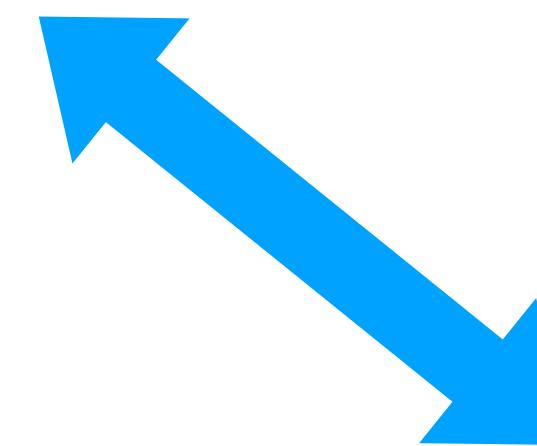
$$\left(O_j + \sum_{i=1}^m \alpha_i A_{ij} \right) \mathbf{y}_j = 0, \quad 1 \leq j \leq m$$



Inputs



Outputs



F. V. Atkinson and A. B. Mingarelli
Multiparameter eigenvalue problems:
Sturm-Liouville theory
(2010)

$$(K_i - \alpha_i K_0) \mathbf{y}_{\otimes} = 0, \quad 1 \leq i \leq m$$

$$\mathbf{y}_{\otimes} = \bigotimes_{j=1}^m \mathbf{y}_j$$

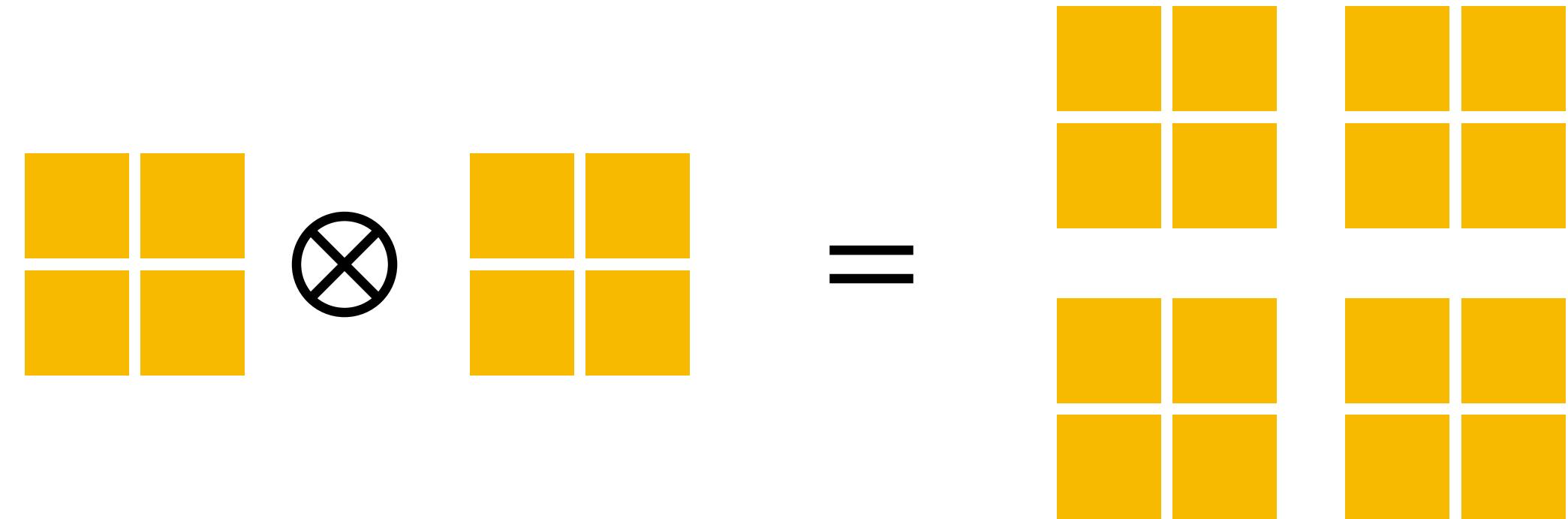
Kronecker Determinant

$$K_0 = \begin{vmatrix} A_{11} & \cdots & A_{1m} \\ \vdots & \ddots & \vdots \\ A_{m1} & \cdots & A_{mm} \end{vmatrix}_{\otimes}$$

$$\begin{vmatrix} G_{11} & \cdots & G_{1m} \\ \vdots & \ddots & \vdots \\ G_{m1} & \cdots & G_{mm} \end{vmatrix}_{\otimes} \equiv \sum_{\sigma \in S_m} \text{sgn}(\sigma) G_{1\sigma(1)} \otimes G_{2\sigma(2)} \cdots \otimes G_{m\sigma(m)}$$

Kronecker Product

$$K_i = \begin{vmatrix} A_{11} & \cdots & A_{1(i-1)} & O_1 & A_{1(i+1)} & \cdots & A_{1m} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ A_{m1} & \cdots & A_{m(i-1)} & O_m & A_{m(i+1)} & \cdots & A_{mm} \end{vmatrix}_{\otimes}$$



$M \times M$

$M^m \times M^m$

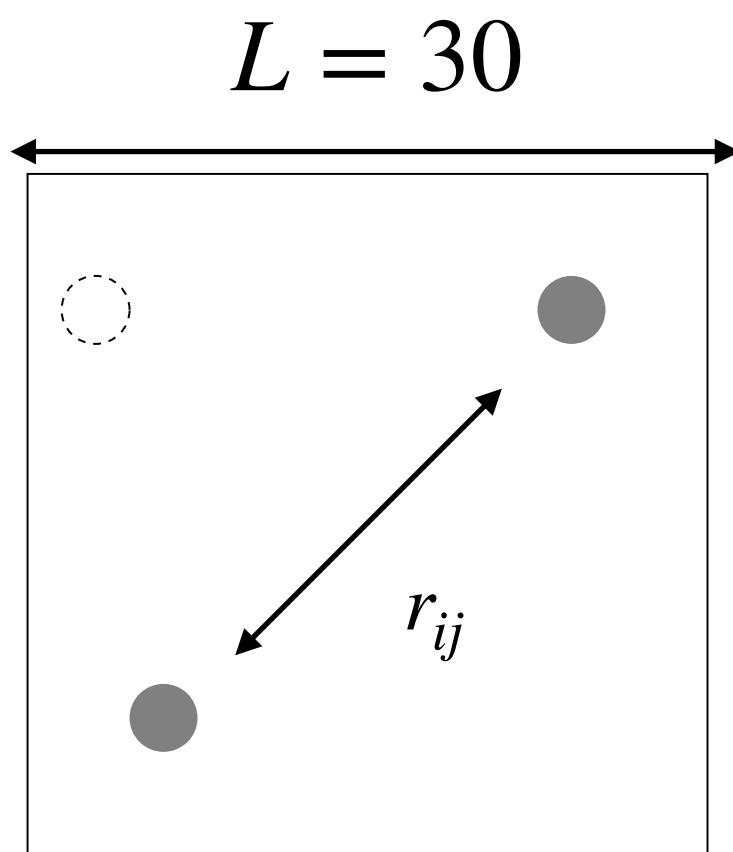
*can be different sizes

A toy model

$$\left(H_0^{[3]} + V_0 \exp(-r_{12}^2/R^2) + V_0 \exp(-r_{13}^2/R^2) + V_0 \exp(-r_{23}^2/R^2) - E_3 N^{[3]} \right) \mathbf{y}^{[3]} = 0$$

$$m = \hbar = c = 1$$

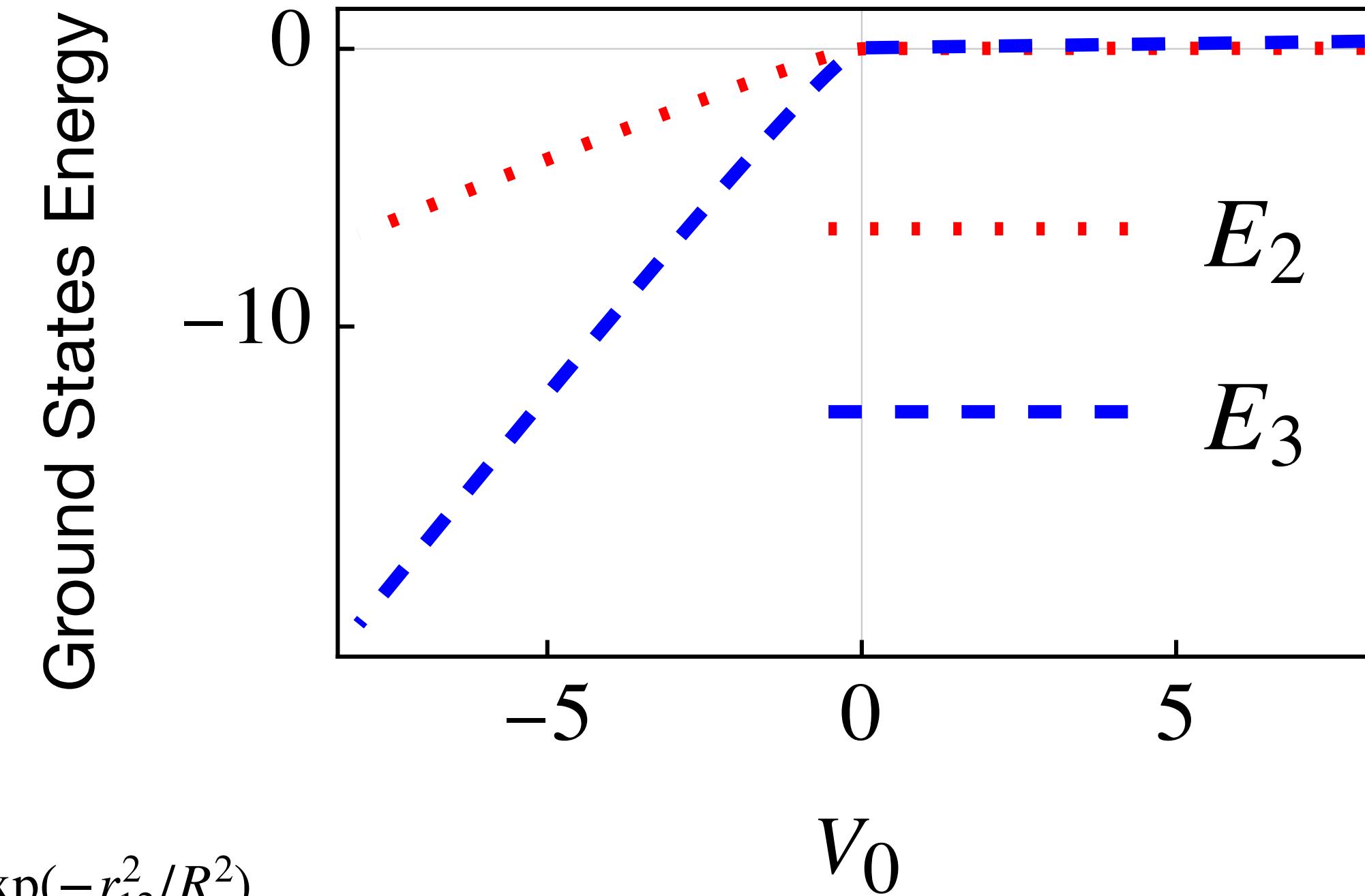
$$H_1^{[3]} \leftarrow \exp(-r_{12}^2/R^2) + \exp(-r_{13}^2/R^2) + \exp(-r_{23}^2/R^2)$$



$$V_{ij} = V_0 \exp(-r_{ij}^2/R^2) \quad R = 2.0$$

$$H_1^{[2]} \leftarrow \exp(-r_{12}^2/R^2)$$

$$\left(H_0^{[2]} + V_0 \exp(-r_{12}^2/R^2) - E_2 N^{[2]} \right) \mathbf{y}^{[2]} = 0$$

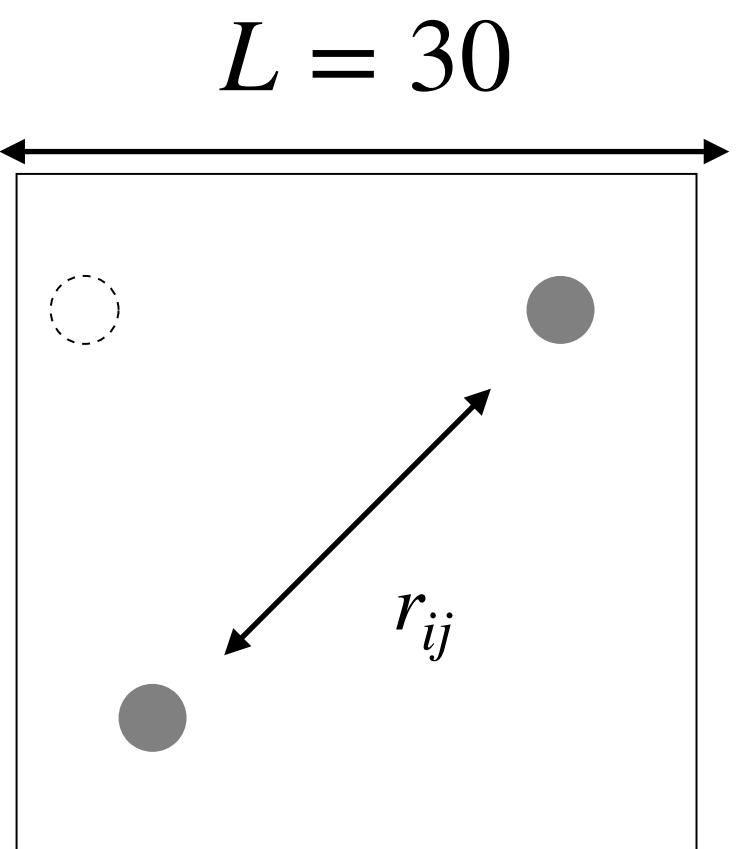


A toy model

$$\left(H_0^{[3]} + V_0 \exp(-r_{12}^2/R^2) + V_0 \exp(-r_{13}^2/R^2) + V_0 \exp(-r_{23}^2/R^2) - E_3 N^{[3]} \right) \mathbf{y}^{[3]} = 0$$

$$m = \hbar = c = 1$$

$$H_1^{[3]} \leftarrow \exp(-r_{12}^2/R^2) + \exp(-r_{13}^2/R^2) + \exp(-r_{23}^2/R^2)$$



$$O_1 = H_0^{[2]} + V_0 H_1^{[2]}$$

$$O_2 = H_0^{[3]} + V_0 H_1^{[3]}$$

$$A_{11} = -N^{[2]} \quad A_{21} = 0$$

$$\left(O_j + \sum_{i=1}^2 \alpha_i A_{ij} \right) \mathbf{y}_j = 0, \quad 1 \leq j \leq 2$$

$$V_{ij} = V_0 \exp(-r_{ij}^2/R^2) \quad R = 2.0$$

$$H_1^{[2]} \leftarrow \exp(-r_{12}^2/R^2)$$

$$A_{12} = 0 \quad A_{22} = -N^{[3]}$$

$$\left(H_0^{[2]} + V_0 \exp(-r_{12}^2/R^2) - E_2 N^{[2]} \right) \mathbf{y}^{[2]} = 0$$

Diagonal MEP

$$\alpha_1 = E_2$$

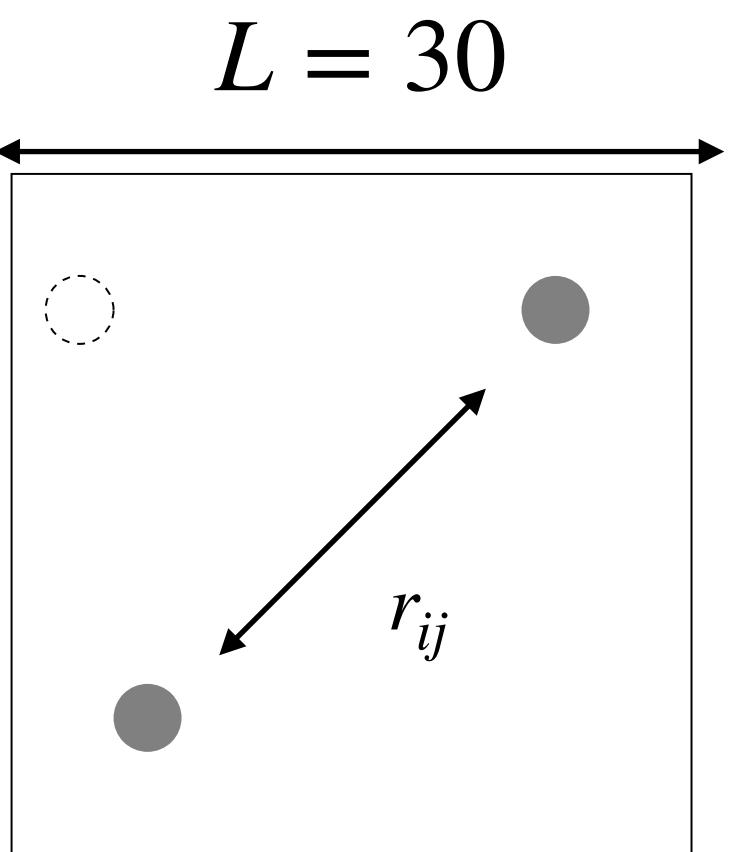
$$\alpha_2 = E_3$$

A toy model

$$\left(H_0^{[3]} + V_0 \exp(-r_{12}^2/R^2) + V_0 \exp(-r_{13}^2/R^2) + V_0 \exp(-r_{23}^2/R^2) - E_3 N^{[3]} \right) \mathbf{y}^{[3]} = 0$$

$$m = \hbar = c = 1$$

$$H_1^{[3]} \leftarrow \exp(-r_{12}^2/R^2) + \exp(-r_{13}^2/R^2) + \exp(-r_{23}^2/R^2)$$



$$O_1 = H_0^{[2]} - E_2 N^{[2]}$$

$$O_2 = H_0^{[3]}$$

$$A_{11} = H_1^{[2]} \quad A_{21} = 0$$

$$\left(O_j + \sum_{i=1}^2 \alpha_i A_{ij} \right) \mathbf{y}_j = 0, \quad 1 \leq j \leq 2$$

$$A_{12} = H_1^{[3]} \quad A_{22} = -N^{[3]}$$

$$V_{ij} = V_0 \exp(-r_{ij}^2/R^2) \quad R = 2.0$$

$$H_1^{[2]} \leftarrow \exp(-r_{12}^2/R^2)$$

$$\left(H_0^{[2]} + V_0 \exp(-r_{12}^2/R^2) - E_2 N^{[2]} \right) \mathbf{y}^{[2]} = 0$$

$$\alpha_1 = V_0 \quad \alpha_2 = E_3$$

MEP

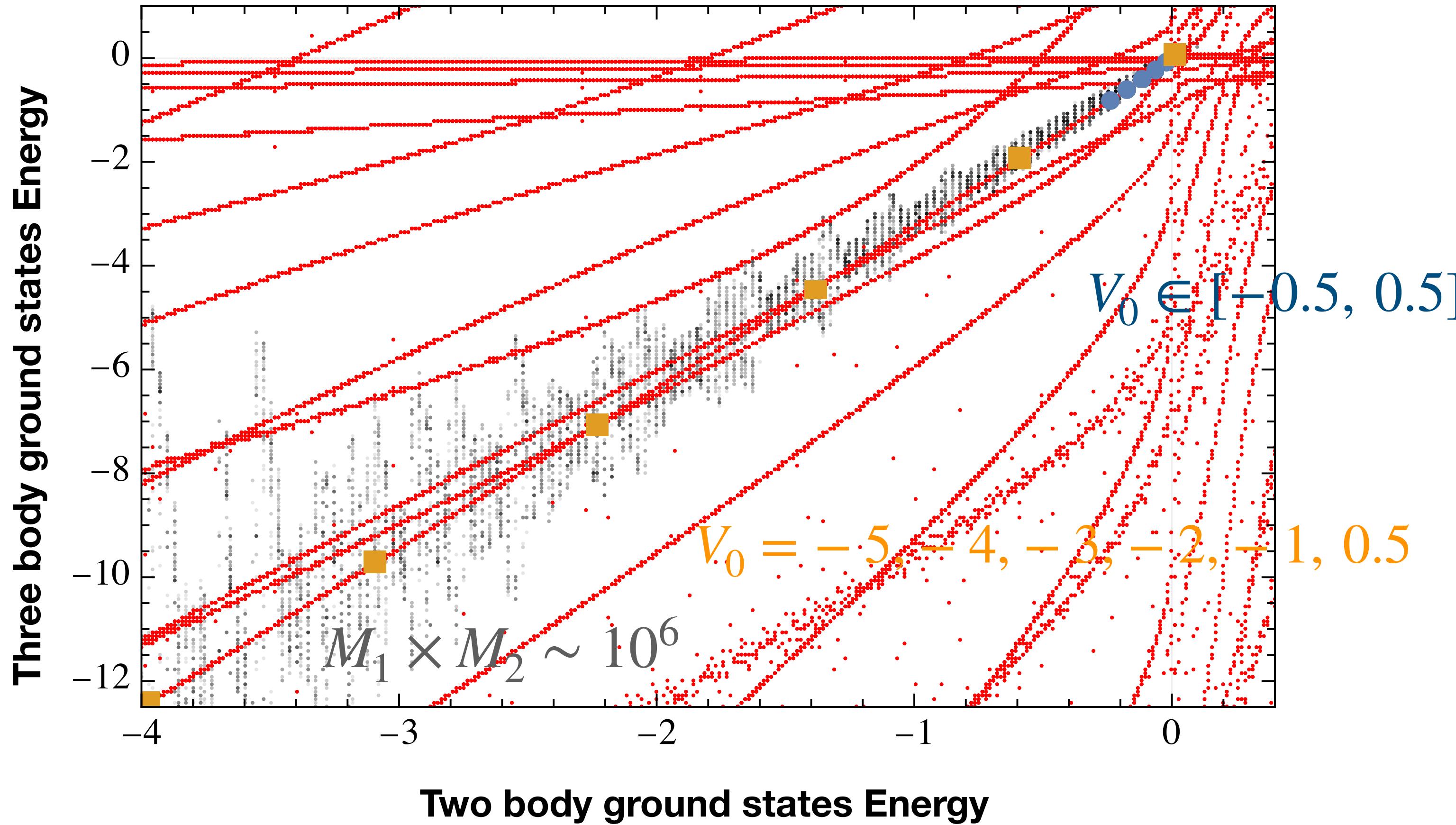
$$K_2 = \begin{vmatrix} H_1^{[2]} & H_0^{[2]} - E_2 N^{[2]} \\ H_1^{[3]} & H_0^{[3]} \end{vmatrix} \otimes$$

$$K_0 = \begin{vmatrix} H_1^{[2]} & 0 \\ H_1^{[3]} & -N^{[3]} \end{vmatrix} \otimes$$

$$M^m = 6^2$$

$$[K_2(E_2) + E_3 K_0]y = 0$$

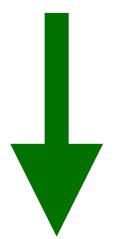
Contamination from G.S. \otimes Fake (Excited)



Kronecker Determinant

$$K_0 = \begin{vmatrix} A_{11} & \cdots & A_{1m} \\ \vdots & \ddots & \vdots \\ A_{m1} & \cdots & A_{mm} \end{vmatrix}_{\otimes}$$

$$(K_i - \alpha_i K_0) \mathbf{y}_\otimes = 0 \quad M^m \times M^m$$



$$(\mathcal{K}_0)_{kl} = (\mathbf{y}_{\otimes,k})^* \cdot K_0 \cdot \mathbf{y}_{\otimes,l} = \begin{vmatrix} (A_{11})_{kl} & \cdots & (A_{1m})_{kl} \\ \vdots & \ddots & \vdots \\ (A_{m1})_{kl} & \cdots & (A_{mm})_{kl} \end{vmatrix}$$

$$(\mathcal{K}_i - \alpha_i \mathcal{K}_0) \mathbf{y} = 0 \quad M \times M$$

Matrix Determinant!

MEP

EC

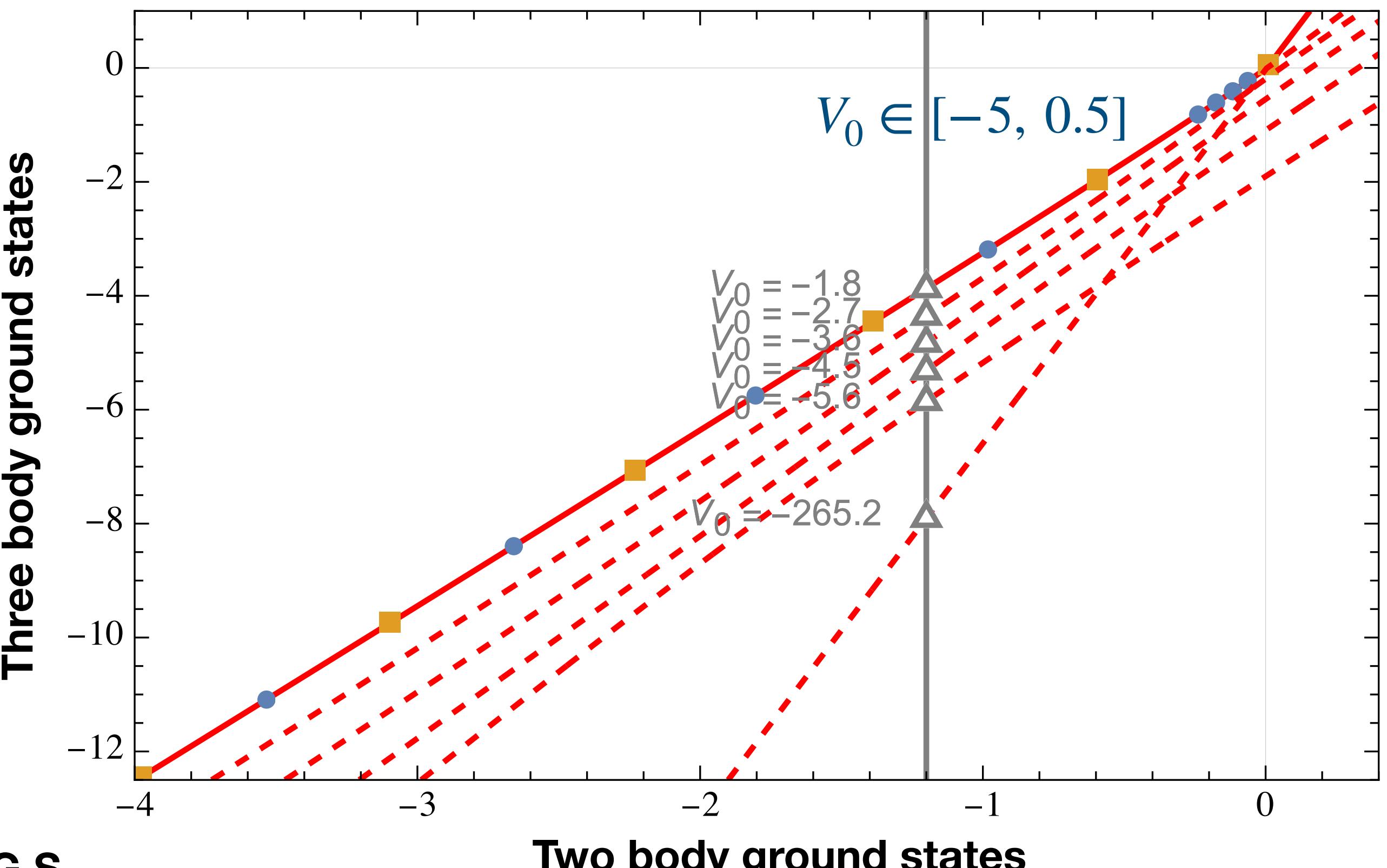
$$(\mathcal{K}_2)_{kl} = - (H_0^{[2]} - E_2 N^{[2]})_{kl} (H_1^{[3]})_{kl} + (H_1^{[2]})_{kl} (H_0^{[3]})_{kl}$$

$$(\mathcal{K}_0)_{kl} = (H_1^{[2]})_{kl} (-N^{[3]})_{kl}$$

$$[\mathcal{K}_2(E_2) + E_3 \mathcal{K}_0] \tilde{\mathbf{y}} = 0$$

Only G.S. \otimes G.S.

$$V_0 = -5, -4, -3, -2, -1, 0.5$$



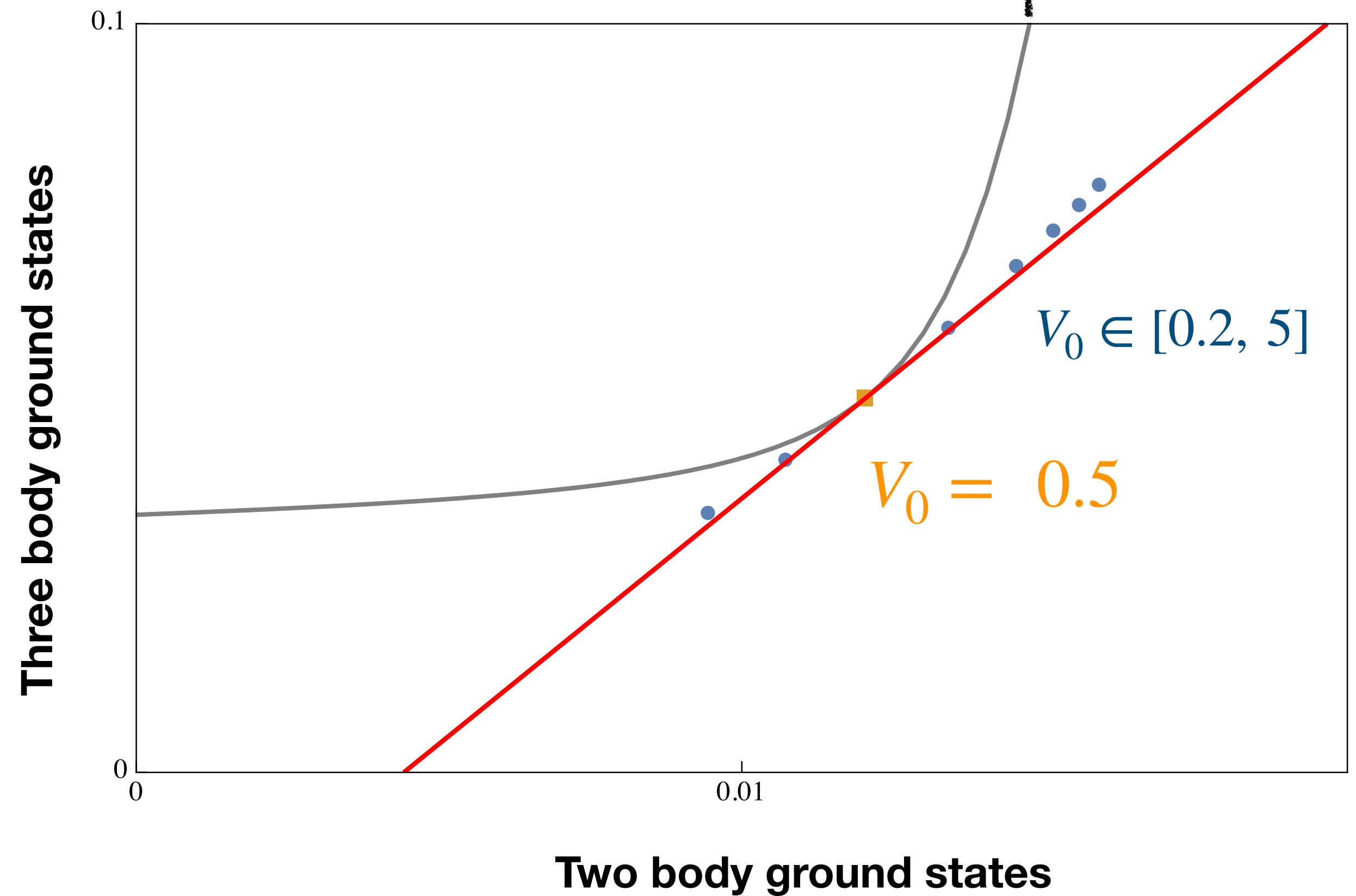
Zooming In

$$(\mathcal{K}_2)_{kl} = - (H_0^{[2]} - E_2 N^{[2]})_{kl} (H_1^{[3]})_{kl} + (H_1^{[2]})_{kl} (H_0^{[3]})_{kl}$$

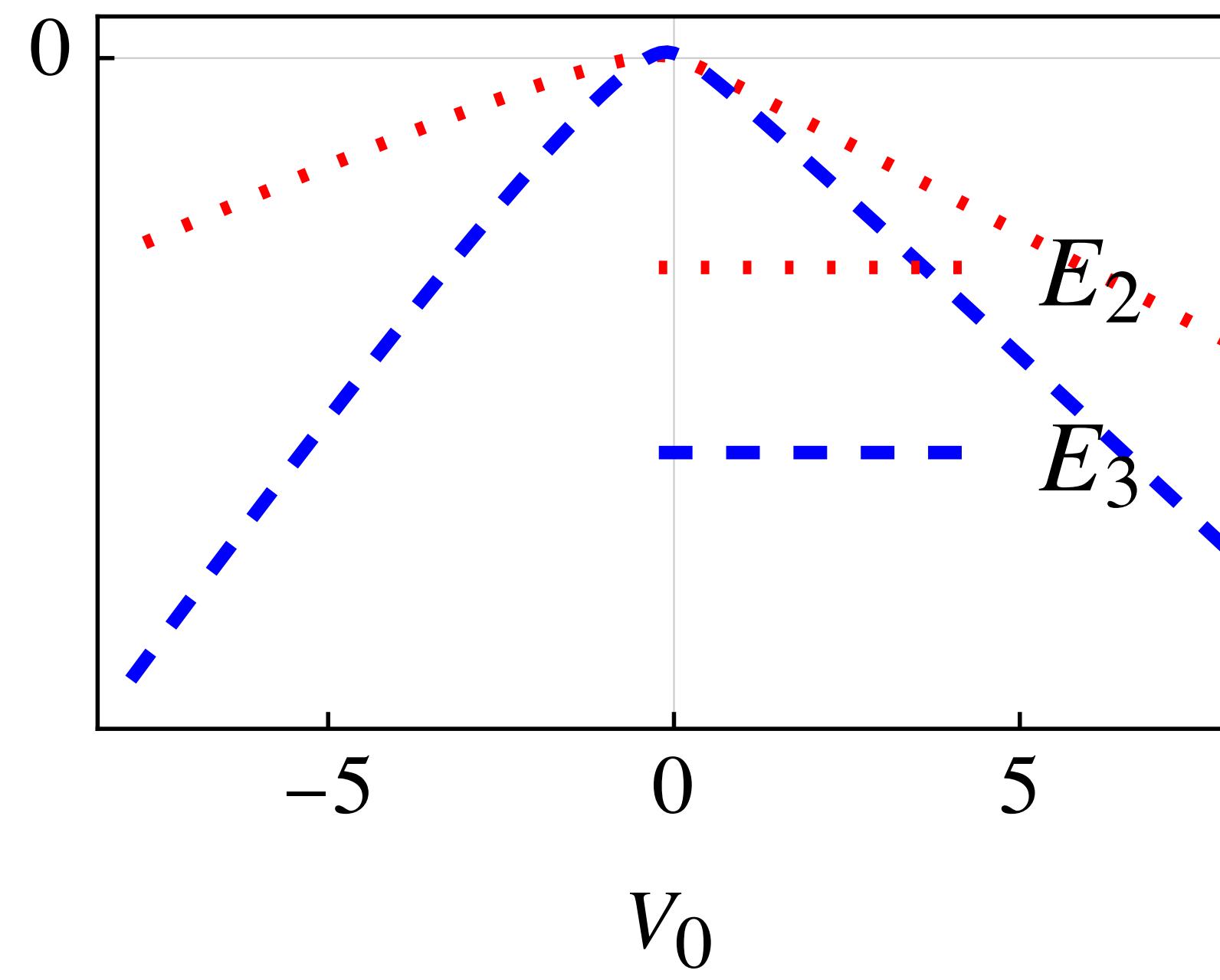
$$(\mathcal{K}_0)_{kl} = (H_1^{[2]})_{kl} (-N^{[3]})_{kl}$$

$$[\mathcal{K}_2(E_2) + E_3 \mathcal{K}_0] \tilde{\mathbf{y}} = 0$$

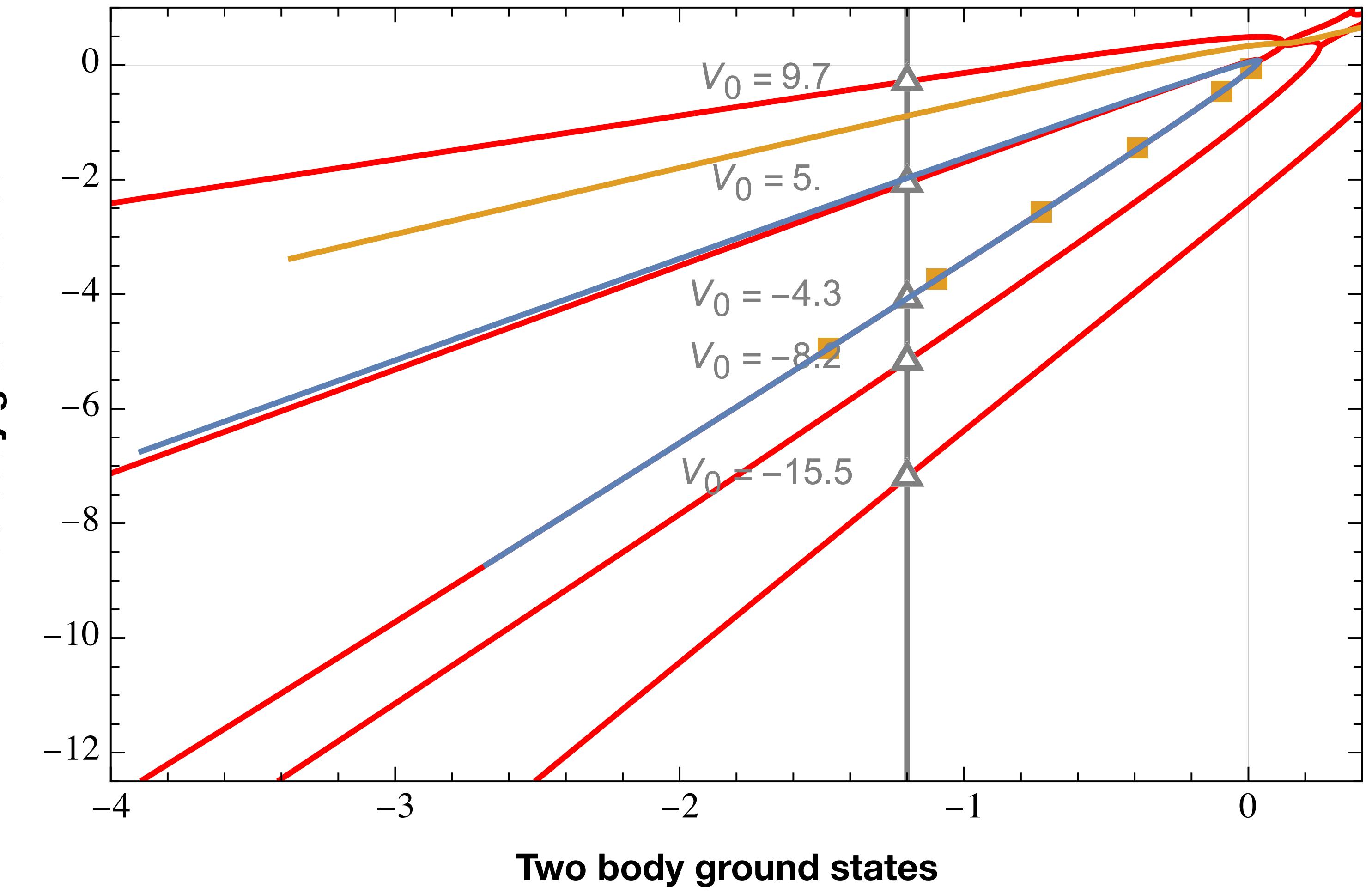
$$H_{n_{\text{par}}}^{-1} \left(H_0 + \sum_{i=1}^{n_{\text{par}}-1} c_i H_i - EN \right) \mathbf{y} + c_{n_{\text{par}}} \mathbf{y} = 0$$



MEP can be Multi(Nil)-Valued



Three body ground states



Data-Driven

Parametric Matrix Model

It could be difficult to access wavefunctions

**IMSRG
NLEFT**

...

EC

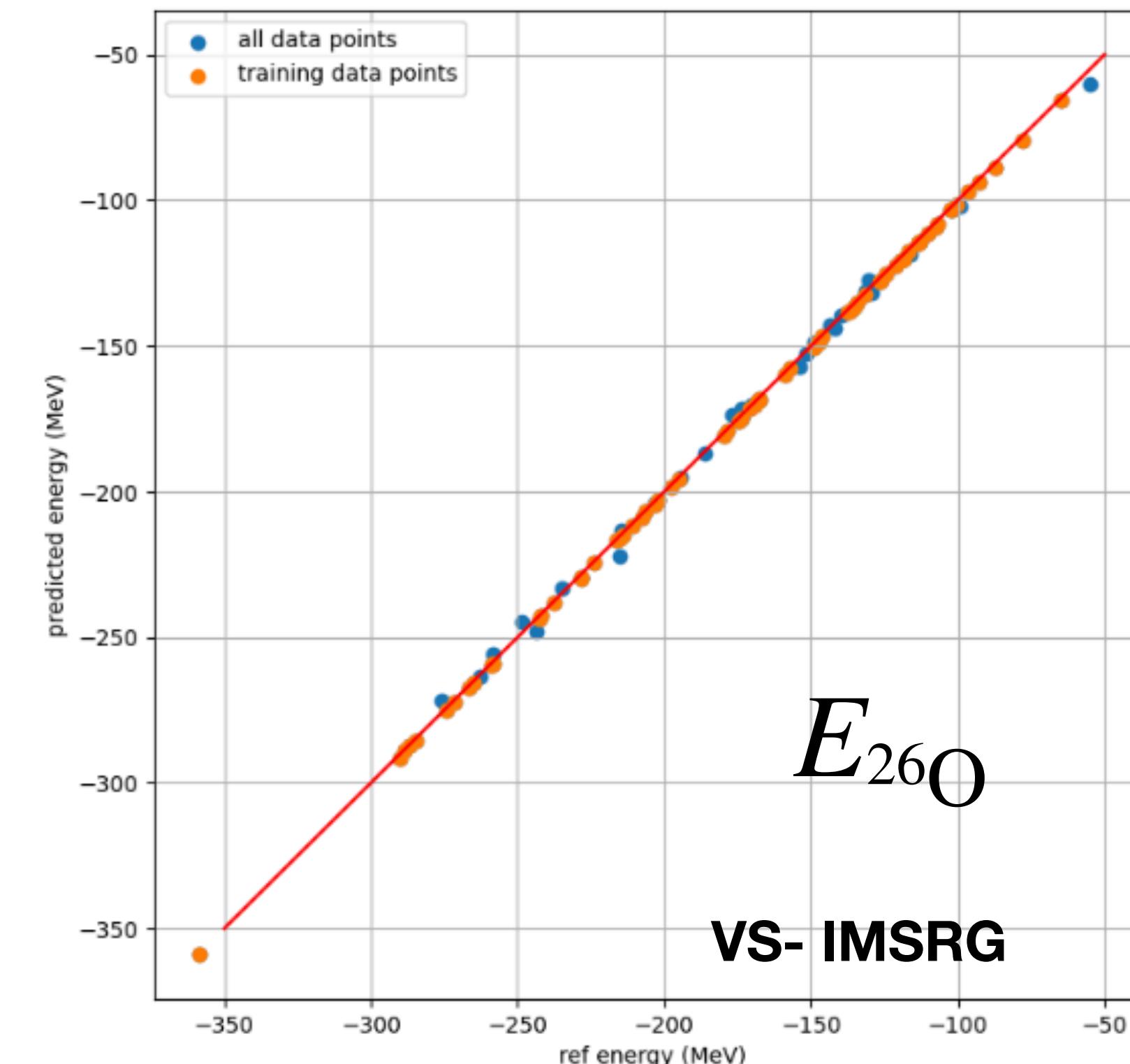
$$\left(H_0 + c_1 H_1 + c_2 H_2 + \dots + c_{n_{\text{par}}} H_{n_{\text{par}}} - EN \right) \mathbf{y} = 0$$

$$H_i^{[a]} \sim [M \times M] \quad \sim 1000 \text{ parameters}$$

We just need to be lucky 1000 times!

Or Gradient Descent (Adam)

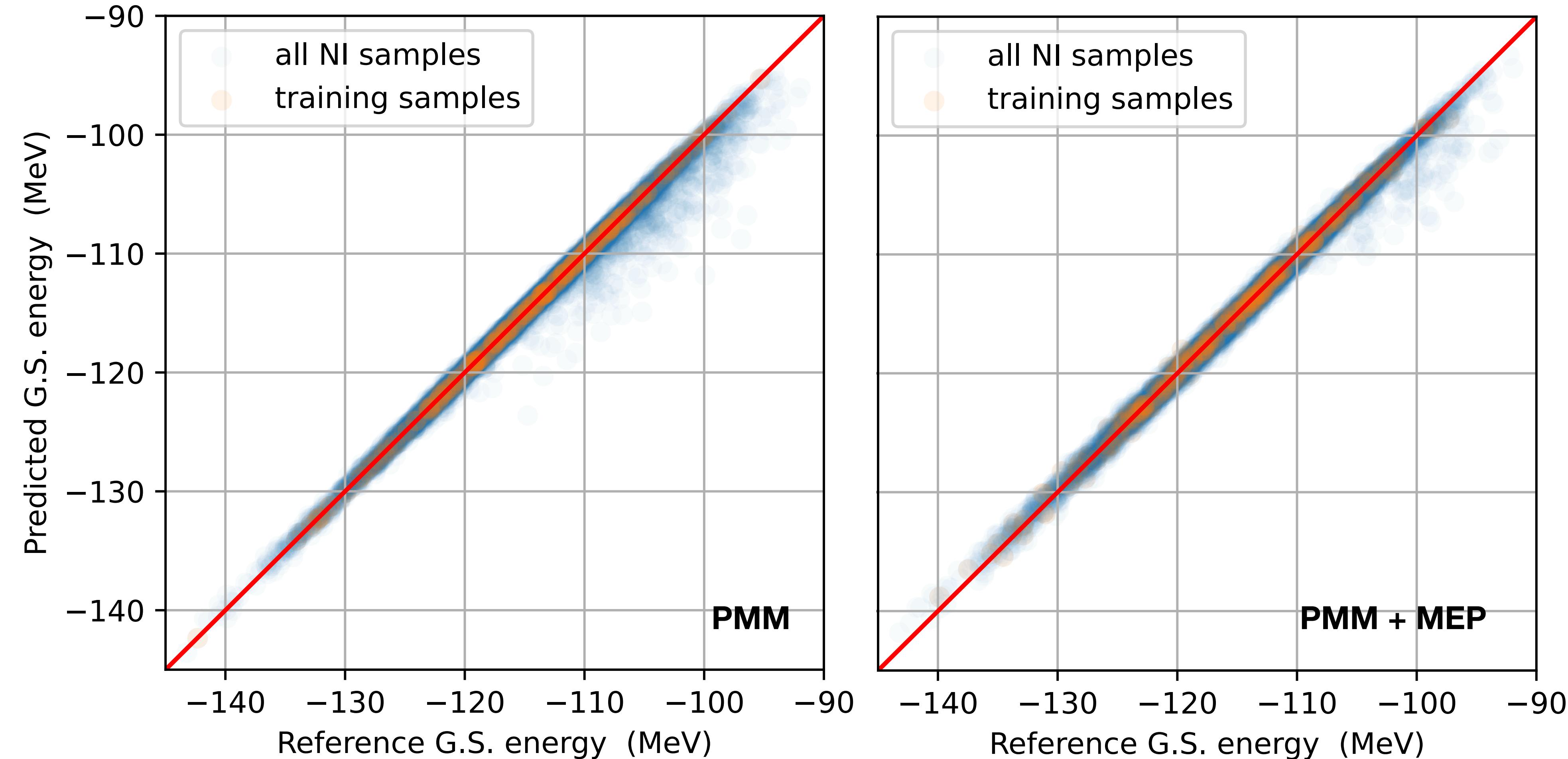
P. Cook et.al. arxiv: 2401.11694 [cs.LG]



$$(\mathcal{K}_i - \alpha_i \mathcal{K}_0) \mathbf{y} = (\mathcal{K}_{i0} + E^{[1]} \mathcal{K}_{i1} + E^{[2]} \mathcal{K}_{i2} + \dots + E^{[n]} \mathcal{K}_{in} - \alpha_i \mathcal{K}_0) \mathbf{y} = 0$$

MEP

PMM



8188 ~ 8218 samples of O16

data from W.G. Jiang et.al., PRC 109 064314 (2024)

Applications

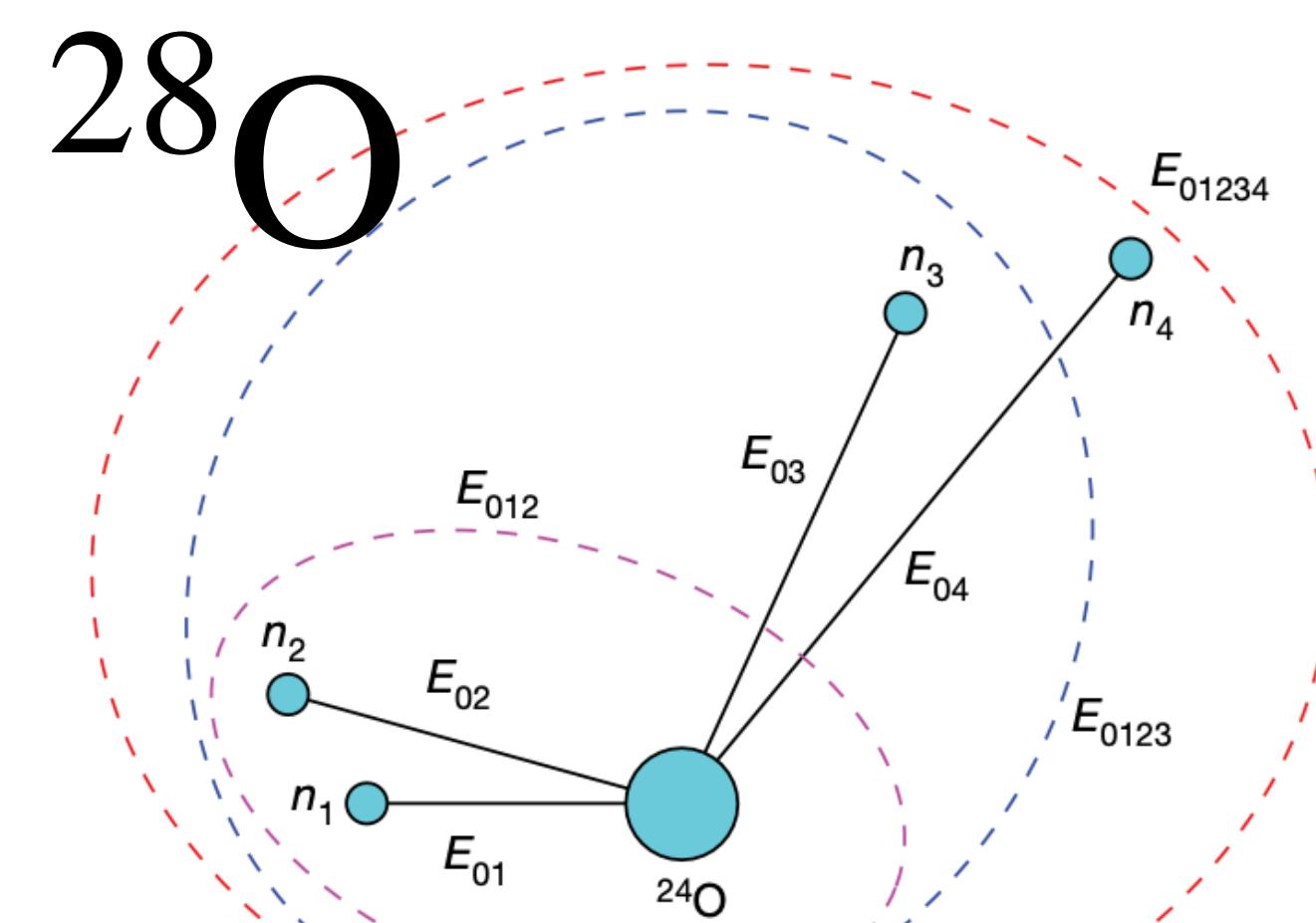
$n - p$ Phase Shifts at Lab Energies 5 MeV and 50 MeV:

$^1S_0, ^3S_1, ^1P_1, ^3P_0, ^3P_1, ^3P_2$

$A = 2 \sim 24$ Observables

G.S. Energies: $^2\text{H}, ^3\text{H}, ^4\text{He}, ^6\text{Li}, ^{16}\text{O}, ^{24}\text{O}$

MEP

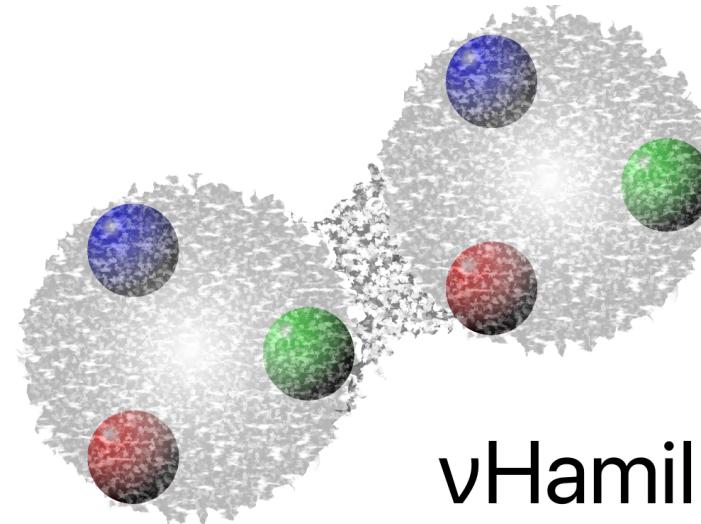


Y. Kondo et.al. Nature 620 (2023) 7976

W.G. Jiang et.al. PRC 109 (2024) 064314
Supplemental materials

Δ – full χ - EFT at NNLO

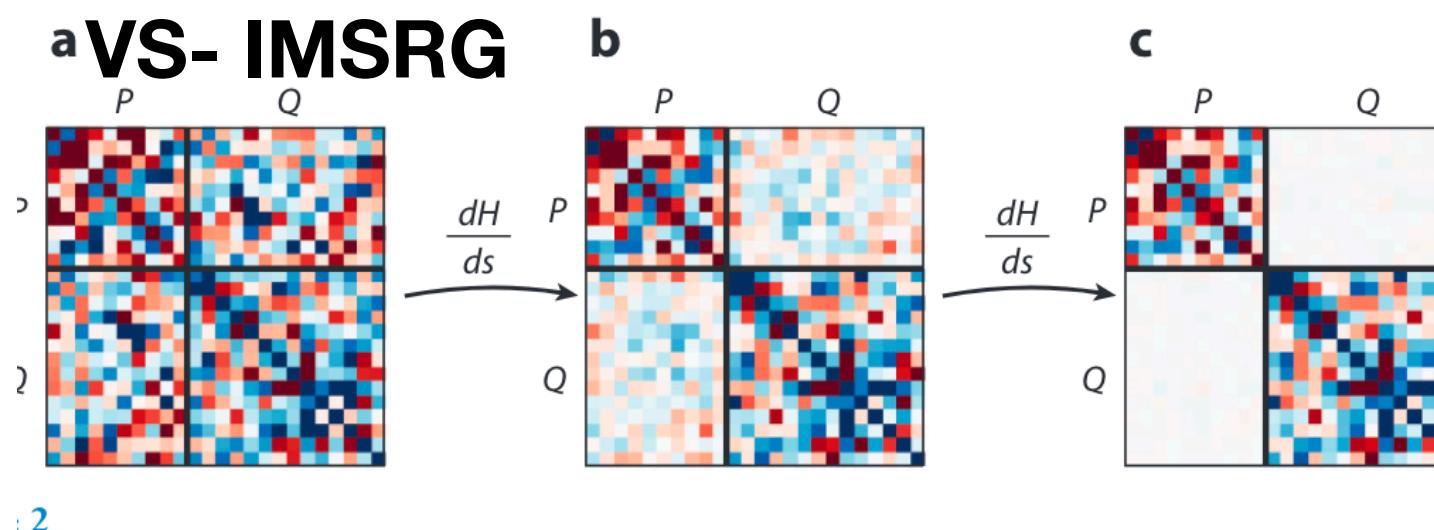
W. G. Jiang, et.al. PRC 102, 054301 (2020)



17 LECs

T. Miyagi, EPJA 59, 150 (2023)

efficiently store three body potentials



S. R. Stroberg et.al., Ann. Rev. Nucl. Part. Sci. 69, 307 (2019)

evolve effective interactions from χ -EFT

KSHELL

N. Shimizu et.al., Comp. Phys. Comm. 244, 372 (2019)

perform diagonalization

Training Samples

E_{24O} E_{27O} E_{28O}

at $\hbar\omega = 16$ MeV, $e_{\max} = 12$, $E_{3\max} = 18$

300 G.S. Energies

~ 2000 Node * Hours on Pegasus Supercomputer @ CCS Tsukuba

~ 700 Node * Hours on Miyabi-G with MIG

MEP

PMM

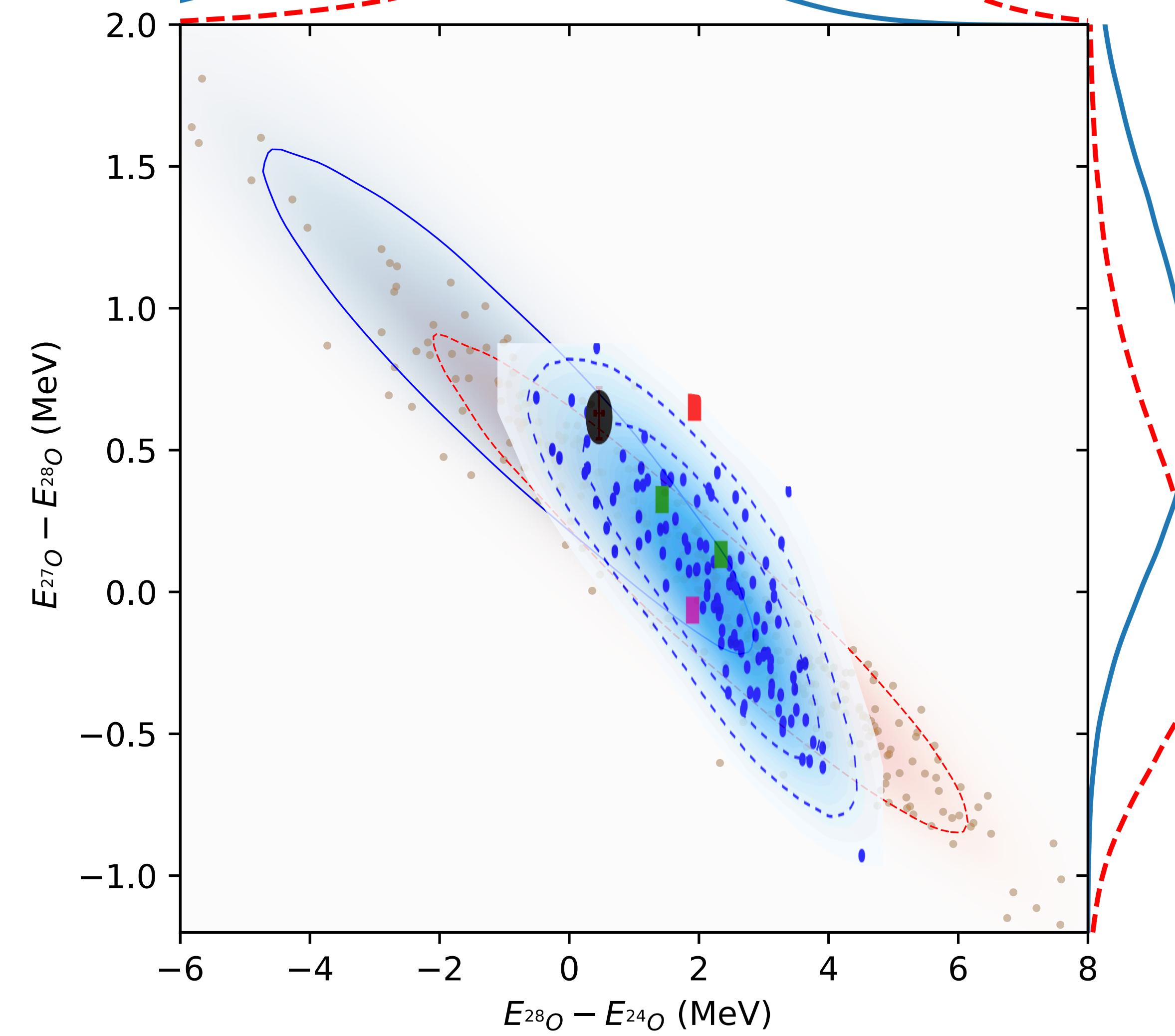
Application

$A = 2 \sim 24$ Observables

According to $A = 2 \sim 4$ NI measure

W.G. Jiang et.al. PRC 109 (2024) 064314

$P(O_1, O_2, \dots) \sim$ Multivariate Gaussian



Y. Kondo et.al. Nature 620 (2023) 797

$A = 2 \sim 24$ Observables

According to $A = 2 \sim 4$ NI measure

Plus constraints of G.S. $^{16}\text{O}, ^{24}\text{O}$

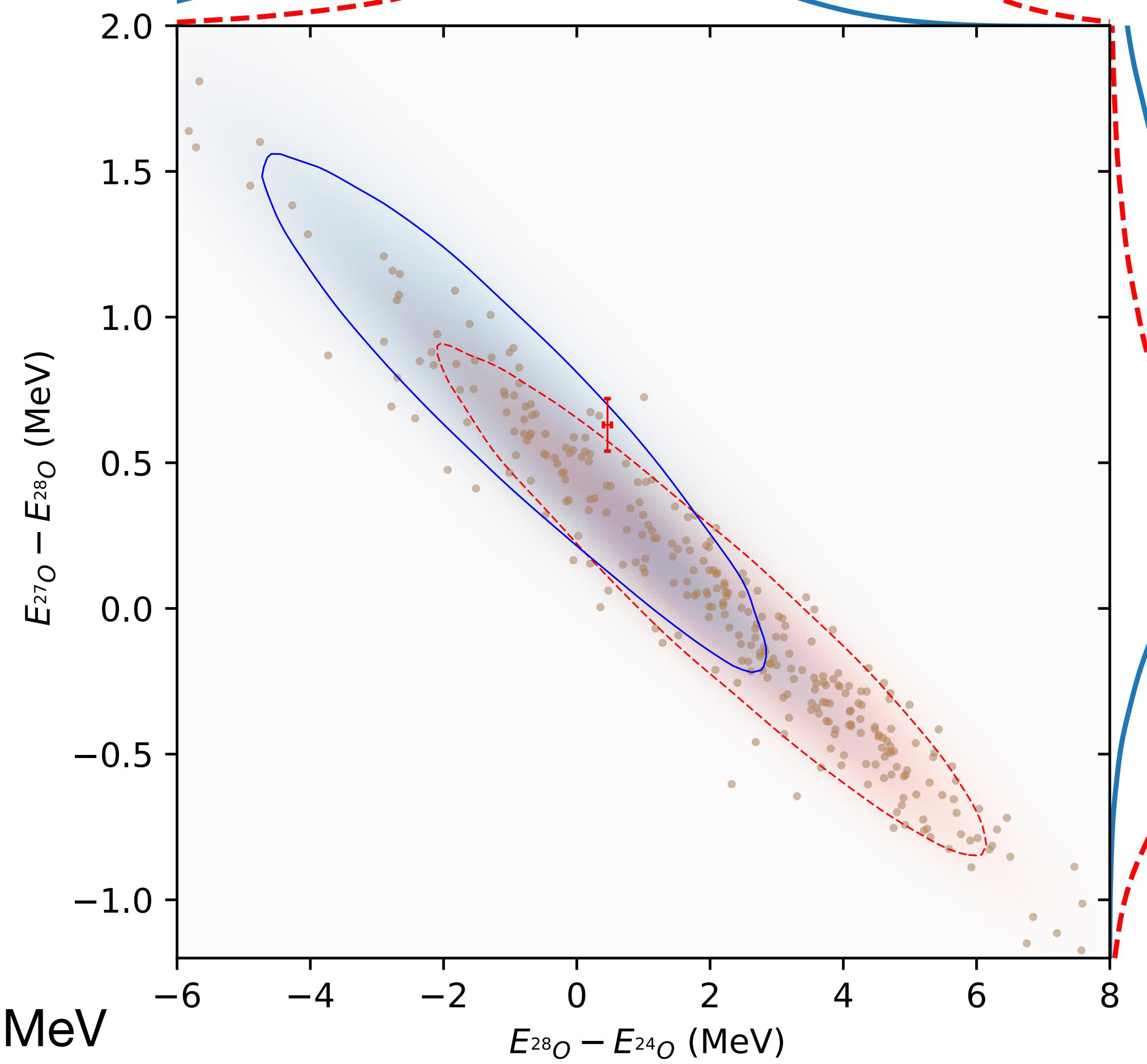
$$E_{^{16}\text{O}} = 127(2)\text{MeV}$$

$$E_{^{24}\text{O}} = 168(3)\text{MeV}$$

PMM error (100 bootstrapping $\times 10^5$ samples): 0.2 MeV

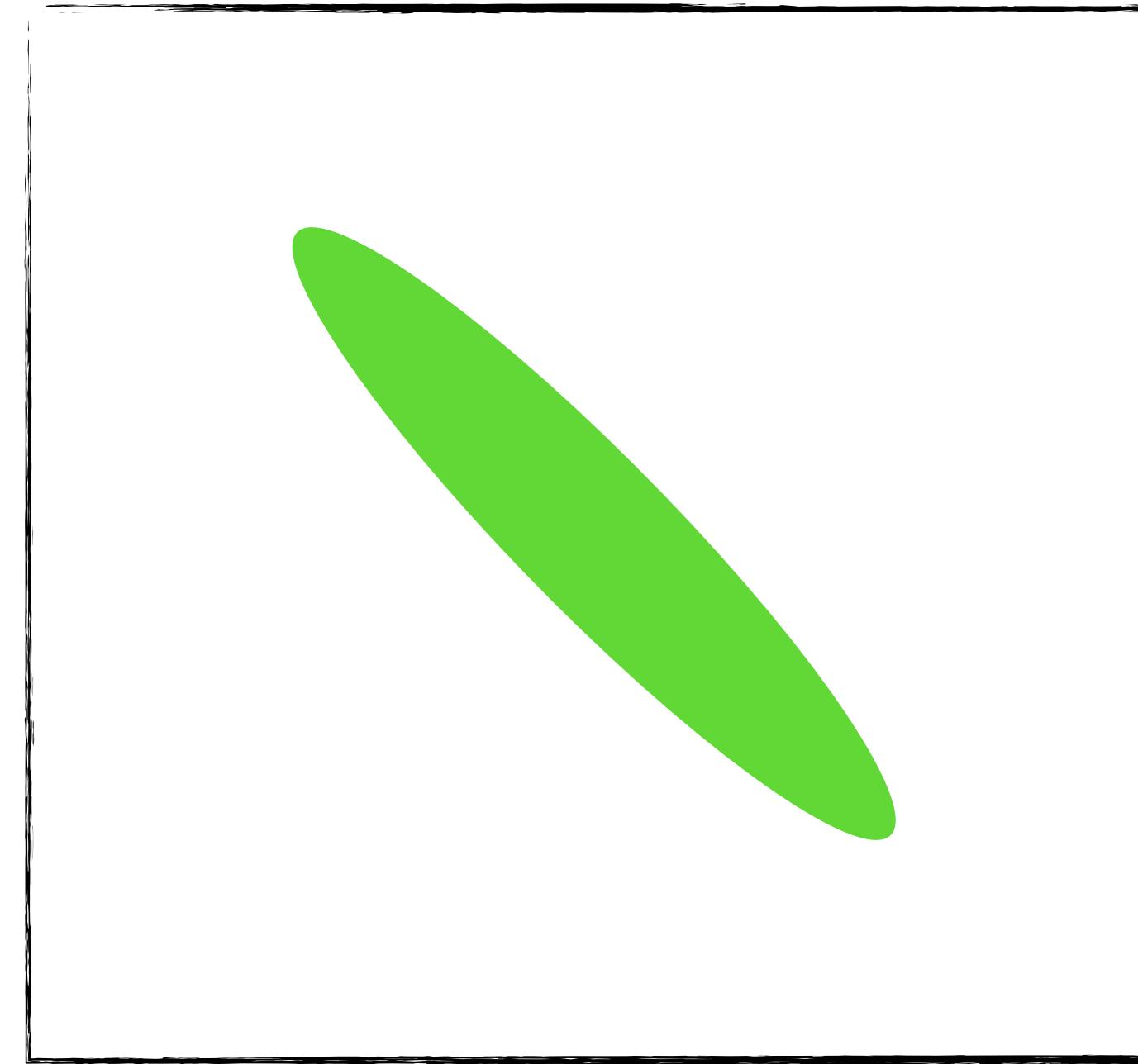
Continuum: 0.5 ~ 1 MeV

G. Hagen, et al. Phys. Scripta 91 063006 (2016)

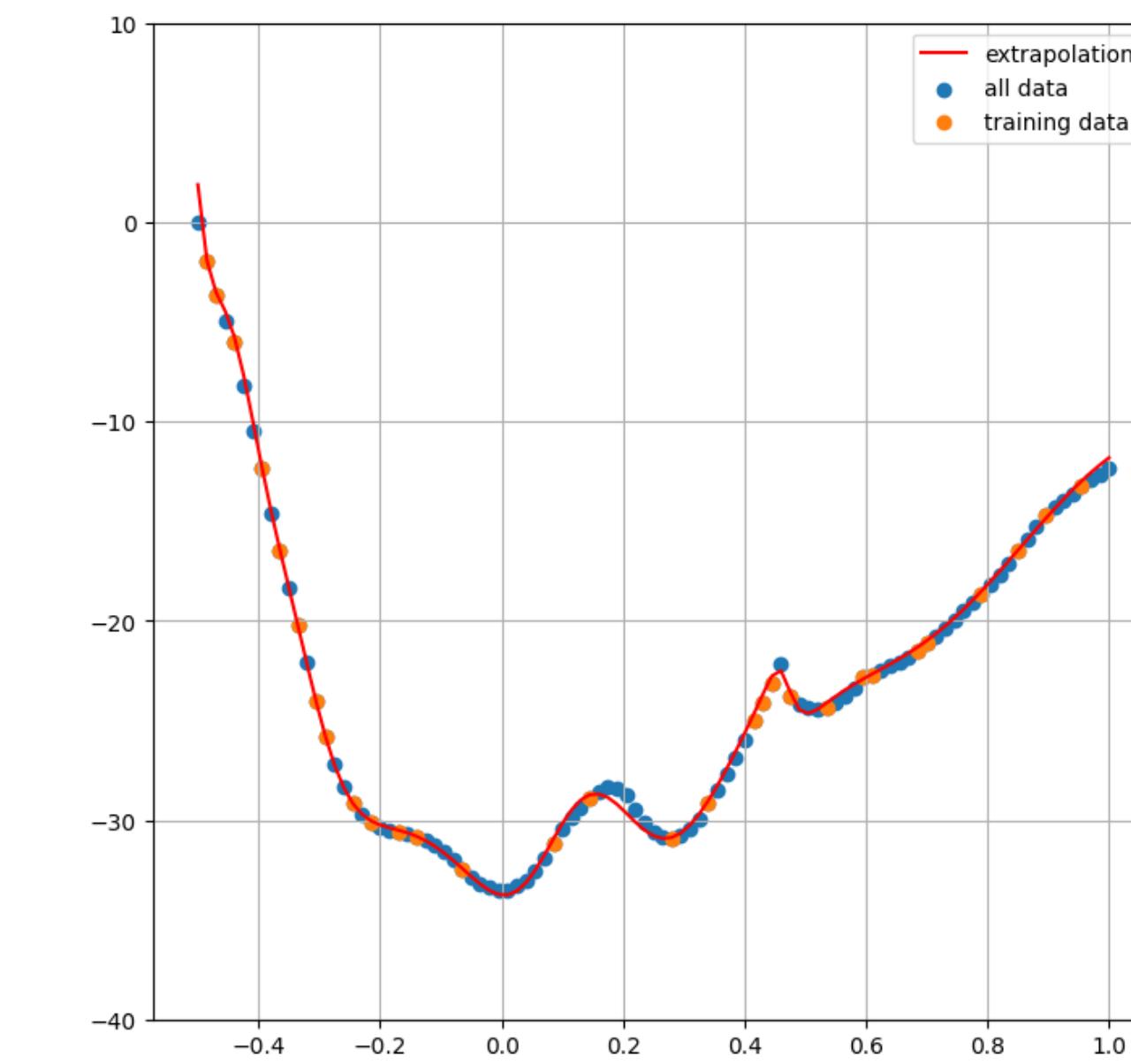


Outlook

Many body observables



Three body observables



Improving Emulator for Potential Energy Surface
Using Augmented Lagrangian Method
[Eur. Phys. J. A 46, 85–90 (2010)]

Summary

- We developed a connection emulator
- An efficient and deterministic substitution of statistical workflow
- Prob critical point in solution space
- A customizable hybrid surrogate model
- Currently: (phase shifts), eigenstates...

Next: General observables via PT/ALM

Thank You!