



中山大學天琴中心

TIANQIN CENTER FOR GRAVITATIONAL PHYSICS, SYSU



超轻轴子的引力波信号

黄发朋 (Fa Peng Huang)

中山大学物理与天文学院

based on recent work with my Ph.D students:

Ning Xie, **FPH**, arXiv:2503.10347;

FPH, arXiv: 2409.19906

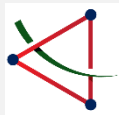
Aidi Yang, **FPH**, arXiv:2404.18703, JCAP (2025) ;

Jing Yang, Ning Xie, **FPH**, JCAP 11 (2024) 045;

Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024).

Axion Dark Matter: Theory and Phenomenology

@山东大学, 青岛, 2025.05.11

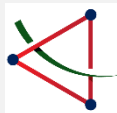


Outline

1. Axion and axion dark matter (DM), Gravitational wave (GW)
2. DFSZ axion and its phase transition GW signals
3. μeV axion and radio signals of **axion** DM (multi-messenger)
4. 10^{-12} - 10^{-17} eV axion: GW and pulsar timing measurement
5. 10^{-21} eV fuzzy axion DM GW
6. Summary and outlook



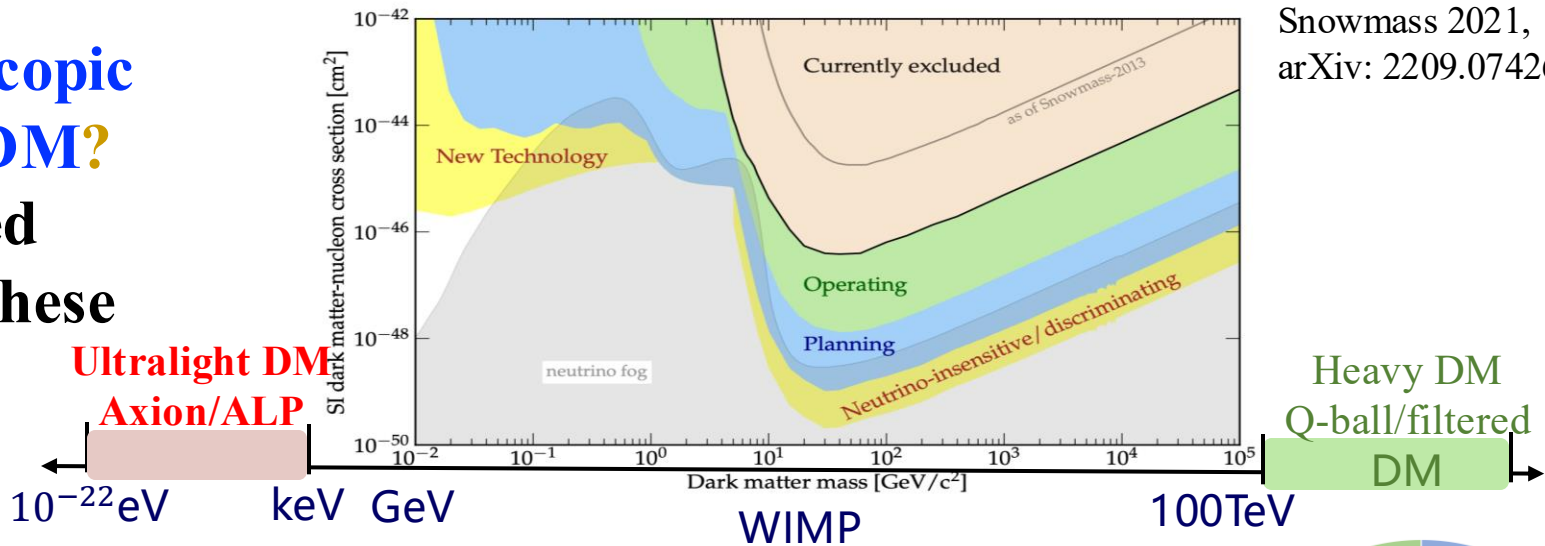
“Ultralight” DM



Motivation

What is
the microscopic
nature of DM?

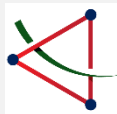
No expected
signals in these
region



arXiv: 1904.07915
Snowmass 2021,
arXiv: 2209.07426

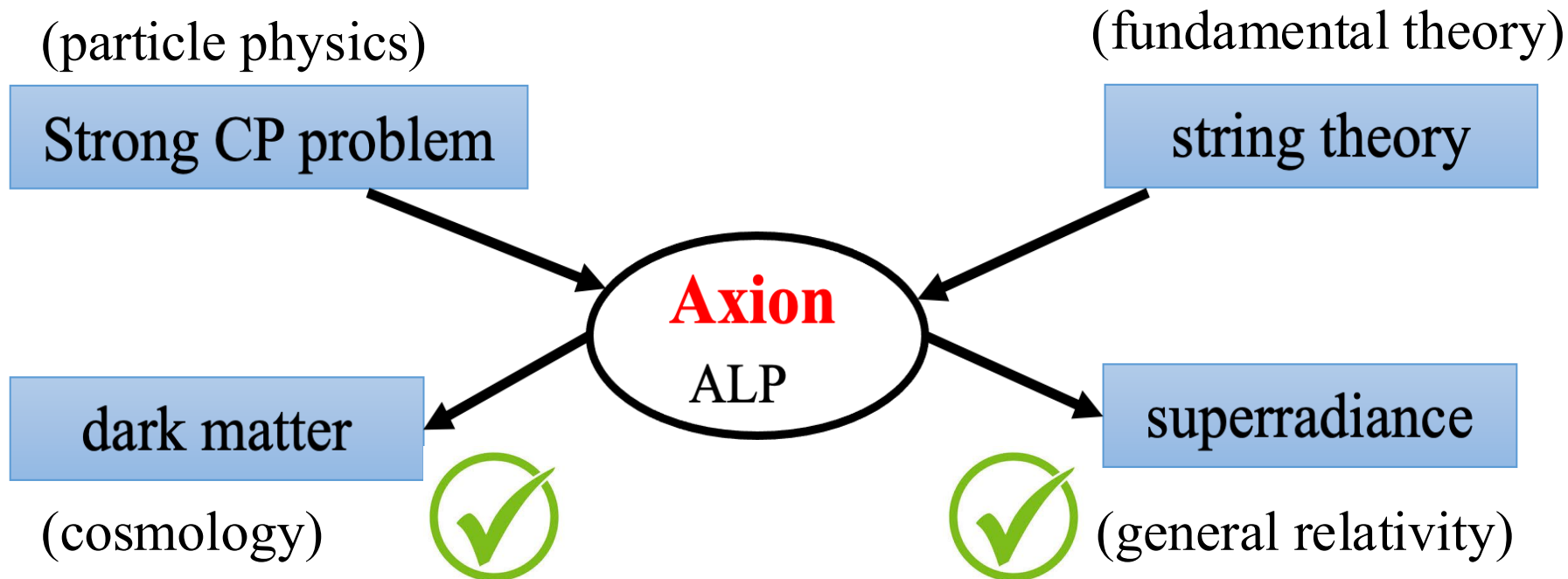
- new DM mechanism beyond freeze out: cosmic phase transition
- new detection method: GW detector (LISA, TianQin, Taiji, aLIGO, FAST, SKA, NanoGrav, Cosmic Explorer...)

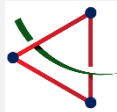




Axion particle cosmology

Ultralight axion is a promising DM candidate.





GW detection of DM

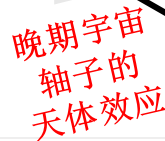
The observation of GW@LIGO initiates a new era of exploring DM and new physics by GW.

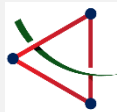
Fundamental Physics and Cosmology with TianQin, arXiv: [2502.20138](https://arxiv.org/abs/2502.20138)

J.Jaekel, V. V. Khoze, M. Spannowsky,
Phys.Rev. D94 (2016) no.10, 103519
Zhaofeng Kang, et.al. arXiv:2101.03795, arXiv:2003.02465
Yan Wang, Chong Sheng Li, and **FPH**, arXiv:2012.03920
FPH, Eibun Senaha Phys.Rev. D100 (2019) no.3, 03501
FPH PoS ICHEP2018 (2019) 397
FPH, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028
FPH, Jiang-Hao Yu, Phys.Rev. D98 (2018) no.9, 095022
FPH, Xinmin Zhang, Phys.Lett. B788 (2019) 288-
Haipeng An, et.al, arXiv: 2208.14857, arXiv:2009.12381, arXiv:2201.05171



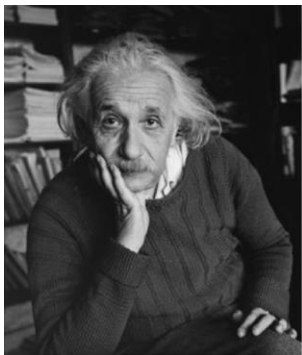
Credit: Gianfranco Bertone et. al.



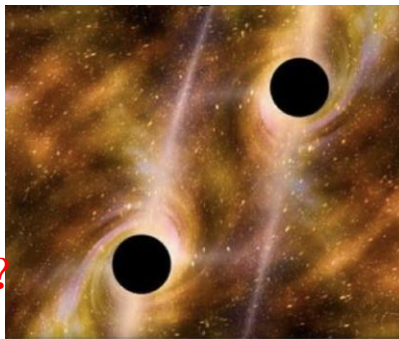


What is GW ?

$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$
 Isolated sources:
 quadrupole radiation
 axion effect?



$$h_{ij} \simeq \frac{2G}{c^4 r} \ddot{Q}_{ij}^{TT}(t - r/c)$$

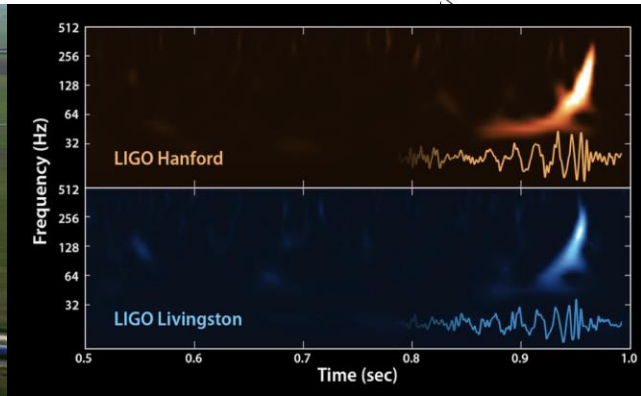
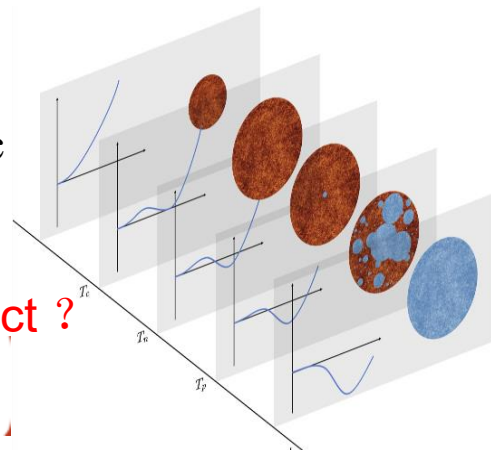


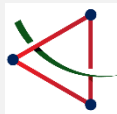
Stochastic sources:

anisotropic stress tensor

axion effect ?

$$\Pi_{ij}(\mathbf{x}, t)$$





GW radiation in a nutshell

The quadruple nature of GW !

EM wave
radiation

$$\ddot{\mathbf{d}} = e\ddot{\mathbf{x}}$$

$$L_{\text{electric quadrupole}} = \frac{1}{20} \ddot{\mathbf{Q}}^2 \equiv \frac{1}{20} \ddot{Q}_{jk} \ddot{Q}_{jk}$$

$$Q_{jk} \equiv \sum_A e_A \left(x_{Aj} x_{Ak} - \frac{1}{3} \delta_{jk} r_A^2 \right)$$

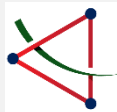
GW
radiation

$$\ddot{\mathbf{d}} = \sum_{\text{particles } A} m_A \ddot{\mathbf{x}}_A = \dot{\mathbf{p}} = 0$$

momentum conservation

$$L_{\text{mass quadrupole}} = \frac{1}{5} \langle \ddot{\mathbf{I}}^2 \rangle \equiv \frac{1}{5} \langle \ddot{I}_{jk} \ddot{I}_{jk} \rangle$$

$$I_{jk} \equiv \sum_A m_A \left(x_{Aj} x_{Ak} - \frac{1}{3} \delta_{jk} r_A^2 \right)$$



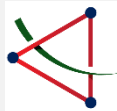
Cosmological source of GW

General GW source from the early universe

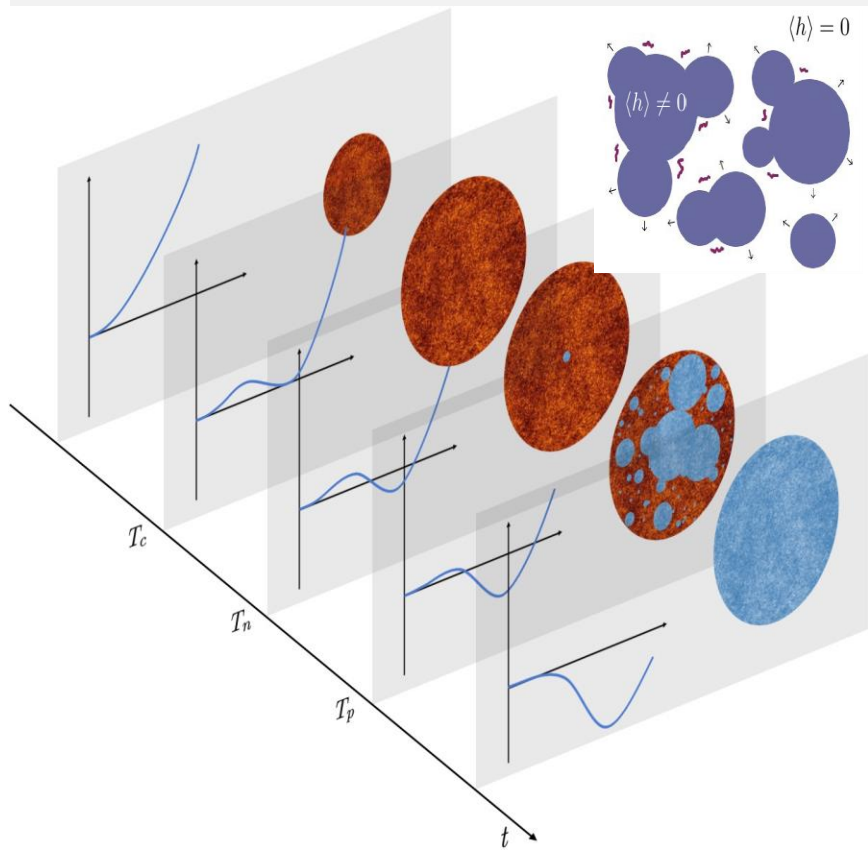
$$\ddot{h}_{ij}(\mathbf{x}, t) + 3H \dot{h}_{ij}(\mathbf{x}, t) - \frac{\nabla^2}{a^2} h_{ij}(\mathbf{x}, t) = 16\pi G \Pi_{ij}(\mathbf{x}, t)$$

各向异性
剪切应力张量

GW sources	Sources of tensor anisotropic stress	General form Π_{ij}
Collisions of bubble walls	scalar field gradients	$[\partial_i \phi \partial_j \phi]^{TT}$
Sound waves and turbulence	bulk fluid motion	$[\gamma^2(\rho + p)v_i v_j]^{TT}$
Primordial magnetic fields	gauge fields	$[-E_i E_j - B_i B_j]^{TT}$
Scalar perturbations	second order scalar perturbations	$\partial_i \Psi, \partial_i \Phi$



Phase transition GW in a nutshell



calculate the finite-temperature effective potential using the thermal field theory: free energy density.

$$V_{\text{eff}}^{(1)}(\bar{\phi}) = \sum_i n_i \left[\int \frac{d^D p}{(2\pi)^D} \ln(p^2 + m_i^2(\bar{\phi})) + J_{\text{B,F}} \left(\frac{m_i^2(\bar{\phi})}{T^2} \right) \right]$$

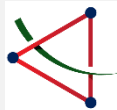
$$S(T) = \int d^4 x \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 + V_{\text{eff}}(\phi, T) \right]$$

$$\Gamma = \Gamma_0 e^{-S(T)}$$

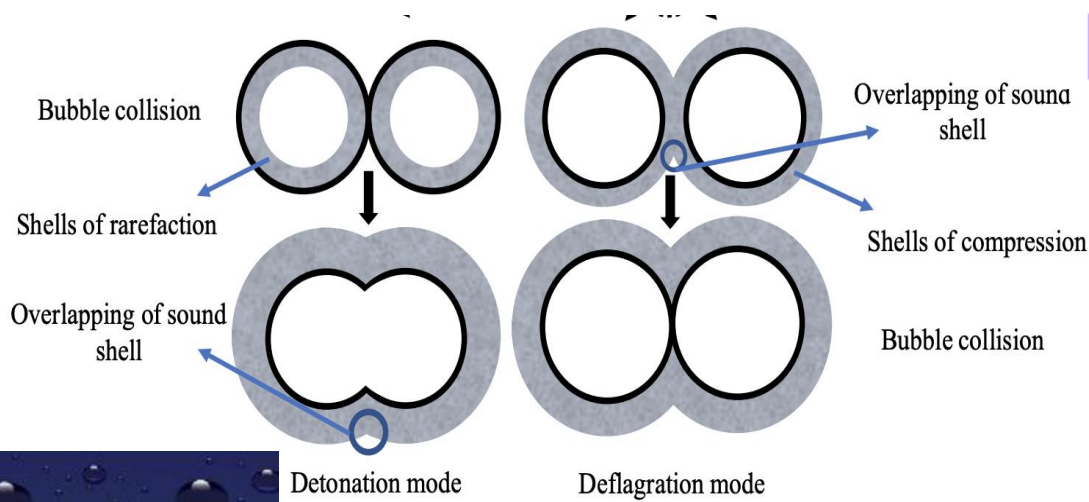
Strong first-order phase transition (SFOPT)

这世上的热闹，源自隧穿

Xiao Wang, **FPH**, Xinmin Zhang, JCAP05(2020)045



Phase transition GW in a nutshell



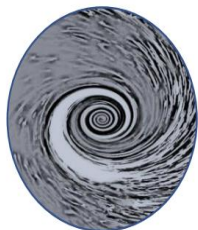
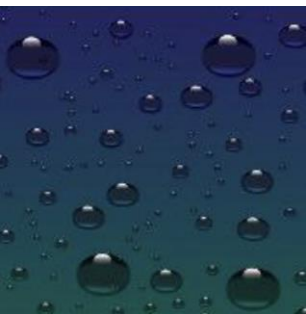
$$[\partial_i \phi \partial_j \phi]^{TT}$$

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

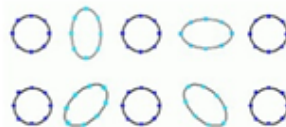
$$h_{ij} \simeq \frac{2G}{c^4 r} \ddot{Q}_{ij}^{TT}(t - r/c)$$

Detonation mode

Deflagration mode



Turbulence



$$[\gamma^2 (\rho + p) v_i v_j]^{TT}$$

E. Witten, Phys. Rev. D 30, 272 (1984)

C. J. Hogan, Phys. Lett. B 133, 172 (1983);

M. Kamionkowski, A. Kosowsky and M. S. Turner, Phys. Rev. D 49, 2837 (1994))

EW phase transition

GW becomes more interesting and realistic after the discovery of

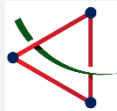
Higgs by LHC and GW by LIGO.

Higgs by LHC and GW by LIGO.

Higgs by LHC and GW by LIGO.

Higgs by LHC and GW by LIGO.

Xiao Wang, **FPH**, Xinmin Zhang, JCAP05(2020)045



Phase transition dynamics

Theory: 相变引力波信号、
相变暗物质、早期宇宙电弱
重子生成机制**最核心却最**
难计算的是泡泡膨胀速度

v_w

Experiment: 实验
上最重要的相变参数
也是泡泡膨胀速度

S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang
, arXiv:2007.10343,
Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith,
arXiv:2009.14295v2
Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903
Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005

Finite-temperature effective potential

$V_{eff}(\phi, T)$

α

T_p

$R_* H_*$

- (1). Daisy resummation problem: **Pawani scheme vs. Arnold scheme**
- (2). Gauge dependence problem: **see Michael J. Ramsey-Musolf's works**
- (3). No perturbative calculations: **lattice calculations**
and dim-reduction method: **by D. Weir, Michael J. Ramsey-Musolf et.al**

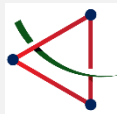
Bubble wall velocity

v_w

Energy budget

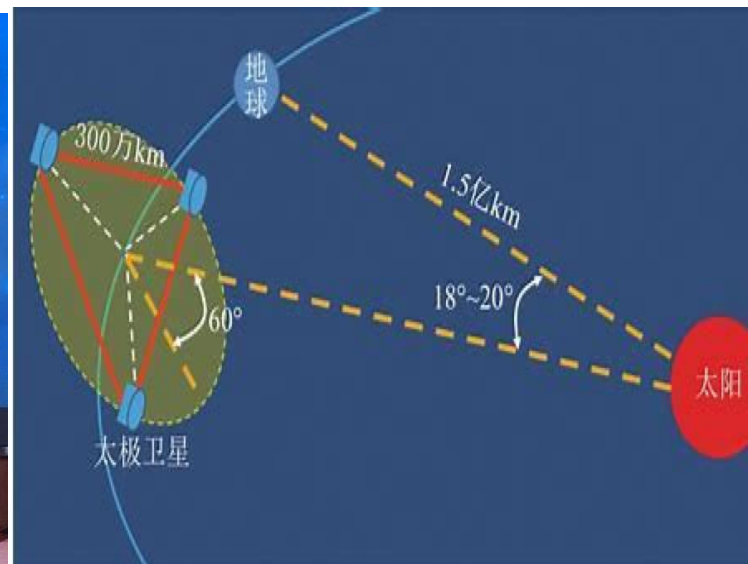
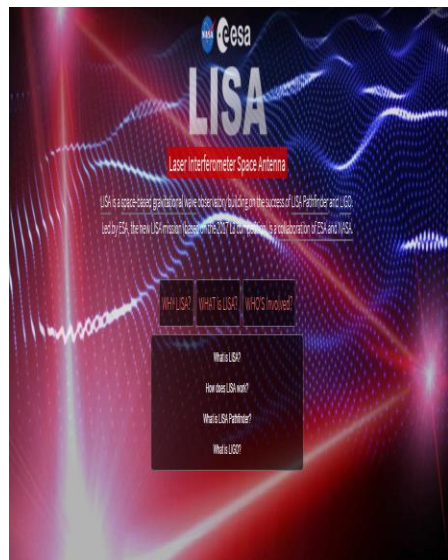
κ

F. Giese, T. Konstandin, K. Schmitz and J. van de Vi
, arXiv:2010.09744
Xiao Wang, **FPH** and Xinmin Zhang,
Phys.Rev.D 103 (2021) 10, 103520
Xiao Wang, Chi Tian, **FPH**, JCAP 07 (2023) 006



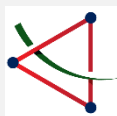
GW experiments

LISA/TianQin/Taiji ~2034



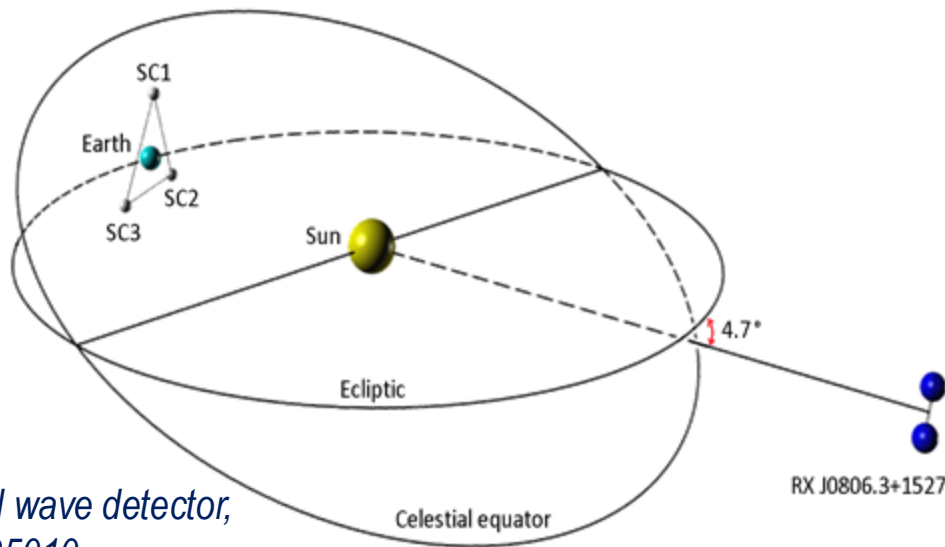
“天琴”

“Harpe in space”

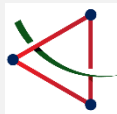


TianQin

- Expected around 2035
- Geocentric orbit, normal triangle constellation, radius $\sim 10^5$ km
- Unique frequency band 10^{-4} -1 Hz, easier for deployment, tracking, control, and communication

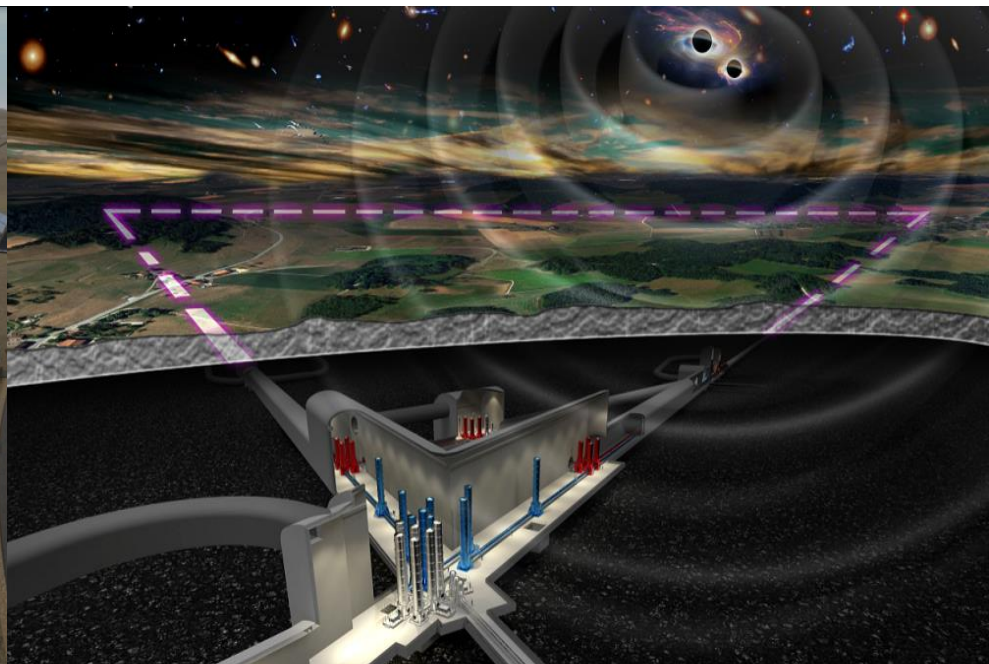


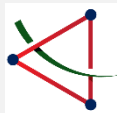
*J. Luo et al. TianQin: a space-borne gravitational wave detector,
Class. Quant. Grav. 33 (2016) no.3, 035010.*



GW experiments

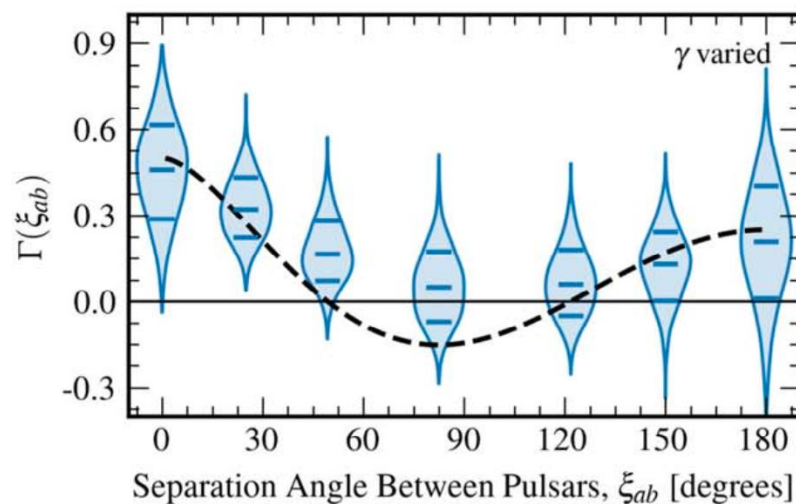
**Next generation: Einstein telescope
Cosmic Explorer**





Radio telescope and pulsar timing array

2023 June 29th: NANOGrav, EPTA, InPTA, Parkes PTA, CPTA



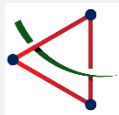
Hellings-Downs correlation curve
First observation of stochastic GW

FAST

High sensitivity sub

SKA

μJy

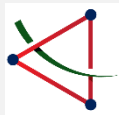


Outline

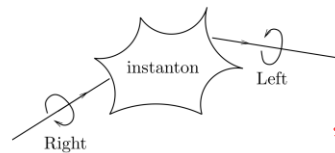
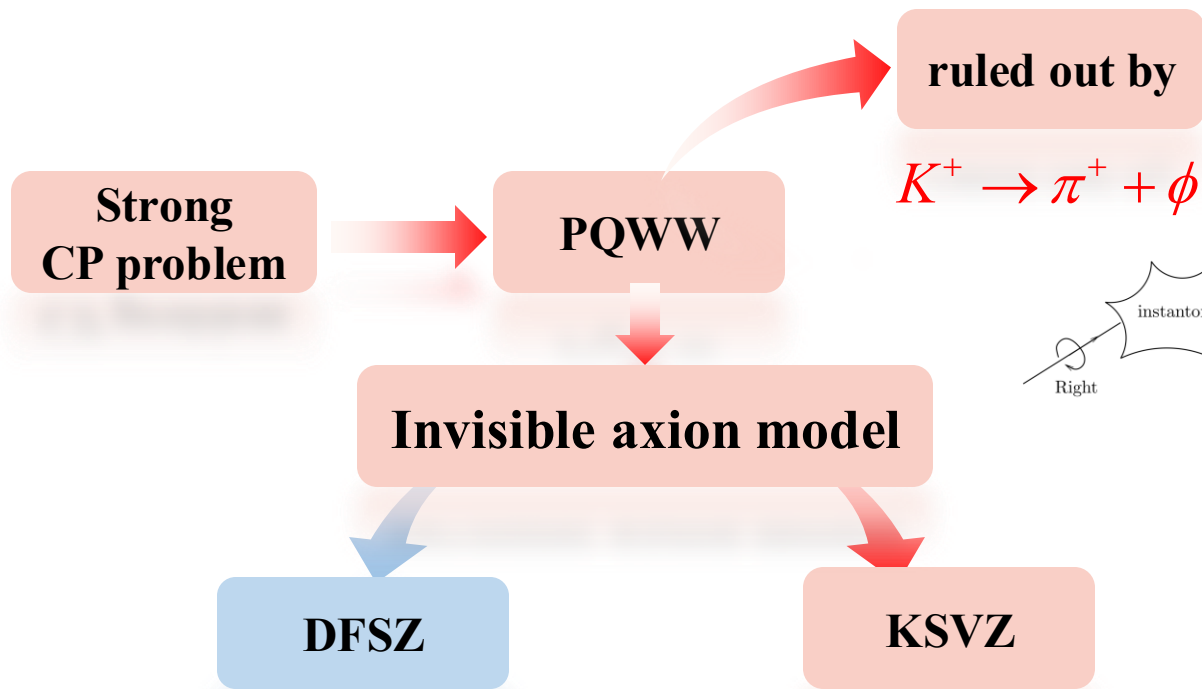
1. Axion and axion dark matter (DM), Gravitational wave (GW)
2. DFSZ axion and its phase transition GW signals
3. μeV axion and radio signals of **axion** DM (multi-messenger)
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“Ultralight” DM

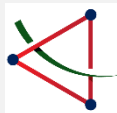


Strong CP problem and QCD axion



't Hooft, G. Phys.
Rev.Lett. 37, 8 (1976)

M. Dine, W. Fischler, and M. Srednicki, Physics letters B 104, 199 (1981).



GW detection of DFSZ axion model

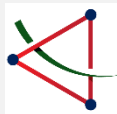
The U(1) Peccei-Quinn symmetry breaking might be a SFOPT process, which could produce detectable phase transition GW.

$$\begin{aligned} V_{\text{tree}} = & -\mu_1^2 |H_u|^2 - \mu_2^2 |H_d|^2 + \lambda_1 |H_u|^4 + \lambda_2 |H_d|^4 + \lambda_4 |H_u^\dagger H_d|^2 \\ & - \mu_3^2 |\sigma|^2 + \lambda_3 |\sigma|^4 + \lambda_{12} |H_u|^2 |H_d|^2 + \lambda_{13} |\sigma|^2 |H_u|^2 \\ & + \lambda_{23} |\sigma|^2 |H_d|^2 + \left(\lambda_5 \sigma^2 \tilde{H}_u^\dagger H_d + \text{h.c.} \right), \\ \sigma = & \frac{1}{\sqrt{2}} \left(v_\sigma + \sigma^0 + i\eta_\sigma^0 \right). \end{aligned}$$

$$V_{\text{eff}}(\sigma^0, T) \equiv V_{\text{tree}}(\sigma^0) + V_{\text{CW}}(\sigma^0) + V_T(\sigma^0, T),$$

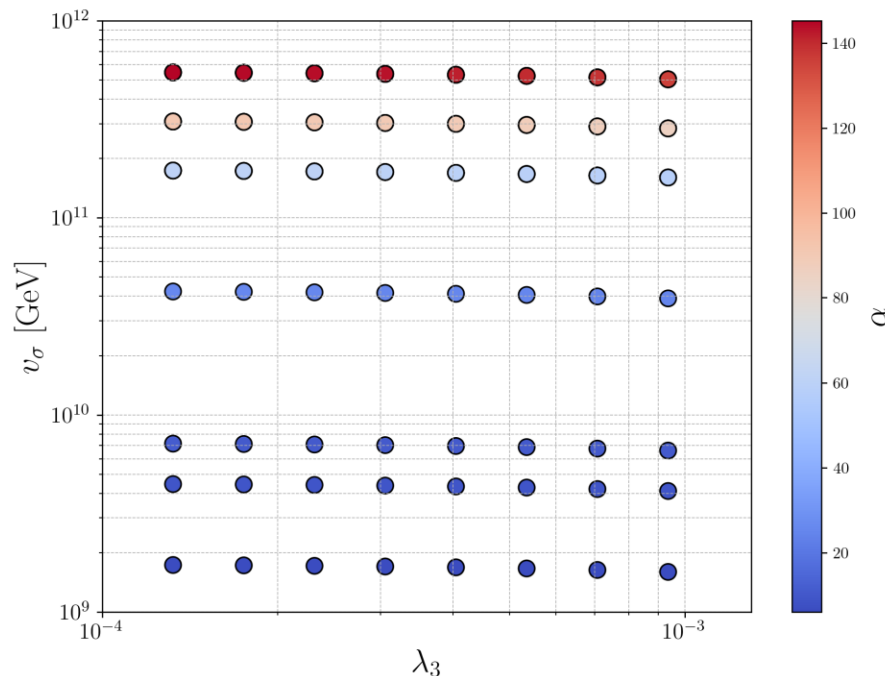


Aidi Yang, FPH*,
arXiv:2404.18703, JCAP (2025)



GW detection of DFSZ axion model

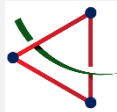
There exists a substantial parameter space for realizing a SFOPT in the DFSZ model.



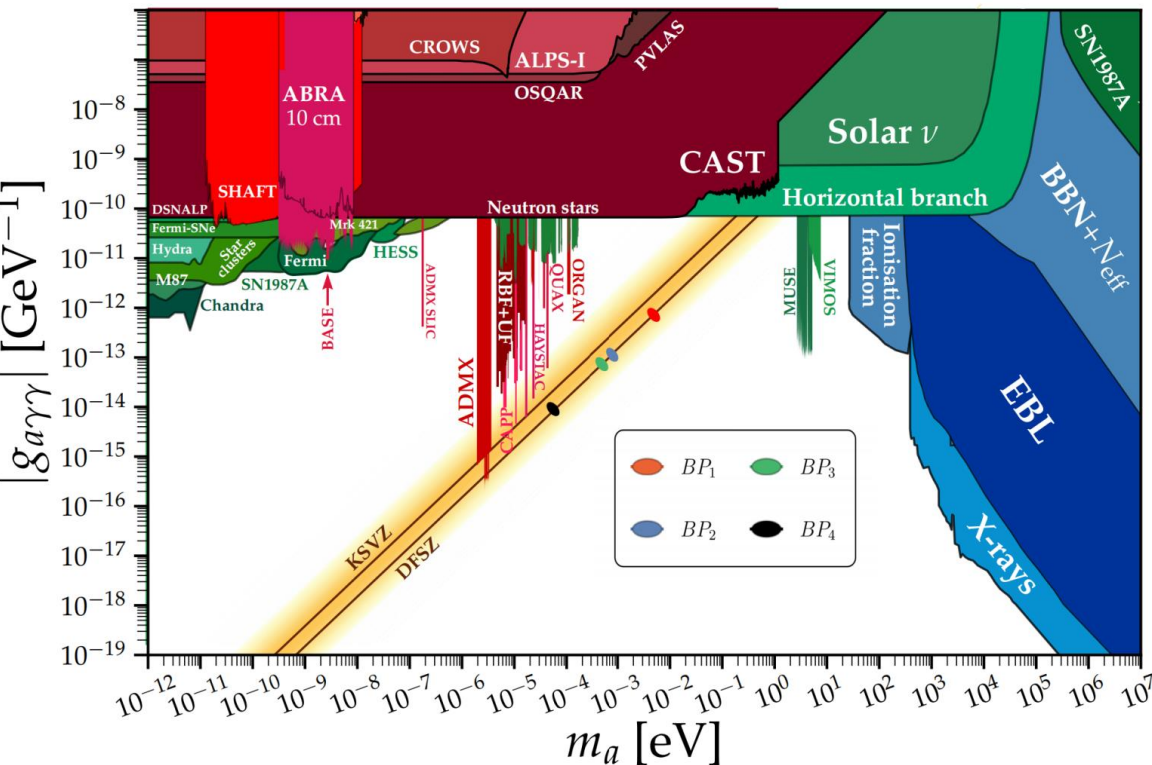
Aidi Yang, FPH*,
arXiv:2404.18703, JCAP (2025)

Each point plotted in the figure represents a set of parameters for which a SFOPT occurs. The color of each point in the figure indicates the value of α , with the color bar ranging from blue to red.

Parameter space exploration results of the SFOPT in the DFSZ axion model.



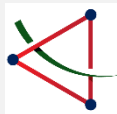
GW detection of DFSZ axion model



Aidi Yang, FPH*,
arXiv:2404.18703, JCAP (2025)

All four points fall within the expected yellow region which represents the range of axion masses and coupling constants that have not yet been excluded by ALP-photon coupling experiments.

The m_a and $g_{a\gamma\gamma}$ obtained from all four benchmark points



GW detection of DFSZ axion model

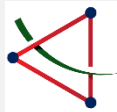
10^{-3} - 10^{-5} eV DFSZ axion

	f_{PQ} [GeV]	m_a [eV]	$g_{a\gamma\gamma}$ [GeV $^{-1}$]
BP ₁	1.24×10^9	4.59×10^{-3}	7.01×10^{-13}
BP ₂	7.01×10^9	8.12×10^{-4}	1.21×10^{-13}
BP ₃	1.62×10^{10}	3.51×10^{-4}	5.37×10^{-14}
BP ₄	1.57×10^{11}	3.63×10^{-5}	5.54×10^{-15}

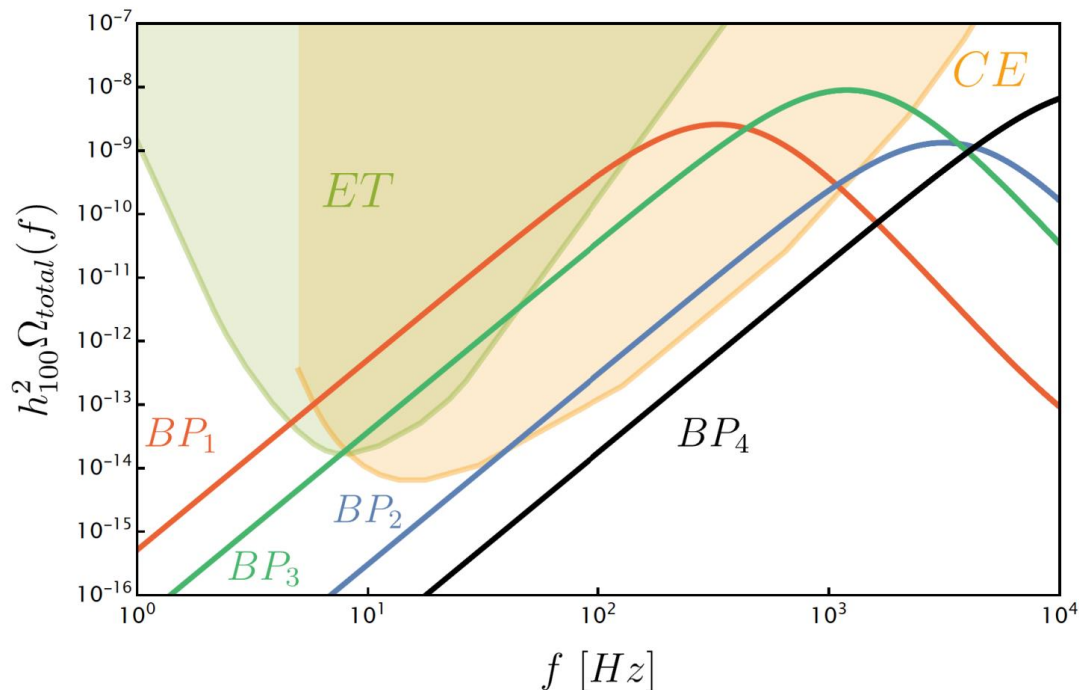
Aidi Yang, FPH*,
arXiv:2404.18703, JCAP

The axion mass and axion-photon coupling constant predicted by the DFSZ model are consistent with current experimental constraints.

The corresponding m_a and $g_{a\gamma\gamma}$ values.



GW detection of DFSZ axion model

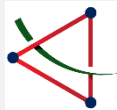


$$\text{SNR}_{\text{gw}} = \sqrt{T_t \int_{f_{\min}}^{f_{\max}} df \left(\frac{h_{100}^2 \Omega_{\text{gw}}}{h_{100}^2 \Omega_{\text{sens}}} \right)^2},$$

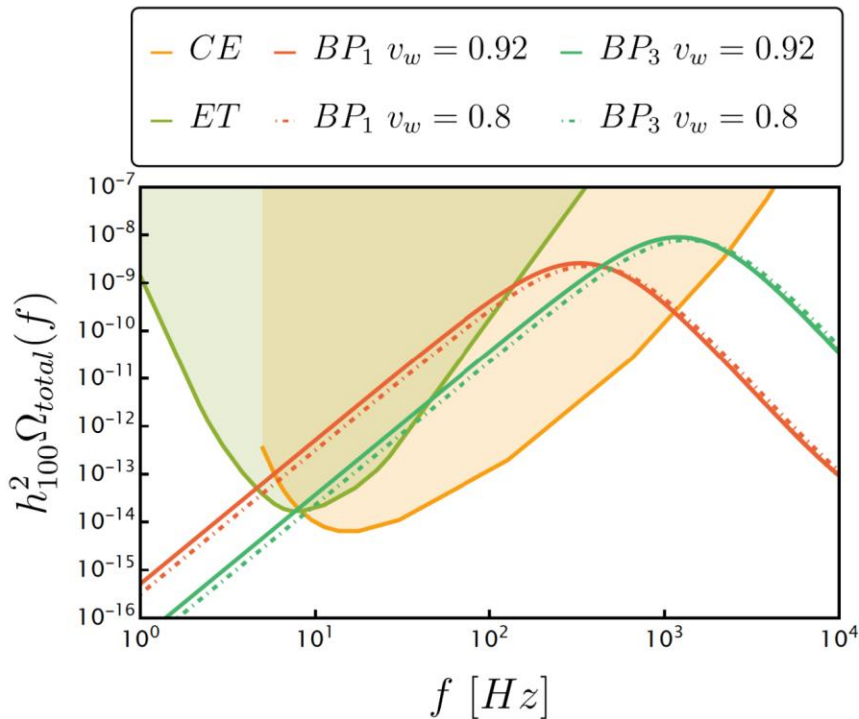
$$\text{SNR : } BP_1(21794.10) \quad BP_2(46.19) \\ BP_3(3964.97) \quad BP_4(0.15)$$

The SNR values of BP_1 , BP_2 , and BP_3 exceed the CE SNR threshold of 8, indicating that they can be detected by the CE detector.

Aidi Yang, FPH*, arXiv:2404.18703, JCAP (2025)



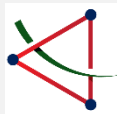
GW detection of DFSZ axion model



To assess the constraining power of CE on these phase transition parameters more precisely, we need to quantify the uncertainties of each parameter. FM analysis is a powerful statistical tool widely used to estimate the precision with which model parameters can be determined from a given set of observations.

Aidi Yang, FPH*,
arXiv:2404.18703, JCAP (2025)

The differences in the frequency and strength of the GW spectra for DFSZ model with $v_w=0.92$ and $v_w=0.8$.

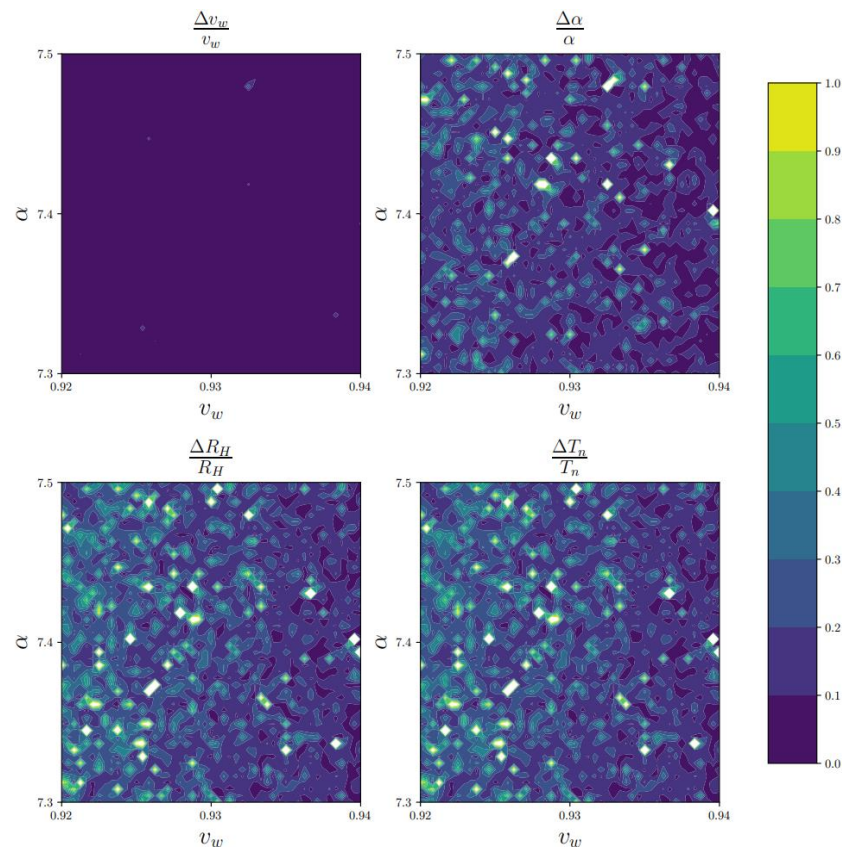


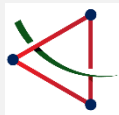
GW detection of DFSZ axion model

Relative uncertainty with phase transition dynamics parameters bubble wall speed v_w , transition strength α , Hubble-scaled mean bubble spacing R_H , and nucleation temperature T_n . From dark blue to light green, the relative uncertainties gradually increase.

By Fisher matrix analysis,
CE will be most sensitive to v_w

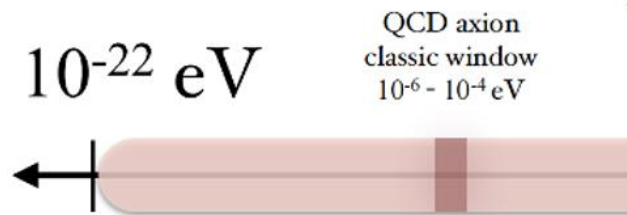
实验上最先确定的相变参数将是
泡泡膨胀速度 v_w



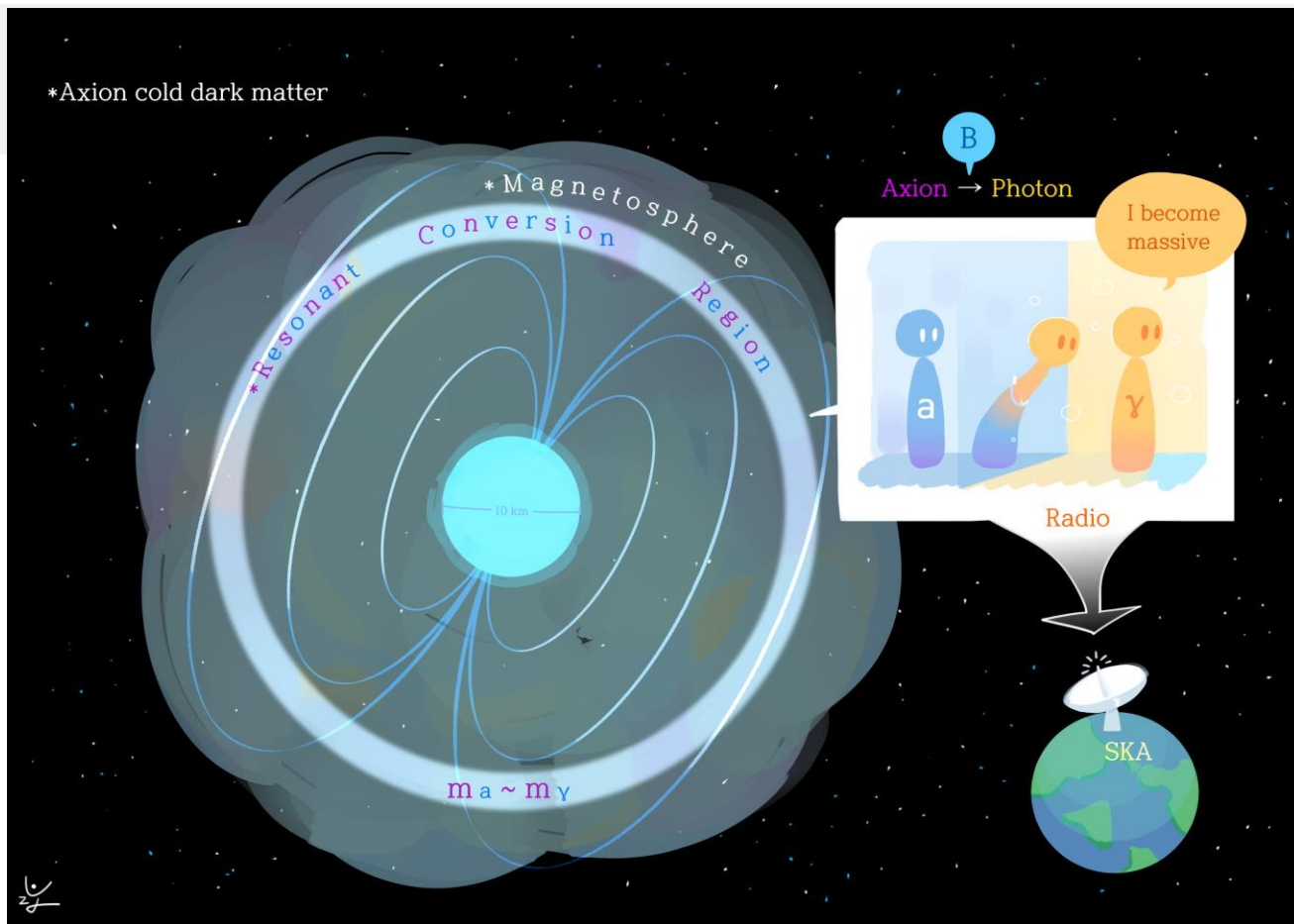
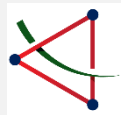


Outline

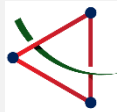
1. Axion and axion dark matter (DM), Gravitational wave (GW)
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3. μeV axion and radio signals of axion DM (multi-messenger)
4. 10^{-12} - 10^{-17} eV axion: GW and pulsar timing measurement
5. 10^{-21} eV fuzzy axion DM GW
6. Summary and outlook



“Ultralight” DM



FPH, K. Kadota, T. Sekiguchi, H. Tashiro, Phys.Rev. D97 (2018) no.12, 123001



天体磁层中的轴子-光子共振过程

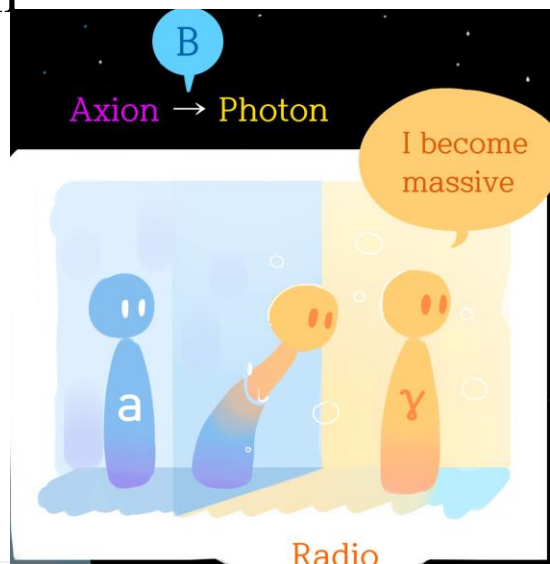
$$L_{\text{int}} = \frac{1}{4} g \tilde{F}^{\mu\nu} F_{\mu\nu} a = -g \mathbf{E} \cdot \mathbf{B} a,$$

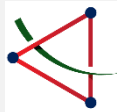
Massive Photon: In the magnetosphere of the neutron star, photon obtains effective mass in the plasma.

$$m_{\gamma}^2 = \omega_{\text{plasma}}^2 = 4\pi\alpha \frac{n_e}{m_e}$$

$$B(r) = B_0 \left(\frac{r}{r_0} \right)^{-3} \quad n_e(r) = n_e^{\text{GJ}}(r) = 7 \times 10^{-2} \frac{1s}{P} \frac{B(r)}{1 \text{ G}} \frac{1}{\text{cm}^3}$$

Thus, the photon mass is location dependent, and within some region





Line-like radio signal from axion DM

$$\nu_{\text{peak}} \approx \frac{m_a}{2\pi} \approx 240 \frac{m_a}{\mu\text{eV}} \text{MHz} \quad 1 \text{ GHz} \sim 4 \mu\text{eV}$$

FAST: 70MHz–3GHz, SKA: 50MHz–14GHz, GBT:0.3–100GHz

Radio telescopes can probe axion mass of 0.2–400 μeV

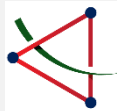
Signal: For a trial parameter set, $S_\gamma \sim 0.51 \mu\text{Jy}$.

Sensitivity: $S_{\min} \sim 0.48 \mu\text{Jy}$ for the SKA1

$S_{\min} \sim 0.016 \mu\text{Jy}$ for SKA2 with 100 hours observation time.

SKA-like experiment can probe the axion DM and the axion mass which corresponds to peak frequency.

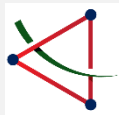
Working in progress on more delicate study.



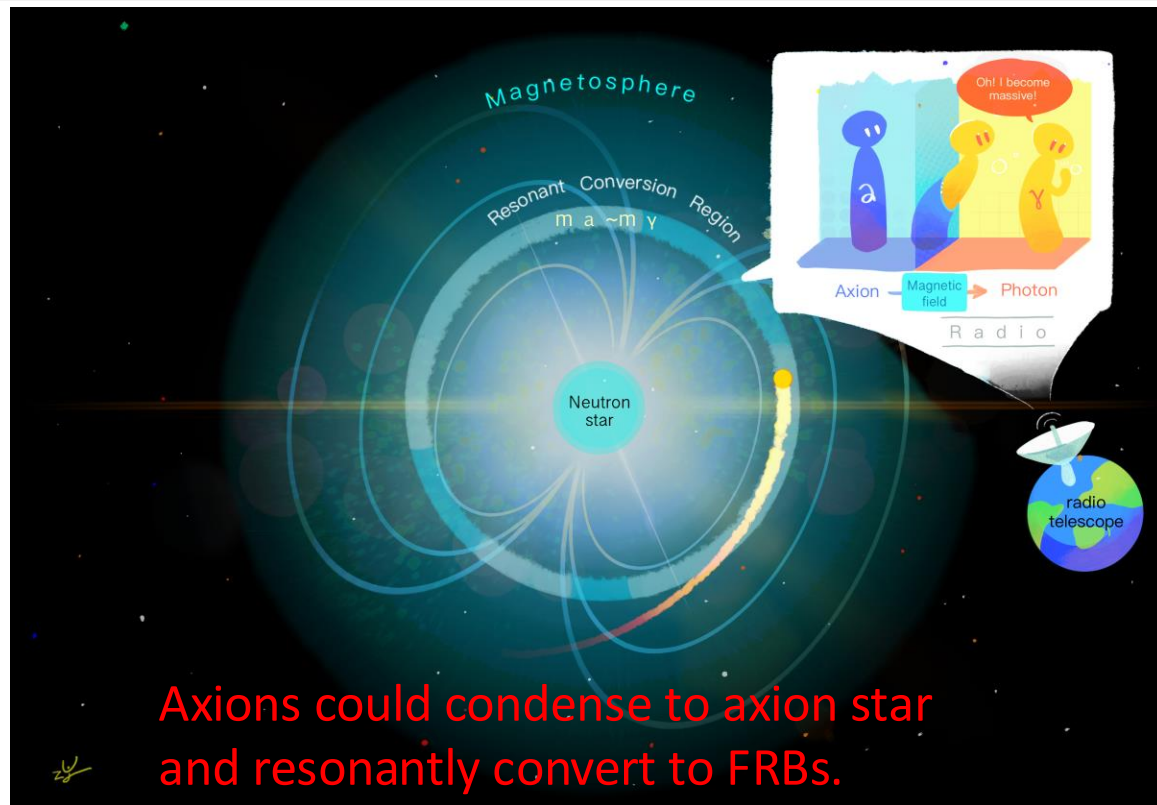
探测暗物质的热门新方法

FPH, K. Kadota, T. Sekiguchi, H. Tashiro, Phys.Rev. D97 (2018) no.12, 123001, arXiv:1803.08230, **Cited by 119 times**

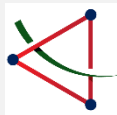
- Promising approaches at SKA&FAST, more and more nice works
- more details see the timely new review papers
- ✓ Physics Briefing Book :
Input for the European Strategy for Particle Physics Update 2020, [arXiv:1910.11775]
- ✓ 2021 white paper by EuCAPT [arXiv:2110.10074]
- ✓ Pierre Sikivie, **Rev.Mod.Phys.93(2021)1,015004**
- ✓ 2022 Snowmass papers: [arXiv:2203.06380, arXiv: 2203.07984]
- ✓ Phys. Rept. 1052(2024)1-48
- ✓ Science Advances Volume 8, Issue 8



Generalize to axion star/FRBs

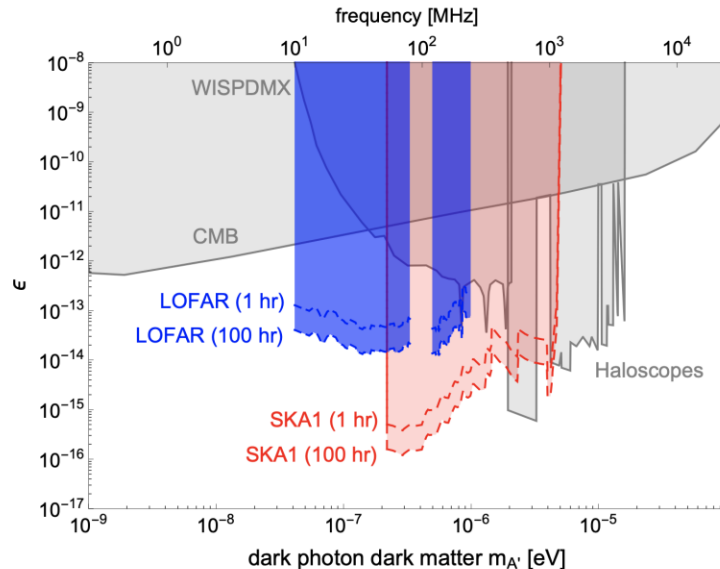
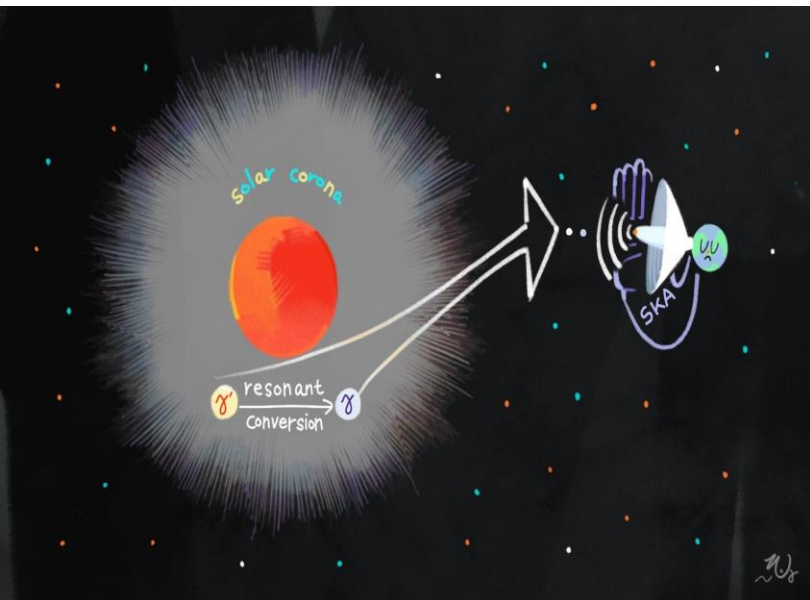


James Buckley, Bhupal Dev, Francesc Ferrer, **FPH**, *Phys.Rev.D* 103 (2021) 4, 043015

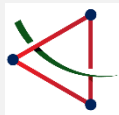


Generalize to dark photon DM case

Haipeng An, **FPH**, Jia Liu, Wei Xue, *Phys. Rev. Lett.* **126**, 181102 (2021)



See Jia Liu, Haipeng An's works for further study with real data



Outline

1. Axion and axion dark matter (DM), Gravitational wave (GW)
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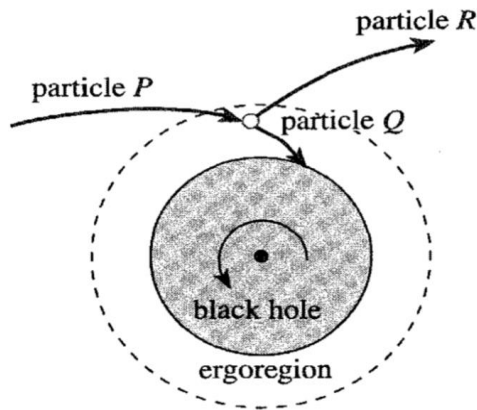
“Ultralight” DM



What is superradiance?

When Klein (-Gordon) meets Kerr——superradiance

$$\Delta \frac{d}{dr} \left(\Delta \frac{dR}{dr} \right) + \left[\omega^2 (r^2 + a^2)^2 - 4aMrm\omega + a^2m^2 - \Delta (m_a^2 r^2 + a^2\omega^2 + \lambda) \right] R = 0$$



Penrose '69 '71
Zel'dovich '72
Starobinsky '73

物体坍缩形成黑洞后，将最后终结于由克尔解描述的一个稳态。

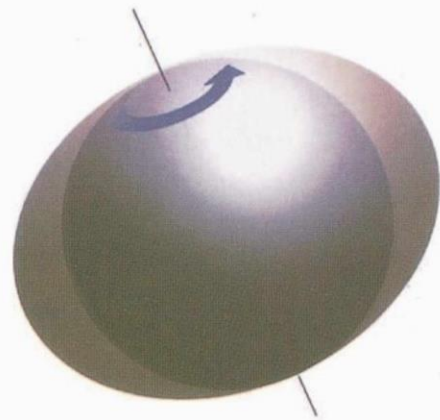


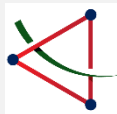
图6.8 当一个“克尔”黑洞的旋转加快时它的赤道附近鼓了出来。而零旋转的黑洞是

书名：时间简史（插图本）

作者：【英】霍金

Exponential growth solution of Klein-Gordon equation due to the boundary condition at the

horizon of Kerr BH. Ultralight axion can form axion cloud around rotating BH, Gravitational atom (GA).



GW of ultralight DM from black hole:

Gravitational atom from superradiance

1. Axions can annihilate to GW
2. Energy-level transition of GA

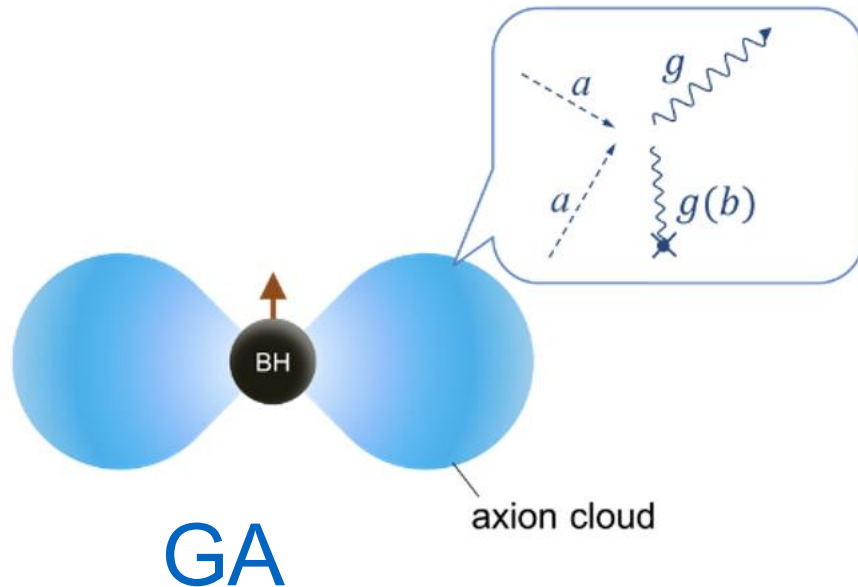
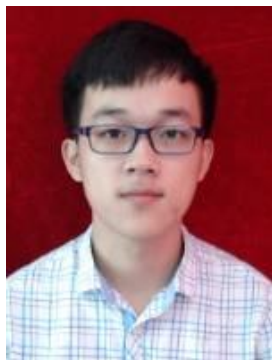


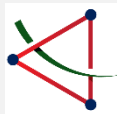
Jing Yang, **FPH**

Jing Yang, Ning Xie, **FPH** arXiv:2306.17113

Ning Xie, **FPH** SCPMA Vol.66, No.1(2024)

Ning Xie, **FPH***, arXiv:2503.10347;





GW of ultralight DM from black hole

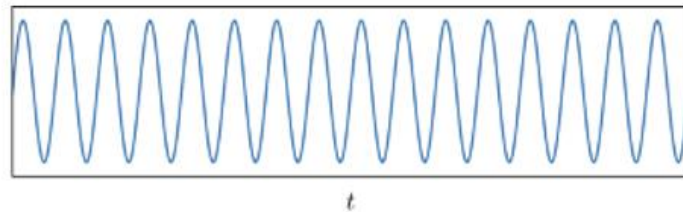
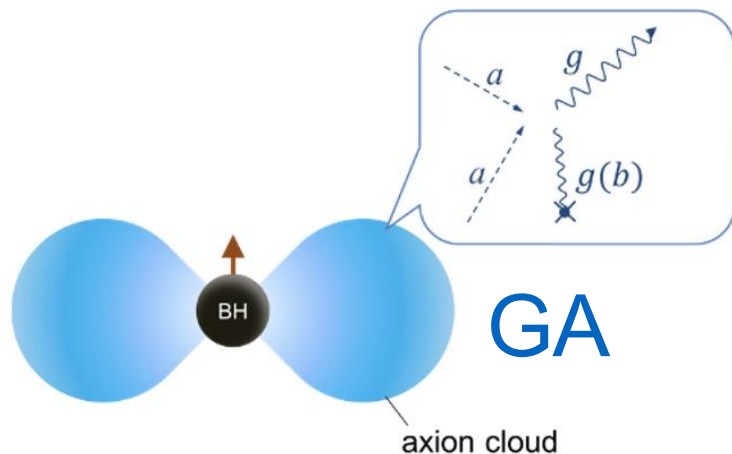
Axions can annihilate to GW

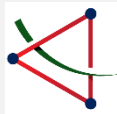
A. Arvanitaki and S. Dubovsky, Phys. Rev. D 83, 044026 (2011)

R. Brito, V. Cardoso and P. Pani, Class. Quant. Grav. 32, no.13, 134001 (2015)

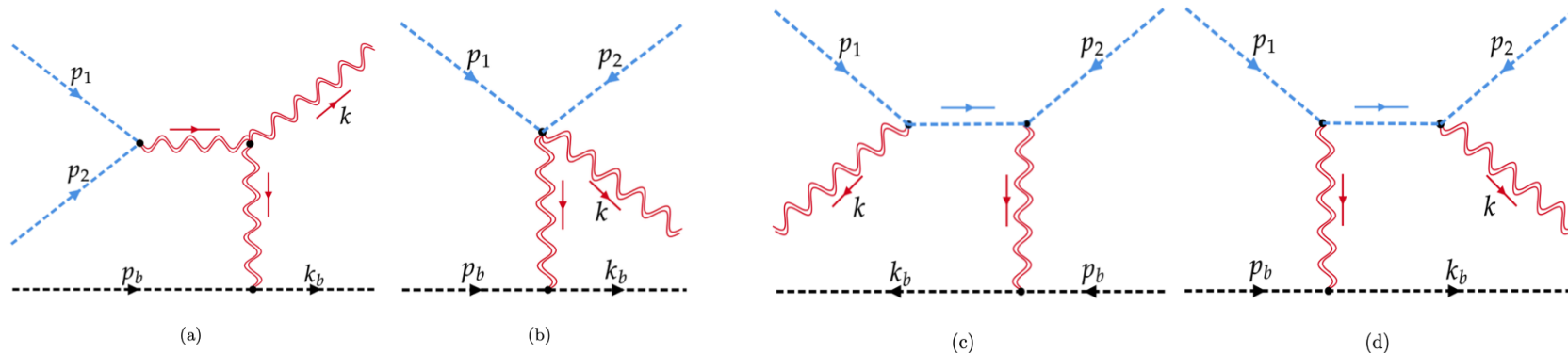
H. Yoshino and H. Kodama, PTEP 2014, 043E02 (2014)

Jing Yang, FPH, Phys.Rev.D 108 (2023) 10, 103002

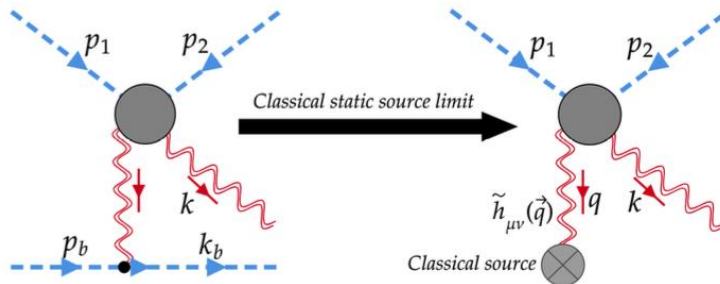




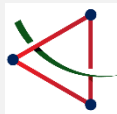
Microscopic physics



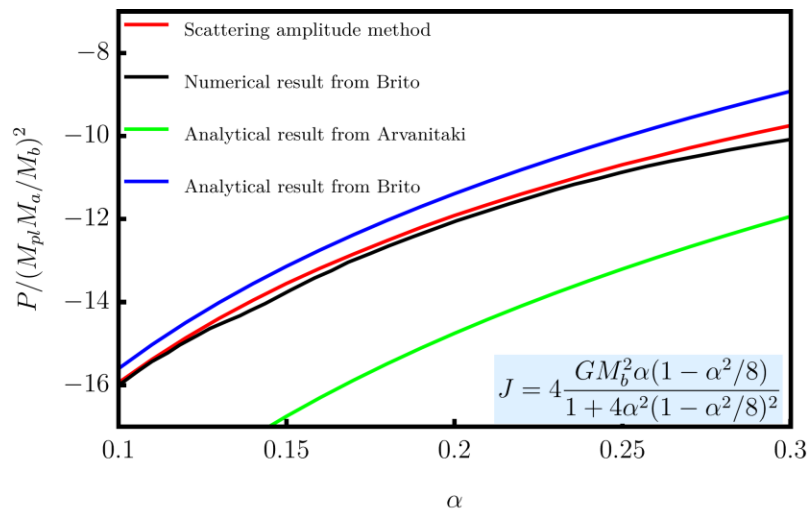
$$M(p_b, p_1, p_2 \rightarrow k, k_b)$$



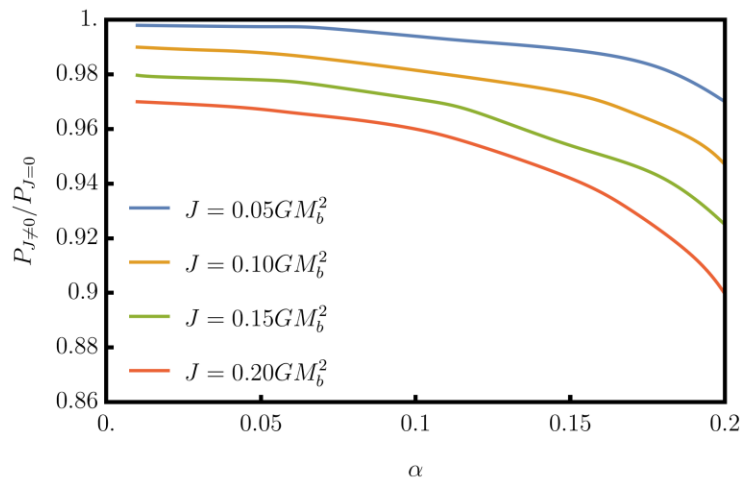
Jing Yang, FPH,
Phys.Rev.D 108 (2023) 10, 103002



GW radiation from axion annihilation

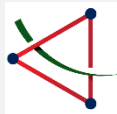


$$\alpha = GM_b m_a \quad M_b = 100 M_{\text{sun}} \quad M_a = M_{\text{sun}}$$

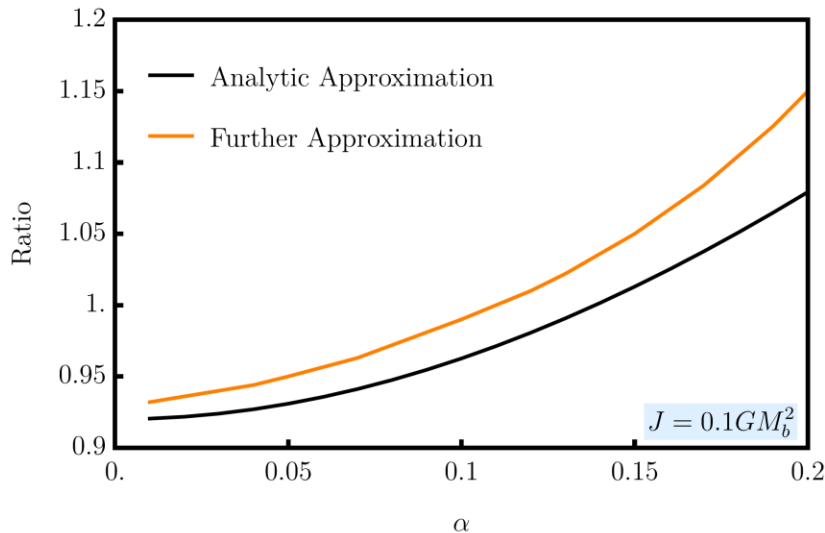


Jing Yang, FPH,
Phys.Rev.D 108 (2023) 10, 103002

- ✓ monochromatic GW signal $\omega_{\text{ann}} \sim 2 m_a$
- ✓ gradually depletion of axion cloud (DC) and reduce GA mass



GW radiation from axion annihilation

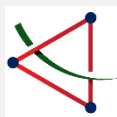


$$P = \frac{(M_a/\text{GeV})^2 \alpha^{14}}{(M_b/\text{GeV})^6 (2 + \alpha^2)^{11} (4 + \alpha^2)^4} \left[(M_b/\text{GeV})^4 (9.671 \times 10^{41} + 5.577 \times 10^{42} \alpha^2 + 1.474 \times 10^{43} \alpha^4 + 2.361 \times 10^{43} \alpha^6) + J(M_b/\text{GeV})^2 \alpha (-3.839 \times 10^{80} - 2.111 \times 10^{81} \alpha^2 - 5.329 \times 10^{81} \alpha^4 - 8.165 \times 10^{81} \alpha^8) + J^2 \alpha^2 (3.809 \times 10^{118} + 2.184 \times 10^{119} \alpha^2 + 5.799 \times 10^{119} \alpha^4 + 9.450 \times 10^{119} \alpha^6) \right] \text{GeV}^2.$$

$$\alpha = GM_b m_a \quad M_b = 100 M_{\text{sun}} \quad M_a = M_{\text{sun}}$$

Jing Yang, FPH,
Phys.Rev.D 108 (2023) 10, 103002

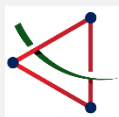
- ✓ monochromatic GW signal $\omega_{\text{ann}} \sim 2 m_a$
- ✓ gradually depletion of axion cloud (DC) and reduce GA mass



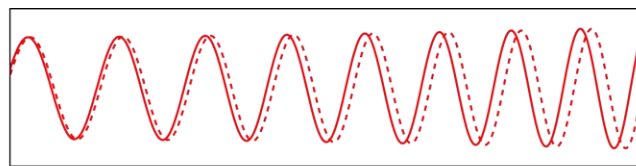
Advantages

- ✓ Simple and straightforward.
- ✓ Easy to include Kerr metric effects.
- ✓ Microscopic physics is intuitive.
- ✓ It is clearly and simple to demonstrate the analytic approximation formulae.

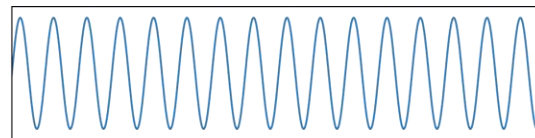
Important for the GW and axion search. More precise calculations and more broad applications are working in progress.



Imprints of axions on GW



t



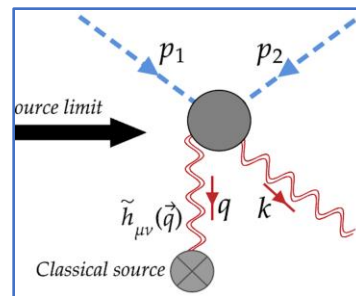
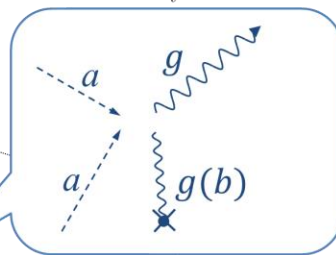
t



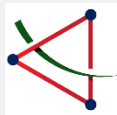
NS

BH

axion cloud



Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024)



Imprints of axions on GW

Without ultralight axions $-\frac{dE_0}{dt} = \mathcal{P}_{\text{GW}} \quad \mathcal{P}_{\text{GW}} = \frac{32}{5} \mu^2 r^4 \omega^6$

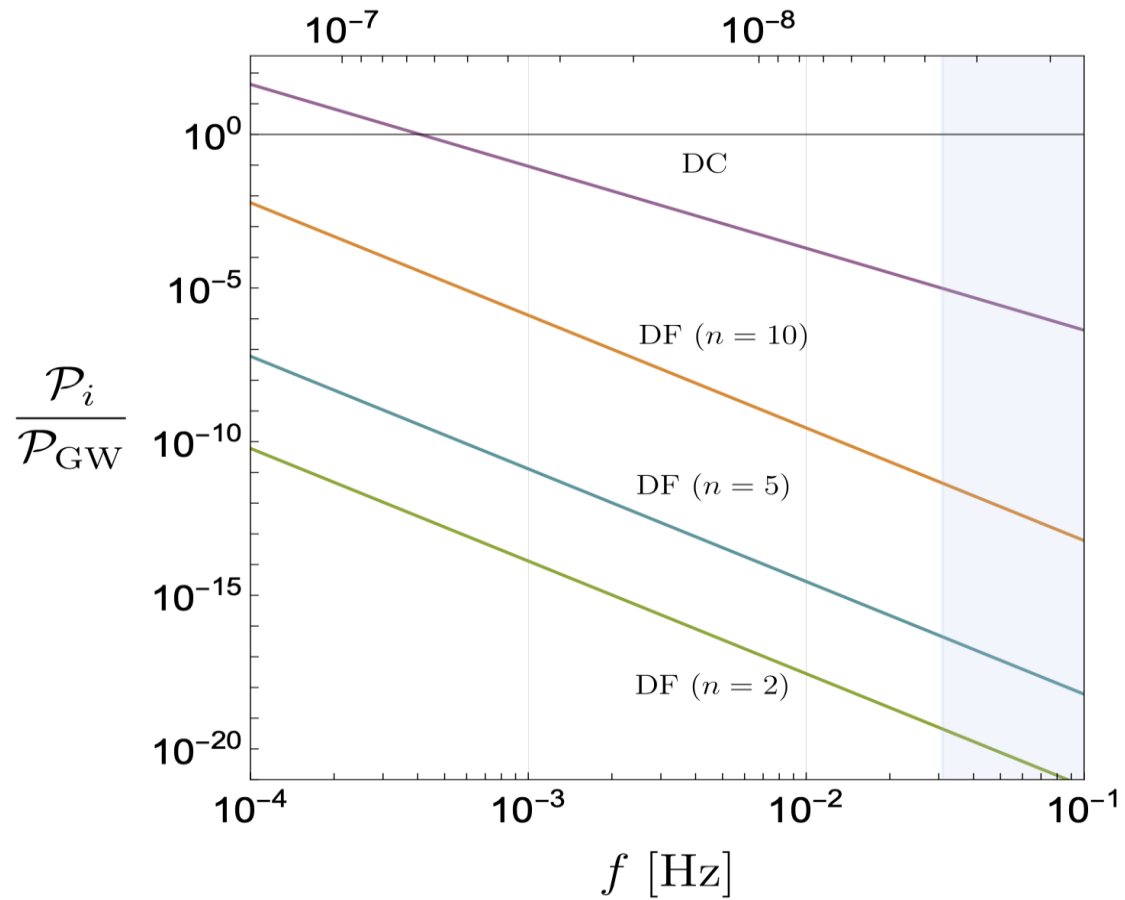
With ultralight axions

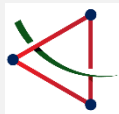
$$-\frac{dE}{dt} = (\mathcal{P}_{\text{GW}} + \boxed{\mathcal{P}_{\text{DC}}} + \mathcal{P}_{\text{DF}} + \mathcal{P}_{\text{DR}})$$

dynamical friction (DF), depletion of axion cloud (DC), dipole radiation(DR)

Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024)

$$M = 100 \text{ M}_\odot, \quad m_{\text{NS}} = 1.5 \text{ M}_\odot$$

 $r \text{ [pc]}$


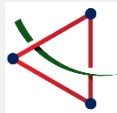


Imprints of axions on GW

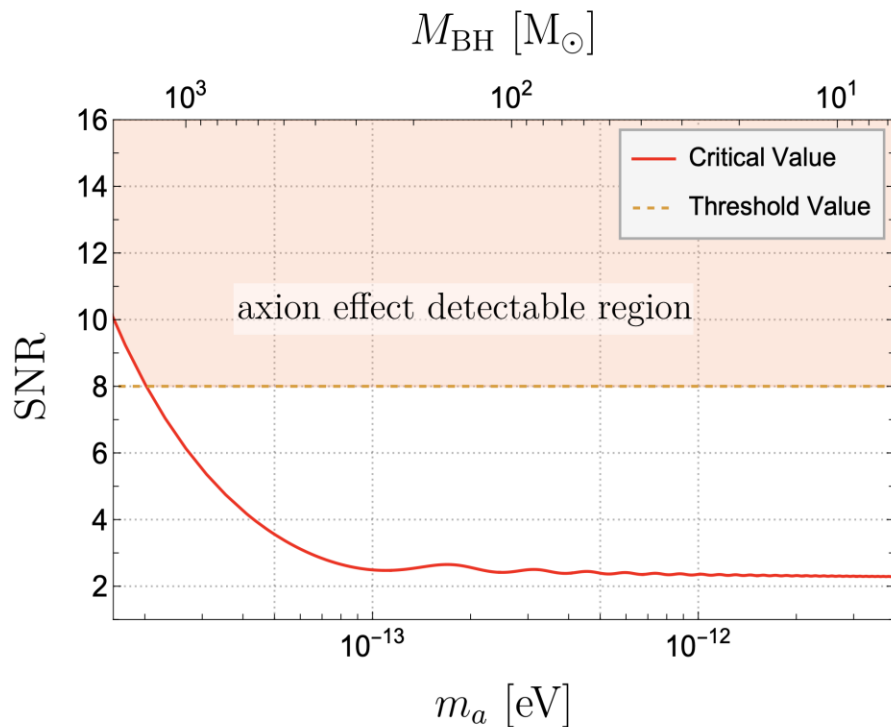
$$\frac{dr}{dt} = \left(-\frac{M m_{\text{NS}}}{2r^2} \right)^{-1} (\mathcal{P}_{\text{GW}} + \mathcal{P}_{\text{DC}} + \mathcal{P}_{\text{DF}} + \mathcal{P}_{\text{DR}})$$

$$\Delta\phi \sim 15\pi \left(\frac{m_a}{10^{-12} \text{ eV}} \right) \left(\frac{f_T}{10^{-2} \text{ Hz}} \right) \left(\frac{T}{5 \text{ yrs}} \right)^2$$

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Complementary search: GW+PTA

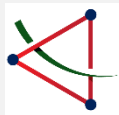


Axions modify the rate of binary period change

$$\Delta\dot{P} = \left| \dot{P} - \dot{P}_{\text{vac}} \right| \approx 10^{-12} \text{ s/s}$$

Future **Pulsar timing measurement** precision, such as SKA

$$10^{-15} \text{ s/s}$$

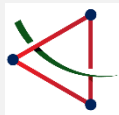


Outline

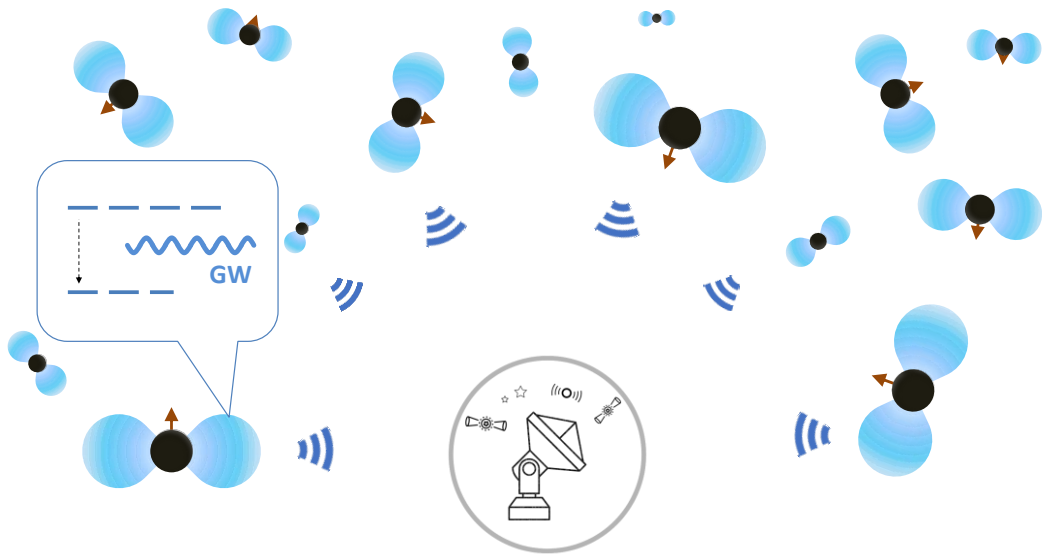
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“Ultralight” DM



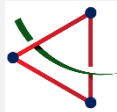
Fuzzy axion (DM) particles



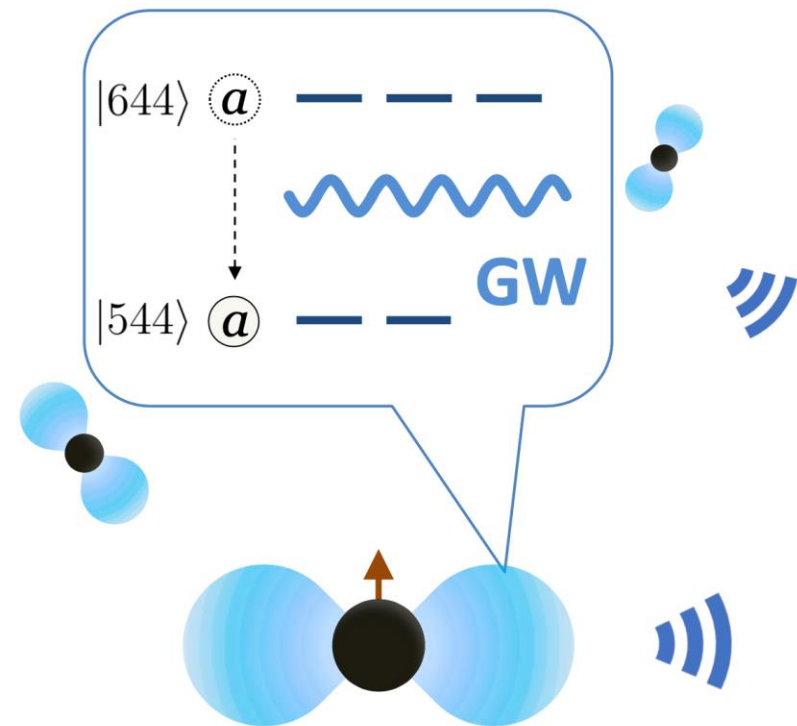
The cosmic populated SMBHs dressed with axion cloud as a natural source of nano-Hertz SGWB. The energy level transition process can radiate GWs continuously, which naturally fall in nano-Hertz frequency band.

Consequently, the PTA could detect this new source which provides a new approach to probe ultralight axion DM and isolated BHs.

Jing Yang, Ning Xie, **FPH***, JCAP 11 (2024) 045

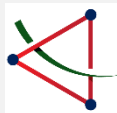


Fuzzy axion (DM) particles



$$\Delta\omega = \frac{11}{1800} \alpha^2 m_a$$

$$P = -\frac{dE}{dt} = \frac{dN_5(t)}{dt} \Delta\omega$$



超轻轴子云的引力波

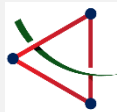
随机引力波背景

$$\Omega_{\text{gw}}(f) = \frac{f}{\rho_c} \int dM dz \frac{dt}{dz} \frac{d\dot{n}}{dM} \frac{dE_s}{df_s},$$

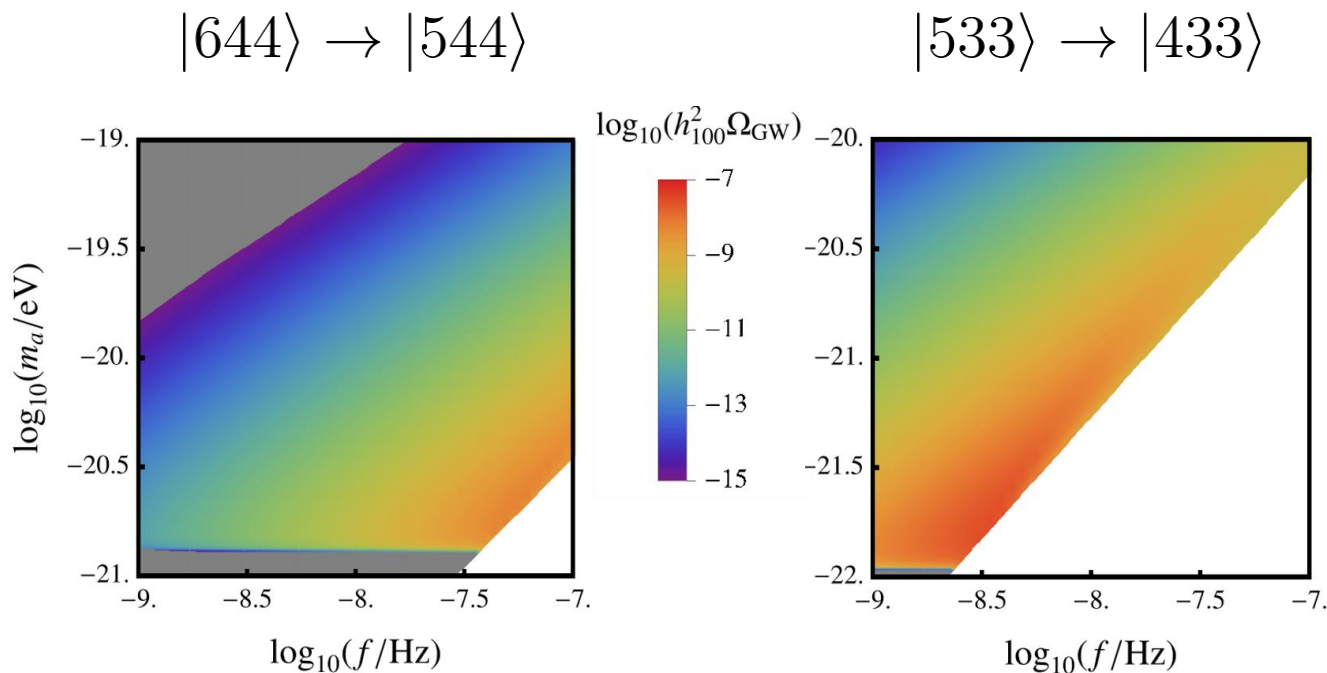
$$dt/dz = \frac{1}{H_0(1+z)\sqrt{\Delta(z)}}$$

$$d\dot{n}/(dM) = (1/t_0)(dn/dM)$$

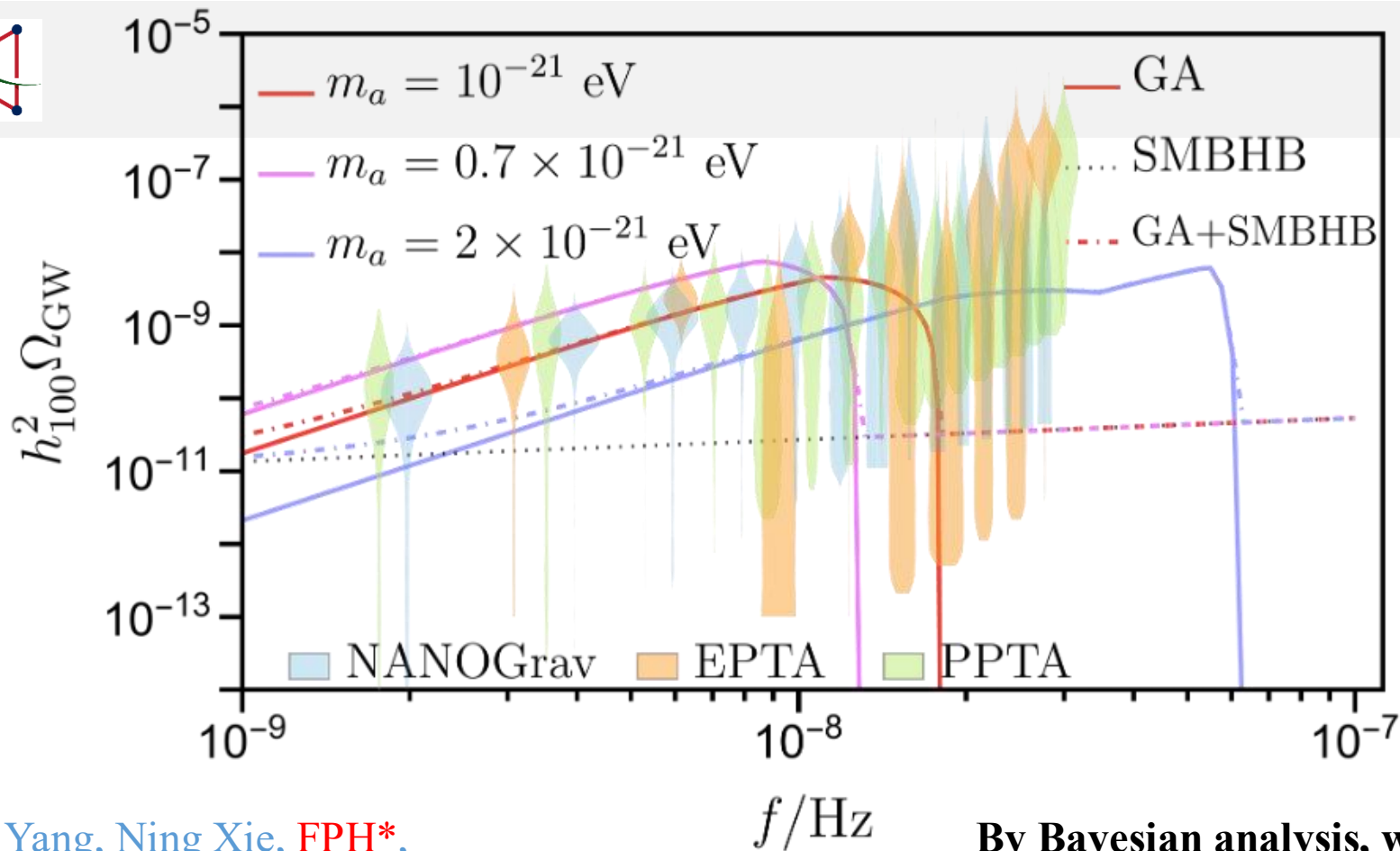
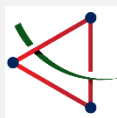
$$\frac{dE_s}{df_s} \approx E_{\text{GW}} \delta(f(1+z) - f_s)$$



随机引力波背景

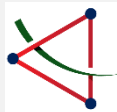


图注：不同超轻轴子质量下随机引力波背景的能量密度



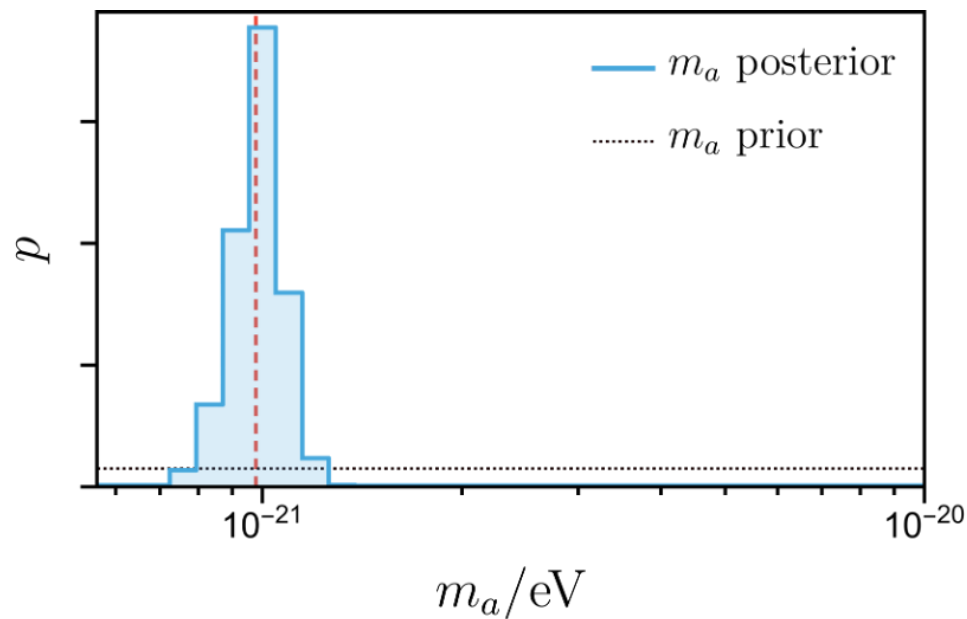
Jing Yang, Ning Xie, FPH*,
JCAP 11 (2024) 045

**By Bayesian analysis, we find
fuzzy DM is favored by the data.**



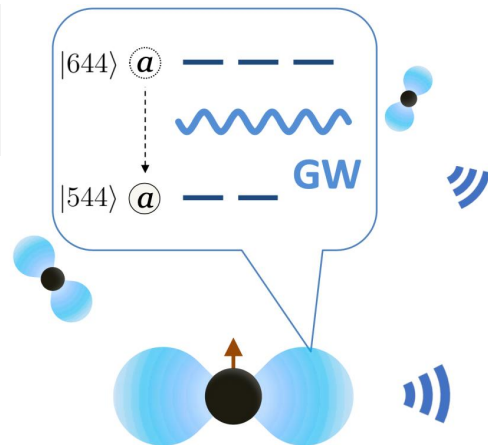
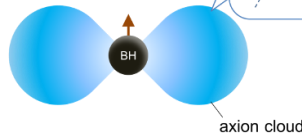
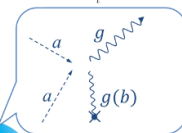
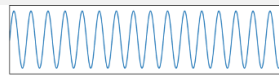
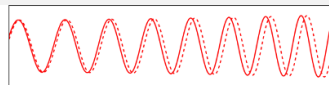
随机引力波背景

贝叶斯分析



贝叶斯因子
(GA+SMBHB/SMBHB) = 7.1

图注：NG15数据给出的超轻轴子质量的后验分布，红色虚线标出了最大概率的质量大小



10⁻²² eV

QCD axion
classic window
10⁻⁶ - 10⁻⁴ eV

The diagram shows a horizontal energy scale. On the left, a vertical tick mark is labeled 10⁻²² eV. To its right is a light red horizontal bar. Within this bar, a darker red vertical segment is labeled 'QCD axion classic window' with '10⁻⁶ - 10⁻⁴ eV' below it. An arrow points to the left from the 10⁻²² eV mark.

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“Ultralight” DM