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Jet modifications in nuclear collisions

- Partons interact with quark-gluon plasma in AA.
- High-*p*_T hardons suppression are governed by parton energy loss. While modifications of jets (as collection of particles) are more complicated.
- This work: employ the LIDO partonic transport model + hydrodynamic description of QGP + simple medium response model to study hadron, jets, and EEC at LHC and RHIC.





Introduction of the LIDO partonic transport model

Nuclear modification factor of hadrons and jets

Substructure: Energy-Energy Correlator

Summary





LIDO Model: Jet Evolution Framework



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Transport Equation in the Incoherent Limit

Combining all relevant processes, the semiclassical transport equation in the incoherent limit with independent collisions can be written as:

$$\frac{df}{dt} = D[f] + C_{1\to 2}[f] + C_{2\to 2}[f] + C_{2\to 3}[f]$$

The distribution function of a hard parton evolves under the influence of:

- Diffusion term: D[f]
- Large-q elastic collisions: $C_{2\rightarrow 2}[f]$
- Diffusion-induced small-q parton splitting/merging: $C_{1\rightarrow 2}[f]$
- Large-q inelastic collisions: $C_{2\rightarrow 3}[f]$

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Including interference from multiple scattering (the LPM effect)

- Energetic splitting has long formation time $\tau_f \sim E/k_{\perp}^2$. All scatterings with medium wihtin τ_f contribute coherently.
- From theory calculation: the rate of medium-induced splittings is suppressed relative to the incoherent limit:

$$\frac{dP_{i\to jk}}{dtdx} \approx \langle \frac{\lambda}{\tau_f} \rangle \frac{dP_{i\to jk}^{\rm incoh}}{dtdx}$$

 $\langle \frac{\lambda}{\tau_{\ell}} \rangle$ is estimate for each parton at run time.

• LIDO is fine tuned to reproduce NLL calculation of medium-induced radiation rates.



Reference(theory):PRC82(2010)064902 Reference(theory):PRC58(1998)1706



Temperature-Dependent Running Coupling

 At finite temperature, we stop the running of α_s below the thermal scale:

 $\mu_{\min} = C_M \times \pi T$

• The $N_f = 3$ strong coupling constant is

$$\alpha_s(Q, T) = \frac{4\pi}{9} \cdot \frac{1}{\ln\left(\frac{\max\left\{Q^2, \mu_{\min}^2\right\}}{\Lambda_{\text{QCD}}^2}\right)}$$





A simple model for medium response

• Energy-momentum deposition to soft sector:

$$rac{d\delta p^{\mu}}{dt}(t,\mathbf{x}) = \int_{p} \Theta(p \cdot u < E_{\min}) p^{\mu} rac{d}{dt} f_{H}(t,\mathbf{x},p)$$

• An ideal-hydro response (no transverse flow)

$$\frac{de}{d\Omega'_k} = \frac{\delta p^0 + \hat{k}' \cdot \delta \vec{p}/c_s}{4\pi}, \quad \frac{d\vec{p}}{d\Omega'_k} = \frac{3(c_s \delta p^0 + \hat{k}' \cdot \delta \vec{p})\hat{k}'}{4\pi}$$

Requires $R_{\text{response}} \gg l_{\text{energy loss}}$.

• Freeze-out to massless particles under a radial transverse flow $v_{\perp} \Rightarrow$ corrects the momentum density in $\eta - \phi$ plane.

$$\frac{d\Delta p_T}{d\phi d\eta} = \int \frac{3}{4\pi} \frac{\frac{4}{3}\sigma u_\mu - \hat{p}_\mu}{\sigma^4} \delta p^\mu(\hat{k}) \frac{d\Omega_{\hat{k}}}{4\pi}$$
$$\sigma = \gamma_\perp [\cosh(\eta - \eta_s - \eta_{\hat{k}}) - v_\perp \cos(\phi - \phi_{\hat{k}})]$$





Jet Definition in LIDO with medium response

• Jets (anti- k_T) are reconstructed from energy bins $E_{T,ij}$, defined by:



• Uncorrelated medium background are assumed to be perfectly subtracted.





Effect of Q_0 (scale for transport initialization) on hadrons R_{AA}



Increasing Q_0 leads to higher R_{AA}



Effect of Q_0 (scale for transport initialization) on jet R_{AA}



Increasing Q_0 leads to higher R_{AA}

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Effect of α_s on Hadrons R_{AA}





Sensitivity to effective jet-medium coupling





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Substructure: Energy-Energy Correlator

$$\frac{\mathrm{d}\Sigma}{\mathrm{d}\theta} = \int \mathrm{d}\vec{n}_1 \,\mathrm{d}\vec{n}_2 \,\frac{\langle\epsilon(\vec{n}_1)\epsilon(\vec{n}_2)\rangle}{Q^2} \,\delta^2(\vec{n}_1 \cdot \vec{n}_2 - \cos(\theta))$$
$$\mathsf{EEC}(\Delta r) = \frac{1}{W_{\mathsf{pairs}}} \frac{1}{\delta r} \sum_{\mathsf{jets} \in [p_{T,1}, p_{T,2}]} \sum_{\mathsf{pairs} \in [\Delta r_a, \Delta r_b]} (p_{T,i}p_{T,j})^n$$

- () Normalize with weighted number of pairs W_{pairs}
- **2** Bin width normalization: $\delta r = \Delta r_b \Delta r_a$
- **③** Exponent values n = 1 used in this analysis



Energy-Energy Correlator Pb+Pb

The medium response contribution $d\Delta p_T/d\phi/d\eta$ has not yet included in the current EEC result (may change large angle behavior)!





Energy-Energy Correlator *p*+Pb





Summary

- Investigated jet and hadron nuclear modification factor R_{AA} in heavy-ion collisions at the LHC using the partonic transport model **LIDO**.
- Highlighted the importance of the correlation between **virtuality evolution** and the **space-time development** of in-medium jet showers.
- Analyzed the **Energy-Energy Correlator (EEC)** in PbPb and pPb collisions, and compared results with **CMS** experimental data.
- **Outlook:** Medium response effects need to be included explicitly for accurate EEC calculations.





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Backup

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Impact of Q_0 on Nuclear Modification Factor R_{AA}



- Jet R_{AA} : Multi-particle jets retain soft radiation inside the cone. Less sensitive to Q_0 .
- Hadron R_{AA} : Single-particle energy loss is directly linked to Q_0 . More sensitive.





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Introduction for LIDO model

We categorize interactions into elastic (hard particle number conserving) and inelastic collisions (hard particle number non-conserving). The inelastic processes are further divided into parton-splitting and parton-fusion contributions.

Elastic collision • $q_{\perp} < Q_{cut}$

 $D[f] = (\eta_{DP} + rac{\hat{q}_s}{4}rac{\partial^2}{\partial p^2})f$

• $q_{\perp} > Q_{cut}$

$$C_{2\leftrightarrow2}^{a}\left(p_{a}\right) = \sum_{bcde}\nu_{c}\int_{q_{\perp}>Q_{cut}}d\Pi_{p_{d}p_{e}}^{p_{b}p_{c}}|M|_{p_{b}p_{c};p_{d}p_{e}}^{2}f_{b}\left(p_{b}\right)f_{0}\left(p_{c}\right)\left(2\pi\right)^{3}\left[-\delta^{(3)}\left(p_{a}-p_{b}\right)\delta^{ab}+\sum_{i=d,e}\delta^{(3)}\left(p_{a}-p_{i}\right)\delta^{ai}\right]$$

 $Q_{cut} = min(\{2m_D, \sqrt{6ET}\})$



Introduction for LIDO model

Incoherent inelastic processes are divided into small-Q diffusion induced radiation/absorption $(1 \leftrightarrow 2)$, and large-Q $2 \leftrightarrow 3$ processes.

Inelastic collision $(1 \leftrightarrow 2)$

 $|q| < Q_{cut}$

$$\mathcal{C}_{1\leftrightarrow2}^{a}(p) = \sum_{bcd} \int dx d\mathbf{k}_{\perp}^{2} \frac{\alpha_{s}(k_{\perp}) P_{cd}^{b}(x)}{(2\pi)(\mathbf{k}_{\perp}^{2} + m_{\infty}^{2})^{2}} \left[-\hat{q}_{g,s}(p) f_{b}(p) \delta^{ab} + \frac{\hat{q}_{g,s}\left(\frac{p}{1-x}\right) f_{b}\left(\frac{p}{1-x}\right)}{1-x} \delta^{ac} + \frac{\hat{q}_{g,s}\left(\frac{p}{x}\right) f_{b}\left(\frac{p}{x}\right)}{x} \delta^{ad} \right]$$

Inelastic collision $(2 \leftrightarrow 3)$

 $|q| > Q_{cut}$

$$\begin{split} C^{a}_{2\leftrightarrow3}(p_{a}) &= \sum_{bcdef} \nu_{c} \int_{q_{\perp} > Q_{cut}} d\Pi^{p_{b}p_{c}}_{p_{d}p_{e}p_{f}} |M|^{2}_{p_{b}p_{c}, p_{d}p_{e}p_{f}} f_{b}(p_{b}) f_{0}(p_{c})(2\pi)^{3} \left[-\delta^{(3)}(p_{a} - p_{b})\delta^{ab} + \sum_{f=d, e, f} \delta^{(3)}(p_{a} - p_{i})\delta^{ai} \right] \\ &+ \left[\text{Absorption term} \right] \end{split}$$