

Studying the $K\Lambda$ strong interaction using femtoscopic correlation functions

Speaker: 刘思威

Cooperator: 谢聚军 研究员

arXiv: 2503.22453

2025年郑州大学强子物理与有效场论前沿讲习班

郑州, 2025年8月22日



CONTENTS

1. *Background*

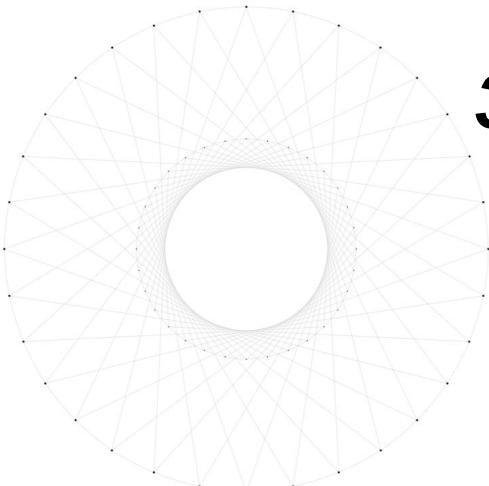
Hadronic Interaction

Progress of correlation function

2. *Methods*

Correlation function

Chiral unitary approach

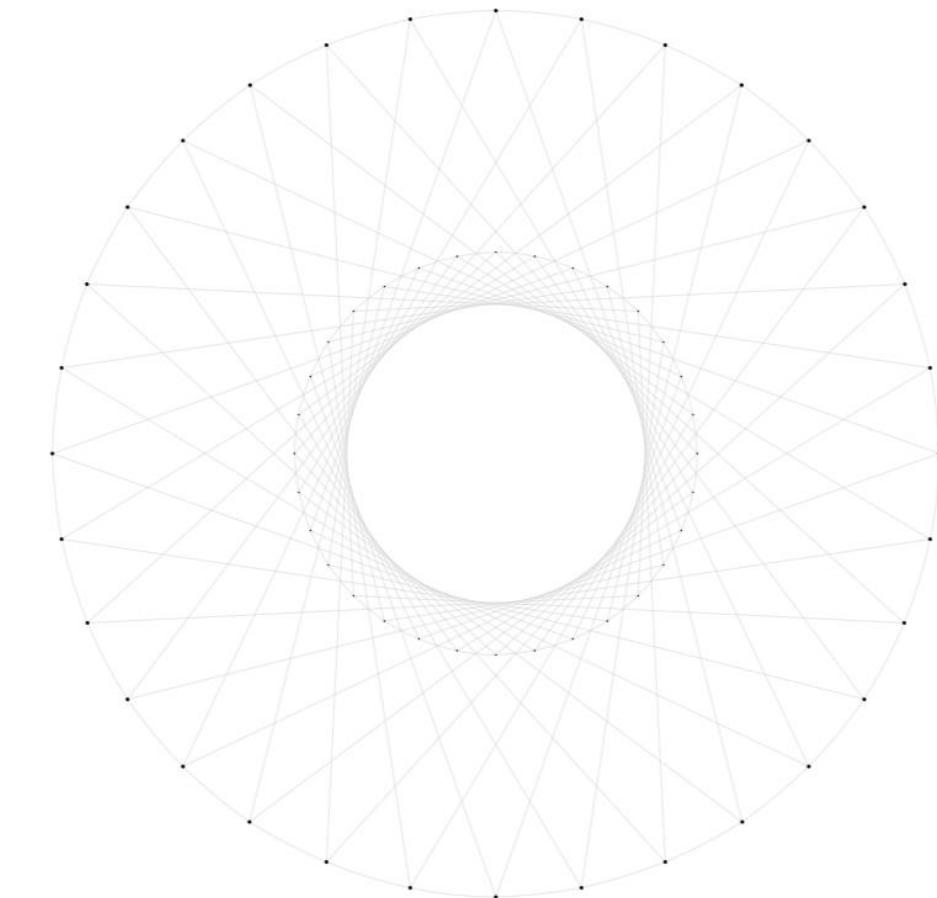


3. *Results*

Coupled channel effects

Correlation function and properties of $N^*(1535)$

4. *Summary*



Background: Hadronic Interactions

Quantum Chromodynamics (QCD)

- Non-perturbation in low-energy region
- Quark confinement
-

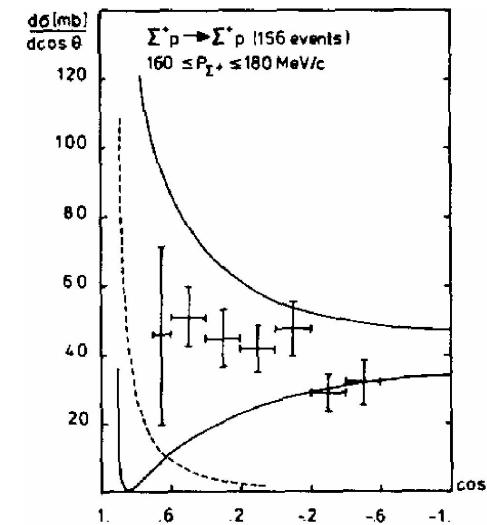
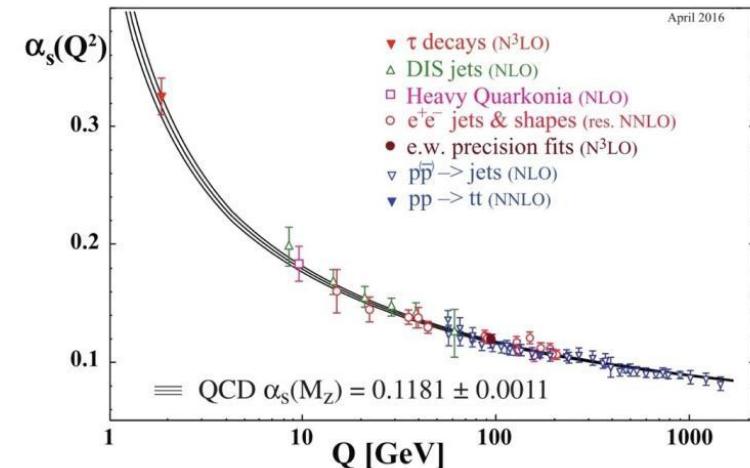
Hadronic Interaction

- Using Lattice QCD and Effective Field Theory
- Hadron molecule

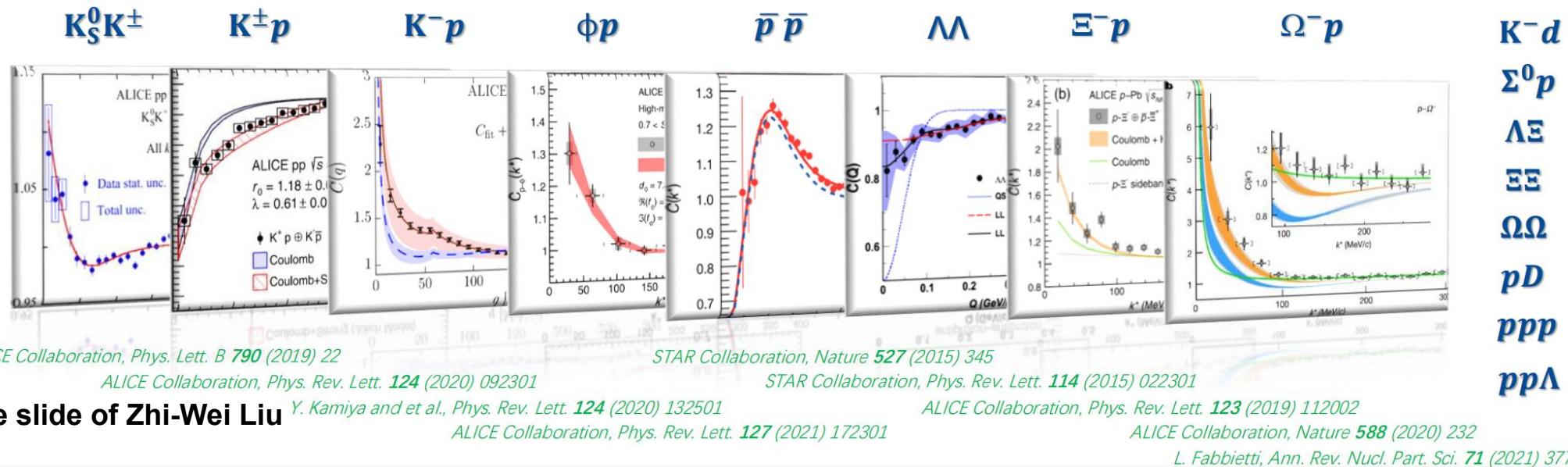
Scattering Experiment

- Lots of hadrons are short-lived
- Scarce and low precision data about hyperons

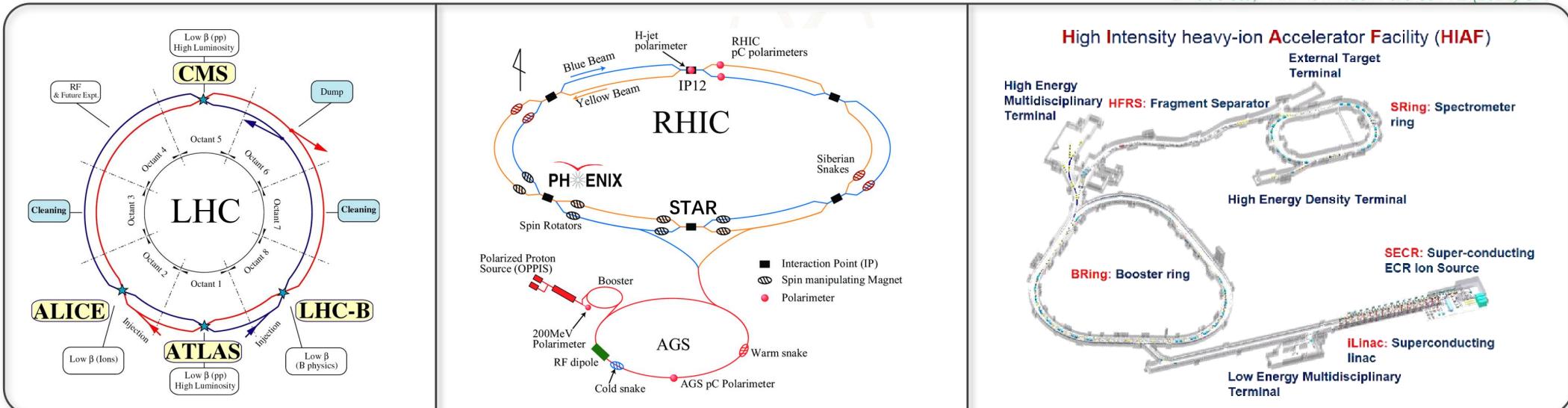
New approach is accquired !!!



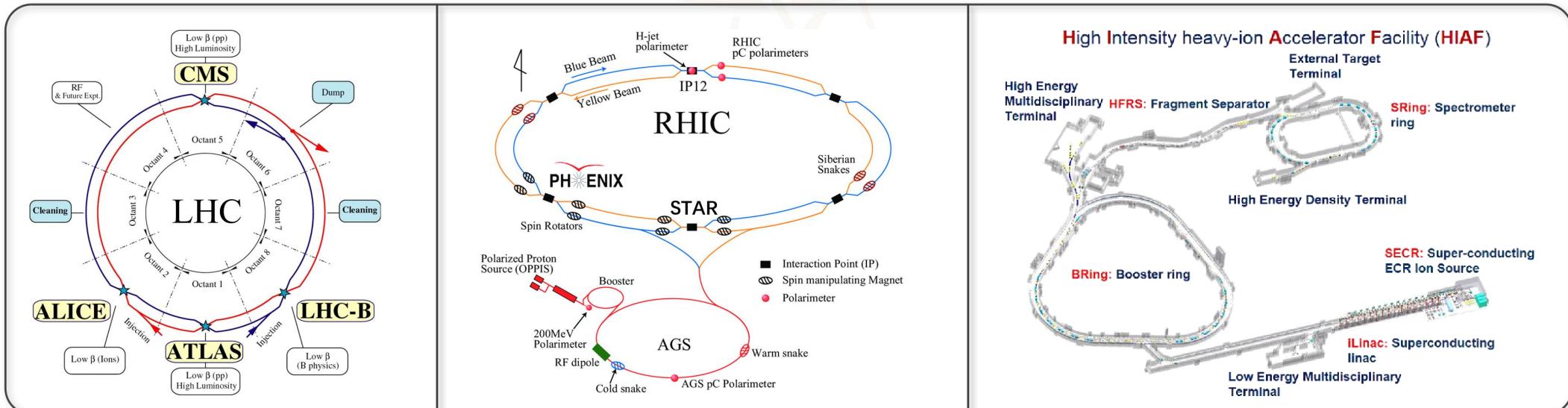
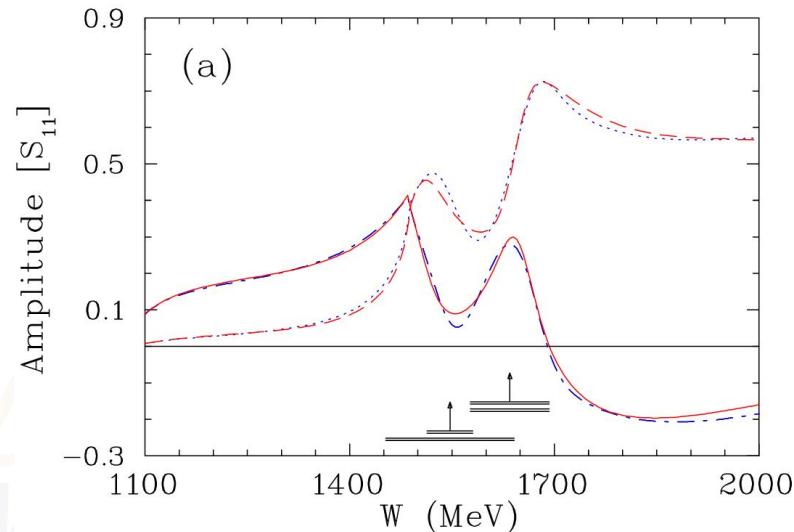
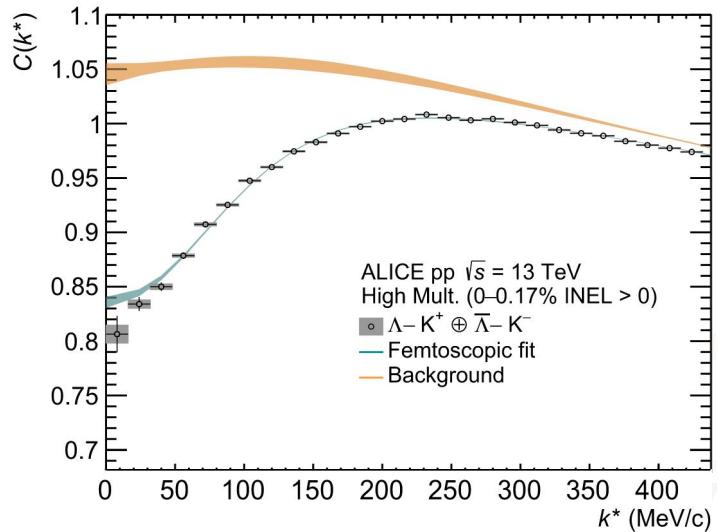
Background: Progress of correlation function



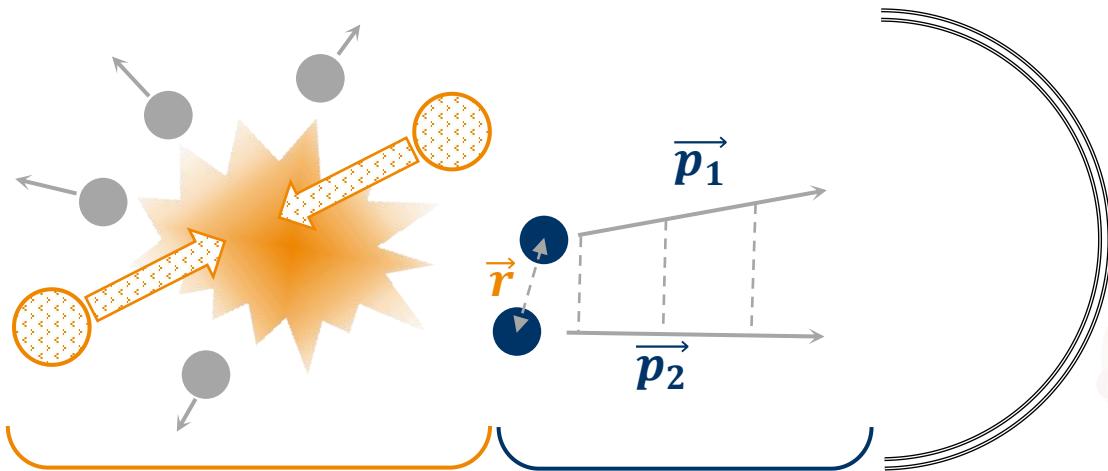
From the slide of Zhi-Wei Liu



Background: Progress of correlation function



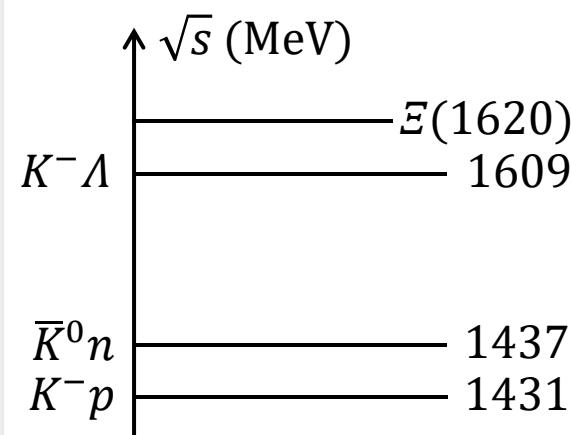
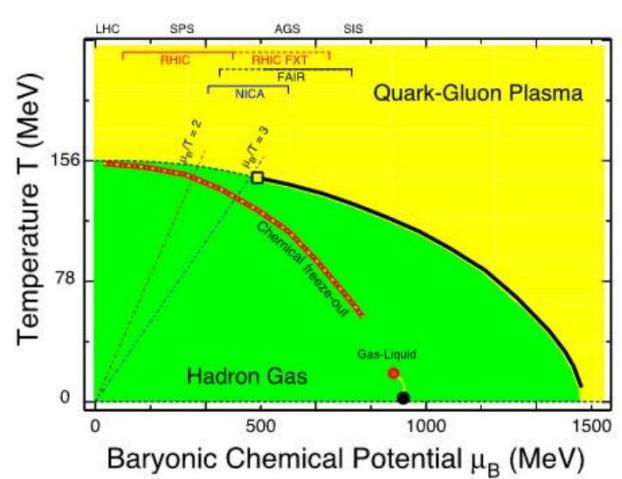
Methods: Correlation function



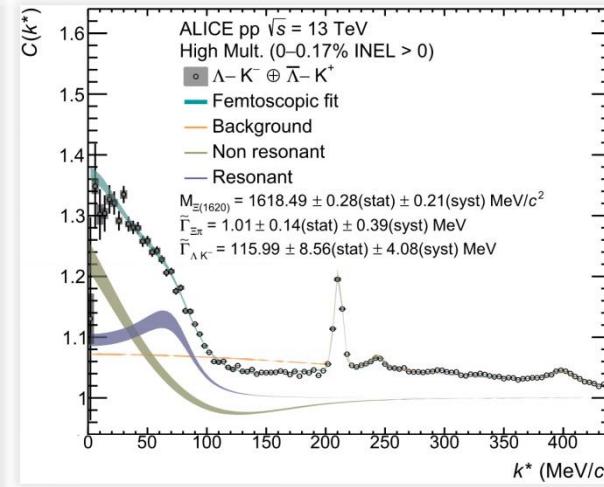
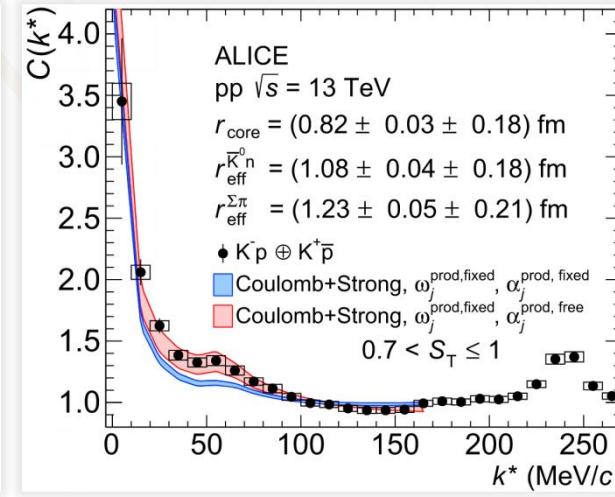
Correlation function
joint probability density

$$C(\vec{p}) = \frac{P(\vec{p}_1, \vec{p}_2)}{P(\vec{p}_1) \cdot P(\vec{p}_2)} \begin{cases} = 1, & \text{no interaction} \\ > 1, & \text{attractive} \\ < 1, & \text{repulsive} \end{cases}$$

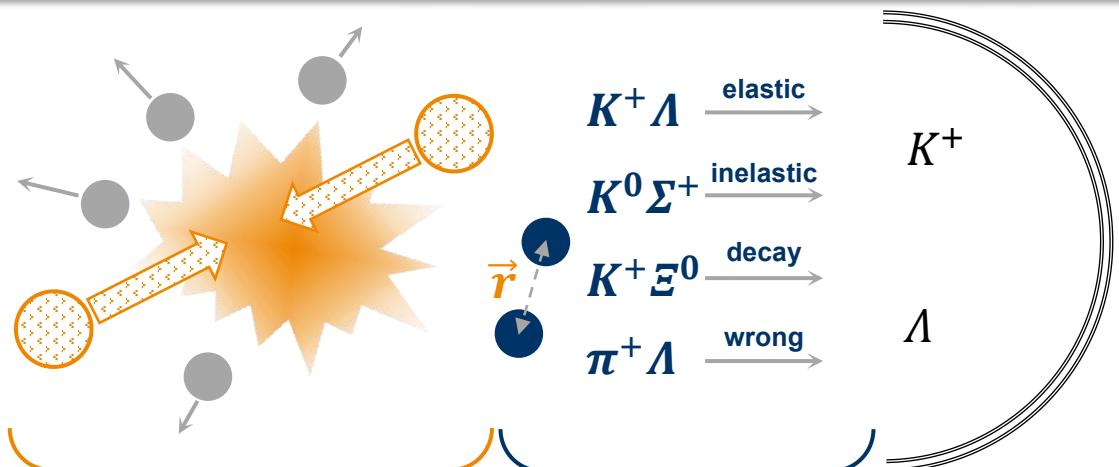
QCD phase transition



Advantages



Methods: Correlation function



Emitting Source
 $S(\vec{r})$

Wave Function
 $\Psi(\vec{r}, \vec{p})$

Detector
 $\xi(\vec{p})$

Emitting Source

$$S(\vec{r}) = \frac{1}{(4\pi R^2)^{3/2}} e^{-\frac{|\vec{r}|^2}{4R^2}}$$

long-lived resonances

$$S_{eff}(\vec{r}) = \omega_S S(\vec{r}, R_1) + (1 - \omega_S) S(\vec{r}, R_2)$$

(ALICE Collaboration), Phys.Lett.B 811 (2020) 135849
(ALICE Collaboration), Phys.Lett.B 845 (2023) 138245

$$C(\vec{p}) = \int d^3\vec{r} S(\vec{r}) |\Psi(\vec{r})|^2$$

Correlation Analysis Tool using the Schrödinger equation (CATS) Formalism

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V(\vec{r}) \right] \psi(\vec{r}) = E\psi(\vec{r})$$

Lednický–Lyuboshits (LL) Formalism

$$\Psi(\vec{r}) = e^{-ikr\cos\theta} + f(k) \frac{e^{-ikr}}{r}, \quad f(k) = \frac{1}{a_0^{-1} + 0.5r_0 k^2 - ik}$$

KooninPratt (KP) Formalism

$$T = V + VGT, \quad \Psi = \phi + GT\phi$$

Methods: Chiral unitary approach

Chiral lagrangian

$$\mathcal{L}_1^{(B)} = \frac{1}{4f^2} \langle \bar{B} i\gamma^\mu [[\phi, \partial_\mu \phi], B] \rangle$$

$$V_{ij} = -\frac{C_{ij}}{4f_i f_j} (k_i + k_j) \sqrt{\frac{(M_i + E_i)(M_j + E_j)}{4M_i M_j}}$$

Lippmann-Schwinger equation

$$T = V + VGT, \quad \Psi = \phi + GT\phi$$

$$G = \int_0^\infty \frac{d^4 q}{(2\pi)^4} \frac{1}{(q^2 - m_j^2 + i\epsilon)((P - q)^2 - m_{j2}^2 + i\epsilon)}$$

Correlation function in coupled channel

$$\begin{aligned} \Psi_{ij}(\vec{r}, \vec{p}) &= \delta_{ij} e^{i\vec{p}\vec{r}} + T_{ij}(\vec{r}, \vec{p}) \int_0^\infty \frac{d^4 q}{(2\pi)^4} \frac{e^{i\vec{p}\vec{r}}}{(q^2 - m_j^2 + i\epsilon)((P - q)^2 - m_{j2}^2 + i\epsilon)} \\ &= \delta_{ij} e^{i\vec{p}\vec{r}} + T_{ij}(\vec{r}, \vec{p}) \widetilde{G}_j(\vec{r}, \vec{p}) \end{aligned}$$

$$C_{theory}(\vec{p}) = 1 + \int_0^\infty d^3 \vec{r} S(\vec{r}) \left(2\text{Re}[T_{ii}(\vec{r}, \vec{p}) \widetilde{G}_i(\vec{r}, \vec{p}) j_0(|\vec{p}| |\vec{r}|)] + \sum_j |T_{ij}(\vec{r}, \vec{p}) \widetilde{G}_j(\vec{r}, \vec{p})|^2 \right)$$

Methods: Properties

Correlation function

$$C_{theory}(\vec{p}) = 1 + \int d^3\vec{r} S(\vec{r}) \left(2\text{Re}[T_{ii}(\vec{r}, \vec{p}) \tilde{G}_i(\vec{r}, \vec{p}) j_0(|\vec{p}| |\vec{r}|)] + \sum_j |T_{ij}(\vec{r}, \vec{p}) \tilde{G}_j(\vec{r}, \vec{p})|^2 \right)$$

$$C_{fit} = \lambda C_{theory} + (1 - \lambda) \times C_{background}$$

R. Molina et al., Eur. Phys. J. C 84 (2024) 328

Observables

$$T = -\frac{8\pi\sqrt{s}}{2M} \frac{1}{a_0^{-1} + 0.5r_0k^2 - ik} \rightarrow a_{0,j} = \frac{2M_j T_{jj}}{8\pi\sqrt{s}} \Big|_{s=s_{th,j}}, \quad r_{0,j} = 2 \frac{\partial}{\partial k^2} \left[-\frac{8\pi\sqrt{s}}{2M_j T_{jj}} + ik_j \right] \Big|_{s=s_{th,j}}$$

R. Molina et al., Phys. Rev. D 109 (2024) 054002

Properties of $N^*(1535)$

$$\text{pole} \quad \text{coupling} \quad \text{decay width}$$

$$\text{Det}[1 - VG]|_{s=s_{\text{pole}}} = 0 \rightarrow T = \frac{g_i g_j}{\sqrt{s} - \sqrt{s_{\text{pole}}}} \rightarrow \Gamma_j = \frac{|g_j|^2 p_j (E_j + M_j)}{8\pi M_{N^*(1535)}}$$

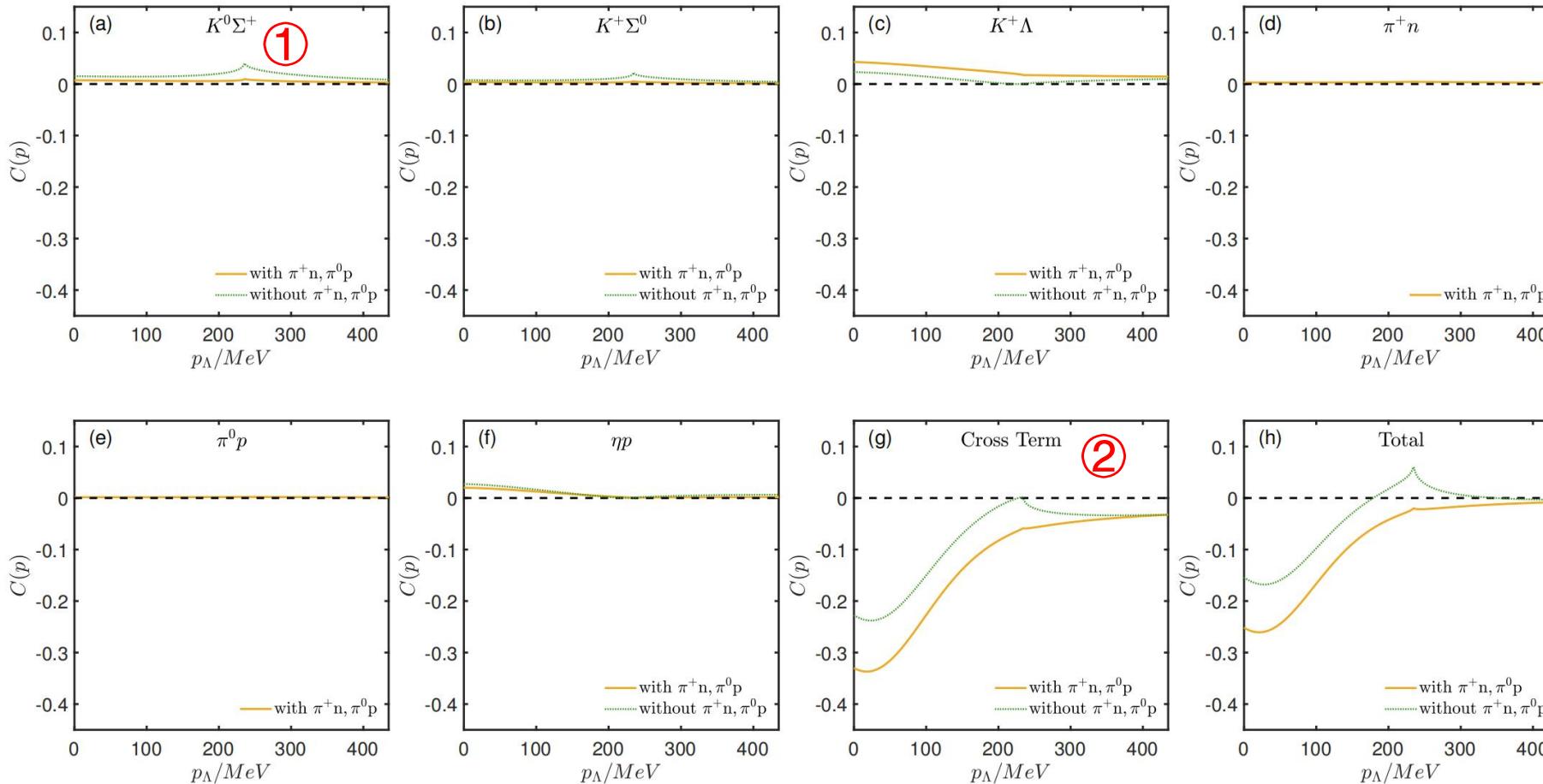
Results: Coupled channel effects

$$C_{theory}(\vec{p}) = 1 + \int_0^\infty d^3\vec{r} S(\vec{r}) \left(2\text{Re}[T_{ii}(\vec{r}, \vec{p}) \tilde{G}_i(\vec{r}, \vec{p}) j_0(|\vec{p}| |\vec{r}|)] + \sum_j |T_{ij}(\vec{r}, \vec{p}) \tilde{G}_j(\vec{r}, \vec{p})|^2 \right)$$



Results: Coupled channel effects

$$C_{theory}(\vec{p}) = 1 + \int_0^\infty d^3\vec{r} S(\vec{r}) \left(2\text{Re}[T_{ii}(\vec{r}, \vec{p}) \tilde{G}_i(\vec{r}, \vec{p}) j_0(|\vec{p}| |\vec{r}|)] + \sum_j |T_{ij}(\vec{r}, \vec{p}) \tilde{G}_j(\vec{r}, \vec{p})|^2 \right)$$



①
a peak associated
higher energy channel

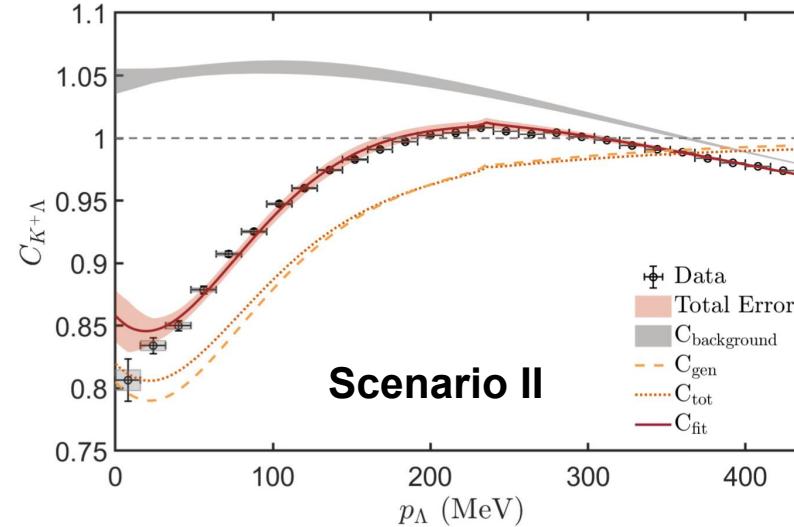
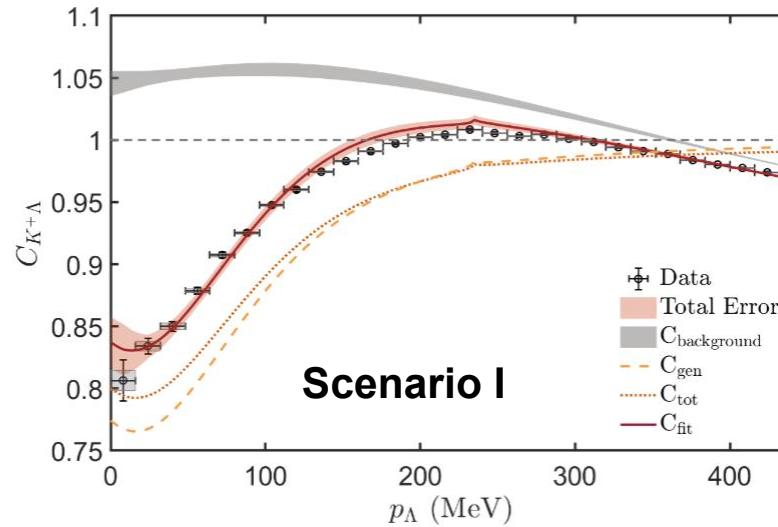
②
the main interaction introduced
by the cross term
 $\text{Re}[T_{ii}\tilde{G}_i j_0(|\vec{p}| |\vec{r}|)]$

Results: Fitting results

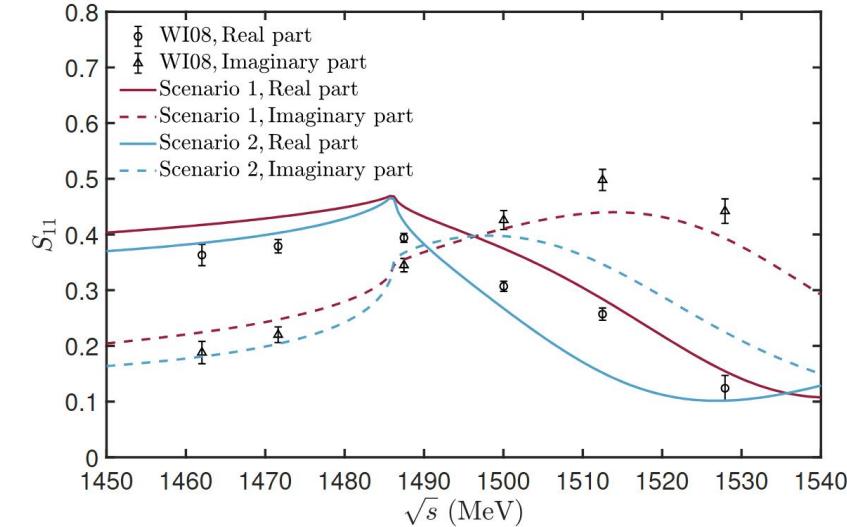
Scenario I: α_i are fixed as in [1],
 N and λ are free parameters

Scenario II: α_i , N , and λ are free parameters

[1]. T. Inoue et al., Phys.Rev.C 65 (2002) 035204



Scenario	I	II
N	0.99554 ± 0.00045	0.99633 ± 0.00032
λ	0.87 ± 0.01	0.91 ± 0.02
$\alpha_{K\Sigma}$	-2.8	-3.17 ± 0.06
$\alpha_{K\Lambda}$	1.6	1.86 ± 0.06
$\alpha_{\eta N}$	0.2	0.76 ± 0.09
$\alpha_{\pi N}$	2.0	1.91 ± 0.04
$\chi^2/\text{d.o.f.}$	1.9	1.2



Results: Fitting results

Scenario	I	II
$\sqrt{s_{\text{pole}}}$ (MeV)	1531 - i 36.5	1512 - i 39.1
Coupling	$g_{K^0\Sigma^+}$	2.21 - i 0.18
	$g_{K^+\Sigma^0}$	1.56 - i 0.13
	$g_{K^+\Lambda}$	1.37 - i 0.10
	g_{π^+n}	0.55 + i 0.32
	g_{π^0p}	0.40 + i 0.22
	$g_{\eta p}$	-1.47 + i 0.43
Partial Width (MeV)	Γ_{π^+n}	19.7
	Γ_{π^0p}	9.9
	$\Gamma_{\eta p}$	41.2
Scattering Length (MeV $^{-1}$)	$a_{K^0\Sigma^+}$	0.28 - i 0.14
	$a_{K^+\Sigma^0}$	0.19 - i 0.13
	$a_{K^+\Lambda}$	0.33 - i 0.14
	a_{π^+n}	-0.08
	a_{π^0p}	0.02
	$a_{\eta p}$	-0.25 - i 0.17
Effective Range (MeV $^{-1}$)	$r_{K^0\Sigma^+}$	1.05 - i 1.55
	$r_{K^+\Sigma^0}$	-1.46 - i 3.69
	$r_{K^+\Lambda}$	-0.52 - i 0.38
	r_{π^+n}	-13.39 - i 18.04
	r_{π^0p}	30.34
	$r_{\eta p}$	-8.60 + i 6.77

1. The mass and width of $N^*(1535)$

	Mass (MeV)	Width (MeV)
PDG (Pole)	1500~1520	80~130
PDG (BW)	1515~1545	125~175

2. The decay widths of $N^*(1535)$

	$\Gamma(N(1535) \rightarrow \eta N)/\Gamma_{tot}$	$\Gamma(N(1535) \rightarrow \pi N)/\Gamma_{tot}$	Γ_1/Γ_2
PDG	0.30~0.55	0.32~0.52	0.99±0.05±0.19
Scenario 1	0.56	0.41	1.4
Scenario 2	0.51	0.42	1.2

3. The scattering length of $K^+\Lambda$

	[1]	Scenario 1	Scenario 2
$a_{K^+\Lambda}$ (MeV $^{-1}$)	0.61±0.03±0.03	0.33-0.14i	0.29-0.17i

Summary

- We successfully calculated the theoretical $K^+\Lambda$ correlation function based on Koonin-Pratt formalism
- From the analysis of coupled channels, two conclusions can be drawn:
 1. High-energy coupled channels will produce **a peak at their thresholds**;
 2. The interference term between $K^+\Lambda$ and plane wave **contributes the main interaction**.
- The correlation function and the scattering experimental data can be **mutually verified in general trends**, but it is difficult to precisely satisfy both sets of data simultaneously at this stage.
- We obtained the mass, width, and couplings of the $N(1535)$, which can well explain the experimental data under the picture of **hadronic molecule**.



中国科学院近代物理所
INSTITUTE OF MODERN PHYSICS

Thanks for your listening!