

# Quark Spaghetti with Glue Balls (SQGB)

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(A simple story that sadly could have been told years ago....)

## **Scientific Question:**

Is there a phase of matter between the Quark  
Gluon Plasma and the Hadron Gas?

Work in part motivated by some  
observations of Lenya Glozman, Tom  
Cohen, and colleagues

## **Motivation from lattice computations:**

Chiral symmetry is restored about 160 MeV, but confinement as measured by Polyakov loop appears to disappear at about 300~MeV. This is the temperature of deconfinement in the pure glue theory

In this intermediate temperature range, entropy is smaller than expected

## Old Problem:

How to characterize confinement:

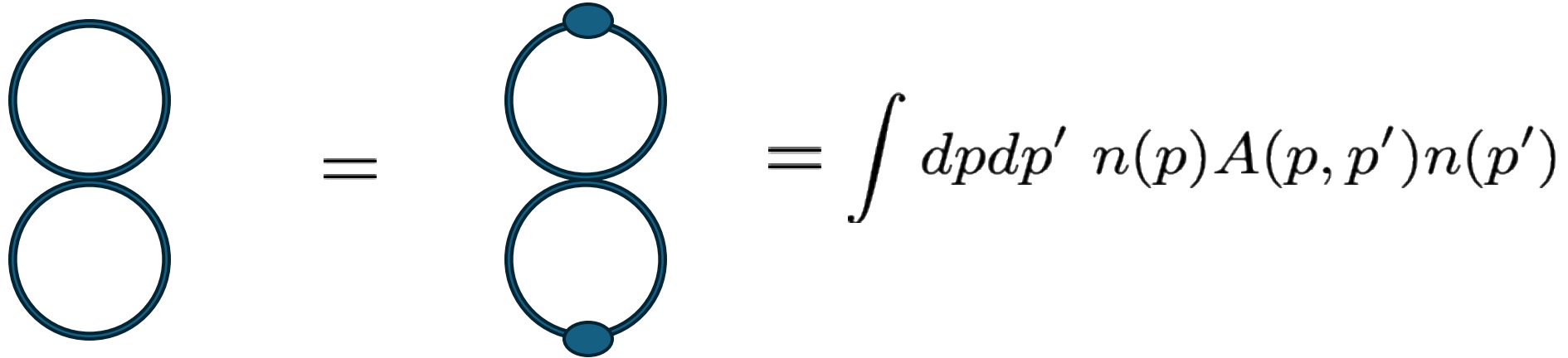
When dynamical quarks are present confinement at finite temperature or density is only approximately characterized by the Polyakov line because strings always break and the free energy of an isolated quark is always finite.

If we go to large  $N_c$ , or in pure gauge theory, there is truly confinement. In pure gauge theory, the deconfinement temperature is about 270 MeV, which is much different from the chiral transition in a theory with dynamical quarks. (The temperature is only 10% different from that in the infinite  $N_c$  limit)

Perhaps the region between 160 MeV and 270 MeV is distinct from that at higher or lower  $T$ . Some evidence from lattice data showing that stringy vortices percolate at around 270 MeV. In this vortex picture, this should be the temperature of deconfinement.

How to characterize these separate regions.

Contribution to the pressure are densities times an amplitude:



$$\begin{array}{c} \bigcirc \\ \bigcirc \end{array} = \begin{array}{c} \bigcirc \\ \bigcirc \end{array} = \int dp dp' n(p) A(p, p') n(p')$$

Kinetic energy of hadrons  $\sim 1$ , quarks  $\sim N_c$ , gluons  $\sim N_c^2$

Amplitude for meson-meson  $\sim 1/N_c$ , glueball-glueball  $\sim 1/N_c^2$

For pressure of hadrons to be of order  $N_c$  the  
hadron density should be of order  $N_c$

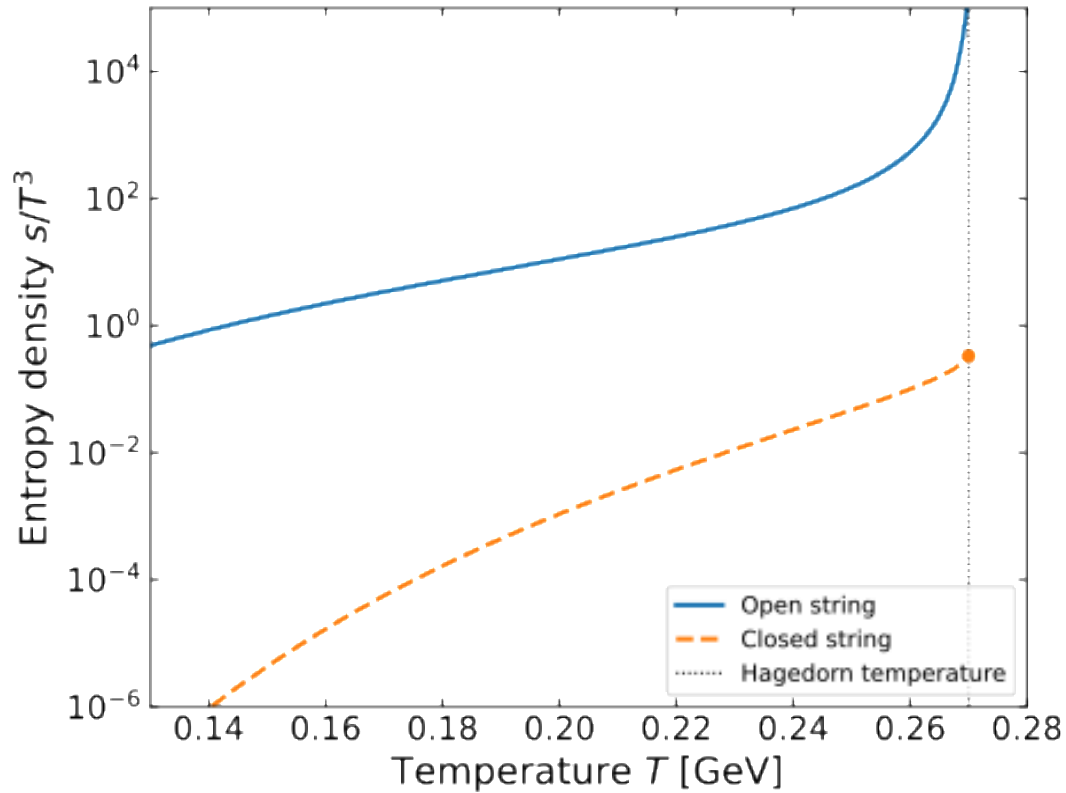
Pressure for glueballs to be of order  $N_c^2$  the  
density of glueballs should be of order  $N_c^2$

Mesons and glueballs, density  $\sim 1$

Quarks and glueballs, density  $\sim N_c$

Quarks and gluons, density  $\sim N_c^2$

## Simple model:



Contribution of glueballs is always very very small.

The closed string, glueball contribution, is non-divergent at the Hagedorn temperature

Mesons: open strings

Glueballs: closed strings

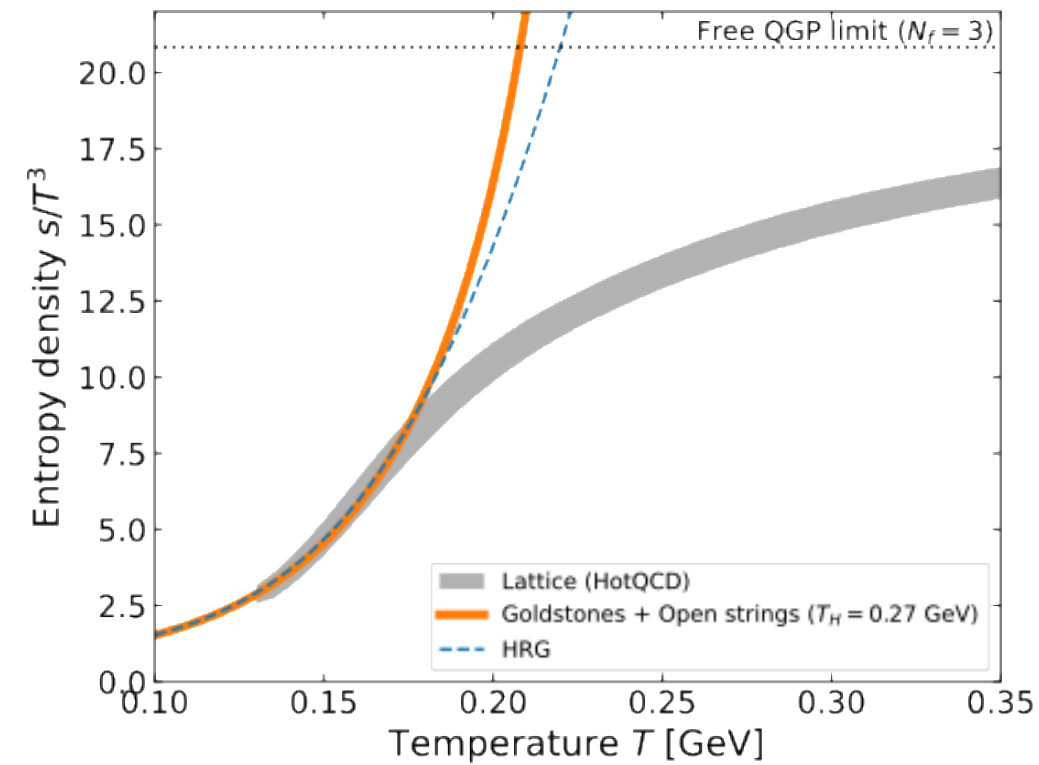
Quarks: Free quarks with mass between constituent mass and current quarks mass

Gluons: Mass between massless and  $\frac{1}{2}$  of glueball mass

Take closed and open string spectrum from string theory. Glueball and mesons must have the same Hagedorn temperature which we take to be

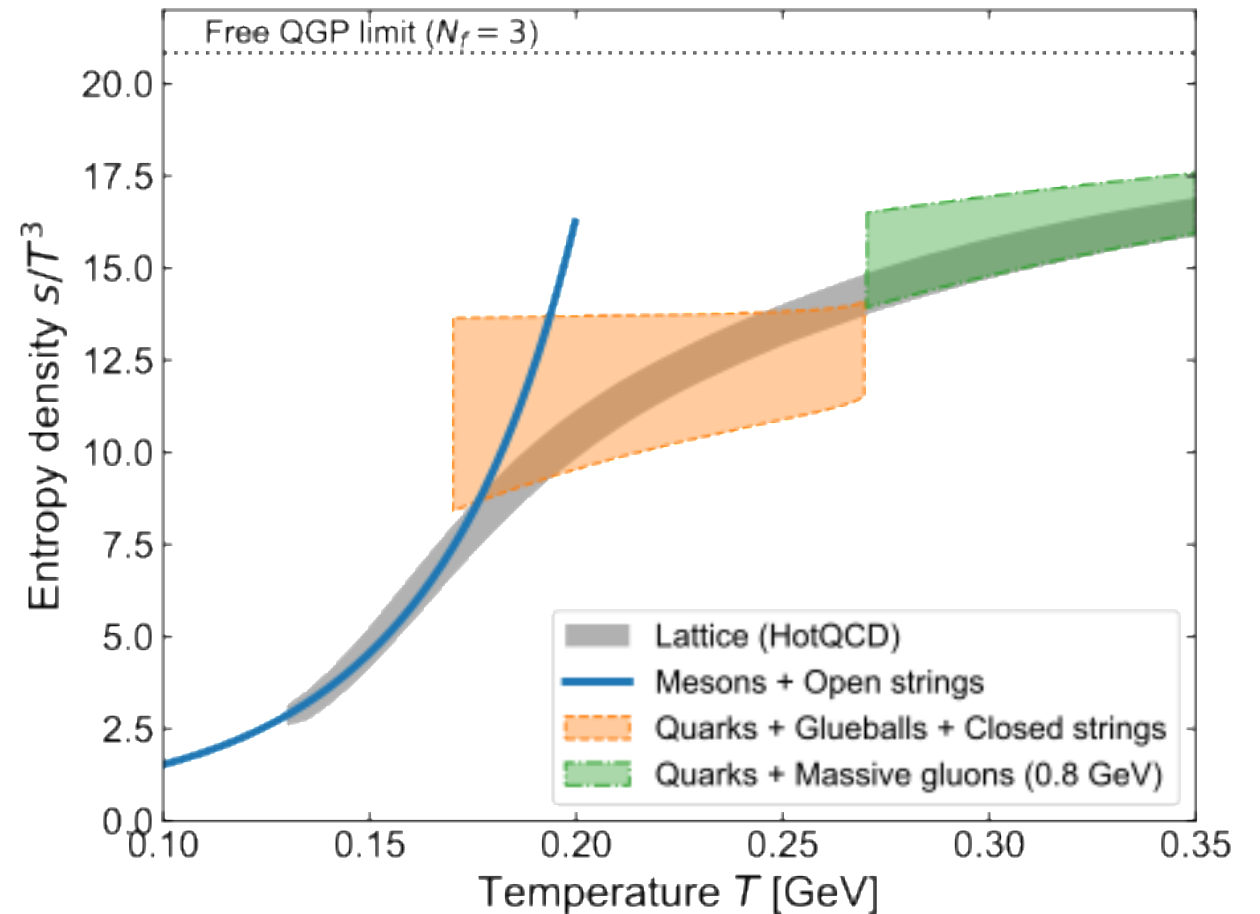
$$T_H = 270 \text{ MeV}$$

We handle the contribution of lowest mass glueball states, and Goldstone boson meson states separately, and then treat higher mass states by a continuum integral. The meson integration begins at 1 GeV, and the glueball at about 3~GeV



Matching on to quarks plus glueball at 160 MeV, and to quarks plus gluons at the Hagedorn temperature, the entropy is reasonably well described

With the string theory resonance model and a Hagedorn temperatures of 270 MeV, describe the lattice data on entropy up until the chiral temperature, where presumably the degrees of freedom become quarks



The quark phase with energy density of order  $N_c$  has to a good approximation, no glueballs, and hence no gluons, in it. This is because glueballs are massive. The Debye screening mass is

$$M_{Debye}^2 = g^2(N_c/3 + N_f/6)T^2$$

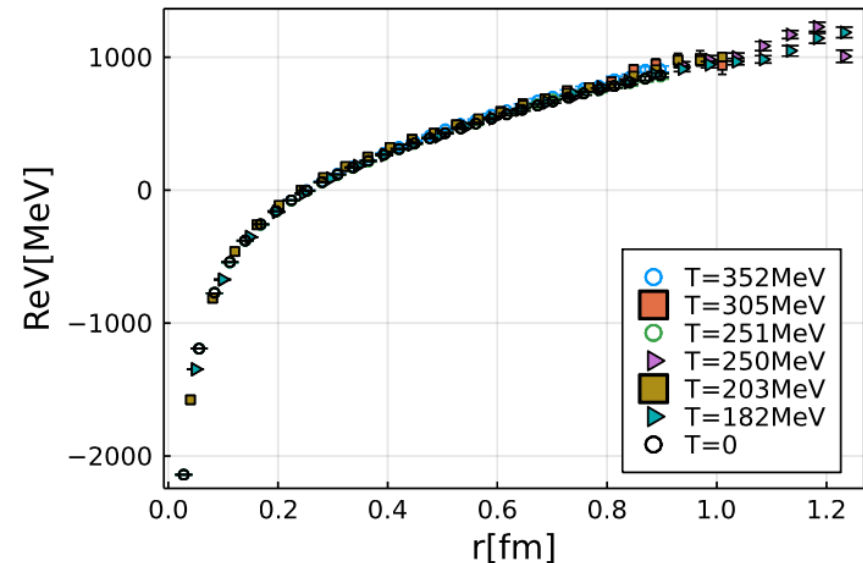
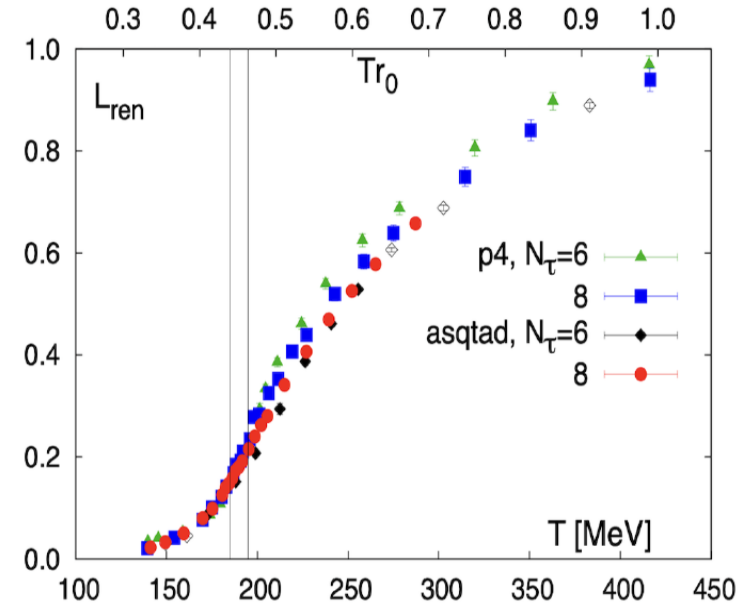
The first term is from gluons and the second from quarks. In the large  $N_c$  limit, one holds the 'tHooft coupling fixed,

$$g_{tHooft}^2 = g^2 N_c$$

If because glueball contribution is small, we ignore the first term, the second term vanishes in the large  $N_c$  limit, so that the new phase of quarks is confined. The quarks are stringy. We call this phase,

### Spaghetti of Quarks with Glueballs (SQGB)

But it is not very expensive **Quark Spaghetti with Glue Balls** since there is much spaghetti and few glueballs!



What are the chiral properties of the SQGB?

The contribution of pions which dilutes the chiral condensate is

$$\frac{\partial P}{\partial m_Q} \sim T^2 / f_\pi^2 \sim \frac{T^2}{N_c \Lambda_{QCD}^2}$$

This vanishes in the large  $N_c$  limit. This is basically the  $1/N_c$  suppression of meson interactions. It is cured by having many mesons so that when the meson density is of order  $N_c$ , the chiral order parameter shrinks to zero. Therefore the transition from hadron matter to the SQGB is associated with a chiral transition.

Not clear how chiral  $U(1)$  is restored. Perhaps this occurs at the higher temperature?

## SQGB and Quarkyonic Matter

Quarkyonic matter is for low temperatures and finite baryon number density. Because there are no gluons the matter is confined. In ordinary hadronic matter there are no baryons for temperature because

$$e^{-M_B/T} \big|_{T \sim \Lambda_{QCD}} \sim e^{-N_c}$$

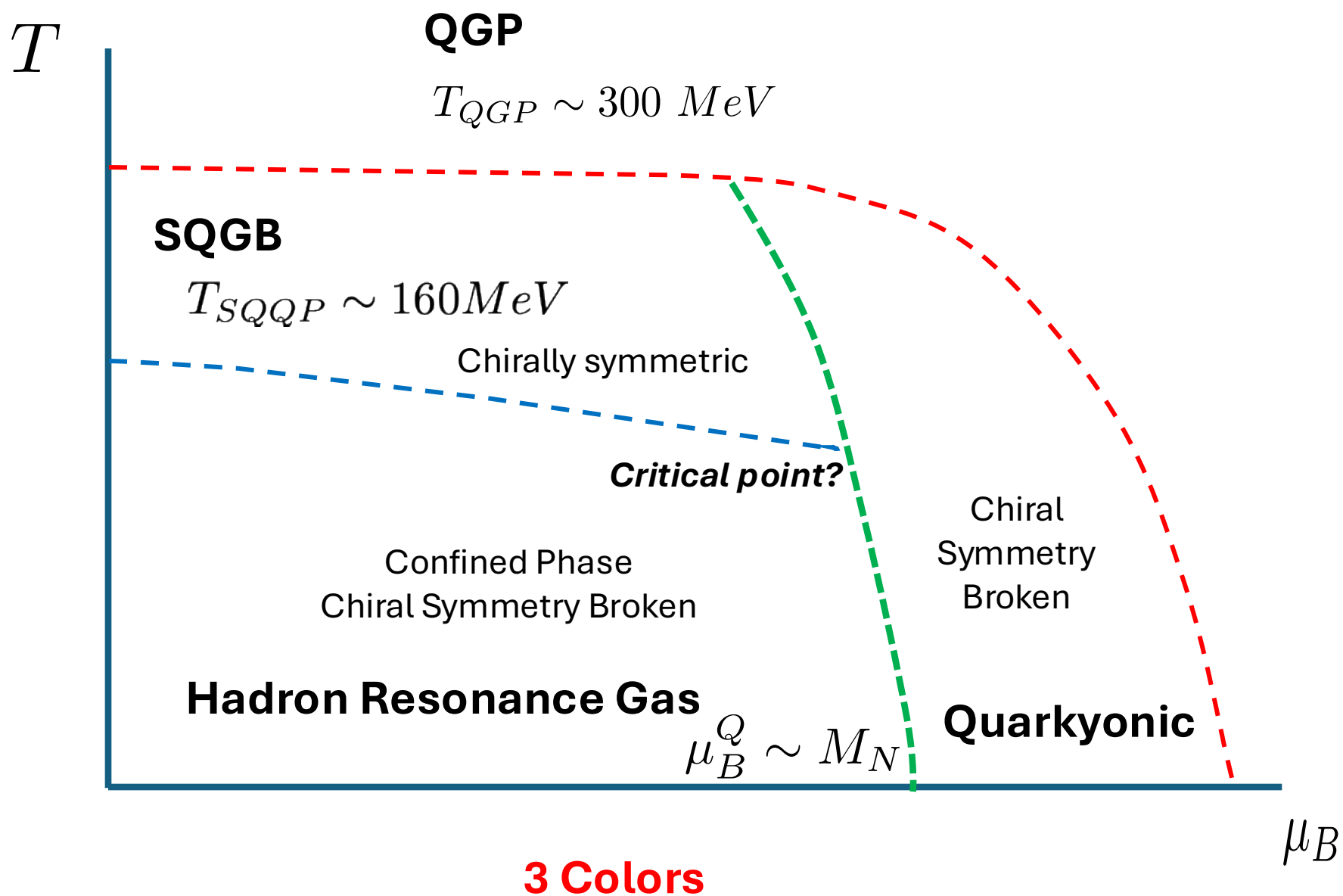
Baryons exist as quark-like degrees of freedom because of a chemical potential

$$\mu_B - M = \kappa T, \kappa \sim 1$$

The density of quarks is of order  $N_c$  in Quarkyonic matter.

This is very similar to the SQGB, except in the SQGP there are antiquarks.

Naively, Quarkyonic Matter has restored chiral symmetry, but condensate forms at the Fermi surface that break chiral symmetry. The weak chiral symmetry breaking effects differentiate the SQGP from Quarkyonic Matter



In the limit if a very large number of colors, because the density of quarks diverges, the QGP and SQGB temperatures should both become the Hagedorn temperature

