Pentagonal Feynman Integrals and Wilson Loop with Lagrangian Insertion at Three-Loop

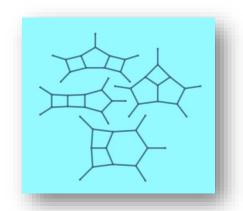
Base on:

Dmitry Chicherin, Yu Wu, Zihao Wu, Yongqun Xu, Shun-Qing Zhang, Yang Zhang [2511.XXXX]

Dmitry Chicherin, Johannes Henn, Yongqun Xu, Shun-Qing Zhang, Yang Zhang [2511.XXXX] (Stay Tune)

and

Yuanche Liu, Antonela Matijašić, Julian Miczajka, Yingxuan Xu, Yongqun Xu, Yang Zhang [2411.18697] (Phys. Rev. D 112, 016021)

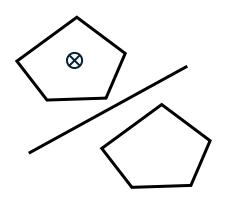


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2025 Oct. 31



Pentagonal Feynman Integrals

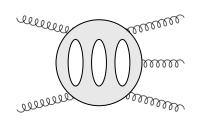
at Three-Loop

Background

Feynman Integrals as the basis building block of perturbative Quantum Field Theory



Collider Physics



 $\mathcal{A}_5^{(3)}(1,2,3,4,5)$

Real world Amplitudes

$$pp \to \gamma \gamma \gamma$$

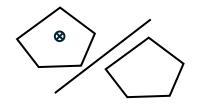
$$pp \rightarrow jjj$$

@NNNLO

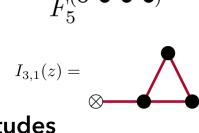
Future Collider: CEPC, FCC, HL-LHC...

Observables in formal theory





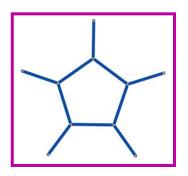




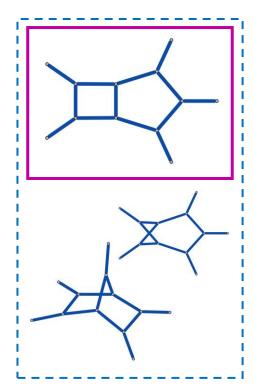
Negative geometries expansion

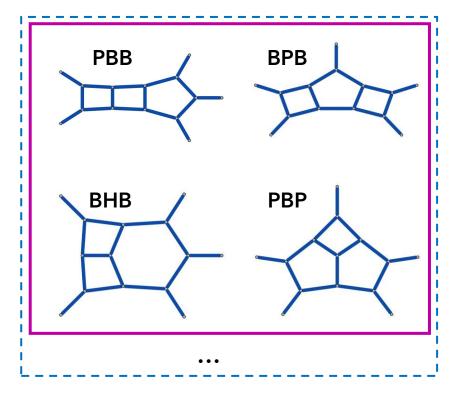
Pentagonal Feynman Integrals at Three-Loop

Planar Part



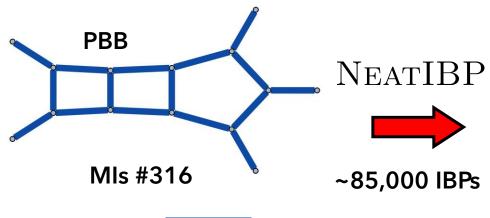
[Bern, Dixon, Kosower ' 92]
[A. Kniehl, V. Tarasov '10]
[Gehrmann, Henn, Lo Presti '18]
[Chicherin, Gehrmann, Henn,
Wasser, Zhang, Zoia '18]
[Abreu, Page, Zeng '18]
[Liu, Matijašić, Miczajka, Xu, Xu,
Zhang '24]
and so on...



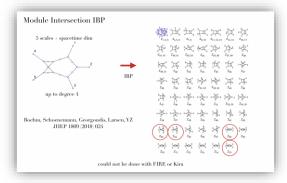


Integration-by-parts (IBP) reduction

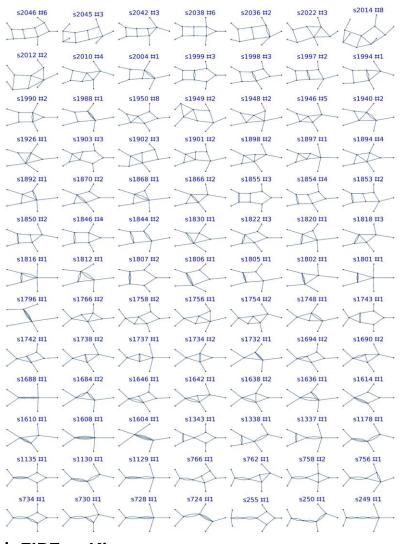
The current bottleneck in perturbative calculation



Picture created by Azurite 1.1.0



Similar Stories at Two Loop Picture from Yang Zhang



Could not be done with FIRE or Kira

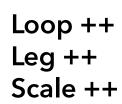
Algebraic Geometry Accelerating IBP

Family	PBB	BPB	BHB	PBP	
Master Integrals	316	367	431	734)
Five Point Integrals	120	164	169	342	Degree 4 and Degree 3 + Dot 1
# IBPs	~85,000	~183,000	Spanning Cuts		

PRD Editor's Suggestion [Liu, Matijašić, Miczajka, Xu, Xu, Zhang '24]



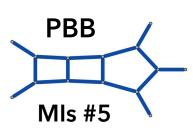
[Wu, Böhm, Ma, Xu, Zhang '23] [Wu, Böhm, Ma, Usovitsch, Xu, Zhang '25]





UT Integrals: Leading Singularities Analysis

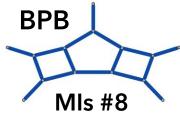
Some expressions..



$$\mathcal{N}_1 = s_{12} s_{23} s_{45}^2 (\ell_1 + k_5)^2$$

$$\mathcal{N}_3 = \frac{s_{45}^2}{\epsilon_5} G(\ell_1, k_1, k_2, k_3, k_4)$$

$$\mathcal{N}_4 = \frac{s_{45}^2}{\epsilon_5} G\left(\ell_1, k_1, k_2, k_3, k_4\right)$$



$$\mathcal{N}_{1} = s_{12}s_{34}s_{45}^{2}(\ell_{2} - k_{1})^{2}$$

$$\mathcal{N}_{5} = \frac{s_{45}s_{12}}{\epsilon_{5}}G(\ell_{2}, k_{1}, k_{2}, k_{3}, k_{4})$$

$$\mathcal{N}_{6} = \frac{s_{45}s_{12}}{\epsilon_{5}}G\begin{pmatrix} \ell_{1}, k_{1}, k_{2}, k_{3}, k_{4} \\ \ell_{2}, k_{1}, k_{2}, k_{3}, k_{4} \end{pmatrix}$$



$$\begin{split} \mathcal{N}_{3} &= -2 \frac{s_{12}s_{15}s_{23}s_{34}s_{45}}{\epsilon_{5}} s_{15} \ell_{2}^{2} \\ &+ \frac{s_{23}s_{34} \left(s_{12}s_{15} + s_{45}s_{15} - s_{12}s_{23} + s_{23}s_{34} - s_{34}s_{45}\right)}{\epsilon_{5}} s_{15} \ell_{2}^{4} \\ &+ \frac{s_{34}s_{45} \left(s_{12}s_{15} - s_{45}s_{15} + s_{12}s_{23} - s_{23}s_{34} + s_{34}s_{45}\right)}{\epsilon_{5}} s_{15} \ell_{2}^{2} \left(\ell_{2} - k_{1}\right)^{2} \\ &- \frac{s_{15}s_{45} \left(s_{12}s_{15} - s_{45}s_{15} - s_{12}s_{23} - s_{23}s_{34} + s_{34}s_{45}\right)}{\epsilon_{5}} s_{15} \ell_{2}^{2} \left(\ell_{2} - k_{1} - k_{2}\right)^{2} \\ &+ \frac{s_{12}s_{15} \left(s_{12}s_{15} - s_{45}s_{15} - s_{12}s_{23} + s_{23}s_{34} + s_{34}s_{45}\right)}{\epsilon_{5}} s_{15} \ell_{2}^{2} \left(\ell_{2} - k_{1} - k_{2} - k_{3}\right)^{2} \\ &- \frac{s_{12}s_{23} \left(s_{12}s_{15} - s_{45}s_{15} - s_{12}s_{23} + s_{23}s_{34} - s_{34}s_{45}\right)}{\epsilon_{5}} s_{15} \ell_{2}^{2} \left(\ell_{2} - k_{1} - k_{2} - k_{3} - k_{4}\right)^{2} \end{split}$$

$$\mathcal{N}_6 = \frac{\varepsilon}{1+2\varepsilon} \frac{s_{12}-s_{45}}{\epsilon_5} G \begin{pmatrix} \ell_1,\ell_2,k_1,k_2,k_3,k_4 \\ \ell_1,\ell_2,k_1,k_2,k_3,k_4 \end{pmatrix} \qquad \qquad \text{By leading singularities analysis in d-dimensional Baikov Rep.}$$



d-dimensional Baikov Rep.

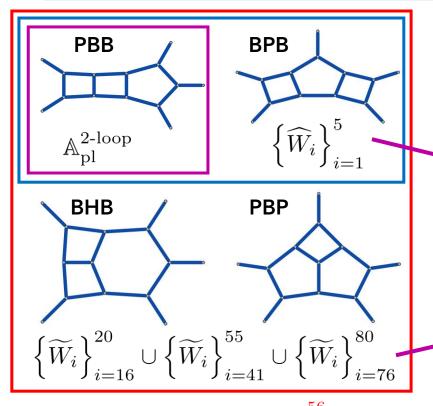
3L4P+1 off-shell leg from

[Di Vita, and Mastrolia, and Schubert, and Yundin '14] [Henn, Lim, Bobadila'23]

DlogBasis [Henn, Mistlberger, Smirnov, Wasser '20]

Many thanks to my incredible collaborators!

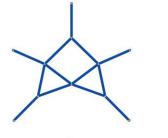
Three Loop Planar Pentagon Alphabet



Alphabet Letters: DNA of Feynman integrals

Conjectured in [Chicherin, Henn, Trnka, Zhang '24]

Originate from two topologies



$$\{\widehat{W}_i\}_{i=1}^5$$

$$W_1 = s_{23}s_{34} - s_{34}s_{45} + s_{45}s_{15}$$

New Square Root!

$$\left\{\sqrt{\Delta_4^{(i)}}\right\}_{i=1}^5$$

$$\begin{split} \Delta_4^{(1)} = & s_{15}^2 s_{12}^2 + s_{23}^2 s_{12}^2 - 2 s_{15} s_{23} s_{12}^2 - 2 s_{23}^2 s_{34} s_{12} \\ & + 2 s_{15} s_{23} s_{34} s_{12} + 2 s_{15} s_{34} s_{45} s_{12} + 2 s_{23} s_{34} s_{45} s_{12} \\ & + s_{23}^2 s_{34}^2 + s_{34}^2 s_{45}^2 - 2 s_{23} s_{34}^2 s_{45} \end{split}$$

56 letters

- = 26 at two-loop.
- +30 more at three-loop
- $d\mathbf{J} = \varepsilon \sum_{i=1}^{50} \mathbf{a}_i d\log W_i \cdot \mathbf{J}$

 $J \in \{PBB,BPB,BHB,PBP\}$

Were found by Baikovletter

[Jiang, Liu, Xu, Yang '24]

Solve

Analytic Boundary Values

by imposing absence of spurious singularities to Differential Equation

$$\operatorname{Res}\left(\frac{\mathrm{d}\widetilde{\mathbf{A}}}{\mathrm{d}x}\right)\bigg|_{y_i}\cdot\mathbf{I}^{(n)}(y_i)=0 \quad \square$$

$$\operatorname{Res}\left(\frac{\mathrm{d}\widetilde{\mathbf{A}}}{\mathrm{d}x}\right)\Big|_{y_i}\cdot\mathbf{I}^{(n)}(y_i) = 0 \quad \square \qquad \qquad \mathbf{I}_i^{(n)}(x_0) = c_i\mathbf{I}_1^{(n)}(x_0) + \sum_{j,k,l}d_{ijk}\int_{\gamma}\mathrm{d}\widetilde{A}_{jl}\cdot\mathbf{I}_l^{(n-1)}(x)$$

spurious singularities:

$$y_1 = \{-2, -1, -1, -1, -1\}$$
 and its cyclic permutation

$$y_6 = \{1/4, -1, -1, -1, -1\}$$

Solve the boundary values up to weight-six

Written in special values of GPL functions

[Gehrmann, Henn, Lo Presti '18] etc.

Symmetric Point in **Euclidean Region**

$$x_0 = y_0 = \{-1, -1, -1, -1, -1\}$$

$$\mathbf{I}_{i}^{(n)}(x_{0}) = c_{i}\mathbf{I}_{1}^{(n)}(x_{0}) + \sum_{j,k,l} d_{ijk} \int_{\gamma} d\widetilde{A}_{jl} \cdot \mathbf{I}_{l}^{(n-1)}(x_{0})$$

Line Integration $y_k \bullet \gamma$

◆ Fast numerical evaluation

***** See also: talk by Li-Hong Huang

more than 6,000 CPU hours by pySecDec for one top sector integral



● Symbol solution up to weight-six

$$[W_{i_1}, W_{i_2} \cdots, W_{i_n}] = \int W_{i_n}(t_{i_n}) \int \cdots \int W_{i_2}(t_2) \int W_{i_1}(t_1) dt_1 dt_2 \cdots dt_{i_n}$$

Three-loop topologies contribute only 2220 weight-six symbol $\ll 5 \times 56^5$



Bootstrapping Wilson loop with Lagrangian insertion

General See also: talk by Yang Zhang

◆ Pentagon functions in Euclidean region up to ...

Pentagonal Wilson Loop with Lagrangian Insertion at Three-Loop

$$\mathcal{N} = 4 \text{ Super Yang-Mills in planar limit}$$

$$F_5 = \frac{\langle W[x_1, x_2, x_3, x_4, x_5] \mathcal{L}(x_0) \rangle}{\langle W[x_1, x_2, x_3, x_4, x_5] \rangle} = \frac{\langle W[x_1, x_2, x_3, x_4, x_5] \mathcal{L}(x_0) \rangle}{\langle W[x_1, \dots, x_n] = \operatorname{tr}_{\operatorname{FP}} \exp\left(\mathrm{i} g_{\mathrm{YM}} t^a \oint A_{\mu}^a \mathrm{d} x^{\mu}\right)} \text{ n-cusp null Wilson loop with a Lagrangian insertion paramalized by Wilson loop without Insertion paramalized by Wilson loop without Insertion$$

n-cusp null Wilson loop with a Lagrangian insertion

normalized by Wilson loop without Insertion

Theoretical Importance

closely related to scattering amplitudes

- Null Wilson Loop ~ amplitudes in N=4 sYM
- Natural in AdS/CFT
- Duality with all-plus pure Yang-Mills amplitudes
- Positivity and geometry
- ..

[Alday, Maldacena '07]

[Drummond, Korchemsky, Sokatchev '07]

[Arkani-Hamed, Henn, Trnka '21]

[Chicherin, Henn '22] [Chicherin, Henn '22]

[Chicherin, Henn, Trnka, Zhang '24] [Brown, Henn, Mazzucchelli, Trnka '24] and so on.

Easy to compute

- ◆ Free of IR divergence
- Uniformly Transcendental of 2L
- Exact dual conformal symmetry
- **◆** All Loop Leading Singularity
- ...



A good application of our Three-Loop Five-Point integrals

(pre) History: A Four-Point Example



Transcendental: Analytic functions
Rational: Leading singularities —

$$\frac{\langle W[x_1, x_2, x_3, x_4] \mathcal{L}(x_5) \rangle}{\langle W[x_1, x_2, x_3, x_4] \rangle} = \frac{2}{\pi^2} \frac{x_{13}^2 x_{24}^2}{x_{15}^2 x_{25}^2 x_{35}^2 x_{45}^2} F(x) + \mathcal{O}(\varepsilon)$$

$$p_i \equiv x_i - x_{i+1}$$

[Alday, Buchbinder, Tseytlin '11] [Engelund, Roiban '11]

$$F(x) = -\frac{1}{4}g^2 + \left(\frac{1}{8}H_{0,0} + \frac{\pi^2}{16}\right)g^4 + \left(-\frac{\pi^2}{16}H_{0,0} - \frac{3}{16}H_{0,0,0,0} + \cdots\right)g^6 + \left(\frac{323\pi^4}{11520}H_{0,0} + \cdots\right)g^8 + \cdots$$

[Alday, Buchbinder, Tseytlin '11]

 $g^2 \equiv \frac{g_{\rm YM}^2 N_{\rm c}}{4\pi^2}$

[Alday, Heslop, Sikorowski '12]

[Alday, Henn, Sikorowski '13]

[Henn, Korchemsky, Mistlberger '19]

Review of perturbative results

	Four-Point	Five-Point	Six-Point
Tree-Level	[Alday, Buchbinder, Tseytlin '11]	[Chicherin, Henn '22] n-point Tree-Level and One-Loop	
One-Loop	[Alday, Heslop, Sikorowski '12]		
Two-Loop	[Alday, Henn, Sikorowski '13]	[Chicherin, Henn '22]	[Carrôlo, Chicherin, Henn Yang, Zhang ' 25]
Three-Loop	[Henn, Korchemsky, Mistlberger '19]	$F_5^{(3)}$	

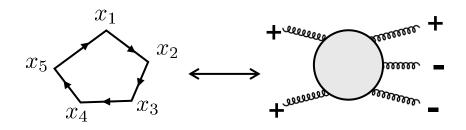
$$F_5 = \sum_{L \ge 1} (g^2)^{L+1} F_5^{(L)} = \lim_{x_0 \to \infty} \frac{\langle W[x_1, x_2, x_3, x_4, x_5] \mathcal{L}(x_0) \rangle}{\langle W[x_1, x_2, x_3, x_4, x_5] \rangle}$$
 We are now Here



See also

[Alday, Buchbinder, Tseytlin '11] Four-Point @ Strong Coupling [Arkani-Hamed, Henn, Trnka '21] Four-Point Negative geometries expansion

Integrand



planar N=4 sYM

from MHV amplitude Wilson Loop duality

$$\log M_n + \mathcal{O}(1/N) \leftrightarrow \log \langle W_n \rangle$$

[Alday, Maldacena '07]

[Drummond, Korchemsky, Sokatchev '07]

[Alday, Roiban '08] [Henn '09]

Loop expansion Lagrangian insertion: $g^2 \partial_{g^2} \langle W_n \rangle = -\mathrm{i} \int \mathrm{d}^d x_0 \langle W_n \mathcal{L}(x_0) \rangle$ MHV amplitudes integrand

$$-i \int d^d x_0 \frac{\langle W[x_1, \cdots, x_n] \mathcal{L}(x_0) \rangle}{\langle W[x_1, \cdots, x_n] \rangle} = g^2 \partial_{g^2} \log \langle W_n \rangle \sim g^2 \partial_{g^2} \log M_n$$

$$W[x_1, \dots, x_n] = \operatorname{tr}_{\mathbf{F}} \mathbf{P} \exp \left(ig_{\mathbf{YM}} t^a \oint A^a_{\mu} dx^{\mu} \right)$$

L-loop Integrand
$$F_n^{(L)} o \int \mathrm{d}^4 y_1 \cdots \int \mathrm{d}^4 y_L \Omega^{(L+1)}$$

[Chicherin, Henn '22]

[Chicherin, Henn, Trnka, Zhang '24]

[Brown, Henn, Mazzucchelli, Trnka '24]

With numerator which suppress all cusp divergence

Direct Computation of $F_5^{(3)}$

 $\mathcal{D}_{16} = \left(\ell_1 - \ell_2\right)^2$

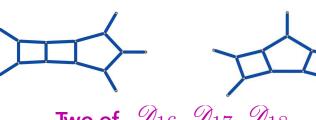
Step1: Separating Targets

$$\begin{array}{|c|c|c|c|c|} \hline \mathcal{D}_{1} = (\ell_{1} - k_{1})^{2} \\ \mathcal{D}_{2} = (\ell_{1} - k_{1} - k_{2})^{2} \\ \mathcal{D}_{3} = (\ell_{1} - k_{1} - k_{2} - k_{3})^{2} \\ \mathcal{D}_{4} = (\ell_{1} - k_{5})^{2} \\ \mathcal{D}_{5} = \ell_{1}^{2} \\ \end{array} \begin{array}{|c|c|c|c|c|} \hline \mathcal{D}_{6} = (\ell_{2} - k_{1})^{2} \\ \mathcal{D}_{7} = (\ell_{2} - k_{1} - k_{2})^{2} \\ \mathcal{D}_{8} = (\ell_{2} - k_{1} - k_{2} - k_{3})^{2} \\ \mathcal{D}_{9} = (\ell_{2} - k_{5})^{2} \\ \mathcal{D}_{10} = \ell_{2}^{2} \\ \hline \end{array} \begin{array}{|c|c|c|c|c|c|} \mathcal{D}_{11} = (\ell_{3} - k_{1})^{2} \\ \mathcal{D}_{12} = (\ell_{3} - k_{1} - k_{2})^{2} \\ \mathcal{D}_{13} = (\ell_{3} - k_{1} - k_{2} - k_{3})^{2} \\ \mathcal{D}_{14} = (\ell_{3} - k_{5})^{2} \\ \mathcal{D}_{15} = \ell_{3}^{2} \\ \hline \end{array}$$

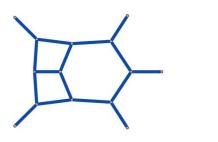
 $\mathcal{D}_{17} = \left(\ell_1 - \ell_3\right)^2$

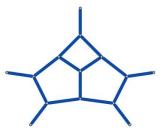
Loop momenta coupling

Free of
$$\mathscr{D}_{16}, \mathscr{D}_{17}, \mathscr{D}_{18}$$
One of $\mathscr{D}_{16}, \mathscr{D}_{17}, \mathscr{D}_{18}$



Two of $\mathcal{D}_{16}, \mathcal{D}_{17}, \mathcal{D}_{18}$





 $\mathcal{D}_{18} = \left(\ell_2 - \ell_3\right)^2$

All of $\mathcal{D}_{16}, \mathcal{D}_{17}, \mathcal{D}_{18}$

Step 1: Separating targets in integrand

$$F_5^{(3)} \to \int d^4 y_1 \int d^4 y_2 \int d^4 y_3 \Omega^{(4)}$$

Step 2: Reduce all integrals into UT basis

NEATIBP



Rational Part: Leading Singularities

$$r_0 = \operatorname{tr}_-((p_1 + p_2)(p_2 + p_3)(p_3 + p_4)(p_4 + p_5))$$

$$r_i = \frac{s_{i+2,i+3}}{s_{i+1,i+4}} \operatorname{tr}_-(p_{i+1}p_{i+2}p_{i+3}p_{i+4}), \ i = 1 \cdots 5$$

Transcendental Part: Weight-Six Symbol

$$g_i^{(3)} = \sum_{i_1, i_2, i_3, i_4, i_5, i_6} t_{i_1, i_2, i_3, i_4, i_5, i_6} [W_{i_1}, W_{i_2}, W_{i_3}, W_{i_4}, W_{i_5}, W_{i_6}]$$

$$[W_{i_1}, W_{i_2} \cdots, W_{i_n}] = \int W_{i_n}(t_{i_n}) \int \cdots \int W_{i_2}(t_2) \int W_{i_1}(t_1) dt_1 dt_2 \cdots dt_{i_n}$$

Cancellation between 3,107,425 terms Results involves 1,373,140 terms...

Dual to four-loop five-point all plus pure Yang-Mills amplitudes?

- ✓ IR finite: All weight 0 ... 5 symbol vanished
- ☑ Respect soft and collinear Limits

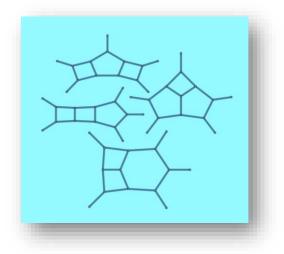
[Wu, Böhm, Ma, Xu, Zhang '23] [Wu, Böhm, Ma, Usovitsch, Xu, Zhang '25]

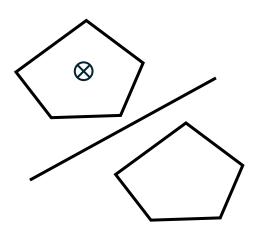
[Chicherin, Henn '22] [Chicherin, Henn '22]

[Chicherin, Henn, Trnka, Zhang '24]

[Brown, Henn, Mazzucchelli, Trnka '24]

Conclusion





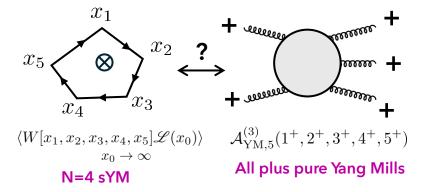
Into the Future

- Boundary value and pentagon function in physical region
- **◆** Function level Wilson Loop with Lagrangian insertion

Positivity properties
$$(-1)^{(L+1)}f_5^{(L)}|_{\text{Eucl}^+} > 0, L=0,1,2$$
 [Chicherin, Henn '22]

Complete monotonicity?
$$(-\partial_x)^n f(x) \ge 0, \forall n \ge 0$$
 [Henn, Raman '24]

 $lacktriang{lacktriang{}}$ Pure Yang-Mills all plus amplitudes. Is it dual to $F_5^{(2)}$?



The expression (6.7) is a special case of the *n*-particle two-loop duality relation that we verified in [25]. However, eq. (6.8) is new, and it predicts the maximally transcendental part of the five-particle three-loop all-plus planar amplitude.

It would be extremely interesting to test our conjecture about the maximally transcendental part of the planar five-particle all-plus amplitude by an explicit three-loop Feynman graph calculation. Furthermore, a three-loop calculation of f_5 will provide a conjecture for the maximally transcendental part of the planar four-loop five-particle all-plus amplitude.

[Chicherin, Henn '22]

Many more other applications...

