



Precision Predictions for Three-Dimensional Nucleon Tomography

方申

Fudan University

第五届量子场论及其应用研讨会

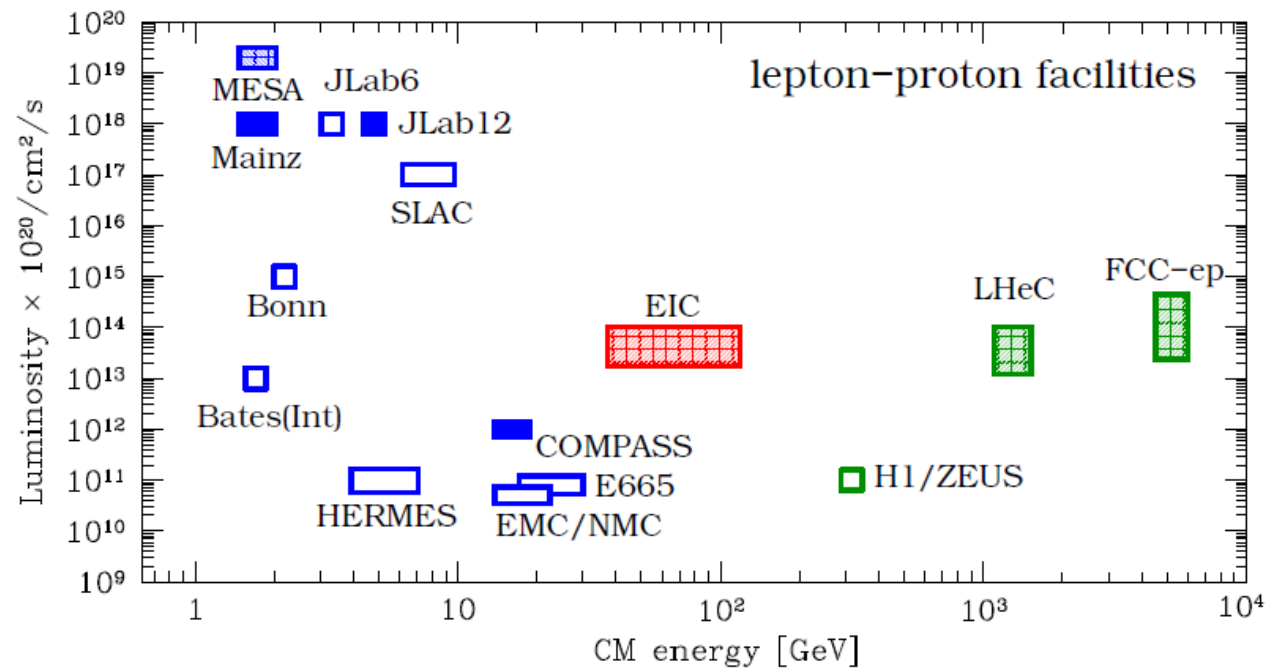
Beijing

Oct 31, 2025

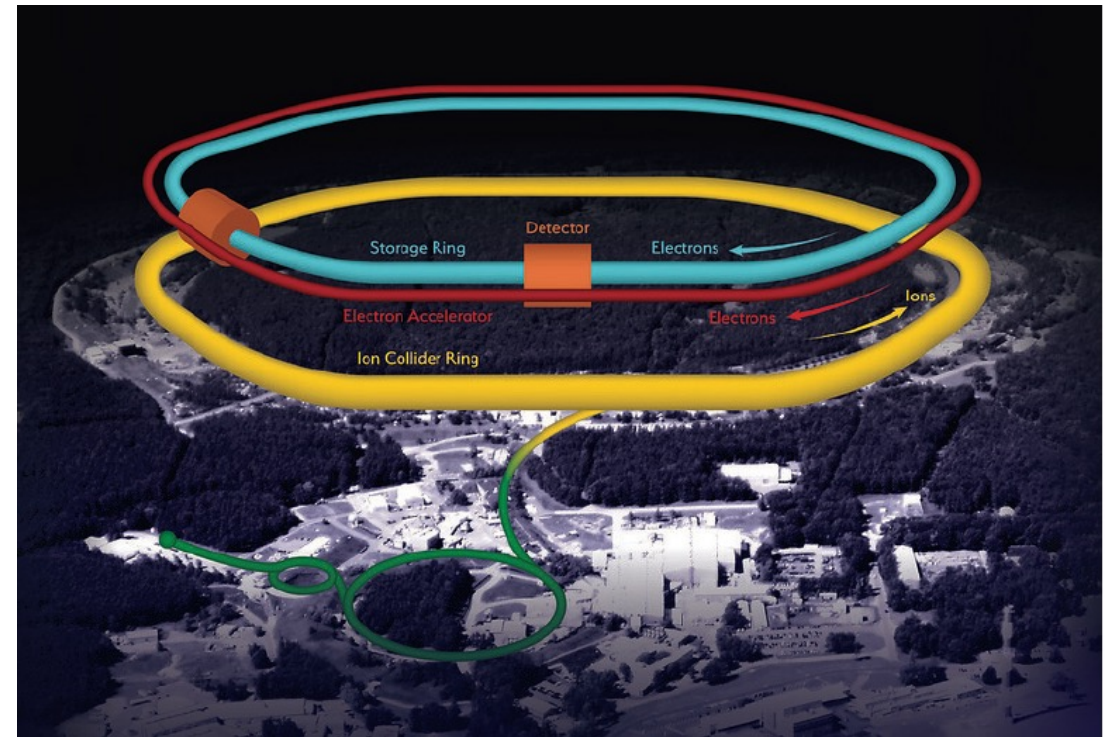
Collaborators: Dingyu Shao, Weiyao Ke, John Terry,
Meisen Gao, Haitao Li

Reference: JHEP05(2024)066, JHEP01(2025)029

Electron-ion collider (EIC)

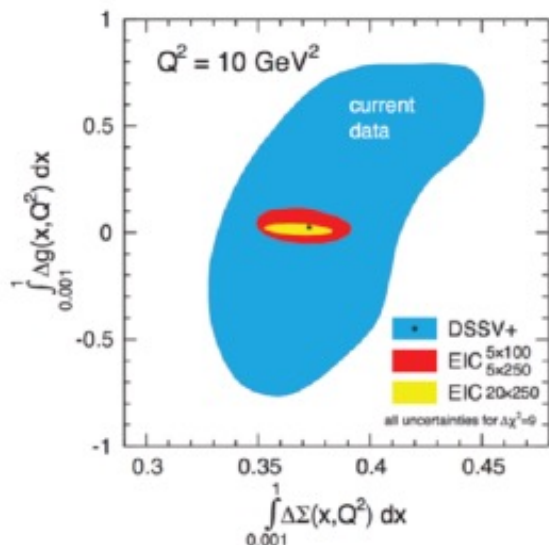


Abdul et al. '22

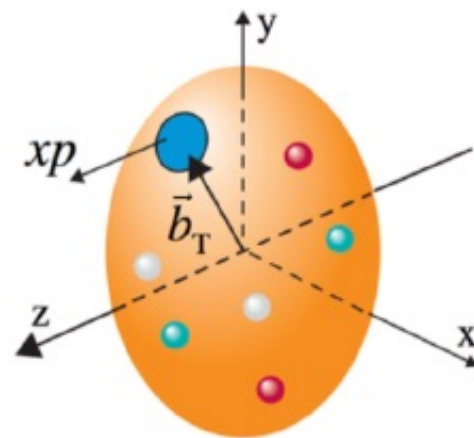


Abdul et al. '22

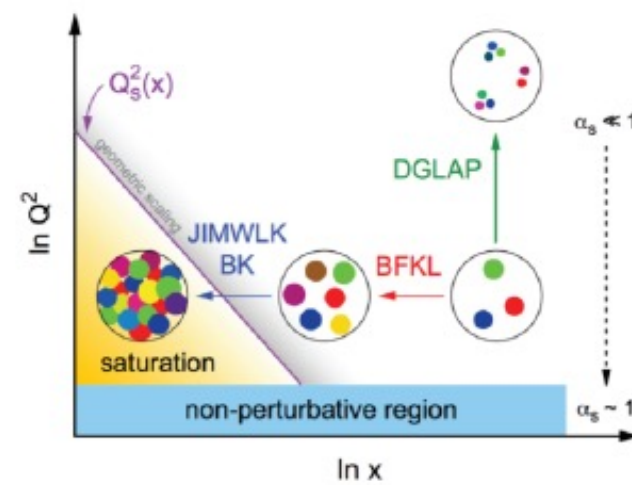
Proton spin



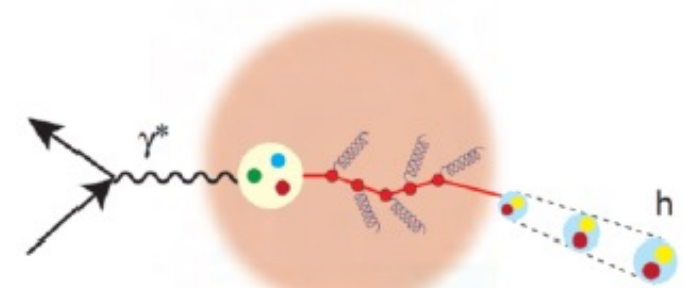
3D nucleon tomography



Gluon saturation



Hadronization in the nucleus

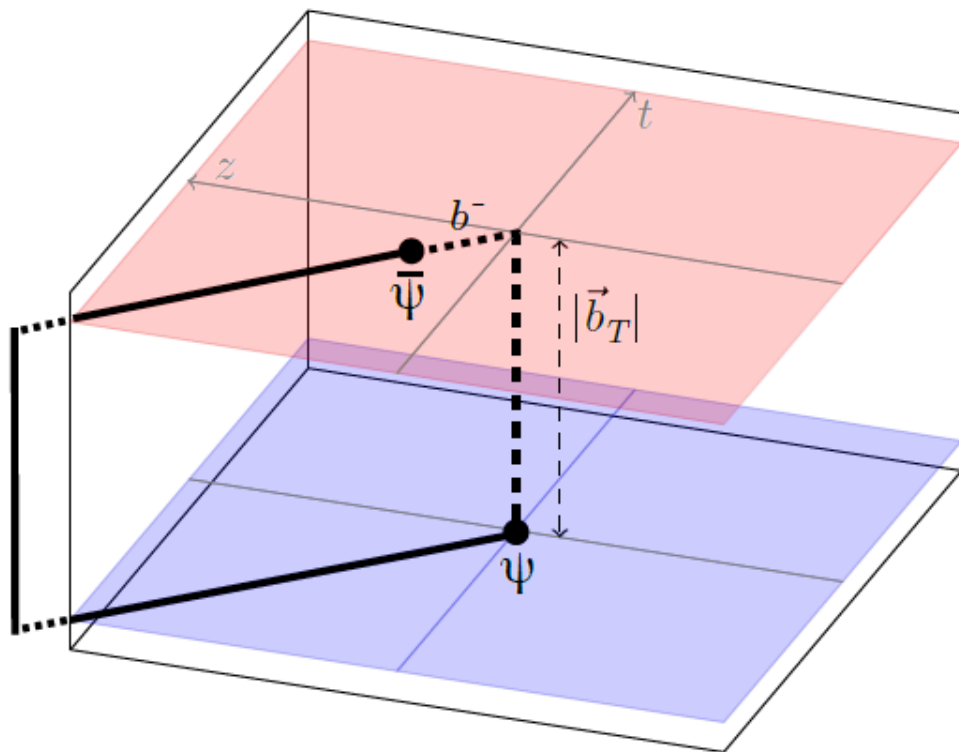


3D imaging of the nucleons

- Both **longitudinal** and **transverse** motion
- Correlation between nucleon spin with parton(quark, gluon) orbital angular momentum

$$\tilde{f}_{i/ps}^{[\Gamma]0(u)}(x, \mathbf{b}_T, \epsilon, \tau, xP^+) = \int \frac{db^-}{2\pi} e^{-ib^-(xP^+)} \left\langle p(P, S) \left| \left[\bar{\psi}^i(b^\mu) W_\square(b^\mu, 0) \frac{\Gamma}{2} \psi^i(0) \right]_\tau \right| p(P, S) \right\rangle$$

- Dirac structures $\Gamma \in \{ \gamma^+, \gamma^+ \gamma_5, i\sigma^{\alpha+} \gamma_5 \}$



Boussarie et al. '23

Leading Quark TMDPDFs

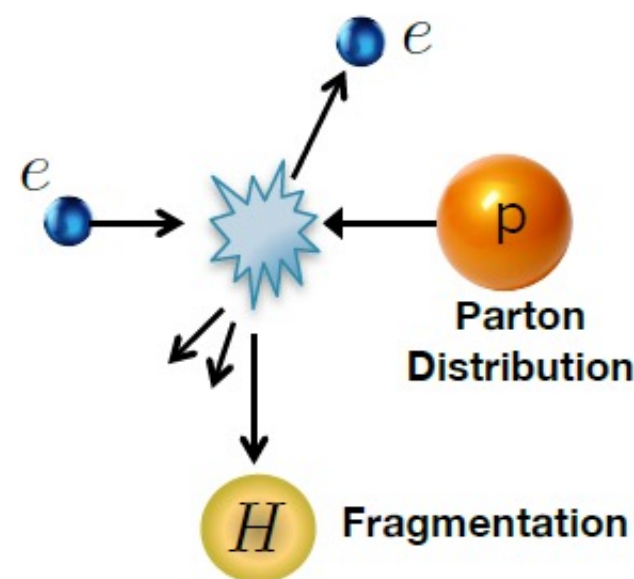
 Nucleon Spin
  Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{Unpolarized}$		$h_1^\perp = \text{Boer-Mulders}$
	L		$g_1 = \text{Helicity}$	$h_{1L}^\perp = \text{Worm-gear}$
	T	$f_{1T}^\perp = \text{Sivers}$	$g_{1T}^\perp = \text{Worm-gear}$	$h_1 = \text{Transversity}$ $h_{1T}^\perp = \text{Pretzelosity}$

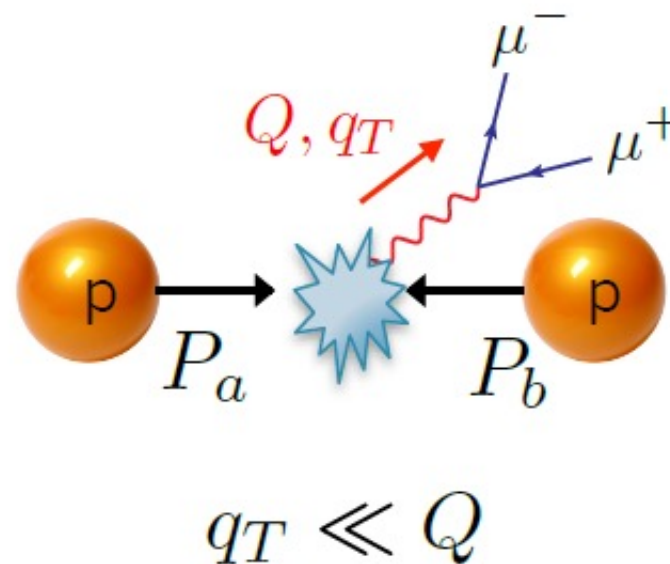
Transverse momentum distributions of quarks

- Three classical processes used to probe quark TMDs:

Semi-Inclusive DIS

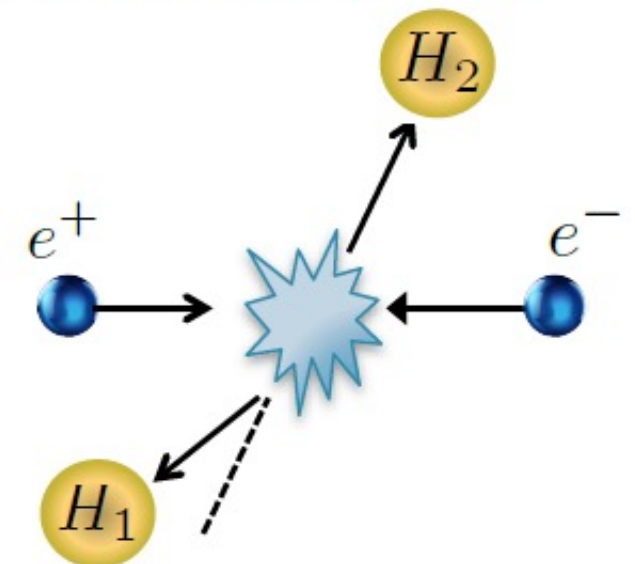


Drell-Yan



Boussarie et al. '23

Dihadron in e^+e^-



- Typical multi-scale problems
- Theory tools: TMD factorization

Collins, Soper & Sterman 1985
Florian & Grazzini '01

(Semi-inclusive) deep inelastic scattering

Factorized SIDIS cross section in the parton model:

Boussarie et al. '23

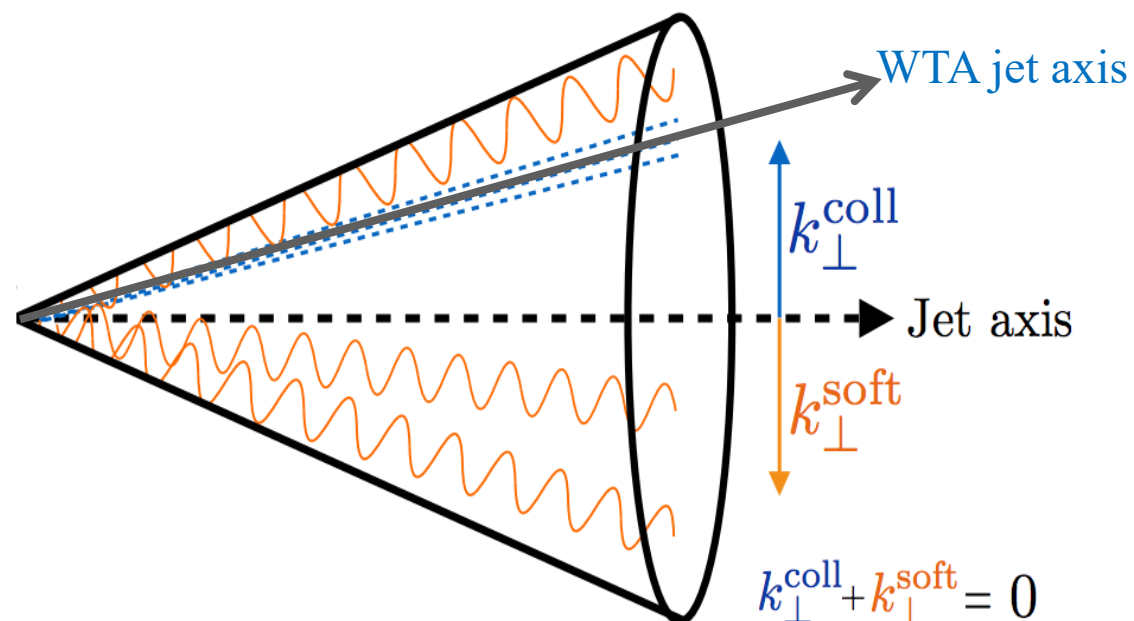
$$\sigma_{\text{SIDIS}} \propto \left| \text{Diagram 1} \right|^2 \approx \left| \text{Diagram 2} \right|^2 \otimes \left| \text{Diagram 3} \right|^2 \otimes \left| \text{Diagram 4} \right|^2$$

Liu, Ringer, Vogelsang & Yuan '18,20

New ideas: $e(\ell) + N(P) \rightarrow e(\ell') + J(P_J) + X$

Winner-Take-All jets/
Recoil-free jets

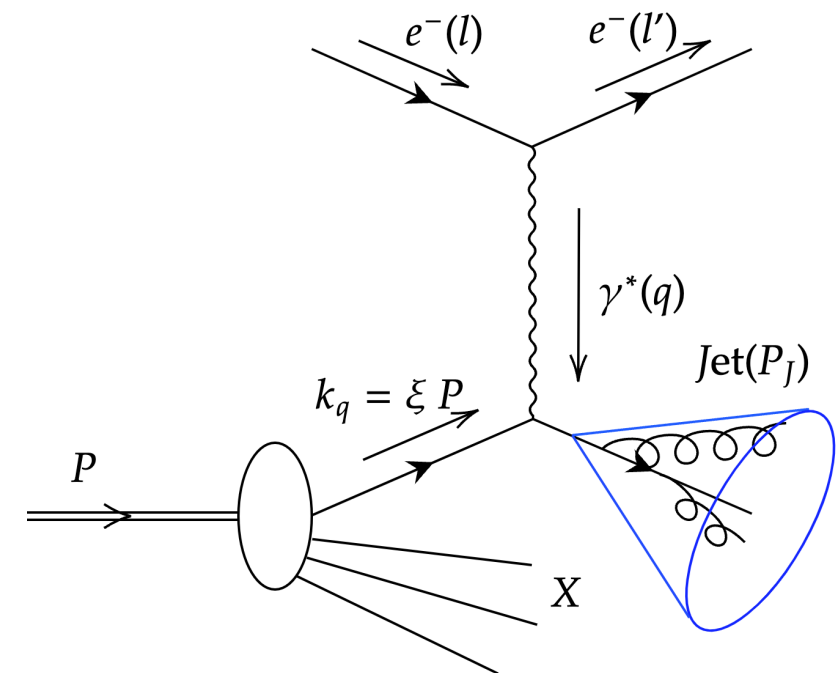
Bertolini, Chan & Thaler '14



$$\text{SJA : } E_{12} = E_1 + E_2 , \\ \mathbf{P}_{12} = \mathbf{p}_1 + \mathbf{p}_2 .$$

$$\text{WTA : } E_{12} = E_1 + E_2 ,$$

$$\mathbf{P}_{12} = E_{12} \left[\frac{\mathbf{p}_1}{|\mathbf{p}_1|} \theta(E_1 - E_2) + \frac{\mathbf{p}_2}{|\mathbf{p}_2|} \theta(E_2 - E_1) \right] .$$



EIC wish list

Regarding DATA

- Measure cross-sections instead of ratios for a more dedicated analysis
- Release both QED corrected and uncorrected data
- Develop method for unbinned cross-sections

Regarding PDFs, FFs and other distributions

- Replication of PDF4LHC and HERA efforts for EIC : PDF4EIC
- Perform global NNLO analysis of polarized PDFs
- Impact of QED corrections on polarized PDFs
- Perform global analysis of DVCS
- Generate threshold resummed PDFs and FFs
- New set of photon PDFs (existing are outdated)

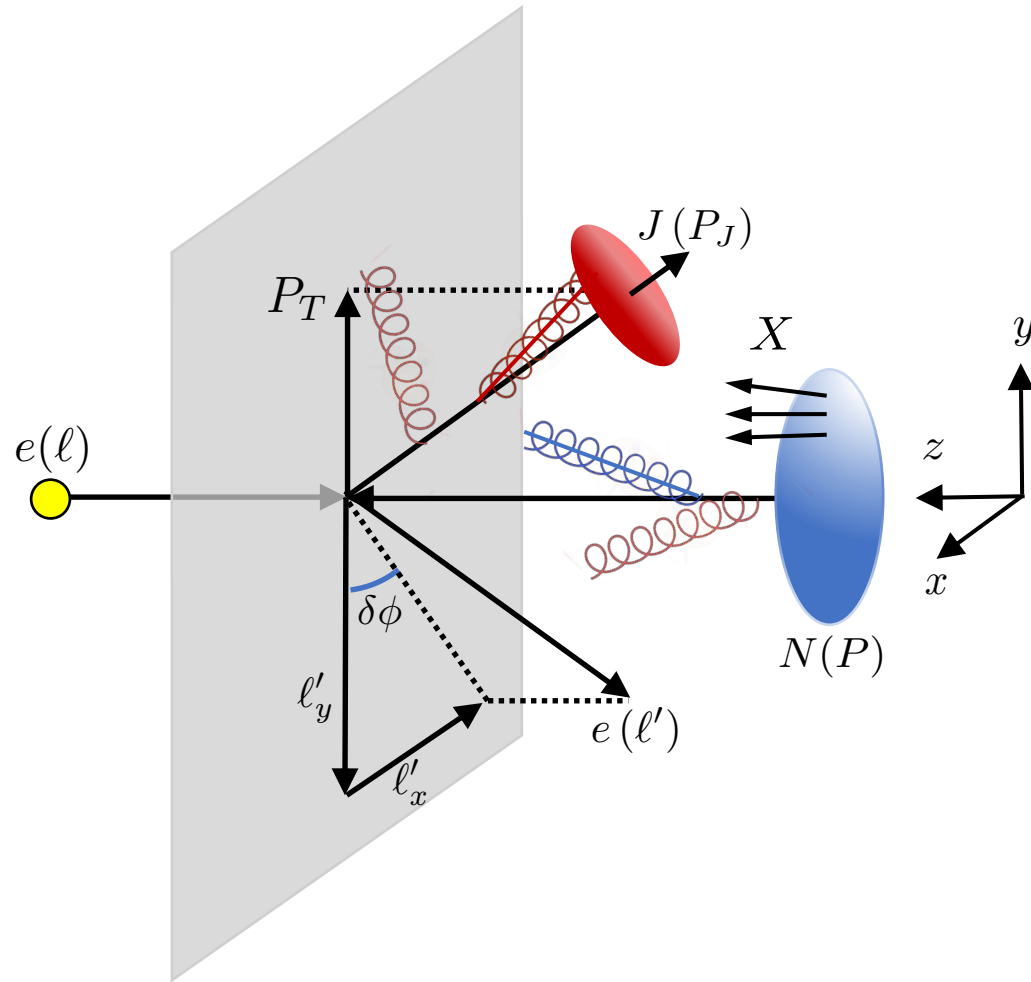
Regarding Perturbative corrections (QCD/QED)

- Jets in DIS: matched NNLO + q_T resummation
- DIS with QED/EW corrections
- Calculations for dihadron production

Regarding Theoretical Issues

- Discuss (non)universality of TMDs
- Search for ideal observables to measure Wigner distribution
- Role of lattice in PDFs (in two slides!)
- Studies for Λ polarization at EIC
- N^* , Δ electro-couplings at $Q^2 > \text{GeV}^2$
- Proton structure functions from transition regime to DIS
- Small- x dipole and quadrupole amplitudes

QCD factorization



$$e(\ell) + N(P) \rightarrow e(\ell') + J(P_J) + X$$

- Factorization formula :

$$\frac{d\sigma}{d^2\ell'_T dy dq_x} = \frac{\sigma_0}{1-y} H(Q, \mu) \mathcal{C}[B \mathcal{J} S] ,$$

$$\mathcal{C}[B \mathcal{J} S] = \sum_q e_q^2 \int \frac{db}{2\pi} \cos(b q_x) B_{q/N}(x_B, b, \mu, \zeta_B/\nu^2) \times \mathcal{J}_q(b, \mu, \zeta_{\mathcal{J}}/\nu^2) S(b, n \cdot n_J, \mu, \nu) .$$

- EFT modes :

hard : $p_h^\mu \sim \ell'_T(1, 1, 1)$,

n-collinear : $p_c^\mu \sim \ell'_T(\delta\phi^2, 1, \delta\phi)$,

soft : $p_s^\mu \sim \ell'_T(\delta\phi, \delta\phi, \delta\phi)$,

n_J-collinear : $p_J^\mu \sim \ell'_T(\delta\phi^2, 1, \delta\phi)_J$.

- Observables:

$$y = 1 - P \cdot \ell' / P \cdot l \quad q_x = \ell'_T \delta\phi$$

- EFT parameters:

$$\delta\phi \ll 1$$

RG evolution

Hierarchy Problem:

$$\mu_b = \frac{2e^{-\gamma_E}}{b} \ll Q \quad L = \ln \frac{Q^2}{\mu_b^2} \gg 1$$

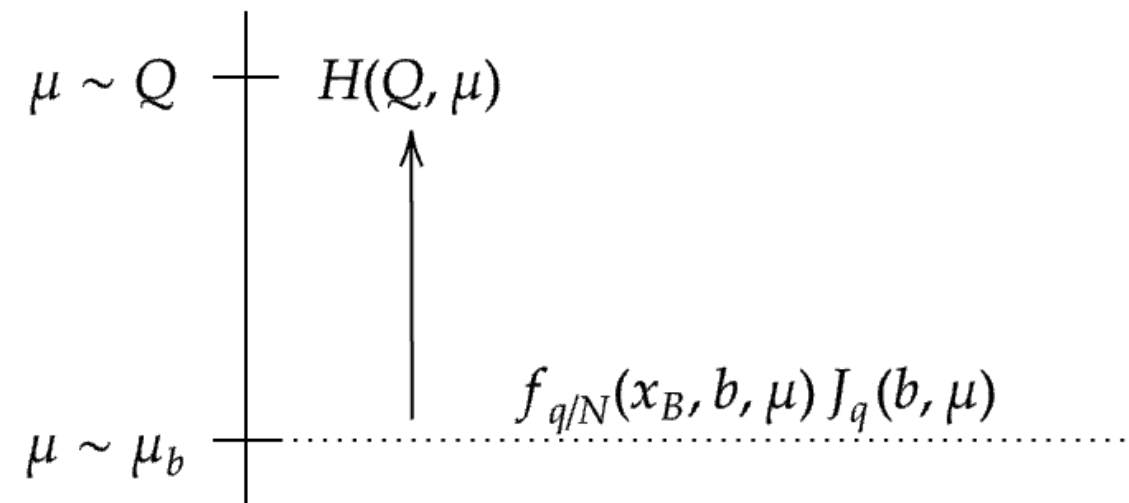
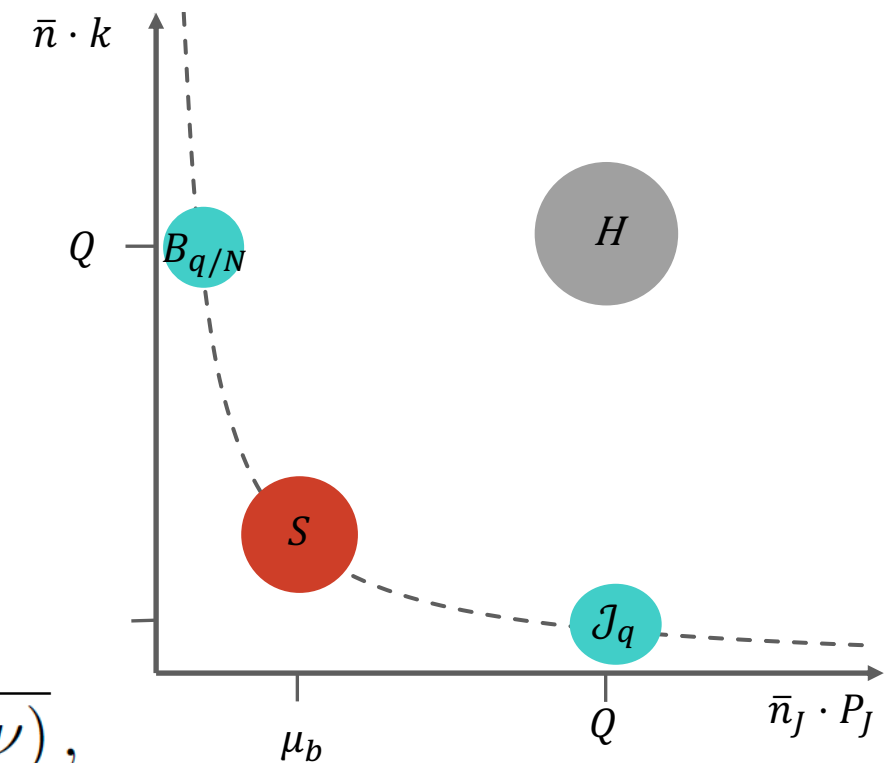
Standard CSS formalism: Collins '13

$$f_{q/N}(x_B, b, \mu, \zeta_f) = B_{q/N}(x_B, b, \mu, \zeta_B/\nu^2) \sqrt{S_{n\bar{n}}(b, \mu, \nu)},$$

$$J_q(b, \mu, \zeta_J) = \mathcal{J}_q(b, \mu, \zeta_J/\nu^2) \frac{S(b, n \cdot n_J, \mu, \nu)}{\sqrt{S_{n\bar{n}}(b, \mu, \nu)}},$$

Scale independence:

$$\frac{d}{d \ln \mu} \ln \sigma_{\text{phsy}}(Q, \mu_b, \mu) = 0$$



Predictions in e-A at NNLL

SF, Ke, Shao, Terry '23

non-perturbative model

$$U_{\text{NP}}^f(x, b, A, Q_0, Q) = \exp \left[-g_1^A b^2 - \frac{g_2}{2} \ln \frac{Q}{Q_0} \ln \frac{b}{b_*} \right],$$

Sun, Isaacson, Yuan & Yuan '14

We apply modified nuclear TMD PDFs

$$g_1^A = g_1^f + a_N(A^{1/3} - 1) \quad a_N = 0.016 \pm 0.003 \text{ GeV}^2$$

Collinear dynamics using EPPS16

Alrashed, Anderle, Kang, Terry & Xing '22

We include LO momentum broadening of the jet within SCET_G

$$J_q^A(b, \mu, \zeta_J) = J_q(b, \mu, \zeta_J) e^{\chi[\xi b K_1(\xi b) - 1]}$$

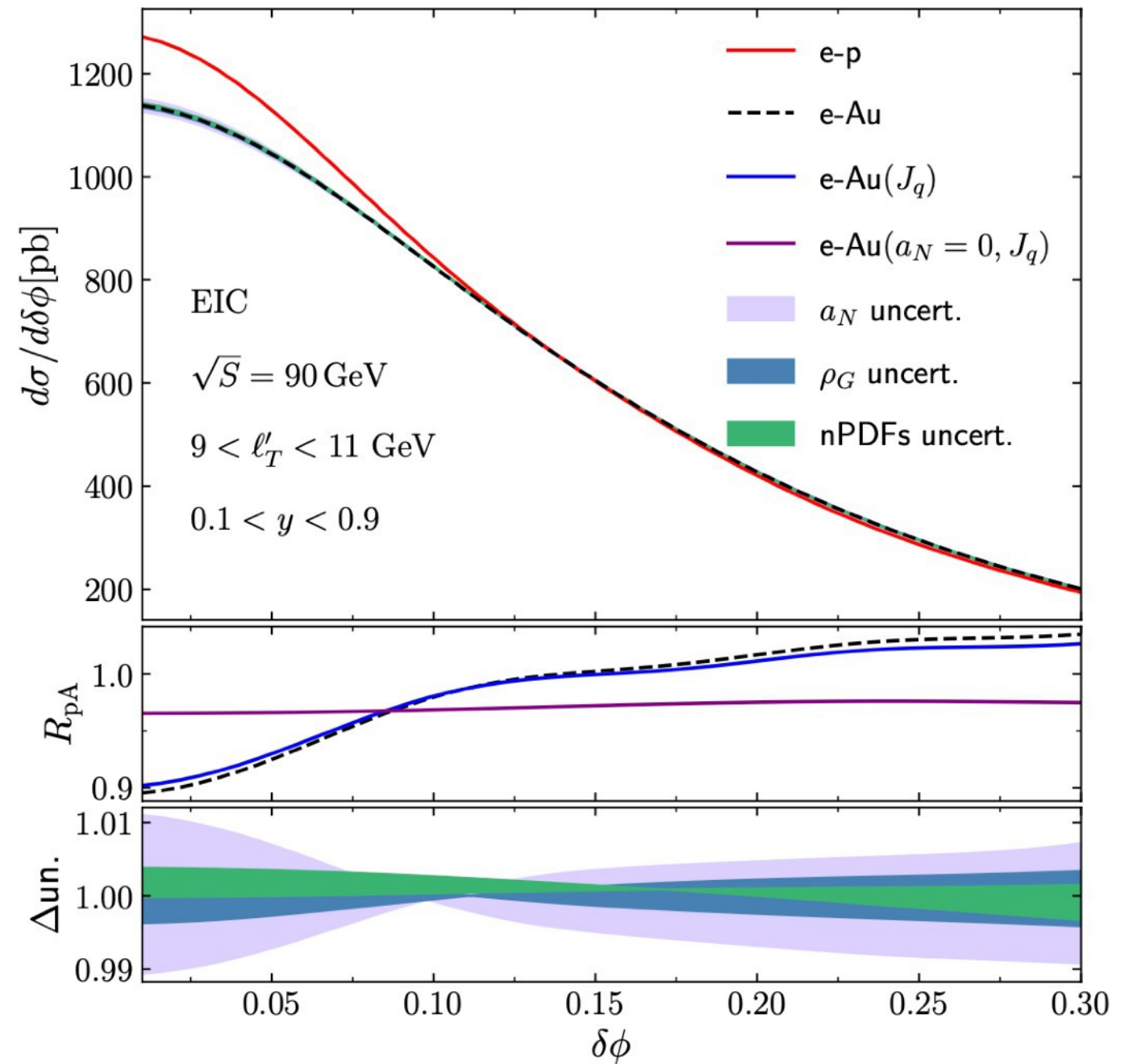
Opacity parameter $\chi = \frac{\rho_G L}{\xi^2} \alpha_s(\mu_{b_*}) C_F$

Gyulassy, Levai & Vitev '02

ρ_G : density of the medium

ξ : the screening mass

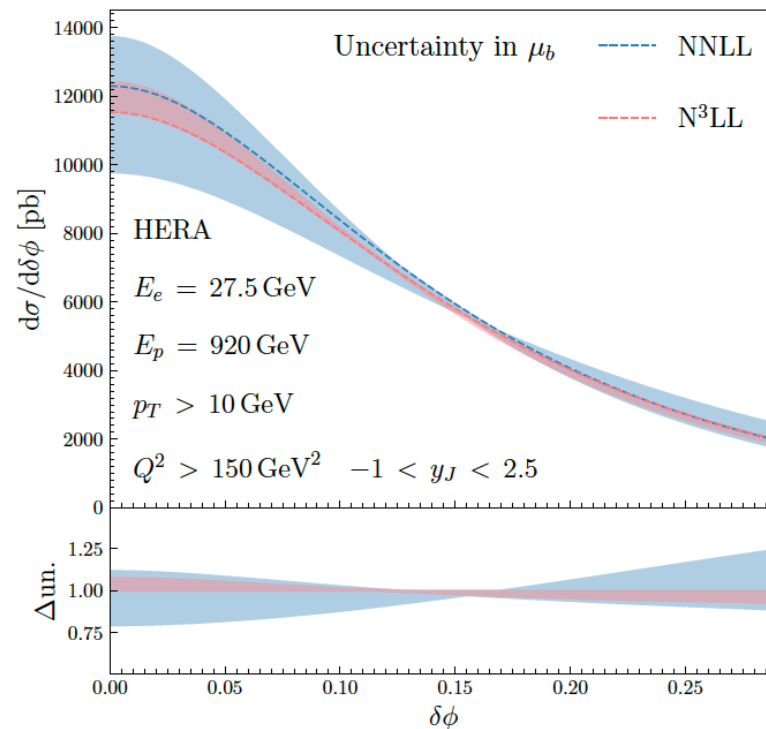
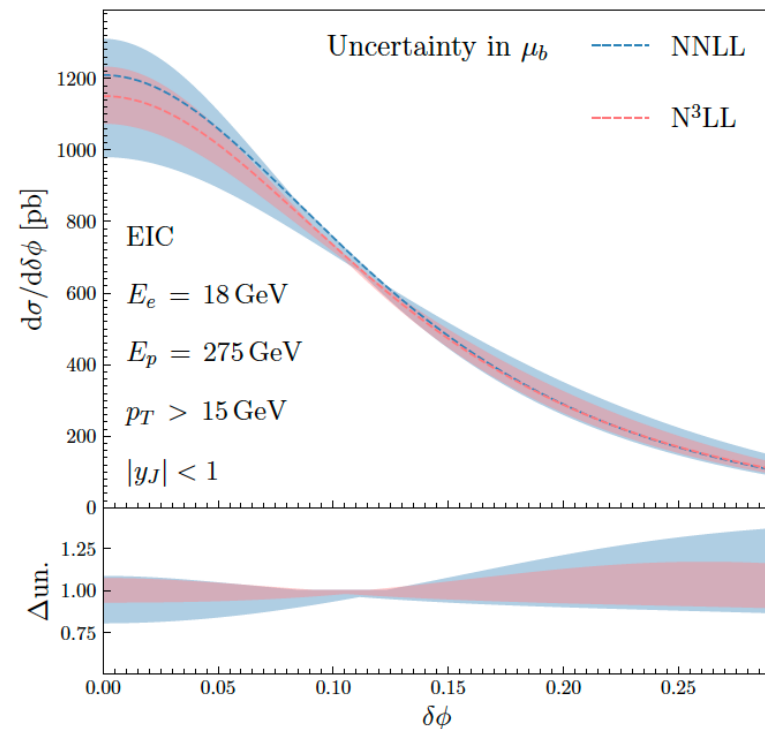
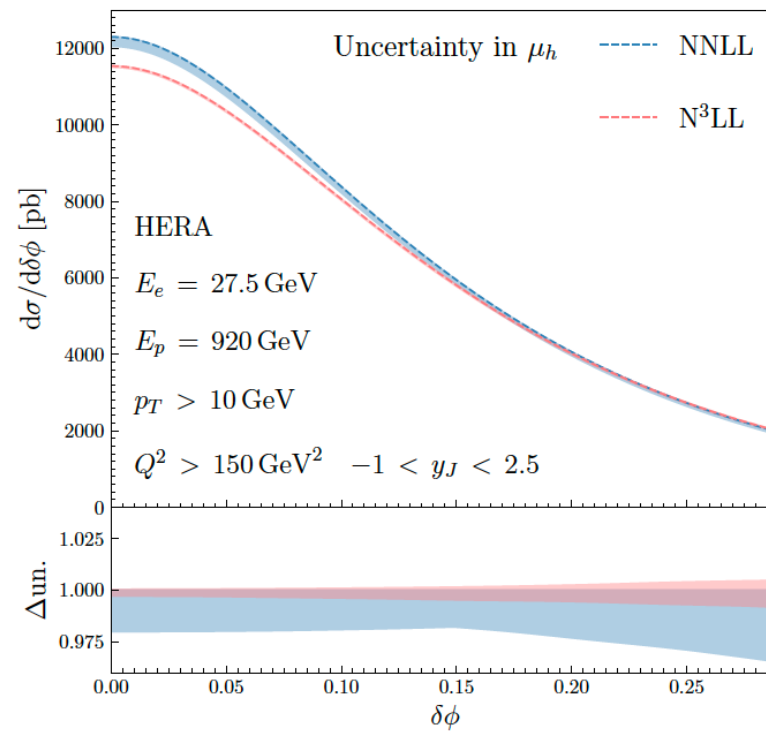
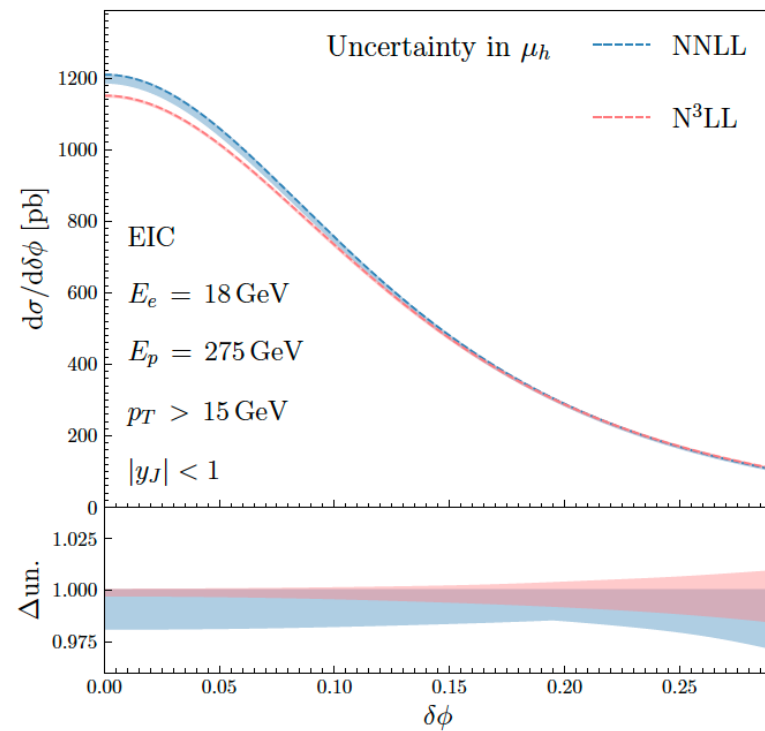
L : the length of the medium



The process is primarily sensitive to the initial state's broadening effects, thereby serving as a clean probe of nTMD PDFs.

Comparison of resummation results at NNLL and N³LL

SF, Gao, Li, Shao '24



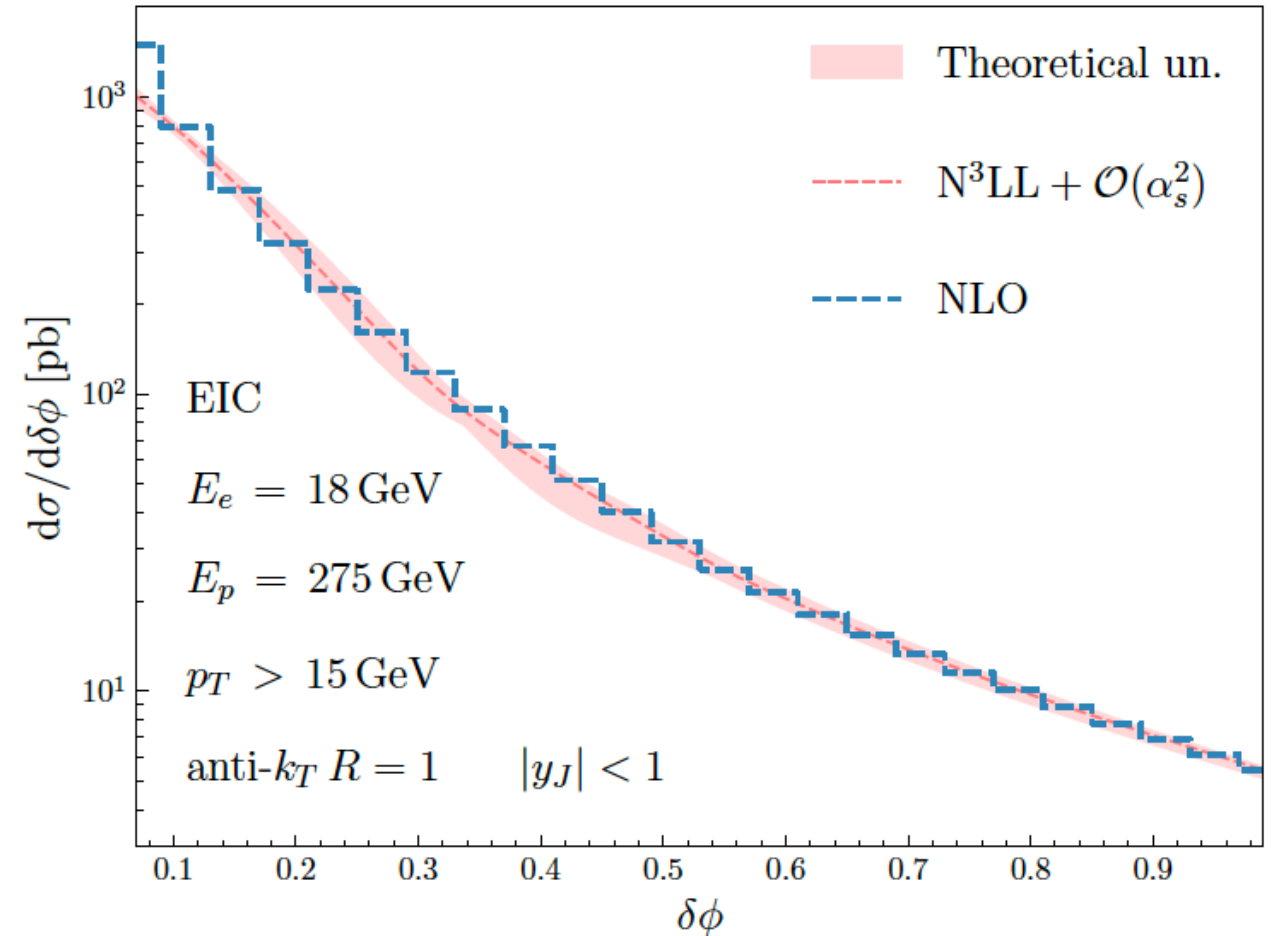
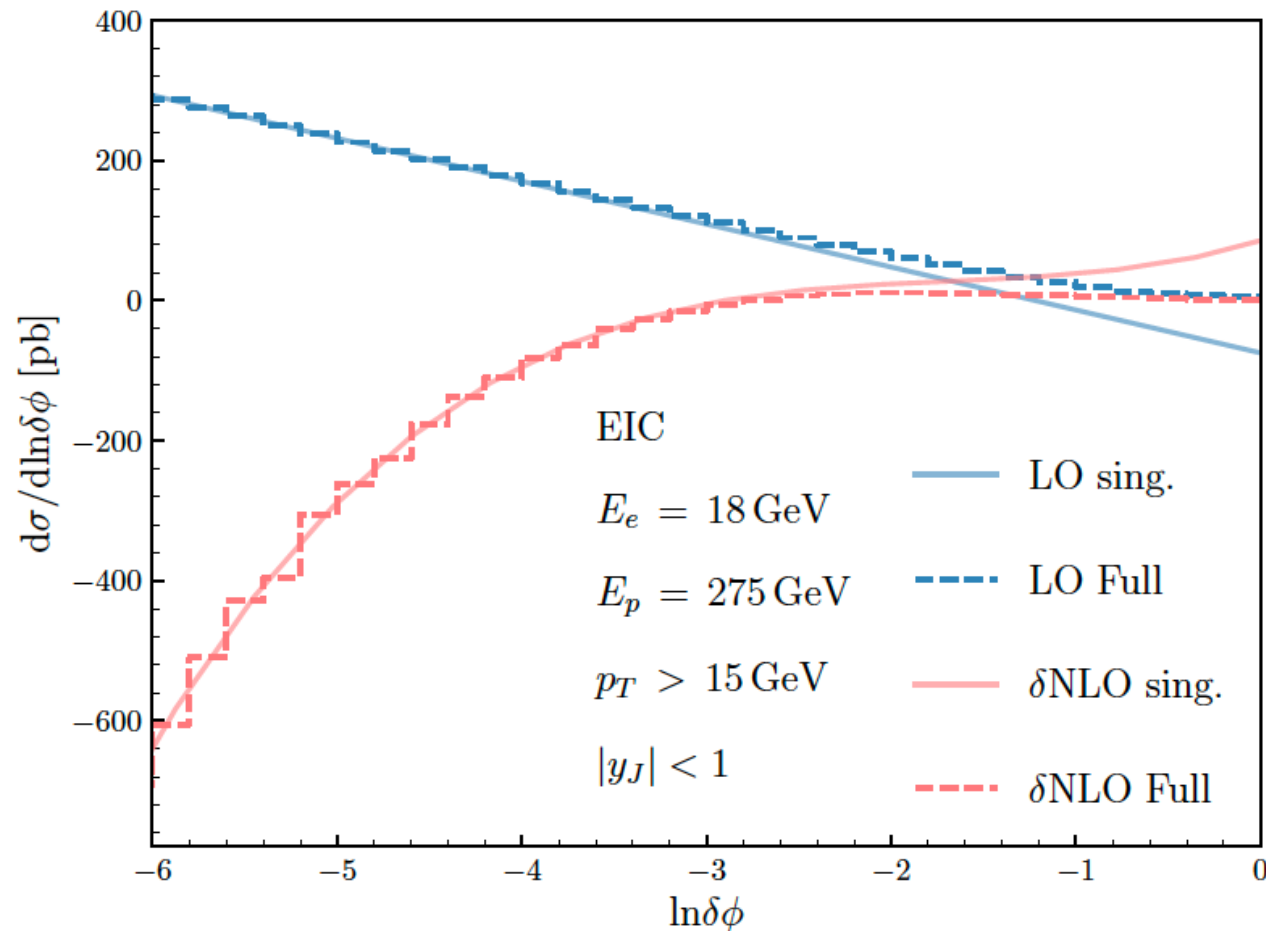
- The uncertainty bands are narrower at N³LL (red) compared to NNLL (blue)
- At N³LL the dominant scale uncertainties are from μ_b variation

$N^3LL + \mathcal{O}(\alpha_s^2)$ predictions on lepton-jet azimuthal correlation

SF, Gao, Li, Shao '24

- In the back-to-back limit ($\delta\phi \rightarrow 0$) the singular contributions in DIS are consistent with the fixed-order results from NLOJET++ up to $\mathcal{O}(\alpha_s^2)$
- In the large $\delta\phi$ region the resummation formula receives significant matching corrections
- It is necessary to switch off the resummation and instead employ fixed-order calculations

$$d\sigma_{\text{add}}(NNNLL + \mathcal{O}(\alpha_s^2)) \equiv d\sigma(NNNLL) + \underbrace{d\sigma(\text{NLO}) - d\sigma(\text{NLO singular})}_{d\sigma(\text{NLO non-singular})}$$



Single spin asymmetry for 1-jettiness

SF, Lin, Shao, Zhou '25

- 1-jettiness:**

$$\tau_1 \equiv \frac{2}{Q^2} \sum_k \min\{q_B \cdot p_k, q_D \cdot p_k\}$$

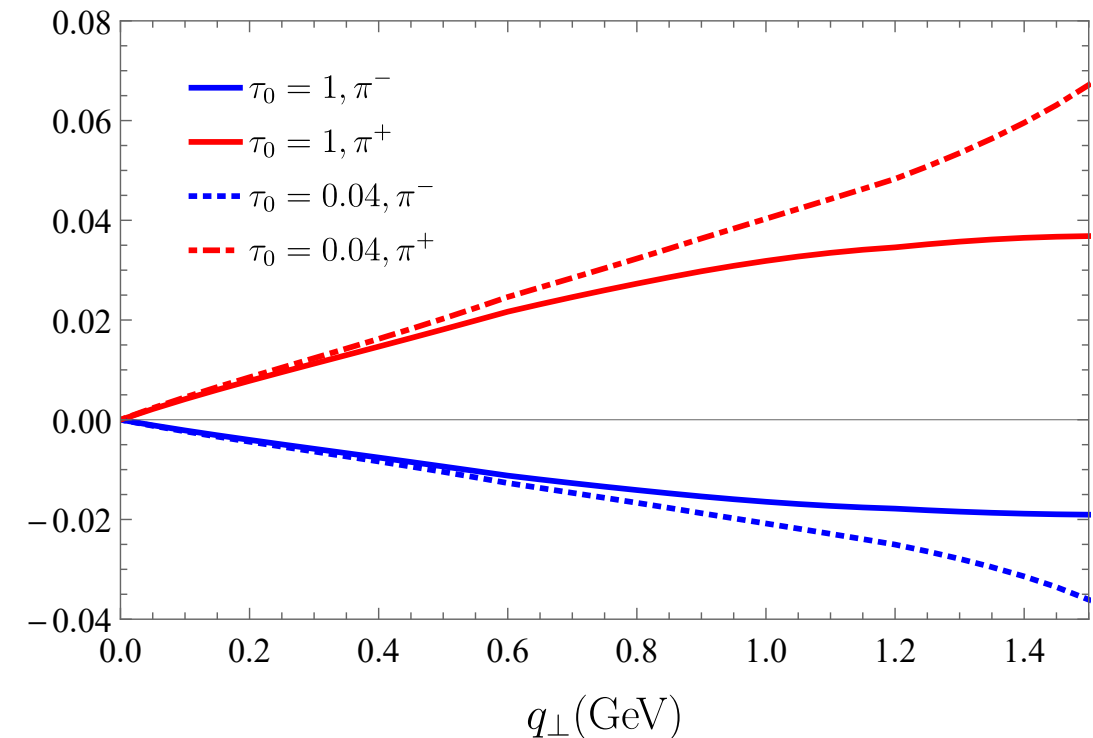
- Single spin asymmetry (SSA) at NLL:**

$$A_N = \frac{F_{UT}}{F_{UU}} = \frac{\int_0^\infty \frac{b^2 db}{4\pi} J_1\left(\frac{bq_\perp}{z_h}\right) \sum_q e_q^2 T_F(x_B, x_B, \mu_b) D_{h/q}(z_h, \mu_b) e^{-S_P(b)}}{\int_0^\infty \frac{b db}{2\pi} J_0\left(\frac{bq_\perp}{z_h}\right) \sum_q e_q^2 f_q(x_B, \mu_b) D_{h/q}(z_h, \mu_b) e^{-S_P(b)}}$$

- Qiu-Sterman function:**

$$T_{F,q}(x, x, \mu) = \mathcal{N}_q(x) f_q(x, \mu)$$

Echevarria, Kang, Terry '14,20



Summary

- We have studied on the lepton-jet correlation in both e-p and e-A collisions. Utilizing SCET, we derived a factorization theorem for back-to-back lepton-jet configurations.
- In e-A collisions, we discussed the utility of our approach in disentangling intrinsic non-perturbative contributions from nTMDs and dynamical medium effects in nuclear environments. We find the process is primarily sensitive to the initial state's broadening effects.
- TMD resummation accuracy has been improved to $N^3\text{LL} + \mathcal{O}(\alpha_s^2)$ accuracy in e-p collisions. It is good to have the measurement at the HERA to make a comparison.
- Our work sets the groundwork for future experiments at the EIC, offering a robust framework for measuring nTMDs and polarized TMDs.

Thank you!