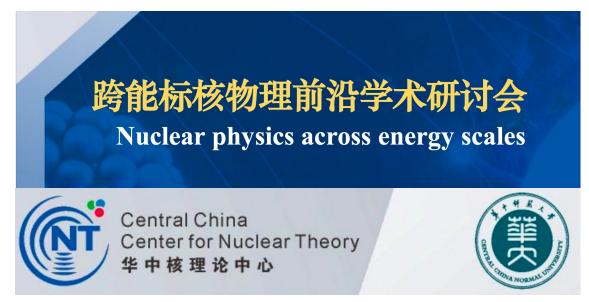
Overview on AMPT and Nuclear Structure

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Outline

- Introduction to A Multi-Phase Transport (AMPT) model
- Selected AMPT results related to nuclear structure
- Outlook
- Summary



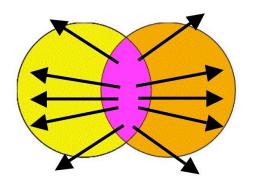


Transport models are suitable for non-equilibrium studies

 Hydro-based hybrid models are now the "standard model" for heavy ion collisions at high energies.
 EoS, η, ζ, ...

For finite/small systems at finite energies,
 or early time & hard probes of large systems at high energies:
 non-equilibrium effects are expected to be important.

 One example is the parton escape mechanism:
 interaction-induced response from kinetic theory
 to the anisotropic spatial geometry alone.



L He et al. PLB 753 (2016); ZWL et al. NPA 956 (2016); HL Li et al. PRC 99 (2019)

Transport models are suitable for non-equilibrium studies

Small system data show similar flow features as large systems:

are they real signals from collectivity? is a parton matter formed in small systems? is the matter close to equilibrium or far off equilibrium (core-corona)? is the matter dominated by gluons or quarks?

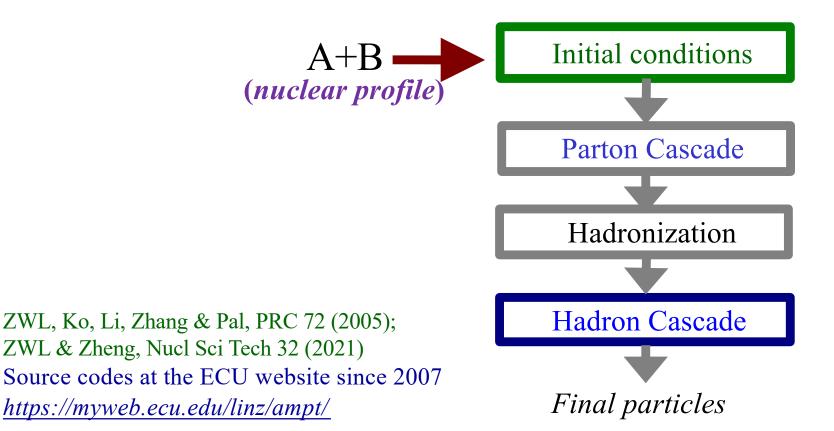
To answer these questions and better study properties of parton matter/QGP, transport models / kinetic theory are crucial as they don't assume equilibrium & can address non-equilibrium dynamics.

B Zhang, CPC 109 (1998); Heiselberg & Levy, PRC 59 (1999); ZWL & Ko, PRC 65 (2002); Z Xu & Greiner, PRC 71 (2005); Bzdak & Ma, PRL 113 (2014); Kurkela et al. PLB 783 (2018) & EPJC 79 (2019); Kurkela, Törnkvist & Zapp, EPJC 84 (2024); ...

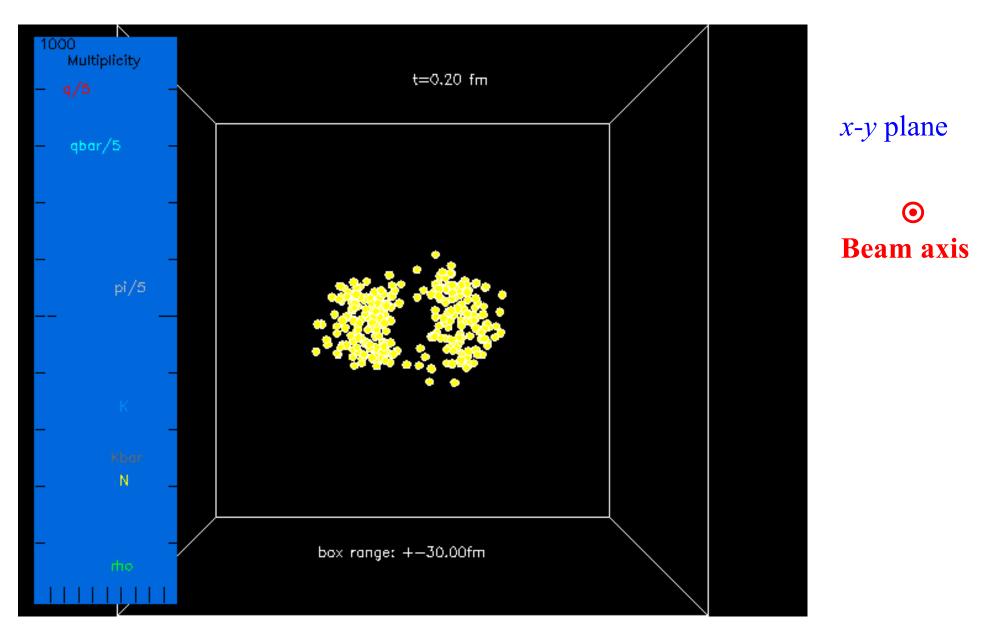
The AMPT model

A multi-phase transport (AMPT) model was constructed as a self-contained kinetic description of nuclear collisions:

- evolves the system from initial condition to final observables
- includes 3D productions of all flavours & conserved charges (B/Q/S/C...)
- includes fluctuating initial conditions & non-equilibrium dynamics/evolution



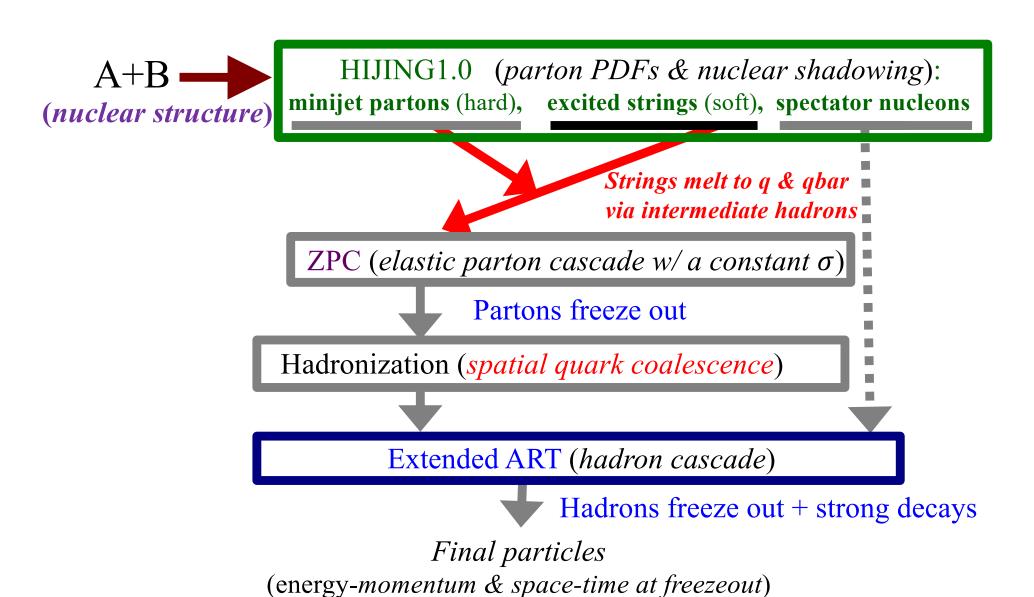
A b=10fm Au+Au event at 200A GeV from string melting AMPT



Initial overlap region has an irregular geometry (will develop into triangular flow v_3)

The AMPT model: string melting version:

applicable when we expect the formation of a parton matter



Main developments of the AMPT model

AMPT-string melting model to describe the bulk matter ZWL, PRC 90 (2014)

New quark coalescence to allow freedom to form M or B

He & ZWL, PRC 96 (2017)

Modern PDFs & b-dependent nuclear shadowing C Zhang, Zheng, Liu, Shi & ZWL, PRC 99 (2019)

Improvements of heavy flavors

Zheng, C Zhang, Shi & ZWL, PRC 101 (2020);
C Zhang, Zheng, Shi & ZWL, PLB 846 (2023)

Local nuclear scaling for self-consistent size dependence C Zhang, Zheng, Shi & ZWL, PRC 104 (2021)

AMPT-HC: pure hadron cascade with mean fields

Yong, Xiao, Gao & ZWL, PLB 820 (2021);
Yong, Li, Xiao & ZWL, PRC 106 (2022)

Include sub-nucleon structure of proton

Zheng, C Zhang, Liu, ZWL, Shou & Yin, EPJC 81 (2021);
XL Zhao, ZWL, Zheng & GL Ma, PLB 839 (2023), ...

Benchmark and improve ZPC parton transport XL Zhao, GL Ma, YG Ma & ZWL, PRC 102 (2020); Mendenhall & ZWL, 2507.23107

Most of these were summarized in a mini-review: ZWL & Zheng, Nucl Sci Tech 32 (2021)

AMPT is also a test-bed of different ideas for the community

- Discovery of triangular flow v₃ Alver & Roland, PRC 81 (2010)
- Early study of flows in small systems Bzdak & Ma, PRL 113 (2014), ...
- Longitudinal decorrelations of flows Pang et al. PRC 91 (2015) & EPJA 52 (2016); ...

Shou, Ma & Ma, PRC 90 (2014);

- Non-equilibrium parton escape for finite systems/energies

 L He et al. PLB 753 (2016);
 ZWL et al. NPA 956 (2016);
 HL Li et al. PRC 99 (2019)
- CME signal & background

 Huang, Ma & Ma, PRC 97 (2018);

 XL Zhao & Ma, PRC 106 (2022);

 Chen, Zhao & Ma, PRC 109 (2024); ...
- Vorticity & polarization observables

 Jiang, ZWL & Liao, PRC 94 (2016);

 Pang et al, PRL 117 (2016);

 H Li et al. PRC 96 (2017); ...
- Effects of nuclear structure

 HJ Xu et al. PRL 121 (2018), CPC 42 (2018);

 HL Li et al. PRL 125 (2020);

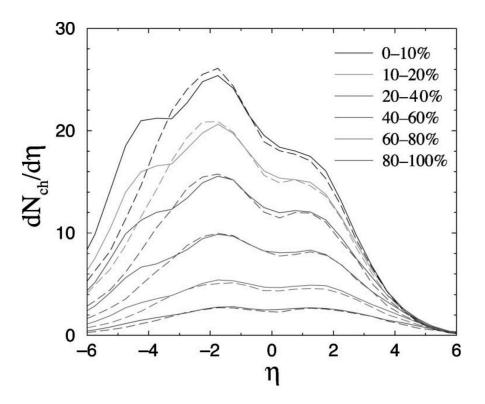
 Jia, Giacalone & Zhang, PRL 131 (2023); ...

The first AMPT study that considered nuclear structure (beyond the Woods-Saxon shape) is on dAu:

ZWL & Ko, PRC 68 (2003)

In this study, we use the Hulthen wave function [29] to model the structure of a deuteron [30]:

$$u(r) = Ce^{-\alpha r}(1 - e^{-\mu r}),$$
 (1)



d+Au at 200A GeV effect of final-state interactions:

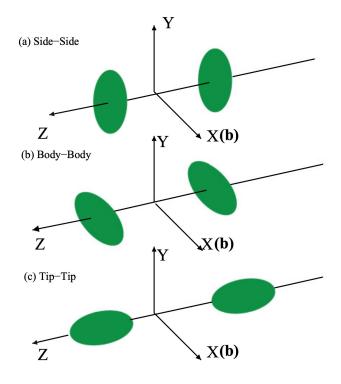
dashed: HIJING

solid: AMPT-default

Nucleon density of ²³⁸U is parametrized as a deformed Woods-Saxon:

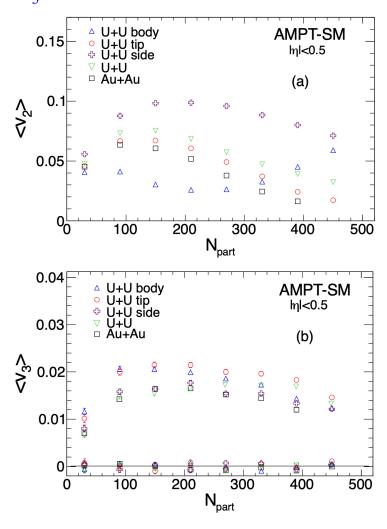
$$\rho = \frac{\rho_0}{1 + \exp([r - R']/a)},$$

$$R' = R[1 + \beta_2 Y_2^0(\theta) + \beta_4 Y_4^0(\theta)]$$



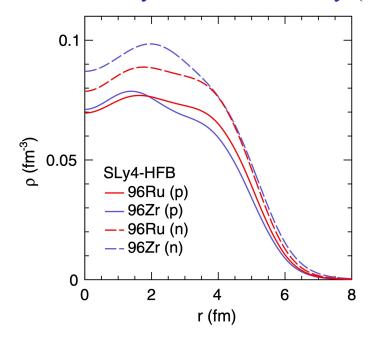
Rihan Haque, ZWL & Mohanty, PRC 85 (2012)

v₂ are very sensitive to orientation; v₃ less so:

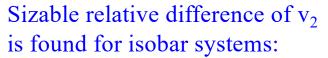


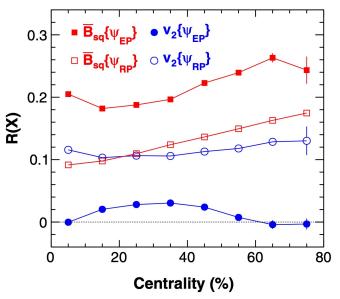
Incorporation of nucleus structure: isobars

For isobar calculations, protons & neutrons are sampled according to radial density distributions from density functional theory (DFT):



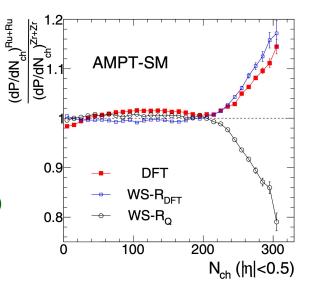
H.J. Xu, Wang, Li, Zhao, ZWL, Shen & Wang, PRL 121 (2018)



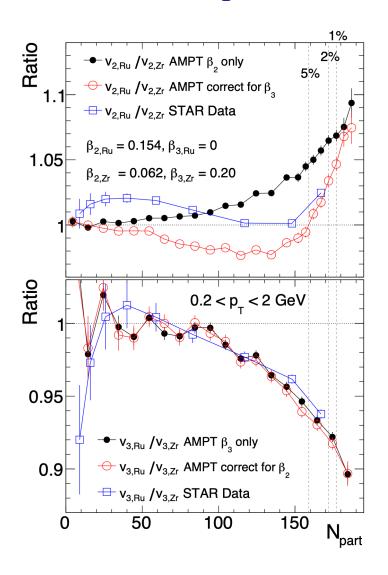


Ratio of N_{ch} distributions probes the nuclear mass radii:

H.L. Li, Xu, Zhao, ZWL, Zhang, Wang, Shen & Wang, PRC 98 (2018)

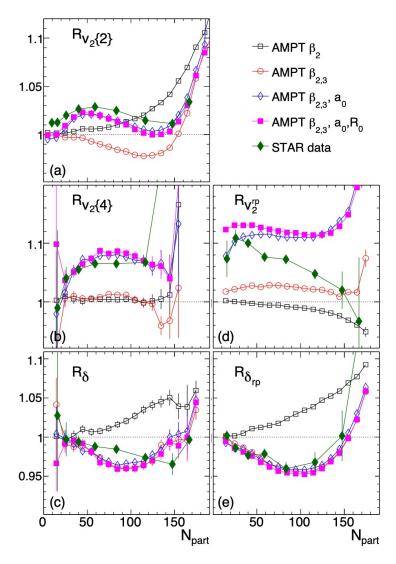


Incorporation of nucleus structure: deformations



Quadrupole and octupole deformations could be extracted from $v_2 \& v_3$

CJ Zhang & Jia, PRL 128 (2022)

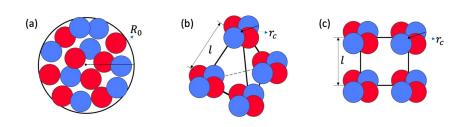


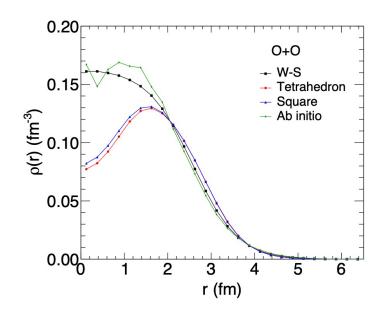
Effects of nuclear skin and nuclear deformations can be disentangled

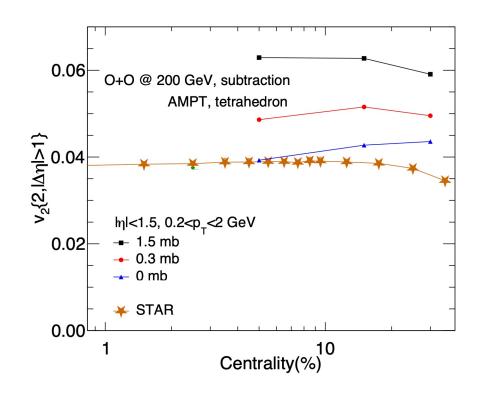
Jia, Giacalone & Zhang, PRL 131 (2023)

Nuclear cluster structures in ¹⁶O:

Xin-Li Zhao, Ma, Zhou, ZWL & Zhang, 2404.09780







A problem:

the model v_2 is higher than STAR data, even at 0 parton cross section σ

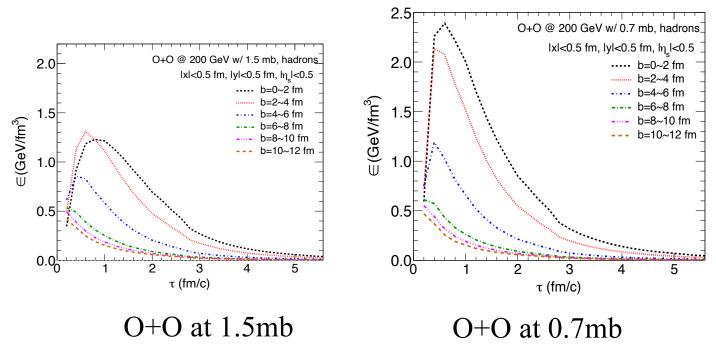
A problem for O+O at 200A GeV:

- → flow from the hadron phase is too big
- → hadron phase starts too early?

Indeed the case:

peak energy density of hadron phase in AMPT-SM

> 1 GeV/fm³ (for the central cell in O+O collisions)



This is because hadronization (quark coalescence) in AMPT-SM starts when parton interactions stop, not at a given energy density.

Qualitative estimate: parton interactions stop when

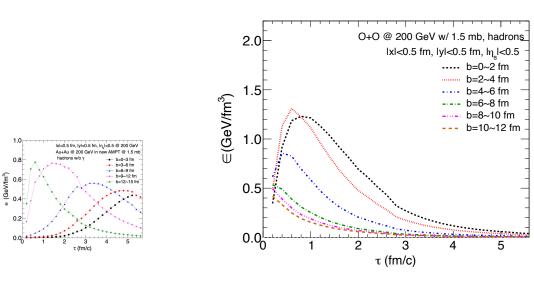
n: parton density;

L: system size.

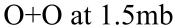
 $\rightarrow n \sim \frac{1}{\sigma L}$ is higher for a smaller system.

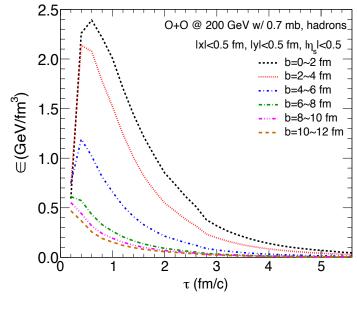
 \rightarrow peak hadron density is higher at $\sigma = 0.7$ mb than at 1.5mb;

and higher in O+O than in Au+Au.



Au+Au at σ =1.5mb

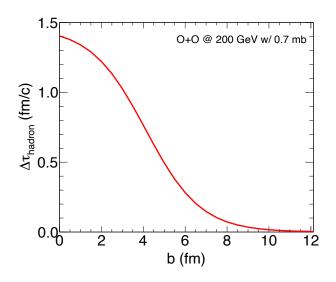


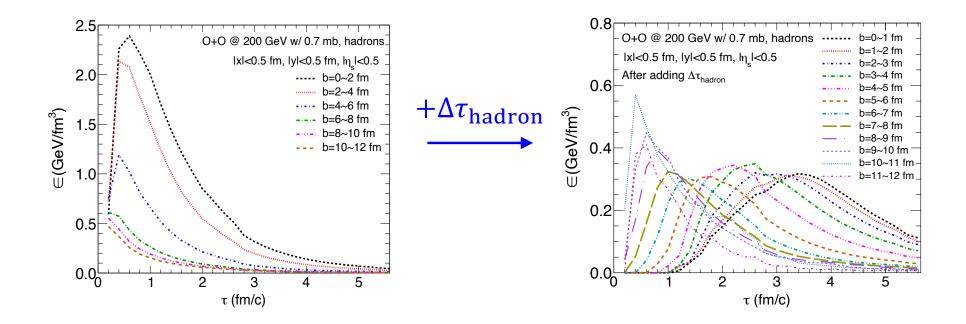


O+O at 0.7mb

Current solution of the problem:

increase hadron proper formation time by $\Delta \tau_{\rm hadron}$ so that the peak hadronic energy density is around a fixed ϵ_c (0.3 GeV/fm³ here)



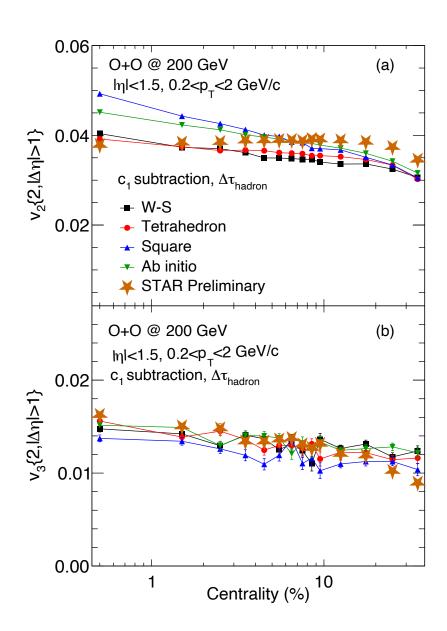


We then calculated flows at RHIC and LHC energies.

O+O at RHIC

Cluster structures in 16 O affect flow, especially v_2

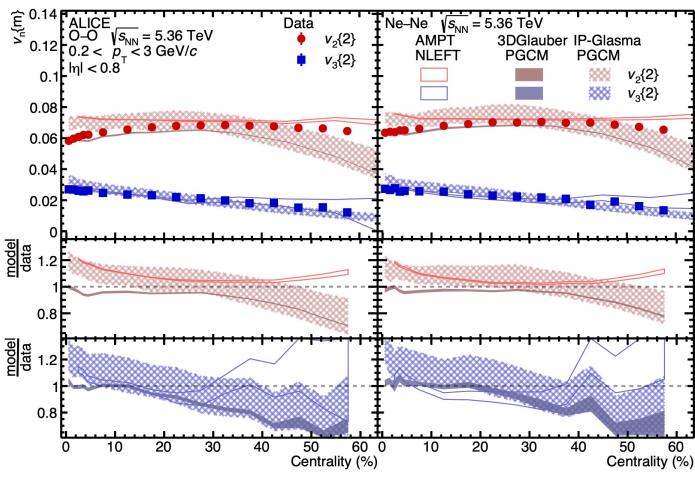
Xin-Li Zhao, Ma, Zhou, ZWL & Zhang, 2404.09780v2



v₂ and v₃ at LHC

O+O

Ne+Ne



AMPT results:

ALICE 2509.06428v1

overall magnitudes are reasonable;

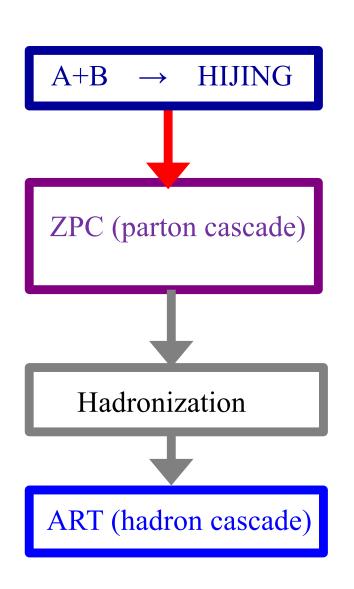
v₂ is too flat versus centrality.

Note: the AMPT band only includes statistical error bars

Outlook

Further developments are needed, with many possible directions:

* = Done



- * Modern PDF & nuclear shadowing
- * Size dependence of initial condition

Nonzero nuclear thickness for lower energies Gluons in initial condition for high energies

Parton cross sections from finite-T pQCD 2 ←→ 3 or 2 ←→ 1 parton interactions QCD equation of state & CEP E&M fields Spin/angular momentum transport

- * Improved quark coalescence
 Hadronization at phase transition density
 Gluons in hadronization
 Include fragmentation (* for heavy flavor)
- * Hadron mean fields
 More resonances

Outlook

ZPC (parton cascade)

Hadronization

Parton cross sections from finite-T pQCD

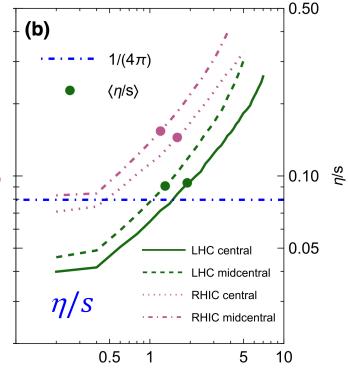
Hadronization at phase transition density

Currently, AMPT treats the parton σ & screening mass μ as constant, independent of density or temperature T:

$$\frac{d\sigma}{dt} \sim \frac{9\pi\alpha_s^2}{2(t-\mu^2)^2}$$

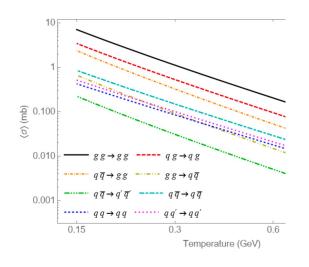
We have calculated parton η/s in AMPT-SM: average value is very small $\sim (1-3)/(4\pi)$, but has T- & t-dependences opposite to pQCD

MacKay & ZWL, EPJC 82 (2022)

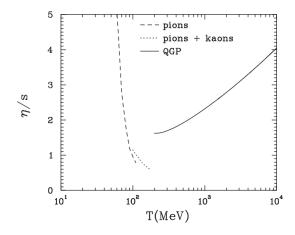


Outlook

Finite-T pQCD $\rightarrow \mu \propto gT$ $\rightarrow \sigma \propto g^4/\mu^2 \propto g^2/T^2$ will be larger at lower T



Ohanaka & ZWL, in preparation; Ross & ZWL, ongoing



$$\rightarrow \eta/s \propto \frac{1}{T^2\sigma} \propto \frac{1}{g^2}$$

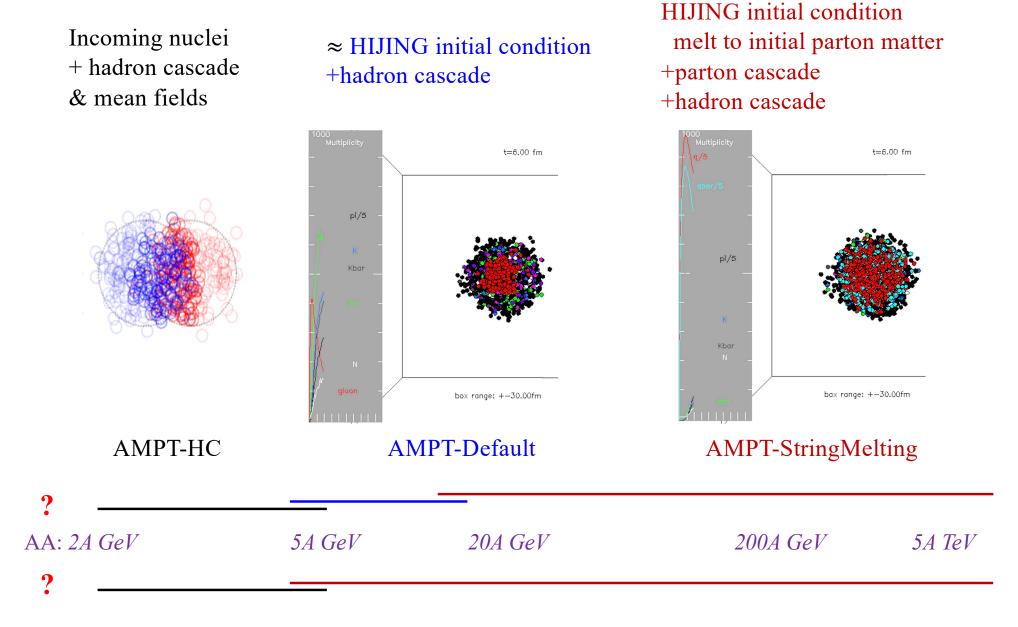
will be \sim the lowest around T_c

Csernai, Kapusta & McLerran, PRL 97 (2006)

Along this direction, we are developing ZPC/AMPT as a dynamical model of finite-T kinetic theory:

- → parton-to-hadron phase transition according to local energy density
- \rightarrow T- and t-dependences of η/s consistent with pQCD
- \rightarrow extract η/s from the non-equilibrium transport approach

The AMPT model across energy scales



Summary

The AMPT model aims to provide a self-contained kinetic description of heavy ion collisions; especially suitable for studies of non-equilibrium dynamics.

Recent improvements have made AMPT more accurate and consistent.

Nuclear structures incorporated into the AMPT initial condition have shown direct effects on certain observables; this opens up new ways to study nuclear structure & origin of collectivity.

Further developments are necessary, including developing ZPC/AMPT into a model for finite-T pQCD kinetic theory.

