

Small scale fluctuations and flow correlations

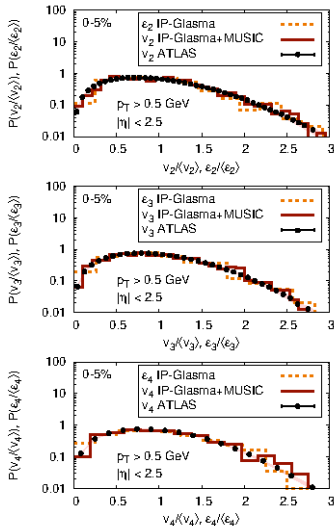
Piotr Bożek

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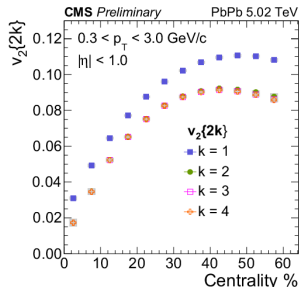
R. Samanta, PB, PRC 109 (2024) 064910



Distribution of flow harmonics



$$v_2\{2\} > v_2\{4\} \simeq v_2\{6\} \simeq v_2\{8\}$$



Higher order cumulants in Fig. 2 behave statistically

CMS Nucl. Phys. A937 (2017) 401

$$V_n = \kappa_n \epsilon_n$$

$$\epsilon_n = \int r^n e^{in\phi} \rho(x, y) d^2r$$

Gale, Jeon, Tribedy, Venugopalan arXiv: 1210.5144

Mapping the flow fluctuations

$$\frac{dN}{d\phi dp d\eta} \propto 1 + 2V_2(p, \eta)e^{-i2\phi} + 2V_3(p, \eta)e^{-i3\phi} + \dots$$

Can we map the flow $V_n(p, \phi)$?

No !

Mapping the flow fluctuations

$$\frac{dN}{d\phi dp d\eta} \propto 1 + 2V_2(p, \eta)e^{-i2\phi} + 2V_3(p, \eta)e^{-i3\phi} + \dots$$

Can we map the flow $V_n(p, \phi)$?

No !

But we could map the covariance

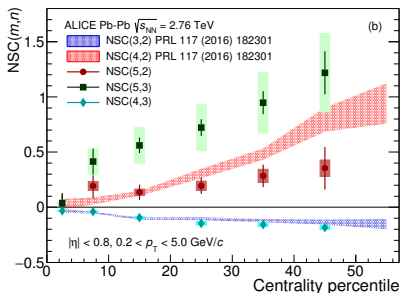
$$\langle V_n(p_1, \eta_1) V_n^*(p_1, \eta_2) \rangle$$

or the correlation

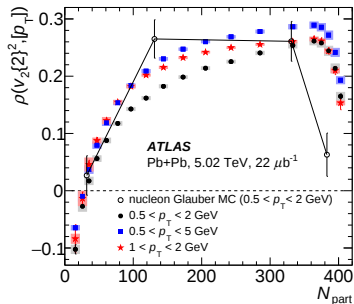
$$\frac{\langle V_n(p_1, \eta_1) V_n^*(p_2, \eta_2) \rangle}{\sqrt{\langle V_n(p_1, \eta_1) V_n^*(p_1, \eta_1) \rangle \langle V_n(p_2, \eta_2) V_n^*(p_2, \eta_2) \rangle}}$$

Fluctuations - covariances ...

► mixed cumulants/covariances



$$\frac{\langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle}{\langle v_n^2 \rangle \langle v_m^2 \rangle}$$

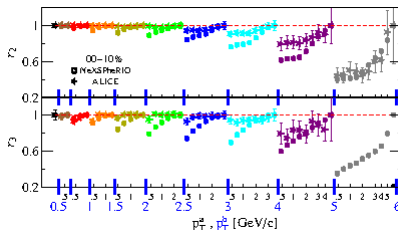


$$\rho(p_T, v_2^2) = \frac{\text{Cov}([p_T], v_n^2)}{\sigma([p_T])\sigma(v_n^2)}$$

Flow decorrelation in p_T

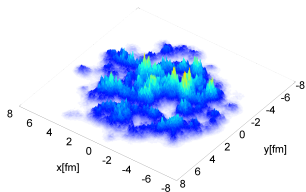
factorization breaking coefficient

$$r_n(p_1, p_2) = \frac{\langle V_n(p_1) V_n^*(p_2) \rangle}{\sqrt{\langle V_n(p_1) V_n^*(p_1) \rangle \langle V_n(p_2) V_n^*(p_2) \rangle}}$$



Gardim, Grassi, Luzum, Ollitrault arXiv: 1211.0989

lumpy structure of the initial density



Schenke, Tribedy, Venugopalan arXiv: 1206.6805

$$r_n(p_1, p_2) \leq 1$$

correlation coefficient

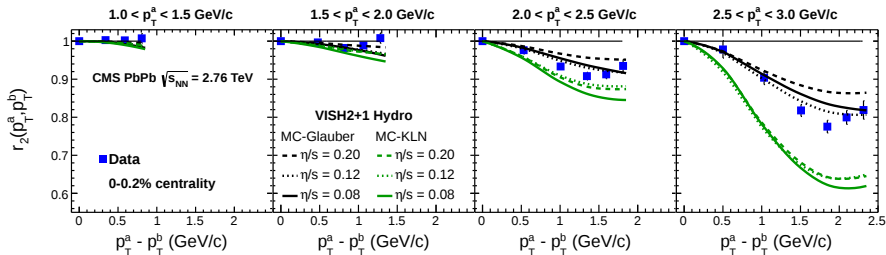
$$r_n(p_1, p_2) = \rho(V_n(p_1), V_n(p_2))$$

other modes in initial density

A. Mazeliauskas, D. Teaney PRC (2015)

$$\epsilon_n \propto \int e^{in\phi} r^n \rho d^2r, \quad \epsilon_{n,k} \propto \int e^{in\phi} r^k \rho \rho d^2r$$

Experiment vs theory

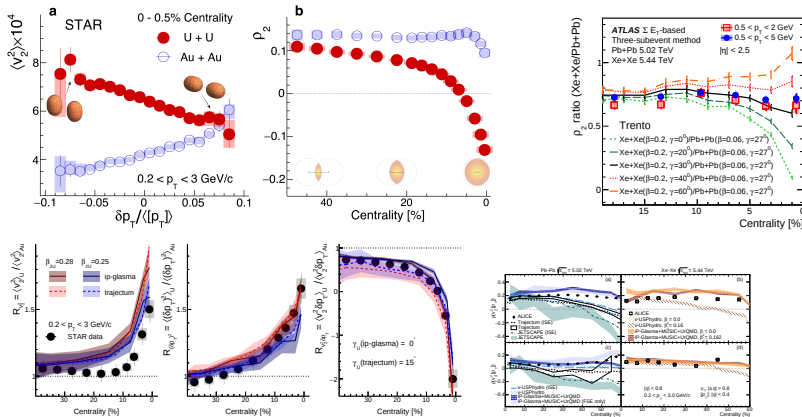


CMS arXiv: 1503.01692

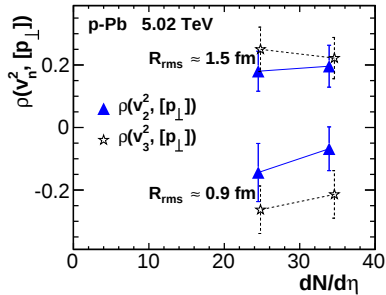
- qualitatively described by models
- can provide constraints on initial fluctuations

$\rho([p_T], v_n^2)$ correlation coefficient

central collisions, sensitive to nuclear deformation: G. Giacalone PRC 2020



Small scales



ρ sensitive sensitive to small scale flc. in p+Pb

Differential measurements

average flow

- $\langle V_n V_n^* \rangle$
- higher cumulants $v_n\{4\}$, $v_n\{6\}$
- scaled cumulants $\langle V_n \dots V_m^* \dots \rangle$
- $\langle \delta p_T v_n^2 \rangle$, $\langle \delta p_T c_n\{4\} \rangle$

+ combination + normalizations

differential flow

- $\langle V_n V_n^*(p) \rangle$ in $v_n(p)$
- $\langle V_n^2 V_n^* V_n^*(p) \rangle$ in $v_n\{4\}(p)$
- factorization breaking, PCA $\langle V_n(p_1) V_n^*(p_2) \rangle$, $\langle V_n^2 V_n^2(p) \rangle$
- + combinations + normalizations

IMPORTANCE

$v_n(p)$ model vs data

PID $v_n(p)$ (quark scaling, intermediate and high p_T , viscosity effects, heavy quarks)

factorization breaking $r_n(p_1, p_2)$ (differential map of fluctuations)

there is more in HI than in $v_v = \kappa_n \epsilon_n$

If you measure higher moments, measure covariances

unknown multidim. probability distribution $P([p_T], V_2, V_3, \dots?)$

measure as many (mixed) moments (cumulants) as possible

e.g. $\langle \delta p_T^2 \rangle$, $\langle \delta p_T^3 \rangle$, $\langle \delta p_T^4 \rangle$, $\rho([p_T], v_n^2)$, $\rho([p_T], v_2^2, v_3^2)$

differential

If we measure

$$\langle \delta p_T^2 \rangle, \langle \delta p_T^3 \rangle, \langle \delta p_T^4 \rangle$$

$$\langle V_n V_n(p)^* \rangle, \langle v_n(p)^2 \rangle, \langle V_n^2 V_n(p)^2 \rangle, \langle V_n^2 V_n^*(p)^2 \rangle$$

why not

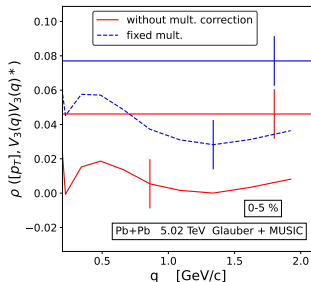
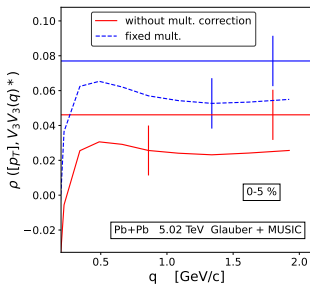
$\text{Cov}(\delta p_T, V_n V_n^*(p))$ or $\text{Cov}(\delta p_T, v_n(p)^2)$

Differential p_T - $v_n(q)$ correlation coefficient

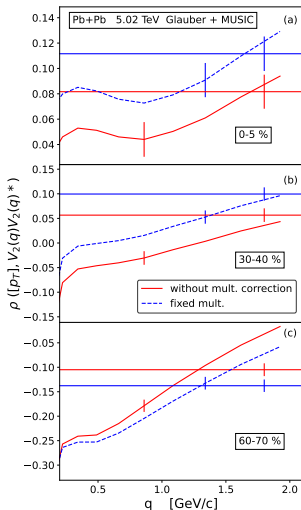
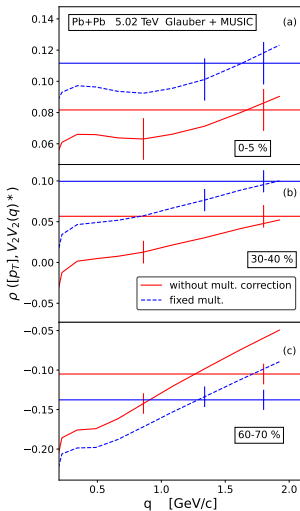
$$\rho(V_n V_n(q)^*, p_T) = \frac{\text{Cov}([p_T], V_n V_n(q)^*)}{\sqrt{\text{Var}([p_T]) \text{Var}(V_n V_n(q)^*)}}$$

$$\rho(V_n(q) V_n(q)^*, p_T) = \frac{\text{Cov}([p_T], V_n(q) V_n(q)^*)}{\sqrt{\text{Var}([p_T]) \text{Var}(V_n(q) V_n(q)^*)}}$$

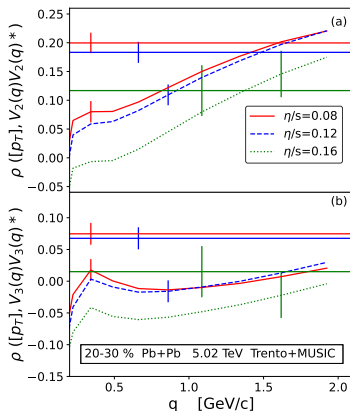
Note: other normalizations possible



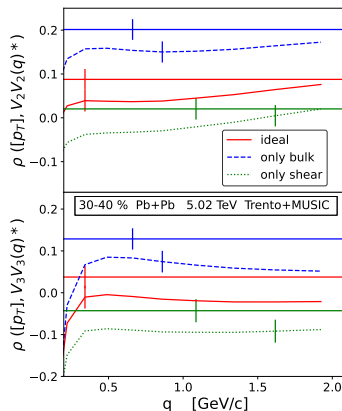
Differential p_T - $v_n(q)$ correlation coefficient



Viscosity dependence



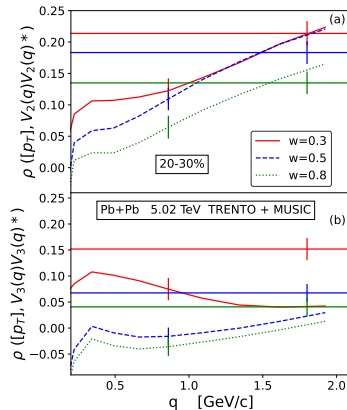
shear viscosity - shift, same shape



bulk viscosity - shift same shape

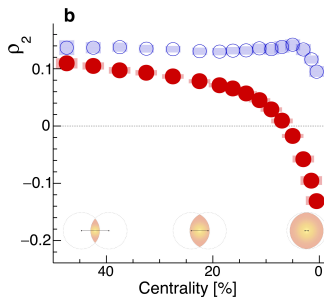
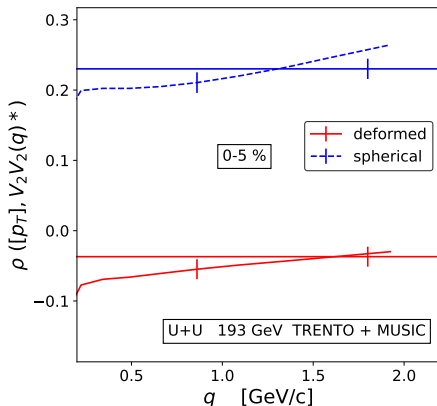
Change in $\rho([p_T], V_n V_n^*)$ same q -dependence of $\rho([p_T], V_n V_n(q)^*)$

granularity in the initial state



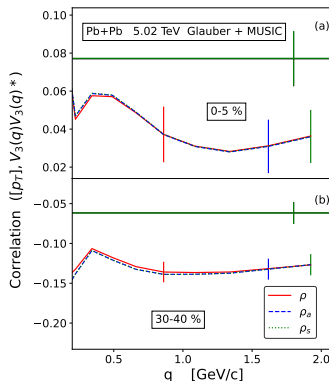
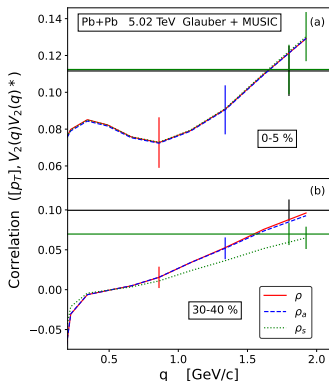
$\rho([p_T], V_n(q)V_n(q)^*)$ sensitive to granularity in the initial state

Deformation



deformation - similar to usual correlation $\rho(v_2^2, p_T)$

simplified expressions - experimentally accessible



$$\rho_a([p_T], V_n(q)V_n(q)^*) = \frac{\text{Cov}([p_T], V_n(q)V_n(q)^*) \langle v_n^2 \rangle}{\sqrt{\text{Var}([p_T]) \text{Var}(v_n^2)} \langle V_n(q)V_n(q)^* \rangle}$$

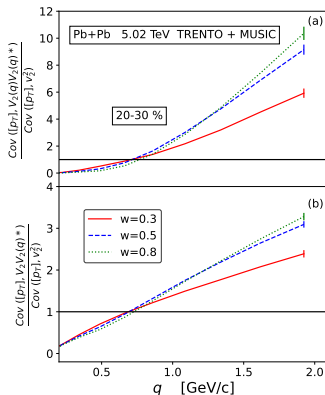
$\text{Var}(v_n^2)$ instead of $\text{Var}(V_n(q)^2)$

Covariance

$$\frac{\text{Cov}([p_T], V_n(q) V_n(q)^*)}{\text{Cov}([p_T], v_n^2)}$$

accessible experimentally

- q dependence sensitive to granularity



Results

- ▶ differential v_n - p_T correlators
- ▶ correlation coefficient $\rho_s = \frac{\text{Cov}([p_T], V_n(q) V_n(q)^*) \langle v_n^2 \rangle}{\sqrt{\text{Var}([p_T]) \text{Var}(v_n^2)} \langle V_n(q) V_n(q)^* \rangle}$
- ▶ scaled covariance $\frac{\text{Cov}([p_T], v_n(q)^2)}{\text{Cov}([p_T], v_n^2)}$
- ▶ q -dependence sensitive to small scale flc. , ...

Outlooks

- ▶ subnucleonic d.o.f., small systems, clustering
- ▶ PID $\rho_s([p_T], V_n(q) V_n(q)^*)$ and $\rho_s([p_T], V_n V_n(q)^*)$
- ▶ mapping of p_T - V_n fluctuations
 $\langle V_n(q_1) V_n^*(q_2) \rangle$, $\langle \delta p_T V_n(q_1) V_n^*(q_2) \rangle$, $\langle (\delta p_T)^2 V_n(q_1) V_n^*(q_2) \rangle$, ...
different radial modes