

# Determination of the Strong Coupling Constant from Inclusive Semi-leptonic B Meson Decays

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arXiv:2412.02480,

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The 2025 International Workshop on the High Energy Circular  
Electron Positron Collider

# The strong coupling constant, $\alpha_s$

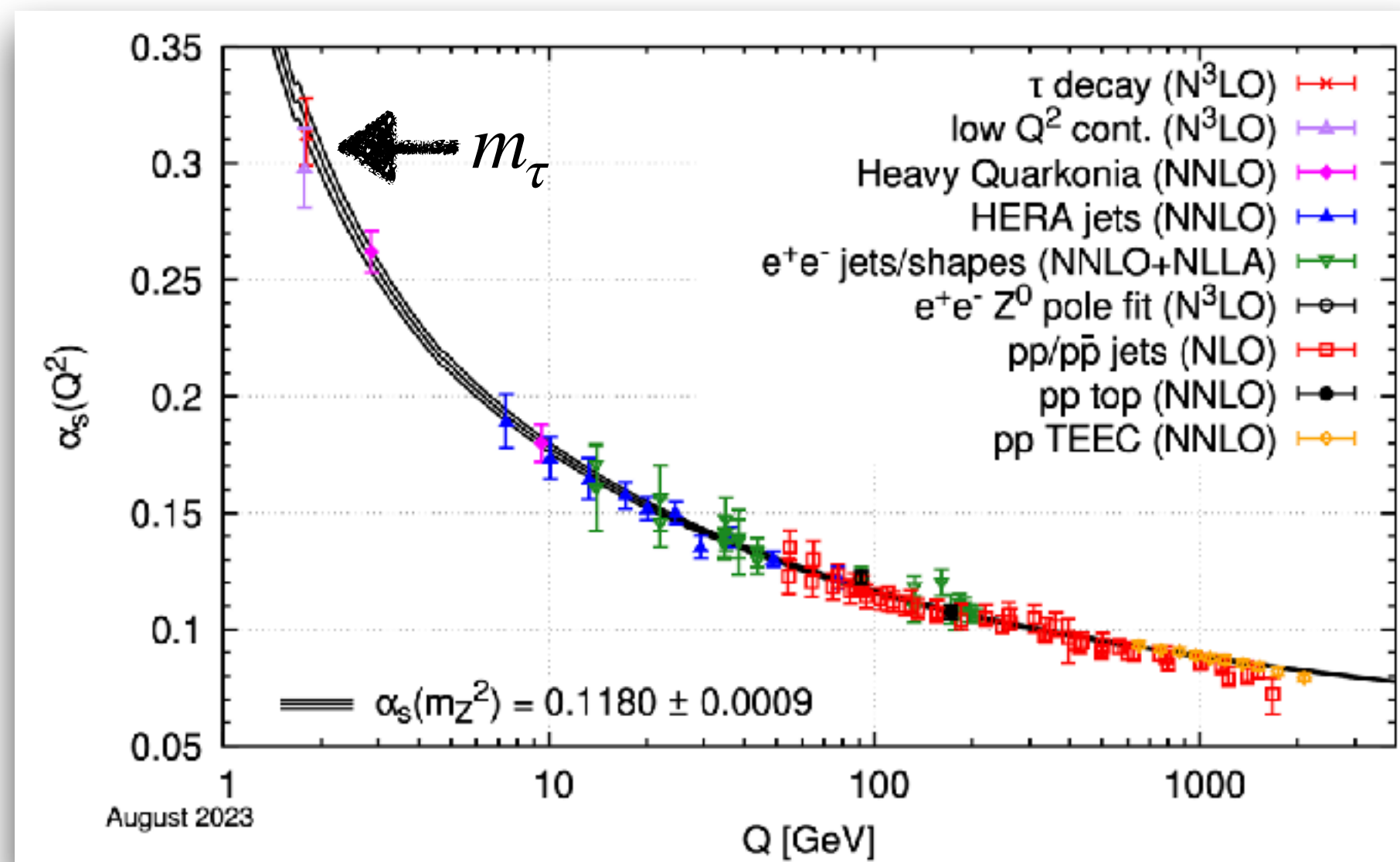
- Strong coupling constant  $\alpha_s$ 
  - encodes the strength of strong interaction,
  - is a fundamental parameter in the Standard Model (SM) and the Quantum Chromodynamics (QCD).
- However, the  $\alpha_s$  is less understood compared with the other couplings.

$\alpha$	$7.2973525693(11) \times 10^{-3}$	$1.5 \times 10^{-10}$
$G_F/(\hbar c)^3$	$1.1663788(6) \times 10^{-5} \text{ GeV}^{-2}$	$5.1 \times 10^{-7}$
$m_Z$	$91.1880(20) \text{ GeV}/c^2$	$2.2 \times 10^{-5}$
$m_W$	$80.3692(133) \text{ GeV}/c^2$	$1.7 \times 10^{-4}$
$\alpha_s(m_Z^2)$	$0.1180(9)$	$7.6 \times 10^{-3}$
$m_H$	$125.20(11) \text{ GeV}/c^2$	$8.7 \times 10^{-4}$
$m_e$	$0.51099895000(15) \text{ MeV}/c^2$	$2.9 \times 10^{-10}$
$m_\mu$	$0.1056583755(23) \text{ GeV}/c^2$	$2.1 \times 10^{-8}$
$m_\tau$	$1.77693(9) \text{ GeV}/c^2$	$5.1 \times 10^{-5}$
$\overline{m}_u(2 \text{ GeV})$	$2.16(7) \text{ MeV}/c^2$	$3.2 \times 10^{-2}$
$\overline{m}_d(2 \text{ GeV})$	$4.70(7) \text{ MeV}/c^2$	$1.5 \times 10^{-2}$
$\overline{m}_s(2 \text{ GeV})$	$93.5(8) \text{ MeV}/c^2$	$8.6 \times 10^{-3}$
$\overline{m}_c(\overline{m}_c)$	$1.2730(46) \text{ GeV}/c^2$	$3.6 \times 10^{-3}$
$\overline{m}_b(\overline{m}_b)$	$4.183(7) \text{ GeV}/c^2$	$1.7 \times 10^{-3}$
$m_t$	$172.57(29) \text{ GeV}/c^2$	$1.7 \times 10^{-3}$

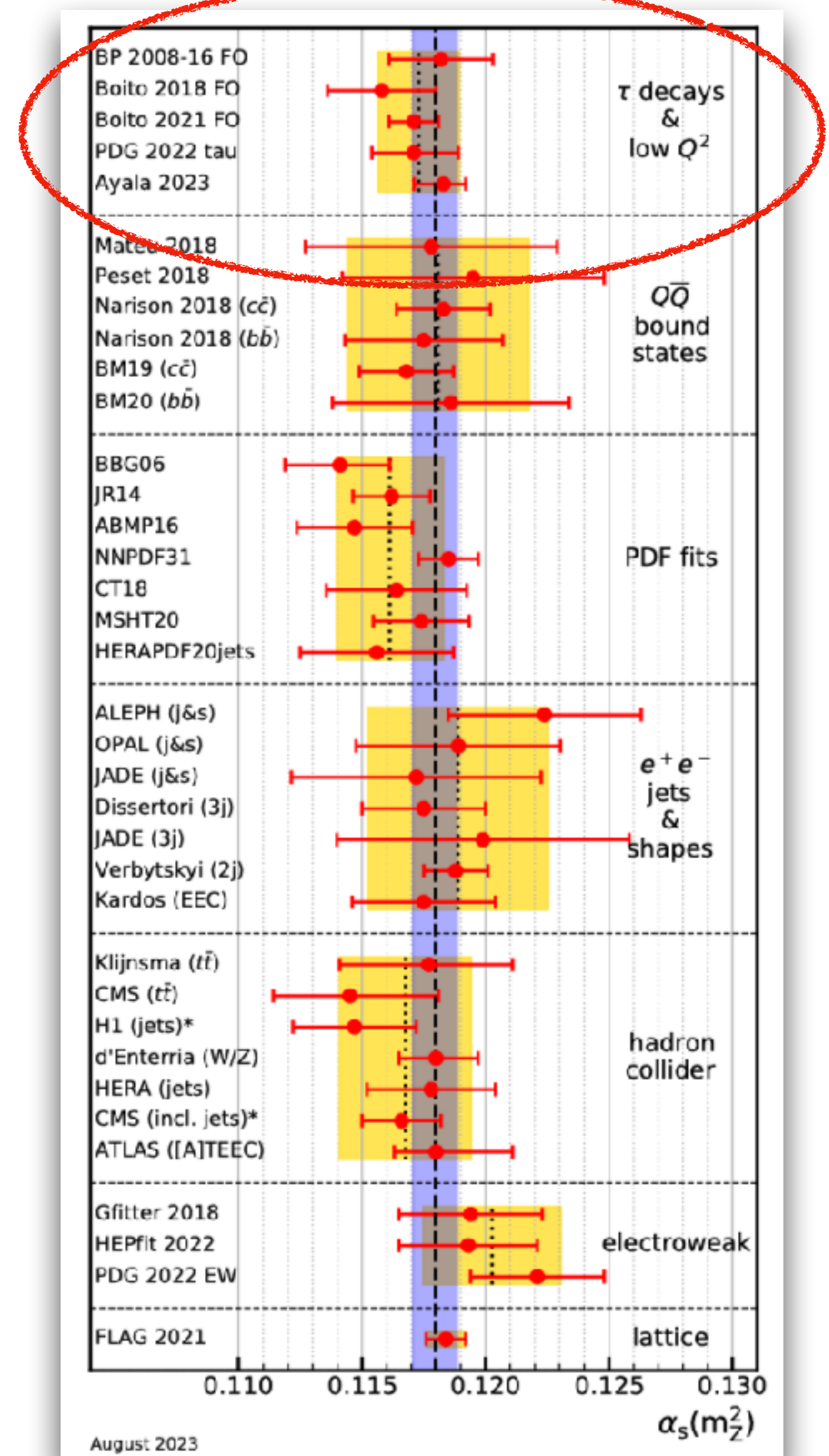


# The strong coupling constant, $\alpha_s$

- The value of  $\alpha_s$  decreases with the increasing energy scale  $\mu$ , described by the Renormalization Group Equation (RGE).
- This running behavior corresponds to fundamental properties of asymptotic freedom and color confinement.
  - Measuring  $\alpha_s(\mu)$  over a wide range of energy scale is crucial for understanding and testing QCD.**
- An uncertainty  $\delta$  on a measurement of  $\alpha_s(\mu)$ , at a scale  $\mu$ , translates to an uncertainty  $\delta' = (\alpha_s^2(m_Z)/\alpha_s^2(\mu)) \cdot \delta$ .
  - Determining  $\alpha_s(\mu)$  in the low energy scale provides precise  $\alpha_s(m_Z)$ .**



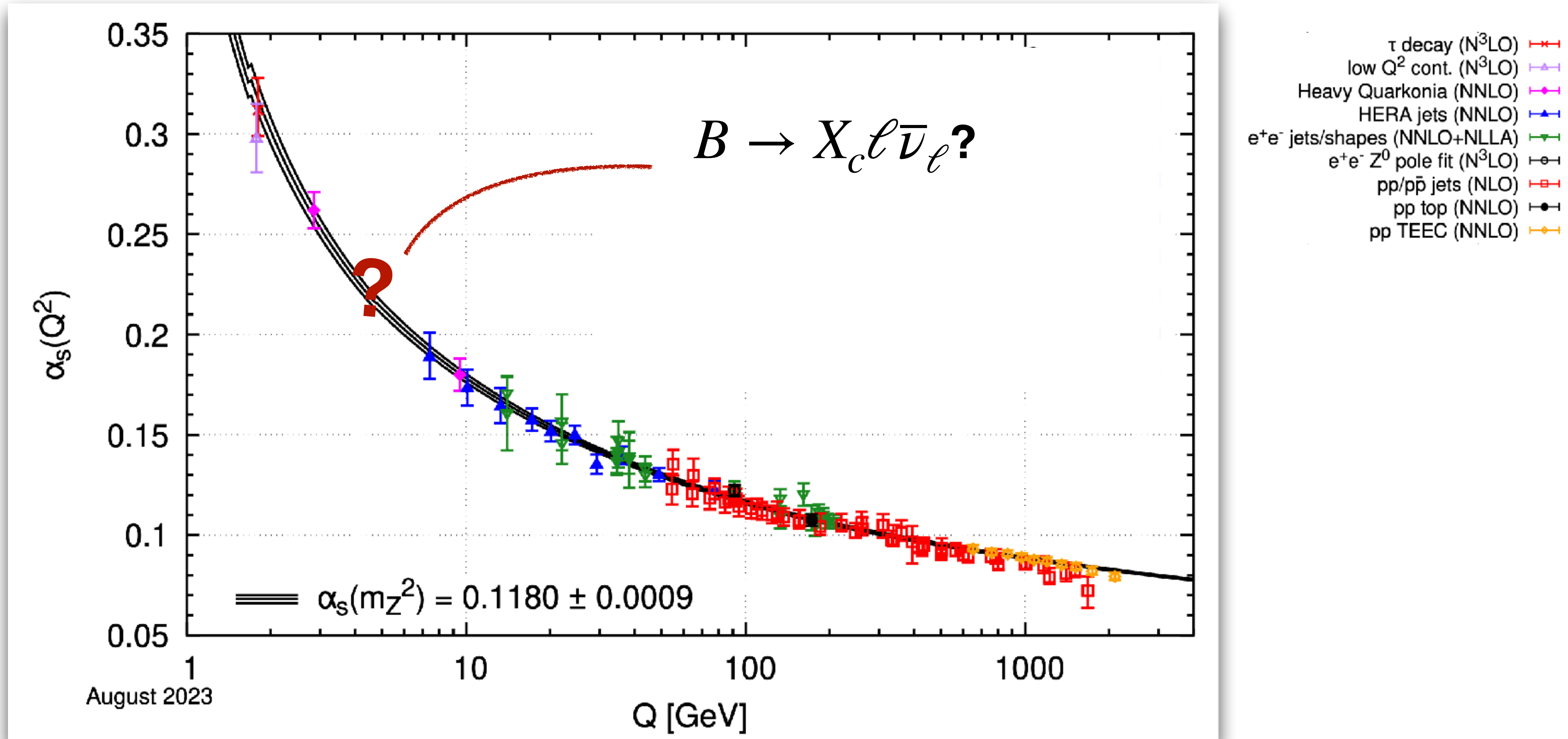
Phys.Rev.D 110 (2024) 3, 030001



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# The strong coupling constant, $\alpha_s$

Is there alternative way for  $\alpha_s$  determination?



# Heavy quark expansion for $B \rightarrow X_c \ell \bar{\nu}_\ell$

In the **Heavy Quark Expansion (HQE)** framework, the inclusive semileptonic  $B$  decay widths can be expressed as double expansions in  $\alpha_s$  and quark masses ( $m_b$ ):

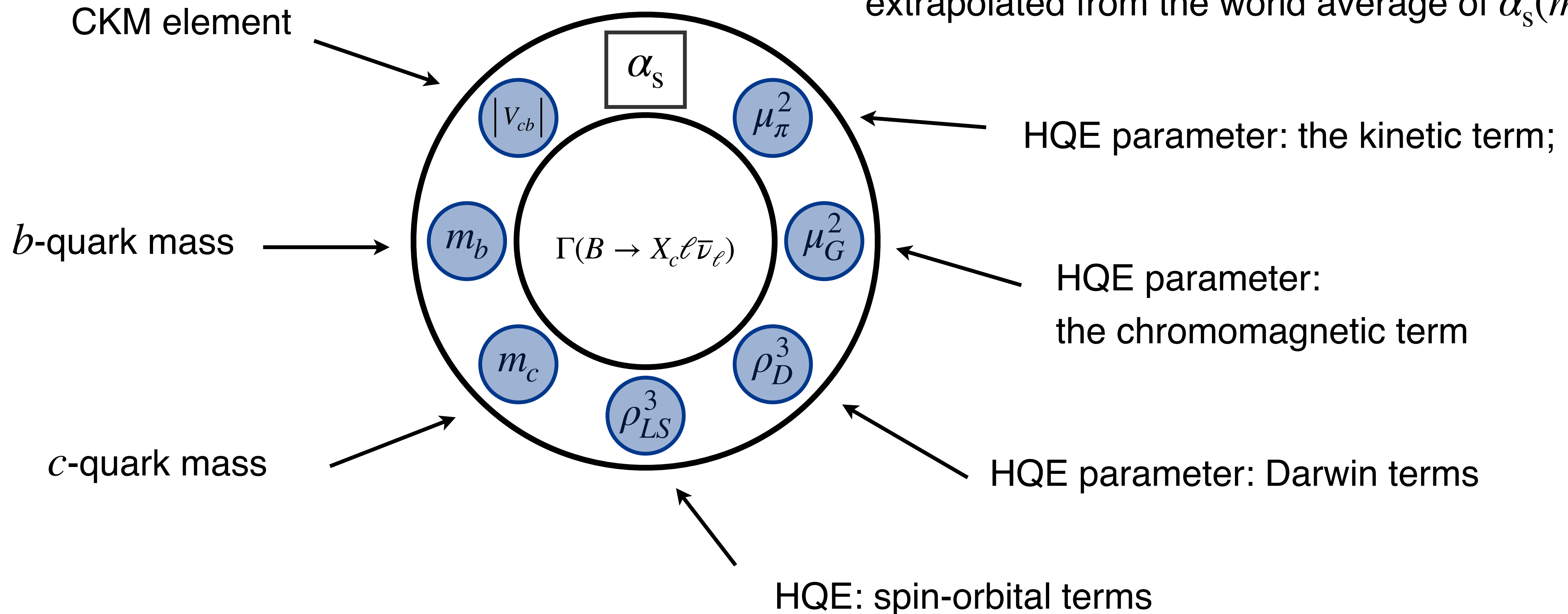
$$\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cb}|^2 m_b^5 A_{\text{ew}}}{192\pi^3} \left[ \begin{aligned} &\mathbf{c}_0 + \mathbf{c}_1 \frac{\alpha_s}{\pi} + \mathbf{c}_2 \left(\frac{\alpha_s}{\pi}\right)^2 + \mathbf{c}_3 \left(\frac{\alpha_s}{\pi}\right)^3 + \mathcal{O}(\alpha_s^4) && \text{..... leading-order power correction} \\ &+ C_{\mu_\pi} \frac{\mu_\pi^2}{2m_b^2} + C_{\mu_G} \frac{\mu_G^2}{2m_b^2} && \text{..... 2nd-order power correction} \\ &+ C_{\rho_D} \frac{\rho_D^3}{2m_b^3} + C_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_b^3} && \text{..... 3rd-order power correction} \\ &+ \dots \end{aligned} \right] \quad \text{..... higher-order power corrections}$$



# Heavy quark expansion for $B \rightarrow X_c \ell \bar{\nu}_\ell$

$$\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cb}|^2 m_b^5 A_{\text{ew}}}{192\pi^3} \left[ \sum_{i=0,1,2,\dots} c_i \left( \frac{\alpha_s}{\pi} \right)^i + C_{\mu_\pi} \frac{\mu_\pi^2}{2m_b^2} + C_{\mu_G} \frac{\mu_G^2}{2m_b^2} + C_{\rho_D} \frac{\rho_D^3}{2m_b^3} + C_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_b^3} + \dots \right]$$

The  $B \rightarrow X_c \ell \bar{\nu}_\ell$  process was used to determine  $|V_{cb}|$ ,  $m_b$  and  $m_c$ , with  $\alpha_s$  **fixed** at the value extrapolated from the world average of  $\alpha_s(m_Z)$



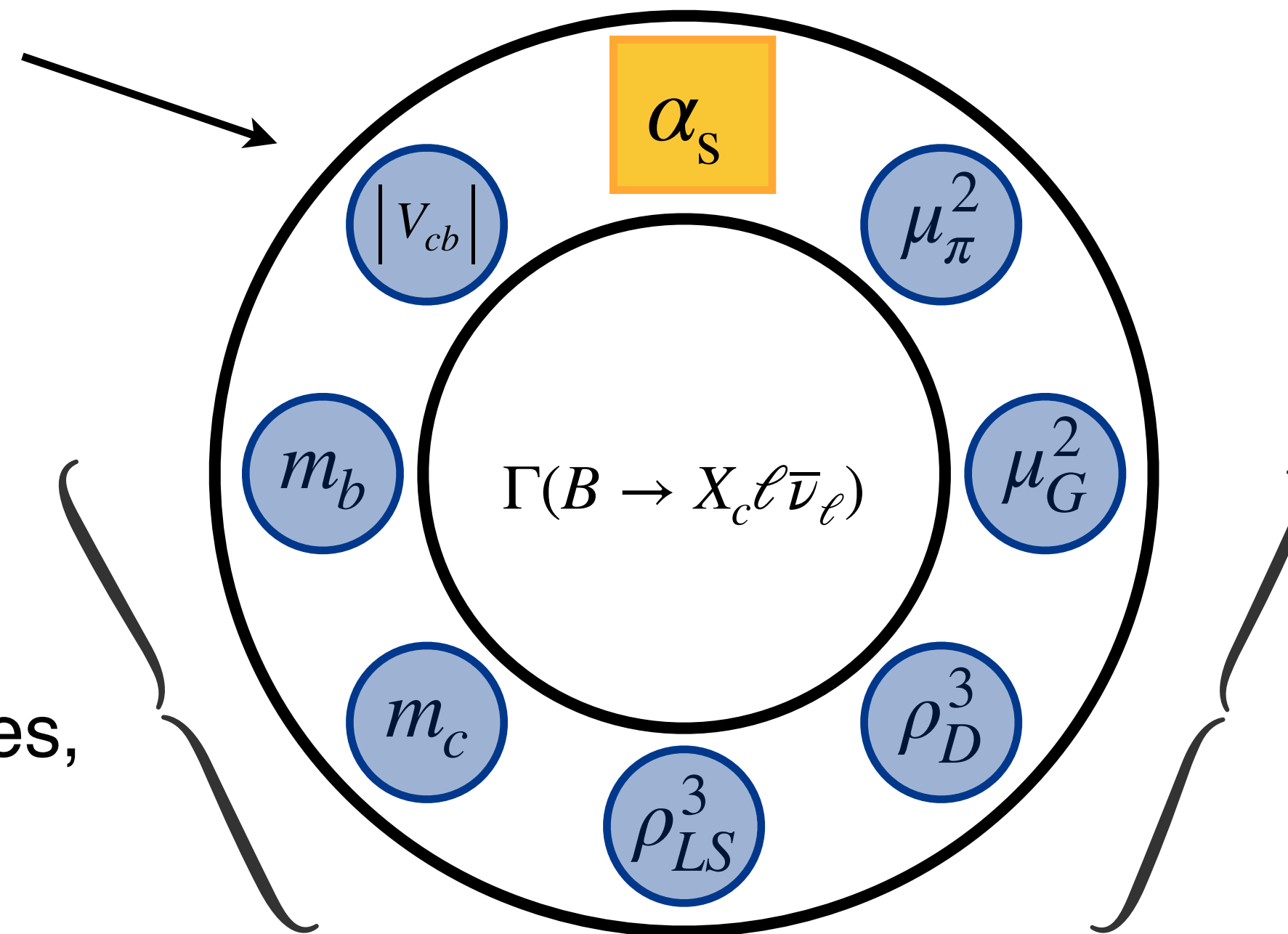
# Heavy quark expansion for $B \rightarrow X_c \ell \bar{\nu}_\ell$ and $D \rightarrow X_s \ell \bar{\nu}_\ell$

$$\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cb}|^2 m_b^5 A_{\text{ew}}}{192\pi^3} \left[ \sum_{i=0,1,2,\dots} c_i \left( \frac{\alpha_s}{\pi} \right)^i + C_{\mu_\pi} \frac{\mu_\pi^2}{2m_b^2} + C_{\mu_G} \frac{\mu_G^2}{2m_b^2} + C_{\rho_D} \frac{\rho_D^3}{2m_b^3} + C_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_b^3} + \dots \right]$$

Exclusive  $B$  decays  
 $W$  decays, etc

Extract  $\alpha_s$  from  $B \rightarrow X_c \ell \bar{\nu}_\ell$ , by fixing  
 the other parameters at external  
 determinations?

B/D meson masses,  
 Lattice QCD,  
 etc.



HQE parameters:  
 Differential decay width of  $B \rightarrow X_c \ell \bar{\nu}$   
 Lattice QCD

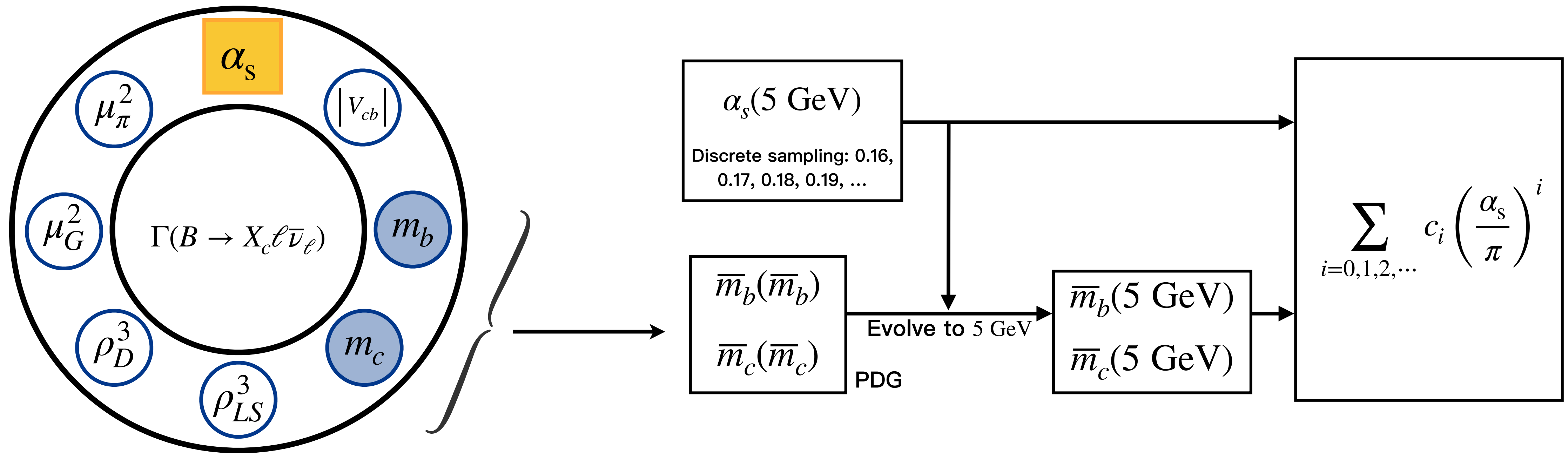
# A numerical relation between $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$ and $\alpha_s$

$$\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cb}|^2 m_b^5 A_{\text{ew}}}{192\pi^3} \left[ \sum_{i=0,1,2,\dots} c_i \left( \frac{\alpha_s}{\pi} \right)^i + C_{\mu_\pi} \frac{\mu_\pi^2}{2m_b^2} + C_{\mu_G} \frac{\mu_G^2}{2m_b^2} + C_{\rho_D} \frac{\rho_D^3}{2m_b^3} + C_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_b^3} + \dots \right]$$

## Leading-order power correction:

$c_i$ : calculated to 4th order, depending on  $m_b, m_c$ .<sup>[1,2]</sup>

reformulated consistently in terms of the  **$\overline{\text{MS}}$ -renormalized** quark masses  $\bar{m}_b(\mu)$ ,  $\bar{m}_c(\mu)$  and  $\alpha_s(\mu)$ , with  $\mu = 5 \text{ GeV}$ .



[1] <https://doi.org/10.1103/PhysRevD.78.114015>

[2] <https://doi.org/10.1103/PhysRevD.104.016003>

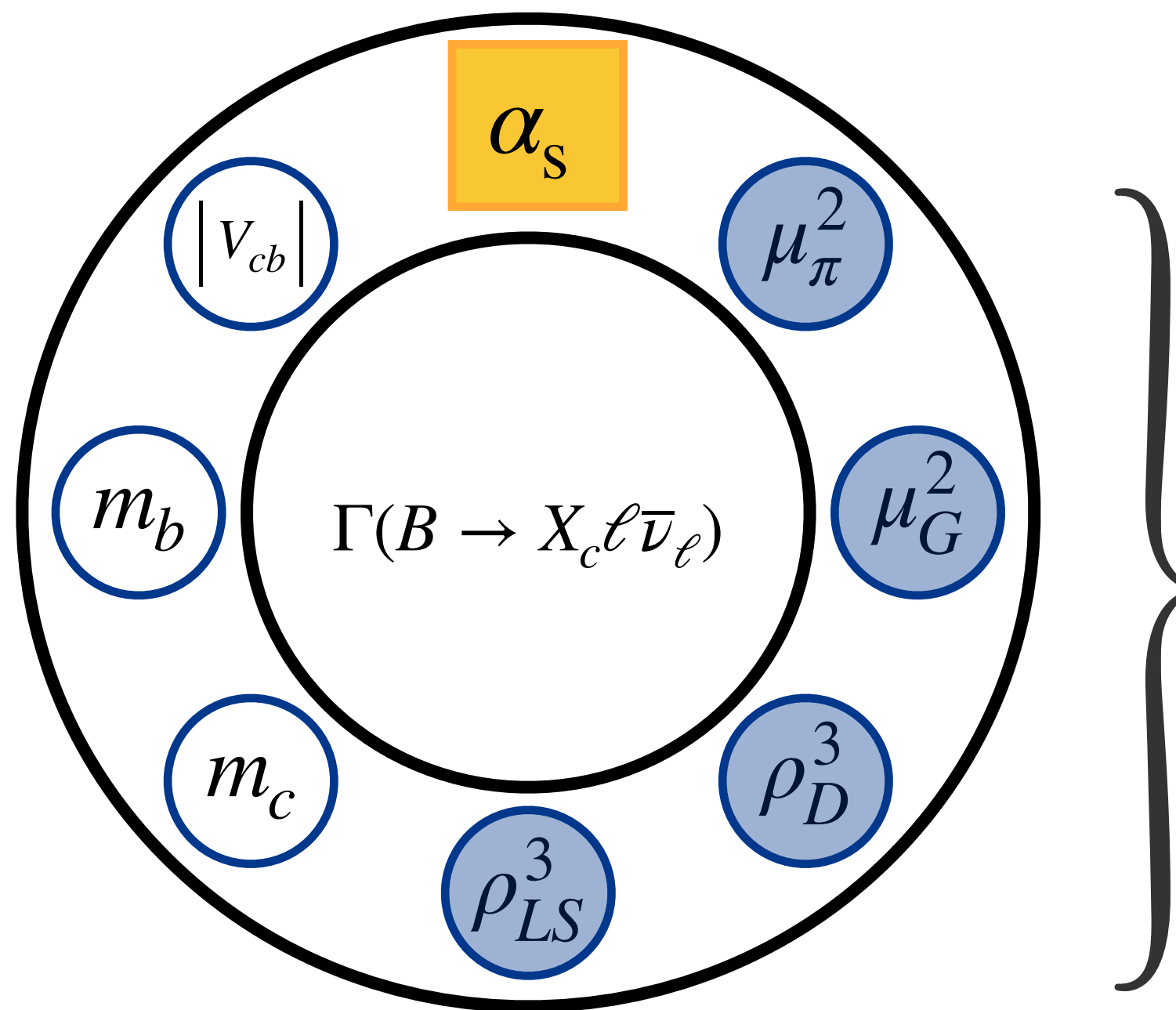


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## • 2nd and 3rd order power corrections:

- decrease the decay width by  $\sim 7\%$  [arXiv:1009.4622].



- The coefficients are also series of  $\alpha_s$ , in the **kinetic scheme**, [1]

$$C_{\mu_\pi} = 2\mathbf{c}_0 \left( \frac{1}{2} - 0.99 \frac{\alpha_s}{\pi} \right), C_{\mu_G} = -2\mathbf{c}_0 \left( 1.94 + 3.46 \frac{\alpha_s}{\pi} \right), \text{ etc.}$$

- The non-perturbative parameters: [2]

$$\mu_\pi^2 = 0.477 \pm 0.056 \text{ GeV}^2, \mu_G^2 = 0.306 \pm 0.050 \text{ GeV}^2, \text{ etc}$$

- The higher order power corrections are estimated around -7%, with the coefficients estimated using their first terms.

- **Truncation error:**  $\mathcal{O}(\alpha_s/m_b^{2,3})$  terms and  $\mathcal{O}(1/m_b^{4,5})$  terms,  $\sim 2.3\%$ .

[1] [https://doi.org/10.1007/JHEP01\(2014\)147](https://doi.org/10.1007/JHEP01(2014)147)

[2] <https://doi.org/10.1016/j.physletb.2021.136679>

# Determining $\alpha_s$ from $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$

$$\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cb}|^2 m_b^5 A_{\text{ew}}}{192\pi^3} \left[ \sum_{i=0,1,2,\dots} c_i \left( \frac{\alpha_s}{\pi} \right)^i + C_{\mu_\pi} \frac{\mu_\pi^2}{2m_b^2} + C_{\mu_G} \frac{\mu_G^2}{2m_b^2} + C_{\rho_D} \frac{\rho_D^3}{2m_b^3} + C_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_b^3} + \dots \right]$$

**Experimental inputs:** the lifetime and the branching ratios,

$$\tau_{B^\pm} = 1.638 \pm 0.004 \text{ ps}, \mathcal{B}(B^\pm \rightarrow X_c \ell \nu) = 10.8 \pm 0.4 \% , \tau_{B^0} = 1.517 \pm 0.004 \text{ ps}, \mathcal{B}(B^0 \rightarrow X_c \ell \nu) = 10.1 \pm 0.4 \% .$$

TABLE I. The parameters used during the construction of the theoretical model.

Parameter	Notation	Value & error	Note
Fermi coupling constant	$G_F$	$1.16637886 \times 10^{-5} \text{ GeV}^{-2}$	[50]
Electroweak correction factor	$A_{\text{ew}}$	1.014	[51]
CKM matrix element	$ V_{cb} $	$0.0398 \pm 0.0006$	[50]
$b$ -quark mass in $\overline{\text{MS}}$	$\bar{m}_b(\bar{m}_b)$	$4.18^{+0.03}_{-0.02} \text{ GeV}$	[50]
$c$ -quark mass in $\overline{\text{MS}}$	$\bar{m}_c(\bar{m}_c)$	$1.27 \pm 0.02 \text{ GeV}$	[50]
HQE parameters	$\mu_\pi^2$	$0.477 \pm 0.056 \text{ GeV}^2$	[28]
	$\mu_G^2$	$0.306 \pm 0.050 \text{ GeV}^2$	[28]
	$\rho_D^3$	$0.185 \pm 0.031 \text{ GeV}^2$	[28]
	$\rho_{LS}^3$	$-0.130 \pm 0.092 \text{ GeV}^2$	[28]
	$b$ -quark mass in kinetic scheme	$m_b^{\text{kin}} = 4.573 \pm 0.012 \text{ GeV}$	[28]

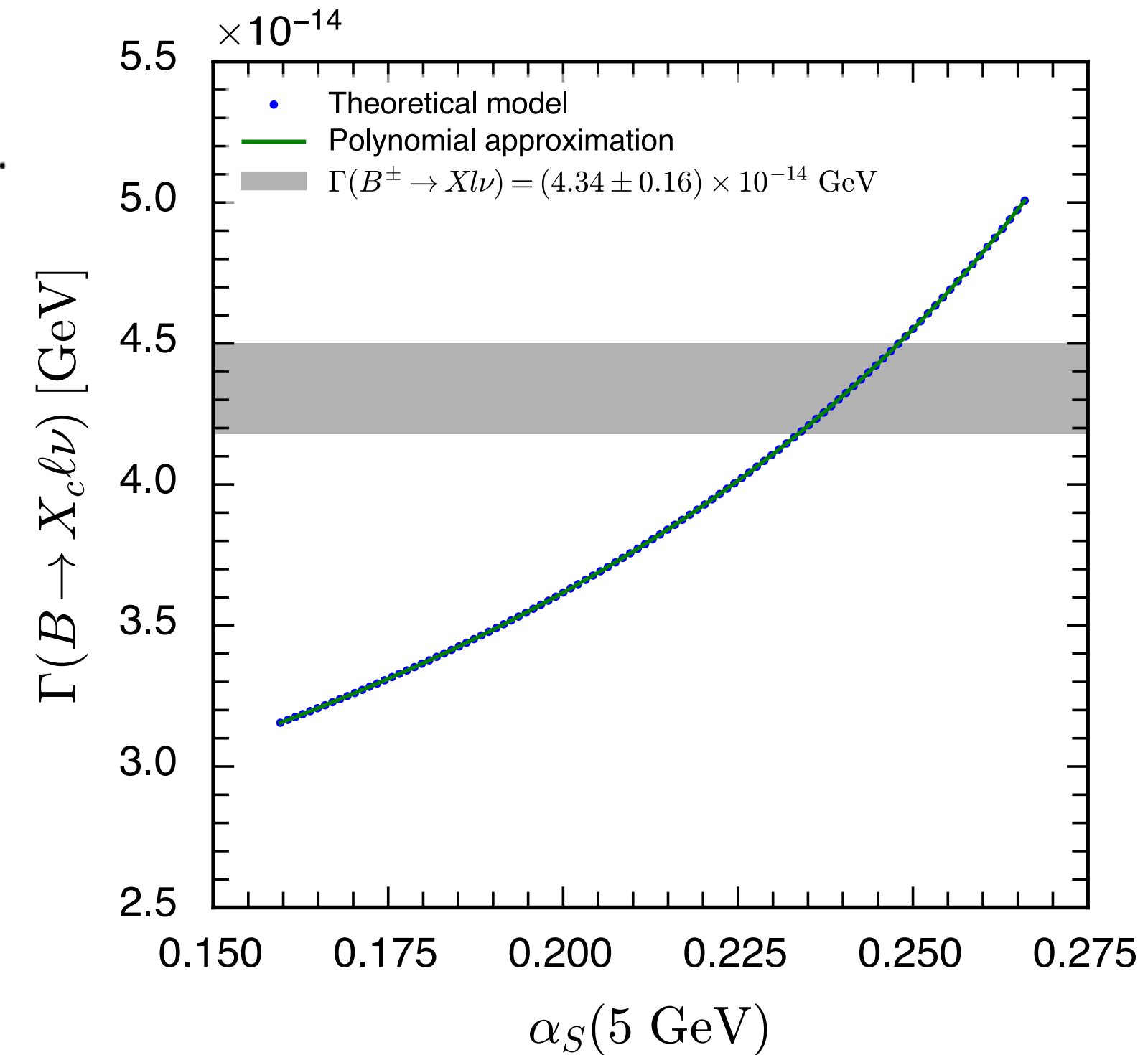


FIG I. The numerical dependence of  $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$  versus  $\alpha_s(5 \text{ GeV})$ , compared with experimental measurement of  $B^\pm$ .

# Determining $\alpha_s$ from $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$

The  $\chi^2$  fit, defined as  $\chi^2(\alpha_s) = \frac{[\Gamma_B - \hat{\Gamma}_B(\alpha_s)]^2}{\sigma_{\Gamma_{B,\text{exp}}}^2 + \sigma_{\Gamma_{B,\text{theo}}}^2}$ , yields

$$\alpha_s(5 \text{ GeV}) = 0.246 \pm 0.013$$

Combining the two fits gives  $\alpha_s(5 \text{ GeV}) = 0.245 \pm 0.009$ , corresponding to  $\alpha_s(m_Z) = 0.1266 \pm 0.0023$ .

***This resulting precision is comparable to the PDG averages of  $\alpha_s(m_Z)$  obtained from other experimental methods.***

TABLE. The relative uncertainty contributions to the theoretical prediction of  $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$  and the  $\alpha_s(5 \text{ GeV})$  fitting result using  $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$ . Values in the parenthesis are the perspective values considering future improvements.

	$\Gamma_{sl}$ prediction [%]	$\alpha_s(5 \text{ GeV})$ [%]
$ V_{cb}  = 0.0398 \pm 0.0006$	3.0 (1.4)	2.1 (1.0)
$\bar{m}_b(\bar{m}_b) = 4.18_{-0.02}^{+0.03} \text{ GeV}$	3.0 (1.1)	2.1 (0.8)
$\bar{m}_c(\bar{m}_c) = 1.27 \pm 0.02 \text{ GeV}$	2.1 (1.4)	1.4 (1.0)
R-scale $\mu = 5_{-2.5}^{+5} \text{ GeV}$	4.4 (2.2)	3.1 (1.6)
High-order power corrections	2.3 (2.3)	1.6 (1.6)
$\tau_{B^\pm} = 1.638 \pm 0.004 \text{ ps}$	-	0.2
$\mathcal{B}(B^\pm \rightarrow X_c \ell \nu) = 10.8 \pm 0.4 \%$	-	3.0 (2.2)
Sum	6.9(3.2)	5.7 (3.5)

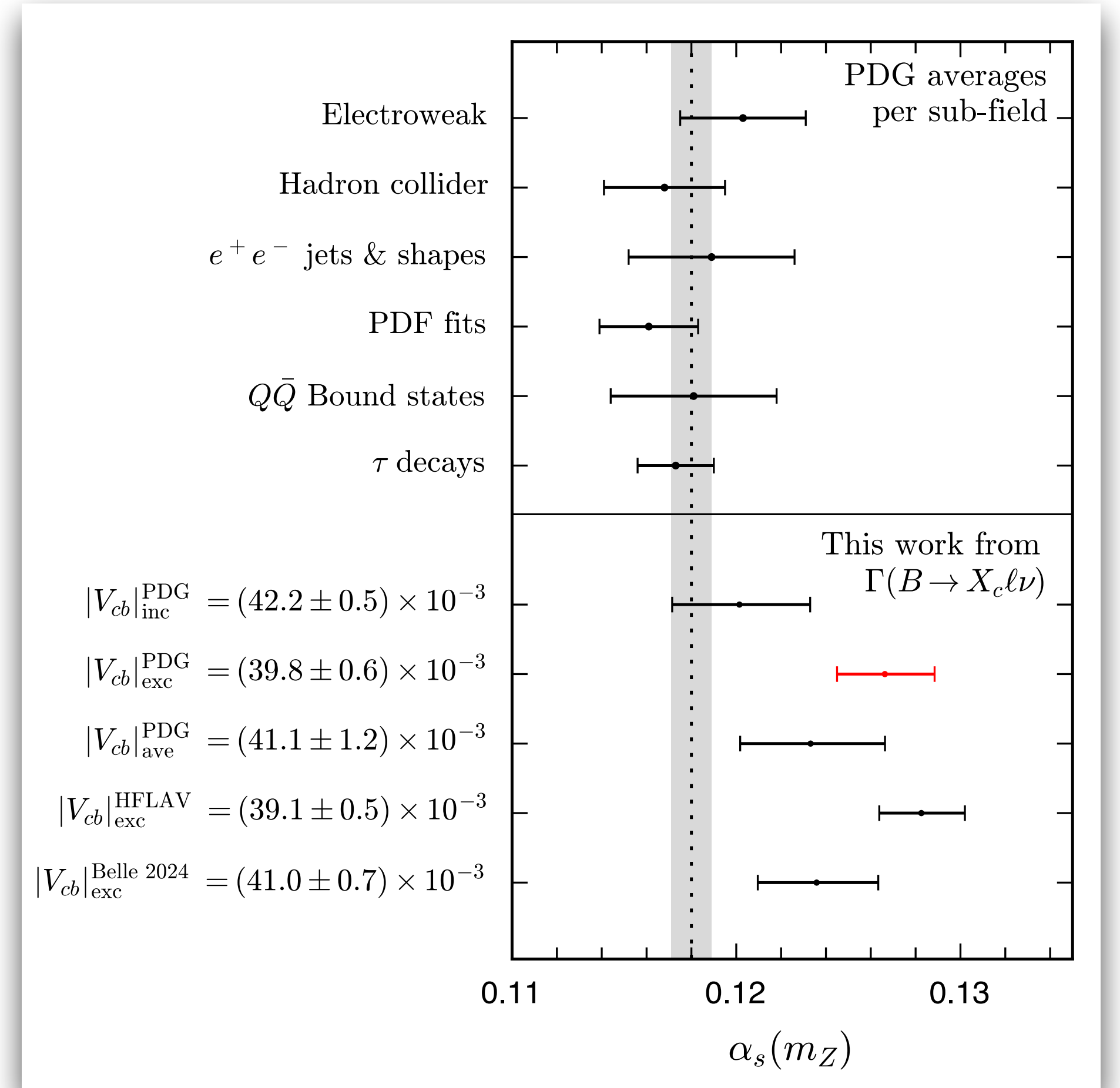


FIG. The comparison of the  $\alpha_s(m_Z)$  pre-averages from six experimental subfields in PDG and the extrapolated values obtained in this work. The running of  $\alpha_s$  along the energy scale is conducted using the RunDec package [arXiv:hep-ph/0004189]



# Determining $\alpha_s$ from $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$

- **Primary uncertainties:**

- the perturbative expansion;
- Branching ratio
- the value of  $|V_{cb}|$

- **The resulting  $\alpha_s(5 \text{ GeV})$  is larger than the world average value by  $\sim 2\sigma$ , which is correlated with the  $|V_{cb}|$  puzzle.**

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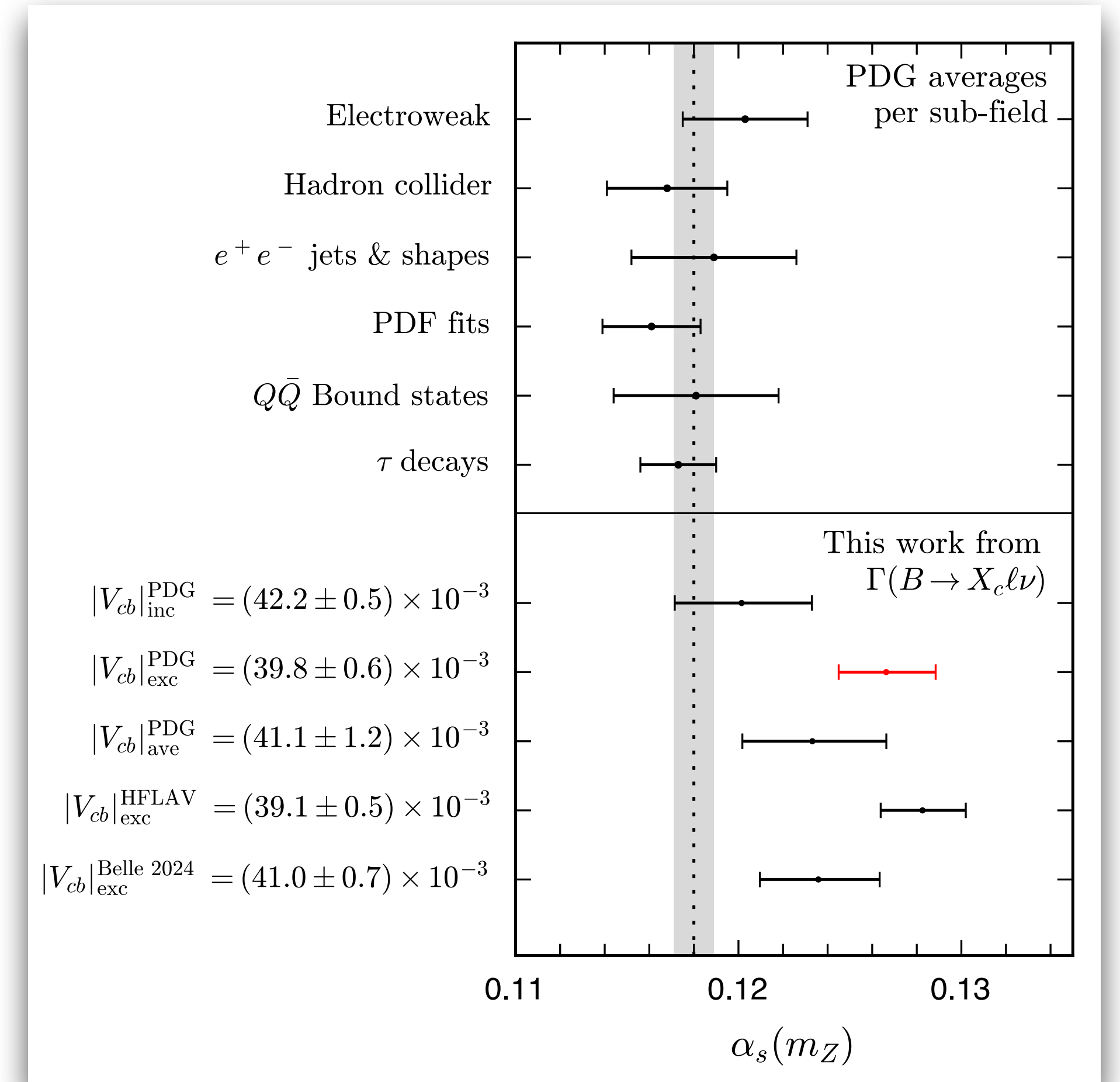


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# Determining $\alpha_s$ from $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$

- **Prospects:**

- $|V_{cb}|$  and B decay should be deeper understood at Belle II and **future Z factories**.
- other improvement from:
  - higher-order perturbation calculation,
  - lattice determinations of b- and c-quark masses.

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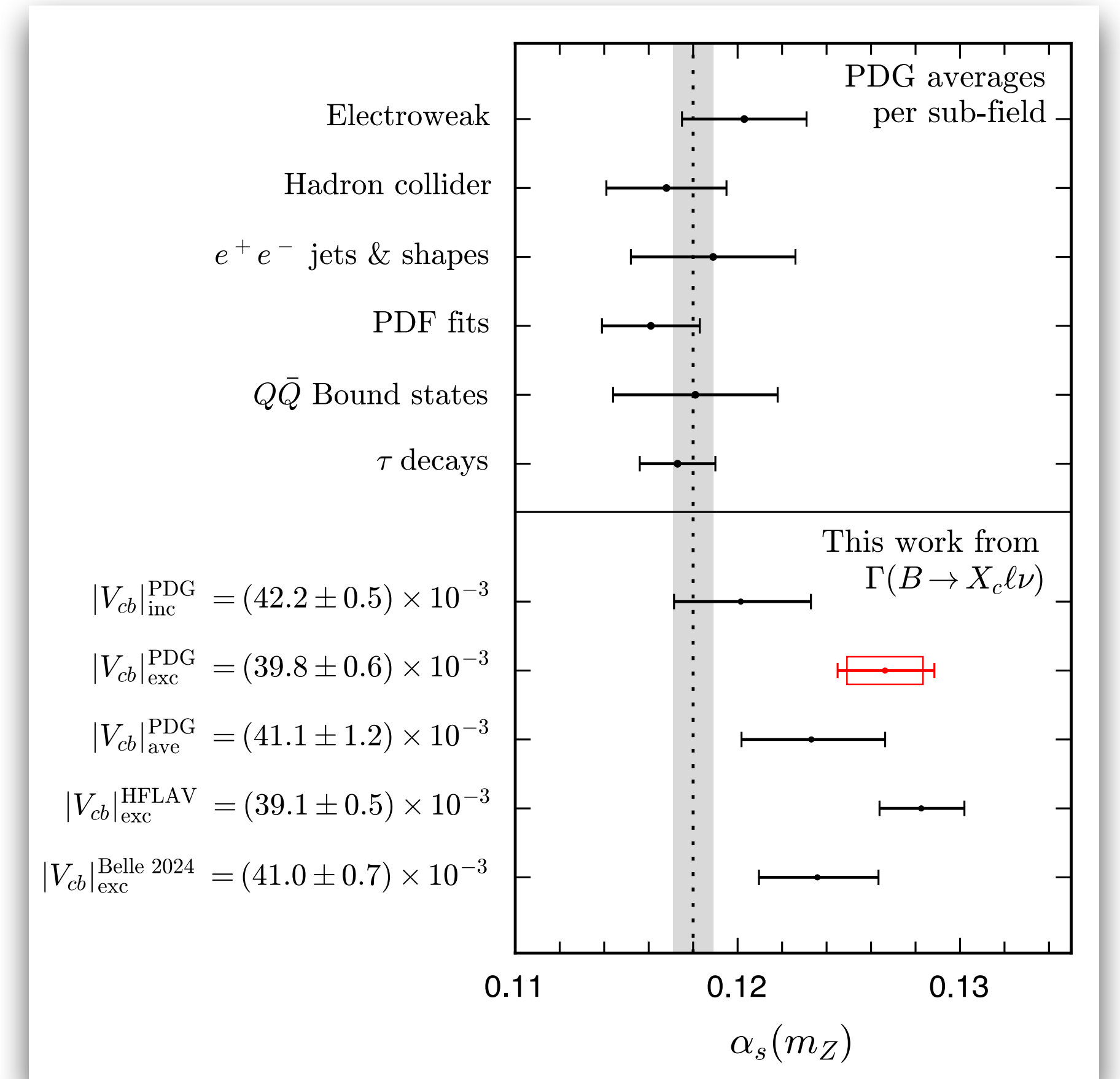


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# Determining $\alpha_s$ from $\Gamma(D \rightarrow X_s \ell \bar{\nu}_\ell)$

Another work on  $D \rightarrow X_c \ell \bar{\nu}$  [arXiv:2406.16119], using

$$\Gamma(D \rightarrow X_s \ell \bar{\nu}_\ell) = \frac{G_F^2 |V_{cs}|^2 m_c^5 A_{\text{ew}}}{192\pi^3} \left[ \begin{aligned} &f_0 + f_1 \frac{\alpha_s}{\pi} + f_2 \left(\frac{\alpha_s}{\pi}\right)^2 + \mathcal{O}(\alpha_s^3) \\ &+ f_{\mu_\pi} \frac{\mu_\pi^2}{2m_c^2} + f_{\mu_G} \frac{\mu_G^2}{2m_c^2} \\ &+ f_{\rho_D} \frac{\rho_D^3}{2m_c^3} + f_{\rho_{LS}} \frac{\rho_{LS}^3}{2m_c^3} + \frac{32\pi^2}{m_c^3} B_{\text{WA}} \\ &+ \dots \end{aligned} \right]$$

yields,

$$\alpha_s(m_c) = 0.445 \pm 0.009_{\text{exp}} \pm 0.081_{m_c} \pm 0.056_{\text{trun}} \pm 0.057_{\text{others}}$$

***The HQE expressions of the semi-leptonic D & B decay widths share the same parameters.***

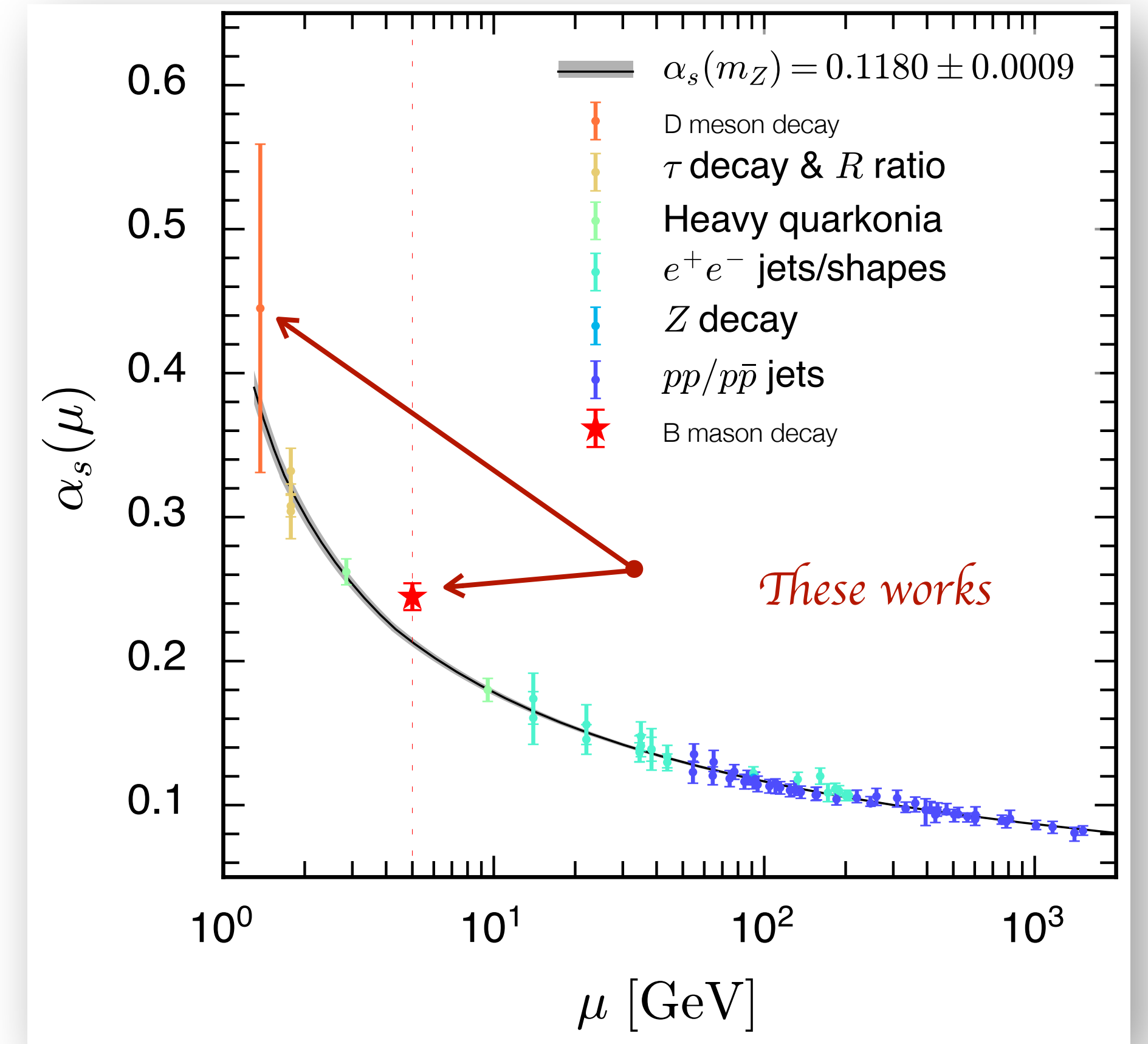


FIG. Comparison of the  $\alpha_s(5 \text{ GeV})$  and  $\alpha_s(m_c)$  results with the  $\alpha_s$  determinations at different energy scales.



# Discussion

Recent studies determine the HQE parameters from  $D \rightarrow X_c \ell \bar{\nu}$ , providing good precision.

[*JHEP* 05 (2025) 061], [*Eur.Phys.J.C* 85 (2025) 9, 1011]

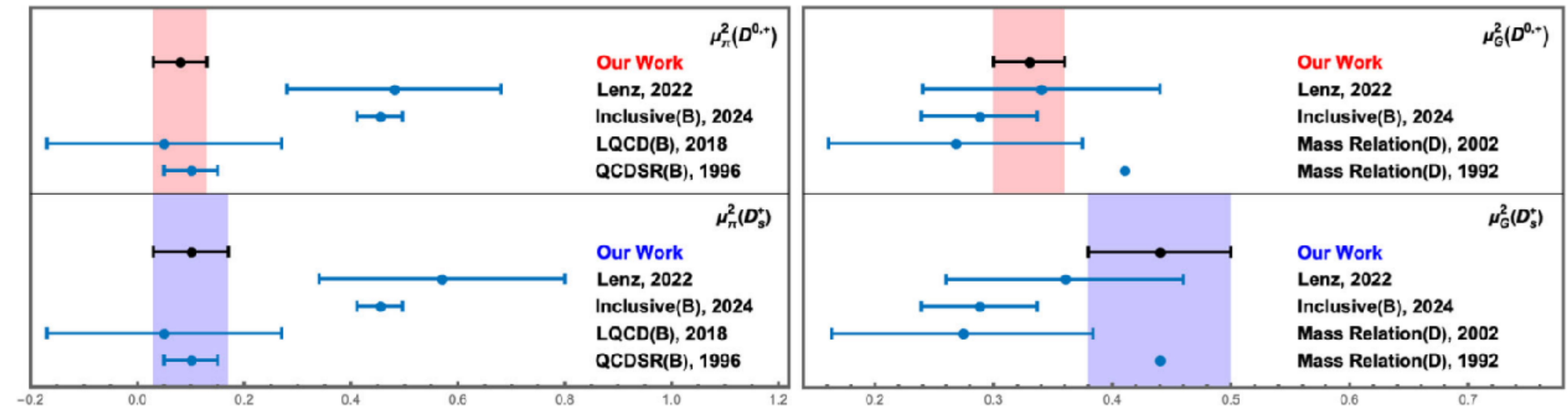
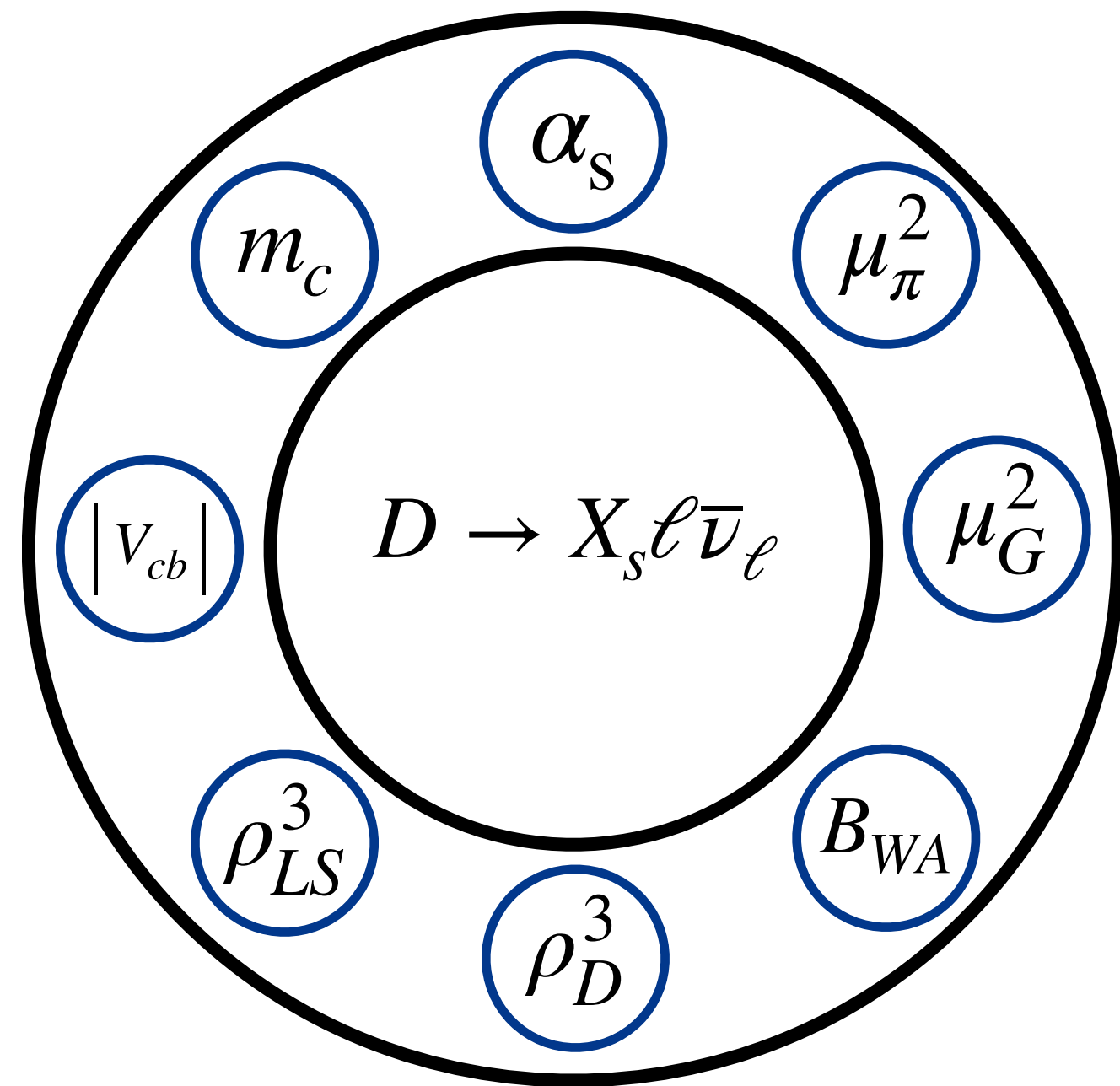


Fig. 1 The comparison between the results for  $\mu_\pi^2$  (left) and  $\mu_G^2$  (right) in “Our work” and those obtained in the literature. The “Inclusive(B), 2024” [38] results are obtained in a global fit to semi-leptonic inclusive  $B$  decay observables for the corresponding  $B$  meson parameters in the kinetic mass scheme, which are related to the  $D$  meson ones by the heavy quark symmetry. The “LQCD(B), 2018” [50] and “QCDSR(B),

1996” [51] results for  $\mu_\pi^2$  are also for the  $B$  meson, calculated by using lattice QCD and QCD sum rules, respectively. The “Mass Relation(D), 2002” [52] and “Mass Relation(D), 1992” [53] results for  $\mu_G^2$  are obtained by using the two versions of mass relations between the  $D_{(s)}$  and  $D_{(s)}^*$  mesons

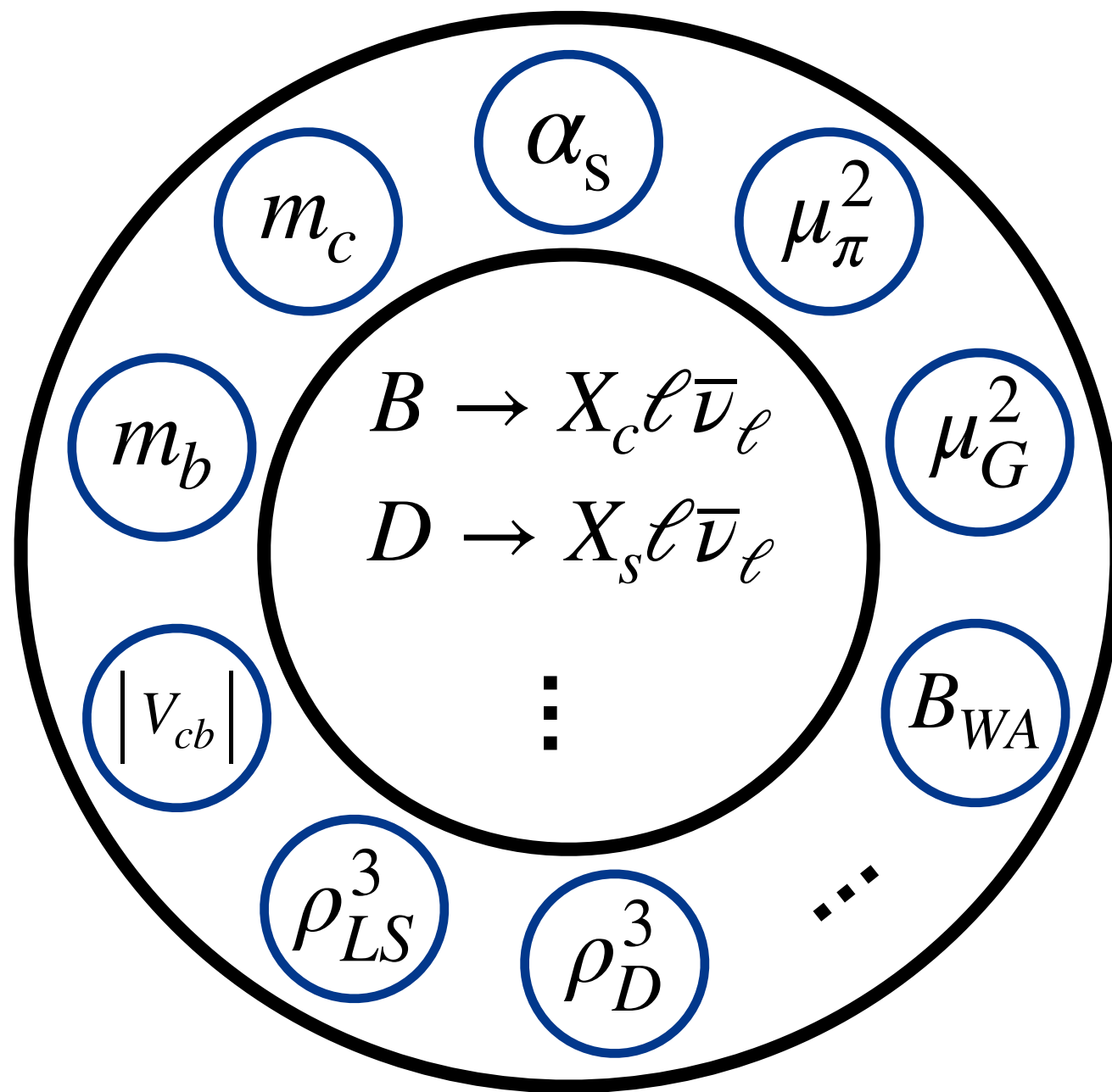
*Eur.Phys.J.C* 85 (2025) 9, 1011

# Discussion

Eventually, a global fit combining multi-processes, multiple observables is expected,

- Simultaneously extract  $\alpha_s$  and other parameters.
- Validate the consistency of the HQE theory.

**CEPC Tera-Z provides valuable experimental input !**



Particle	BESIII	STCF (1 ab <sup>-1</sup> )	Belle II (50 ab <sup>-1</sup> on $\Upsilon(4S)$ )	LHCb (300 fb <sup>-1</sup> )	CEPC (TDR)
$B^0, \bar{B}^0$	-	-	$5.4 \times 10^{10}$	$3 \times 10^{13}$	$4.8 \times 10^{11}$
$B^\pm$	-	-	$5.7 \times 10^{10}$	$3 \times 10^{13}$	$4.8 \times 10^{11}$
$B_s^0, \bar{B}_s^0$	-	-	$6.0 \times 10^8$ (5 ab <sup>-1</sup> on $\Upsilon(5S)$ )	$1 \times 10^{13}$	$1.2 \times 10^{11}$
$B_c^\pm$	-	-	-	$1 \times 10^{11}$	$7.2 \times 10^8$
$\Lambda_b^0, \bar{\Lambda}_b^0$	-	-	-	$2 \times 10^{13}$	$1 \times 10^{11}$
$D^0, \bar{D}^0$	$1.2 \times 10^8$	$7.2 \times 10^9$	$4.8 \times 10^{10}$	$7 \times 10^{14}$	$8.3 \times 10^{11}$
$D^\pm$	$1.2 \times 10^8$	$5.6 \times 10^9$	$4.8 \times 10^{10}$	$3 \times 10^{14}$	$4.9 \times 10^{11}$
$D_s^\pm$	$1 \times 10^7$	$1.8 \times 10^9$	$1.6 \times 10^{10}$	$1 \times 10^{14}$	$1.8 \times 10^{11}$
$\Lambda_c^\pm$	$0.3 \times 10^7$	$1.1 \times 10^9$	$1.6 \times 10^{10}$	$1 \times 10^{14}$	$6.2 \times 10^{10}$
$\tau^+ \tau^-$	$3.6 \times 10^8$	$3.6 \times 10^9$	$4.5 \times 10^{10}$		$1.2 \times 10^{11}$

**[2025 Chinese Phys. C 49 103003]** Expected yields of b-hadrons, c-hadrons, and  $\tau$  leptons at BESIII, STCF, Belle II, LHCb Upgrade II, and CEPC.

# Summary

As a fundamental parameter of QCD,  $\alpha_s$  widely entangled in the observables of semileptonic decays of heavy mesons.

- The HQE expression of  $\Gamma(B \rightarrow X_c \ell \bar{\nu}_\ell)$  shows enough sensitivity to support a determination of  $\alpha_s$  with a precision comparable to the existing approaches.
  - A  $2\sigma$  larger  $\alpha_s(m_Z)$  is obtained, which is correlated to the  $|V_{cb}|$  puzzle.
- Eventually, a global fit combining multi-processes, multiple observables is expected.
- CEPC provides valuable experimental input:
  - Precise  $|V_{cb}|$  determination.
  - Provide experimental data of semileptonic B/D decays.

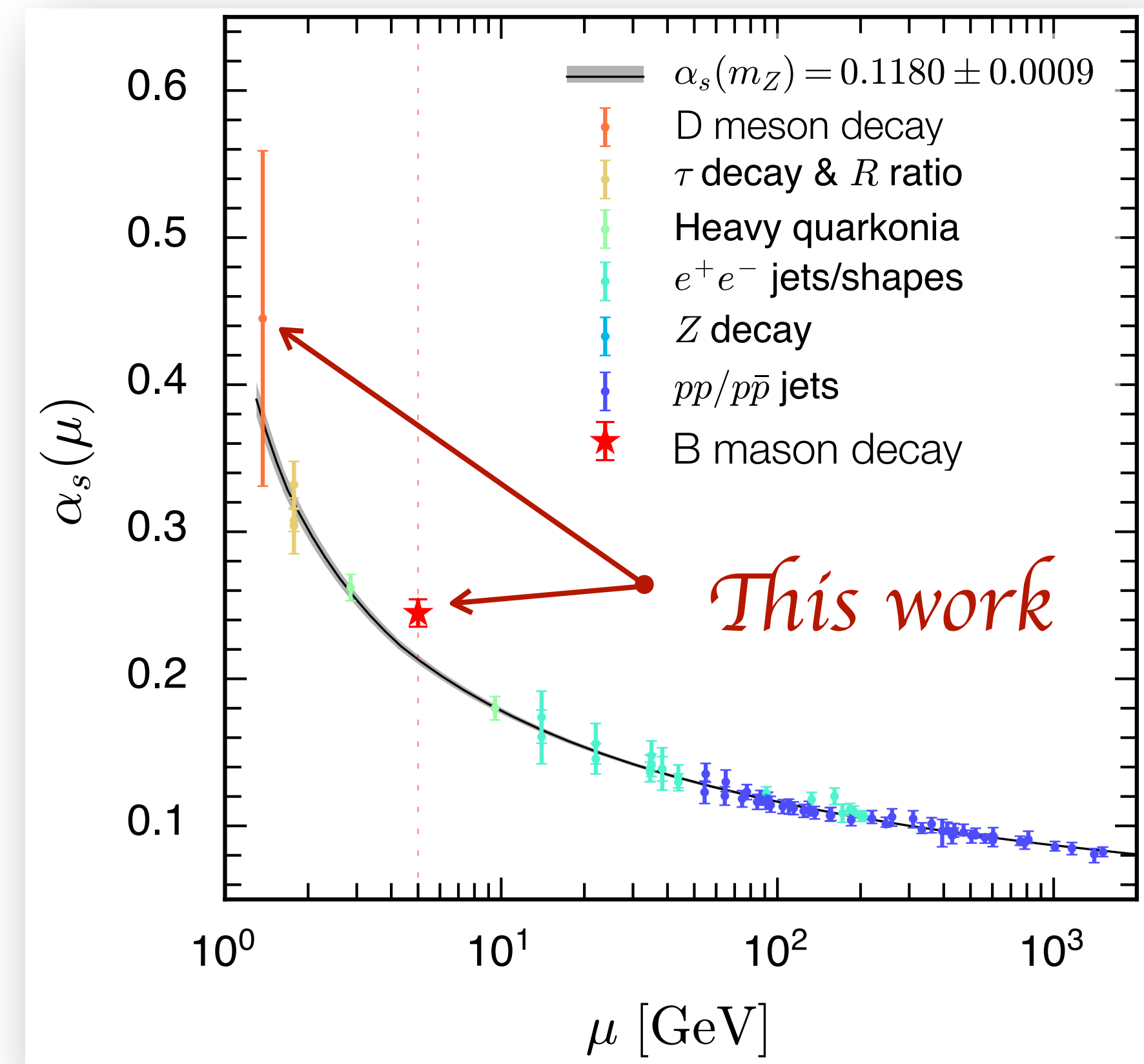


FIG. The  $\alpha_s(5 \text{ GeV})$  and  $\alpha_s(m_c)$  result compared with the  $\alpha_s$  measurements at other energy scales.



# Thanks for your attention

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arXiv:2412.02480,

with Long Chen, Jinfei Wu, Xinchou Lou, Xiang Chen, Xin Guan, Yan-Qing Ma, Manqi Ruan.

The 2025 International Workshop on the High Energy Circular  
Electron Positron Collider