

# Electromagnetic deflection effects in the integrated luminosity measurement at the CEPC

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#### Overview



- Introduction
- Very forward region at CEPC
- Electromagnetic deflection of the initial state (EMD1)
- Electromagnetic deflection of the final state (EMD2)
- Experimental corrections
- Summary



#### Introduction

Challenging control of luminosity systematics at 10<sup>-4</sup> at the Z<sup>0</sup> pole

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- Most sources from mechanics and MDI have been studied and documented in CEPC CDR, CEPC TDR, JINST 17 P09014(2022) and PTEP 10 103H02(2024), including experimental determination of the beam energy spread and its impact on integrated luminosity and precision EW observables
- Electromagnetic fields of opposite bunches will impact initial states of particles in the bunches of opposite charges prior to interaction – EMD1
- In a similar manner, e<sup>-</sup> and e<sup>+</sup> in the final state, whose four-momenta are already altered by EMD1, will be deflected after interaction by the fields of the opposite-charge bunches, which impacts their final states – EMD2
- Both EMD1 and EMD2 will affect the Bhabha count in the luminometer and, consequently, impact the luminosity measurement
- EMD1 and EMD2 haven't been experimentally measured yet, but they can be estimated from simulation



## Very forward region at CEPC (assumed in this study)

- The luminometer is positioned at 95 cm distance from the interaction point, covering the polar angles from 30 mrad to 105 mrad
- Fiducial volume 53 mrad to 79 mrad, where the energy resolution of a high-energy electron (positron) will be constant due to the shower containment
- EMD effects are discussed assuming head-on collision geometry as if the luminometer's halves would be positioned at the outgoing beams (s-frame), that is 16.5 mrad with respect to the z-axis in the laboratory frame
- linstantaneous luminosity at the Z<sup>0</sup> resonance  $\mathcal{L}_0$ =1.15·10<sup>35</sup> cm<sup>-2</sup> s<sup>-1</sup>
- Beam parameters taken from the post-CDR studies (arXiv:2102.09627v1)

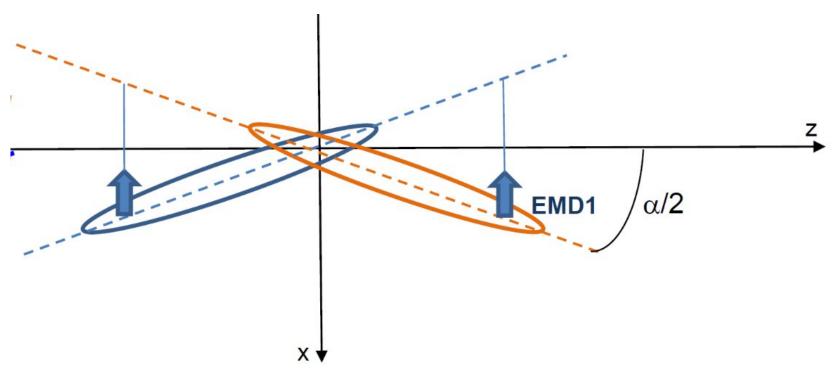
$N (10^{10})$	$\beta_x^*$ (m)	$\beta_y^*$ (mm)	$\sigma_x$ ( $\mu$ m)	$\sigma_y$ ( $\mu$ m)	$\sigma_z$ (mm)
15	0.2	1.0	6.0	0.036	8.7



# Electromagnetic deflection of the initial state

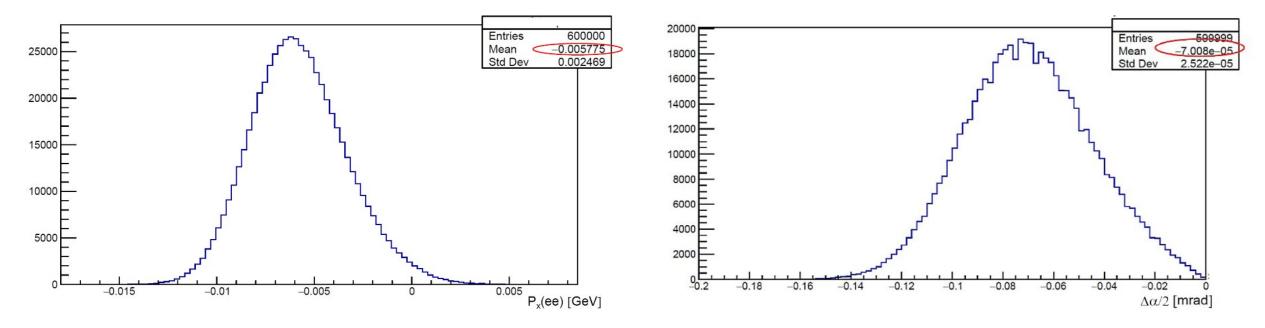
#### EMD1 – $p_x$ -kick of the initial state (s-axis)

- GuineaPig C++ V.1.2.2 (impact of EM fields of outgoing bunches on initial states)
- BHLUMI V4.04 (Low Angle Bhabha Scattering LABS)
- E<sub>beam</sub>= 45.5 GeV
- ~6·10<sup>5</sup> e<sup>+</sup>e<sup>-</sup> interacting pairs





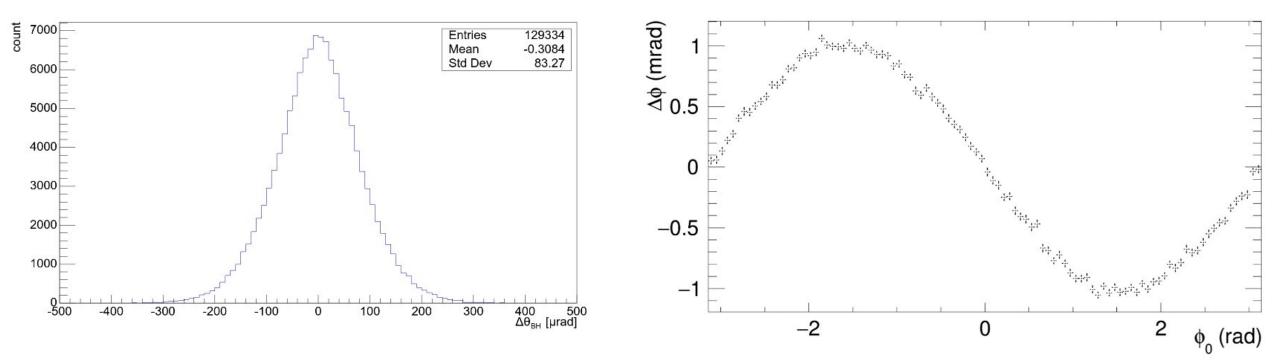
#### Electromagnetic deflection of the initial state



- $e^+e^-$  system receives kick of ~5.8 MeV in x-direction, or ~2.9 MeV per particle in average
- No shift along y-axis
- x-angle effectively reduced for 140 μrad, 70 μrad per beam
- Knowing that energy of the interacting particles (E) depends on change of the crossing angle ( $\Delta\alpha$ ),  $\Delta E$  is found to be ~52 keV per beam



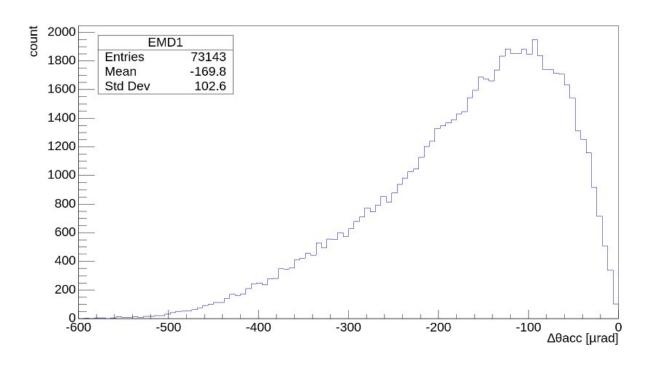
## Electromagnetic deflection of the initial state – impact on final state



- The EMD1 effect results in smearing of the polar angle of the final state particles ( $\Delta\theta_{BH}$ ) with RMS of ~83 µrad
- Maximal  $\Delta\theta_{BH}$  occurs for Bhabha events emitted along the x-axis
- Deviation in azimuthal angles of the Bhabha final states are maximal ( $\Delta \phi \sim 1$  mrad) for LABS events emitted along y-axis ( $\phi_0 = \pi/2$ )



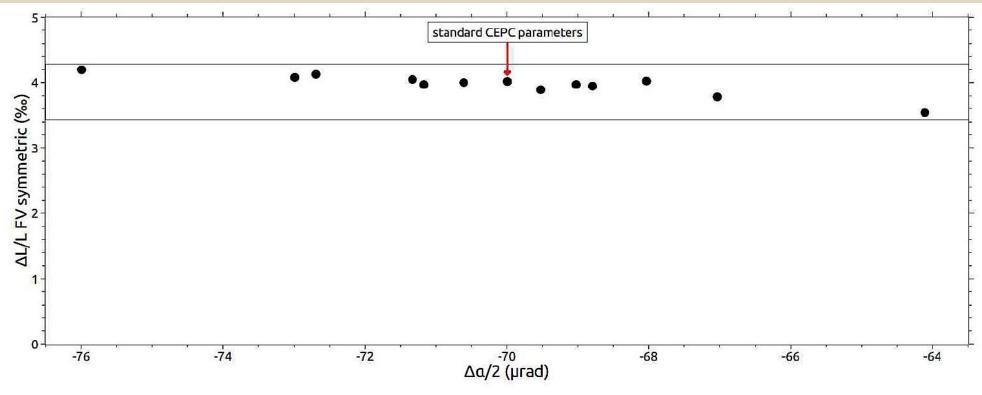
#### Electromagnetic deflection of the initial state – impact on final state



- EMD1 does not bias polar angles of final state e<sup>-</sup> and e<sup>+</sup>, but changes the angle between them, influencing their back-to-back propagation
- Change of angle between final state electrons and positrons ( $\Delta\theta_{acc}$ ) is in average  $^{\sim}$  170 µrad, affecting LABS count in the luminometer due to the loss of collinearity
- The relative loss of count is found to be  $\sim 4\cdot 10^{-3}$  what is  $\sim 40$  times larger than the  $\mathcal{L}_{int}$  precision goal of  $10^{-4}$
- This loss can be taken as correction to the measured integrated luminosity.



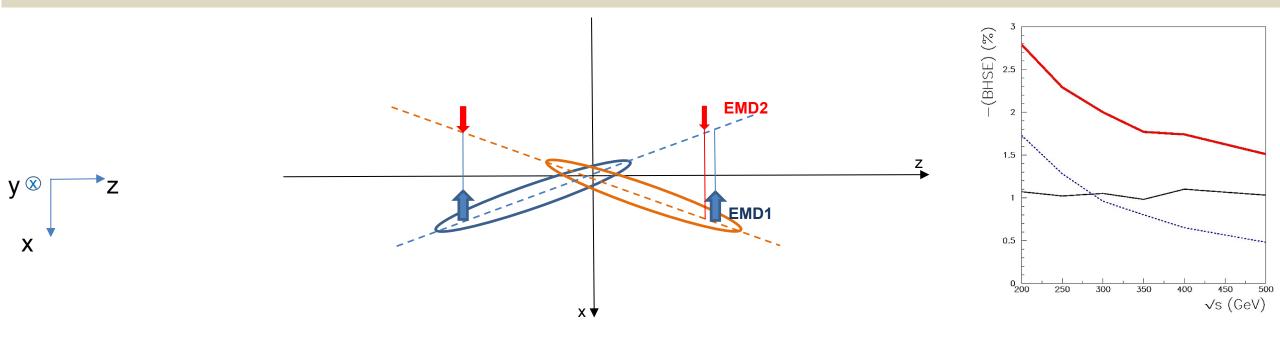
#### Electromagnetic deflection of the initial state – impact on final state



- Variation of beam parameters (bunch charges and dimensions) in the  $\pm 10\%$  range w.r.t. the nominal beam size produce deviation of  $\Delta \mathcal{L}/\mathcal{L} \simeq 2 \cdot 10^{-4}$
- Asymmetric counting (55-77 mrad counting angles in one half of the luminometer and 53-79 mrad in the other, subsequently applied to the left and right arm of the detector) reduces the uncertainty to  $\Delta \mathcal{L}/\mathcal{L} \sim 6\cdot 10^{-5}$



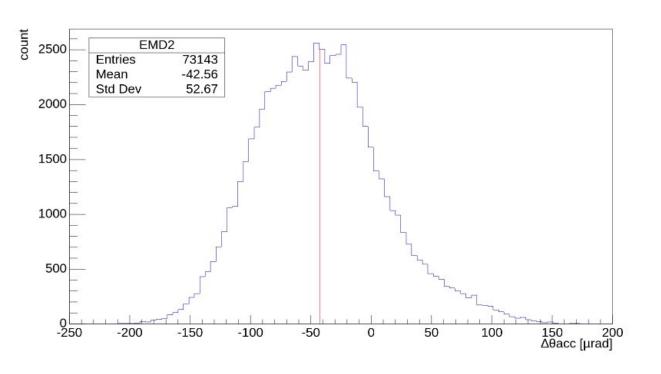
#### Electromagnetic deflection of the final state

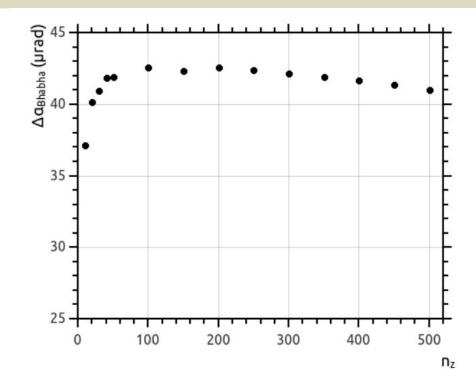


- Final LABS states will also be influenced by electromagnetic fields of the incoming bunches of opposite charges – EMD2
- LABS final states will experience focusing effect towards z-axis
- EMD2 rises with decreasing center-of-mass energy, thus being of relevance at the Z<sup>0</sup> resonance (<u>C. Rimbault</u> et al., 2007 JINST 2 P09001)



#### Electromagnetic deflection of the final state

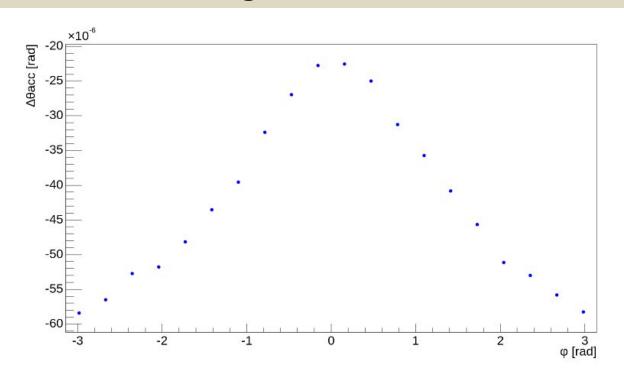


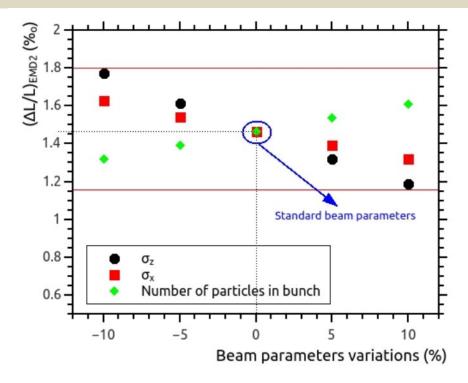


- · In the applied setup, EMD2 induced additional change in Bhabha collinearity is ~43 μrad in average
- The effect is in principal sensitive to the simulation settings (GuineaPig); our results are obtained with 250 longitudinal slices (saturation region) and 7 grids



#### Electromagnetic deflection of the final state

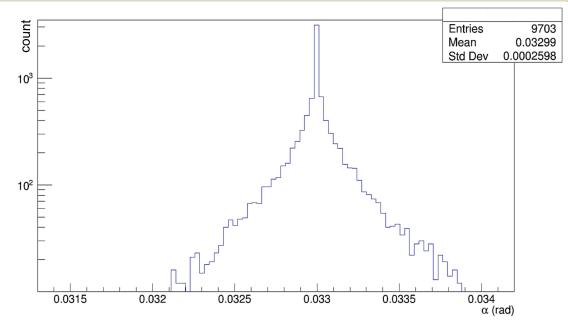




- Change in collinearity of final states  $\Delta\theta_{acc}$  caused by EMD2 also depends on the azimuthal angle ( $\phi$ ) of the final state particles, being largest for particles emitted along the x axis, similarly to the EMD1
- EMD2-induced count loss in the luminometer is ~1.4·10<sup>-3</sup>
- Final state electromagnetic deflection is also relatively insensitive to 10% variation of the bunch parameters and can be corrected from simulation with the relative uncertainty no larger than 3·10<sup>-4</sup>
- The largest  $\Delta \mathcal{L}/\mathcal{L}$  sensitivity is induced by the bunch length variations



#### Experimental corrections – EMD1



- The crossing angle at CEPC can be precisely determined from the central processes like s-channel di-muon production
- Kinematics of di-muon final states will be precisely determined in the central tracker, where momenta of charged high-energy leptons would be measured with the resolution of  $\Delta p_T/p_T^2 \sim 10^{-5}$  GeV<sup>-1</sup>.
- ~10 $^4$  di-muon events generated at the Z $^0$  pole with WHIZARD V2.8.3 event generator; muons' momenta smearing of order of ~10 $^{-5}$  GeV $^{-1}$ , the crossing angle can be measured with the standard error of ~260 $\mu$ rad/ $\sqrt{N_{\mu\mu}}$
- In other words, crossing angle precision of order of ~1  $\mu$ rad can be achieved with as little as 70 pb<sup>-1</sup>  $\mathcal{L}_{int}$  (a couple of minutes)



#### Experimental corrections – EMD2

- With the Si-wafers placed in front of the luminometer, ~µrad precision is achievable in angular measurements of LABS electron and positron
- EMD2 induced loss of count can be corrected in a semi-dependent way from simulation, presumably on a
  basis of experimentally measured acolinearity of LABS final states in the in the luminometer and/or other
  observables ongoing study
- Additional considerations: possible changes of the luminometer fiducial volume towards smaller polar angles; inclusion of radiative processes



#### Summary

- The effects of electromagnetic deflection of initial (EMD1) and final Bhabha states (EMD2) are quantified in simulation with the nominal post-CDR CEPC beams
- EMD1 and EMD2, if not corrected, will cause the relative loss of count in the luminometer of  $\Delta \mathcal{L}/\mathcal{L} \sim 6.10^{-3}$
- EMD1 will lead to reduction of the crossing angle of ~140  $\mu$ rad. Measurement of the crossing angle with  $\mu$ rad precision will be possible with the central-tracker reconstruction of di-muon production. From simulation, relative correction of the integrated luminosity can be determined once the crossing angle is known. The correction is of order of 4 ·10<sup>-3</sup>. Uncertainty of the correction is not larger than 2·10<sup>-4</sup>
- EMD2 can be also corrected from simulation if the angle between LABS electron and positron in known with a sufficient precision, presumably of the order of several  $\mu$ rad provided by the Si-wafer placed in front of the luminometer. The correction induced by the EMD2 effect is smaller than  $2\cdot 10^{-3}$ . Variations of the nominal bunch sizes and population not larger than 10%, introduce uncertainty of this correction not larger than  $3\cdot 10^{-4}$

(Results presented here are submitted to the Progress of Theoretical and Experimental Physics, IF 8.6, arXiv:2511.00687v1 [hep-ex])



# **THANK YOU!**



