

Recent Results on Light Nuclei and Hypernuclei Production and Directed Flow in Au+Au Collisions at RHIC

Physics of High Net-Baryon Densities (PHD2025)

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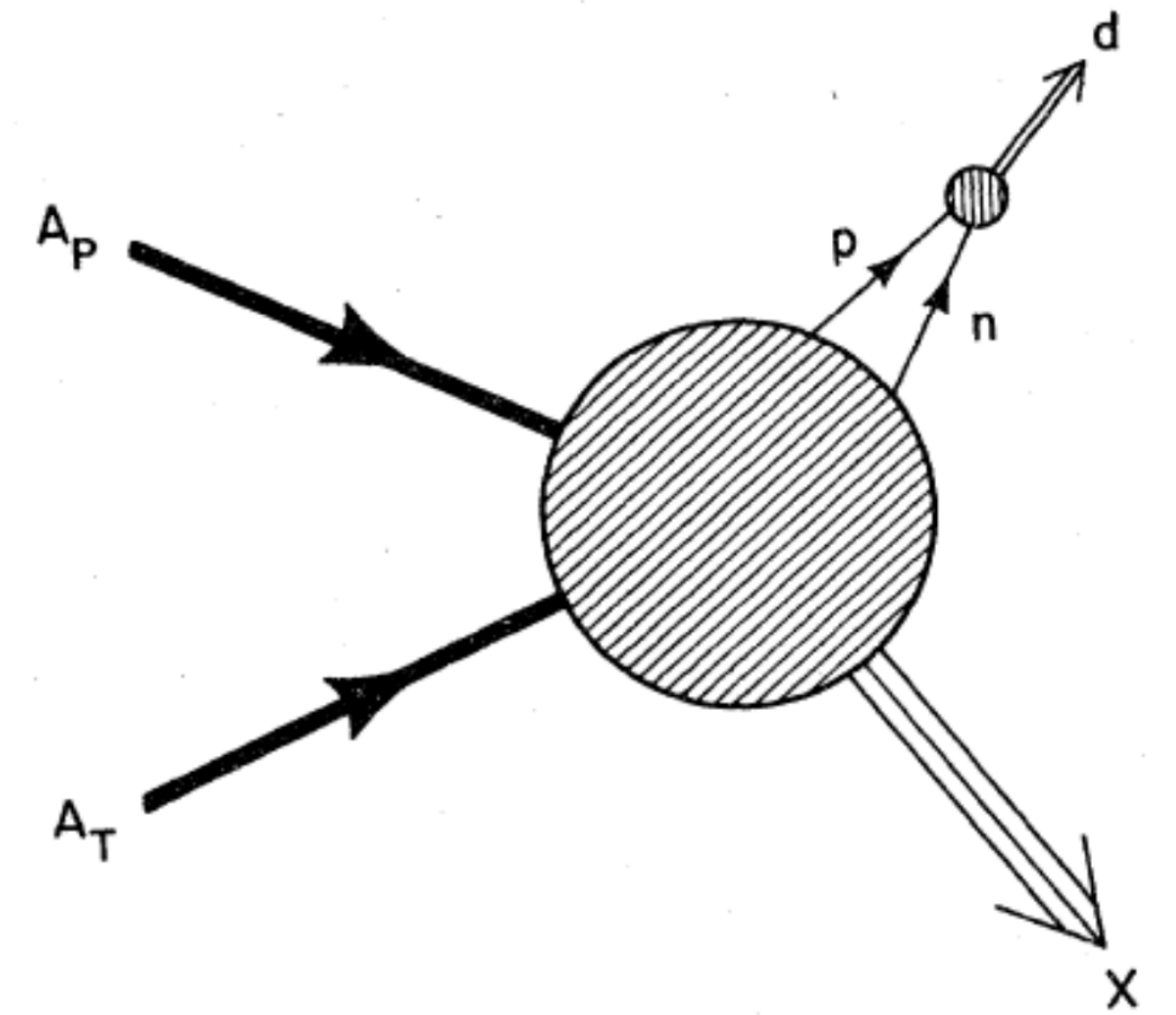
Outline

- Introduction
- Results
 - Production Yields
 - Transverse Momentum Distributions
 - Rapidity Distributions
 - Directed Flow
- Summary and outlook

Introduction

0. Production mechanisms

- Thermal model
 - Particle yields predicted with thermal equilibrium assumptions
 - Chemical freeze-out temperature T_{ch}
 - Baryo-chemical potential μ_B
 - Volume V
 - Strangeness correlation length r_C
- Coalescence model
 - Formed by nucleons (and hyperons) which are nearby in phase space after kinetic freeze-out
 - Nucleon distributions at freeze-out
 - Nuclear wave function
- Others (hadronic scattering, potential mechanism, spectator fragmentation, etc.)



J. I. Kapusta, PRC 21, 1301 (1980)

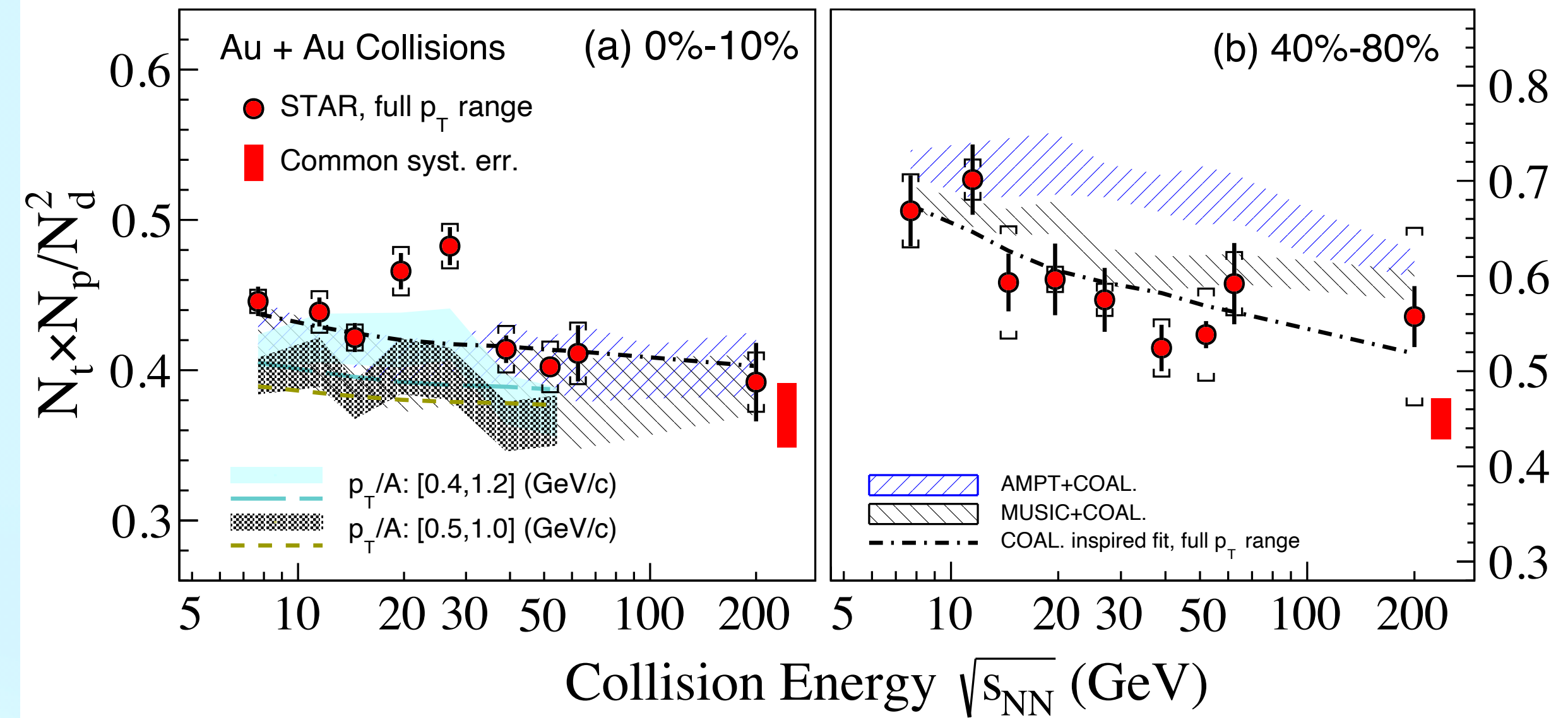
Motivations

1. Probing critical point/phase transition

- The compound ratio $\frac{N_t N_p}{N_d^2}$ has been

proposed to be sensitive to **neutron density fluctuations**

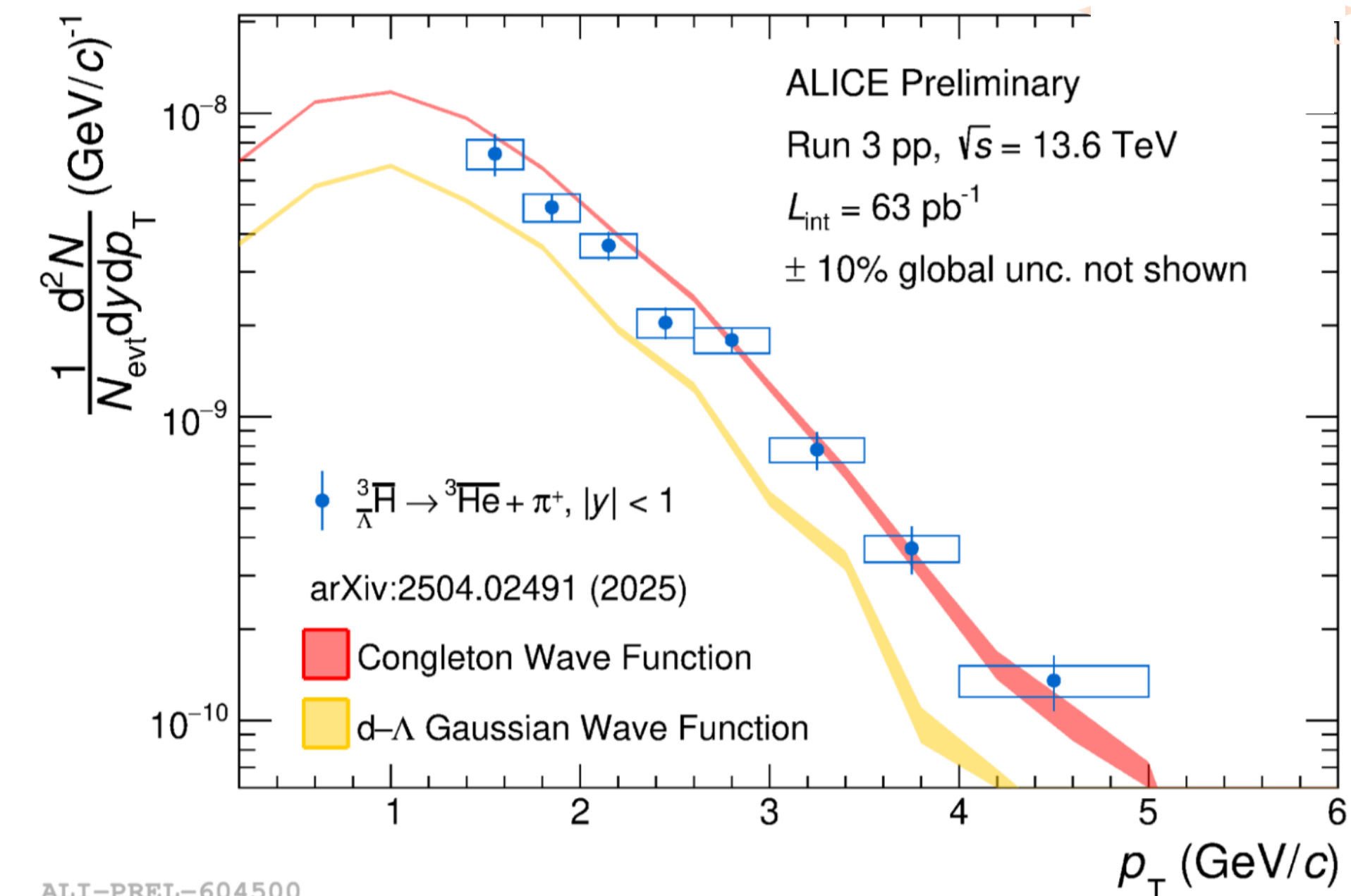
- Enhancements at 19.6 GeV and 27 GeV (4.1σ combined) from BES-I



STAR, Phys.Rev.Lett. 130, 202301 (2023)

2. Probing hypernuclei wave functions

- Structure of hypernuclei are connected to **YN** and **YNN** interactions
- Hypernuclei yield in **small systems**
 - probe short-distance structure of hypernuclear wavefunctions

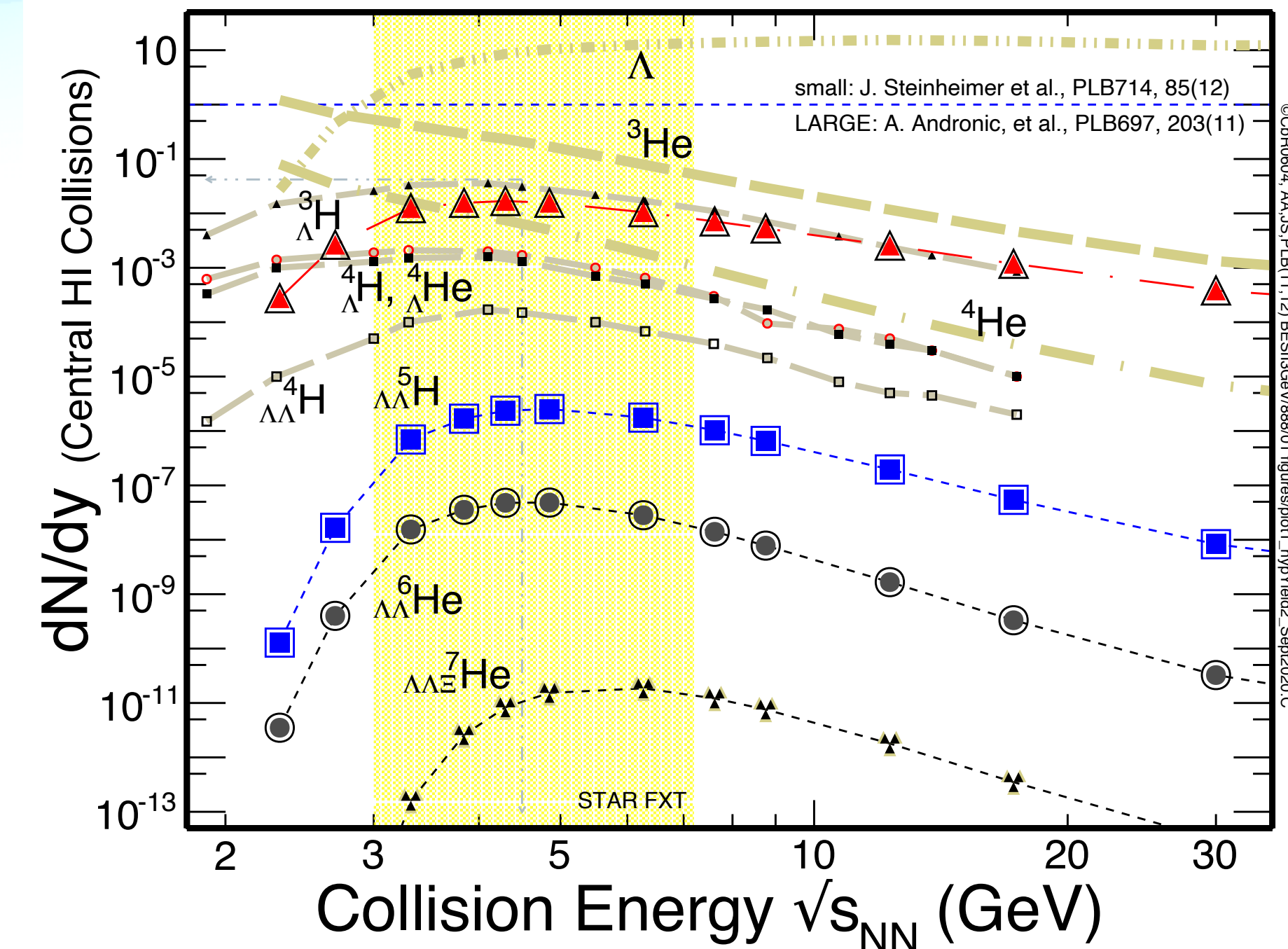
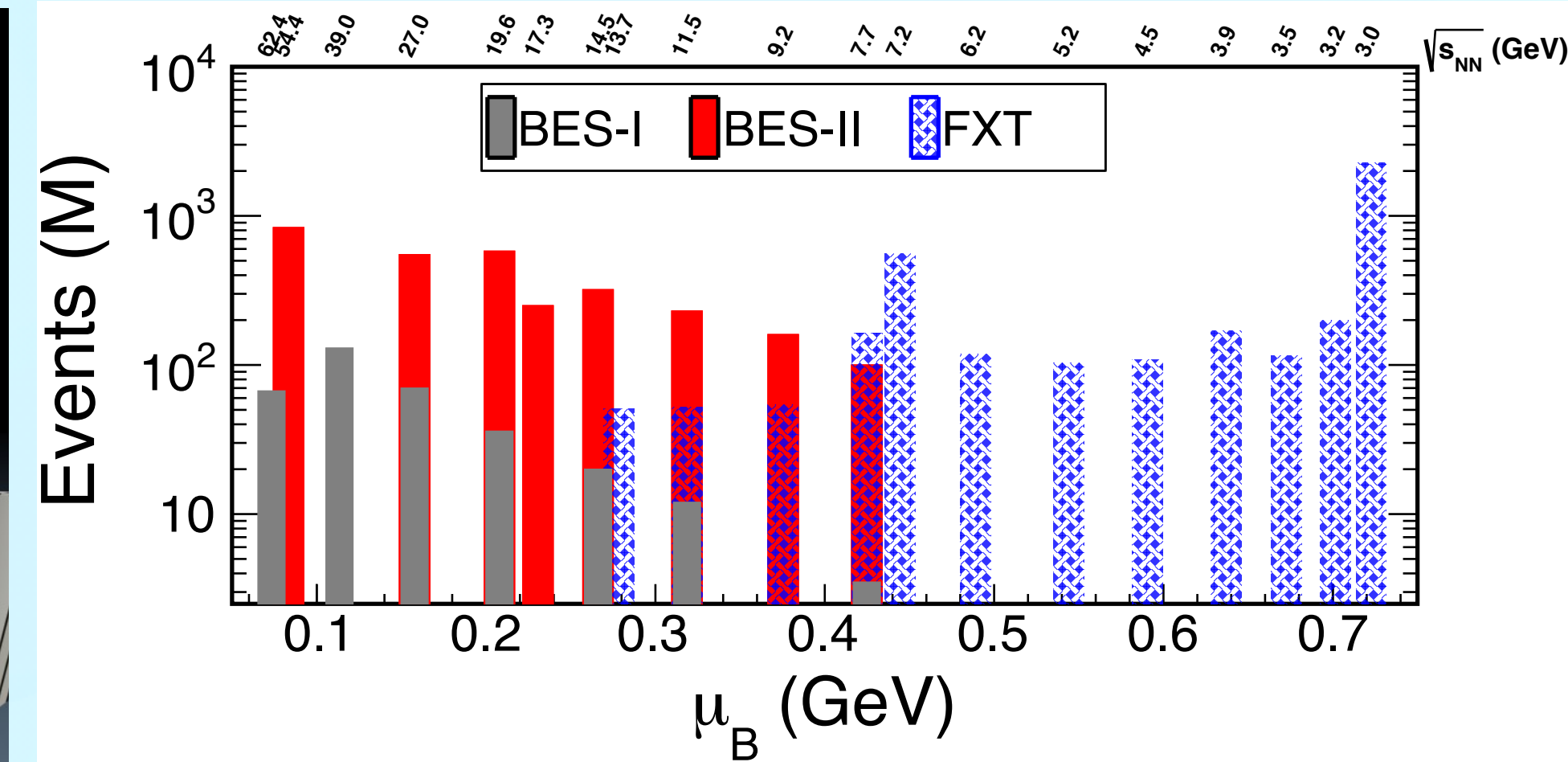
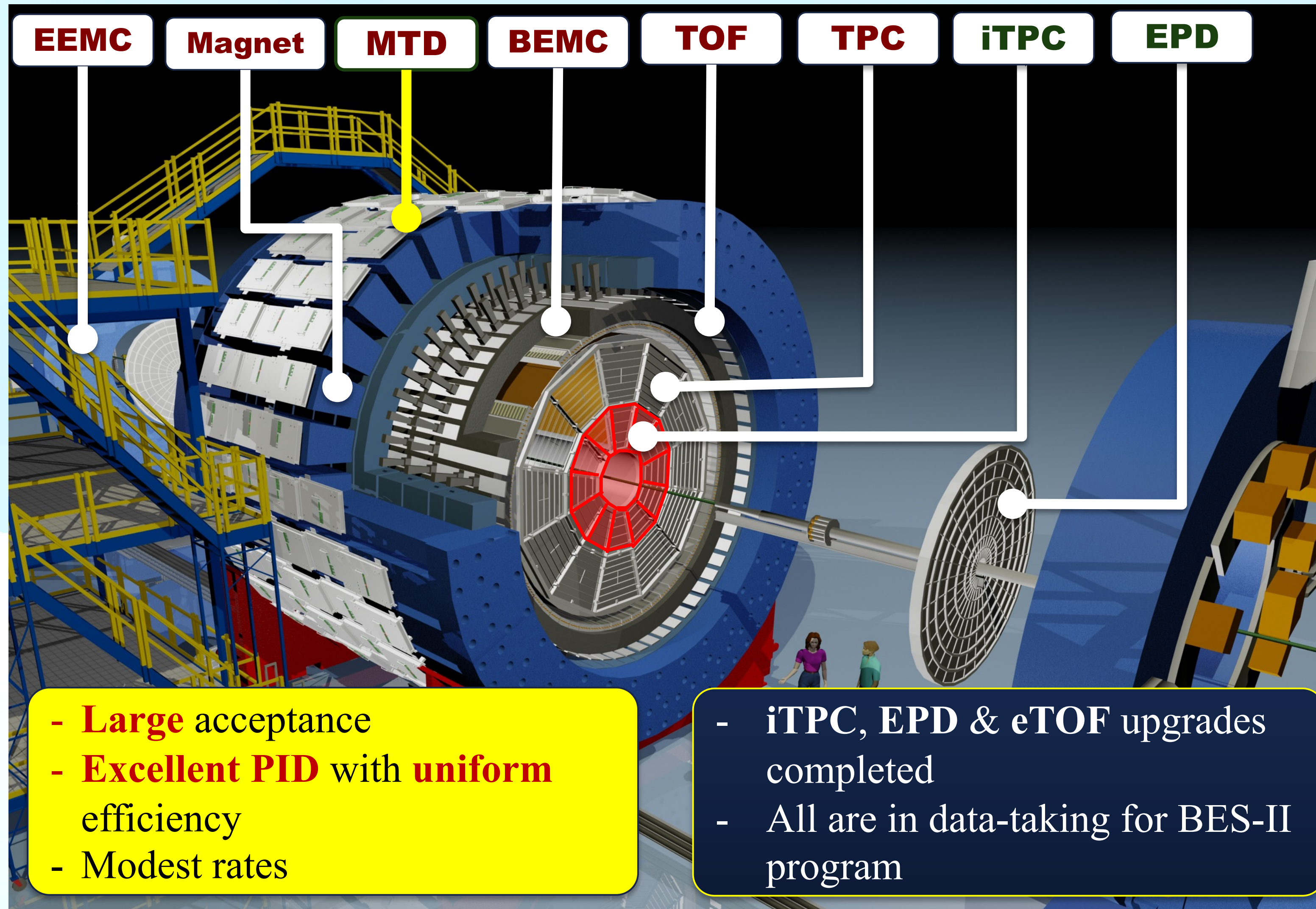


ALI-PREL-604500

Phys.Rev.C 103, 014907 (2021)

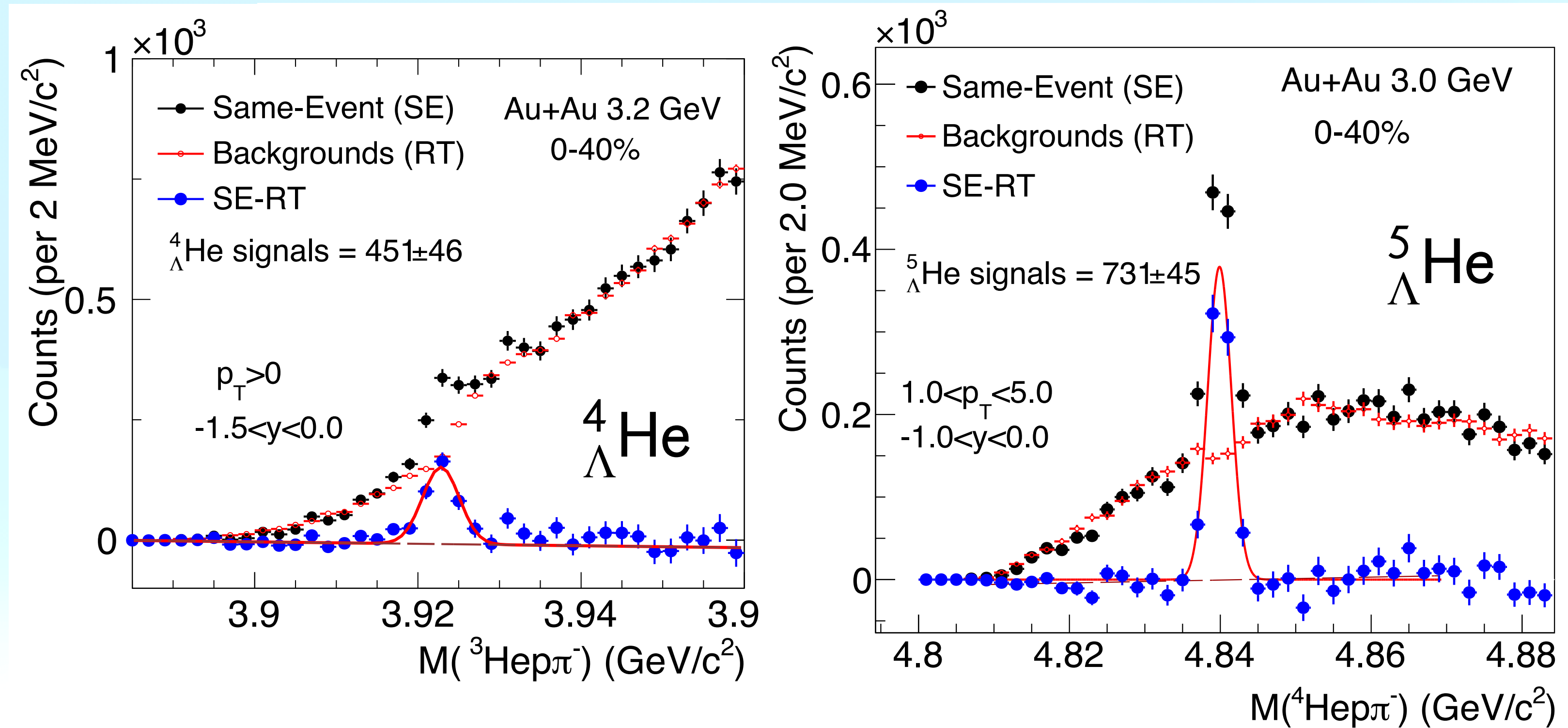
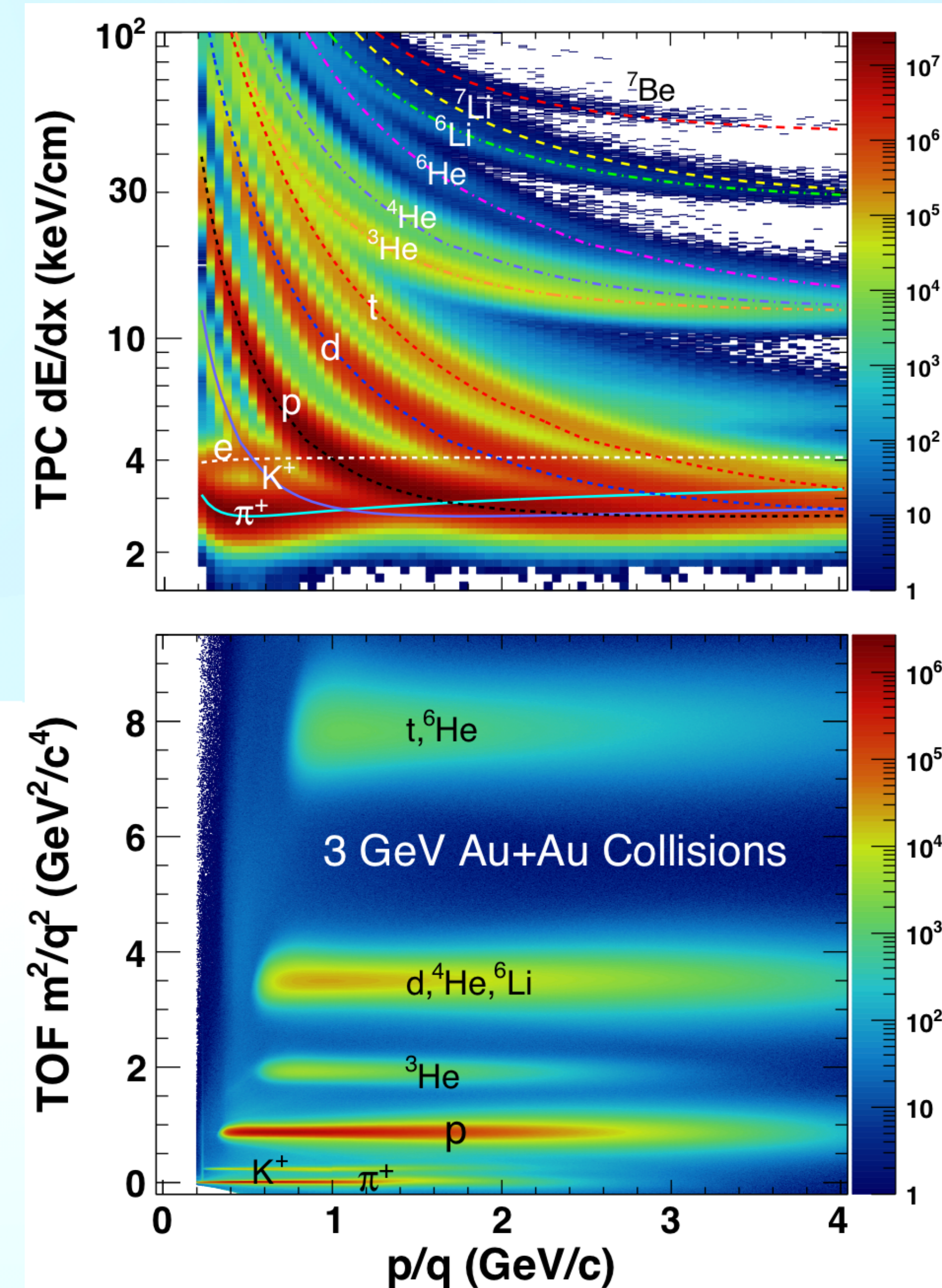
Janik Ditzel (ALICE), QM2025 3

STAR Beam Energy Scan II (2018-2021)



- Mid-rapidity coverage from 3-4.5 GeV in FXT

Particle Identification and Hypernuclei Reconstruction



- Hypernuclei reconstructed via weak decay channels
e.g.: $^4_\Lambda\text{He} \rightarrow ^3\text{He} + p + \pi^-$, $^5_\Lambda\text{He} \rightarrow ^4\text{He} + p + \pi^-$
- KFParticle* package utilizing information from error matrices applied to maximize significance

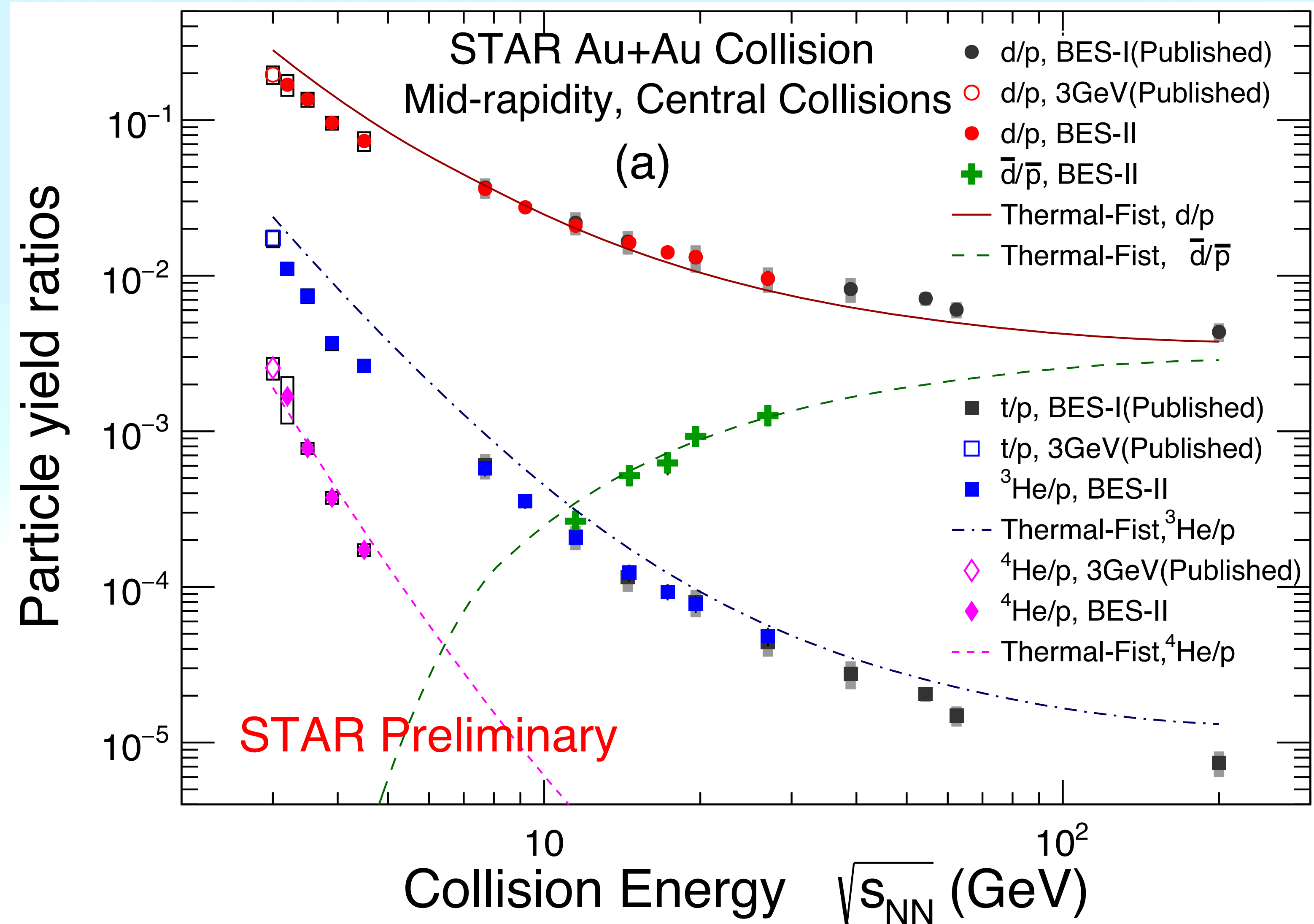
Nucl.Sci.Tech. 34 (2023) 10, 158

- TPC and (e)TOF used for particle identification

Nuclei-to-Nucleon Yield Ratios

STAR, Phys.Rev.C 110, 054911 (2024)
STAR, Phys.Rev.Lett. 130, 202301 (2023)

- d/p fairly well described by thermal model for $\sqrt{s_{NN}} \geq 7.7$ GeV, but overestimated ≤ 7.7 GeV
- t/p and ${}^3\text{He}/p$ overestimated by factor of 2 for all energies
- Only slight overestimation of ${}^4\text{He}/p$
- $\bar{d}/\bar{p}$ well described for $\sqrt{s_{NN}} \geq 7.7$ GeV



A=2,4: fairly well described; A=3: strongly overpredicted

Energy Dependence of Compound Ratio

STAR, Phys.Rev.C 110, 054911 (2024)

STAR, Phys.Rev.Lett. 130, 202301 (2023)

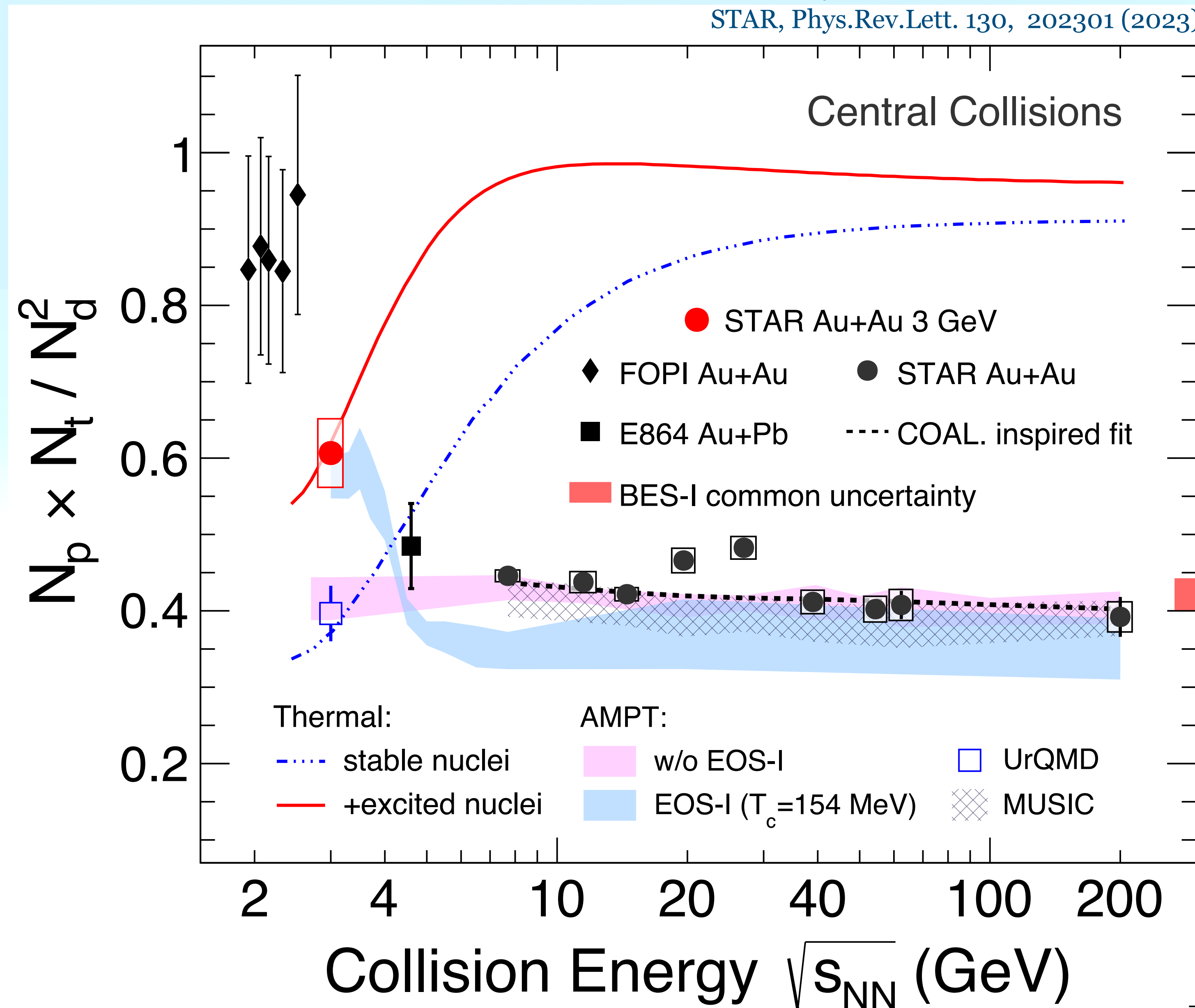
- $N_t N_p / N_d^2$ at $\sqrt{s_{NN}} = 3$ GeV is higher compared to 7.7 GeV and above

- Coalescence calculations incorporating **1st order phase transition** describes the ratio at 3 GeV

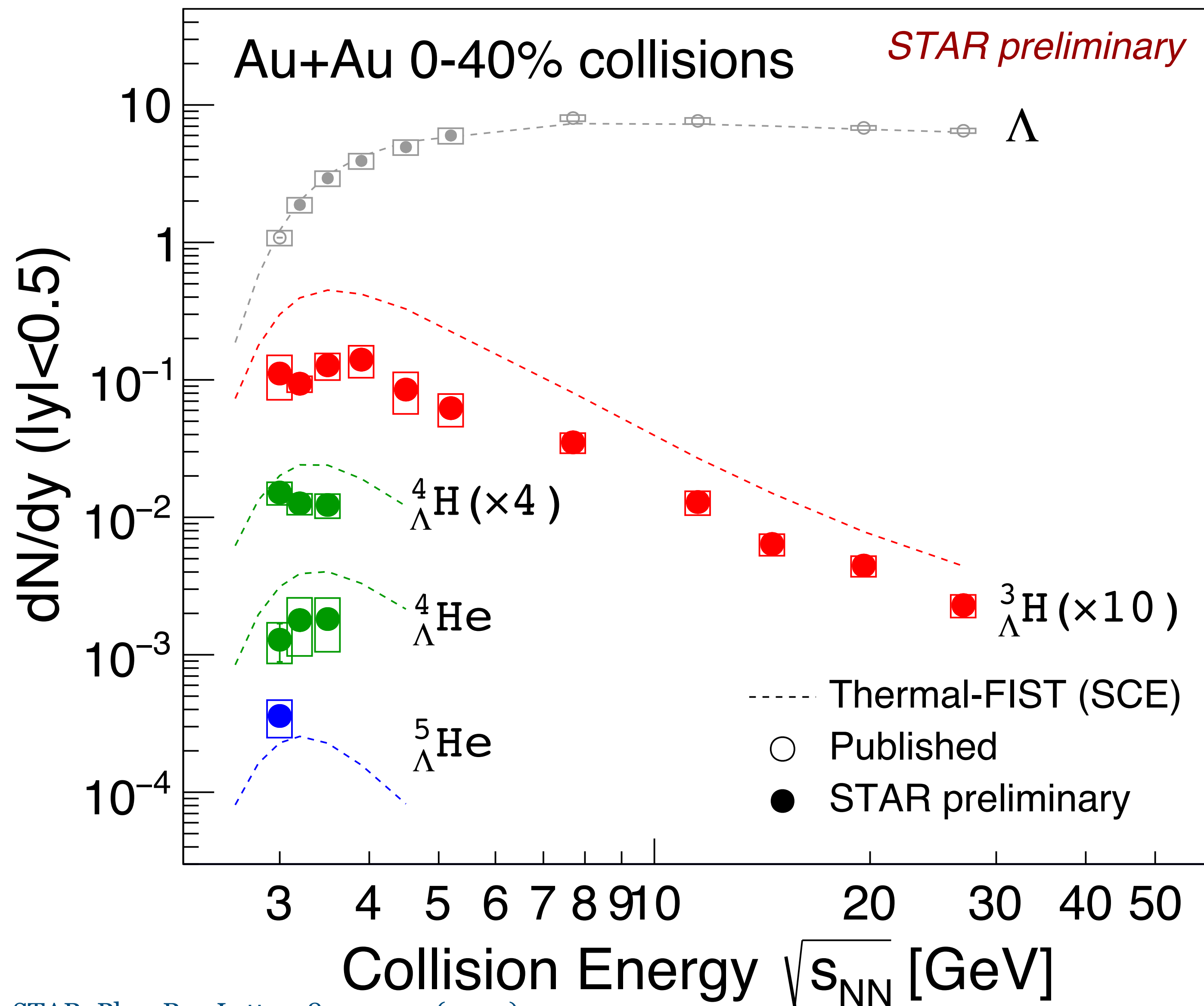
arXiv:2205.11010

- Enhancement could be partially arising from **decay of unstable nuclei**

e.g.: ${}^4\text{He}^* \rightarrow p + t$



Hypernuclei Excitation Functions



- Thermal model

- strongly overestimates ${}^3_{\Lambda}\text{H}$
- slightly overestimates ${}^4_{\Lambda}\text{H}$, ${}^4_{\Lambda}\text{He}$
- slightly underestimates ${}^5_{\Lambda}\text{He}$

**A=3 strongly overpredicted,
A=4 show milder deviations,
similar to the trend for
normal nuclei**

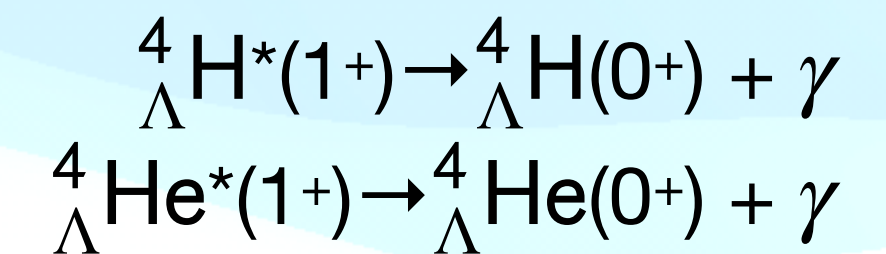
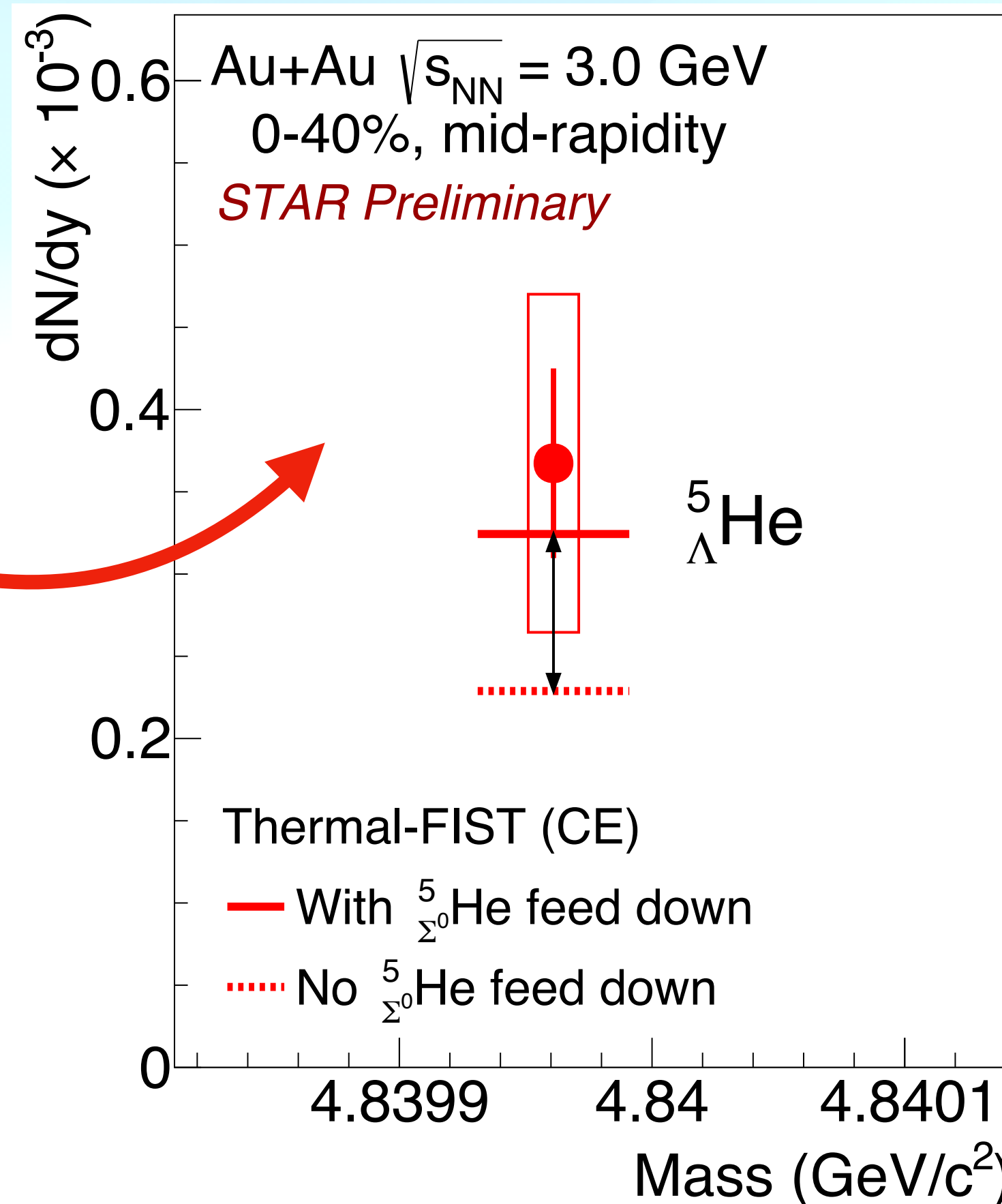
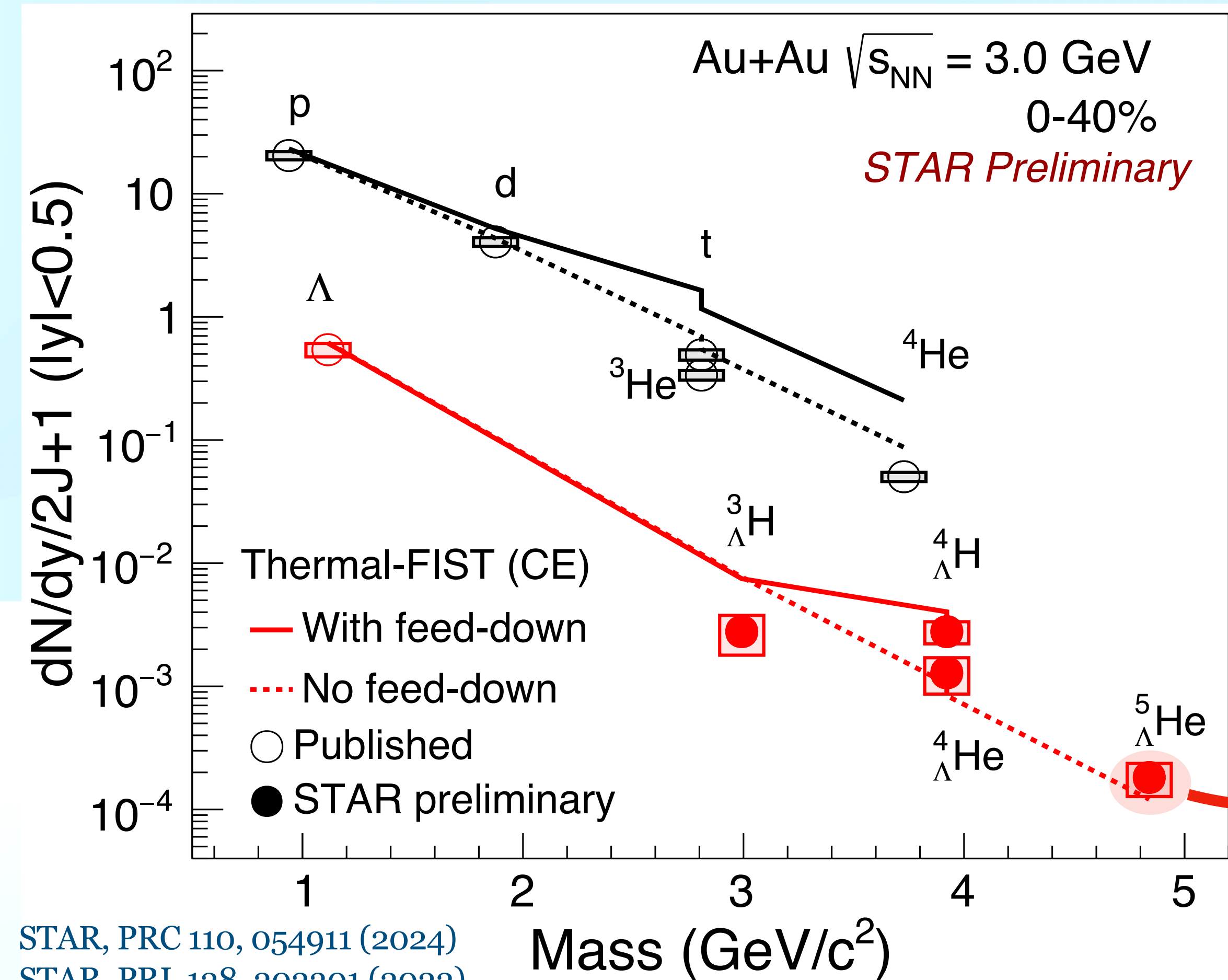
STAR, Phys.Rev.Lett. 128, 202301 (2022)

STAR, Phys.Rev.C 102, 034909 (2020)

STAR, JHEP 10, 139 (2024)

(Hyper)nuclei Mid-Rapidity Yields at $\sqrt{s_{NN}} = 3$ GeV

- Inclusion of excited hypernuclear state production is crucial for models
- Thermal model overestimate ${}^4_{\Lambda}\text{H}$, ${}^4_{\Lambda}\text{He}$ after including feed-down from excited states



Hint of possible feed-down from

$${}^5_{\Sigma^0}\text{He} \rightarrow {}^5_{\Lambda}\text{He} + \gamma$$

J. Johnstone et al., JPG 8, L105 (1982)

STAR, PRC 110, 054911 (2024)
STAR, PRL 128, 202301 (2022)
STAR, JHEP 10 (2024) 139

Probing the Hypertriton (${}^3_{\Lambda}\text{H}$) Wave Function

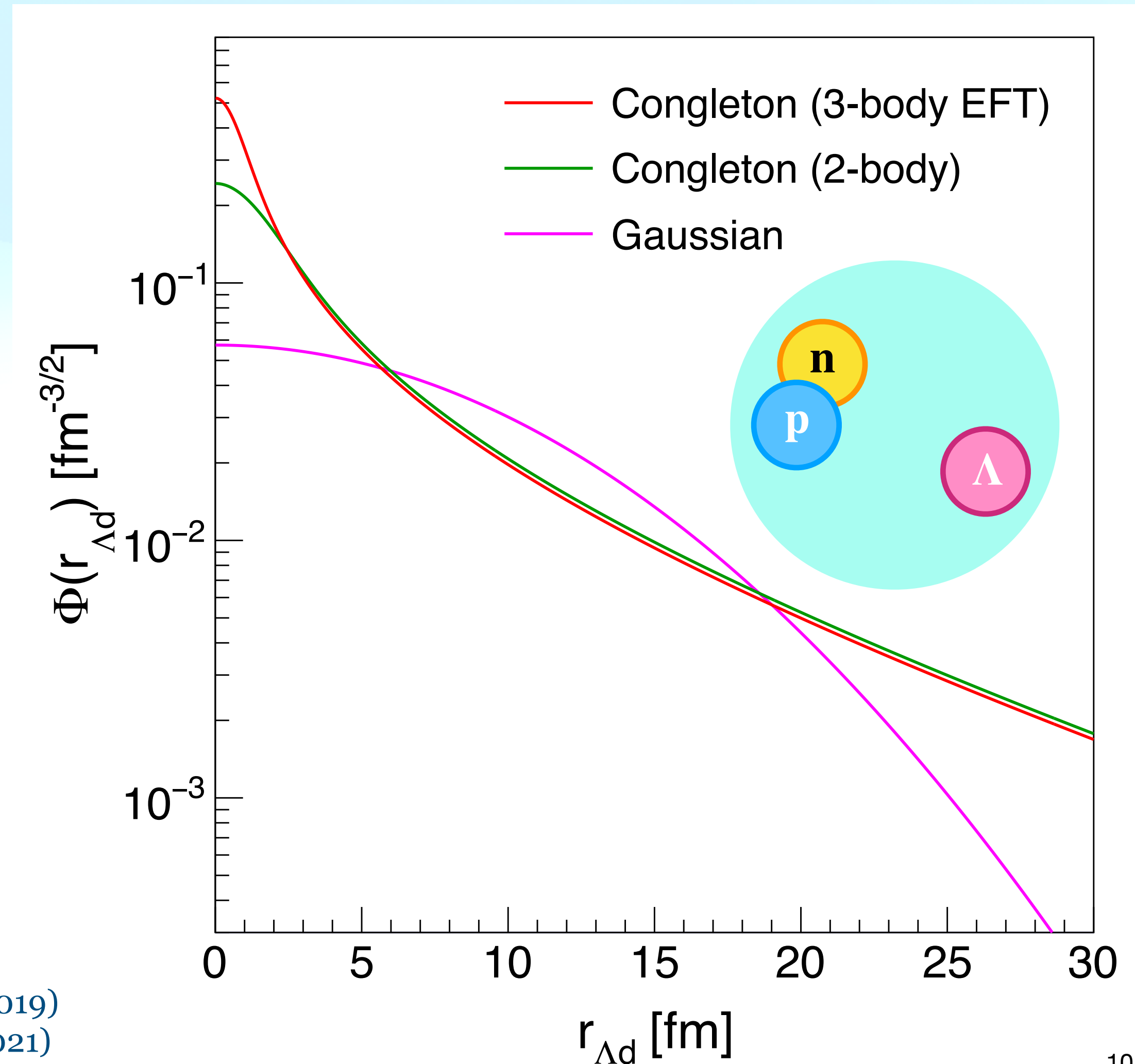
- **Wigner-function coalescence**: coalescence probability given by overlap between ${}^3_{\Lambda}\text{H}$ wave function and nucleon/hyperon distributions at freeze out

- ${}^3_{\Lambda}\text{H}$: loosely bound hypernuclei
($B_{\Lambda} = 0.164 \pm 0.043$ MeV)

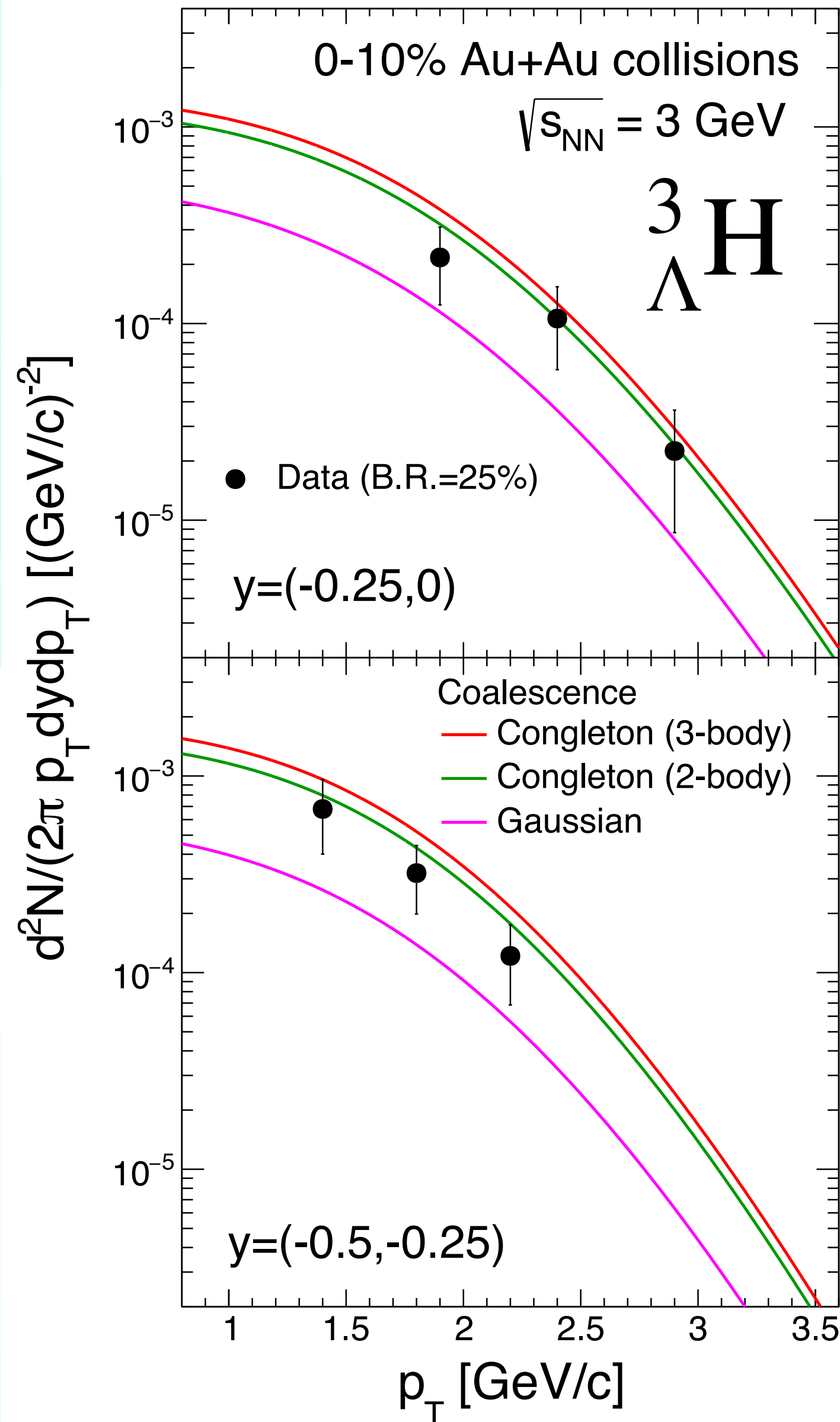
- Gaussian: $\Phi_{{}^3_{\Lambda}\text{H}}(r_{pn}, r_{\Lambda}) = \left(\frac{1}{3\pi^2 b_{pn}^2 b_{\Lambda}^2} \right)^{\frac{3}{4}} e^{-\frac{r_{pn}^2}{4b_{pn}^2} - \frac{r_{\Lambda}^2}{3b_{\Lambda}^2}}$
 $b_{\Lambda} = 7.2\text{fm}$ Phys.Rev.C 103, 014907 (2021)

- Congleton: $\hat{\Phi}_{{}^3_{\Lambda}\text{H}(d\Lambda)}(q) = A \frac{e^{-\frac{q^2}{Q_{\Lambda}^2}}}{q^2 + \alpha_{\Lambda}^2}$
 - 2-body model of ${}^3_{\Lambda}\text{H}$ J.Phys.G 18, 339-357 (1992)
($Q_{\Lambda}, \alpha_{\Lambda}$) = (1.17, 0.068) fm⁻¹

- 3-body effective field theory
($Q_{\Lambda}, \alpha_{\Lambda}$) = (2.5, 0.068) fm⁻¹ Phys.Rev.C 100, 034002 (2019)
Phys.Rev.C 103, 014907 (2021)



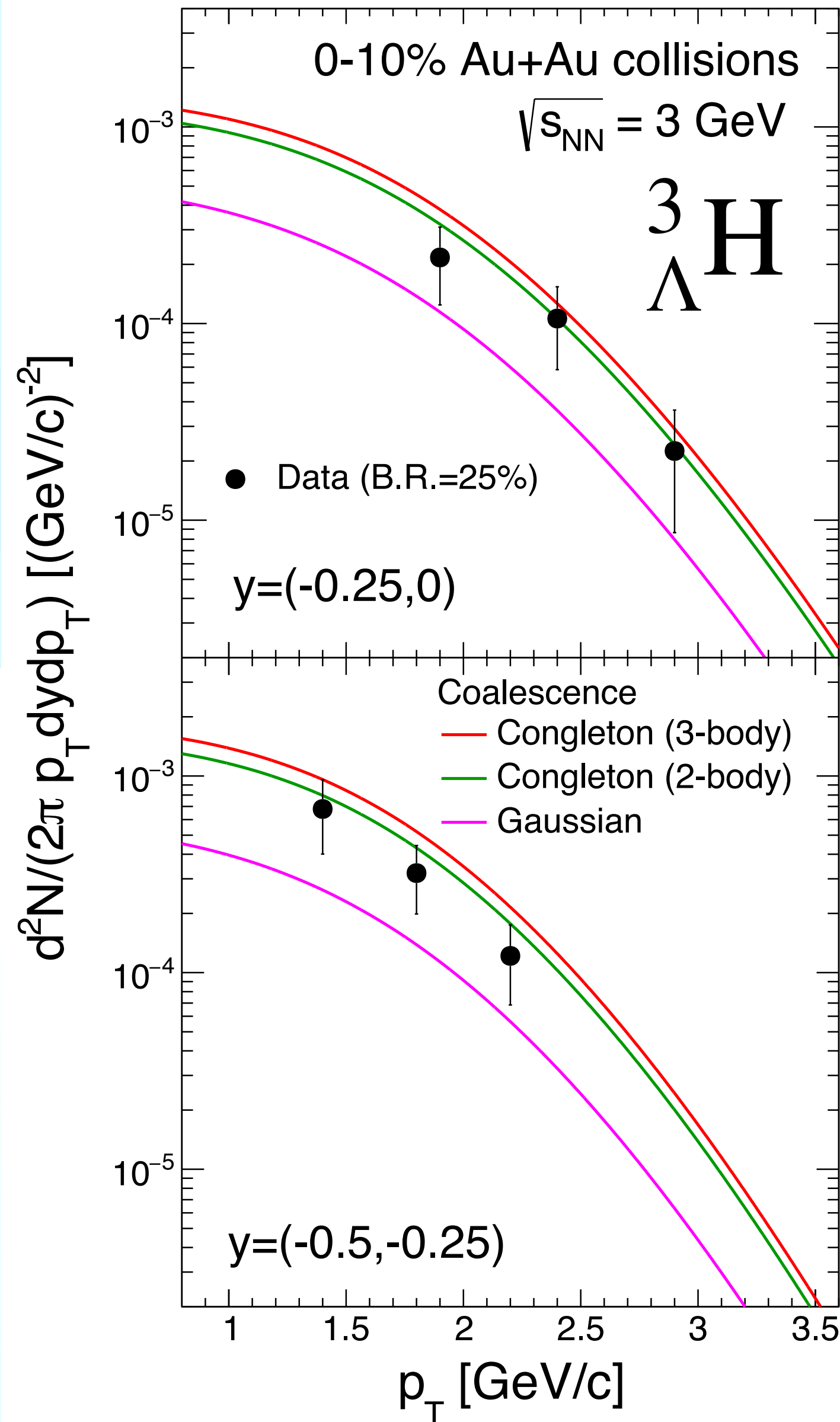
${}^3_{\Lambda}\text{H}$ Spectra at 3 GeV



- Coalescence using Congleton wave function describe ${}^3_{\Lambda}\text{H}$ spectra at $\sqrt{s_{NN}} = 3 \text{ GeV}$

STAR, Phys.Rev.Lett. 128 (2022) 20, 202301
 Coalescence: arXiv: 2510.06758

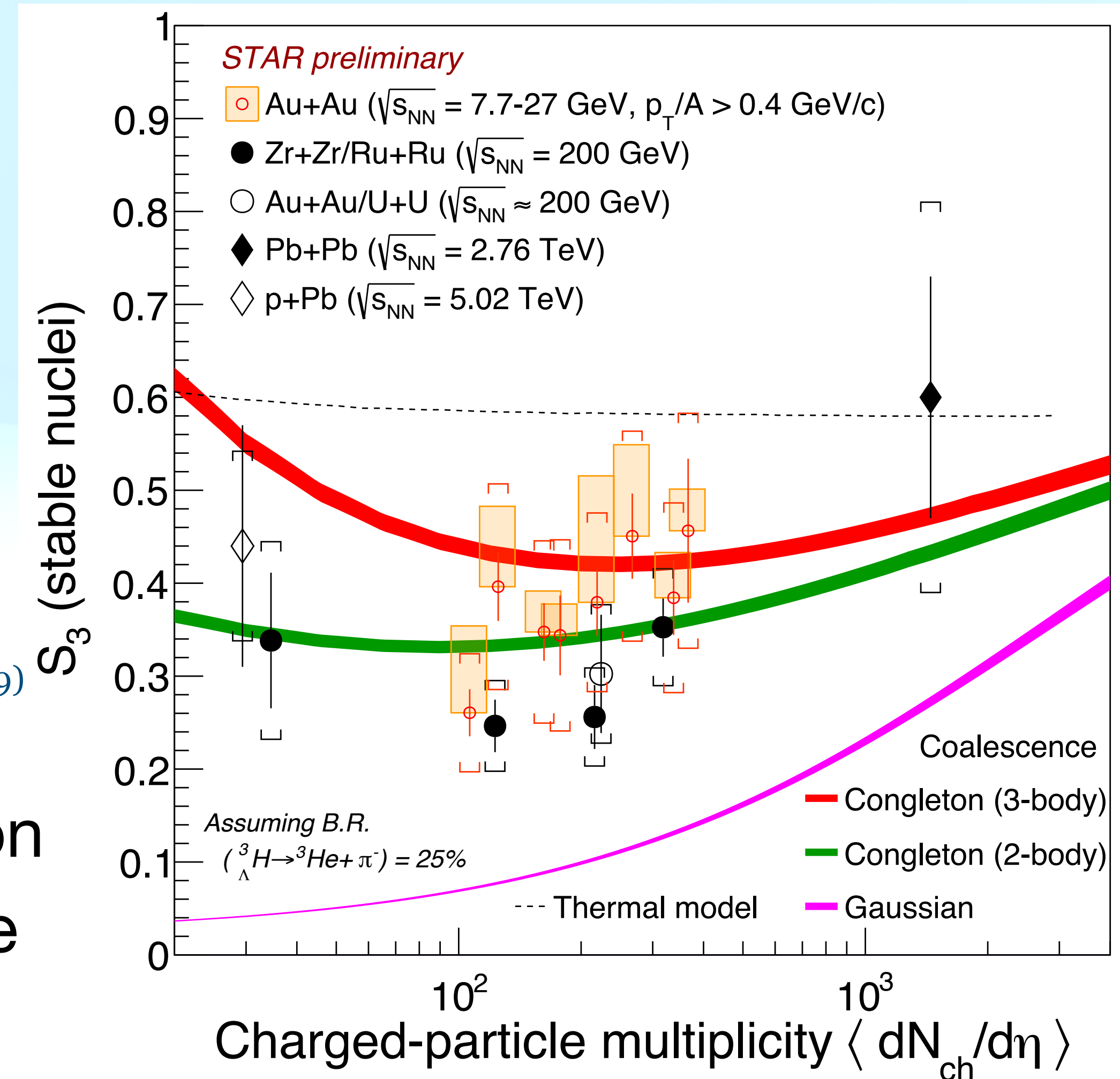
Strangeness Population Factor $S_3 = \frac{{}^3_{\Lambda}\text{H}}{{}^3\text{He} \times \frac{\Lambda}{p}}$



- Coalescence using Congleton wave function describe ${}^3_{\Lambda}\text{H}$ spectra at $\sqrt{s_{NN}} = 3 \text{ GeV}$

STAR, Phys.Rev.Lett. 128 (2022) 20, 202301
 STAR, Nature 2024, 77 (2024)
 ALICE, Phys.Rev.Lett. 128, 252003 (2022)
 ALICE, Phys.Lett.B 754 (2016) 360-372
 Coalescence: arXiv: 2510.06758
 Thermal: Comput.Phys.Commun. 244, 295 (2019)

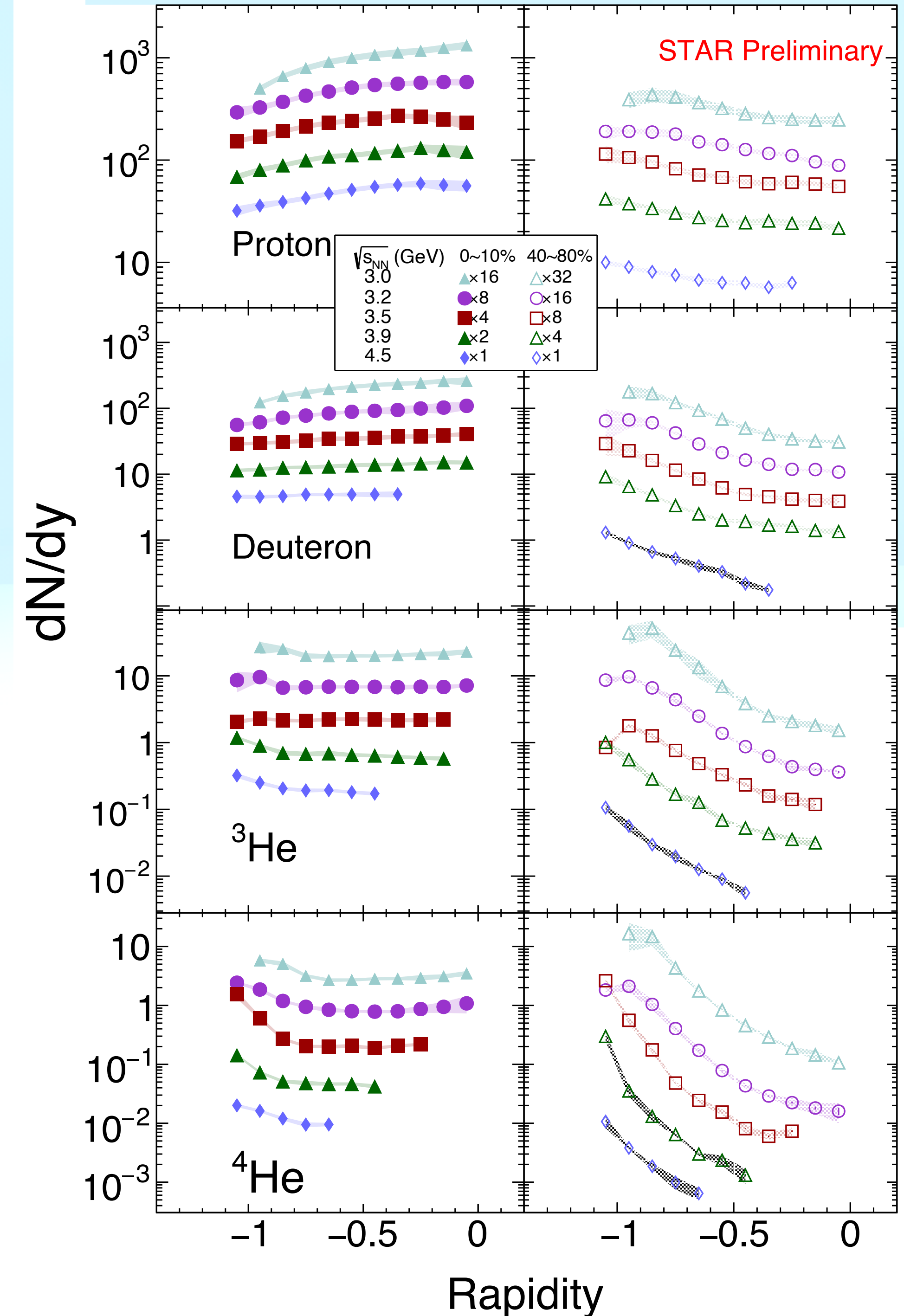
- Coalescence using Congleton wave function describe S_3 across wide range of energies



Sensitivity of ${}^3_{\Lambda}\text{H}$ yields to short-distance structure of the wave function

Rapidity Dependence of Nuclei Production

- Peak structures at target rapidity are more prominent
 - in more peripheral collisions
 - for nuclei with higher mass
- Primarily due to the **fragmentation of spectators**



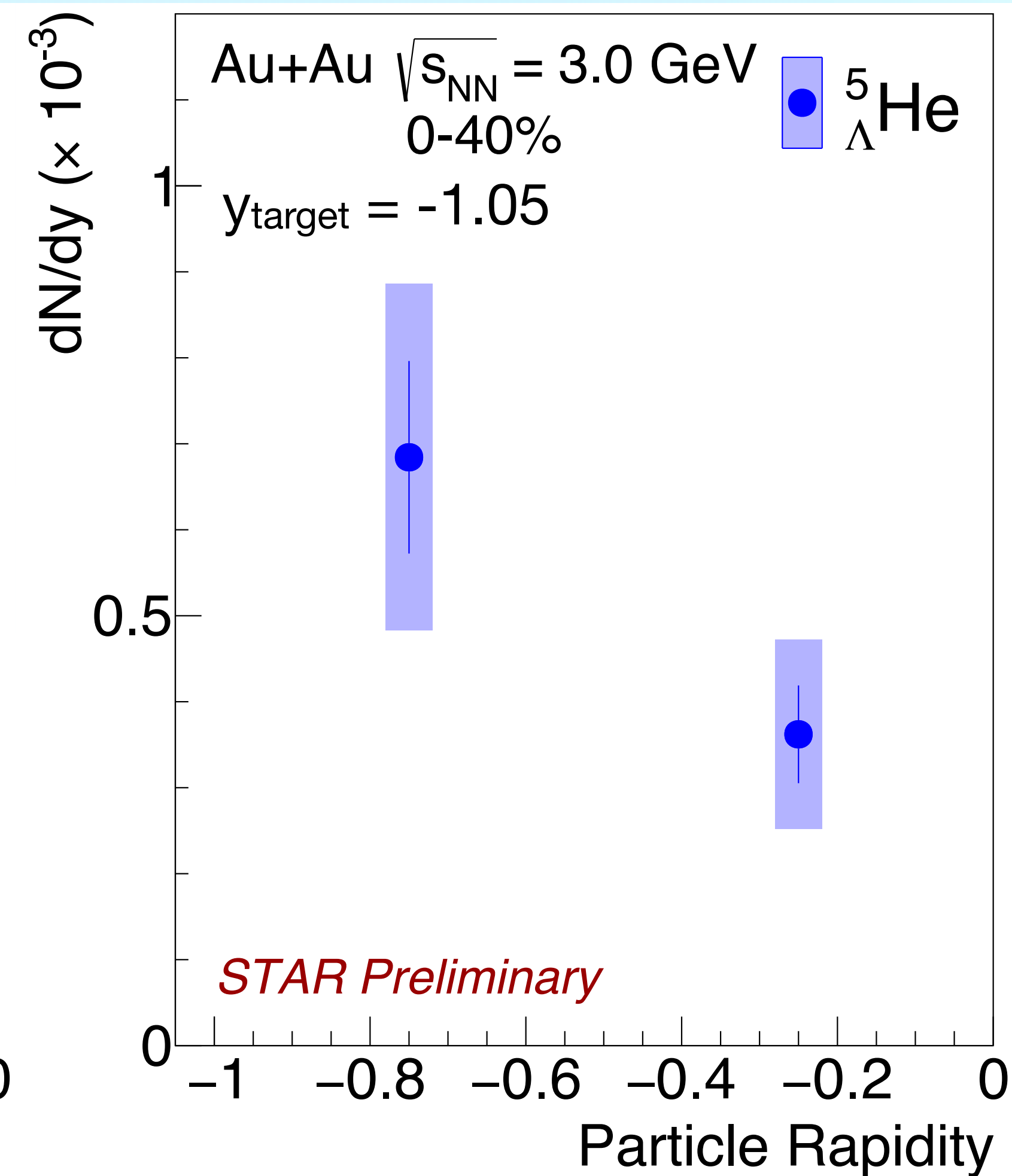
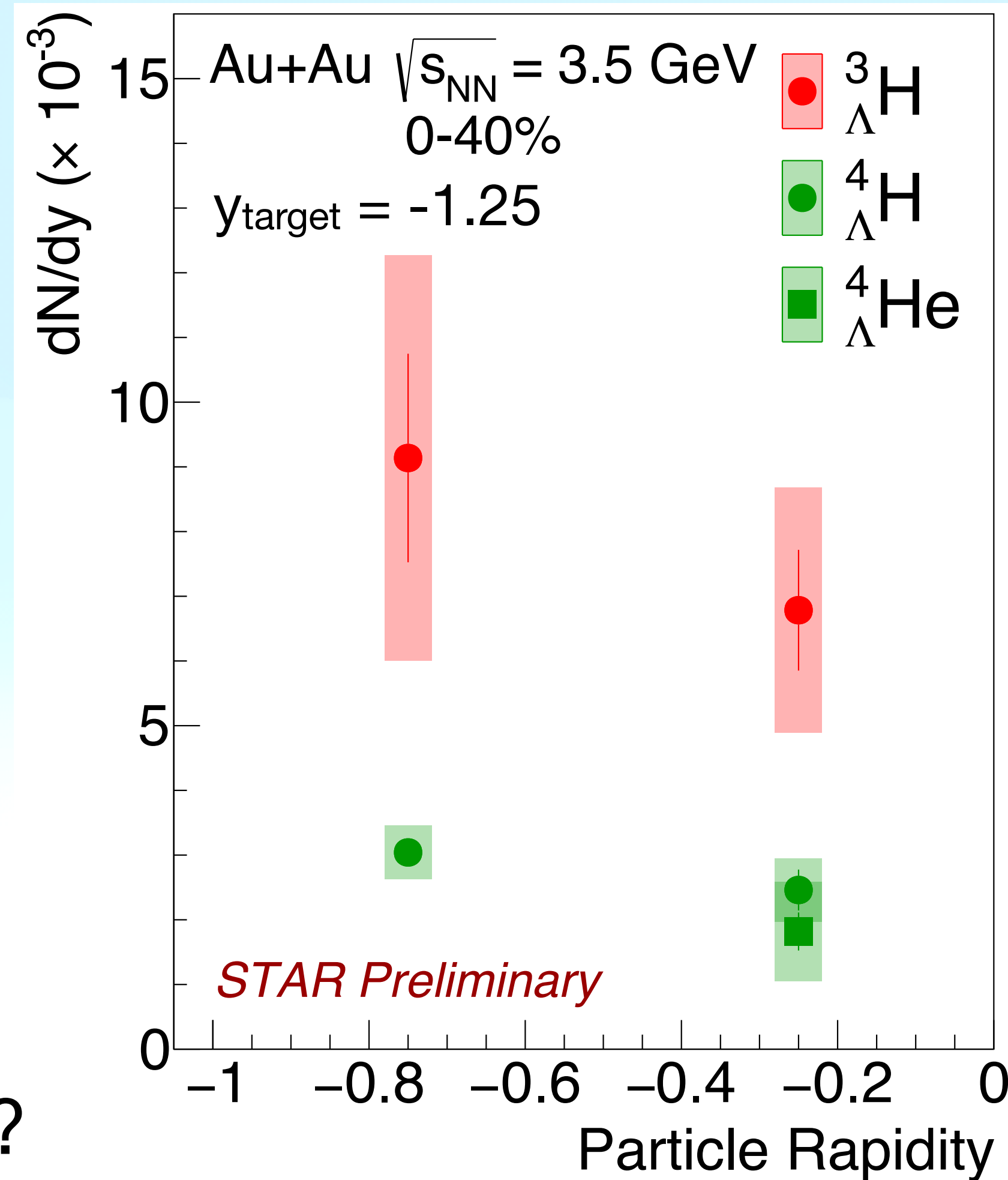
Rapidity Dependence of Hypernuclei Production

- Significant hypernuclei production near target rapidity

Spectators also play an important role in hypernuclei production

- Not “thermal”-like
- Coalescence? or
- Λ capture + fragmentation?

Phys.Rev.C 76, 024909 (2007)



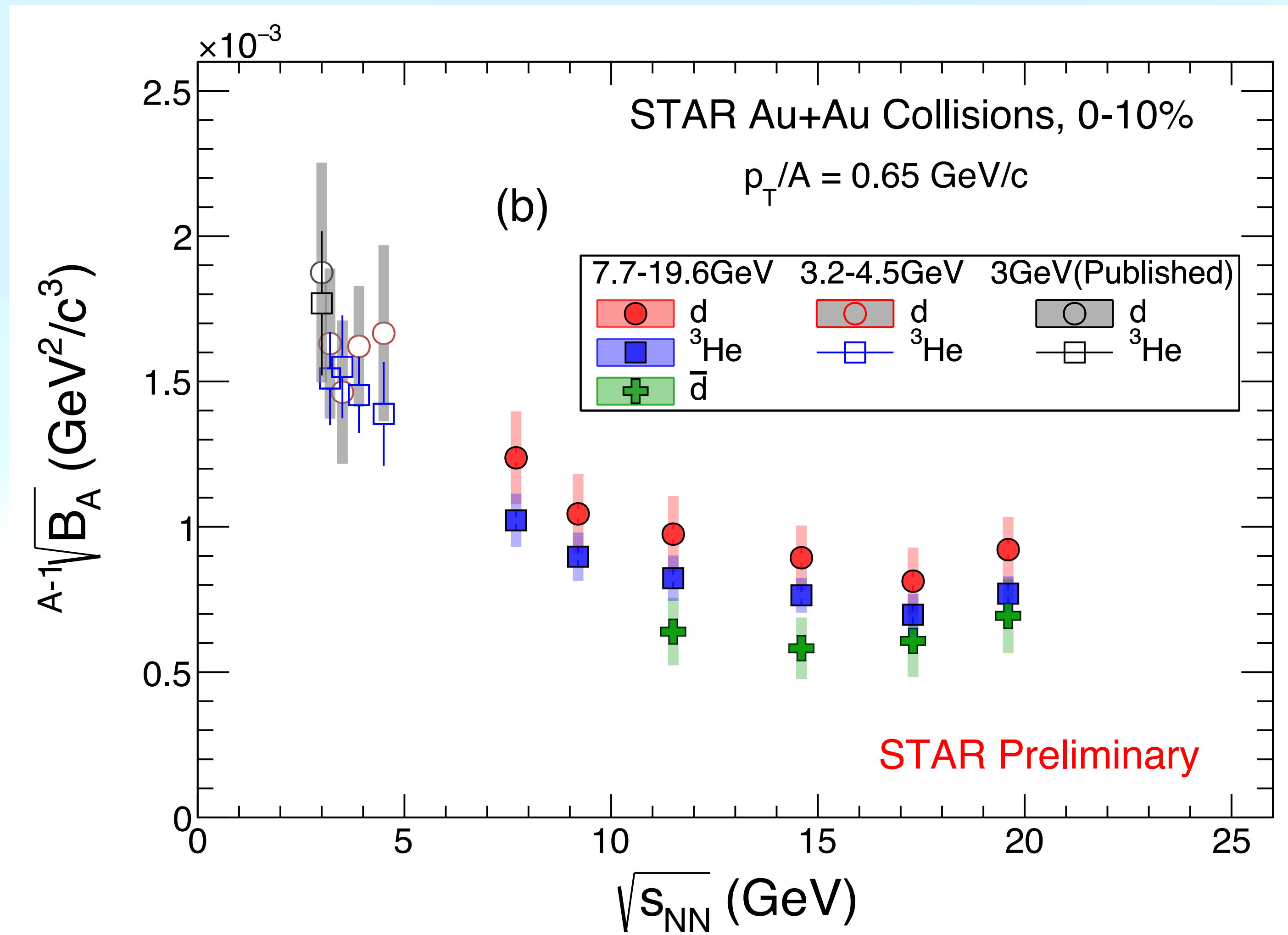
- More differential measurements underway!

Coalescence Parameters (B_A)

$$B_A = \frac{E_A \frac{d^3 N_A}{d^3 p_A}}{\left(E_p \frac{d^3 N_p}{d^3 p_p}\right)^Z \left(E_n \frac{d^3 N_n}{d^3 p_n}\right)^{A-Z}}$$

$p_A/A = p_p = p_n$

- Coalescence parameters (B_A) reflect the coalescence probability
- $A^{-1}\sqrt{B_A}$ decrease with increasing energy
 - effective volume increases with increasing energy



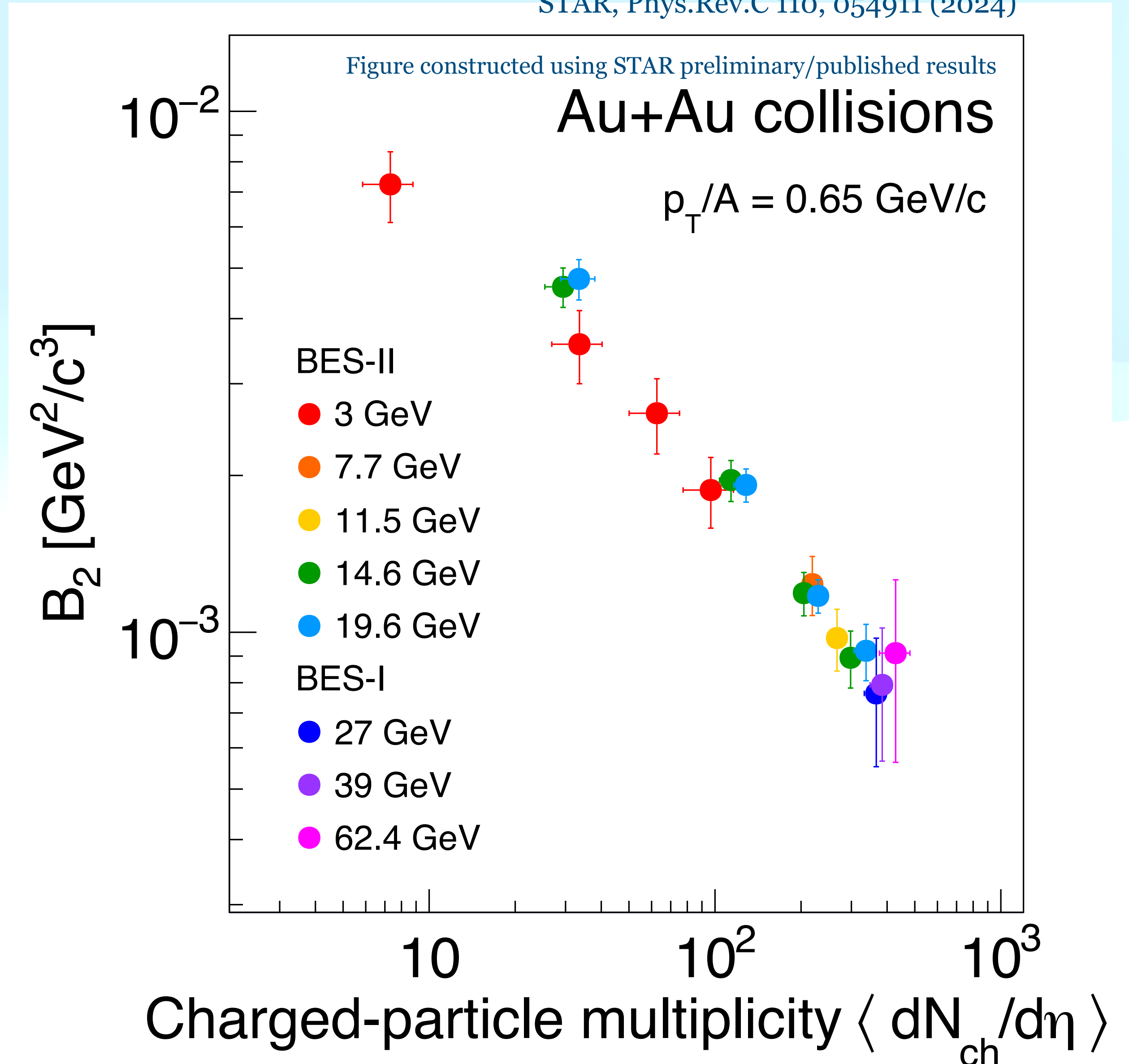
STAR, Phys.Rev.C 110, 054911 (2024)

Coalescence Parameters (B_A)

- $dN_{\text{ch}}/d\eta$ is a proxy of the volume of the system
- B_2 decreases with $dN_{\text{ch}}/d\eta$
 - Different collision energies approx. follow a global trend

Probability of coalescence is anti-correlated with the size of the system

STAR, Phys. Rev. C 99, 64905 (2019)
STAR, Phys. Rev. C 110, 054911 (2024)

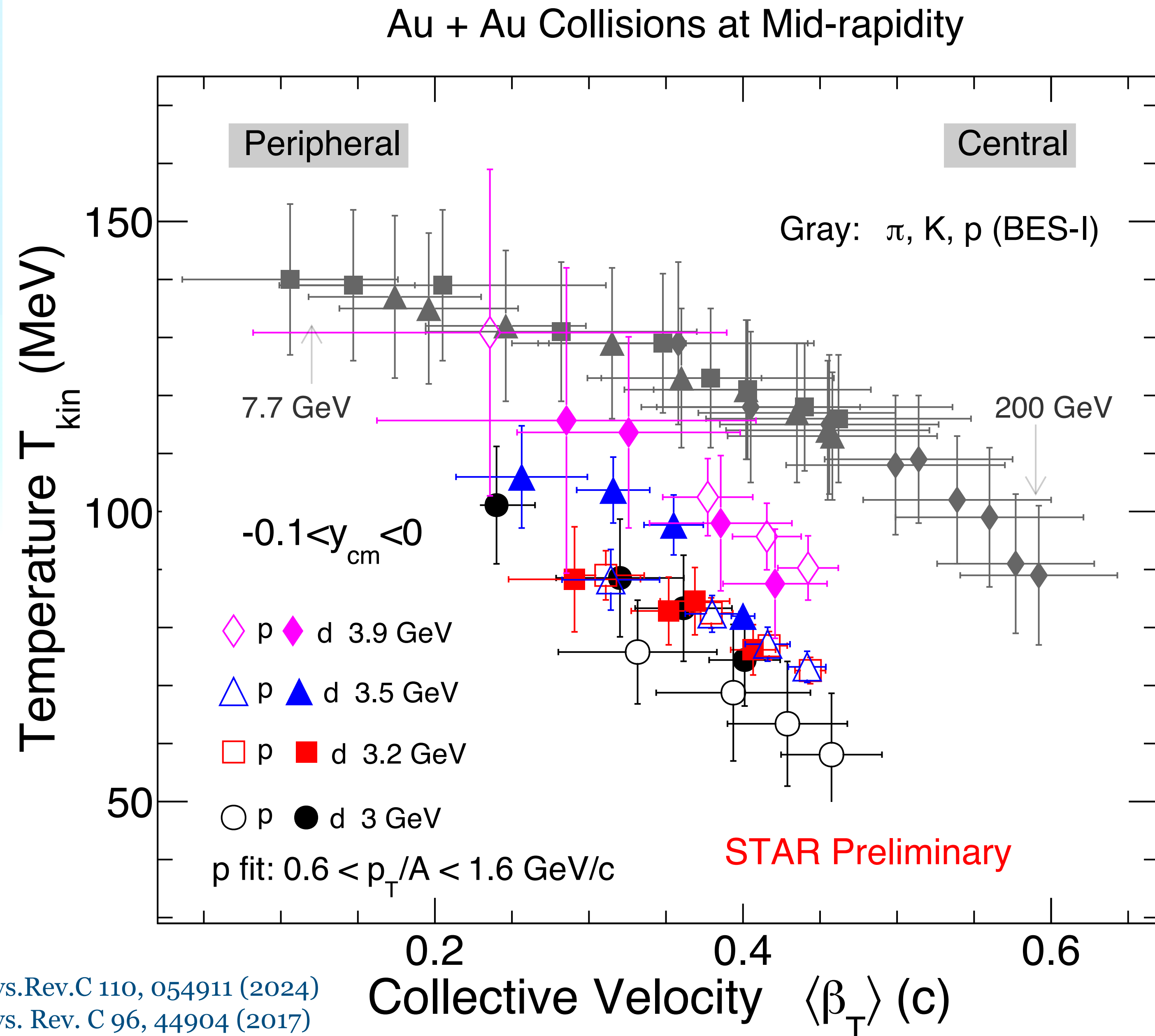


Deuteron Kinetic Freeze-out Parameters

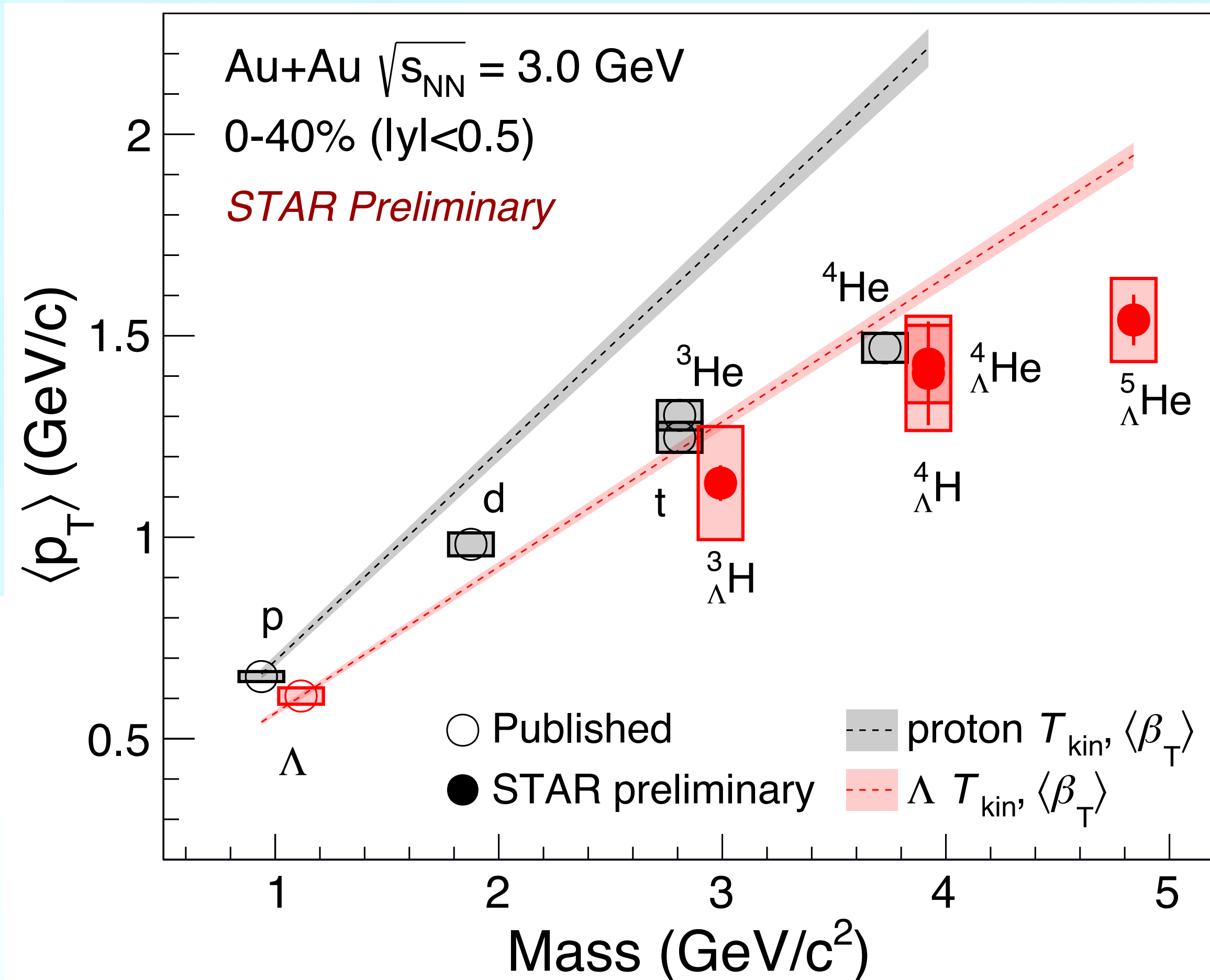
- Hydrodynamic-inspired Blast-Wave model: assumes particles are emitted **thermally** from an expanding source with a common T_{kin} and $\langle \beta_T \rangle$
- Fit p and d spectra separately to obtain kinetic freeze-out parameters
- Distinct kinetic freeze-out surfaces for p and d

Deuterons deviate from the hydrodynamical blast-wave picture

STAR, Phys.Rev.C 110, 054911 (2024)
STAR, Phys. Rev. C 96, 44904 (2017)



Mean Transverse Momentum $\langle p_T \rangle$ at 3 GeV



- Vary mass to construct blast-wave predictions using $p(\Lambda)$ freeze-out parameters

	p	Λ
$\langle \beta_T \rangle$ (c)	0.43	0.33
T_{kin} (GeV)	0.065	0.076

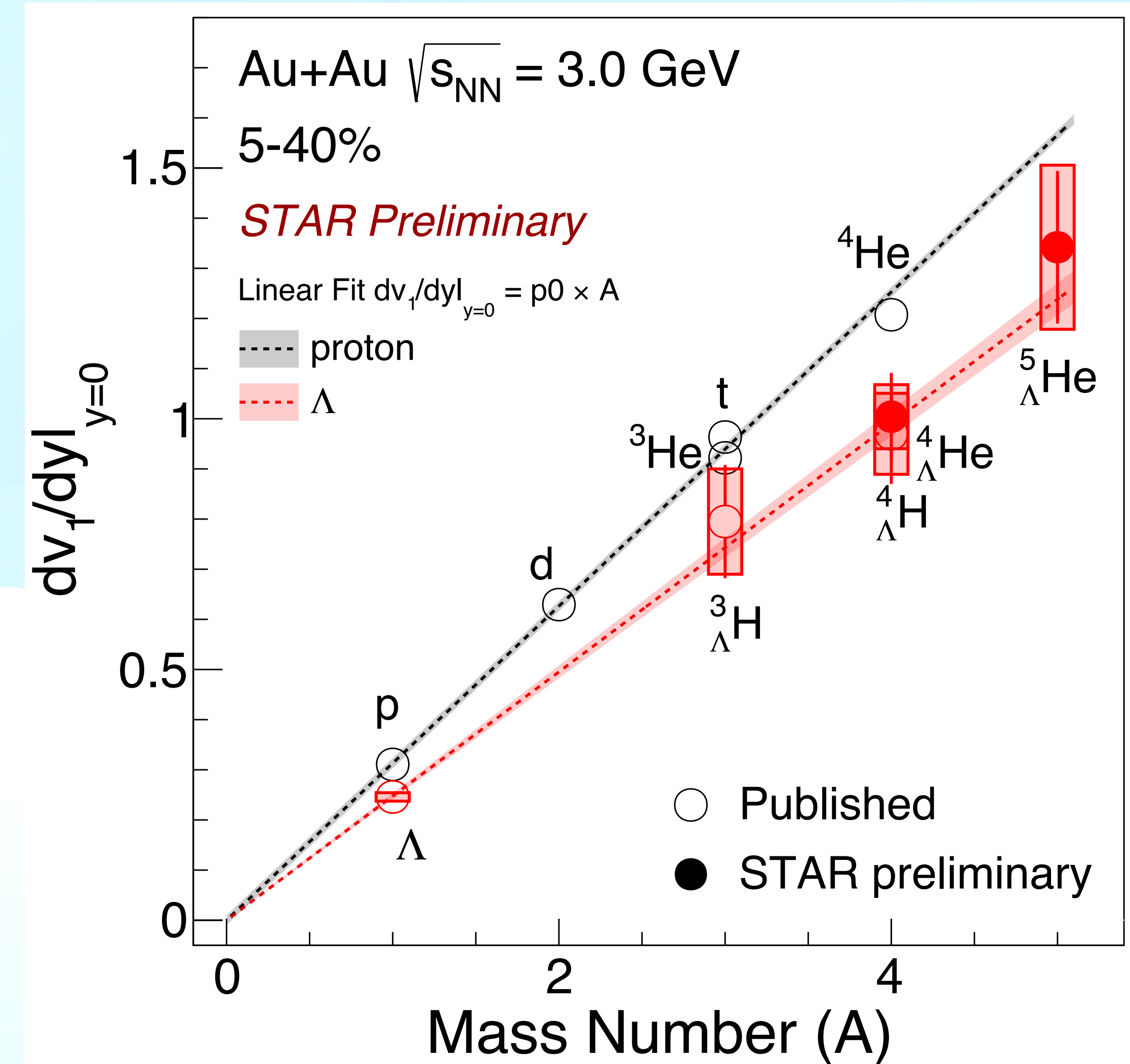
E. Schnedermann et al., PRC 48, 2462 (1993)
ALICE, PLB 794, 50–63 (2019)

- Hypernuclei closer to Λ blast-wave ansatz

Kinematics of (hyper)nuclei deviate from the collective expansion of the bulk matter

STAR, PRC 110, 054911 (2024)
STAR, PRL 128, 202301 (2022)
STAR, JHEP 10 (2024) 139

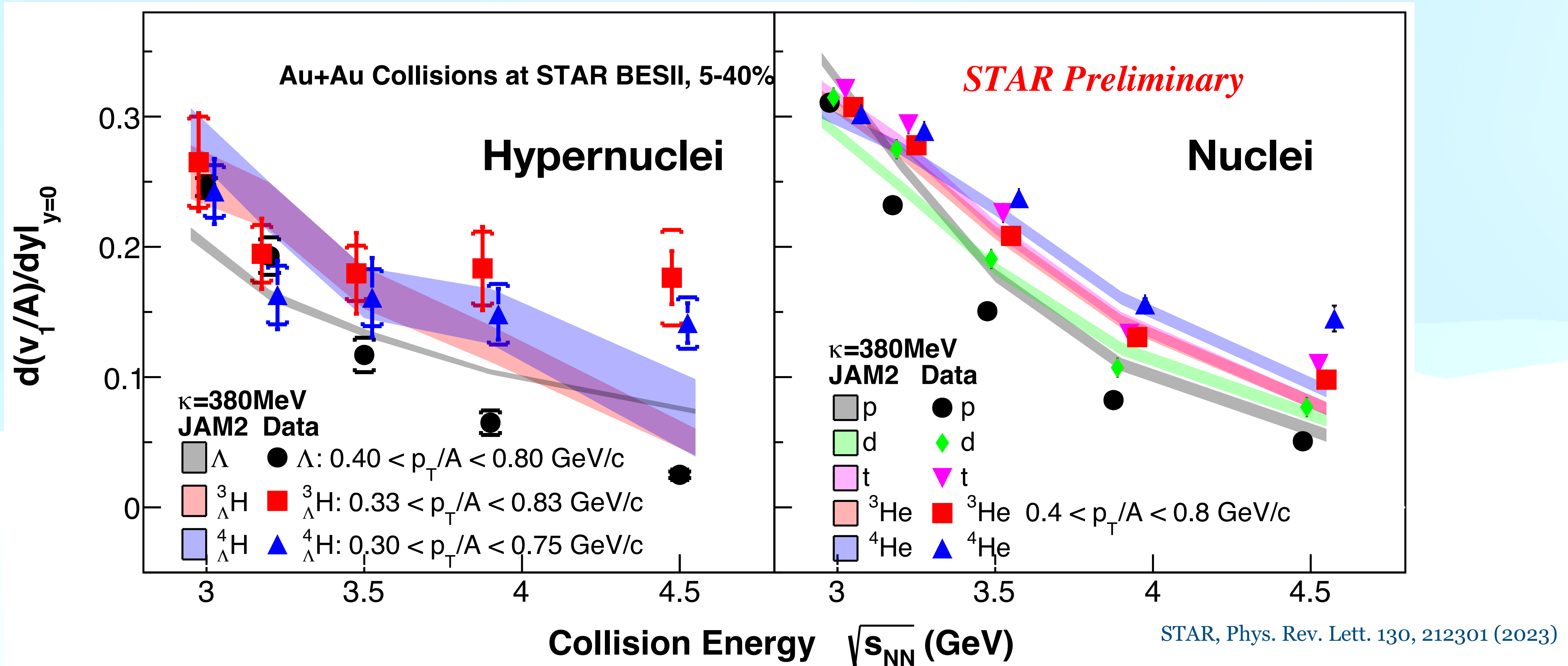
Mass Number Scaling of Directed Flow at 3 GeV



- Mid-rapidity v_1 slope of nuclei increase linearly with atomic mass number A
- Hypernuclei v_1 slope
 - systematically lower than light nuclei of similar A
 - compatible with Λ atomic mass number scaling, up to $A=5$

A scaling of v_1 , qualitatively consistent with coalescence

Collision Energy Dependence of (Hyper)nuclei Directed Flow



- Larger hypernuclei flow at 4.5 GeV than expected by JAM + coalescence

Agreement with A scaling deteriorates as the collision energy increases

Summary

- Nuclei (d , t , ^3He , ^4He) and hypernuclei ($^3_{\Lambda}\text{H}$, $^4_{\Lambda}\text{H}$, $^4_{\Lambda}\text{He}$, and $^5_{\Lambda}\text{He}$) yield and directed flow measurements from STAR BES-II at $\sqrt{s_{\text{NN}}} = 3\text{-}27$ GeV

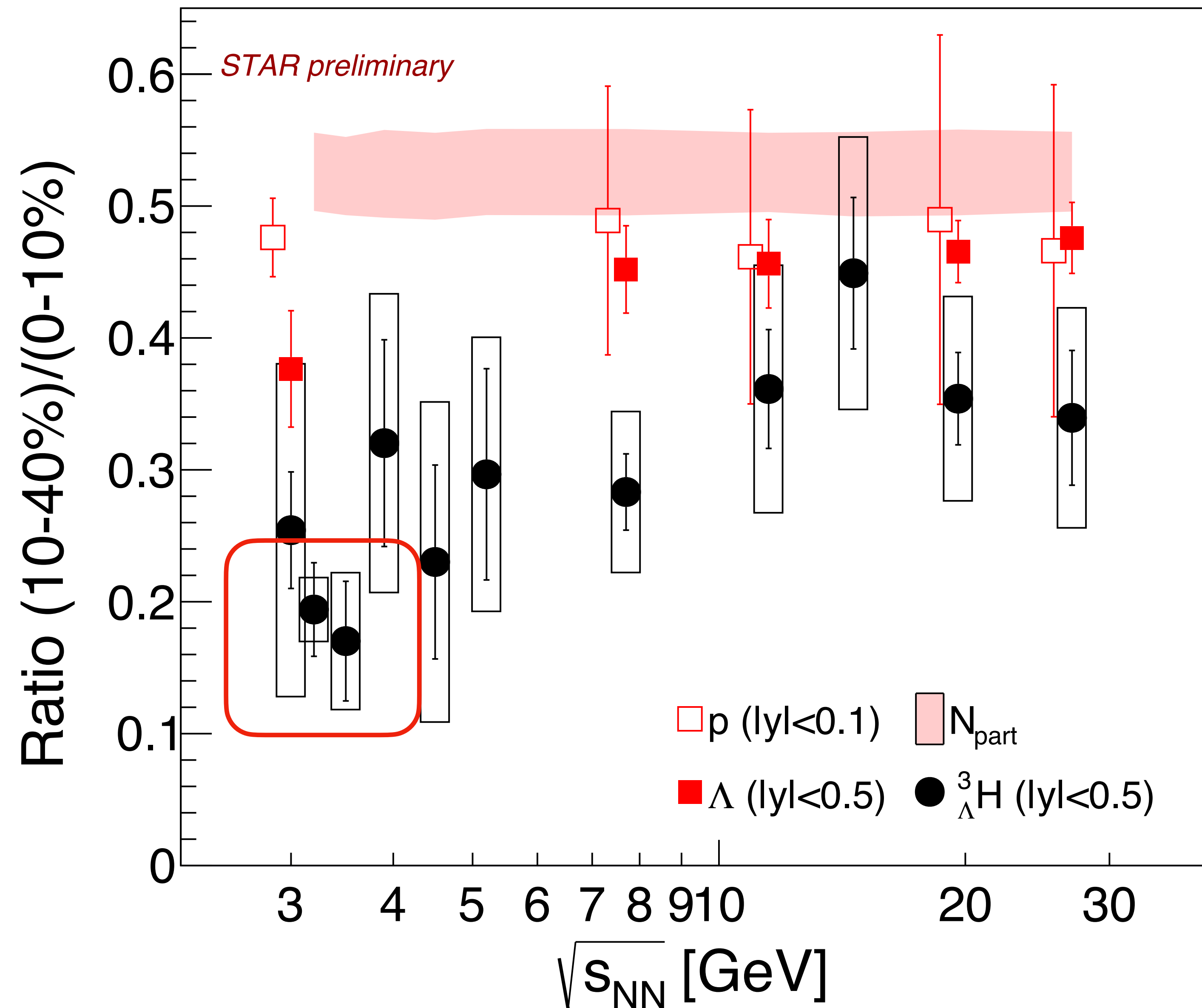
- Thermal model
 - Describes $A=2, 4, 5$ fairly well
 - **Strongly overpredicts $A=3$, $N_t N_p / N_d^2$, S_3**
- Blast-wave model
 - **Overpredicts $\langle p_T \rangle$ of (hyper)nuclei**
- Coalescence model
 - Describe $N_t N_p / N_d^2$ fairly well
 - **Describe S_3 well with Congleton wave function**
 - Qualitatively consistent with v_1 measurements, tension with hypernuclei data at 3.9-4.5 GeV

- **Data favors coalescence as the dominant production mechanism at midrapidity**
- **Coalescence calculations indicate sensitivity of $^3_{\Lambda}\text{H}$ yields to its wave function in low multiplicity environments $\rightarrow$ opportunities to constrain $YN(N)$ interactions**

Thank you for listening!!

Backup slides follow

Centrality Dependence

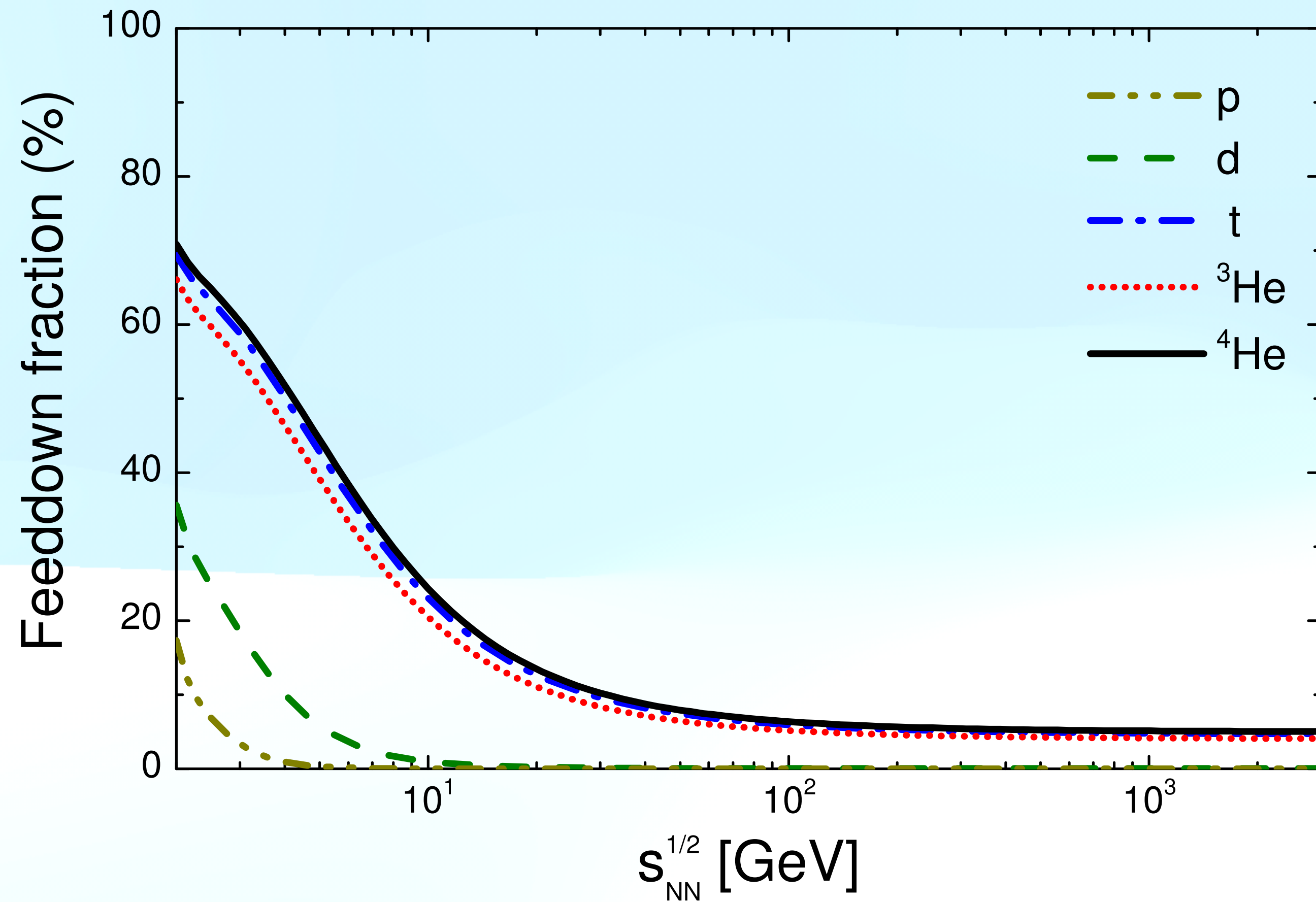


- Below $\sqrt{s_{NN}} \sim 4$ GeV, yield ratio (10-40%)/(0-10%) for ${}^3_{\Lambda}H$ is smaller compared to proton, Λ , and N_{part}

Suppression of ${}^3_{\Lambda}H$ production in more peripheral collisions compared to proton and Λ below $\sqrt{s_{NN}} \sim 4$ GeV

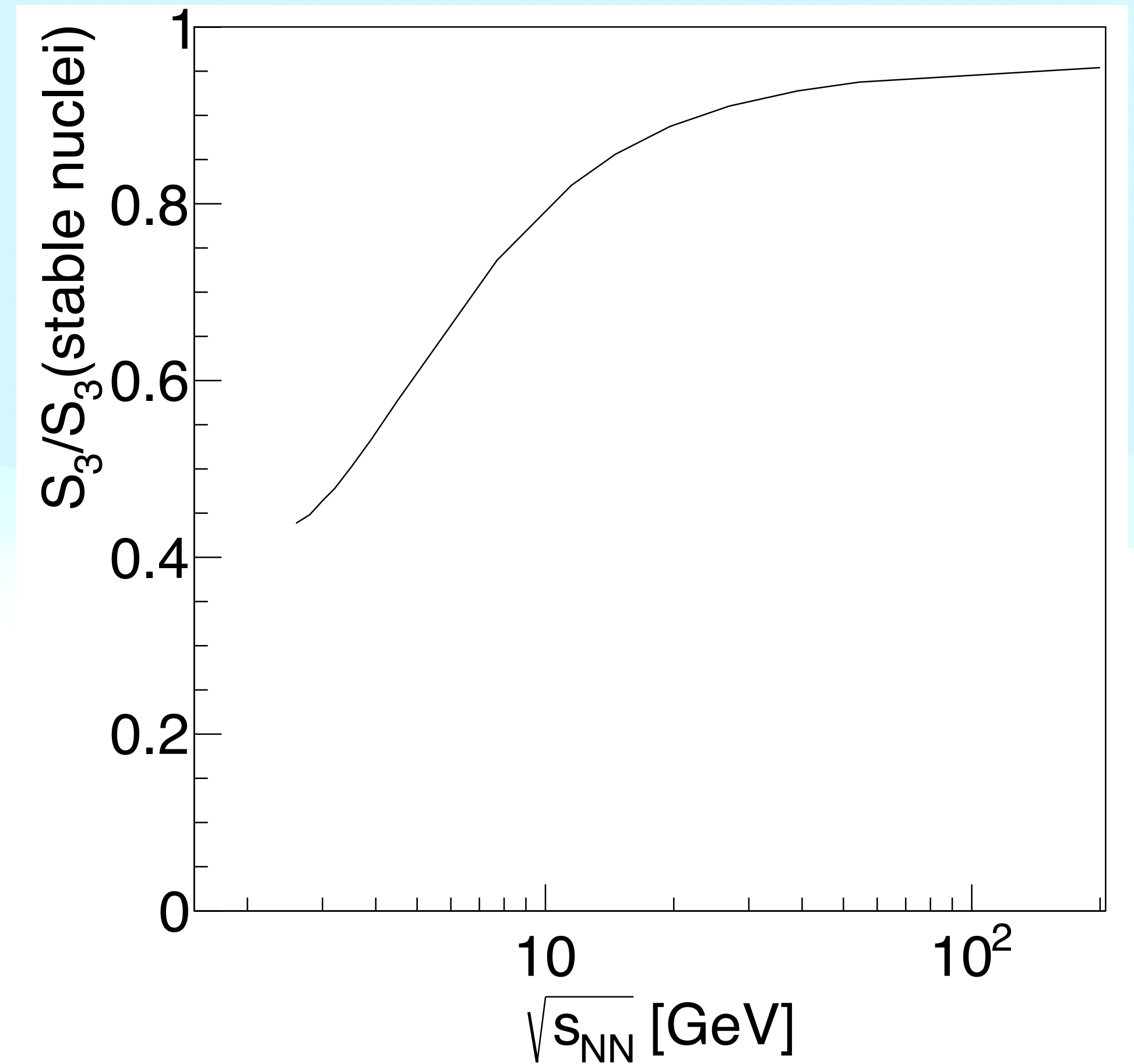
- Steep decline may be related to coalescence production and smaller source size in more peripheral collisions

Feed-down from Unstable Nuclei



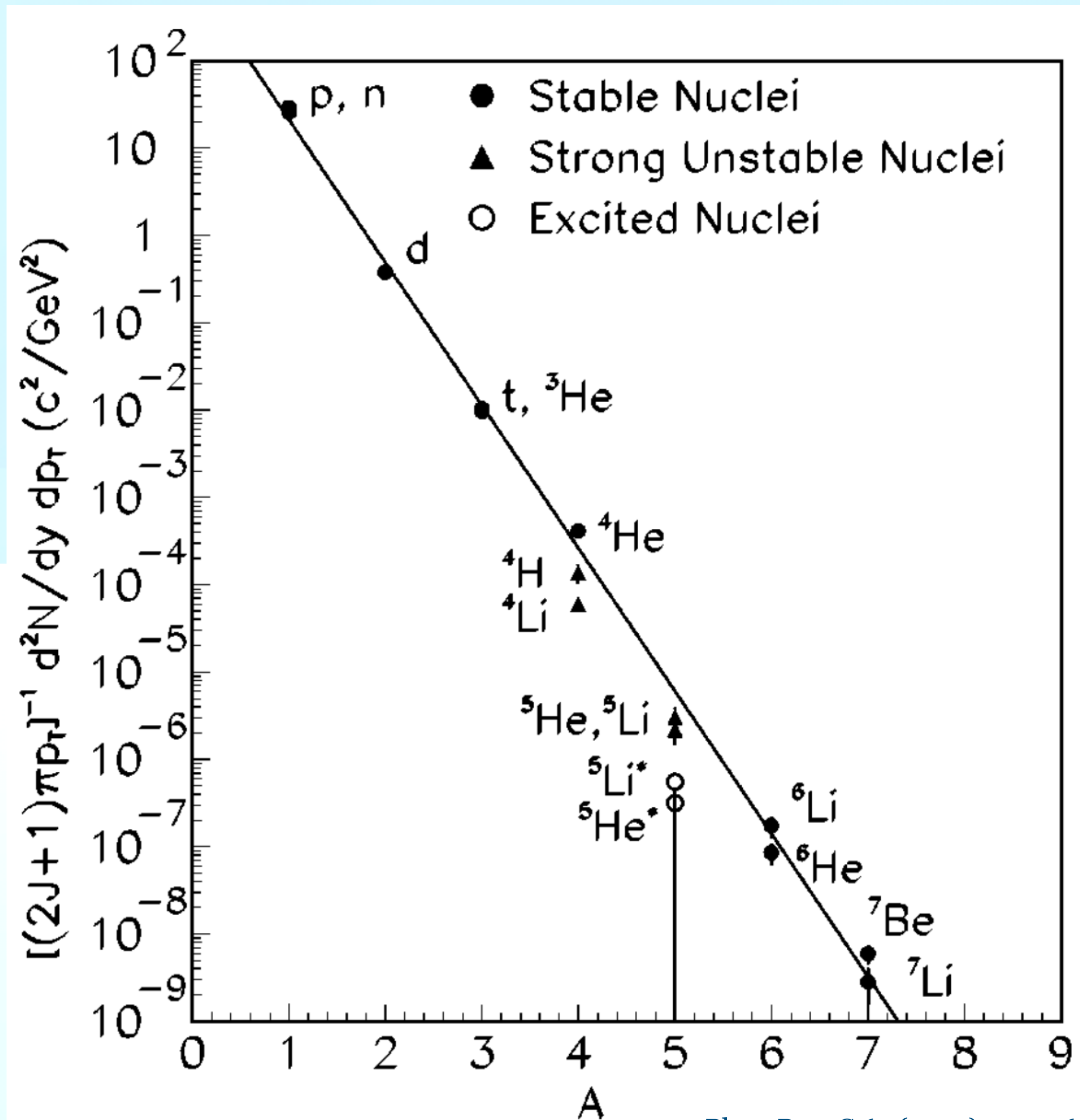
- Feed-down fractions estimated using Thermal-FIST

V. Vovchenko et al, Phys. Lett. B 809 (2020) 135746



- Feed-down correction of S_3 from unstable nuclei estimated using Thermal-FIST

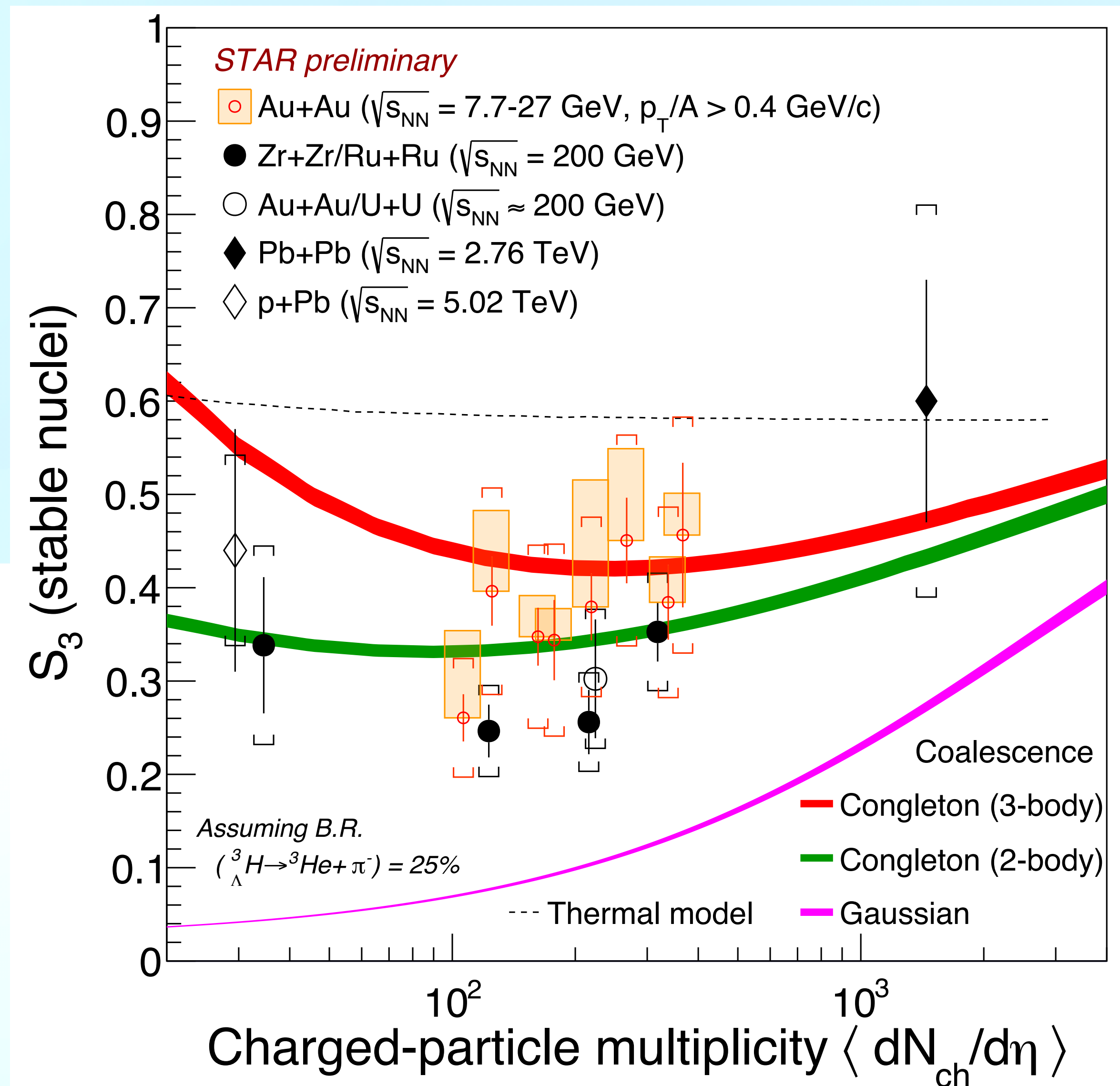
Feed-down from Unstable Nuclei



Phys. Rev. C 65 (2001) 014906

- Suppression of $A=4$ unstable states compared to ${}^4\text{He}$ ground state observed at E864

Strangeness Population Factor



Bracket:
syst. unc.

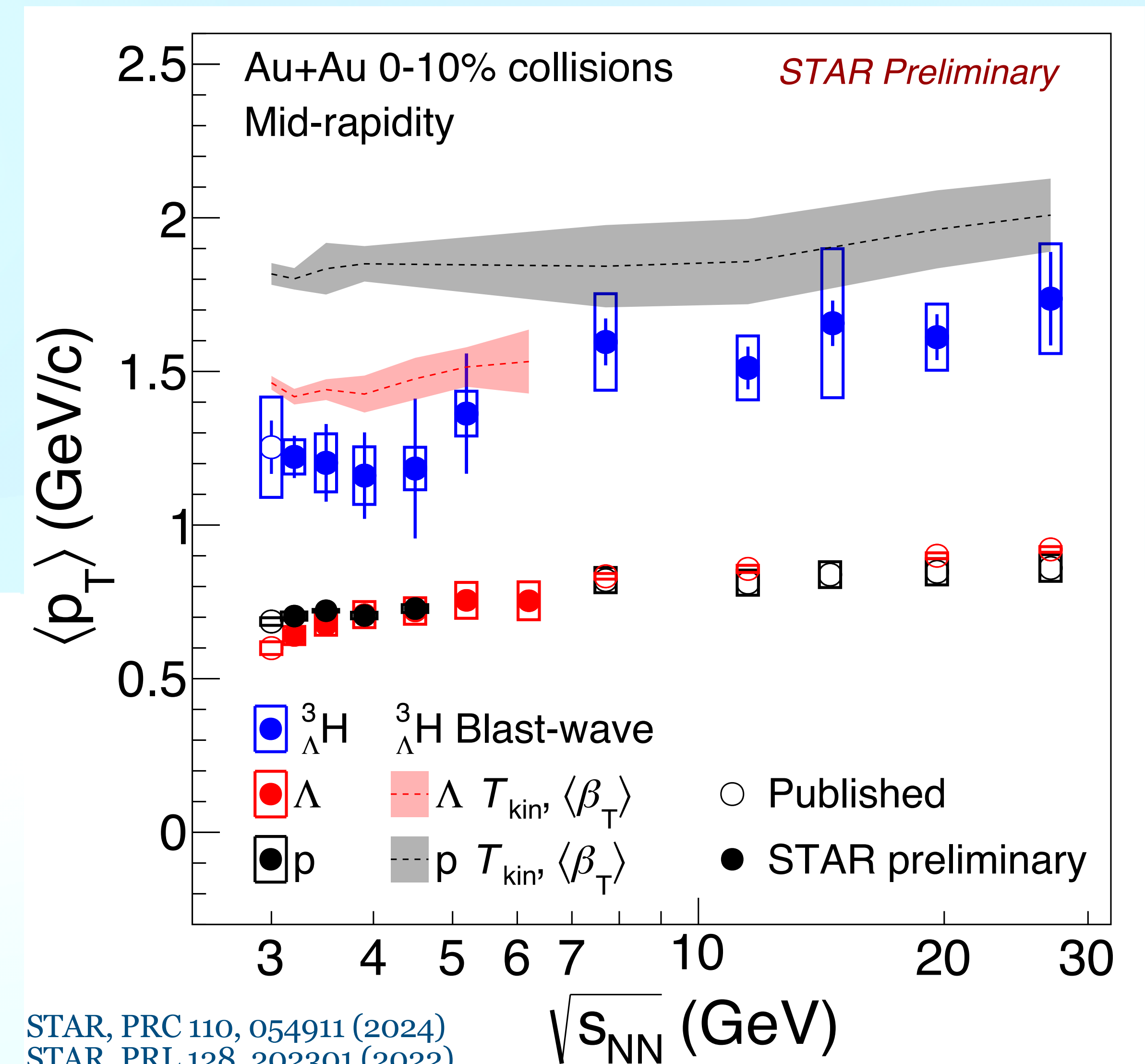
Line:
stat. unc.

Upper edge of box:
 S_3 value **corrected**
for feed-down from
thermal model

Lower edge of box:
 S_3 value **uncorrected**
for feed-down

- Unstable nuclei production are suppressed relative to stable nuclei (see backup)
- The true value of S_3 (stable nuclei) very likely lies between the upper and lower limits

Insight from Blast-Wave Comparison



- At $\sqrt{s_{NN}} \geq 7.7$ GeV, ${}^3_\Lambda\text{H}$ $\langle p_T \rangle$ tends to approach the blast-wave prediction with proton freeze-out parameters
 - May be due to increasing effective volume for coalescence with increasing energy

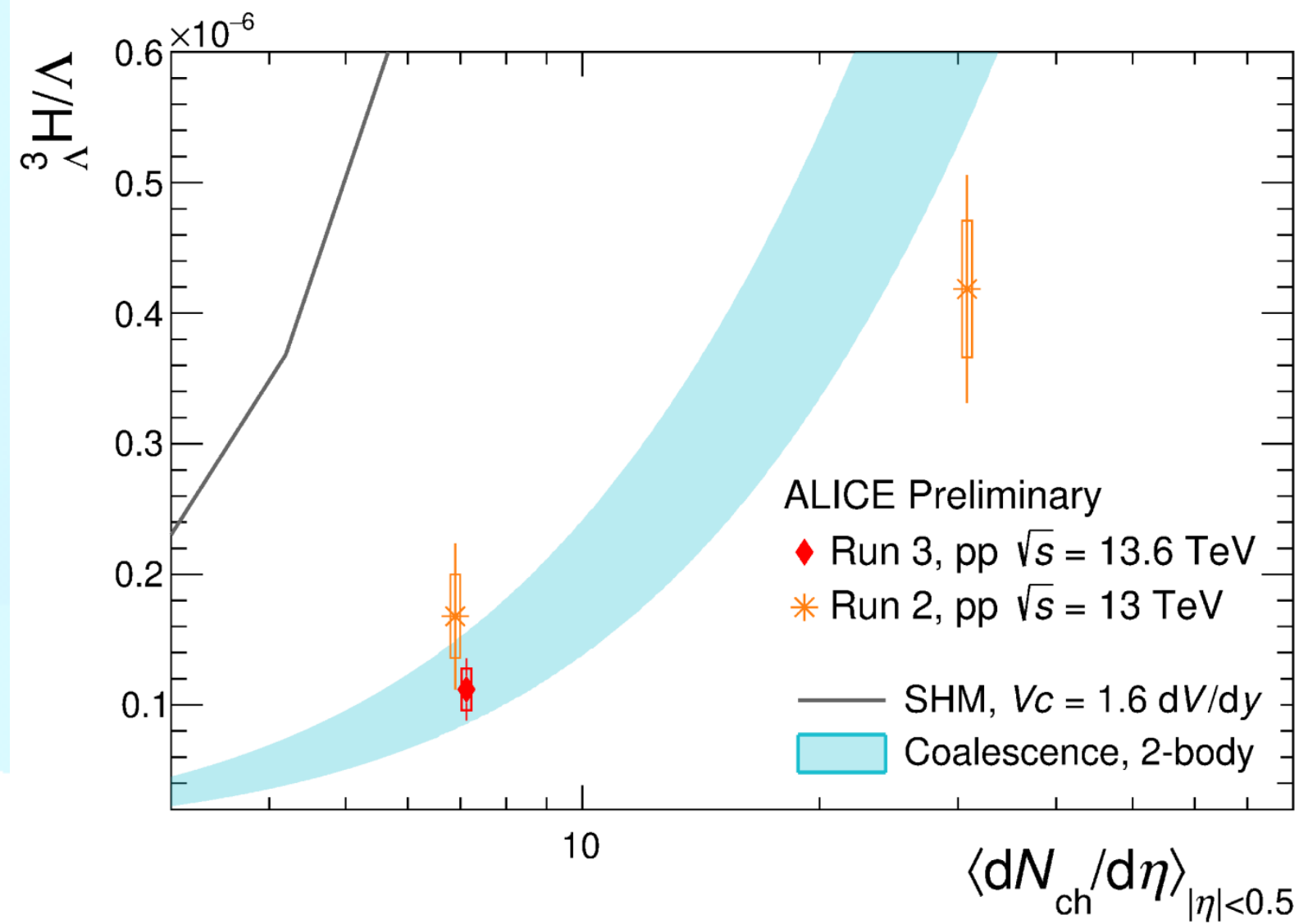
D.-N. Liu et al., PLB 855, 138855 (2024)

- Noticeable change between 4.5 and 7.7 GeV

STAR, PRC 110, 054911 (2024)
 STAR, PRL 128, 202301 (2022)
 STAR, JHEP 10 (2024) 139

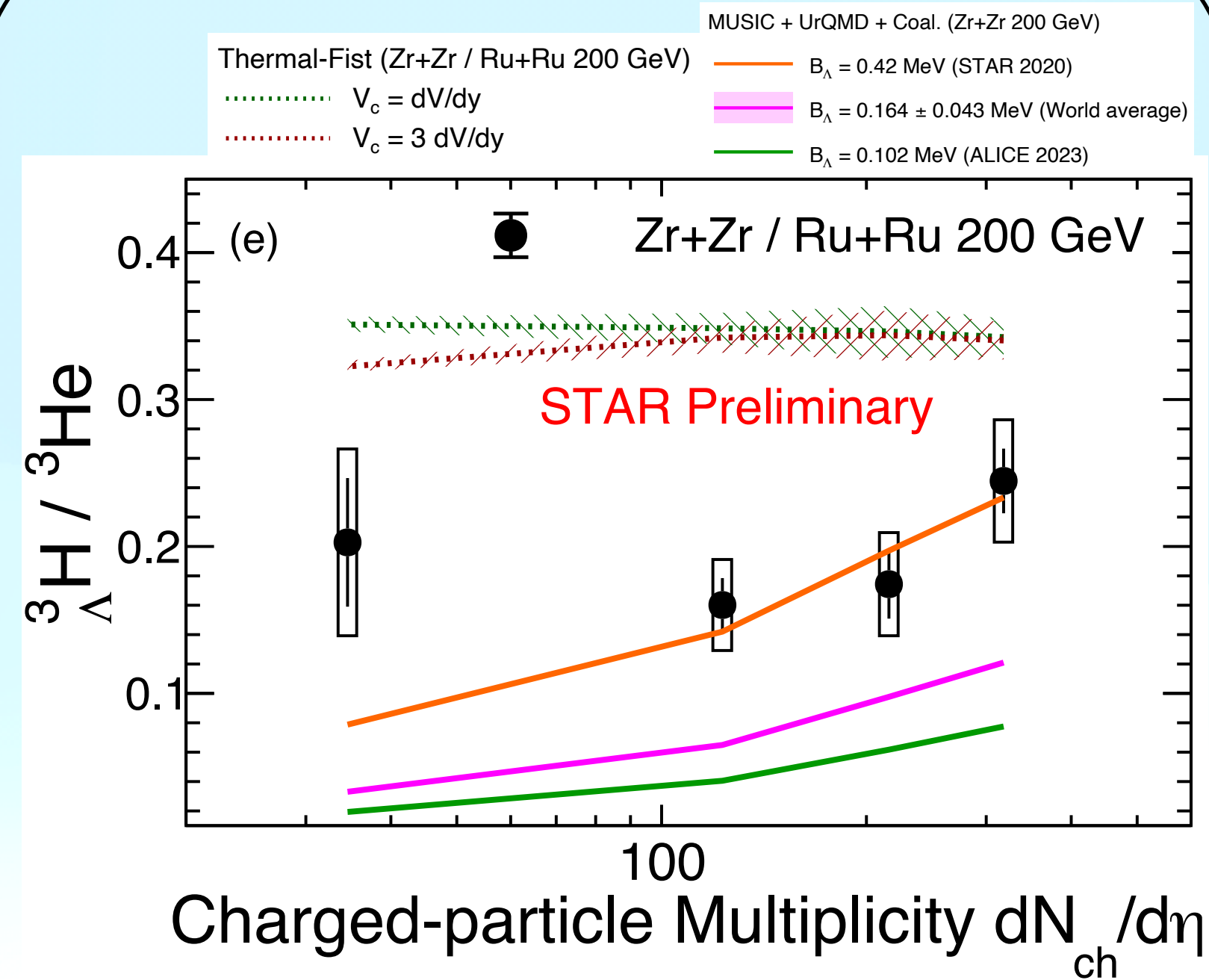
System Size Dependence at High Collision Energies

J. Ditzel (ALICE), QM2025



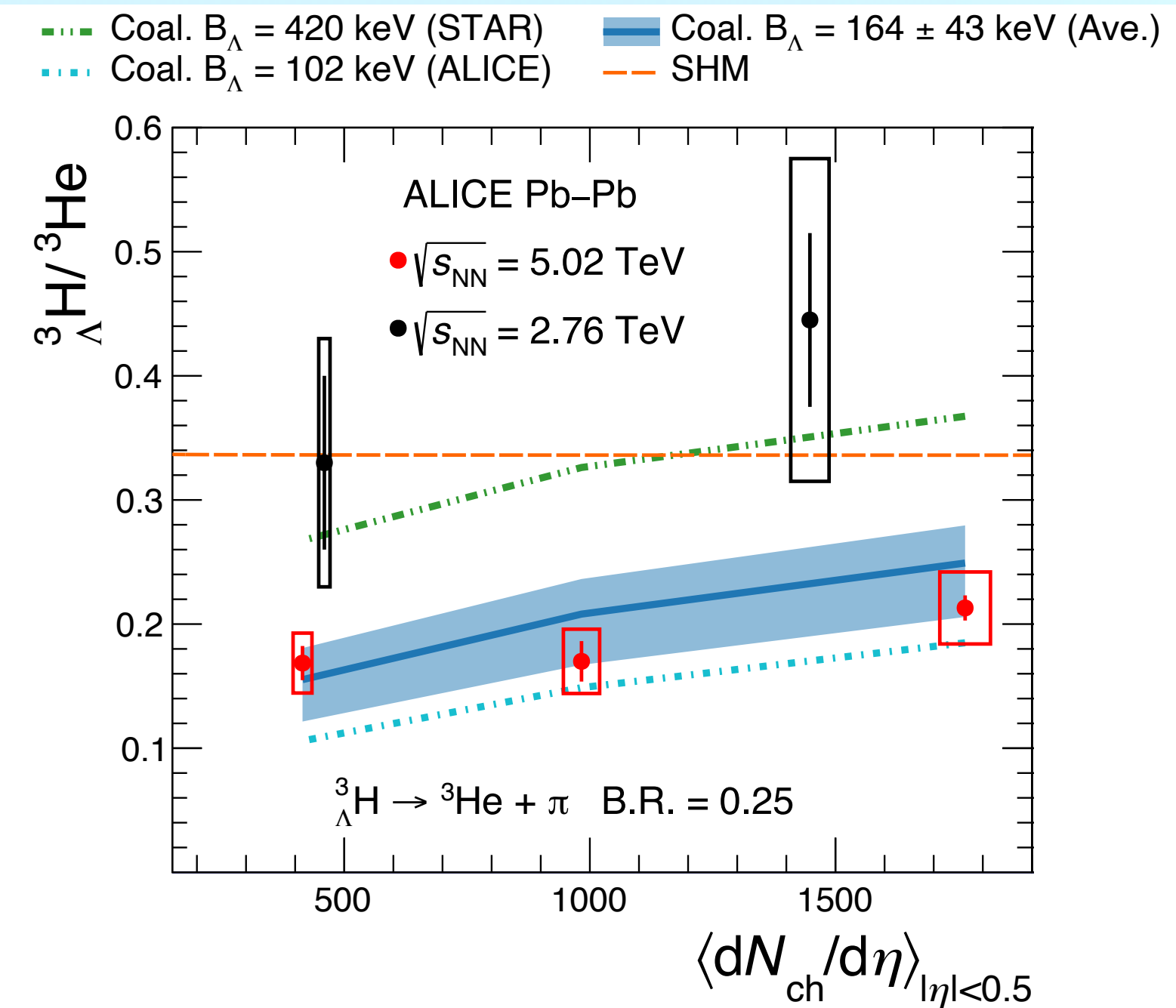
p+p @ $\sqrt{s} = 13, 13.6$ GeV
 $dN_{ch}/d\eta \sim O(10)$

D. Li (STAR), SQM2024



Ru/Zr+Ru/Zr @ $\sqrt{s_{NN}} = 200$ GeV
 $dN_{ch}/d\eta \sim O(100)$

ALICE, Phys.Lett.B 860 (2025) 139066



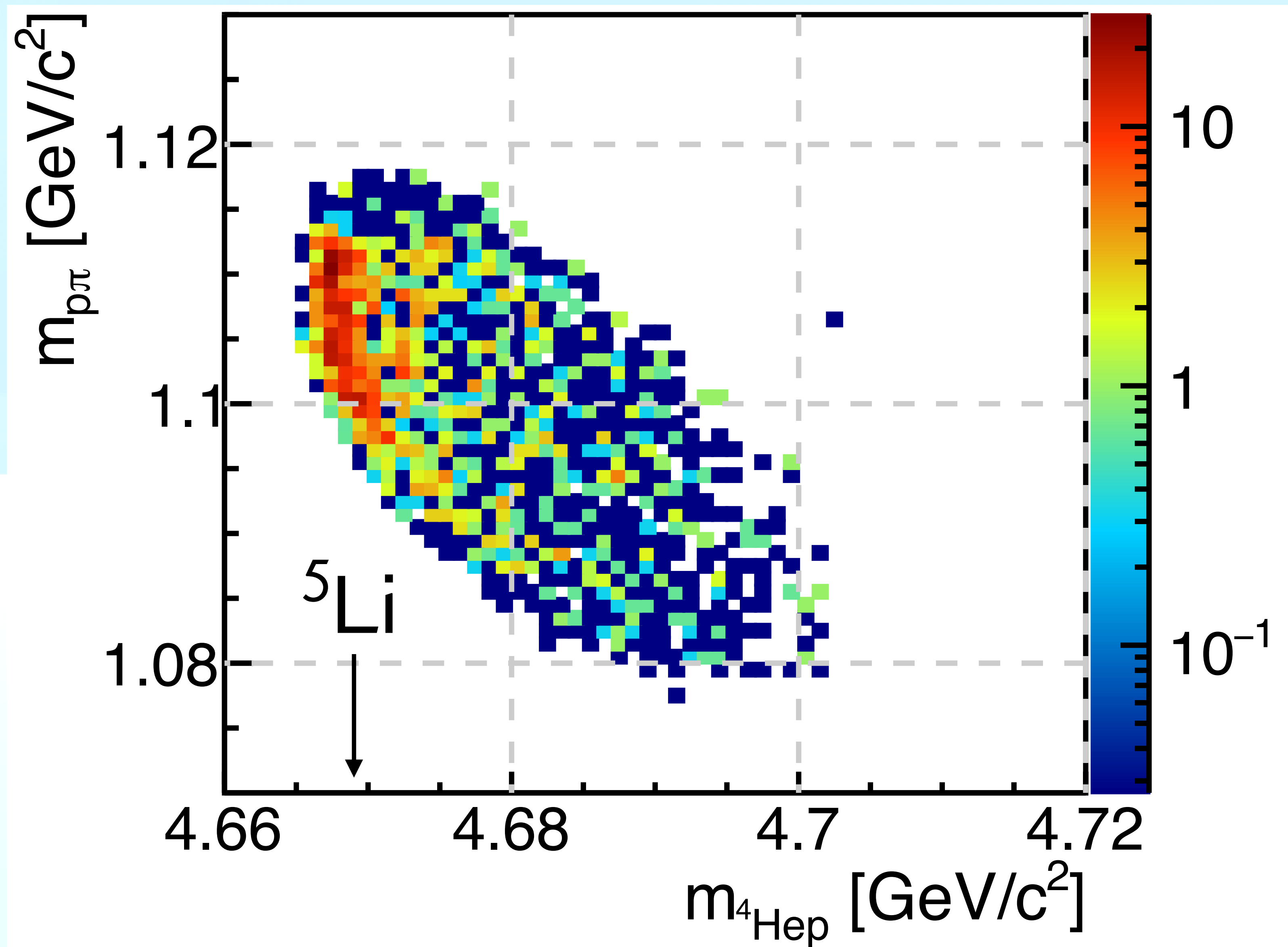
Pb+Pb @ $\sqrt{s_{NN}} = 2.76, 5.02$ TeV
 $dN_{ch}/d\eta \sim O(1000)$

Small systems

Large systems

Across all systems, thermal model overestimates $^3_\Lambda\text{H}$ production

Dalitz plot for ${}^5_{\Lambda}\text{He} \rightarrow {}^4\text{He} + \text{p} + \pi^{-}$



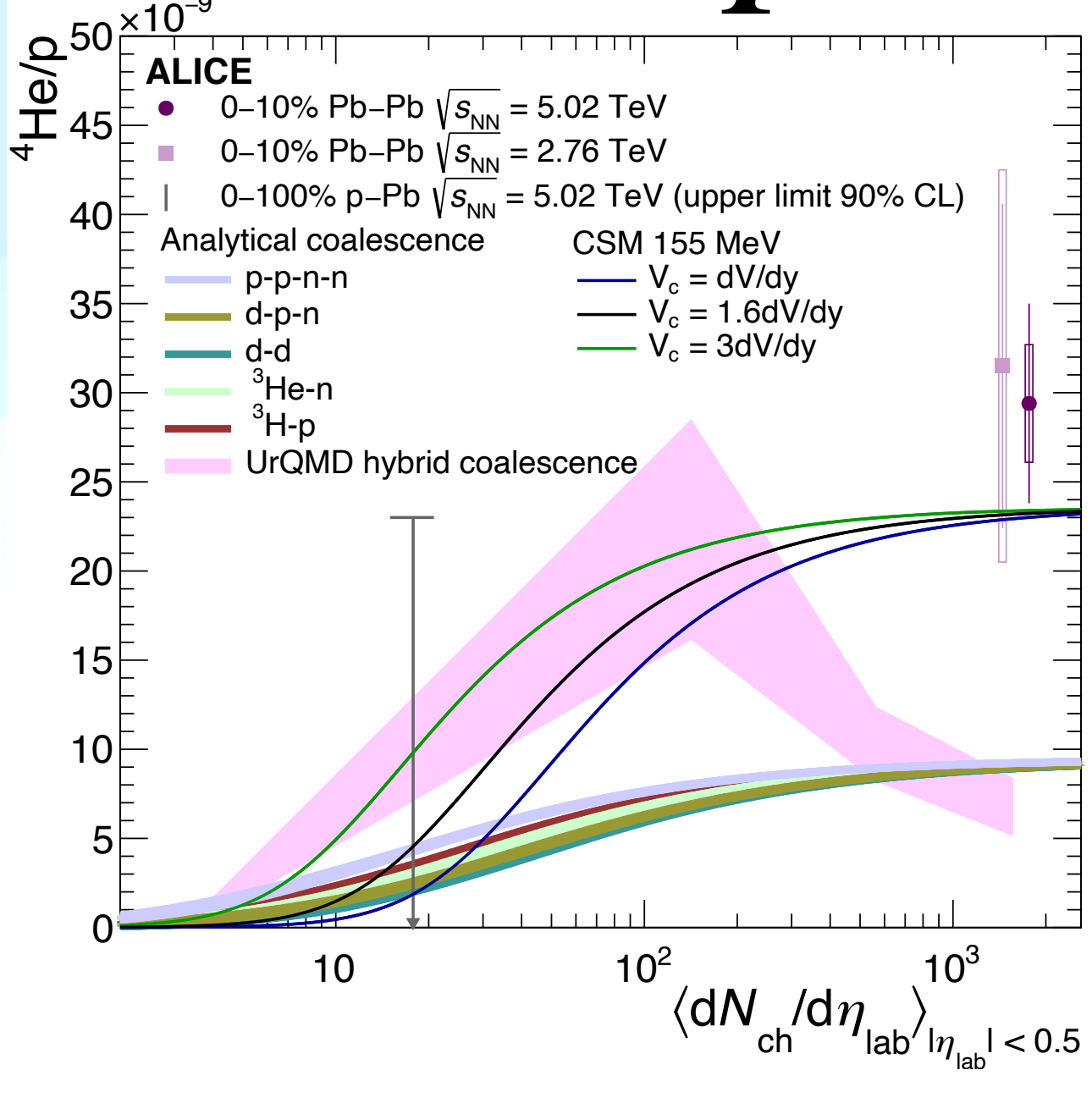
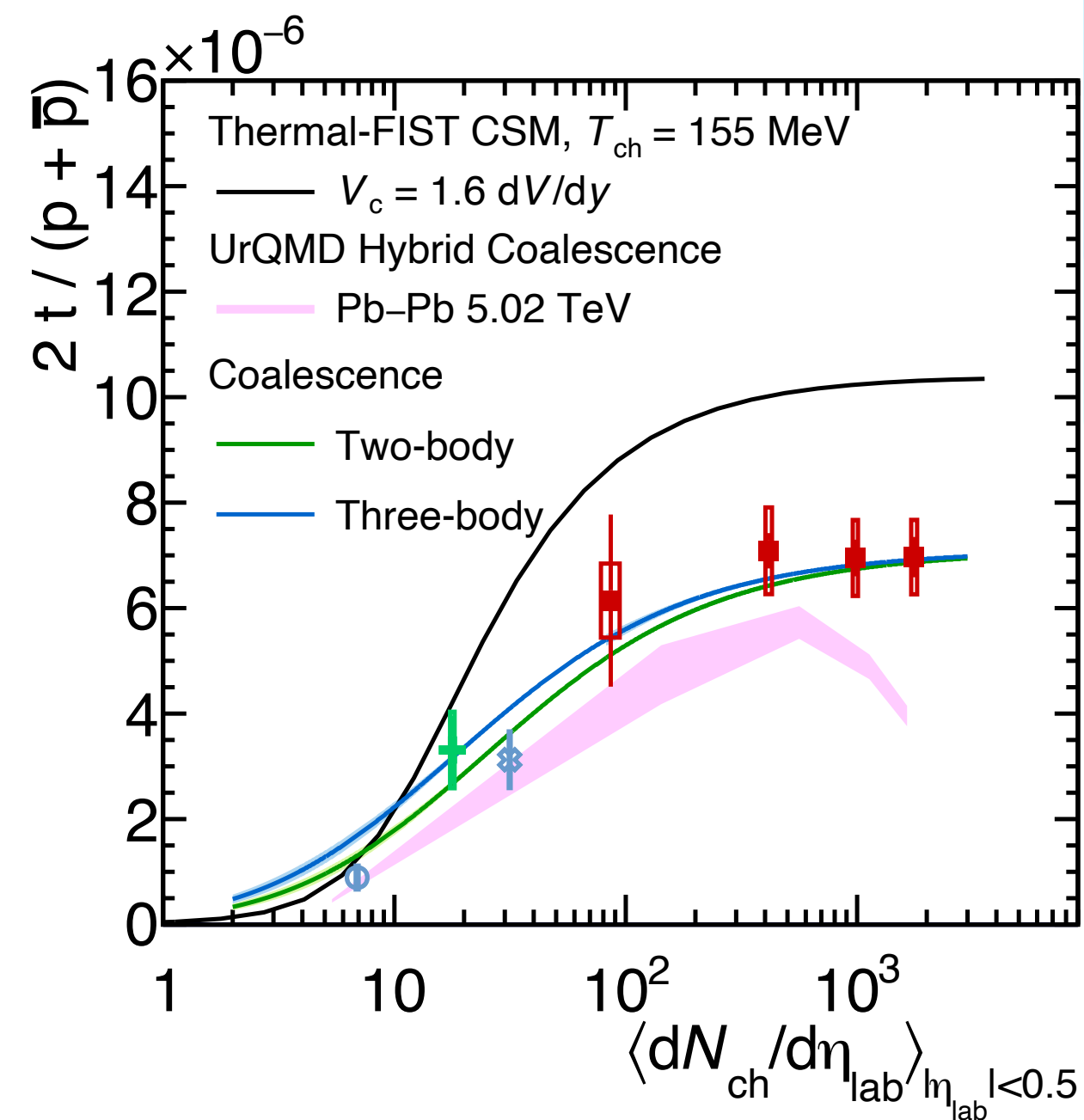
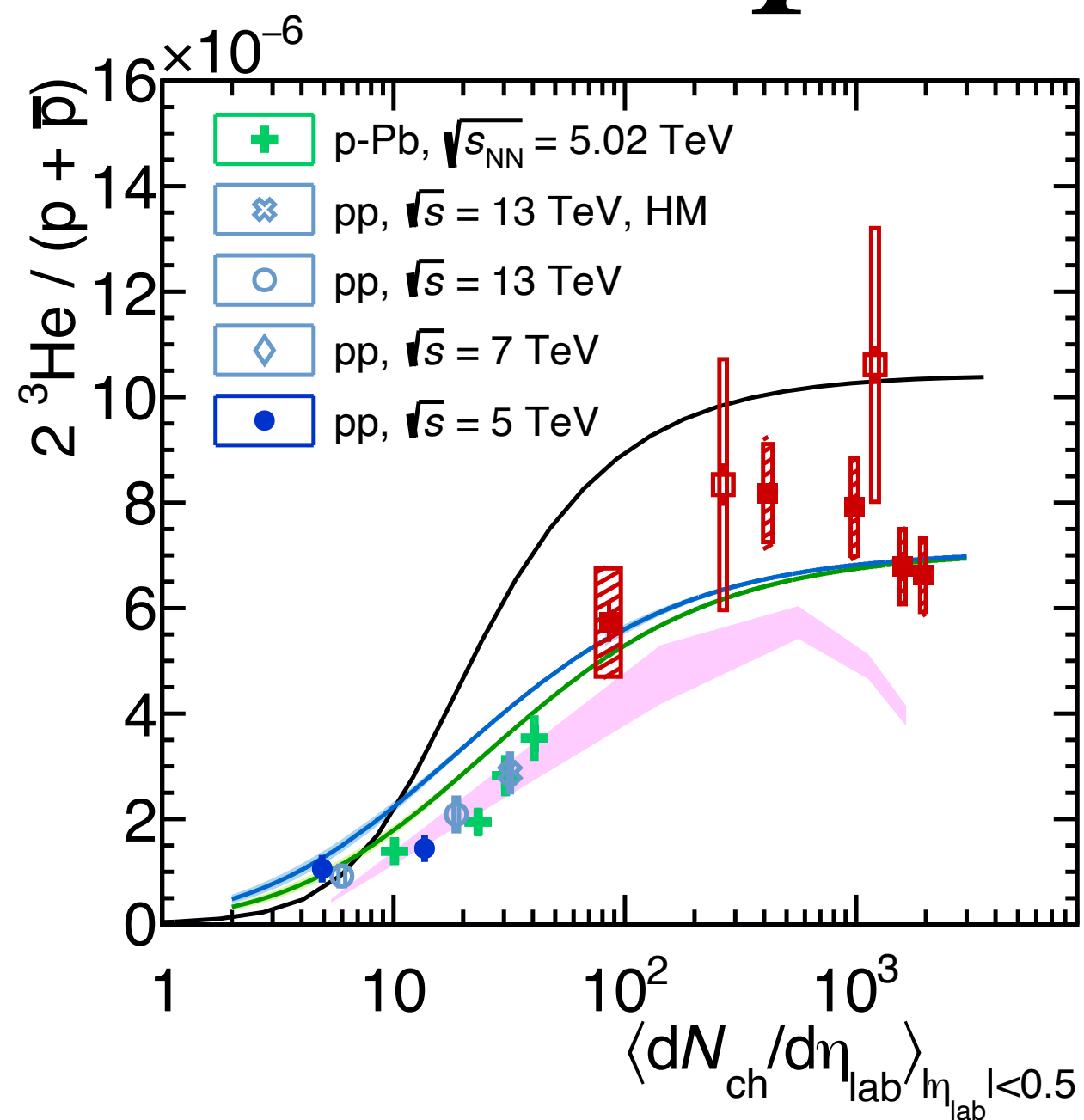
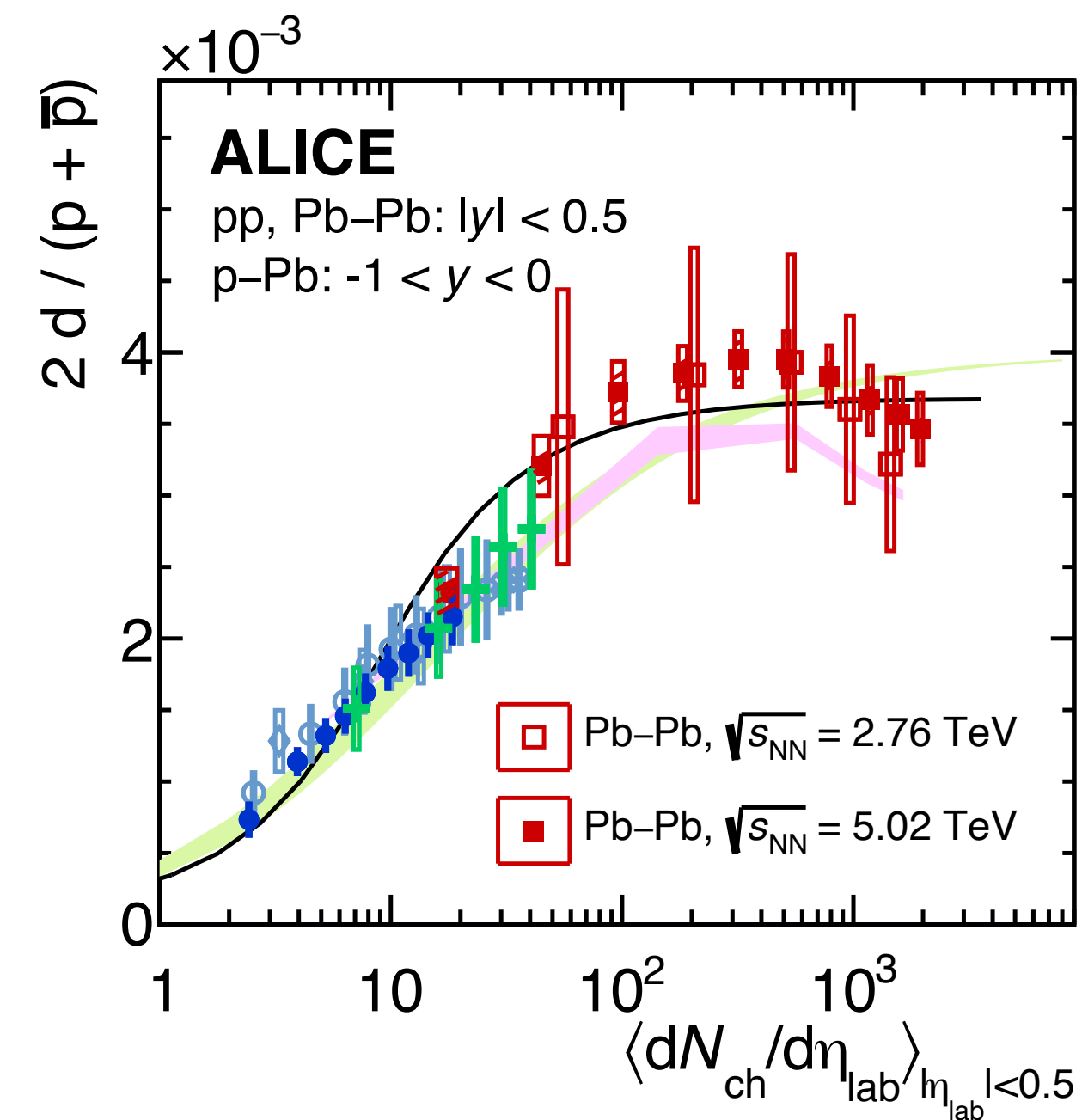
Nuclei Production at the LHC

d/p

${}^3\text{He}/p$

t/p

${}^4\text{He}/p$



A=2,4: fairly well described; A=3 nuclei: strongly overpredicted

Coalescence calculations (arXiv:2510.06758)

Phys.Rev.C 103, 014907 (2021)

$$B_3({}_\Lambda^3\text{H}) \approx \frac{3}{m^2} \frac{2s_{\Lambda^3\text{H}} + 1}{(2s_N + 1)^3} (2\pi)^6 \int d^3r_{pn} \int d^3r_\Lambda |\Phi_{\Lambda^3\text{H}}(r_{pn}, r_\Lambda)|^2 \mathcal{S}_3(r_{pn}, r_\Lambda)$$

Wave function

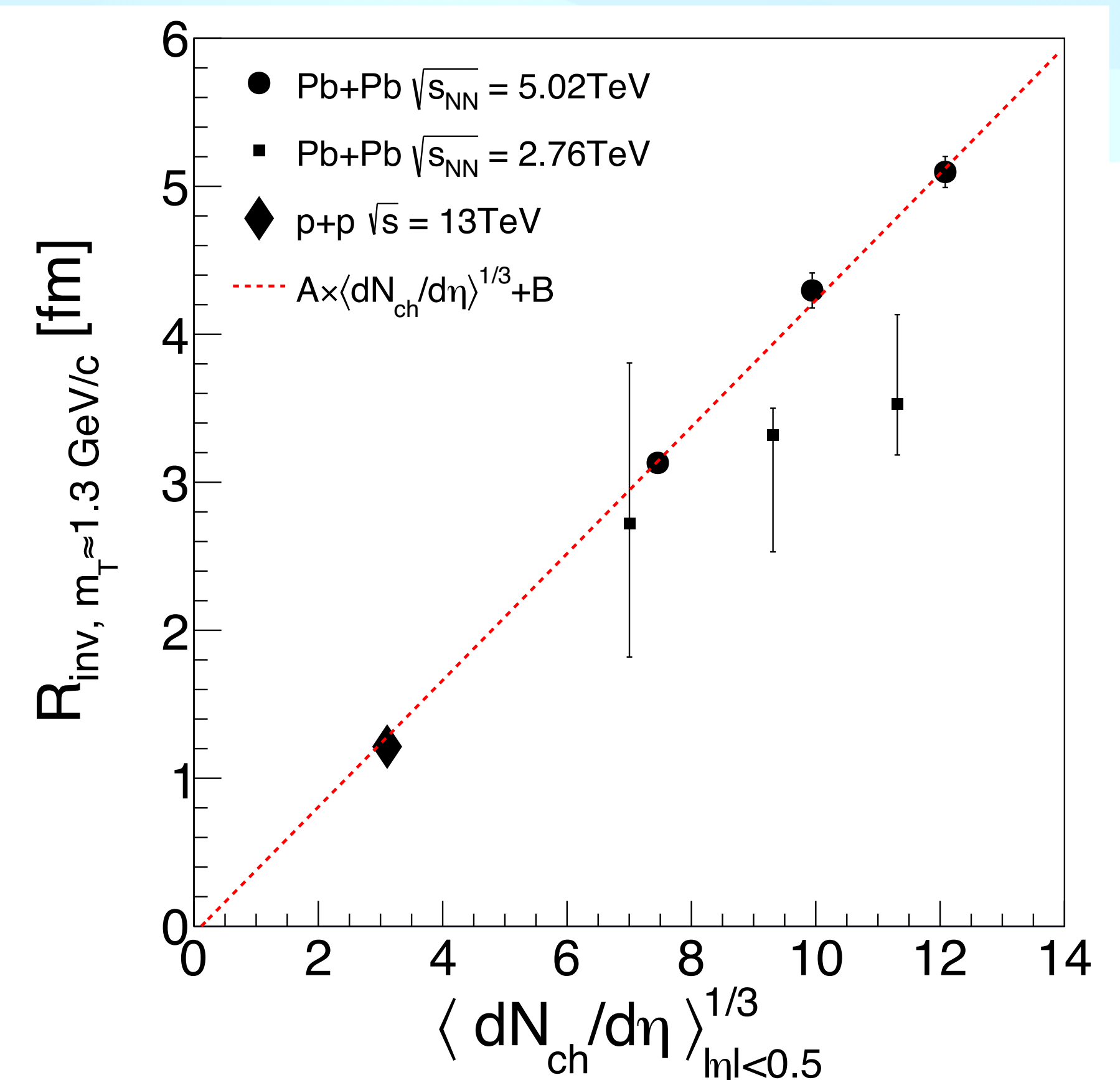
Nucleon
source

- Assuming an isotropic Gaussian nucleon source

$$\mathcal{S}_3(r_{12}, r_3) = \frac{1}{(12\pi^2 R_{\text{inv}}^4)^{3/2}} \exp\left(-\frac{r_{12}^2 + \frac{4}{3}r_3^2}{4R_{\text{inv}}^2}\right)$$

- Use ALICE baryon-baryon **correlation function data** to estimate R_{inv} for **high energy collisions**

ALICE, arXiv:2505.01061

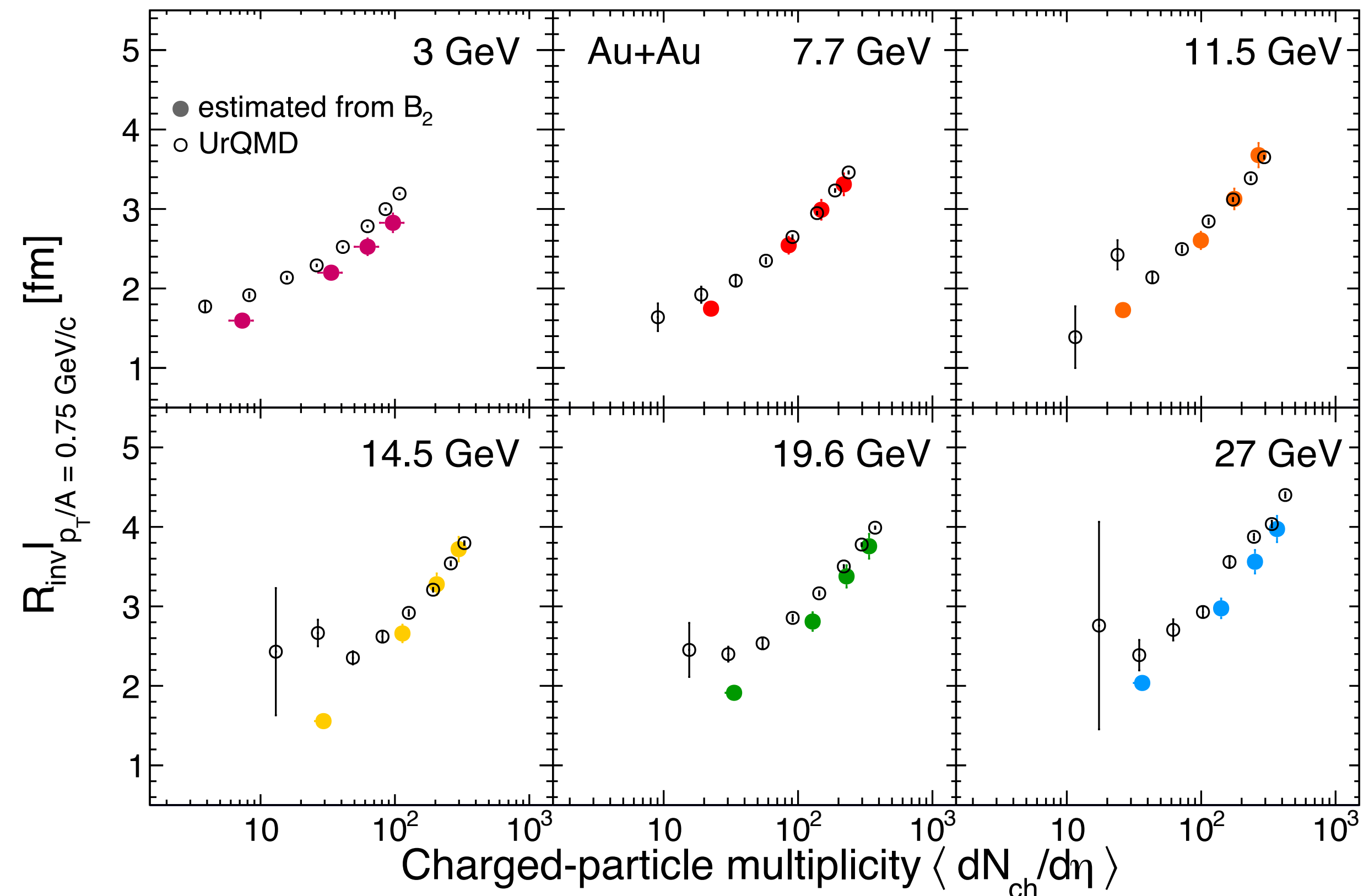
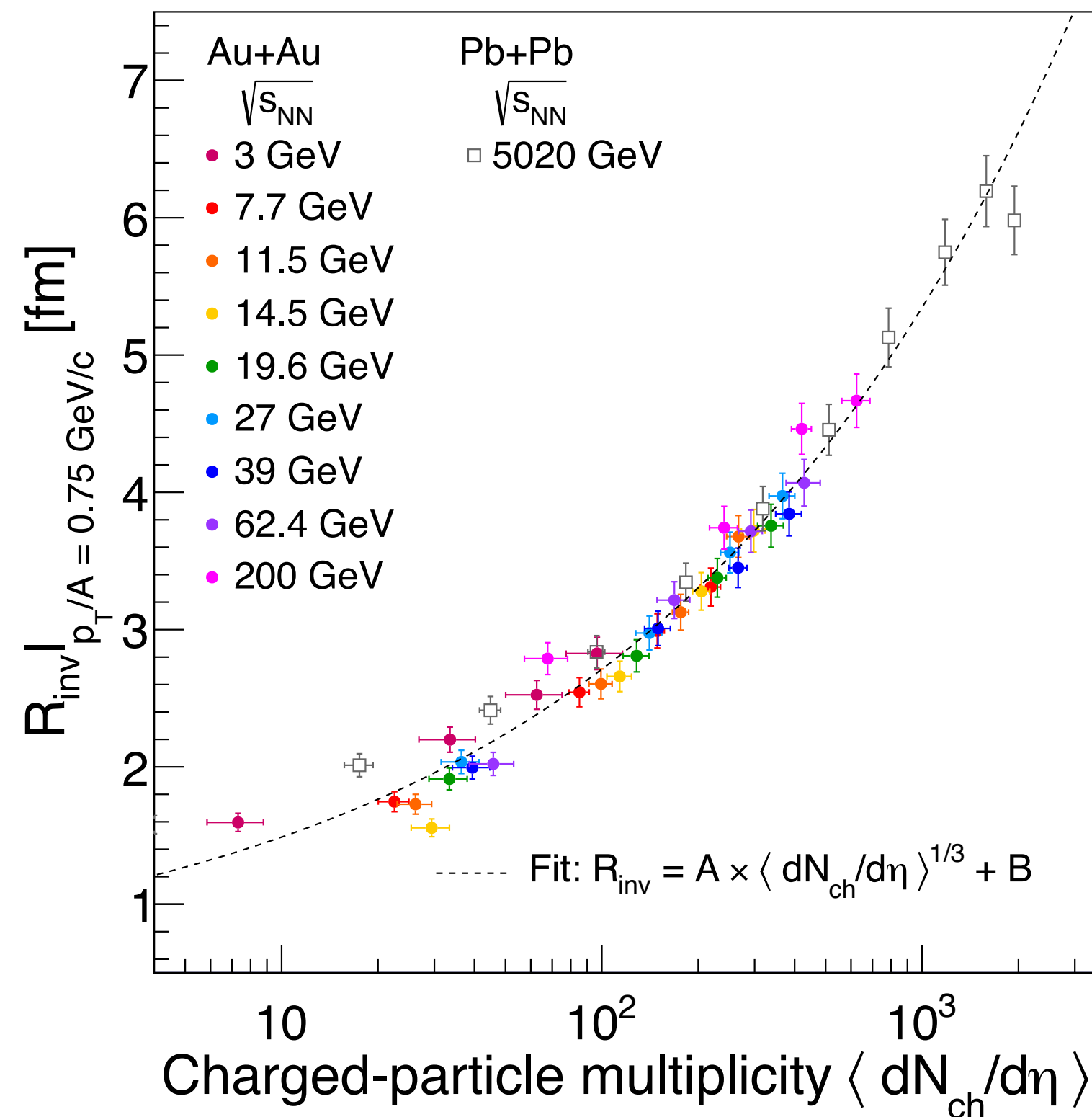


Coalescence calculations (arXiv:2510.06758)

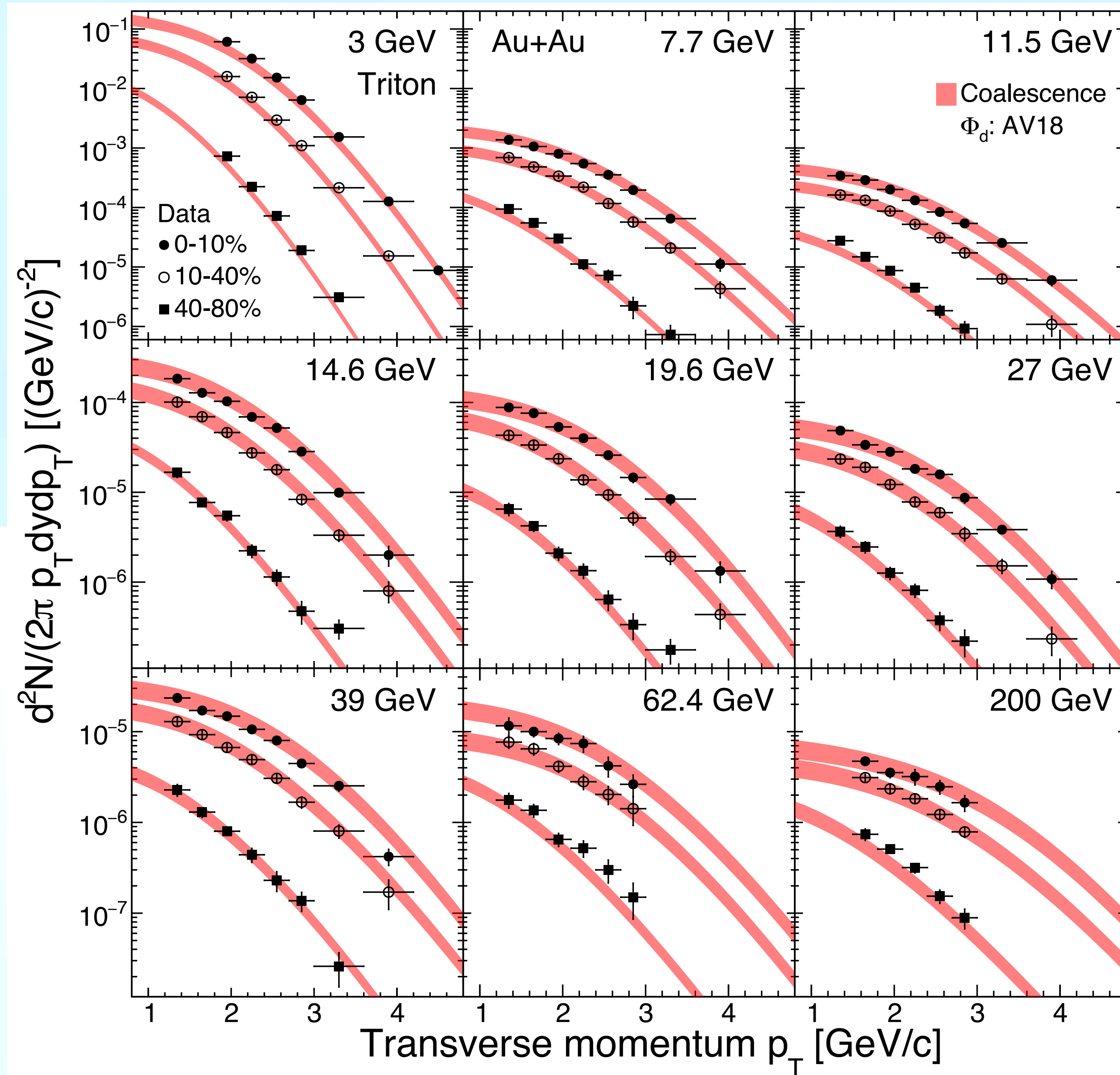
- Use deuteron data to estimate R_{inv} for **BES energies**
- Calculate B_2 as function for pT and centrality using data

- Use coalescence formula with Argonne v_{18} wave function to obtain R_{inv}

$$B_2(d) \approx \frac{2(2s_d + 1)}{m(2s_N + 1)^2(2\pi)^3} \int d^3r |\Phi_d(r)|^2 \mathcal{S}_2(r)$$



Coalescence calculations (arXiv:2510.06758)



- Calculations describe triton data well

Overview of (Hyper)Nuclei Results from BES-II

$\sqrt{s_{NN}}$ [GeV]	d	t	${}^3\text{He}$	${}^4\text{He}$	$\bar{d}$	${}^3_{\Lambda}\text{H}$	${}^4_{\Lambda}\text{H}$	${}^4_{\Lambda}\text{He}$	${}^5_{\Lambda}\text{He}$
3	✓	✓	✓	✓		✓	✓	✓	✓
3.2	✓	✓	✓	✓		✓	✓	✓	
3.5	✓	✓	✓	✓		✓	✓	✓	
3.9	✓	✓	✓	✓		✓	✓		
4.5	✓	✓	✓	✓		✓			
5.2						✓			
7.7	✓		✓		✓	✓			
9.2	✓		✓		✓				
11.5	✓		✓		✓	✓			
14.6	✓		✓		✓	✓			
17.3	✓		✓		✓				
19.6	✓		✓		✓	✓			
27	✓		✓		✓	✓			

FXT: mid-rapidity to target
rapidity from 3-4.5 GeV

COL: mid-rapidity

✓ : Published
✓ : Preliminary