### 第十届手征有效场论研讨会

# Discontinuity calculus

analytic continuation and topological structure of two-body scattering amplitudes

Speaker: Hao-Jie Jing (SXU)

NJNU · Nanjing · 2025-10-20

Based on: HJJ, Xiong-Hui Cao and Feng-Kun Guo [arXiv:2507.06175]

[ To be published in Frontiers of Physics ]

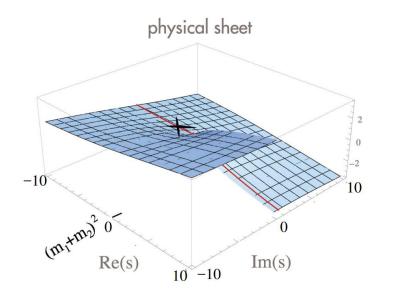
#### Contents

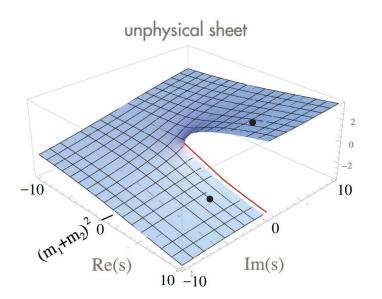
- Motivation
- Discontinuity Calculus
- Discontinuity Analysis of the Two-Point Green's Function
- Discontinuity Analysis of the Partial-Wave Scattering Matrix
- Uniformization of the Partial-Wave Scattering Matrix
- Summary and Outlook



#### Motivation

- S-matrix theory is a cornerstone tool for studying scattering processes
  - Bound states/Resonances correspond to poles in scattering amplitudes. The pole positions and residues encode crucial particle properties.
  - Causality constraint: Unstable resonance poles reside on unphysical Riemann sheets of the complex energy plane.
  - Discontinuity analysis of the S-matrix is indispensable: it extracts information about poles/zeros (discreteness) and branch cuts (discontinuity).



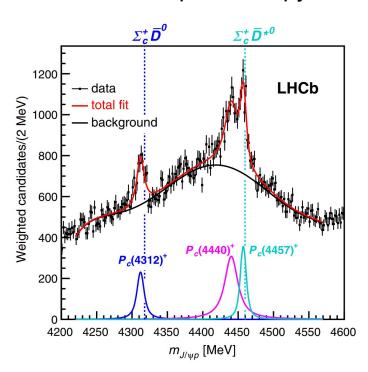


Feng-Kun Guo et al, Rev.Mod.Phys. 90 (2018).

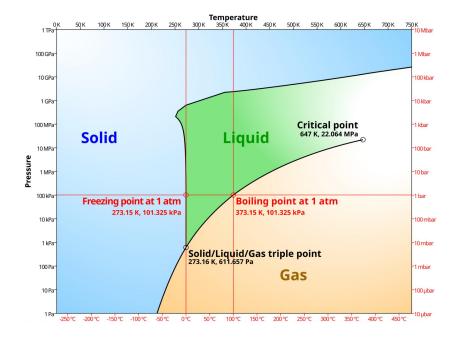
#### Motivation

• Discreteness and discontinuity are ubiquitous in phenomena:

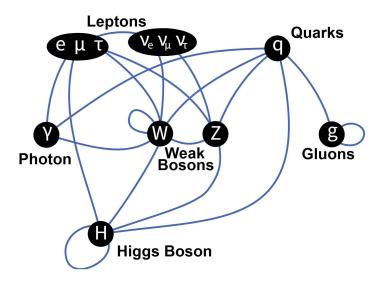
#### Hadron spectroscopy



#### Statistical physics



Standard model



LHCb, Phys. Rev. Lett. 122(2019) 22, 222001

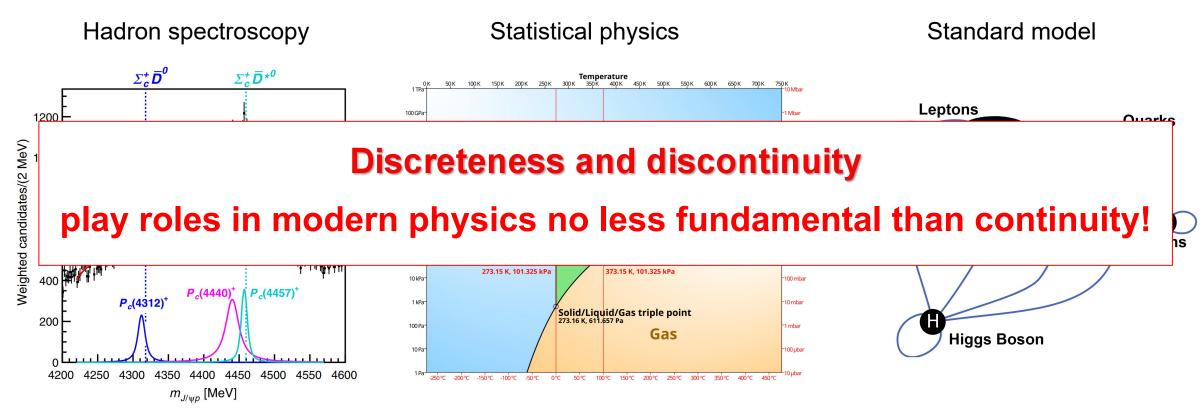
Phase diagram of water

**Elementary particles mass gaps** 



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## Calculus vs Discontinuity Calculus

- Differential operator
  - 1. Constant kernel: dC = 0
  - 2. Linearity:

$$d [\alpha_1 f_1(z) + \alpha_2 f_2(z)] = \alpha_1 df_1(z) + \alpha_2 df_2(z)$$

3. Leibniz law:

$$d[f_1(z)f_2(z)] = df_1(z)f_2(z) + f_1(z)df_2(z)$$

4. Chain rule:

$$dF[f(z)] = \frac{dF(w)}{dw}\Big|_{w=f(z)} df(z)$$

5. Newton-Leibniz formula:

$$f(z) = C(z_0) + \int_{z_0}^z \mathrm{d}f(z')$$

## Calculus vs Discontinuity Calculus

- Differential operator
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- Discontinuity operator
  - 1. Holomorphic kernel:  $\mathbb{D}h(z) = 0$
  - 2. Linearity:

$$\mathbb{D}\left[\alpha_1 f_1(z) + \alpha_2 f_2(z)\right] = \alpha_1 \mathbb{D}f_1(z) + \alpha_2 \mathbb{D}f_2(z)$$

3. Leibniz law:

$$\mathbb{D}[f_1(z)f_2(z)] = \mathbb{D}f_1(z) \ f_2(z) + f_1(z) \ \mathbb{D}f_2(z) - \mathbb{D}f_1(z) \ \mathbb{D}f_2(z)$$

4. Chain rule:

$$\mathbb{D}F\left[f(z)\right] = F\left[f(z)\right] - \left[F\left(\omega\right) - \mathbb{D}F(\omega)\right]_{\omega = f(z) - \mathbb{D}f(z)}$$

5. Dispersion relation:

$$f(z) = h(z) + \frac{1}{2\pi i} \int \frac{\mathbb{D}f(z')}{z' - z} dz'$$

ullet Discontinuity of the square root function  $f(z)=\sqrt{z} \ (z\in\mathbb{C})$ 

$$\begin{cases} \mathbb{D}(\sqrt{z})^2 = \mathbb{D}z = 0\\ \mathbb{D}(\sqrt{z})^2 = 2\sqrt{z} \ \mathbb{D}\sqrt{z} - (\mathbb{D}\sqrt{z})^2 \end{cases}$$

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$$\mathbb{D}\sqrt{z} = 2\sqrt{z} \times \theta(z) + 0 \times \theta(-z) = 2\sqrt{z} \ \theta(z)$$

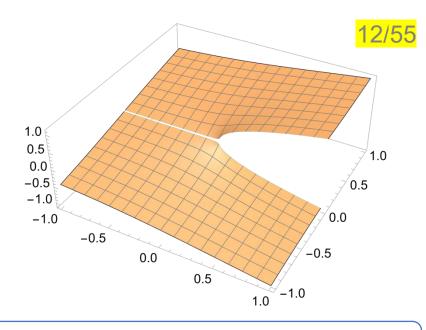
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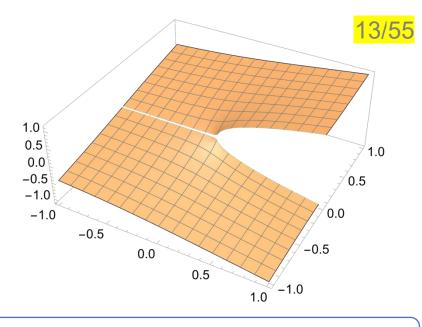
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$$\mathbb{K}_1\sqrt{z} = 2\sqrt{z} \qquad \mathbb{K}_2\sqrt{z} = 0$$

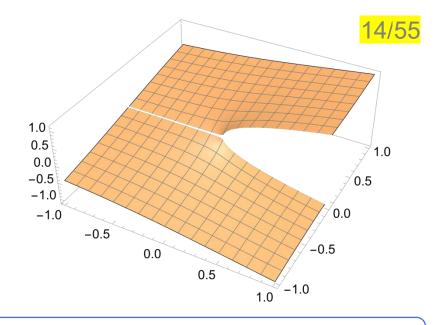
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$$\begin{bmatrix} \mathbb{T}_1 \sqrt{z} = (1 - \mathbb{K}_1) \sqrt{z} = -\sqrt{z} \\ \mathbb{T}_1^2 \sqrt{z} = -\mathbb{T}_1 \sqrt{z} = \sqrt{z} \end{bmatrix}$$

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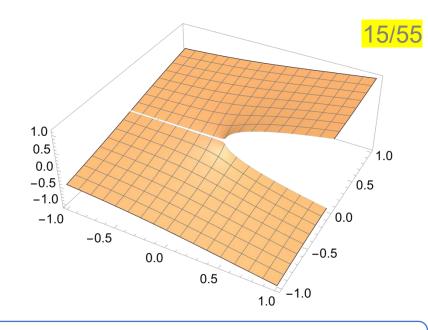
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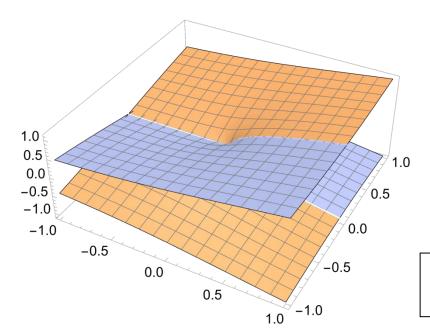


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$$\mathbb{K}_1 \sqrt{z} = 2\sqrt{z} \qquad \mathbb{K}_2 \sqrt{z} = 0$$



$$RS{\sqrt{z}} = \mathbb{C} \times {\mathbb{T}_1 \mid \mathbb{T}_1^2 \sim 1} \cong \mathbb{C} \times \mathbb{Z}_2$$

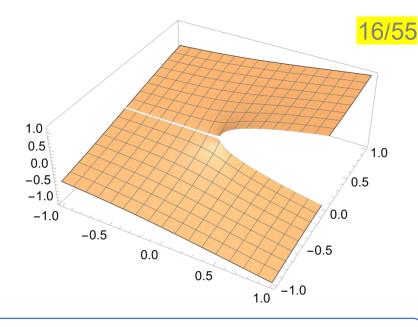
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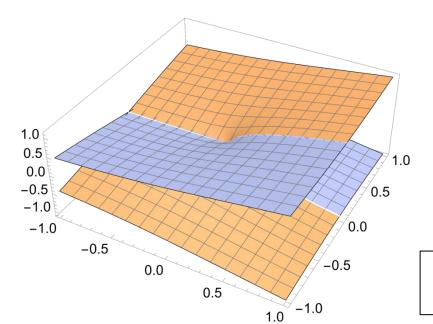


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$$T_1\sqrt{z} = (1 - \mathbb{K}_1)\sqrt{z} = -\sqrt{z}$$
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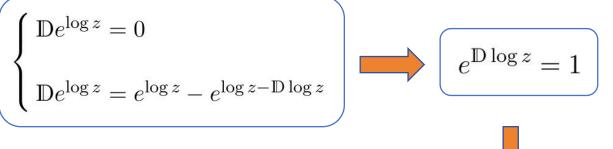


$$\mathbb{K}_1 \sqrt{z} = 2\sqrt{z} \qquad \mathbb{K}_2 \sqrt{z} = 0$$

The Riemann surface for the n-th root function:

$$RS\{\sqrt[n]{z}\} \cong \mathbb{C} \times \mathbb{Z}_n \text{ for } n \in \mathbb{N}^*$$

• Discontinuity of the logarithmic function  $f(z) = \log z \ (z \in \mathbb{C}^* \equiv \mathbb{C} \setminus \{0\})$ 

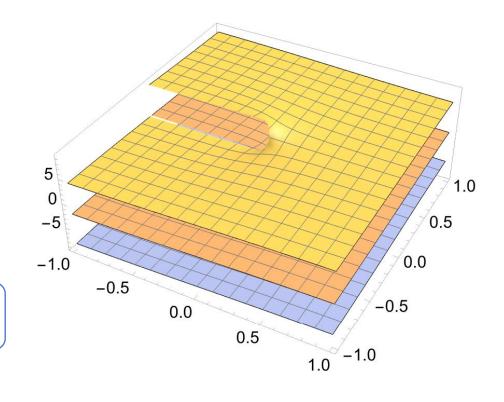




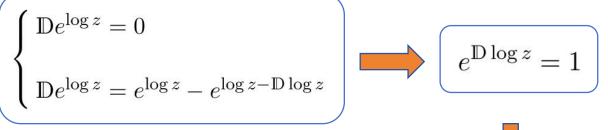
$$\mathbb{D}\log z = \pm 2\pi i\,\theta(-z)$$



$$RS\{\log z\} = \mathbb{C}^* \times \{\mathbb{T}_1, \mathbb{T}_{-1} \mid \mathbb{T}_1 \mathbb{T}_{-1} \sim 1\} \cong \mathbb{C}^* \times \mathbb{Z}$$



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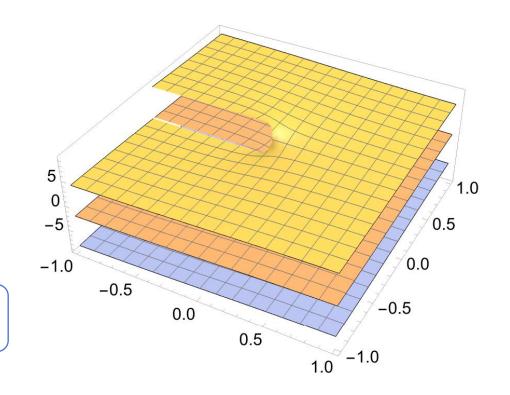




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Discontinuity of meromorphic functions

$$\mathbb{D}\left(\frac{1}{z}\right) = \mp \, 2\pi i \, \delta(z)$$

$$\mathbb{D}\left(\frac{1}{z^{n+1}}\right) = \mp 2\pi i \, \frac{(-1)^n}{n!} \frac{\mathrm{d}^n}{\mathrm{d}z^n} \delta(z)$$

• Discontinuity of complex-valued functions

$$\mathbb{D}f(z) = \sum_{i=1}^{n} \mathbb{K}_{i}f(z) \ \theta_{i}(z) + \sum_{j=1}^{n'} \alpha_{j}\delta_{j}(z)$$

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$$\mathbb{D}f(z) = \sum_{i=1}^{n} \mathbb{K}_{i}f(z) \ \theta_{i}(z) + \sum_{j=1}^{n'} \alpha_{j}\delta_{j}(z)$$

Continuation of complex-valued functions

 $\mathbb{K}_i f(z)$ : the **continuation kernel** of f(z)

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$$\mathbb{D}\left[\mathbb{T}_i f(z)\right] = \sum_{j=1}^n \mathbb{K}_j \left[\mathbb{T}_i f(z)\right] \theta_j(z) + \sum_{j=1}^{n'} \alpha_j \delta_j(z)$$

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 $\mathbb{T}_i \equiv 1 - \mathbb{K}_i$ : the continuation generator

$$f(z) \rightarrow \mathbb{T}_i f(z) \rightarrow \mathbb{T}_j \mathbb{T}_i f(z) \rightarrow \mathbb{T}_k \mathbb{T}_j \mathbb{T}_i f(z) \rightarrow \cdots$$

• Discontinuity of complex-valued functions

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Continuation of complex-valued functions

**Monodromy group** 

 $\mathbb{K}_i f(z)$ : the **continuation kernel** of f(z)

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: the **continuation generator**

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$$f(z) \rightarrow \mathbb{T}_i f(z) \rightarrow \mathbb{T}_j \mathbb{T}_i f(z) \rightarrow \mathbb{T}_k \mathbb{T}_j \mathbb{T}_i f(z) \rightarrow \cdots$$

A discontinuity relation

$$\begin{cases} \mathbb{D}\left[\frac{1}{f(z)} \cdot f(z)\right] = 0 \\ \mathbb{D}\left[\frac{1}{f(z)} \cdot f(z)\right] = \mathbb{D}\left[\frac{1}{f(z)}\right] \cdot f(z) + \frac{1}{f(z)} \cdot \mathbb{D}f(z) - \mathbb{D}\left[\frac{1}{f(z)}\right] \cdot \mathbb{D}f(z) \end{cases}$$



$$\mathbb{D}f(z) = -f(z) \cdot \mathbb{D}[1/f(z)] \cdot [f(z) - \mathbb{D}f(z)]$$

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The scalar two-point one-loop Green's function

$$G(p^{2}) \equiv \int_{\mathbb{R}^{4}} \frac{i \, d^{4}q}{(2\pi)^{4}} \frac{1}{\left[ (p/2 + q)^{2} - m_{1}^{2} + i\epsilon \right] \left[ (p/2 - q)^{2} - m_{2}^{2} + i\epsilon \right]}$$

• Discontinuity analysis with explicit expression  $(s=p^2)$ 

$$G(s) = \frac{1}{(4\pi)^2} \left[ a(\mu) + \log\left(\frac{m_1 m_2}{\mu^2}\right) + \frac{(m_1^2 - m_2^2)}{s} \log\frac{m_1}{m_2} + 8\pi\rho(s)\varphi(s) \right]$$



direct calculation 
$$\rho(s) = \frac{1}{16\pi s} \sqrt{(s-s_+)(s-s_-)} \qquad s_{\pm} = (m_1 \pm m_2)^2$$

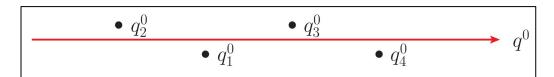
$$\varphi(s) = \log [c - s - 16\pi s \rho(s)] - \log [c - s + 16\pi s \rho(s)]$$

$$\mathbb{D}G(s) = -2i\rho(s)\theta(s - s_{+})$$

• Discontinuity analysis from the definition

$$D_1 \equiv q^0 + E/2 - \omega_1$$
  $D_3 \equiv q^0 - E/2 + \omega_2$   $D_2 \equiv q^0 + E/2 + \omega_1$   $D_4 \equiv q^0 - E/2 - \omega_2$   $(E = \sqrt{s})$ 

$$G(s) = \int_{\mathbb{R}^4} \frac{id^4q}{(2\pi)^4} \frac{1}{(D_1 + i\epsilon)(D_2 - i\epsilon)(D_3 - i\epsilon)(D_4 + i\epsilon)}$$



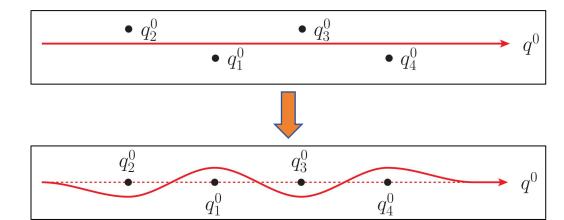


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$$G(s) = \int_{\mathcal{C} \times \mathbb{R}^3} \frac{id^4q}{(2\pi)^4} \frac{1}{D_1 D_2 D_3 D_4}$$





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$$G(s) = \int_{\mathcal{C}(n_G) \times \mathbb{R}^3} \frac{id^4q}{(2\pi)^4} \frac{1}{D_1 D_2 D_3 D_4}$$

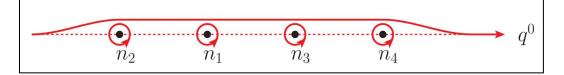
$$g_2^0 \qquad q_3^0 \qquad q_4^0$$

$$g_1^0 \qquad q_1^0 \qquad q_1^0 \qquad q_1^0$$

$$g_2^0 \qquad q_1^0 \qquad q_2^0$$

Discontinuity analysis from the definition

$$I(s, \boldsymbol{n}) = \int_{\mathcal{C}(\boldsymbol{n}) \times \mathbb{R}^3} \frac{id^4q}{(2\pi)^4} \frac{1}{D_1 D_2 D_3 D_4}$$





$$\mathbb{D}_E I(s, \boldsymbol{n}) = \int_{\mathcal{C}(\boldsymbol{n}) \times \mathbb{R}^3} \frac{i d^4 q}{(2\pi)^4} \, \mathbb{D}_E \left( \frac{1}{D_1 D_2 D_3 D_4} \right)$$



$$\mathbb{D}_E I(s, \boldsymbol{n}) = \int_{\mathcal{C}(\boldsymbol{n}) \times \mathbb{R}^3} \frac{i d^4 q}{(2\pi)^4} \left[ (2\pi i)^2 \left( \frac{\delta_1 \delta_3}{D_2 D_4} + \frac{\delta_1 \delta_4}{D_2 D_3} + \frac{\delta_2 \delta_3}{D_1 D_4} + \frac{\delta_2 \delta_4}{D_1 D_3} \right) \right]$$

#### Homotopy class symbol

$$\boldsymbol{n} = (n_1, n_2, n_3, n_4) \in \mathbb{Z}^4$$

$$\mathbb{D}_{E}\left(\frac{1}{D_{1}}\right) = -2\pi i \delta\left(q^{0} + E/2 - \omega_{1}\right) \equiv -2\pi i \delta_{1}$$

$$\mathbb{D}_{E}\left(\frac{1}{D_{2}}\right) = -2\pi i \delta\left(q^{0} + E/2 + \omega_{1}\right) \equiv -2\pi i \delta_{2}$$

$$\mathbb{D}_E\left(\frac{1}{D_3}\right) = +2\pi i \delta\left(q^0 - E/2 + \omega_2\right) \equiv +2\pi i \delta_3$$

$$\mathbb{D}_E\left(\frac{1}{D_4}\right) = +2\pi i \delta\left(q^0 - E/2 - \omega_2\right) \equiv +2\pi i \delta_4$$

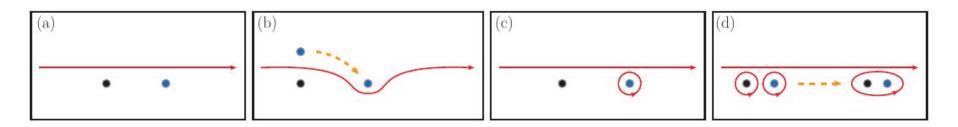


Discontinuity analysis from the definition

$$\mathbb{D}_{E}I(s, \boldsymbol{n}) = \int_{\mathcal{C}(\boldsymbol{n}) \times \mathbb{R}^{3}} \frac{id^{4}q}{(2\pi)^{4}} \left[ (2\pi i)^{2} \left( \frac{\delta_{1}\delta_{3}}{D_{2}D_{4}} + \frac{\delta_{1}\delta_{4}}{D_{2}D_{3}} + \frac{\delta_{2}\delta_{3}}{D_{1}D_{4}} + \frac{\delta_{2}\delta_{4}}{D_{1}D_{3}} \right) \right]$$

Homotopy class symbol

$$\boldsymbol{n}=(n_1,n_2,n_3,n_4)\in\mathbb{Z}^4$$



$$\mathbb{D}_{E}I(s,\boldsymbol{n}) = (2\pi i)^{2} \int_{\mathbb{R}^{4}} \frac{id^{4}q}{(2\pi)^{4}} \left[ c_{13}(\boldsymbol{n}) \frac{\delta_{1}\delta_{3}}{D_{2}D_{4}} + c_{14}(\boldsymbol{n}) \frac{\delta_{1}\delta_{4}}{D_{2}D_{3}} + c_{23}(\boldsymbol{n}) \frac{\delta_{2}\delta_{3}}{D_{1}D_{4}} + c_{42}(\boldsymbol{n}) \frac{\delta_{2}\delta_{4}}{D_{1}D_{3}} \right] \qquad \text{Winding number difference of the properties of the properties$$

Winding number difference

Physical case 
$$n_G = (0, 1, 1, 0)$$

$$\mathbb{D}_{E}G(s) = \mathbb{D}_{E}I(s, \mathbf{n}_{G}) = (2\pi i)^{2} \int_{\mathbb{R}^{4}} \frac{id^{4}q}{(2\pi)^{4}} \left( -\frac{\delta_{1}\delta_{3}}{D_{2}^{+}D_{4}^{+}} - \frac{\delta_{2}\delta_{4}}{D_{1}^{+}D_{3}^{+}} \right) = -2i\rho(s) \theta(s - s_{+})$$

The **Cutkosky rules** for computing discontinuities is a special case within discontinuity analysis. R.E.Cutkosky, J.Math.Phys. 1 (1960) 429-433

Continuation of the two-point Green's function

$$\mathbb{D}G(s) = -2i\rho(s)\theta(s - s_{+})$$



$$\mathbb{T}_1 G(s) = (1 - \mathbb{K}_1)G(s) = G(s) + 2i\rho(s)$$



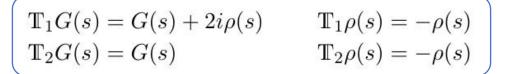
$$\mathbb{K}_1 \rho(s) = \mathbb{K}_2 \rho(s) = 2\rho(s)$$

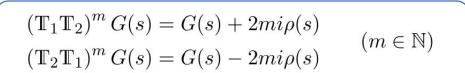
$$\mathbb{K}_1 G(s) = -2i\rho(s) \quad \mathbb{K}_2 G(s) = 0$$

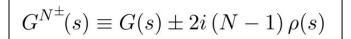


$$\mathbb{D}\rho(s) = 2\rho(s)\theta \left[ (s - s_+)(s - s_-) \right] - \frac{(m_1^2 - m_2^2)}{8i} \delta(s)$$

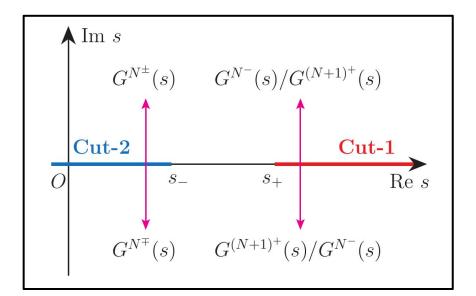
$$\theta[(s - s_{+})(s - s_{-})] = \theta(s - s_{+}) + \theta(s_{-} - s)$$







$$(N = I, II, III, \cdots)$$
  
 $G^{I^{\pm}}(s) \equiv G(s)$ 





Topological structure of the two-point Green's function

$$RS\{G(s)\} = \mathbb{C} \times \{\mathbb{T}_1, \mathbb{T}_2 | \mathbb{T}_1^2 \sim 1, \mathbb{T}_2^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}$$

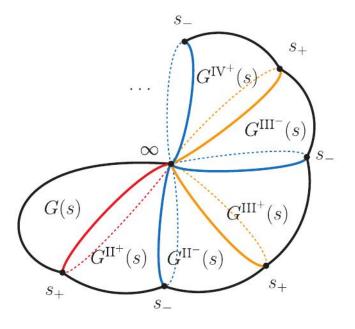
$$\mathbb{T}_1 G^{N^-}(s) = G^{(N+1)^+}(s), \quad \mathbb{T}_1 G^{(N+1)^+}(s) = G^{N^-}(s).$$
 $\mathbb{T}_2 G^{N^+}(s) = G^{N^-}(s), \quad \mathbb{T}_2 G^{N^-}(s) = G^{N^+}(s).$ 



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 $\mathbb{T}_2 G^{N^+}(s) = G^{N^-}(s), \quad \mathbb{T}_2 G^{N^-}(s) = G^{N^+}(s).$ 



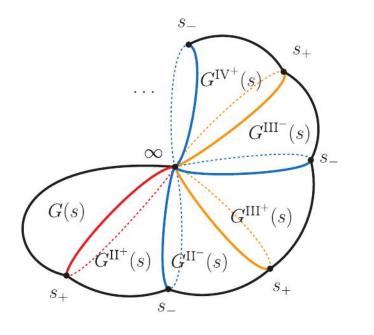
 $\overline{\mathrm{RS}}\{G(s)\}$ 

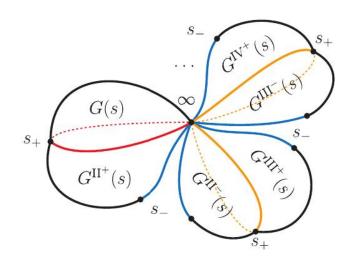


Topological structure of the two-point Green's function

$$RS\{G(s)\} = \mathbb{C} \times \{\mathbb{T}_1, \mathbb{T}_2 | \mathbb{T}_1^2 \sim 1, \mathbb{T}_2^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}$$

$$\mathbb{T}_1 G^{N^-}(s) = G^{(N+1)^+}(s), \quad \mathbb{T}_1 G^{(N+1)^+}(s) = G^{N^-}(s).$$
 $\mathbb{T}_2 G^{N^+}(s) = G^{N^-}(s), \quad \mathbb{T}_2 G^{N^-}(s) = G^{N^+}(s).$ 





 $\overline{RS}\{G(s)\}\$ 

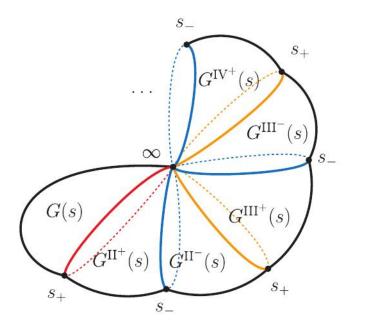
 $RS\{G(s)\}\setminus \mathbb{T}_2 = \mathbb{C} \times \{\mathbb{T}_1 \mid \mathbb{T}_1^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}_2$ 

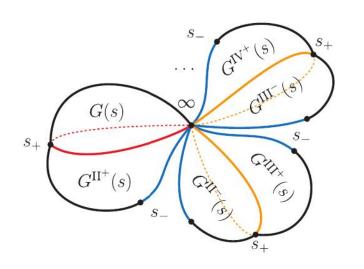
#### Discontinuity Analysis of the Two-Point Green's Function

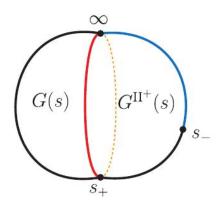
Topological structure of the two-point Green's function

$$RS\{G(s)\} = \mathbb{C} \times \{\mathbb{T}_1, \mathbb{T}_2 | \mathbb{T}_1^2 \sim 1, \mathbb{T}_2^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}$$

$$\mathbb{T}_1 G^{N^-}(s) = G^{(N+1)^+}(s), \quad \mathbb{T}_1 G^{(N+1)^+}(s) = G^{N^-}(s).$$
 $\mathbb{T}_2 G^{N^+}(s) = G^{N^-}(s), \quad \mathbb{T}_2 G^{N^-}(s) = G^{N^+}(s).$ 







$$\overline{\mathrm{RS}}\{G(s)\}$$

$$RS\{G(s)\}\setminus \mathbb{T}_2 = \mathbb{C} \times \{\mathbb{T}_1 \mid \mathbb{T}_1^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}_2$$

$$\overline{\mathrm{RS}}\{G(s)\}\backslash \mathbb{T}_2\cong S^2$$

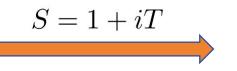
## Contents

- Motivation
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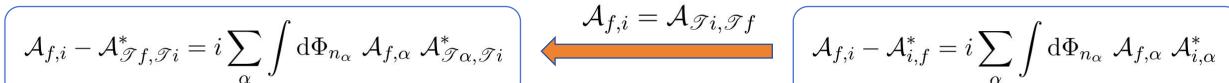
Unitarity condition

$$SS^{\dagger} = 1$$



$$T - T^{\dagger} = iT \, T^{\dagger}$$

Analytical condition



$$\mathcal{A}_{f,i} = \mathcal{A}_{\mathscr{T}i,\mathscr{T}f}$$

$$\mathcal{A}_{f,i} - \mathcal{A}_{i,f}^* = i \sum_{\alpha} \int d\Phi_{n_{\alpha}} \, \mathcal{A}_{f,\alpha} \, \mathcal{A}_{i,\alpha}^*$$

Unitarity condition for the partial-wave scattering matrix

$$T(s) - T^*(s) = i T(s) \cdot \varrho(s) \cdot T^*(s)$$



$$\left| \mathbf{T}(s) - \mathbf{T}^*(s) = i \; \mathbf{T}(s) \cdot \varrho(s) \cdot \mathbf{T}^*(s) \right| \leftarrow \left( \mathcal{A}_{fi} = \sum_{a,b} F_{fa}^* F_{ib} \; \mathbf{T}_{ab} \right)$$

$$\int d\Phi_{n_i} F_{ia}^* F_{ib} = 2\rho_a \delta_{ab}$$

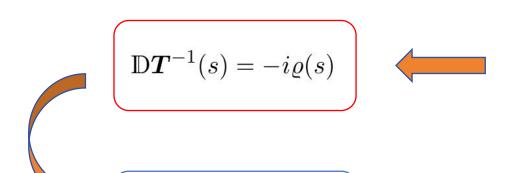
$$\varrho(s) = \sum_{a} \varrho_{a}(s) \, \theta(s - s_{a})$$

$$[\varrho_a(s)]_{bc} = 2\rho_a(s)\delta_{ab}\delta_{ac}$$

Discontinuity

$$T^*(s+i0^+) = T(s-i0^+) = T(s) - \mathbb{D}T(s)$$





$$\mathbb{D}\boldsymbol{T}(s) = i \; \boldsymbol{T}(s) \cdot \varrho(s) \cdot [\boldsymbol{T}(s) - \mathbb{D}\boldsymbol{T}(s)]$$

$$\mathbb{D}f(z) = -f(z)\,\mathbb{D}[1/f(z)]\,[f(z) - \mathbb{D}f(z)]$$

$$\mathbb{D}\boldsymbol{T}^{-1}(s) = \mathbb{D}\boldsymbol{G}(s)$$

$$G(s) \equiv \operatorname{diag}\{G_1(s), G_2(s), \cdots, G_{n_c}(s)\}$$



$$T(s) = [h(s) + G(s)]^{-1}$$

Continuation

$$\mathbb{T}_i \boldsymbol{T}(s) = \left[ \boldsymbol{h}(s) + \mathbb{T}_i \boldsymbol{G}(s) \right]^{-1}$$



$$RS\{T(s)\} \cong RS\{G(s)\}$$

$$\mathbb{D}G(s) = -i\sum_{a=1}^{n_c} \theta_a(s) \sum_{b=1}^{a} \varrho_b(s)$$

$$\mathbb{D}G(s) = -i\sum_{a=1}^{n_c} \theta_a(s) \sum_{b=1}^{a} \varrho_b(s) \qquad \qquad \theta(s - s_a) = \sum_{b=a}^{n_c} \theta_{[s_b, s_{b+1}]}(s) \equiv \sum_{b=a}^{n_c} \theta_b(s)$$

$$\mathbb{K}_{1a}\boldsymbol{G}(s) = -i\sum_{b=1}^{a} \varrho_b(s) \qquad (a = 1, \dots, n_c) \qquad \mathbb{T}_{1a}\boldsymbol{G}(s) = \boldsymbol{G}(s) + i\sum_{b=1}^{a} \varrho_b(s)$$

$$\mathbb{T}_{1a}\mathbf{G}(s) = \mathbf{G}(s) + i\sum_{b=1}^{a} \varrho_b(s)$$

$$\mathbb{D}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = i\sum_{b=1}^{n_c} \theta_b(s) \left[\sum_{c=1}^{\min\{a,b\}} \varrho_c(s) - \sum_{c>\min\{a,b\}}^{b} \varrho_c(s)\right] + \text{"Cut-2 terms"} + \text{"Pole-terms"}$$

$$\mathbb{K}_{1b}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = i\left[\sum_{c=1}^{\min\{a,b\}} \varrho_c(s) - \sum_{c>\min\{a,b\}}^{b} \varrho_c(s)\right] \qquad \mathbb{T}_{1b}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = \boldsymbol{G}(s) + i\sum_{c>\min\{a,b\}}^{\max\{a,b\}} \varrho_c(s)\right]$$

$$\mathbb{T}_{1b}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = \boldsymbol{G}(s) + i \sum_{c > \min\{a,b\}}^{\max\{a,b\}} \varrho_c(s)$$

HJJ, Xiong-Hui Cao and Feng-Kun Guo [arXiv:2507.06175]

$$\mathbb{T}_{1b}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = \mathbb{T}_{1a}\left[\mathbb{T}_{1b}\boldsymbol{G}(s)\right]$$

$$\mathbb{T}_{1a}\left[\mathbb{T}_{1a}\boldsymbol{G}(s)\right] = \boldsymbol{G}(s)$$

$$RS\{\boldsymbol{G}(s)\} = \mathbb{C} \times \{\mathbb{T}_{11}, \cdots, \mathbb{T}_{1n_c} | \mathbb{T}_{1a}^2 \sim 1, \mathbb{T}_{1a} \mathbb{T}_{1b} \sim \mathbb{T}_{1b} \mathbb{T}_{1a} \text{ (for any } a, b)\} \cong \mathbb{C} \times \mathbb{Z}_2^{n_c}$$

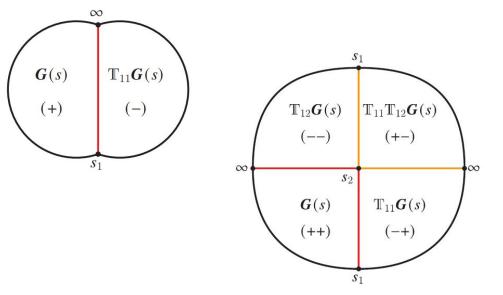


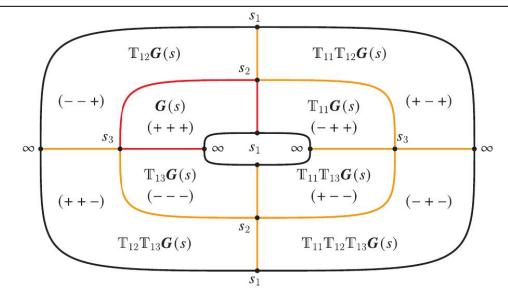
• The topological structure of the Riemann surface

$$n_c = 1$$
 
$$RS\{G(s)\} = \mathbb{C} \times \{T_{11} | T_{11}^2 \sim 1\} \cong \mathbb{C} \times \mathbb{Z}_2$$

$$n_c = 3$$

$$RS\{G(s)\} = \mathbb{C} \times \{\mathbb{T}_{11}, \mathbb{T}_{12}, \mathbb{T}_{13} | \mathbb{T}_{1a}^2 \sim 1, \mathbb{T}_{1a} \mathbb{T}_{1b} \sim \mathbb{T}_{1b} \mathbb{T}_{1a} (a, b = 1, 2, 3)\} \cong \mathbb{C} \times \mathbb{Z}_2^3$$



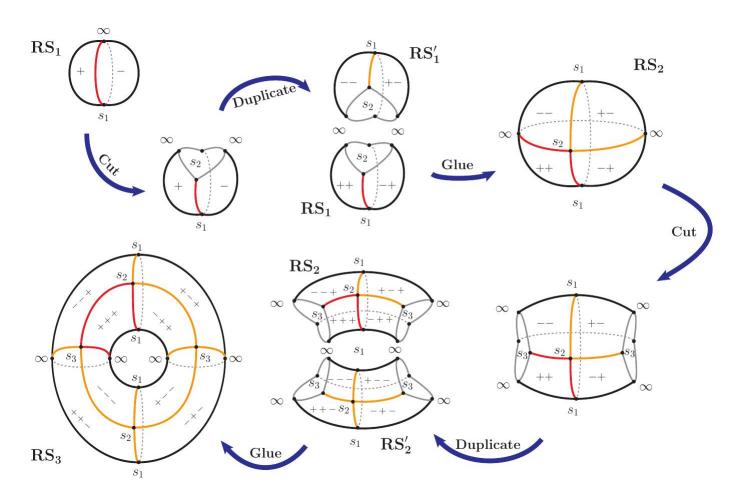


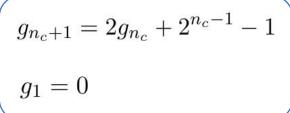
$$n_c = 2$$

$$RS\{G(s)\} = \mathbb{C} \times \{\mathbb{T}_{11}, \mathbb{T}_{12} | \mathbb{T}_{1a}^2 \sim 1 \ (a = 1, 2), \mathbb{T}_{11}\mathbb{T}_{12} \sim \mathbb{T}_{12}\mathbb{T}_{11}\} \cong \mathbb{C} \times \mathbb{Z}_2^2$$



• The topological structure of the Riemann surface







$$g_{n_c} = (n_c - 3)2^{n_c - 2} + 1$$



$n_c$	1	2	3	4	5	
$g_{n_c}$	0	0	1	5	17	

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- Scattering amplitudes are multivalued functions of invariant masses.
  - □ In the complex energy plane, right-hand cuts, left-hand cuts, circular cuts, and other complex analytic structures pose challenges for amplitude analysis.
  - □ Origin of multivaluedness: The Riemann surface for scattering amplitudes forms a multiple covering of the complex plane.

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- Uniformization

Conformally map the Riemann surface onto the interior of the Riemann sphere



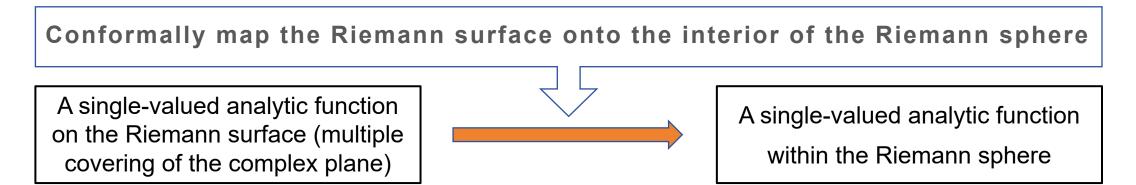
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Conformally map the Riemann surface onto the interior of the Riemann sphere

A single-valued analytic function on the Riemann surface (multiple covering of the complex plane)

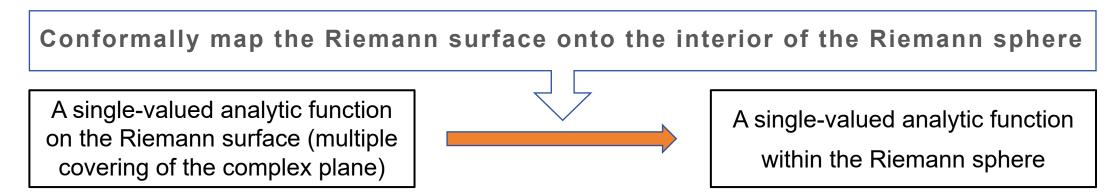


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- Uniformization

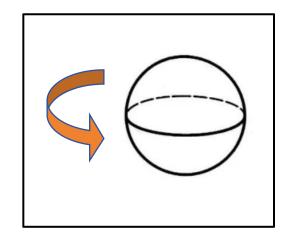


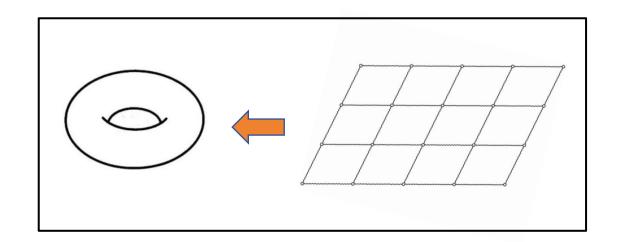
• This is mathematically equivalent to finding the **simply connected covering space** of the Riemann surface for the scattering amplitude — the **Universal Covering**.

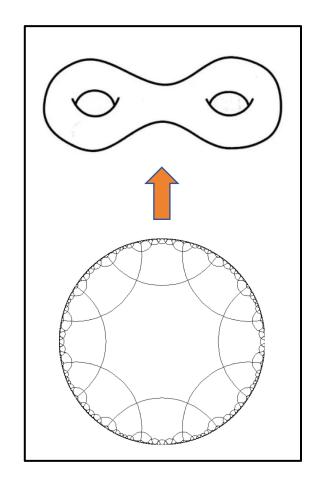


• Uniformization theorem for closed orientable Riemann surfaces

$\overline{\text{Genus }g}$	Universal covering	Covering mapping		
0	$\text{Riemann sphere } \hat{\mathbb{C}}$	Rational function		
1	Complex plane $\mathbb C$	Elliptic function		
$\geq 2$	$\text{Unit disk } \mathbb{D}$	Automorphic function		



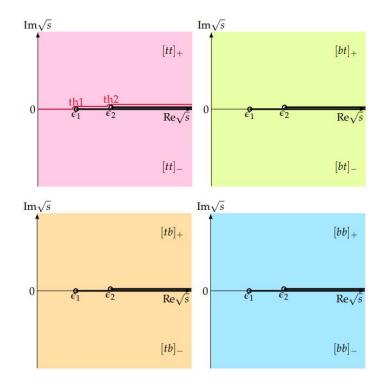


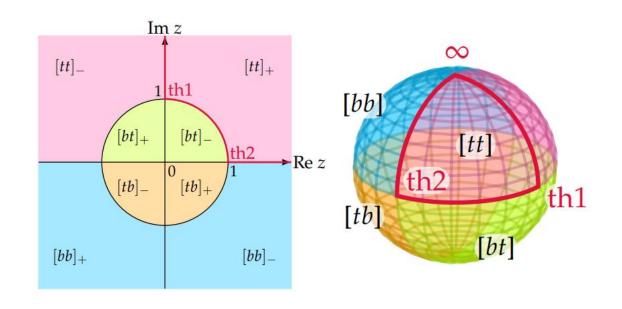


• 2-channel case g=0

M.Kato, Annals Phys. 31 (1965) 1, 130-147

$$q_1 = \frac{\Delta_{12}}{2} \left( z + \frac{1}{z} \right)$$
  $q_2 = \frac{\Delta_{12}}{2} \left( z - \frac{1}{z} \right)$   $\Delta_{ij} = \sqrt{|s_i - s_j|}$   $q_i = \sqrt{s - s_i}$ 





W.Yamada, O.Morimatsu, and T.Sato, Phys. Rev. Lett. 129, 192001 (2022)

Rez

# Uniformization of the Partial-Wave Scattering Matrix

• 3-channel case g=1

W.Yamada, O.Morimatsu, and T.Sato, Phys. Rev. Lett. 129, 192001 (2022)

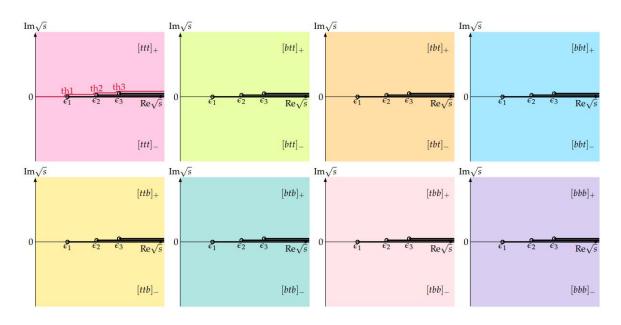
$$q_{1} = \frac{\Delta_{12}}{2} \left( \frac{\sin(4K(1/\gamma^{2})z, 1/\gamma^{2})}{\gamma} + \frac{\gamma}{\sin(4K(1/\gamma^{2})z, 1/\gamma^{2})} \right)$$

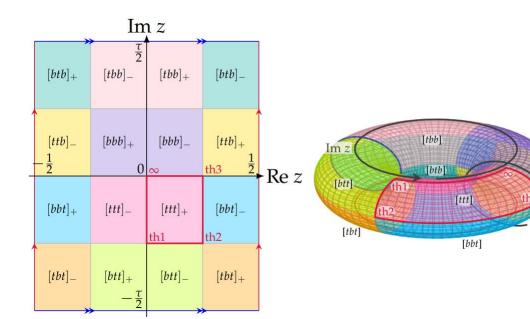
$$q_{2} = \frac{\Delta_{12}}{2} \left( \frac{\sin(4K(1/\gamma^{2})z, 1/\gamma^{2})}{\gamma} - \frac{\gamma}{\sin(4K(1/\gamma^{2})z, 1/\gamma^{2})} \right)$$

$$q_{3} = \frac{\Delta_{12}}{2} \frac{\gamma \sin'(4K(1/\gamma^{2})z, 1/\gamma^{2})}{\sin(4K(1/\gamma^{2})z, 1/\gamma^{2})}$$

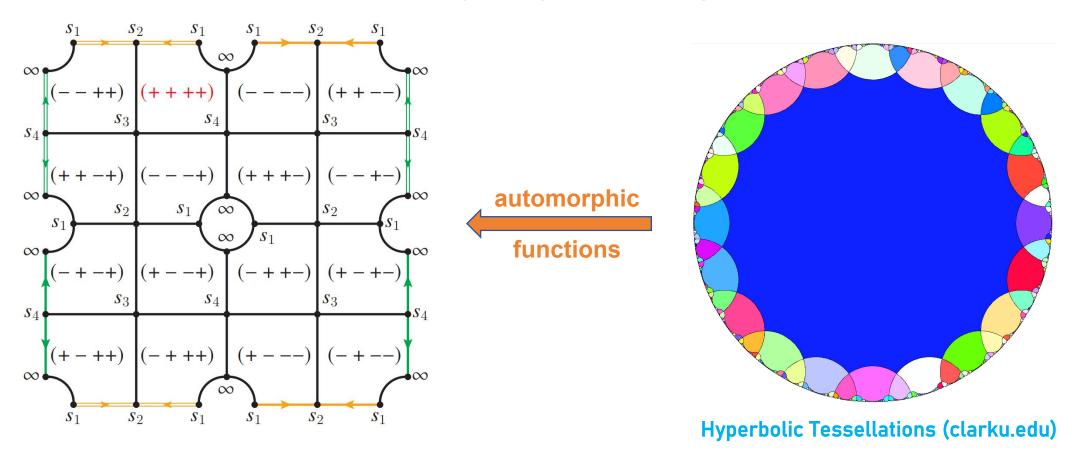
$$\operatorname{sn}^{-1}(v,k) = \int_0^v \frac{\mathrm{d}v'}{\sqrt{(1-v'^2)(1-k^2v'^2)}}$$

$$\gamma = \frac{\Delta_{13} + \Delta_{23}}{\Delta_{12}} \qquad K(k) = \int_0^1 \frac{dv}{\sqrt{(1 - v^2)(1 - k^2 v^2)}}$$





- 4-channel case g=5
  - The Riemann surface comprises 16 sheets and can be conformally mapped via automorphic functions onto a regular hyperbolic 20-gon within the unit disk



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# Summary and Outlook

- **Discontinuity Calculus**: A systematic framework for handling analytic continuation of complex functions and the topological structure of Riemann surfaces.
- Discontinuity analysis applied to scattering amplitudes in two-body scattering problems.

- Crossing symmetry: Investigate discontinuities in scattering amplitudes within unphysical regions.
- Three- and four-body processes: Extend discontinuity analysis to multivariate complex functions.
- Multi-loop amplitudes: Analyze analytic continuation and topological structures of Riemann surfaces.
- ... and beyond

# Thank you for your attention!