



Hunt for $X(17)$

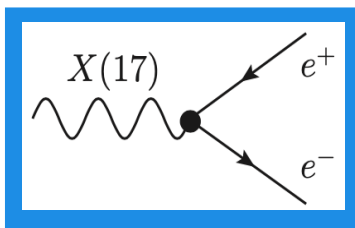
Jun Jiang, Cong-Feng Qiao (UCAS), Yu-Han Zhao (SDU),
arXiv: 2511.xxxx

Jun Jiang (蒋 军)
Shandong University
2025-10-20 @ 南京师范大学, 南京

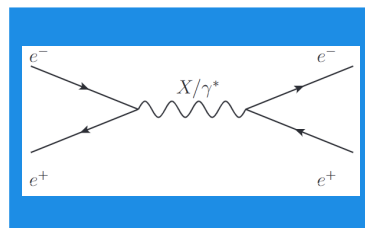
CONTENTS



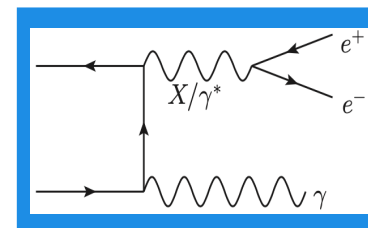
Motivation



Constraints on
 Xe^+e^- coupling



$X(17)$
@ PADME



Revisit $X(17)$
@ BESIII



Summary



1 Motivation

Observation of Anomalous Internal Pair Creation in Be8 : A Possible Indication of a Light, Neutral Boson

Citations per year

A. J. Krasznahorkay (Debrecen, Inst. Nucl. Res.), M. Csatlós (KLTE-ATOMKI), L. Csige (KLTE-ATOMKI), Z. C. ...
Apr 7, 2015

5 pages

Published in: *Phys.Rev.Lett.* 116 (2016) 4, 042501

Published: Jan 26, 2016

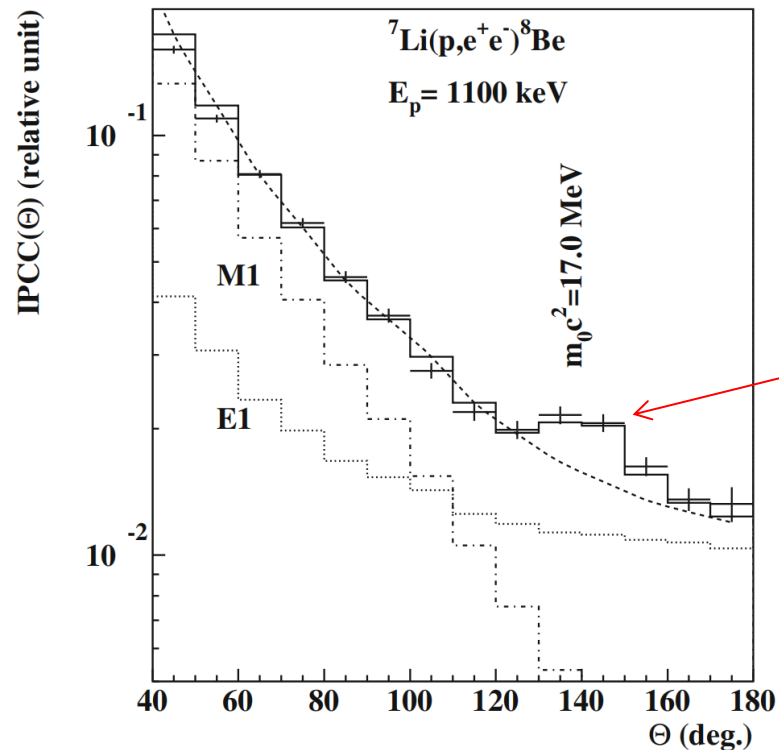
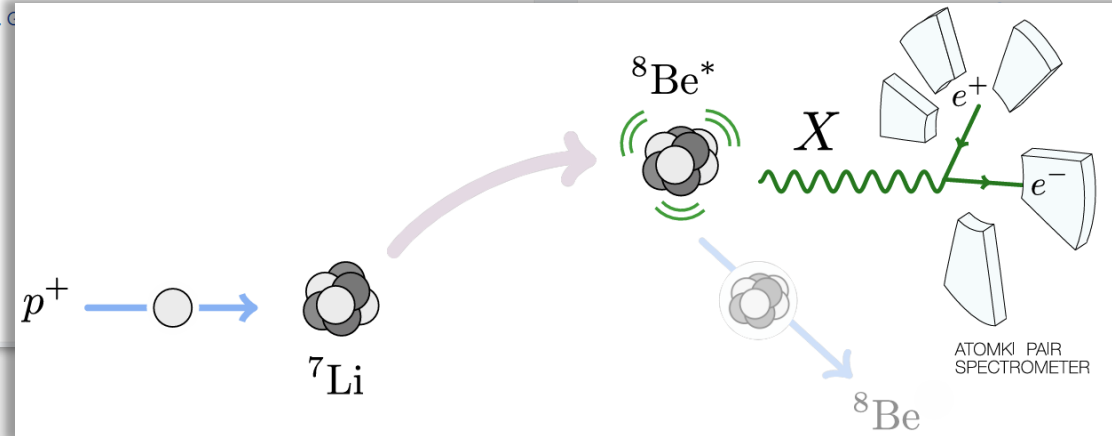
e-Print: 1504.01527 [nucl-ex]

DOI: 10.1103/PhysRevLett.116.042501

PDG: Heavy Particle Production Cross Section

View in: CERN Document Server, ADS Abstract Service, Nuclear Science References

[pdf](#) [links](#) [cite](#) [claim](#)



electron-positron invariant mass and opening angle are related by

$$m_{ee}^2 = 2E_{e^+}E_{e^-} - 2\sqrt{E_{e^+}^2 - m_e^2}\sqrt{E_{e^-}^2 - m_e^2}\cos\theta + 2m_e^2$$

$$= (1 - y^2)E^2 \sin^2 \frac{\theta}{2} + 2m_e^2 \left(1 + \frac{1 + y^2}{1 - y^2} \cos\theta \right) + \mathcal{O}(m_e^4)$$

where

$$E \equiv E_{e^+} + E_{e^-} \quad \text{and} \quad y \equiv \frac{E_{e^+} - E_{e^-}}{E_{e^+} + E_{e^-}}$$

Phys. Rev. Lett. 116, 042501 (2016)

Eur. Phys. J. C (2023) 83 :230

Phys.Rev.D 95 (2017) 3, 035017



1 Motivation

4

Nucleus (MeV)	m_X (MeV)	Experiment	Ref.
$^8\text{Be}^*(18.15)$	$16.86 \pm 0.06 \pm 0.50$	Atomki	[2, 6]
$^8\text{Be}^*(18.15)$	$17.17 \pm 0.07 \pm 0.20$	Atomki	[6]
$^4\text{He}^*(20.21/21.01)$	$16.94 \pm 0.12 \pm 0.21$	Atomki	[9]
$^{12}\text{C}^*(17.23)$	$17.03 \pm 0.11 \pm 0.20$	Atomki	[10]
$^8\text{Be}^*(\text{GDR})$	$16.95 \pm 0.48 \pm 0.35$	Atomki	[11, 12]
$^8\text{Be}^*(18.15)$	$16.66 \pm 0.47 \pm 0.35$	VNU-UoS	[13]
$^8\text{Be}^*(17.64/18.15)$	$< 16.81 [R_{\text{Be}} = 6 \cdot 10^{-6}]$	MEG II	[17]
$e^+e^- \rightarrow X_{17}$	$16.90 \pm 0.02 \pm 0.05$	PADME	[20, 21]

More experiments are on the way:

Mu3e experiment

Montreal Tandem accelerator

New JEDI experiment at GANIL

nTOF collaboration at CERN

PRad experiment at JLab

...

Light Dark Matter 2025

2025年4月8日至11日
Europe/Rome 时区

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Wifi password: INFN4All!

概览

日程表

报告列表

Registration

Venue

Scientific Committees

Accommodation

Registration and fee

分会

Searches for X17 at accelerators

2025年4月10日 09:00

报告材料

还没有材料

报告列表 日程表

周四 10/04

打印 PDF 全屏 详细视图 过滤

09:00

The X17 boson: status and prospects
Claudio Toni
09:00 - 09:25

Kinematical evidence for the X17 particle
Attila Krasznahorkay
09:25 - 09:50

Recent X17 results from the MEG experiment
Francesco Renga
09:50 - 10:15

The X17 experiment at LNL
Tommaso Marchi
10:15 - 10:40

Discussion
10:40 - 11:00

11:00

arXiv:2504.11439

Universe 10 (2024) 11, 409

Phys. Rev. C 14, 28 (1976)

<https://agenda.infn.it/event/43758/sessions/32329/#20250410>



1 Motivation

5



J. L. Feng et al., “Protophobic **Fifth-Force Interpretation** of the Observed Anomaly in Be^8 Nuclear Transitions”

Phys. Rev. Lett. 117(2016), 071803
Phys.Rev.D 95 (2017), 035017
Phys.Rev.D 102 (2020), 036016



A. C. Hayes et al., “Angular correlations in the e^+e^- decay of excited states in Be^8 ”

Phys.Rev.C 105 (2022), 055502

- this resonance is dominated by M1 and E1 decay
- the existence of a ‘bump’ in the measured angular distribution is strongly dependent on the assumed M1/E1 ratio, which is a strong function of energy, while it was assumed to be a constant in the Atomki experiments.



Xilin Zhang et al., “Can nuclear physics explain the anomaly observed in the internal pair production in the Beryllium-8 nucleus?” “Can a protophobic vector boson explain the ATOMKI anomaly?”

- study the possibility of using the nuclear transition form factor to explain the anomaly
- The net result is that X production is dominated by direct transitions induced by $E1^X$ and $L1^X$ (transverse and longitudinal electric dipoles) and $C1^X$ (charge dipole) without going through any nuclear resonance (i.e. Bremsstrahlung radiation) with a smooth energy dependence that occurs for all proton beam energies above threshold. This contradicts the experimental observations and invalidates the protophobic vector boson explanation.

Phys.Lett.B 773 (2017) 159-165
Phys.Lett.B 813 (2021) 136061

More Info.

Eur. Phys. J. C (2023) 83:230
<https://doi.org/10.1140/epjc/s10052-023-11271-x>

THE EUROPEAN
PHYSICAL JOURNAL C



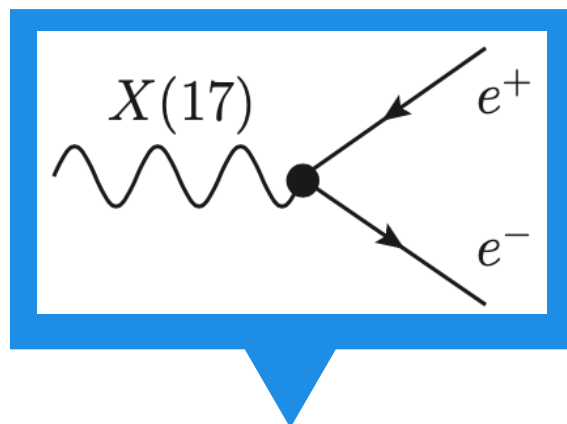
Review

Shedding light on X17: community report

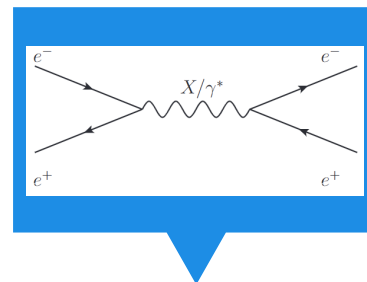
CONTENTS



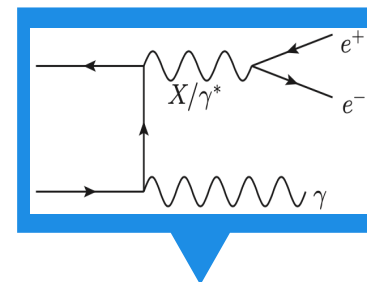
Motivation



Constraints on
 Xe^+e^- coupling



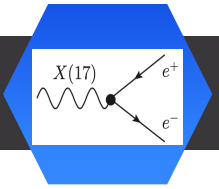
$X(17)$
@ PADME



Revisit $X(17)$
@ BESIII



Summary

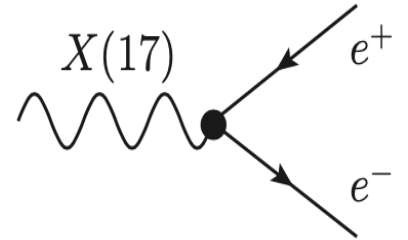


2 Constraints on Xe^+e^- Coupling

7

■ Mixing of V +/- A model:

$$\mathcal{L}_X = -\frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{1}{2}m_X^2 X_\mu X^\mu - \sum_f e\bar{f}\gamma_\mu(\varepsilon_f^v - \varepsilon_f^a\gamma_5)fX^\mu$$



With such model, we suggest identifying X(17) in **(1) $e^+e^- \rightarrow X(17) + g$** ,
(2) $J/\psi \rightarrow X(17) + g$ at BESIII, where X(17) is reconstructed by e^+e^- pairs.

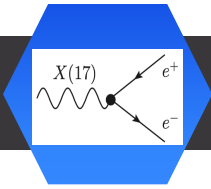
Eur. Phys. J. C 78, 456 (2018)
 Eur. Phys. J. C 79, 404 (2019)

- Scalar X(17) hypothesis is excluded by parity conservation in ^8Be experiment; while a pseudoscalar state can explain only the ^8Be and ^4He anomaly, but not the ^{12}C one.

JHEP 02 (2023) 154,
 JHEP 07 (2023) 168 (erratum)

- N.V. Krasnikov, “Implications of last NA64 results and the electron g_e-2 anomaly for the X(16.7) boson survival”
- The combination of latest NA64 experiment and recent value of the electron anomalous magnetic momentum exclude the pure vector and axial-vector couplings of X(17) boson to electrons at the 90% C.L.,
 - but the model with “V $\pm$ A” interaction still survives and can further explains both the electron and muon anomalous magnetic momentum.

Mod.Phys.Lett.A 35 (2020), 2050116



2 Constraints on Xe^+e^- Coupling

8

■ Atomki experiments

Decay length

$$L = \frac{\hbar \times c}{\Gamma(X \rightarrow e^+e^-)} \times \frac{v}{\sqrt{1-v^2}},$$

$$\Gamma(X \rightarrow e^+e^-) \approx \frac{\alpha m_X}{3} ((\varepsilon_e^v)^2 + (\varepsilon_e^a)^2),$$

$$\mathcal{O}\left(\frac{m_e^2}{m_X^2}\right) \approx 0.1\%$$

2016, ^8Be ; 2019/2021, ^4He ; 2022, ^{12}C

Requiring $L < 1 \text{ cm}$

$$\frac{m_X = 16.7 \pm 0.35(\text{stat.}) \pm 0.5(\text{sys.}) \text{ MeV},}{\frac{\Gamma(^8\text{Be}^* \rightarrow ^8\text{Be}X) \text{Br}(X \rightarrow e^+e^-)}{\Gamma(^8\text{Be}^* \rightarrow ^8\text{Be}\gamma)} = 5.8 \times 10^{-6}.}$$

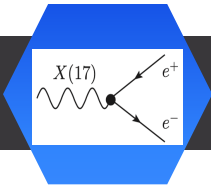
$$(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \gtrsim 1.8 \times 10^{-10}.$$

$$\frac{m_X = 16.94 \pm 0.12(\text{stat}) \pm 0.21(\text{sys}) \text{ MeV},}{\frac{\Gamma(^4\text{He}^* \rightarrow ^4\text{He}X) \text{Br}(X \rightarrow e^+e^-)}{\Gamma(^4\text{He}^* \rightarrow ^4\text{He}\gamma)} = (5.1 \pm 0.13) \times 10^{-6}.}$$

$$(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \gtrsim 3.5 \times 10^{-10}.$$

$$\frac{m_X = 17.03 \pm 0.11(\text{stat}) \pm 0.20(\text{sys}) \text{ MeV},}{\frac{\Gamma(^{12}\text{C}^* \rightarrow ^{12}\text{C}X) \text{Br}(X \rightarrow e^+e^-)}{\Gamma(^{12}\text{C}^* \rightarrow ^{12}\text{C}\gamma)} = (3.6 \pm 0.3) \times 10^{-6}.}$$

$$(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \gtrsim 7.7 \times 10^{-11}.$$



2 Constraints on Xe^+e^- Coupling

9

■ Electron anomalous magnetic momentum

- In 2022, an improved determination of the electron anomalous magnetic momentum (AMM) is obtained, $a_e^{\text{exp}} = 0.00115965218059(13)$, which is 2.2 times more accurately than the value in 2008.
- In comparison with the SM prediction in 2018, $a_e^{\text{SM}} = 0.00115965218161(23)$, we have the negative 3.9σ discrepancy between the measurement and the SM prediction

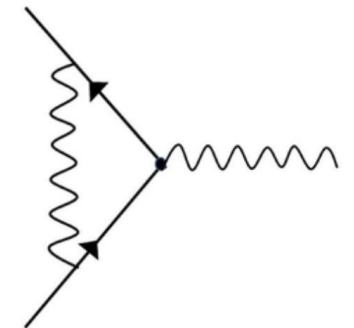
$$\Delta a_e \equiv a_e^{\text{exp}} - a_e^{\text{SM}} = (-1.02 \pm 0.26) \times 10^{-12}.$$

Phys.Rev.Lett. 130 (2023), 071801
Science 360 (2018), 191

- $X(17)$ boson contributes to Δa_e through the leading one-loop triangle diagram

$$a_e^X = \frac{\alpha}{3\pi} \left(\frac{m_e}{m_X} \right)^2 [(\varepsilon_e^v)^2 - 5(\varepsilon_e^a)^2].$$

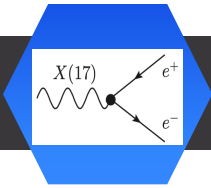
Nucl.Phys.B 137 (1978) 63-76



To diminish the 3.9σ discrepancy between the measurement and SM prediction,

$$(-1.5 \pm 0.4) \times 10^{-6} \lesssim (\varepsilon_e^v)^2 - 5(\varepsilon_e^a)^2 \lesssim 0.$$

The pure vector hypothesis of $X(17)$ boson will enlarge the discrepancy of electron AMM!



2 Constraints on Xe^+e^- Coupling

10

■ Beam dumb experiments

- Bremsstrahlung reaction $e^- Z \rightarrow e^- Z X$, X(17) is produced by initial or final state radiation off a single electron.
- Within the errors of $\mathcal{O}\left(\frac{m_e^2}{m_X^2}\right) \approx 0.1\%$, the production rate depends on the coupling parameters only through $(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2$

➤ Recently, NA64 collaboration searched for X(17) twice and set limits on the Xee coupling parameter at the 90% C. L.

Excluded region: $1.3 \times 10^{-4} < \sqrt{(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2} < 4.2 \times 10^{-4}$

Phys. Rev. Lett. 120 (2018), 231802

Phys. Rev. D 101 (2020) 7, 071101

Excluded region: $1.2 \times 10^{-4} < \sqrt{(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2} < 6.8 \times 10^{-4}$
(2017&2018 data combined)

➤ Around 17 MeV, E137, E144 Beam dumb experiments

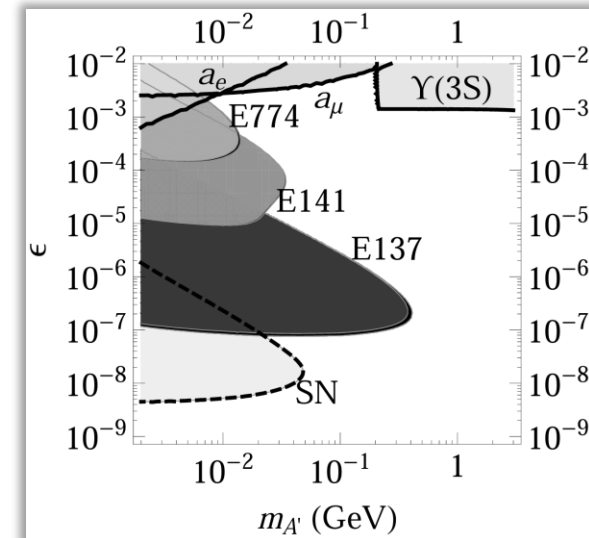
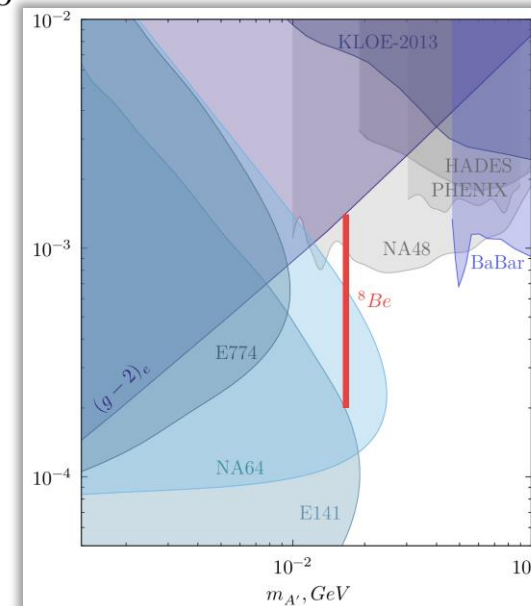
Excluded region:

$$8 \times 10^{-8} < \sqrt{(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2} < 3 \times 10^{-4}$$

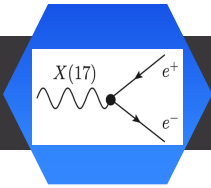
Survival region:

$$(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \gtrsim 4.6 \times 10^{-7}.$$

We do not discuss the $< 6.4 \times 10^{-15}$ region.



Phys.Rev.D 80 (2009) 075018



2 Constraints on Xe^+e^- Coupling

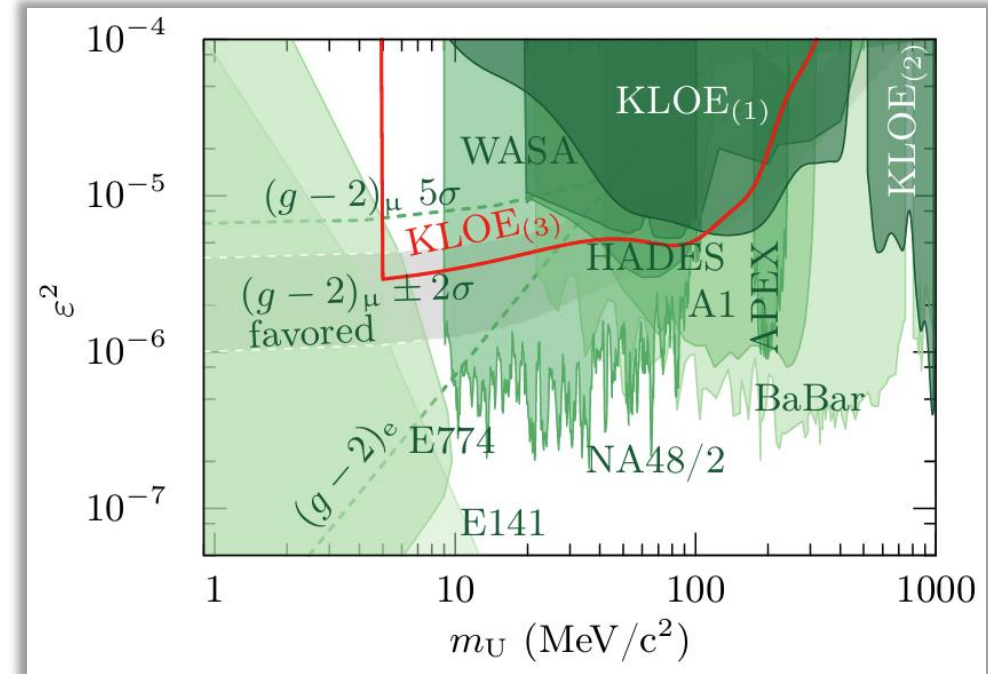
11

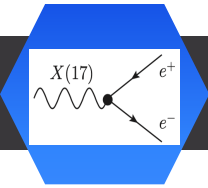
■ KLOE-2 experiment

- The KLOE-2 experiment searched for a low-mass vector boson U in $e^+e^- \rightarrow U\gamma$, $U \rightarrow e^+e^-$.
- Within the errors of $\mathcal{O}\left(\frac{m_e^2}{m_X^2}\right) \approx 0.1\%$, the production rate depends on the coupling parameters only through $(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2$

➤ An upper bound on Xee coupling around 17 MeV can be set at the 90% C. L.,

$$(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \lesssim 4 \times 10^{-6}.$$





2 Constraints on Xe^+e^- Coupling

12

■ Survival regions for Xe^+e^- coupling

Beam dumb
experiments

KLOE-2
experiment

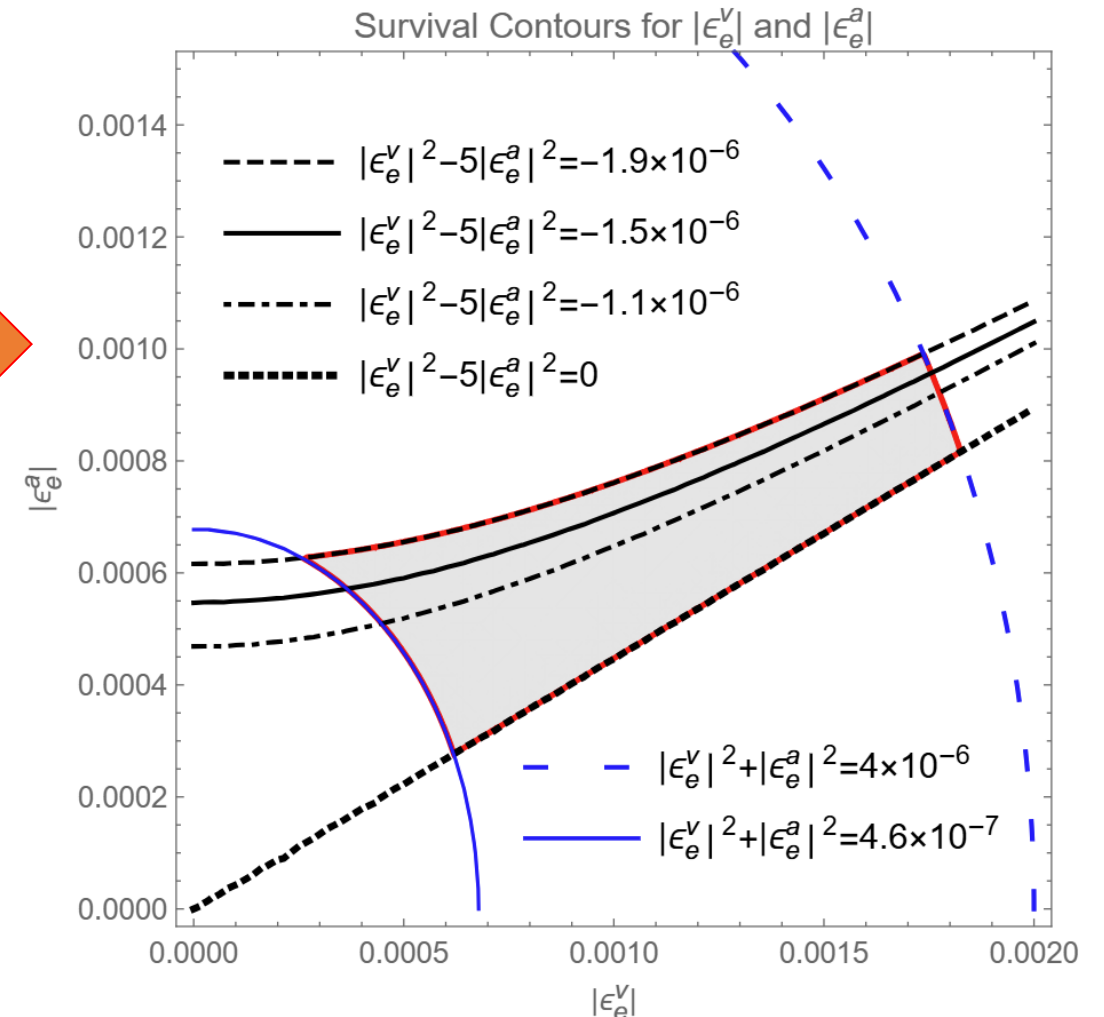
$$4.6 \times 10^{-7} \lesssim (\epsilon_e^v)^2 + (\epsilon_e^a)^2 \lesssim 4 \times 10^{-6},$$

$$(-1.5 \pm 0.4) \times 10^{-6} \lesssim (\epsilon_e^v)^2 - 5(\epsilon_e^a)^2 \lesssim 0,$$

Electron AMM

- Pure V hypothesis: disfavored by Electron AMM.
- Pure A hypothesis: $6.8 \times 10^{-4} \lesssim |\epsilon_e^a| \lesssim (5.4 \pm 0.7) \times 10^{-4}$.

Pure axial-vector coupling is excluded.

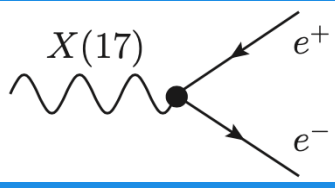


CONTENTS

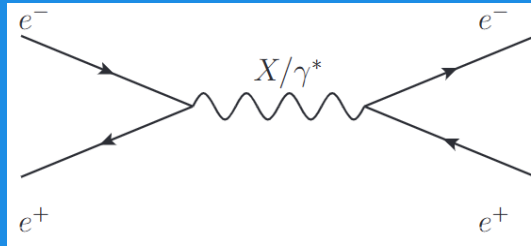
MOTIVATION



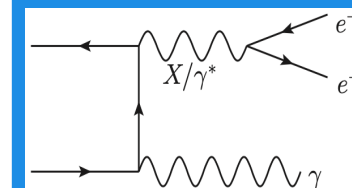
Motivation



Constraints on
 Xe^+e^- coupling



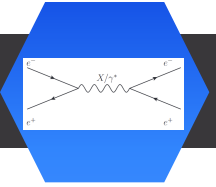
X(17) @ PADME



Revisit X(17)
@ BESIII

SUMMARY

Summary



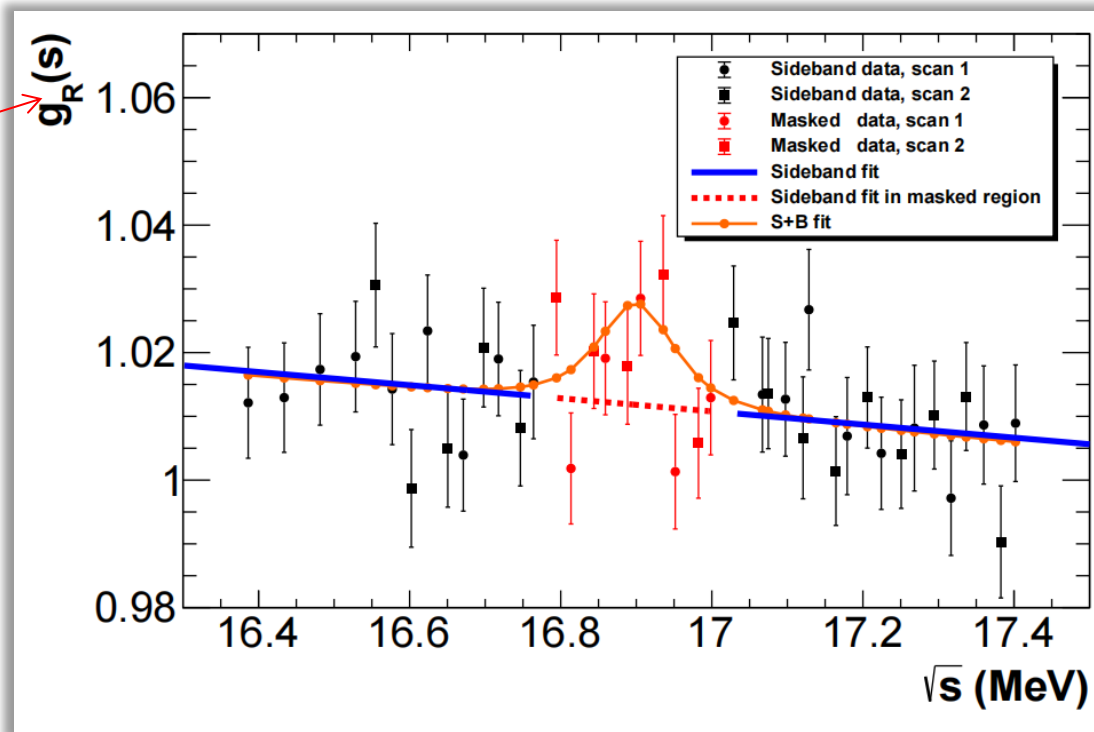
3 X(17) @ PADME

14

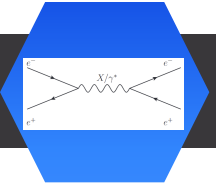
- PADME Experiment at the Frascati DAΦNE LINAC searched for X(17) using a positron beam incident on a fixed target
- The production rates of the Bhabha scattering $e^+e^- \rightarrow e^+e^-$ and the $e^+e^- \rightarrow \gamma\gamma$ process have been measured within the beam energy between 262 and 296 MeV, corresponding to center-of-mass energies $16.4 < \sqrt{s} < 17.4$ MeV

$$g_R(s) = K(s) \left[1 + \frac{S(s; M_X, g_{ve}) \varepsilon_{\text{sig}}(s)}{B(s)} \right]$$

Terms linearly dependent on $\sqrt{s}$ and are collectively referred to as $K(s)$



- Significant deviation is found at $\sqrt{s} \simeq 16.90$ MeV, with local significance of 2.5σ and global significance of $(1.77 \pm 0.15) \sigma$.
- By assuming the pure vector coupling, the upper limit $|\varepsilon_e^v| \lesssim 5.6 \times 10^{-4}$ is obtained at 90% C.L.



3 X(17) @ PADME

15

■ $e^+e^- \rightarrow X(17)$

$$\sigma = \frac{2\pi^2\alpha}{s}((\varepsilon_e^a)^2 + (\varepsilon_e^v)^2)\delta(1 - \frac{m_X}{\sqrt{s}}) + \mathcal{O}((\frac{m_e}{\sqrt{s}})^2).$$

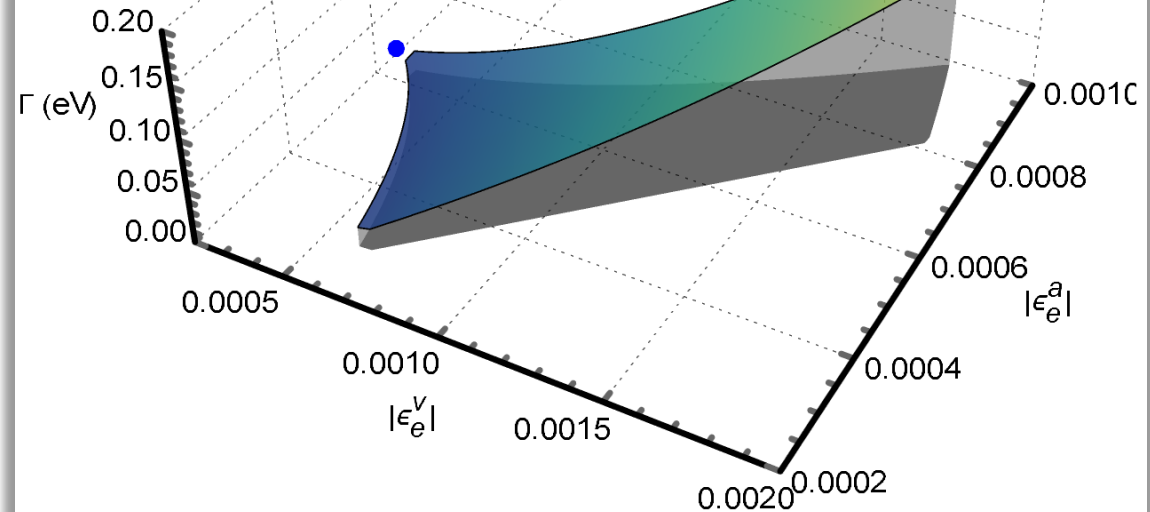
Cross section: 89 ~ 777 nb

■ $X(17) \rightarrow e^+e^-$

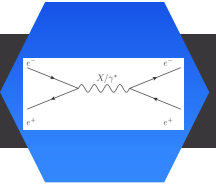
$$0.018 \text{ eV} \lesssim \Gamma(X \rightarrow e^+e^-) \lesssim 0.16 \text{ eV},$$

$$(|\varepsilon_e^v|, |\varepsilon_e^a|)_{\text{max}}^{\Gamma \rightarrow e^+e^-} \simeq (1.8, 0.84) \times 10^{-3},$$

$$(|\varepsilon_e^v|, |\varepsilon_e^a|)_{\text{min}}^{\Gamma \rightarrow e^+e^-} \simeq (2.4, 6.3) \times 10^{-4}.$$



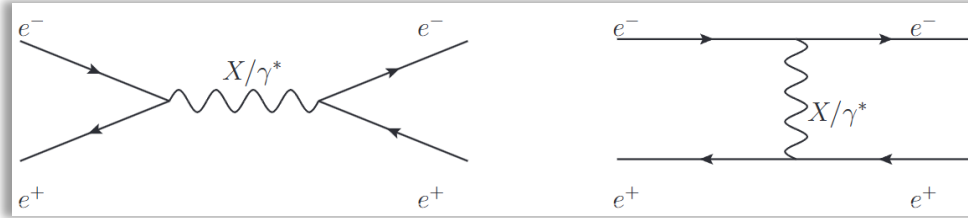
Lifetime of X(17) at rest frame: 4~36 fs.



3 X(17) @ PADME

16

■ $e^+e^- \rightarrow X(17)/\gamma^* \rightarrow e^+e^-$



Background (B)

$$\frac{d\sigma_B}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left(\frac{3 + \cos^2\theta}{1 - \cos\theta} \right)^2 + \mathcal{O}\left(\left(\frac{m_e}{\sqrt{s}}\right)^2\right).$$

$$\sigma_B = \frac{\pi\alpha^2}{2s} (I(\theta_2) - I(\theta_1)) + \mathcal{O}\left(\left(\frac{m_e}{\sqrt{s}}\right)^2\right),$$

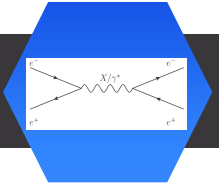
$$I(\theta) \equiv \frac{\cos^3\theta}{3} + \cos^2\theta + 9\cos\theta + \frac{16}{1 - \cos\theta} + 16\log(\cos\theta - 1) - \frac{31}{3},$$

divergent at $\cos\theta=1$ ($t=0$)

Signal (S)

$$\begin{aligned} \frac{d\sigma_S}{d\cos\theta} = & \frac{\pi\alpha^2 s}{2(s - m_X^2)^2(s + 2m_X^2 - s\cos\theta)^2} \left(((\varepsilon_e^v)^4 + (\varepsilon_e^a)^4) \right. \\ & \times (s^2(\cos^2\theta + 3)^2 - 6m_X^2 s(\cos^2\theta + 3)(\cos\theta + 1) + 2m_X^4(\cos\theta(5\cos\theta + 6) + 9)) \\ & + 2(\varepsilon_e^v \varepsilon_e^a)^2 (s^2(\cos\theta(\cos\theta(\cos\theta + 8) + 6) + 8) - 7) \\ & \left. - 2m_X^2 s(\cos\theta + 1)(\cos\theta(5\cos\theta + 16) - 1) + 2m_X^4(\cos\theta(11\cos\theta + 26) + 7) \right) \\ & + \mathcal{O}\left(\left(\frac{m_e}{\sqrt{s}}\right)^2\right). \end{aligned} \quad (30)$$

$$\begin{aligned} \sigma_S = & \frac{4\pi\alpha^2}{3m_X^2 s(s + m_X^2)(s - m_X^2)^2} \\ & (s((\varepsilon_e^v)^4 + (\varepsilon_e^a)^4)(3s^3 + 4m_X^2 s^2 + m_X^4 s - 6m_X^6) \\ & + 2(\varepsilon_e^v)^2(\varepsilon_e^a)^2(3s^3 + 22m_X^2 s^2 - 5m_X^4 s - 18m_X^6)) \\ & + 6((\varepsilon_e^v)^4 + (\varepsilon_e^a)^4 + 6(\varepsilon_e^v)^2(\varepsilon_e^a)^2) m_X^2(m_X^2 - s)(m_X^2 + s)^2 \log\left(\frac{m_X^2 + s}{m_X^2}\right) \\ & + \mathcal{O}\left(\left(\frac{m_e}{\sqrt{s}}\right)^2\right). \end{aligned}$$



3 X(17) @ PADME

17

■ $e^+e^- \rightarrow X(17)/\gamma^* \rightarrow e^+e^-$

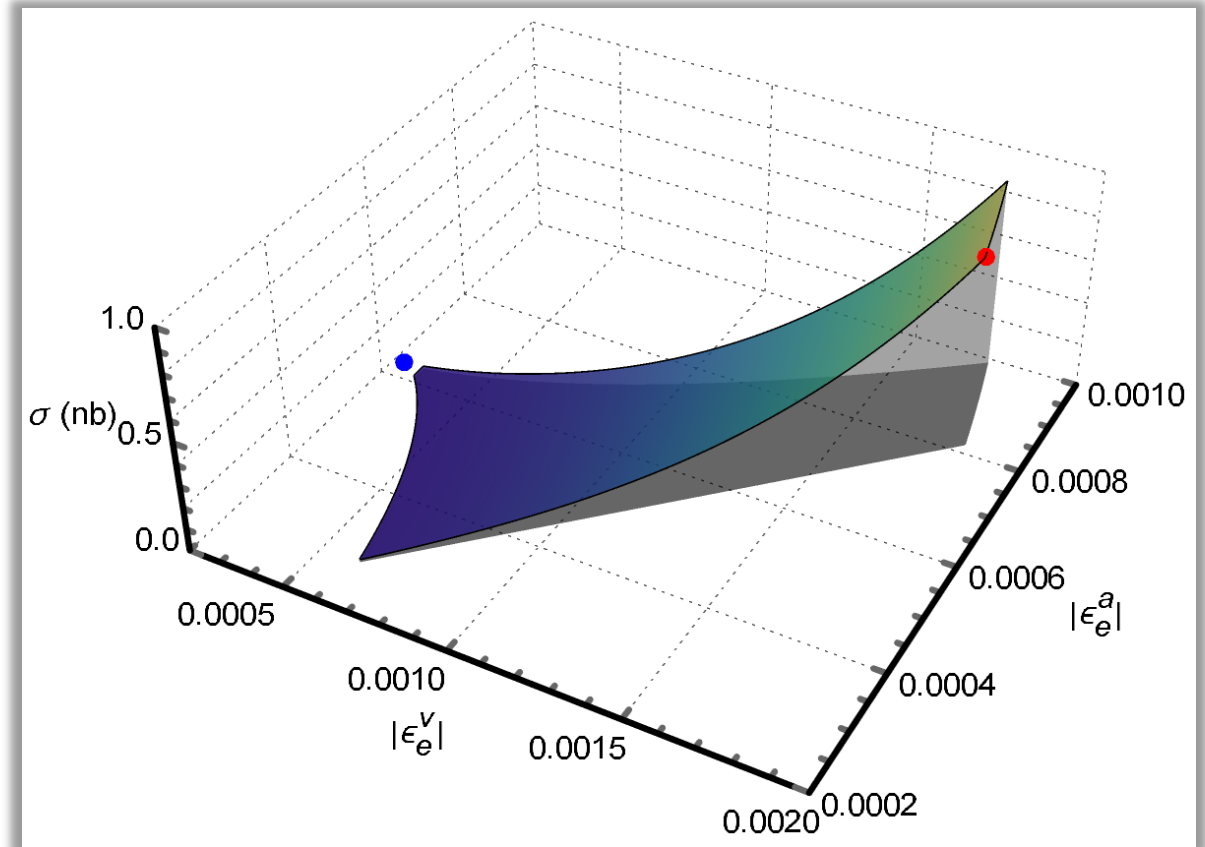
$m_X = 17 \text{ MeV}, \sqrt{s} = 17.02 \text{ MeV}, \Gamma_X = 0.089 \text{ eV}$

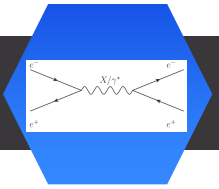
$0.011 \text{ nb} \lesssim \sigma(e^+e^- \rightarrow X \rightarrow e^+e^-) \lesssim 0.87 \text{ nb},$

$(|\epsilon_e^v|, |\epsilon_e^a|)_{\text{max}}^{e^+e^- \rightarrow X \rightarrow e^+e^-} \simeq (1.8, 0.82) \times 10^{-3},$
 $(|\epsilon_e^v|, |\epsilon_e^a|)_{\text{min}}^{e^+e^- \rightarrow X \rightarrow e^+e^-} \simeq (2.4, 6.3) \times 10^{-4},$

Cross sections depends on $|\sqrt{s} - m_X|$ strongly.

$\sqrt{s} \text{ (MeV)}$	$\sigma_{\text{min}} \text{ (nb)}$	$\sigma_{\text{max}} \text{ (nb)}$
17.02	0.011	0.87
17.05	0.0017	0.14
17.10	0.00043	0.035
17.20	0.00011	0.0086





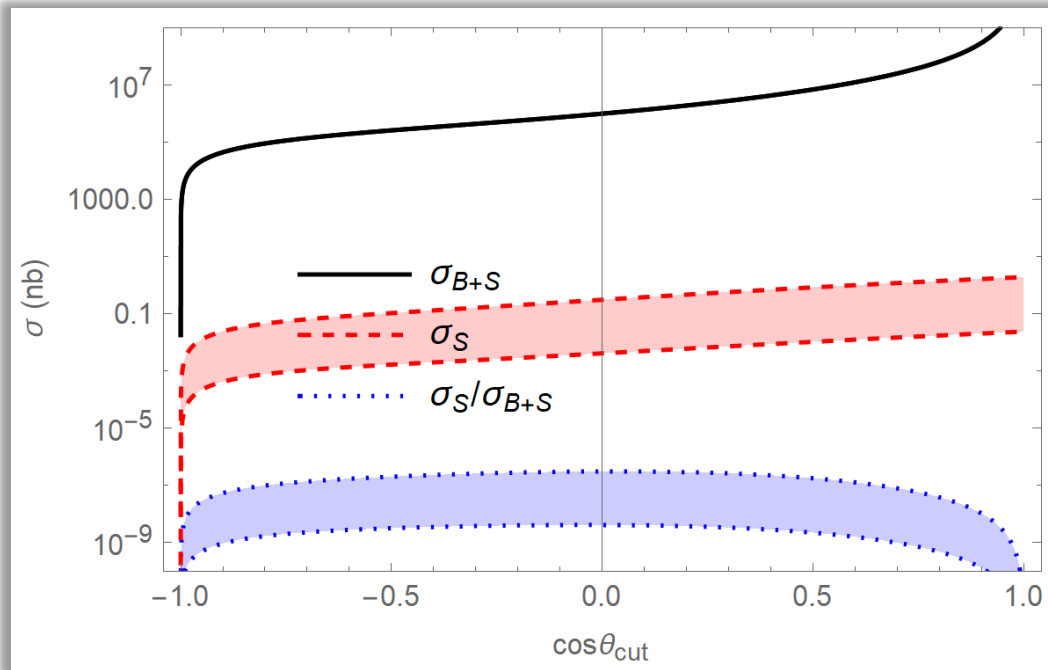
3 X(17) @ PADME

18

■ $e^+e^- \rightarrow X(17)/\gamma^* \rightarrow e^+e^-$

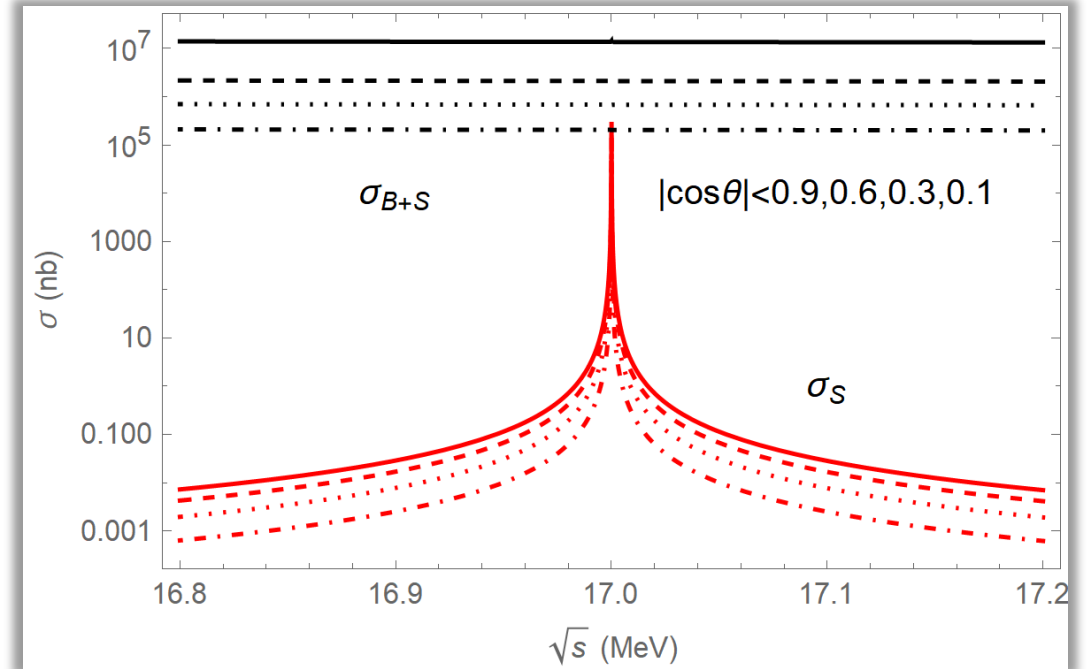
$m_X = 17 \text{ MeV}$, $\sqrt{s} = 17.02 \text{ MeV}$, $\Gamma_X = 0.089$

The range of integration is $\cos\theta \in [-1, \cos\theta_{\text{cut}}]$



Proper cuts on $|\cos\theta|$ improve the signal-to-noise ratio

Cross section $\sigma(e^+e^- \rightarrow X(17))$: 89 ~ 777 nb



PADME: Scan energy steps 1.5 MeV; beam energy spread 0.75 MeV; ECal hit reconstruction threshold 1 MeV

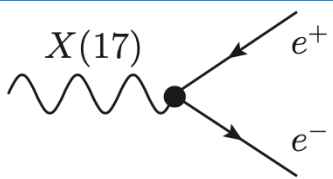
Due to the extremely low signal-to-noise ratio, I am pessimistic about detecting X(17) in the PADME experiment (Bhabha scattering).

CONTENTS

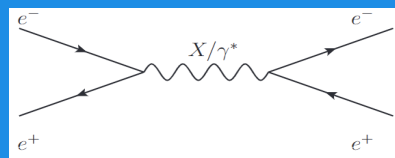
MOTIVATION



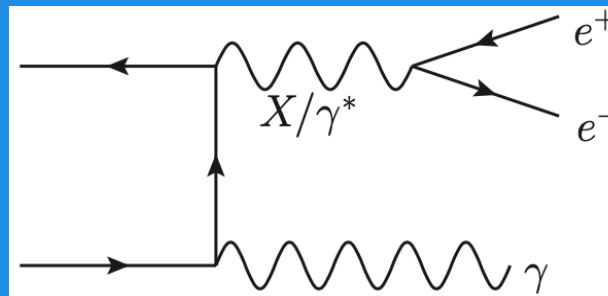
Motivation



Constraints on
 Xe^+e^- coupling



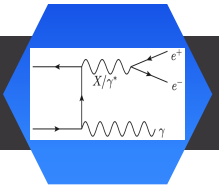
$X(17)$
@ PADME



Revisit $X(17)$ @ BESIII

SUMMARY

Summary

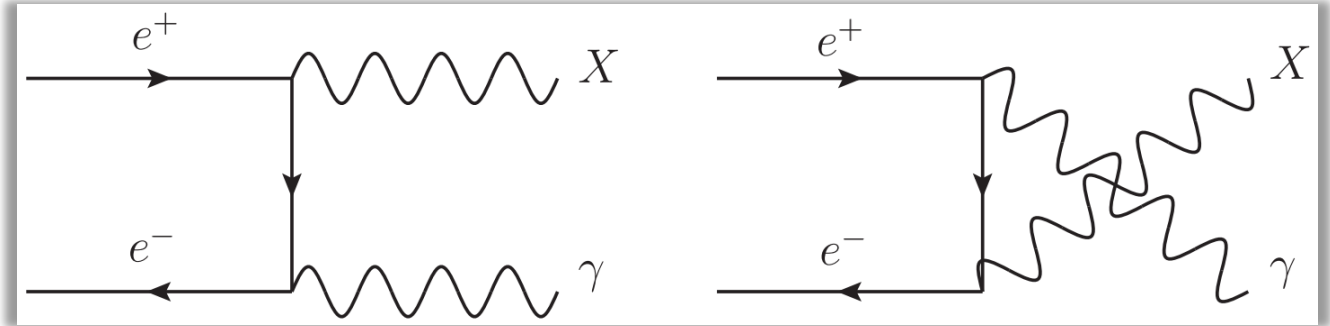


4 Revisit X(17) @ BESIII

20

■ $e^+e^- \rightarrow X(17)+\gamma$

Within the errors of $\mathcal{O}\left(\frac{m_e^2}{m_X^2}\right) \approx 0.1\%$, the differential cross section depends on the coupling parameters only through the argument $(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2$

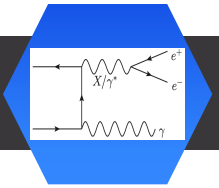


$$\frac{d\sigma}{d\cos\theta} = \frac{2\pi\alpha^2}{s} \left((\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \right) \left(\frac{1 + \cos^2\theta}{1 - \cos^2\theta} \right) + \mathcal{O}\left(\left(\frac{m_e}{m_X}\right)^2, \left(\frac{m_X}{\sqrt{s}}\right)^2, \left(\frac{m_e}{\sqrt{s}}\right)^2\right).$$

The integrated cross section:

$$\sigma = \frac{4\pi\alpha^2}{s} \left((\varepsilon_e^v)^2 + (\varepsilon_e^a)^2 \right) \left(\text{Log}\left[\frac{s}{m_e^2}\right] - 1 \right) + \mathcal{O}\left(\left(\frac{m_e}{m_X}\right)^2, \left(\frac{m_X}{\sqrt{s}}\right)^2, \left(\frac{m_e}{\sqrt{s}}\right)^2\right).$$

The leading order results are the argument $(\varepsilon_e^v)^2 + (\varepsilon_e^a)^2$ times the differential and total cross sections of $e^+e^- \rightarrow \gamma\gamma$



4 Revisit X(17) @ BESIII

21

■ $e^+e^- \rightarrow X(17)+\gamma$

At $\sqrt{s} = 3.7$ GeV, the integrated cross section:

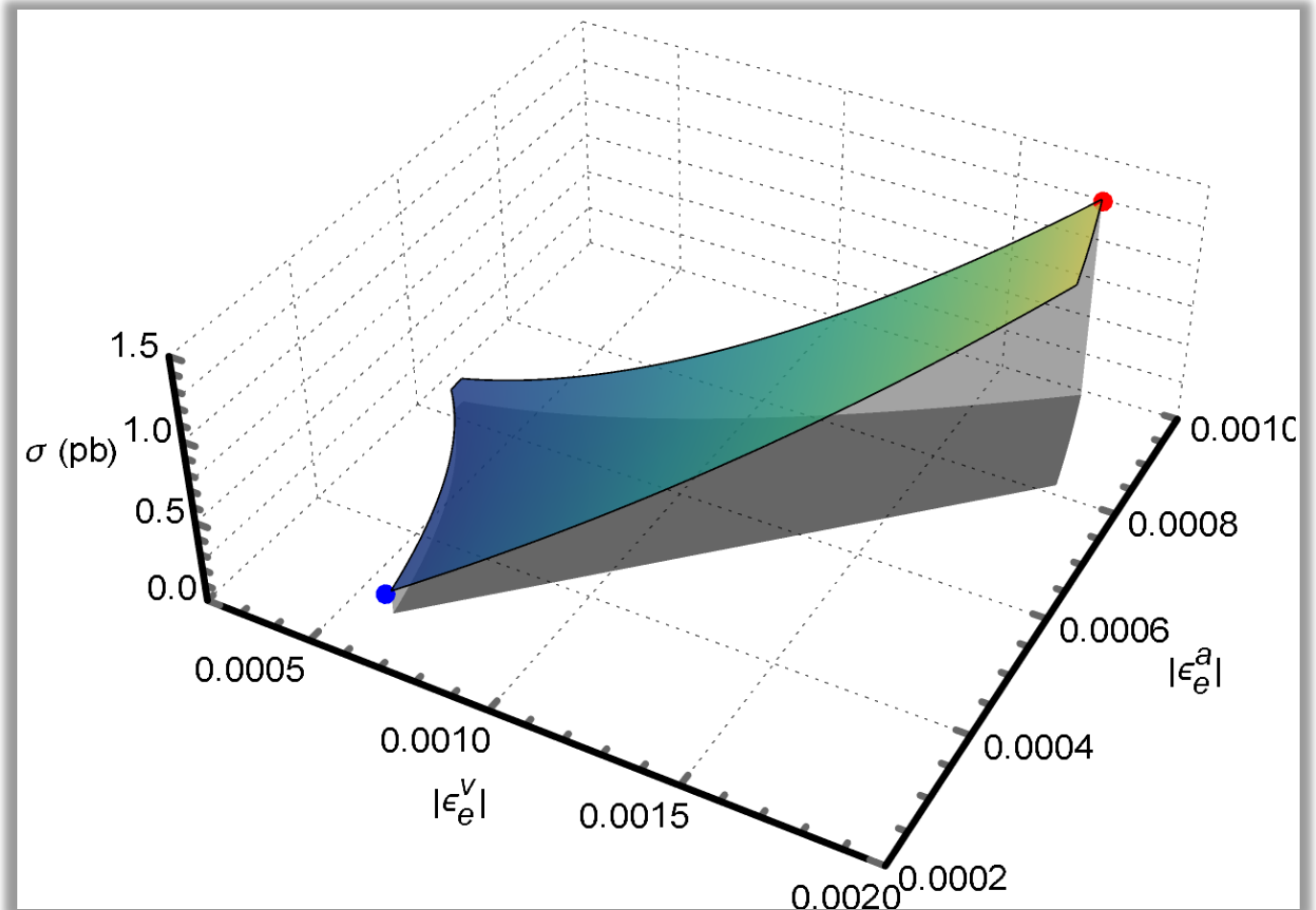
$$0.144 \text{ pb} \lesssim \sigma(e^+e^- \rightarrow X\gamma) \lesssim 1.28 \text{ pb},$$

with maximum and minimum coupling points

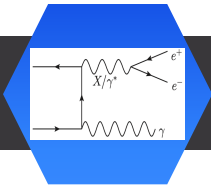
$$\begin{aligned} (|\varepsilon_e^v|, |\varepsilon_e^a|)_{\max}^{e^+e^- \rightarrow X\gamma} &\simeq (1.7, 0.99) \times 10^{-3}, \\ (|\varepsilon_e^v|, |\varepsilon_e^a|)_{\min}^{e^+e^- \rightarrow X\gamma} &\simeq (6.1, 2.7) \times 10^{-4}. \end{aligned}$$

Given the luminosity of BESIII at $\sqrt{s} = 3.7$ GeV
 $\mathcal{L} = 10^{33} \text{ cm}^{-2} \text{ s}^{-1} \approx 10^4 \text{ pb}^{-1} \text{ year}^{-1}$ and
 BESIII cover 93% solid angle,
 we roughly estimate X(17) events per year:

$$(1.3 \sim 12) \times 10^3$$



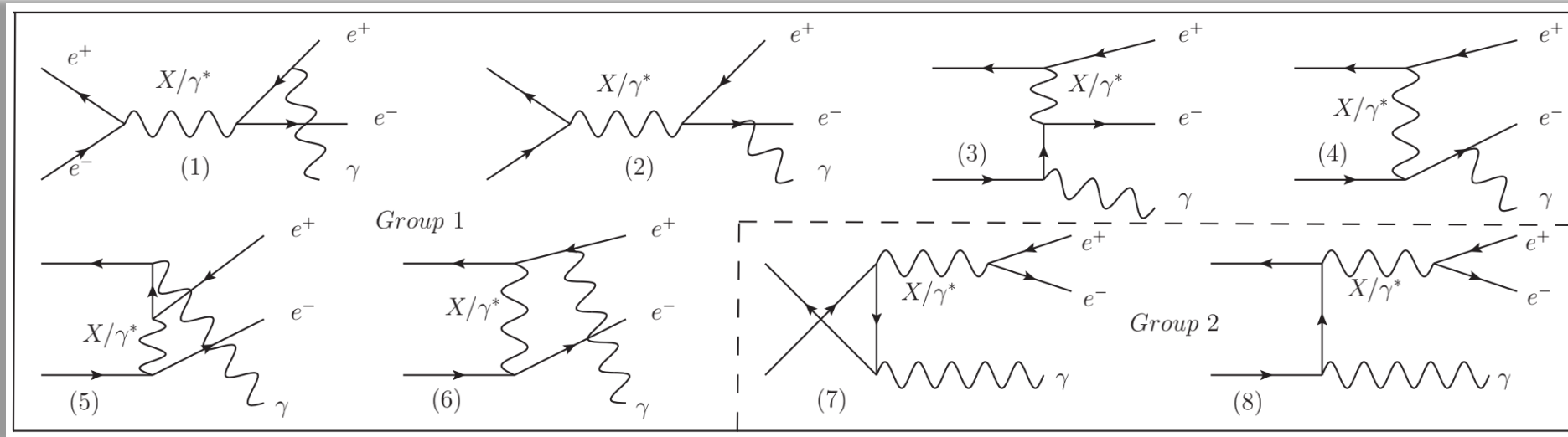
In e^+e^- c.m.s., the lifetime of X(17) at $\sqrt{s} = 3.7$ GeV is boosted by about two orders to 0.44~3.9 ps.
 Decay length of X(17): $L \approx 0.13 \sim 1.2$ mm, which implies X(17) created at the primary vertex in chamber.



4 Revisit X(17) @ BESIII

22

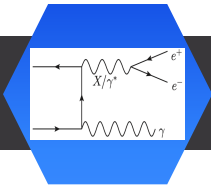
■ Identify X in $e^+e^- \rightarrow e^+e^-g$



The differential cross section of $e^+(p_a)e^-(p_b) \rightarrow e^+(p_1)e^-(p_2)\gamma(p_3)$ can be formulated as

$$d\sigma = \frac{(2\pi)^4}{4\sqrt{(p_a \cdot p_b)^2 - m_e^4}} \sum_{\text{spins}} |\mathcal{M}_{\text{sg}} + \mathcal{M}_{\text{bg}}|^2 d\Phi_3, \quad (39)$$

where $\sum$ stands for the average on initial spins and sum over final spins, $\mathcal{M}_{\text{sg}}$ and $\mathcal{M}_{\text{bg}}$ are the amplitudes of the signal and background Feynman diagrams respectively, and $d\Phi_3$ is the three-particle phase space.

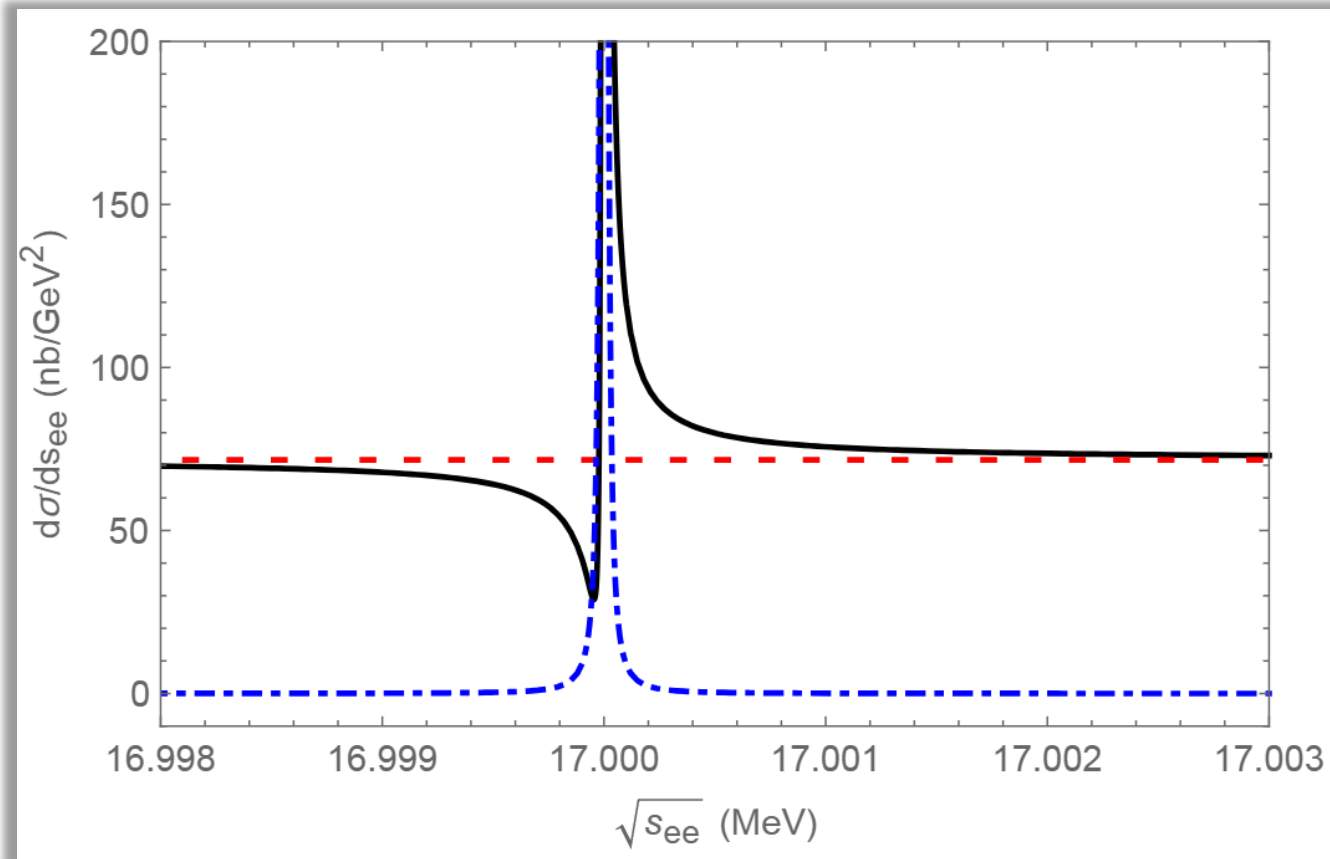


4 Revisit X(17) @ BESIII

23

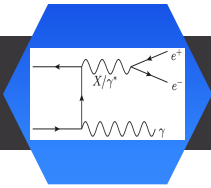
- Identify X in $e^+e^- \rightarrow e^+e^- g$

Differential distribution of e^+e^- invariant mass spectrum



$$(|\varepsilon_e^v|, |\varepsilon_e^a|)_{\max}^{e^+e^- \rightarrow X\gamma} \simeq (1.7, 0.99) \times 10^{-3},$$

$|\cos\alpha| < 0.93$ for electrons/positrons,
 $|\cos\alpha| < 0.8$ and $E_\gamma > 25$ MeV for photons,



4 Revisit X(17) @ BESIII

24

■ Identify X in $e^+e^- \rightarrow e^+e^-g$

Energy resolution at 20 MeV: $\delta(E)/E = 9\% \Rightarrow \delta(17\text{MeV}) \approx 1.5\text{MeV}$

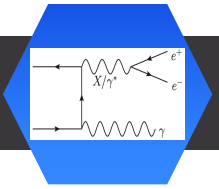
Nucl. Instrum. Meth. A 614 (2010), 345-399

Energy resolution (MeV)	1	1.5	3
Mass ranges (MeV)	[14, 20]	[12.5, 21.5]	[8, 26]
$B (\times 10^4 \text{ year}^{-1})$	14.8	22.5	48.8
$S (\times 10^4 \text{ year}^{-1})$	1.2 (0.13)		
$S/\sqrt{S+B}$	30 (3.4)	25 (2.7)	17 (1.9)

For the Gaussian distribution, 3σ deviation away from the expectation account for 99.73% of the set.

Values outside (inside) the parentheses are for the maximum (minimum) events in $e^+e^- \rightarrow Xg$.

The number of events and signal-to-noise ratio data indicate that it is feasible to detect/exclude X(17) in the BESIII experiment.

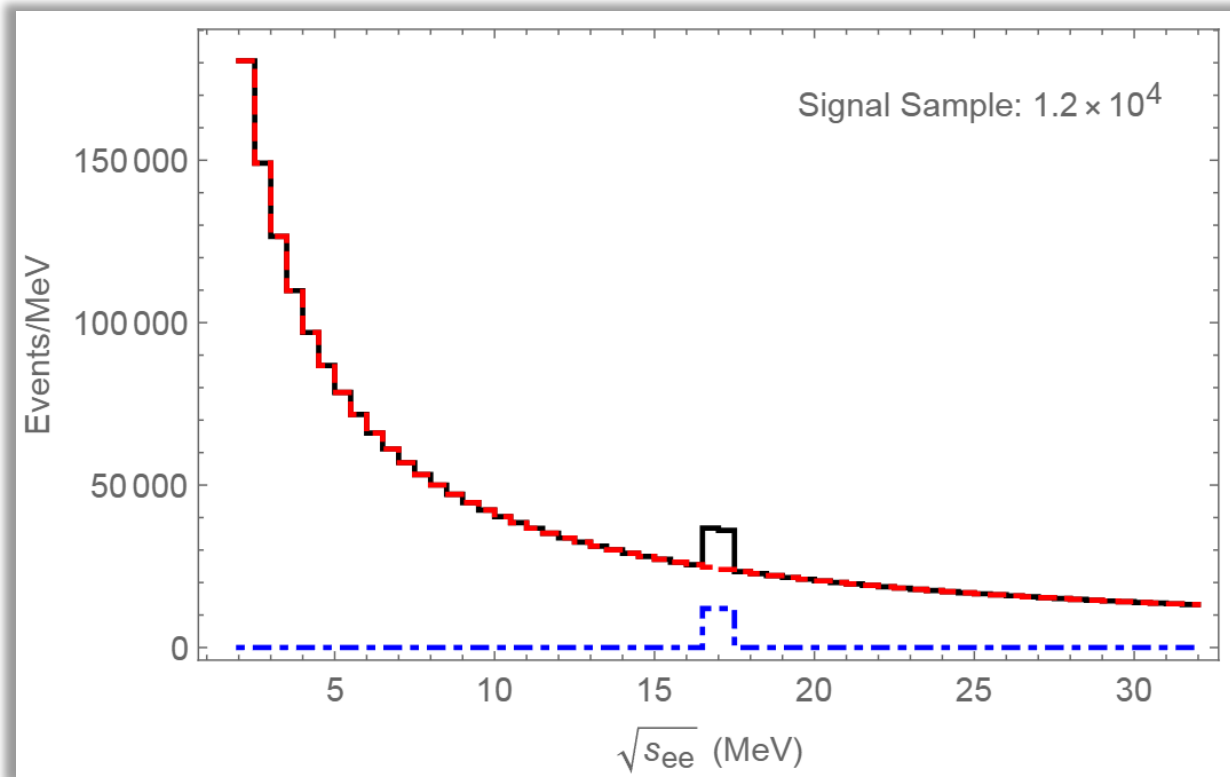


4 Revisit X(17) @ BESIII

25

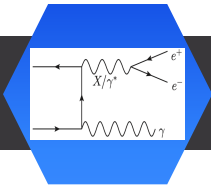
■ Identify X in $e^+e^- \rightarrow e^+e^-g$

Invariant mass spectrum distribution **before smearing**



- The bin size: 0.5 MeV
- Maximum Sample: 1.2×10^4

If the minimum sample (0.13×10^4) was adopted, the height of X(17) bump will be about 10% of that and become a shoulder unless we adopt smaller bin size.

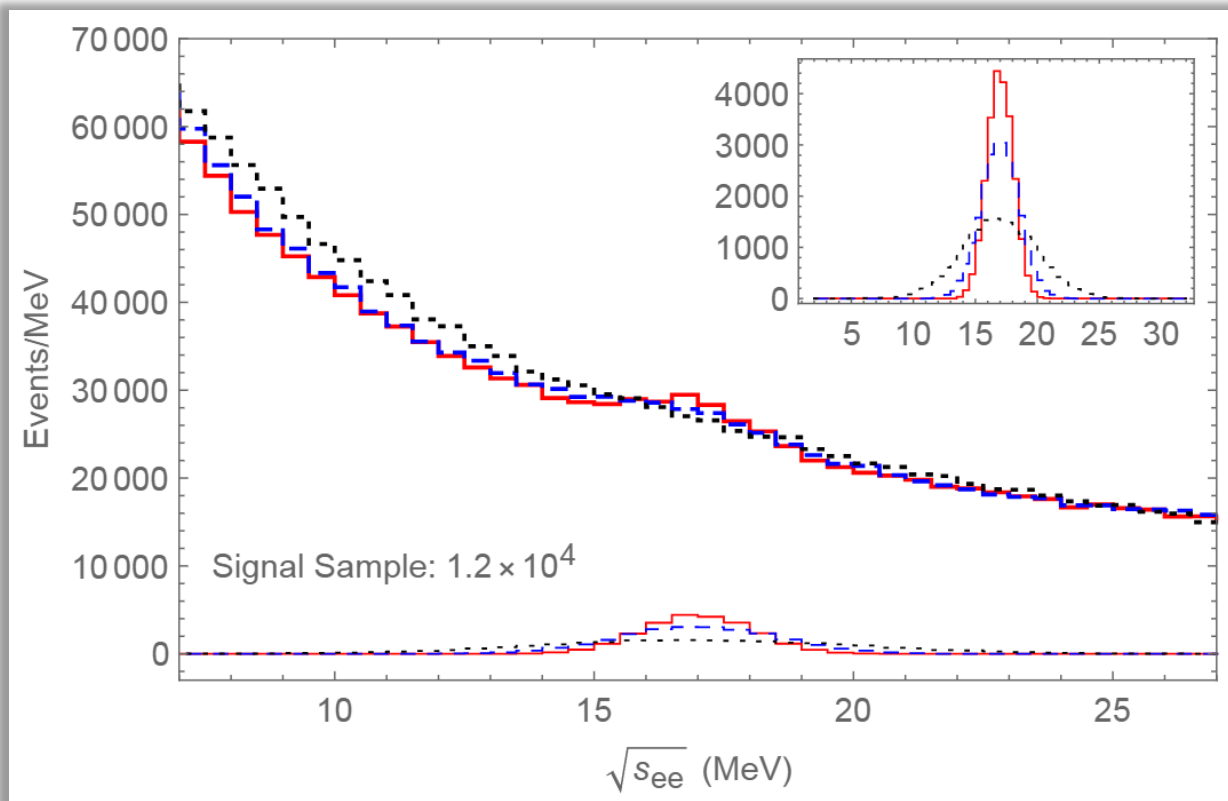


4 Revisit X(17) @ BESIII

26

■ Identify X in $e^+e^- \rightarrow e^+e^-g$

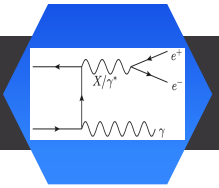
Invariant mass spectrum distribution **After smearing**



- The bin size: 0.5 MeV
- Maximum Sample: 1.2×10^4
- Energy resolutions: 1, 1.5, 3 MeV for different legends

93% coverage of solid angle is considered, but efficiencies etc. are not taking into account.

Modifying the bin size will not make the bump sharper after smearing.

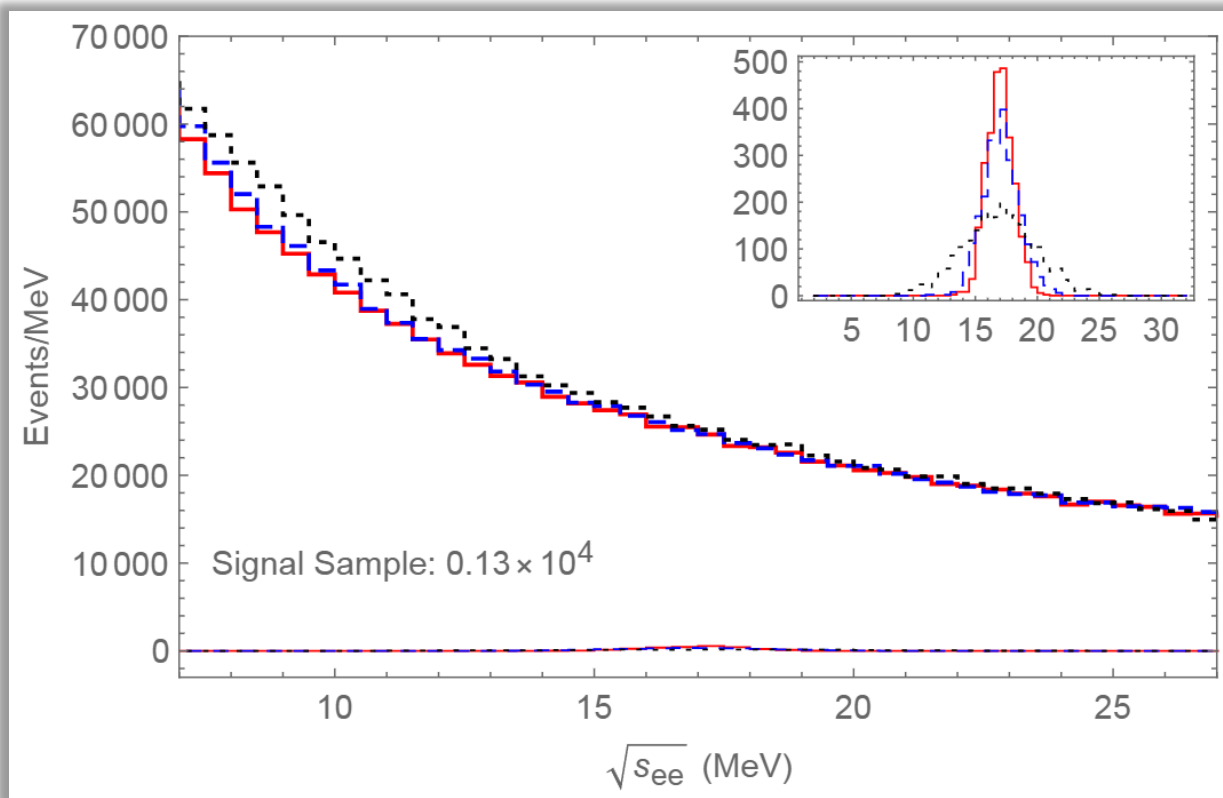


4 Revisit X(17) @ BESIII

27

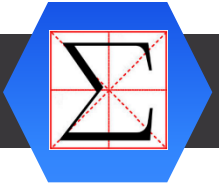
■ Identify X in $e^+e^- \rightarrow e^+e^-g$

Invariant mass spectrum distribution **After smearing**



- The bin size: 0.5 MeV
- Maximum Sample: 0.13×10^4
- Energy resolutions: 1, 1.5, 3 MeV for different legends

Height of the bump decreases by about 90% in comparison with the maximum sample, and we can not identify the X(17) signals over the background events.



- ◆ The Atomki anomalies can be explained by introducing the $X(17)$ boson with $V \pm A$ interaction
- ◆ Based on existing experimental data, **constraints on the couplings of $X(17)$ to electrons within this $V \pm A$ model** are obtained, while pure V and A couplings are disfavored.
- ✓ Pessimistic for $e^+e^- \rightarrow X(17) \rightarrow e^+e^-$ at PADME, but optimistic for $e^+e^- \rightarrow X(17)+\gamma \rightarrow e^+e^-\gamma$ at BESIII
- ✓ **BESIII experiment has the ability to give a decisive conclusion on $X(17)$**

Thank You!

Hunt for X(17)

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2025-10-20 @ 南京师范大学, 南京