

# Searching for the $T_{bc}$ molecular candidate via lineshape analysis

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- 3 Lineshapes in  $\bar{B}^0 D^0$  and  $\bar{B}_s^0 D^0$  invariant mass distributions
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# Motivation: predictions on the $T_{bc}$ with $J^P = 0^+$

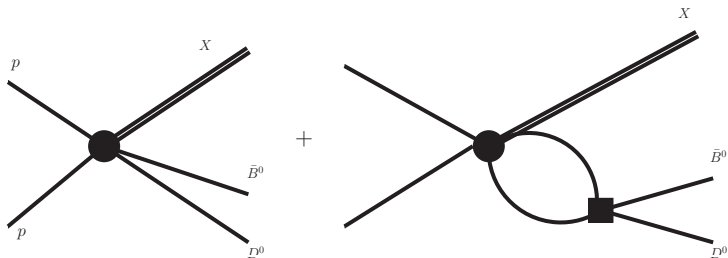
## Deep bound state

- 11 MeV below  $\bar{B}^0 D^0$  threshold [M. Karliner et al, PRL 119 \(2017\)](#)
- 23 MeV below  $\bar{B}^0 D^0$  threshold in QM [T. F. Caramés et al, PRD99 \(2019\)](#)
- $39 \left( \begin{smallmatrix} +4 \\ -6 \end{smallmatrix} \right) \left( \begin{smallmatrix} +8 \\ -18 \end{smallmatrix} \right)_{latt.}$  MeV below  $\bar{B}^0 D^0$  threshold [A. Radhakrishnan et al, PRD 110\(2024\)](#)
- 136 – 178.4 MeV below  $\bar{B}^0 D^0$  threshold [C.R. Deng et al EPJA 56 \(2020\)](#); [Y. Tan et al, EPJ.Plus 135 \(2020\)](#);

## Shallow bound state

- $B_2 = 0.42 - 31.35$  MeV in OBE, [Ning Li et al, PRD 88\(2013\)](#); [M. Z. Liu et al, PRD 110\(2024\)](#)
- $B_2 = 0.75 - 2.01$  MeV in HEFT, [Luciano M. Abreu, NPB 985 \(2022\)](#)
- $B_2 = 1.7 - 8.4$  MeV in V-exchange, [S. Sakai et al, PRD96\(2017\)](#)
- 4 MeV below  $\bar{B}^0 D^0$  threshold in QM, [G. Yang et al, PRD101 \(2020\)](#)
- A few MeV below  $\bar{B} D$  threshold, [C. Alexandrou et al, PRL132 \(2024\)](#)
- $B_2 = 0.8 - 3.6$  MeV in ChEFT, [Zhe Liu et al, PRD112\(2025\)](#)

# $\bar{B}D$ FSI in $pp$ collision



$$th_2 - th_1 = 4.5 \text{ MeV}$$

- $|\bar{B}D\rangle_{I=0} = -\frac{1}{\sqrt{2}}|\bar{B}^0 D^0\rangle + \frac{1}{\sqrt{2}}|B^- D^+\rangle,$   
 $|\bar{B}D\rangle_{I=1} = -\frac{1}{\sqrt{2}}|\bar{B}^0 D^0\rangle - \frac{1}{\sqrt{2}}|B^- D^+\rangle$
- $\bar{B}D$  FSI similar to the one of  $ND$  in  $\Lambda_b \rightarrow \pi^- p D^0$ . S. Sakai et al, PLB808 (2020)
- Threshold effect in  $\{\pi^0 \pi^0, \pi^+ \pi^-\}$  scattering, P. Budini et al, PRL6(1961); U.G.

Meissner, PLB406(1997); J. Gasser et al, NPB 850 (2011); X.H. Liu et al, EPJC73(2013) ...

## ERE in $\{\bar{B}^0 D^0, B^- D^+\}$ scattering

$$t_{ij} = \frac{2\pi}{\det} \begin{pmatrix} \frac{1}{\mu_1} \left( -\frac{1}{a_{22}} + ip_2 \right) & -\frac{1}{a_{12}\sqrt{\mu_1\mu_2}} \\ -\frac{1}{a_{12}\sqrt{\mu_1\mu_2}} & \frac{1}{\mu_2} \left( -\frac{1}{a_{11}} + ip_1 \right) \end{pmatrix}_{ij},$$

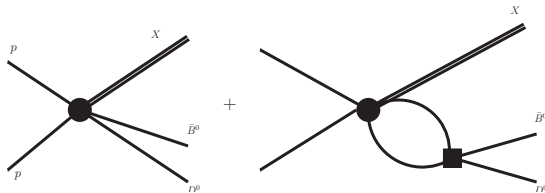
$$\det = \left( \frac{1}{a_{12}} \right)^2 - \left( \frac{1}{a_{11}} - ip_1 \right) \left( \frac{1}{a_{22}} - ip_2 \right),$$

$$\frac{1}{a_{22,\text{eff}}} = \frac{1}{a_{22}} - \frac{a_{11}}{a_{12}^2 (1 + a_{11}^2 p_1^2)} - i \frac{a_{11}^2 p_1}{a_{12}^2 (1 + a_{11}^2 p_1^2)},$$

with  $p_1 = \sqrt{2\mu_1(w - th_1)}$  and  $p_2 = \sqrt{2\mu_2(w - th_2)}$ .

- $\kappa = \sqrt{2\mu_2\delta} = O(\delta^{1/2})$  with  $\delta = m_{B^-} + M_{D^+} - m_{\bar{B}^0} - M_{D^0}$
- the isospin breaking difference  $a_{22} - a_{11}$  is of  $O(\delta)$  and neglected.

# Production amplitude



$$\mathcal{A} = V_P^\Lambda \left( 1 + G_{\bar{B}^0 D^0}^\Lambda t_{\bar{B}^0 D^0, \bar{B}^0 D^0} + G_{B^- D^+}^\Lambda t_{B^- D^+, \bar{B}^0 D^0} \right)$$

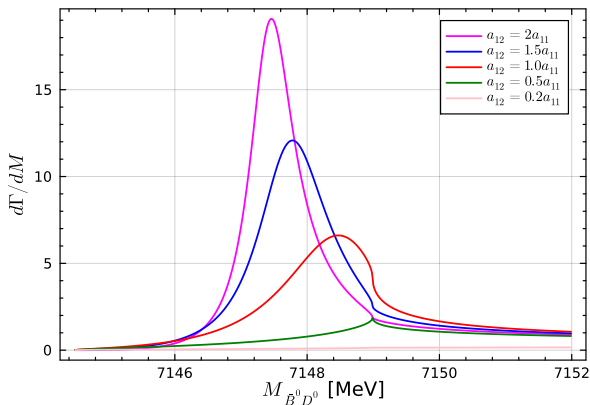
- $V_P^\Lambda$  is scale (or cut-off) dependent and  $G_i \propto \Lambda$
- $\mathcal{A}$  being cut-off independent leads to  $V_P^\Lambda \propto 1/\Lambda$ .

$$\mathcal{A} = V_P \left( t_{\bar{B}^0 D^0, \bar{B}^0 D^0} + t_{B^- D^+, \bar{B}^0 D^0} \right)$$

$$\frac{d\Gamma_{pp \rightarrow X \bar{B}^0 D^0}}{dM_{\bar{B}^0 D^0}} = \mathcal{N} p_X p_{D^0} \left| t_{\bar{B}^0 D^0, \bar{B}^0 D^0} + t_{B^- D^+, \bar{B}^0 D^0} \right|^2,$$

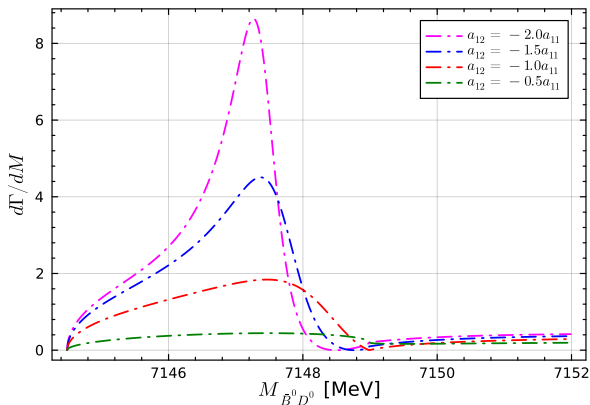
- $G_{\bar{B}^0 D^0}^\Lambda \simeq G_{B^- D^+}^\Lambda$ ;  $\mathcal{N}$  is an overall normalization constant,  $p_X$  ( $p_{D^0}$ ) is the momentum of the  $X$  ( $D^0$ ) in the  $pp$  ( $\bar{B}^0 D^0$ ) rest frame

# $\{\bar{B}^0 D^0, B^+ D^-\}_B$ coupled scattering in ERE



- the pink, green, red, blue, and magenta lines corresponding to  $a_{22,eff} = 0.10 + i0.17$  fm,  $0.18 + i1.42$  fm,  $-2.66 + i1.77$  fm,  $-2.91 + i0.69$  fm and  $-2.83 + i0.35$  fm respectively.

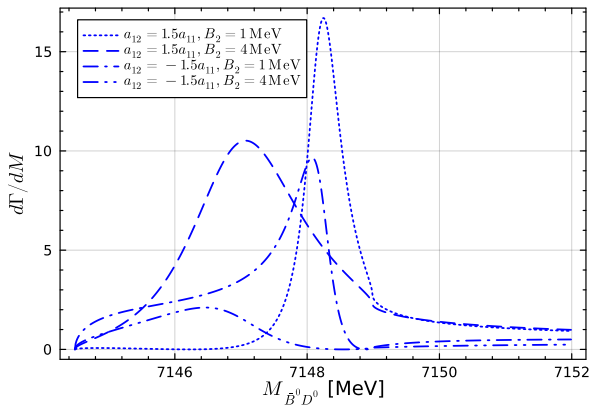
# $\{\bar{B}^0 D^0, B^+ D^-\}_B$ coupled scattering in ERE



- the green, red, blue, and magenta lines corresponding to  $a_{22,eff} = -0.18 + i1.42 \text{ fm}, -2.66 + i1.77 \text{ fm}, -2.91 + i0.69 \text{ fm}, -2.83 + i0.35 \text{ fm}$  respectively.

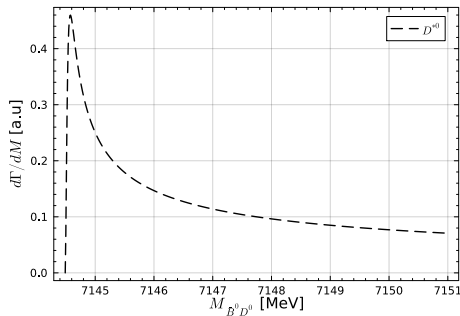
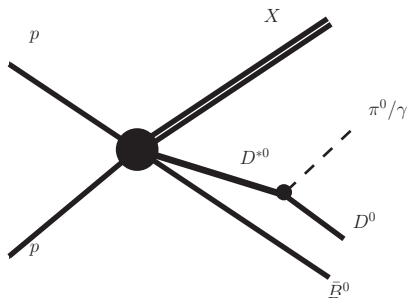


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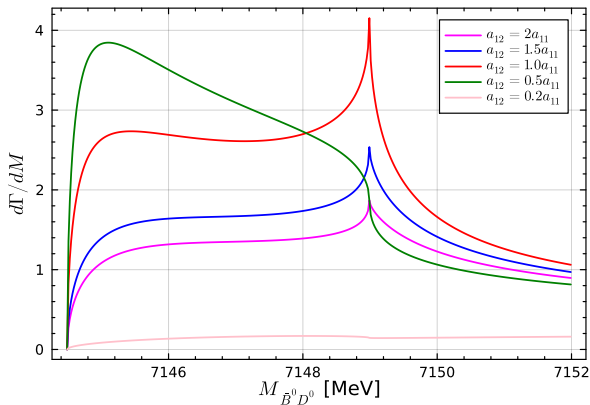
- the binding energy varies from 1 to 4 MeV, where  $a_{22,eff}$  is a function of  $a_{12}^2$  and subsequently evaluated as  $-3.95 + i0.74$  fm and  $-2.20 + i0.62$  fm corresponding to the binding energies of 1 and 4 MeV, respectively

# The feed-down effect from $D^*$ near $\bar{B}^0 D^0$ threshold



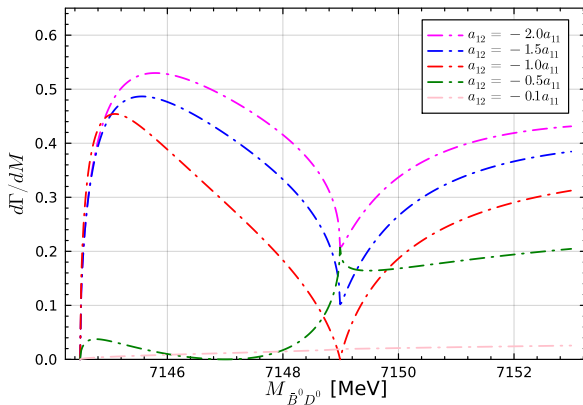
- The feed-down effect from  $D^*$  causes a narrow enhancement near  $\bar{B}^0 D^0$  threshold.
- The feed-down effect pollutes the enhancement from the pole in  $\bar{B} D$  scattering.

# $\{\bar{B}^0 D^0, B^+ D^-\}_V$ coupled scattering in ERE



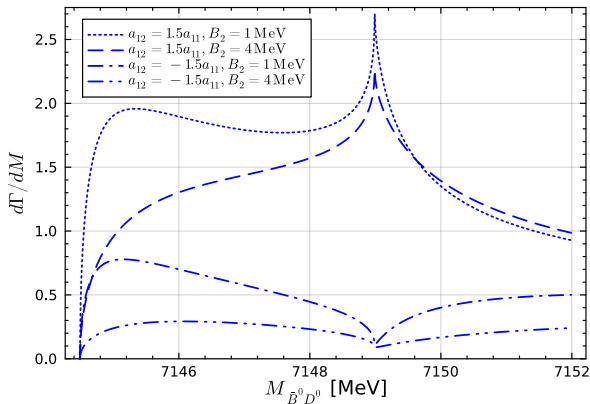
- $a_{22,eff}$  are  $-0.10 + i0.17$  fm,  $-0.18 + i1.42$  fm,  $2.66 + i1.77$  fm,  $2.91 + i0.69$  fm and  $2.83 + i0.35$  fm.
- Green:  $-1/a_{22} + ip_2 - 1/a_{12} \simeq 2(-1/a_{22} - \kappa) @ th_2$

# $\{\bar{B}^0 D^0, B^+ D^-\}_V$ coupled scattering in ERE



- $a_{22,eff}$  are  $-0.03 + i0.04$  fm,  $-0.18 + i1.42$  fm,  $2.66 + i1.77$  fm,  $2.91 + i0.69$  fm, and  $2.83 + i0.35$  fm.
- Green:  $-1/a_{22} + ip_2 - 1/a_{12} \simeq \mathcal{O}(1/a_{12}) @ th_2$

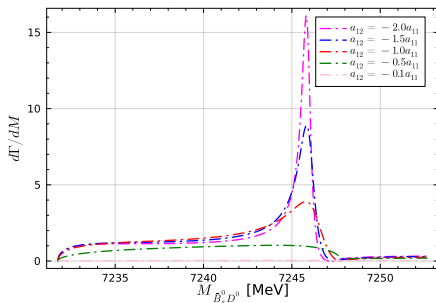
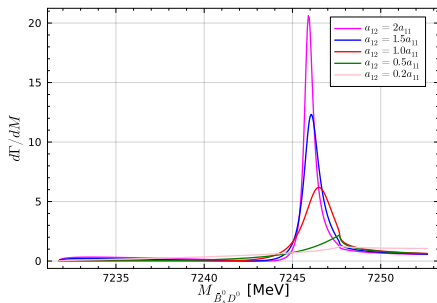
# $\{\bar{B}^0 D^0, B^+ D^-\}_V$ coupled scattering in ERE



- $a_{22,eff}$  are  $3.95 + i0.74 \text{ fm}$  and  $2.20 + i0.62 \text{ fm}$ .

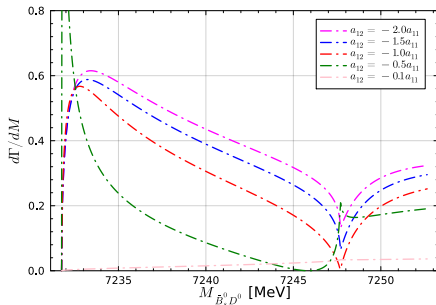
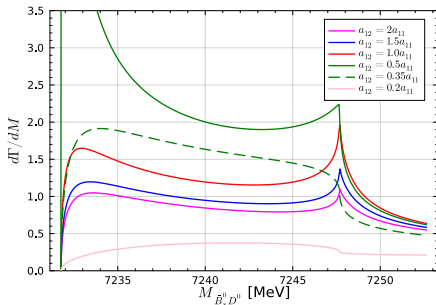
# $\{\bar{B}_s^0 D^0, B^- D_s^+\}_B$ coupled scattering in ERE

$$|B^- D_s^+\rangle = \frac{1}{\sqrt{2}}|\bar{3}\rangle + \frac{1}{\sqrt{2}}|6\rangle, |B_s^0 D^0\rangle = -\frac{1}{\sqrt{2}}|\bar{3}\rangle + \frac{1}{\sqrt{2}}|6\rangle$$



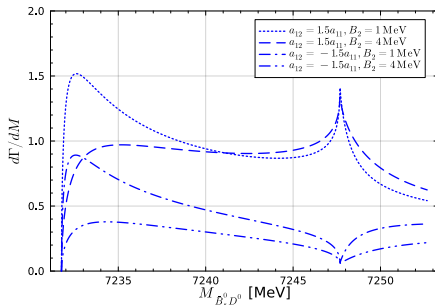
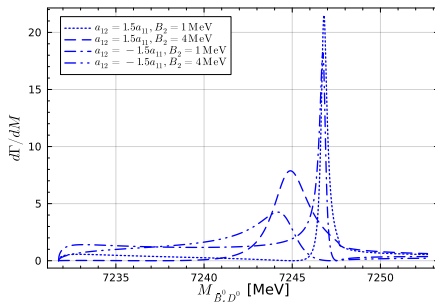
- $a_{22,eff} = 0.07 + i0.32 \text{ fm}, -0.78 + i1.76 \text{ fm}, -2.65 + i0.94 \text{ fm}, -2.73 + i0.40 \text{ fm}$  and  $-2.71 + i0.22 \text{ fm}$ , respectively

# $\{\bar{B}_s^0 D^0, B^- D_s^+\}_V$ coupled scattering in ERE



- $a_{22,eff} = 0.07 + i0.32 \text{ fm}, -0.78 + i1.76 \text{ fm}, -2.65 + i0.94 \text{ fm}, -2.73 + i0.40 \text{ fm}$  and  $-2.71 + i0.22 \text{ fm}$ , respectively

# $\{\bar{B}_s^0 D^0, B^- D_s^+\}_V$ coupled scattering in ERE



- $a_{22,eff} = -3.81 + i0.41$  fm, and  $-1.98 + i0.39$  (left);  
 $a_{22,eff} = 3.81 + i0.41$  fm and  $1.98 + i0.39$  fm (right)
- a near threshold enhancement in  $\bar{B}^0 D_s^+$  invariant mass distribution in  $\Upsilon \rightarrow D^- \bar{B}^0 D_s^+$  W.Y Liu et al; 2508.19007



# Summary

- $\{\bar{B}^0 D^0, B^- D^+\}$  coupled channel scattering is considered to search for the scalar  $T_{bc}$  via the lineshape in  $\bar{B}^0 D^0$  invariant mass distribution in pp collision.
  - In the isospin limit with  $a_{11} = a_{22}$ , the interference between  $t_{11}$  and  $t_{21}$  plays a crucial role in generating the lineshapes.
  - A new lineshape, differing from peak and cusp
  - The feed-down effect from  $D^*$  leads to a narrow enhancement around  $th_1$ .
- 
- Extending to  $\{pK^0, nK^+\}$  and  $\{pD^0, nD^+\}$  scattering.
  - $\{\Xi_c^+ K^-, \Xi_c^0 \bar{K}^0\}$   $\{D^0 \Xi^0, D^+ \Xi^-\}$  adopted to understand the  $\Omega_c^*$ .
  - To understand the shoulder around 4.26 GeV in  $e^+e^- \rightarrow J/\psi \pi^+ \pi^-$ .

# Summary

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  - To understand the shoulder around 4.26 GeV in  $e^+e^- \rightarrow J/\psi \pi^+ \pi^-$ .

*Thanks !*

$$|\bar{B}D\rangle_{I=0} = -\frac{1}{\sqrt{2}}|\bar{B}^0D^0\rangle + \frac{1}{\sqrt{2}}|B^-D^+\rangle,$$

$$|\bar{B}D\rangle_{I=1} = -\frac{1}{\sqrt{2}}|\bar{B}^0D^0\rangle - \frac{1}{\sqrt{2}}|B^-D^+\rangle,$$

$$|\bar{B}^0D^0\rangle = \frac{1}{\sqrt{2}}|6\rangle + \frac{1}{\sqrt{2}}|\bar{3}\rangle \text{ and } |B^-D^+\rangle = \frac{1}{\sqrt{2}}|6\rangle - \frac{1}{\sqrt{2}}|\bar{3}\rangle.$$

$$V_{11} = \langle \bar{B}^0D^0 | \mathcal{H} | \bar{B}^0D^0 \rangle = \frac{1}{2}V_6^{I=1} + \frac{1}{2}V_3^{I=0},$$

$$V_{12} = \langle \bar{B}^0D^0 | \mathcal{H} | B^-D^+ \rangle = \frac{1}{2}V_6^{I=1} - \frac{1}{2}V_3^{I=0},$$

$$V_{22} = \langle B^-D^+ | \mathcal{H} | B^-D^+ \rangle = \frac{1}{2}V_6^{I=1} + \frac{1}{2}V_3^{I=0},$$

$$G_i^\Lambda = \int^\Lambda \frac{d^3\mathbf{q}}{(2\pi)^3} \frac{2\mu_i}{\mathbf{q}^2 - p_i^2 - i\epsilon} = i\frac{\mu_i}{2\pi}p_i + \frac{\mu_i}{\pi^2}\Lambda \left[ 1 + O\left(\frac{p_i^2}{\Lambda^2}\right) \right].$$

$$\begin{aligned}
 t_{11}^{\text{th}} &= \frac{2\pi}{\mu_1} \left( -\frac{1}{a_{22}} - \kappa \right) \left[ \frac{1}{a_{12}^2} - \frac{1}{a_{11}} \left( \frac{1}{a_{22}} + \kappa \right) \right]^{-1} \\
 &\simeq \frac{2\pi}{\mu_1} \left[ \frac{1}{a_{22}} - \frac{1}{a_{12}^2 (a_{11}^{-1} + \kappa)} \right]^{-1}, \\
 t_{12}^{\text{th}} &= -\frac{2\pi}{\sqrt{\mu_1 \mu_2}} \frac{1}{a_{12}} \left[ \frac{1}{a_{12}^2} - \frac{1}{a_{11}} \left( \frac{1}{a_{22}} + \kappa \right) \right]^{-1} \\
 &\simeq -\frac{2\pi}{\sqrt{\mu_1 \mu_2}} \frac{a_{11}}{a_{12}} \left( \frac{a_{11}}{a_{12}^2} - \frac{1}{a_{11}} - \kappa \right)^{-1},
 \end{aligned}$$

where  $\kappa = \sqrt{2\mu_2\delta}$  with  $\delta = m_{B^-} + M_{D^+} - m_{\bar{B}^0} - M_{D^0}$ . One sees that  $\kappa = O(\delta^{1/2})$ .

