



# Neural network extraction of chromo-electric and chromo-magnetic gluon masses

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## Introduction

- Understanding the thermodynamics of QCD at finite temperature is a key goal in high-energy nuclear physics.
- lacktriangle Around the crossover transition (  $T_c \sim 150\text{--}170$  MeV), different degrees of freedom are liberated:

Quarks: dominate entropy and particle number. Gluons: dominate pressure and energy density.

 Disentangling chromo-electric and chromo-magnetic gluons is crucial for understanding the QCD medium.

## Introduction

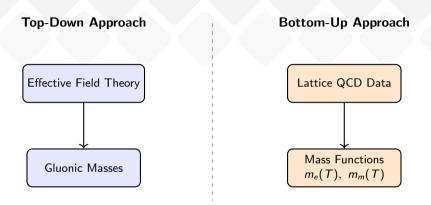
- At finite temperature, the Euclidean O(4) symmetry is reduced to spatial O(3): Longitudinal (electric) and transverse (magnetic) gluons become physically distinct.
- In the high- *T* limit, perturbative and effective field theory approaches suggest:

$$m_e \sim gT, \quad m_m \sim g^2T$$

 Magnetic screening masses are absent at any finite order in perturbation theory; they arise from non-perturbative dynamics and are captured by dimensionally reduced effective theories or lattice QCD.



## Introduction



**Essence:** Extracting gluon properties directly from thermodynamic data without specifying a Lagrangian — a data-driven inverse problem.

## Quasiparticle Model: Thermodynamics with Gluon Masses

- In the Quasiparticle model (QPM), gluons are treated as massive bosons with effective thermal masses  $m_e(T)$  and  $m_m(T)$ .
- The total thermodynamic potential is:

$$\ln Z(T) = \ln Z_e(T) + \ln Z_m(T)$$

Each follows Bose-Einstein statistics:

$$\ln Z_i(T) = -rac{d_i V}{2\pi^2} \int_0^\infty dp \ p^2 \ln \left(1 - e^{-rac{\sqrt{
ho^2 + m_i^2(T)}}{T}}
ight)$$

Key observables computed:

$$P(T), \quad \epsilon(T), \quad \Delta(T) = \frac{\epsilon - 3P}{T^4}$$

#### Degeneracy

$$d_e = 8$$
,  $d_m = 16$  (chromo-electric and -magnetic gluons)

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# Thermal Mass vs. Screening Mass

#### Thermal Mass

- Arises from medium interactions and enters the energy-momentum dispersion relation
- lacksquare Related to the pole of the gluon propagator:  $E^2=
  ho^2+m_{
  m thermal}^2$
- Governs real-time dynamics and transport properties

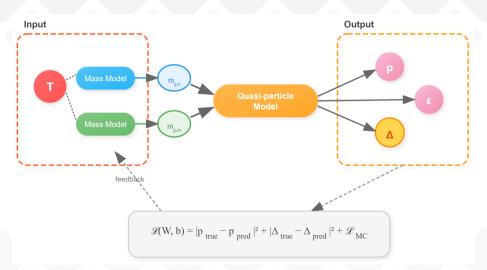
#### Screening Mass

- Characterizes exponential decay of spatial two-point correlations
- Defined from the asymptotic behavior:  $\langle A(x)A(0)\rangle \sim e^{-m_{\text{screening}}|x|}$
- Probed in lattice QCD via Polyakov loop correlators
- lacktriangle Does not necessarily coincide with thermal mass, especially near  $T_c$

These two mass definitions reflect different physical aspects of the medium



## Framework



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# Neural Network Structure and Physics-Informed Regularization

Each gluonic mass m<sub>i</sub>(T) is modeled by a separate ResNet:

ResNet structure helps alleviate the vanishing gradient problem in deep networks. Softplus output ensures m(T) > 0.

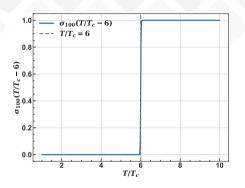
Loss function:

$$\mathcal{L} = \mathsf{MSE}(p/T^4, \Delta) + \lambda \cdot \omega(T) \left( rac{m_e(T)}{m_m(T)} - 2 
ight)^2$$

■ Regularization encodes the high-*T* theoretical prior:

$$rac{m_e}{m_m} 
ightarrow 2$$
 as  $T \gg T_c$ 

•  $\omega(T)$  is a sigmoid weight centered around  $T = 6T_c$ .

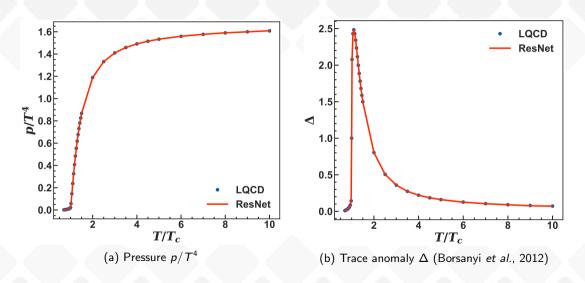


Regularization in High-T

ensures stable optimization and prevents the training instabilities often induced by hard, localized constraints.

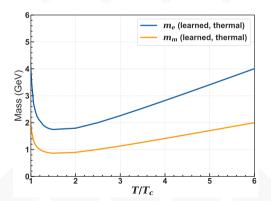
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## Results and Discussion



## Extracted Gluon Mass Functions

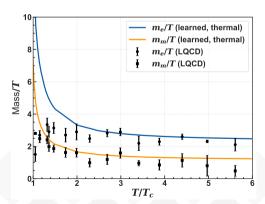
- $\bullet$   $m_e(T)$  and  $m_m(T)$  both show rapid drop near  $T_c$ , reflecting liberation of degrees of freedom.
- At high T, both masses increase linearly, consistent with perturbative behavior.



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# Comparison with Lattice Screening Masses

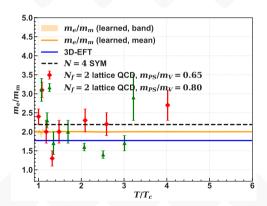
- Thermal masses from the quasiparticle model differ from screening masses (Nakamura; Saito; Sakai, 2004).
- Screening masses probe exponential fall-off of correlators; thermal masses affect real-time propagation.



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# High-T Behavior of $m_e/m_m$

■ The regularization term effectively guides the network toward correct asymptotic behavior (Maezawa *et al.*, 2010).



# Shear Viscosity from the Quasiparticle Model

• Using the extracted thermal masses, we compute the shear viscosity  $\eta$  via the relaxation time approximation (RTA):

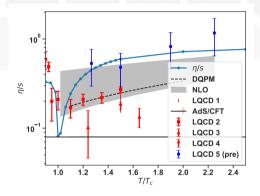
$$\eta = \frac{1}{15T} \sum_{i=e,m} d_i \int \frac{d^3p}{(2\pi)^3} \frac{p^4}{E_i^2} \tau_i f_i (1+f_i)$$

■ The relaxation time is given by:

$$au_i^{-1} = rac{N_c}{4\pi} g^2(T) T \ln \left(rac{2k}{g^2(T)}
ight)$$

with  $g^2(T)$  extracted from  $m_e^2=rac{N_c}{6}g^2T^2$ 

■ Entropy density s is obtained from lattice data via  $s = (\epsilon + p)/T$ 



DQPM: PRC 88, 045204 (2013). NLO: JHEP 03, 179 (2018). LQCD1: PRD 108, 014503 (2023). LQCD2: JHEP 04, 101 (2017). LQCD3: PRD 76, 101701 (2007). LQCD4: PRD 98, 014512 (2018). LQCD5: Zhang, et. al. Spigy Glupn 2025

## Summary

- Developed a neural-network-based framework to extract temperature-dependent chromo-electric and chromo-magnetic gluon masses in pure SU(3) gauge theory.
- Used a quasiparticle model with two separate gluonic mass functions:

$$m_e(T), \quad m_m(T)$$

trained to reproduce lattice pressure and trace anomaly data.

■ Introduced a soft regularization term to guide the model towards the high- *T* theoretical limit:

$$\frac{m_e}{m_m} o 2$$

- lacktriangle Successfully captured the non-perturbative features of gluon mass behavior near  $T_c$ , while reproducing perturbative scaling at high temperatures.
- Extracted thermal masses differ significantly from screening masses, highlighting the distinction between thermal and spatial correlation scales.
- Calculated shear viscosity  $\eta/s$  shows a minimum near  $T_c$  approaching the KSS bound, consistent with lattice QCD results.

# Remarks on High-Temperature Regularization

#### Current Constraint from 3D-EFT and $\mathcal{N}=$ 4 SYM

The regularization enforces  $m_e/m_m \rightarrow 2$  based on:

- 3D effective field theory calculations
- $ightharpoonup \mathcal{N} = 4$  Super Yang-Mills theory results

#### Subtle Issues Worth Noting

- Both 3D-EFT and SYM predictions refer to **screening mass ratios**, while our work extracts **thermal masses**—these need not share identical high- *T* behaviors
- lacktriangle The current temperature range  $T/T_c \in [1,10]$  remains pre-asymptotic

More discussions and comments are welcomed to refine the high-temperature constraints



#### References

BORSANYI, S.; ENDRODI, G.; FODOR, Z.; KATZ, S. D.; SZABO, K. K. Precision SU(3) lattice thermodynamics for a large temperature range. **JHEP**, v. 07, p. 056, 2012.

MAEZAWA, Y.; AOKI, S.; EJIRI, S.; HATSUDA, T.; ISHII, N.; KANAYA, K.; UKITA, N.; UMEDA, T. Electric and Magnetic Screening Masses at Finite Temperature from Generalized Polyakov-Line Correlations in Two-flavor Lattice QCD. **Phys. Rev. D**, v. 81, p. 091501, 2010.

NAKAMURA, A.; SAITO, T.; SAKAI, S. Lattice calculation of gluon screening masses. **Phys. Rev. D**, v. 69, p. 014506, 2004.