

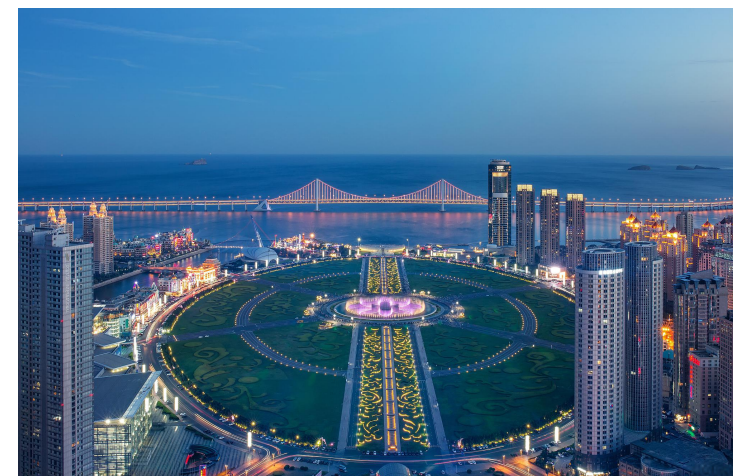
Finite-volume analysis and universal scaling near chiral phase transition in (2+1)-flavor QCD

Sabarnya Mitra

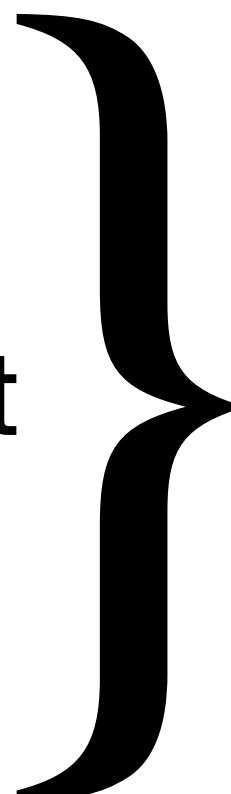
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The QCD phase diagram : From theory to experimental signatures

8-11 October , Dalian workshop 2025



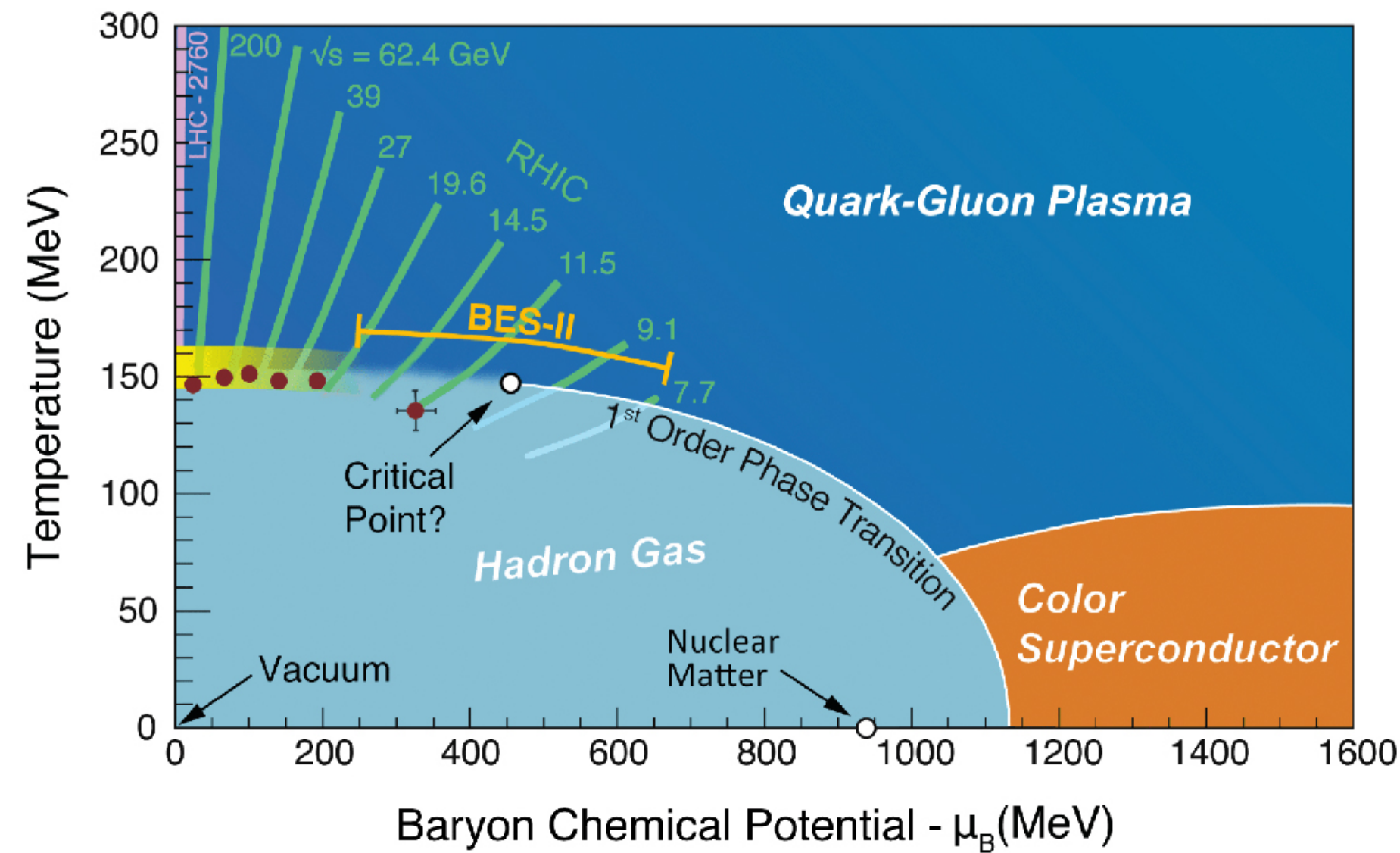
THANKS to my collaborators

- Frithjof Karsch
 - Christian Schmidt
 - Jishnu Goswami
- 
- Bielefeld University, Germany
- Mugdha Sarkar (NTU , Taiwan)
 - Sipaz Sharma (TuM Munich, Germany)

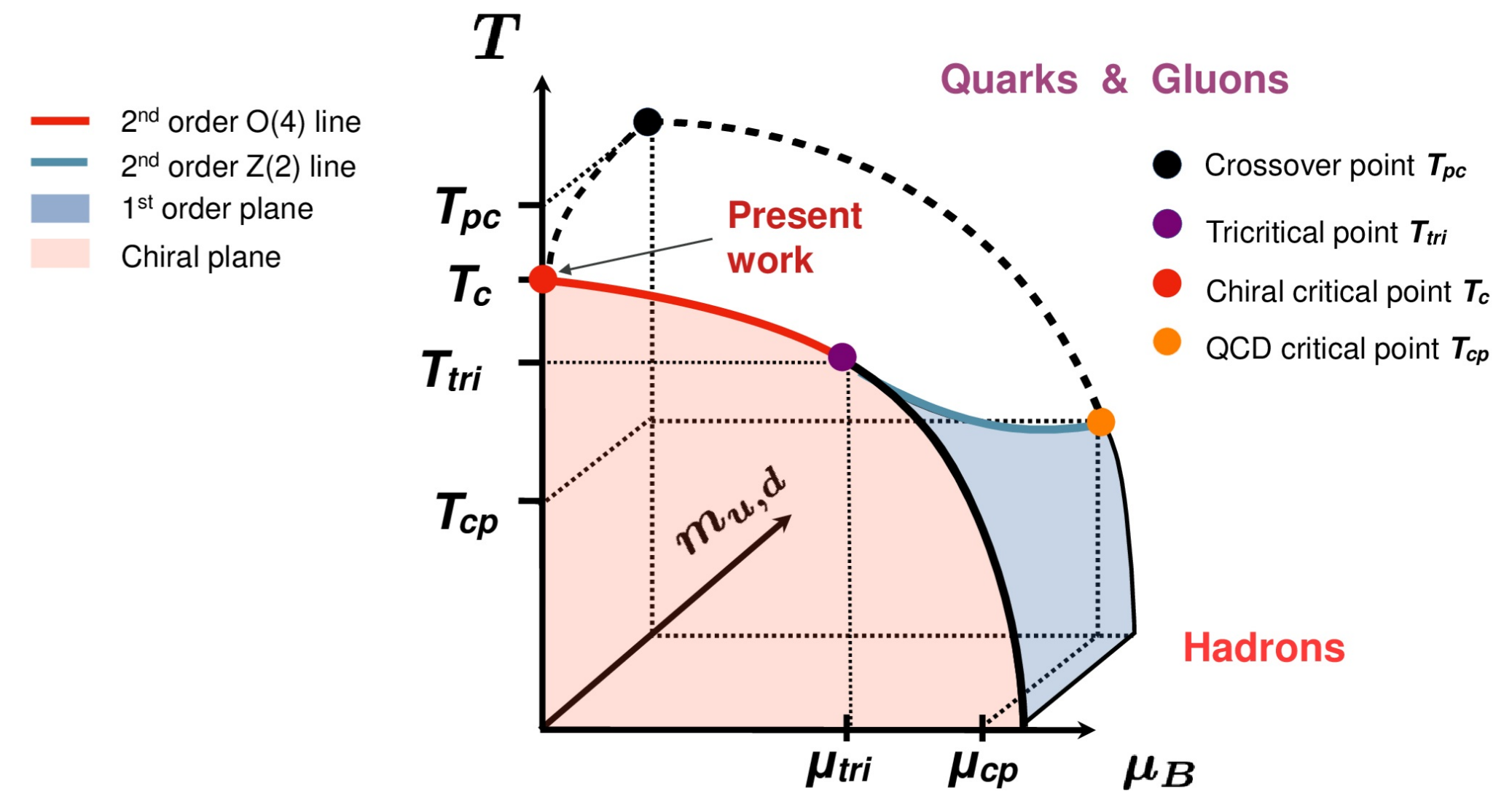
Contents

- Introduction and Motivation
- Improved order Parameter and scaling function
- Uniqueness and universality class
- Volume dependence and finite-volume effects
- Lattice setup and Results
- Conclusion and Outlook

Motivation and Introduction



Credit : CERN courier

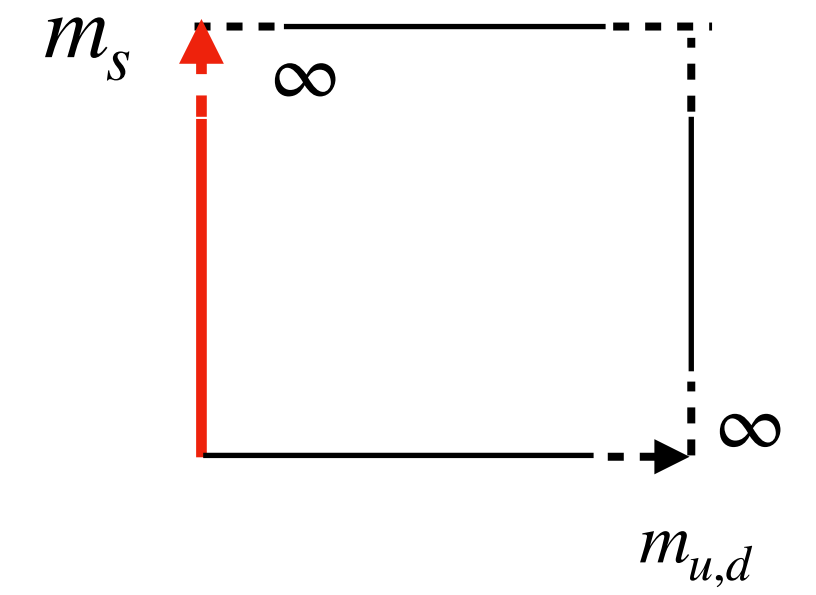


Numerous interesting talks in this workshop

- For **narrowing** onto the **QCD critical point**, one needs knowledge of **chiral phase transition** (cPT) .
- It is also important for constructing early universe QCD, explain heavy-ion experiments etc.
- This provides an **upper bound** on the QCD critical point temperature T_{cp} .

So, something brief about cPT

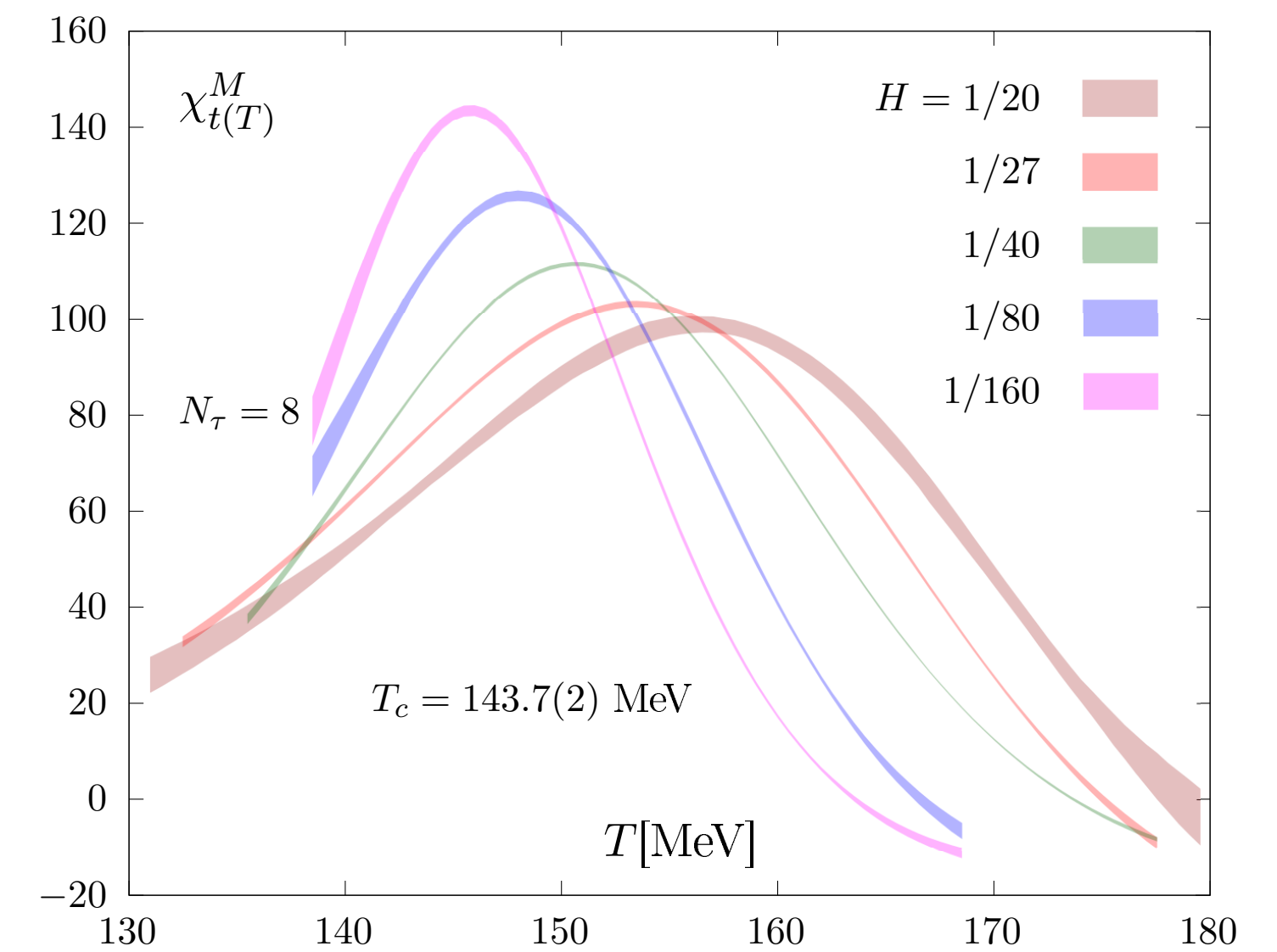
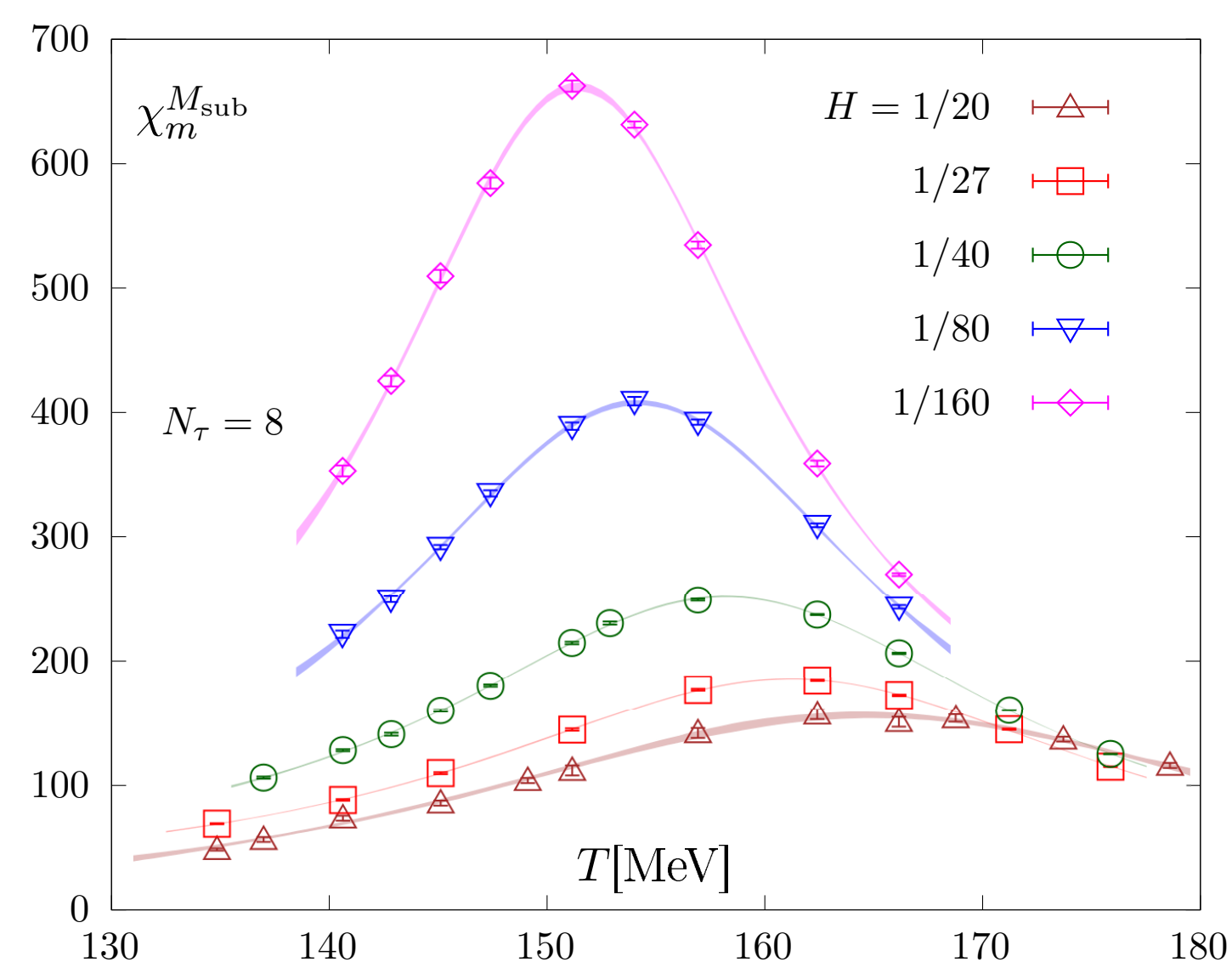
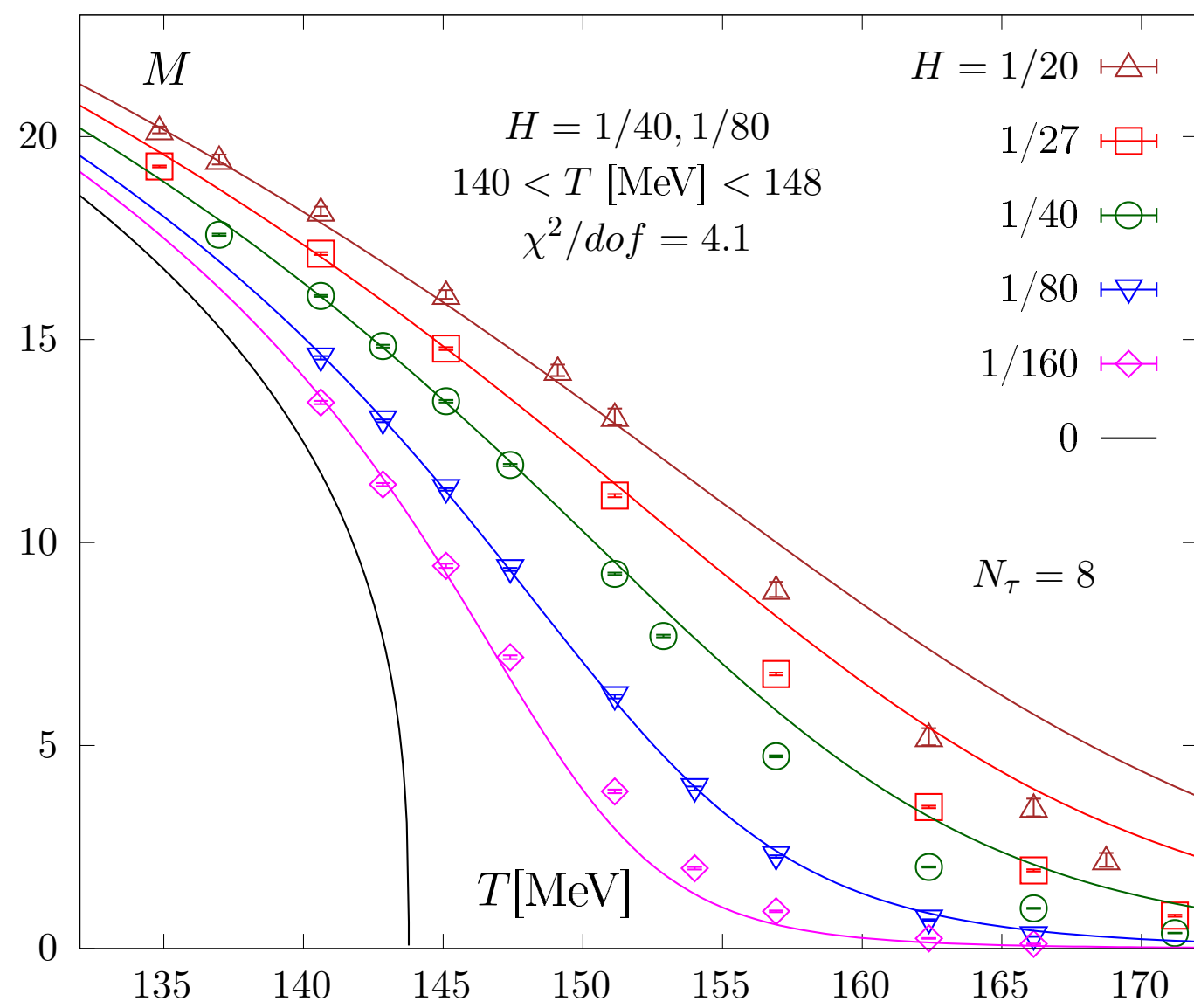
Chiral Phase Transition (cPT) : Brief



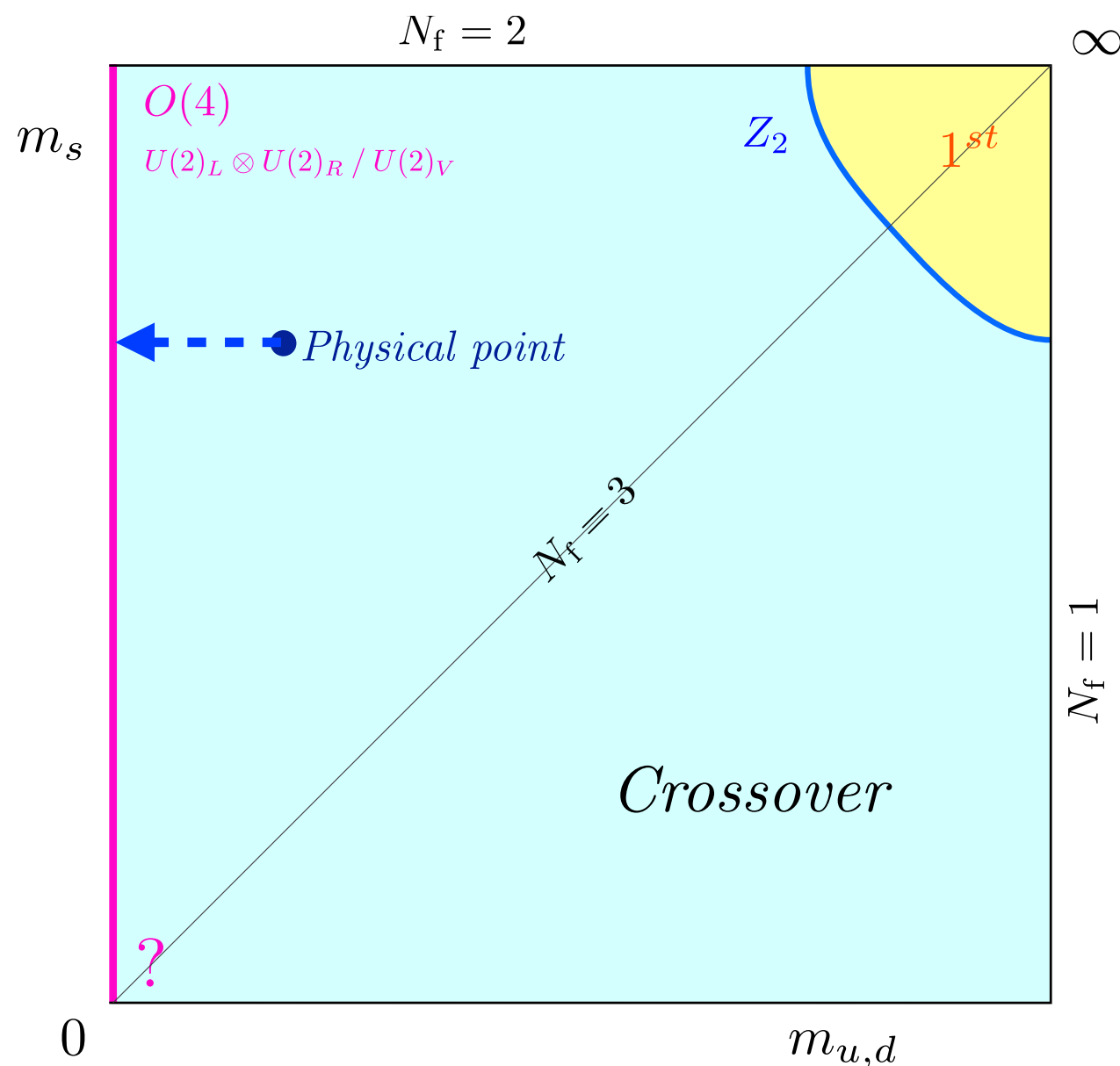
The “phase transition” between the high T - quark-gluon plasma (QGP) and the low T - hadronic phase

- In principle, this happens in the **exact chiral limit** ($m_u = m_d = m_{u,d} = 0$)
- This is a **second order** phase transition, with spontaneous symmetry breaking of the chiral symmetry group $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$
- The **order parameter** $M_{u,d} \sim \partial_{m_{u,d}} f$ remains **continuous, but non-differentiable** across $T = T_c$, for which there are **divergences in the chiral susceptibility** $\chi_{u,d} \sim \partial_{m_{u,d}}^2 f$
- The universality class (UC) of this phase transition determines the **fate of axial $U(1)_A$ symmetry** (restored / broken) **at $T = T_c$** .

The present state-of-the-art status ?



Ding et. al. , PRD 109 (2024) 114516



Cuteri et. al. , JHEP 11 (2021) 141

Ding, Cuteri

Lattice QCD (LQCD) $\xrightarrow[M, \chi]{\text{behaviour}}$ 2^{nd} order nature of cPT (increasingly)

- However , **finite-volume effects (FVE's)** are **yet to be controlled**, specially **close to the chiral limit** . Thus,
- Demanding a **detailed finite-volume analysis (FVA)** here, and re-estimate the **cPT temperature** T_c and the **UC** .

Some quick basic theory first ...

The focus of this work

Theory

Close to $(T = T_c, m_l = m_{u,d} = m_u = m_d = 0)$ for finite volumes, one can write :

$$M_\ell(T, H, L) = h_0^{-1/\delta} H^{1/\delta} f_G(z, z_L) + M_{\ell, reg} \quad \chi_\ell(T, H, L) = h_0^{-1/\delta} f_\chi(z, z_L) H^{1/\delta-1} + \chi_{\ell, reg}$$

Non-analytic singular and analytic regular terms

Scaling functions

Critical exponent δ

Scaling variables z, z_L : on $N_\sigma^3 \cdot N_\tau$ lattice

$$\left. \begin{aligned} z &= z_0 \left(\frac{T - T_c}{T_c} \right) H^{-1/\beta\delta} \\ z_L &= z_{L,0} \frac{N_\tau}{N_\sigma} H^{-\nu/\beta\delta} \end{aligned} \right\} \begin{aligned} H &= m_\ell / m_s \\ \text{Non-universal constants } &h_0, z_0, z_{L,0} \end{aligned}$$

Lattice volume V :

$$V \sim L^3, L = N_\sigma / N_\tau$$

Heng-Tong's talk

Due to additive divergences in $M_\ell, \dots$

Improved order parameter and scaling function

We use an

→ “Improved” order parameter

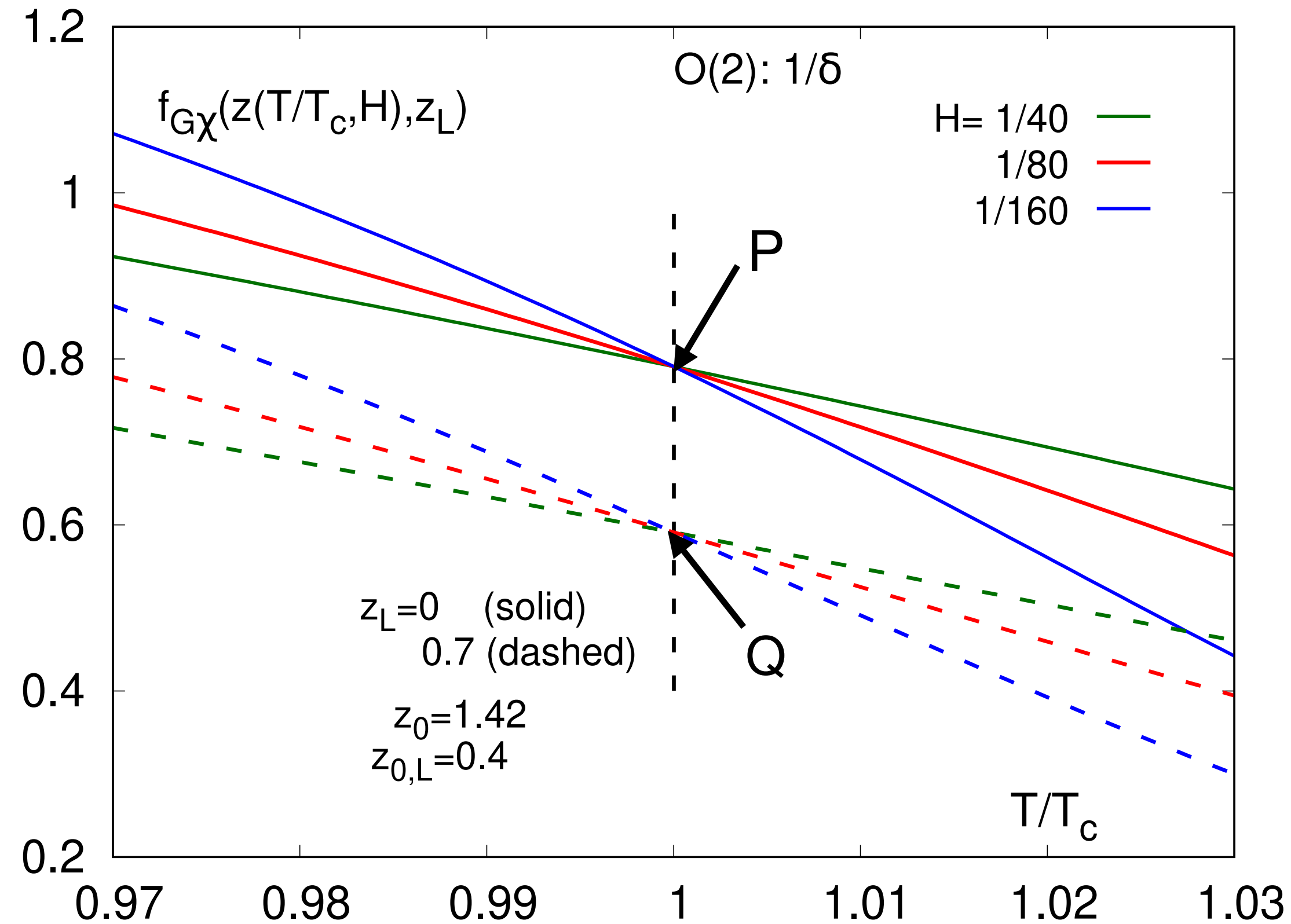
$$M = M_\ell - H\chi_\ell = h_0^{-1/\delta} H^{1/\delta} f_{G\chi}(z, z_L) + M_{reg}$$

↓
“Improved” scaling function

$$f_{G\chi} = f_G - f_\chi$$

- This has **no** additive divergences .
- Therefore, it is **well-defined** in the **continuum and chiral limits** .
- This also has **reduced regular** (*reg*) contribution in H , $\mathcal{O}(H^3) < \mathcal{O}(H)$
- **Augmenting singular behaviour** → making scaling analysis more **convenient** (relatively **straightforward** extraction of T_c and exponent δ)

Let's explore this new scaling function



$$f_{G\chi} = \left(1 - \frac{1}{\delta}\right) f_G + \frac{z}{\beta\delta} f_G^{(z)} + \frac{z_L \nu}{\beta\delta} f_G^{(z_L)}$$

**A unique intersection point for
all H lines at $T = T_c$**

$V \rightarrow \infty$
 $z_L = 0$
P

Obtain δ estimate
directly

$z_L \neq 0$
Q $V < \infty$

Not exact δ , finite-volume
 $\equiv \mathcal{O}(z_L)$ corrections

**Encountered in
Lattice calculations**

So ...

$$f_{G\chi}(z=0, z_L=0) = (1 - 1/\delta)$$

Here, we used δ of $O(2)$

- **On lattice**, it is **not possible** to **extract δ directly** $\leftarrow$ **finite-volume corrections**
- Hence, a careful finite-volume analysis is very much required .
- To enable reliable infinite V – extrapolation , and thus, determine T_c and δ
- Following a **parameter-free approach directly from LQCD calculations** (**no prior assumption of the universality class required** before conducting LQCD estimates)

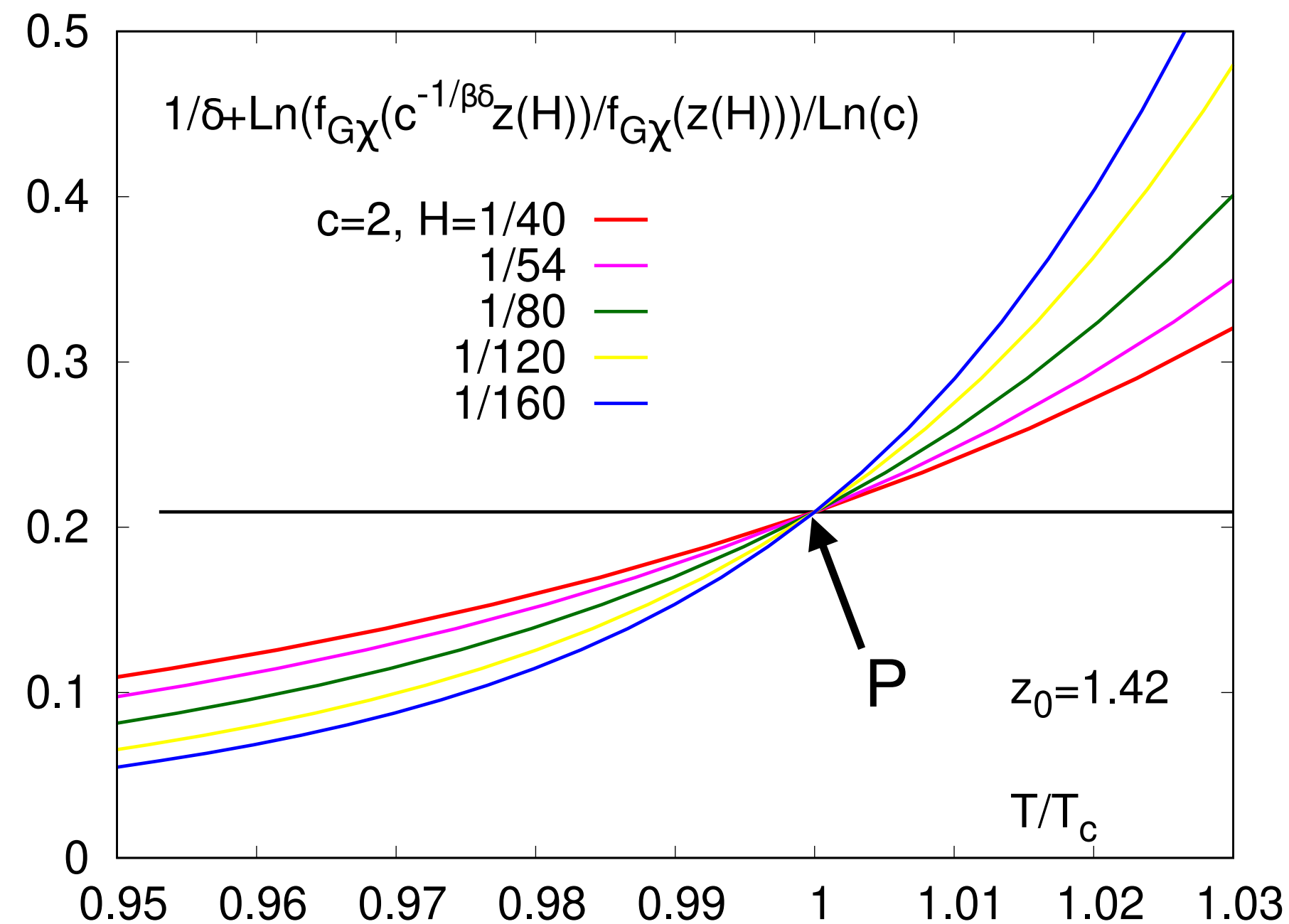
A quick snippet into this parameter-free method

Mitra et. al. , PoS Lattice2024 (2025) 187

Parameter-free method

LATTICE 2024 talk

$\ln V \rightarrow \infty$, construct



$$B(T, H_1, H_2) = \frac{\ln \left[\frac{M(T, H_1)}{M(T, H_2)} \right]}{\ln \left[\frac{H_1}{H_2} \right]} = \frac{1}{\delta} + \ln \left[\frac{f_{G\chi}(z_1(T, H_1))}{f_{G\chi}(z_2(T, H_2))} \right]$$

Fully from lattice

In chiral limit ($H \rightarrow 0$)

◦ $B(T_c, H_1, H_2) = 1/\delta$ with $H_2 = cH$, $H_1 = H$.

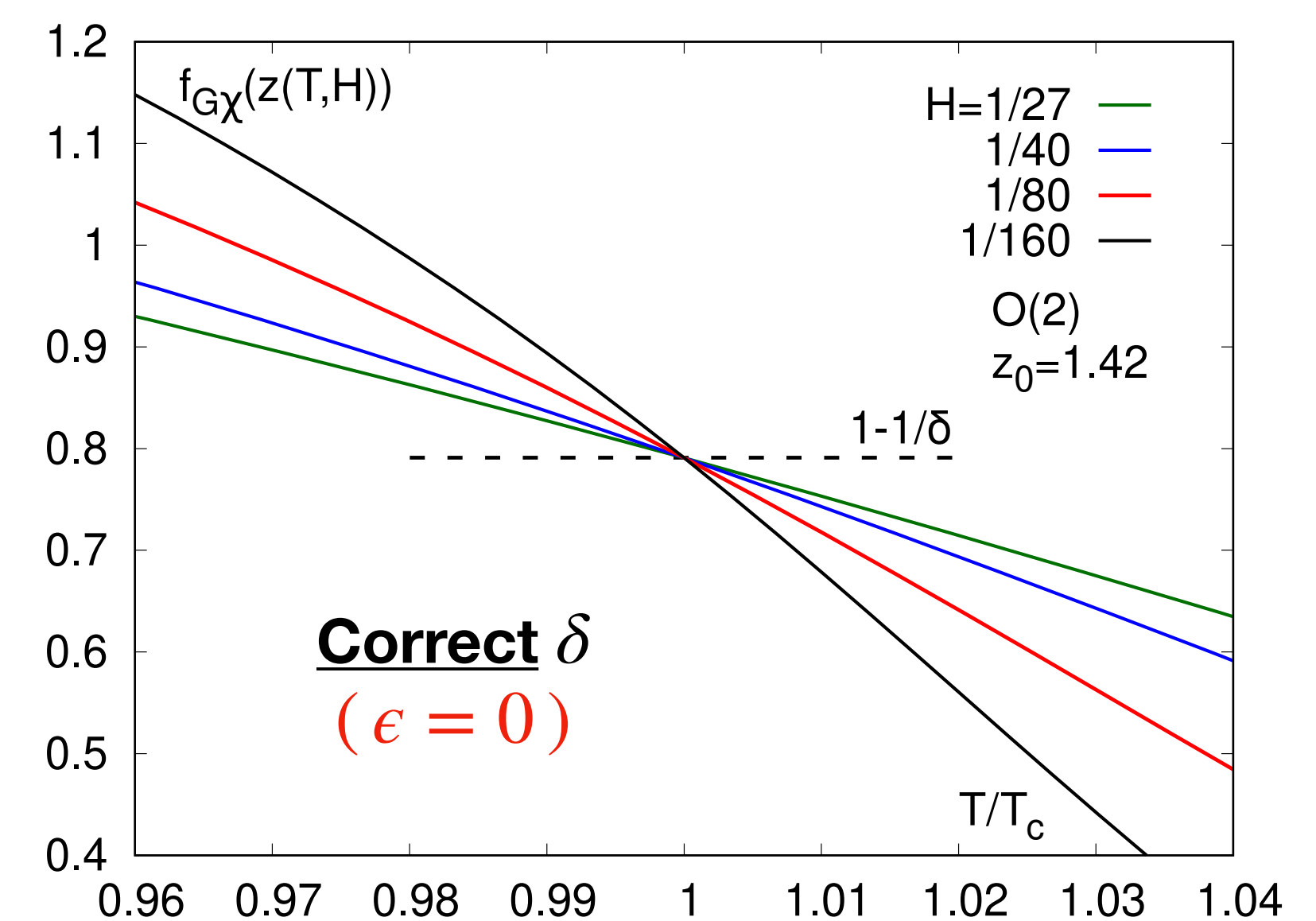
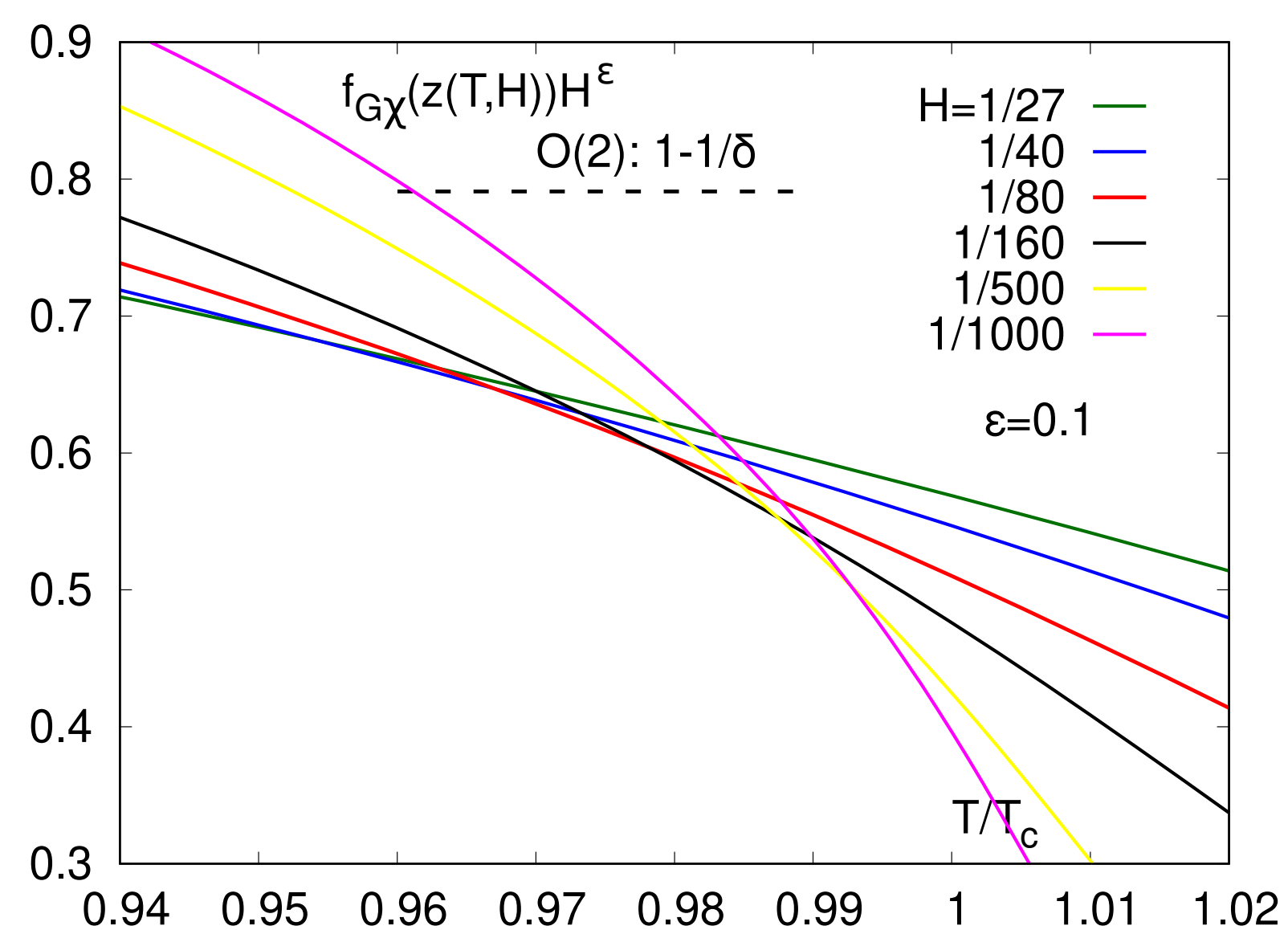
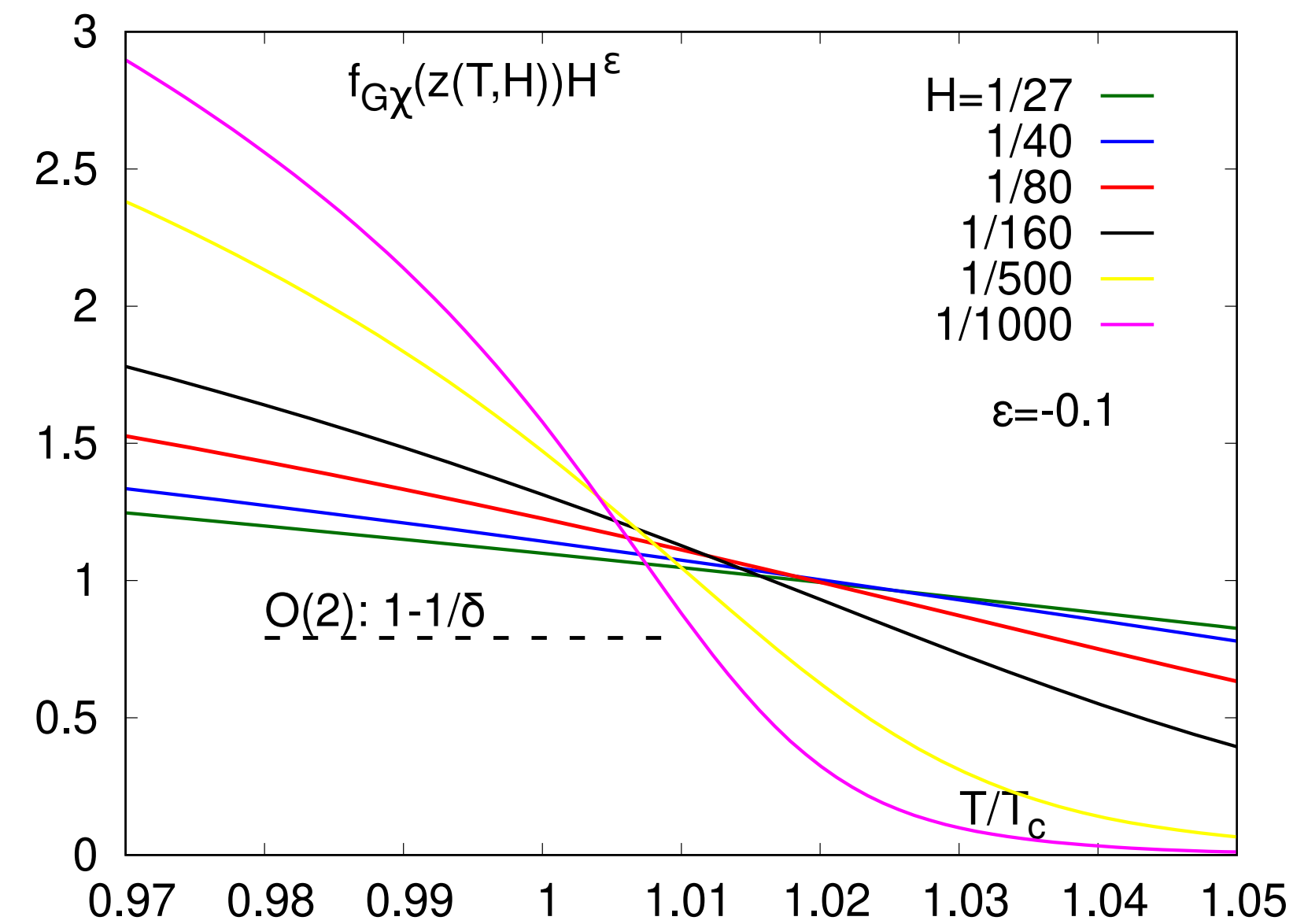
◦ We see the unique intersection point $\mathbf{P} \left(T_c, \frac{1}{\delta} \right)$ in B vs T plane

A way to extract T_c and δ from **direct LQCD calculations**

Feature of uniqueness
important

As we will see now ...

Uniqueness as indicator



We compute here

Wrong δ ($\epsilon \neq 0$)

$$\widetilde{M} = H^{-1/\delta} M = h_0^{-1/\delta} H^\epsilon f_{G\chi}(z, z_L) + M_{reg, H}$$

calculated on lattice

**Exploit this feature
of uniqueness of
intersection in
LQCD calculations**

We use δ of $O(2)$ in our LQCD estimates (educated ansatz)

But before that, ...

Volume dependence in scaling function

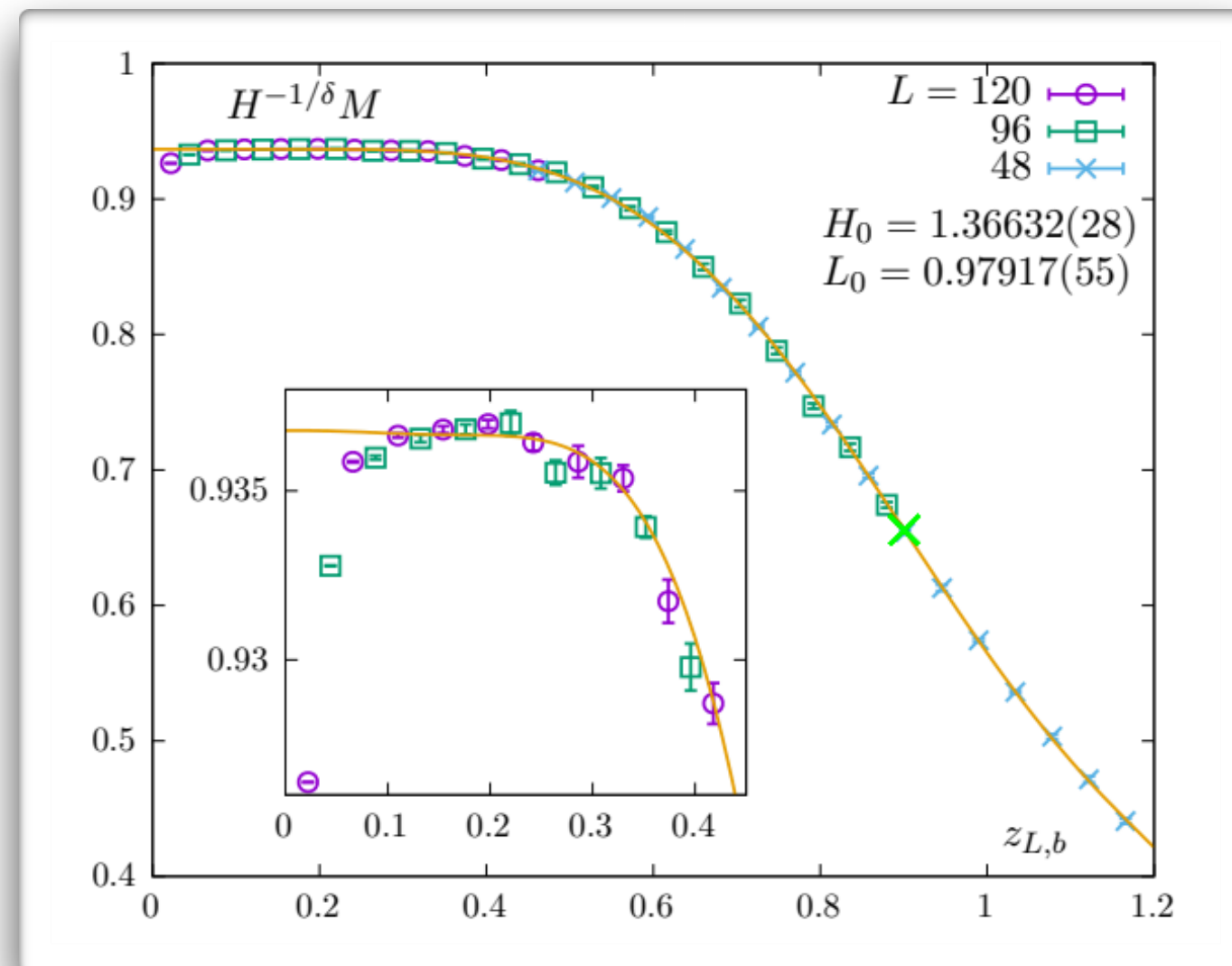
Taylor expansion of $f_{G\chi}$ around $z_L = 0$ [Karsch et. al. , PRD 108, 014505](#)

$$f_{G\chi}(z, z_L) = f_{G\chi}(z, 0) + \sum_{n=0}^{n_u} \sum_{m=3}^{m_u} \left(1 - \frac{1}{\delta} + \frac{n + m\nu}{\beta\delta} \right) a_{nm} z^n z_L^m$$

$$z_{L,b} = L^{-1} H^{-\nu/\beta\delta}$$

Infinite volume
scaling function

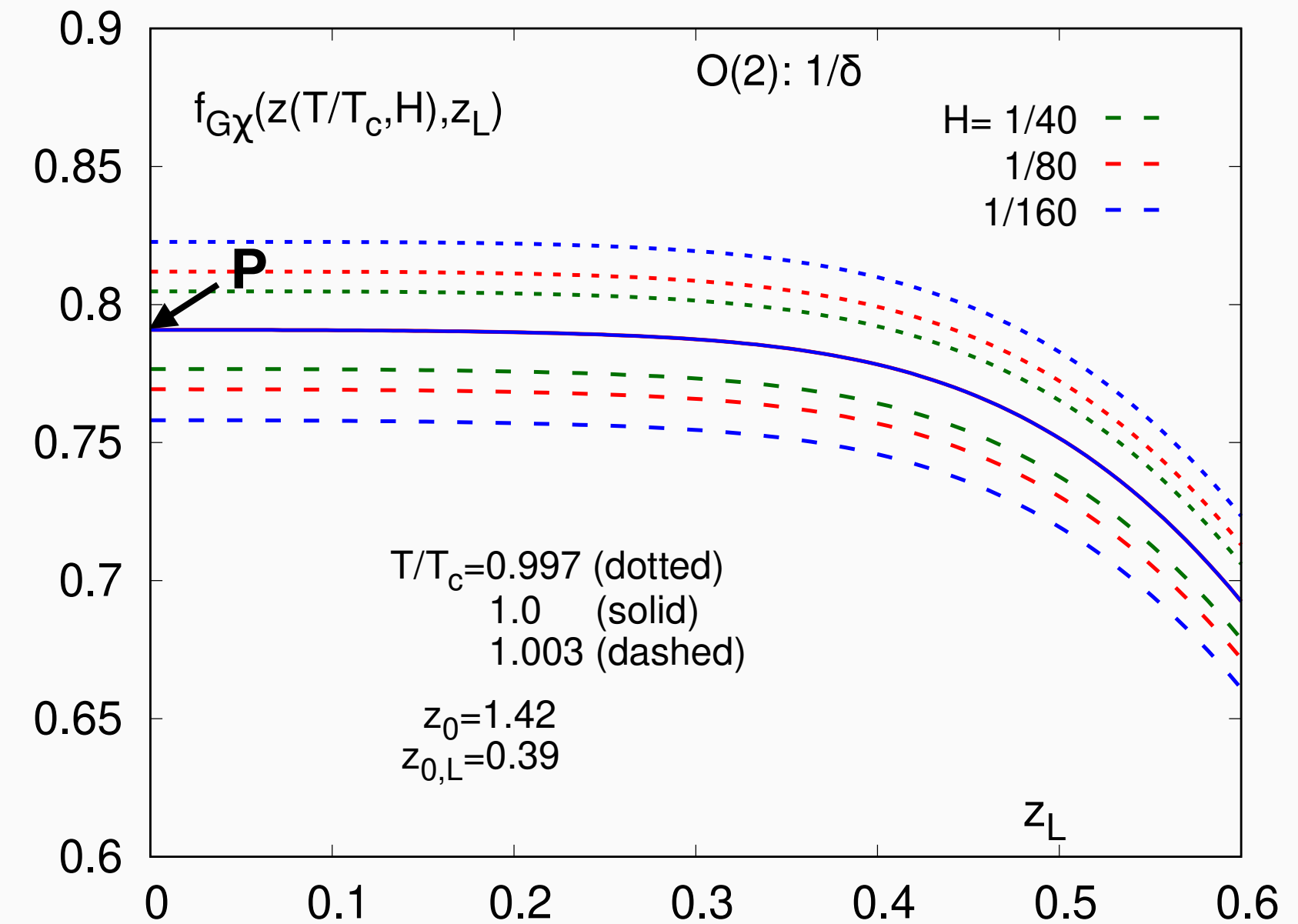
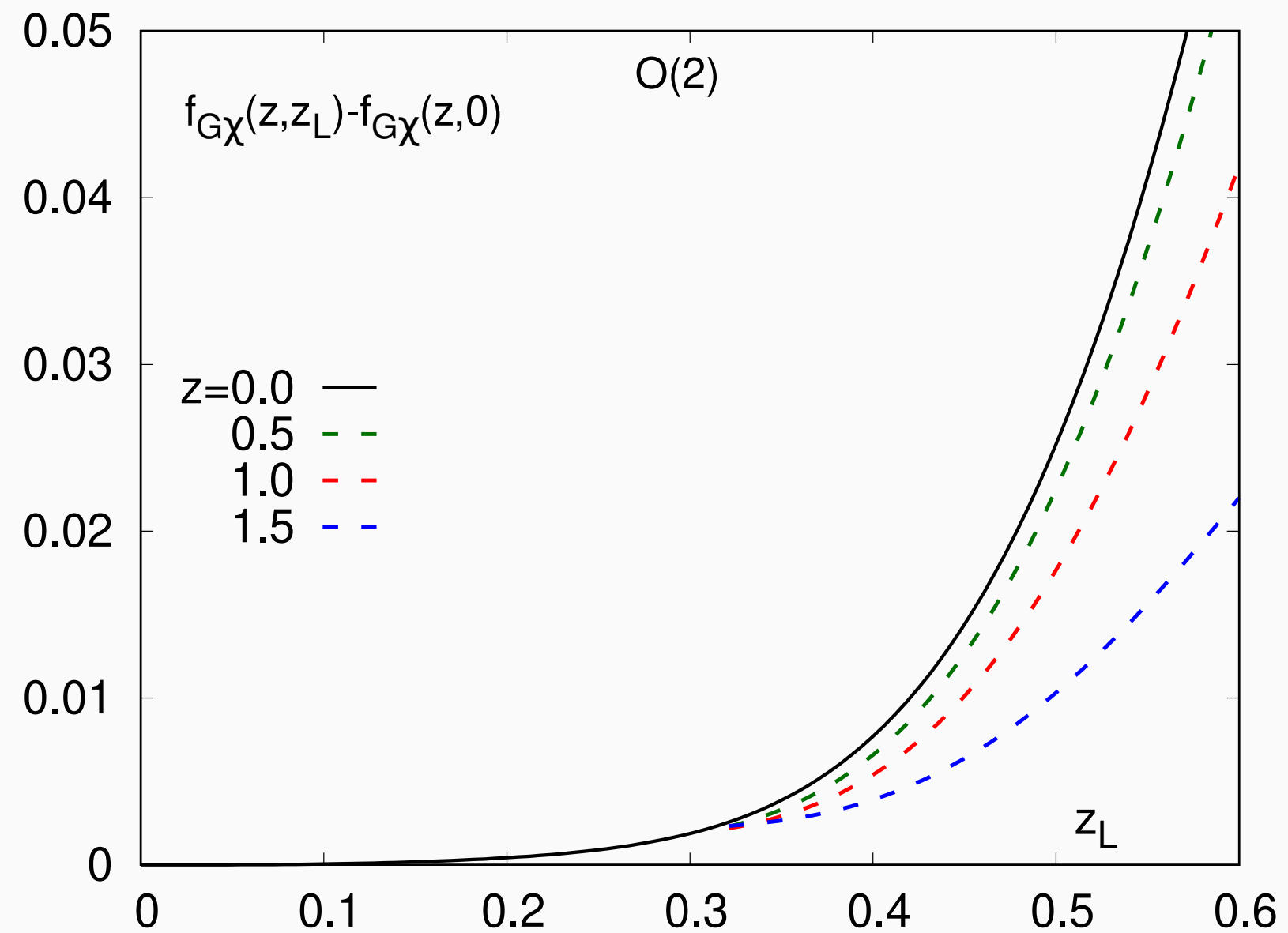
- $m_l = 3$, as Leading order correction $\sim \mathcal{O}(1/V)$
- $n_u = 5$, $m_u = 8$ $\leftarrow$ best fit to Monte-Carlo data on different lattices
- $\beta \sim 0.35$, $\delta \sim 4.78$, $\nu = \beta(\delta + 1)/3 \in 3\text{-d } O(2)$ universality class
 - $f_{G\chi}(0,0) = (1 - 1/\delta)$, at $T = T_c$, $V^{-1} = 0$



[Karsch et. al. , PRD 108, 014505](#)

Some FVE observations ...

Finite-volume improved scaling function



- FVE's are **more severe (sharp rise)** as one goes **towards T_c** i.e. **lower z (greater sensitivity)**
- **Well-controlled** finite-volume effects (FVE's) for $z_L \leq 0.4$, for $T \rightarrow T_c$ [Karsch et. al., PRD 108, 014505](#)
- **One unique line** for $T = T_c$ ($z = 0$), whereas **no such uniqueness** for $T \neq T_c$.
- This translates to **unique intersection point** in T plane for $T = T_c$.
- H — dependence for $T \neq T_c$: **ordering reversal** for $z < 0$ ($T < T_c$) $\iff z > 0$ ($T > T_c$)

Note : No LQCD here, these are model-driven results,

What does LQCD say ... ?

Going towards the lattice : Plan forward

- Use δ of $O(2)$ in the pre-factor $H^{-1/\delta}$
- Use the finite-volume dependent form of the scaling function $f_{G\chi}$ (shown before)
- Thus, compare the lattice data with these benchmarks, and
- Try to improve the finite-volume analysis and its systematics .

Lattice setup ...

Lattice setup

- We use **highly improved staggered quarks / fermions** (HISQ).
- All our simulations are done **presently on $N_\tau = 8$ lattices**.
- We use δ of 3-d $O(2)$ universality class , for finite $V \sim L^3$ calculations here
- We use $T_c = 145.1$ MeV

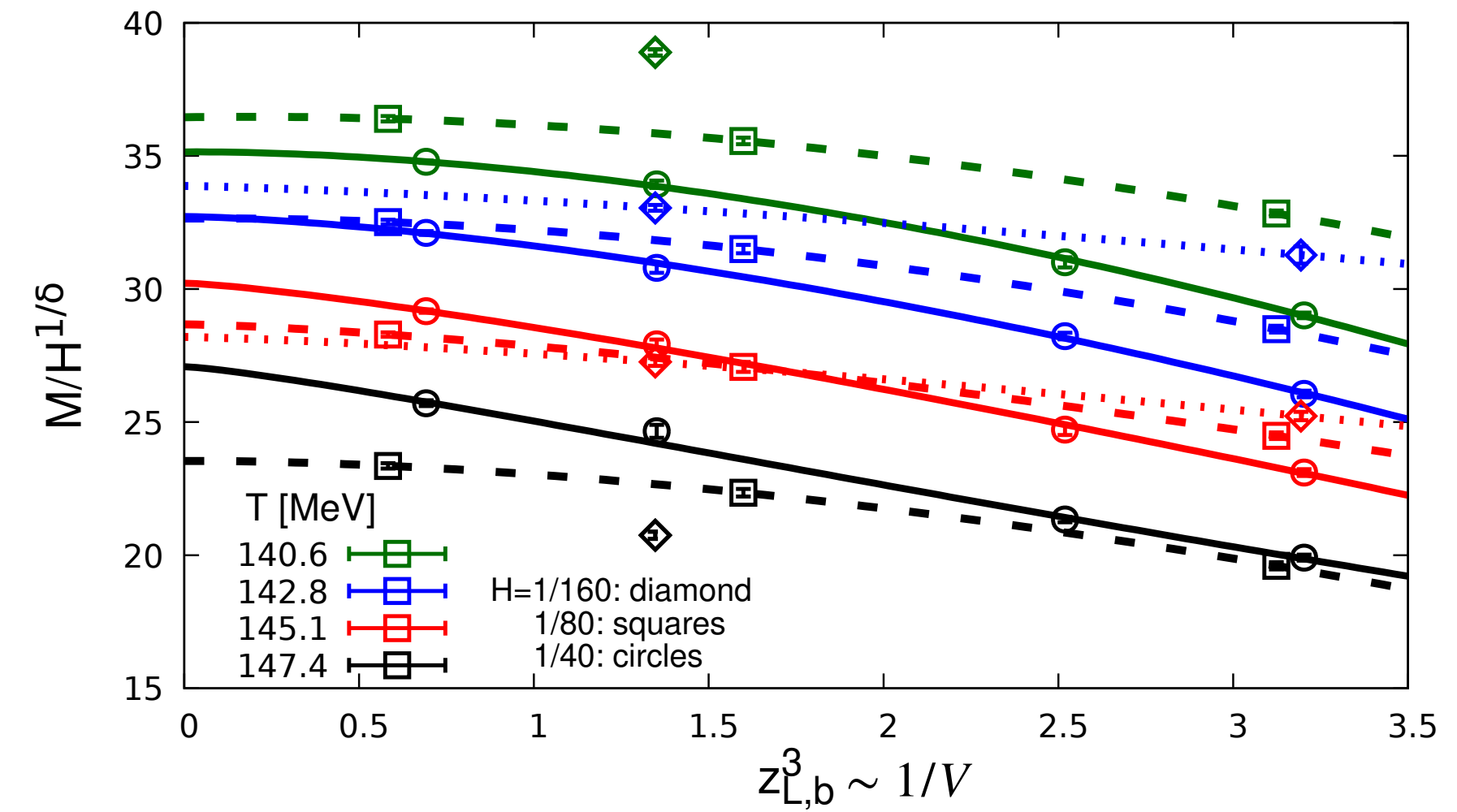
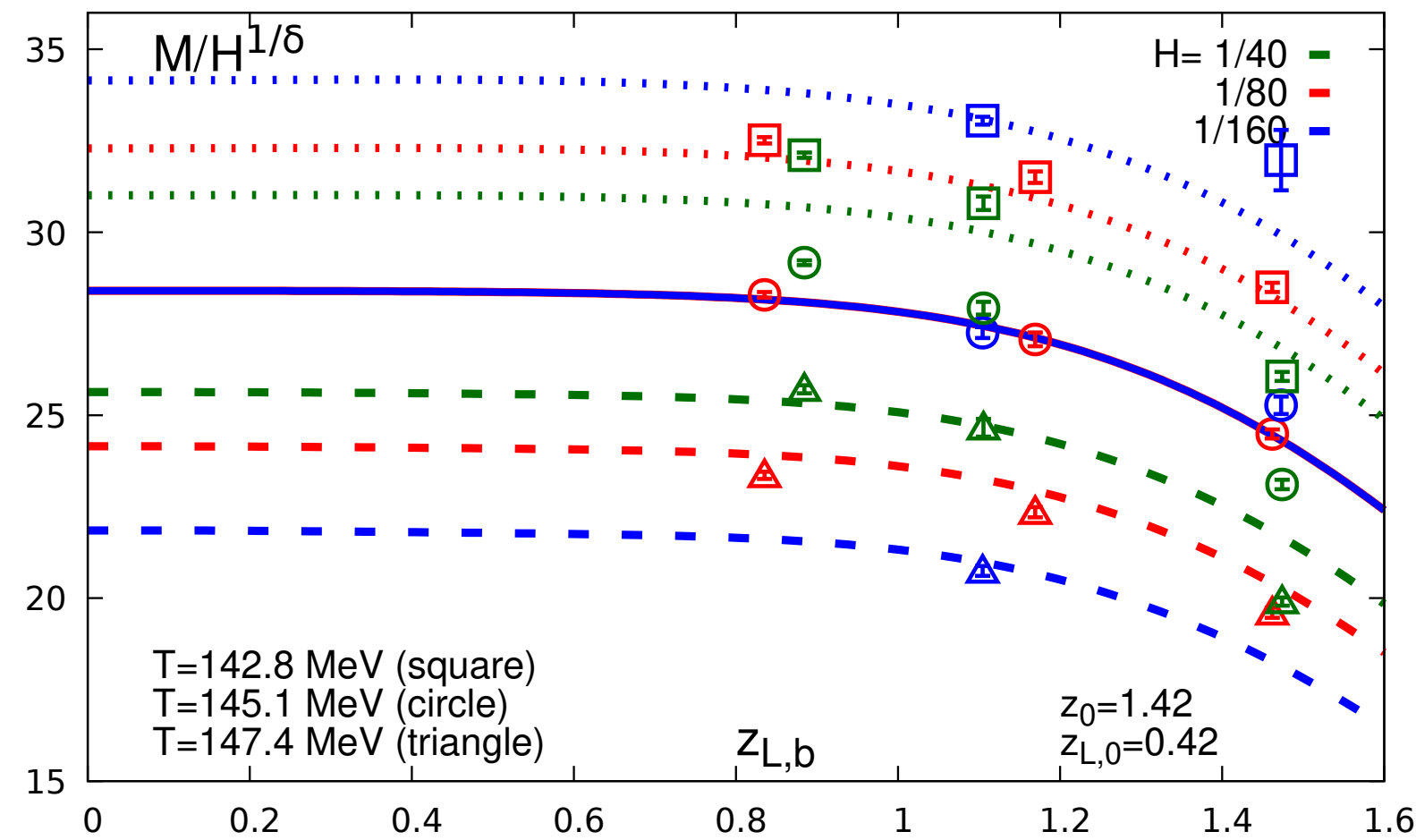
Staggered fermions retain only $O(2)$ part of the full chiral symmetry

The parameters T, H, L are tuned as follows :

- We vary T close to T_c , following $T \in (0.99 T_c , 1.01 T_c)$
- With physical m_s , m_ℓ is varied with $1/240 \leq H \leq 1/27$
- Different spatial volumes **on $N_\tau = 8$** : $L = N_\sigma / N_\tau \in (3 , 10)$

$$\begin{aligned} |z| &< 0.5 , \\ 0.4 &< z_L < 1 \end{aligned}$$

Lattice Results



Better agreement between Lattice data $\iff O(2)$ benchmarks :

- For **lower** $z_{L,b}$ $\iff$ **closer to infinite V limit** ($z_{L,b} = 0$)
- For **lower** H values $\iff$ **approaching chiral limit** ($H \rightarrow 0$)

Bare scaling variable

$$z_{L,b} = z_L / z_{L,0} = \frac{N_\tau}{N_\sigma} H^{-\nu/\beta\delta}$$

Ongoing $\longrightarrow H = 1/160$, on $78^3 \cdot 8$ lattices for $T = 142.8, 145.1, 147.4$ MeV

A third blue point for three T

$$z_{L,b} \sim 0.8, z_{L,b}^3 \sim 0.5$$

Closer

Conclusions

- We work with an **improved order parameter** M having **well-defined continuum and chiral limits**, **reduced regular contributions** and **no additive divergences**.
- **Uniqueness is an important signature** to identify correct δ (essence of a unique point)
- We find **good agreement** between the **expected $O(2)$ singular behaviour** and **the Lattice QCD results** .
(An **increasingly convincing indication** about the $O(2)$ **nature of cPT universality class**) .
- This agreement **improves** as expected for **smaller H** (chiral limit), and **smaller $z_{L,b}$** (infinite- V limit)
- **More well-controlled** finite-volume effects for $z_{L,b} < 1 \Rightarrow$ need for **higher N_σ** , for **smaller H**
(demanding more computational expansivity)

Still a lot going on, and a lot to do

Outlook

- Come up with the results of the new $H = 1/240$ data point
- Quantify regular contributions and analyse their volume dependence .
- Continue with our finite-volume work $\Rightarrow$ infinite V extrapolation $\Rightarrow$ parameter-independent way of extraction of δ from direct LQCD calculations **Mitra et. al. , PoS Lattice2024 (2025) 187**
- Add more (T, m_ℓ) data points to attain higher precision, with the goal to distinguish between δ of $O(2)$ and that of $O(4)$ (in continuum) .
- Approach continuum limit and re-estimate volume corrected, continuum estimate of T_c (how different is it ??) .
 - **Ding et. al. , PRL123 (2019) 6, 062002**

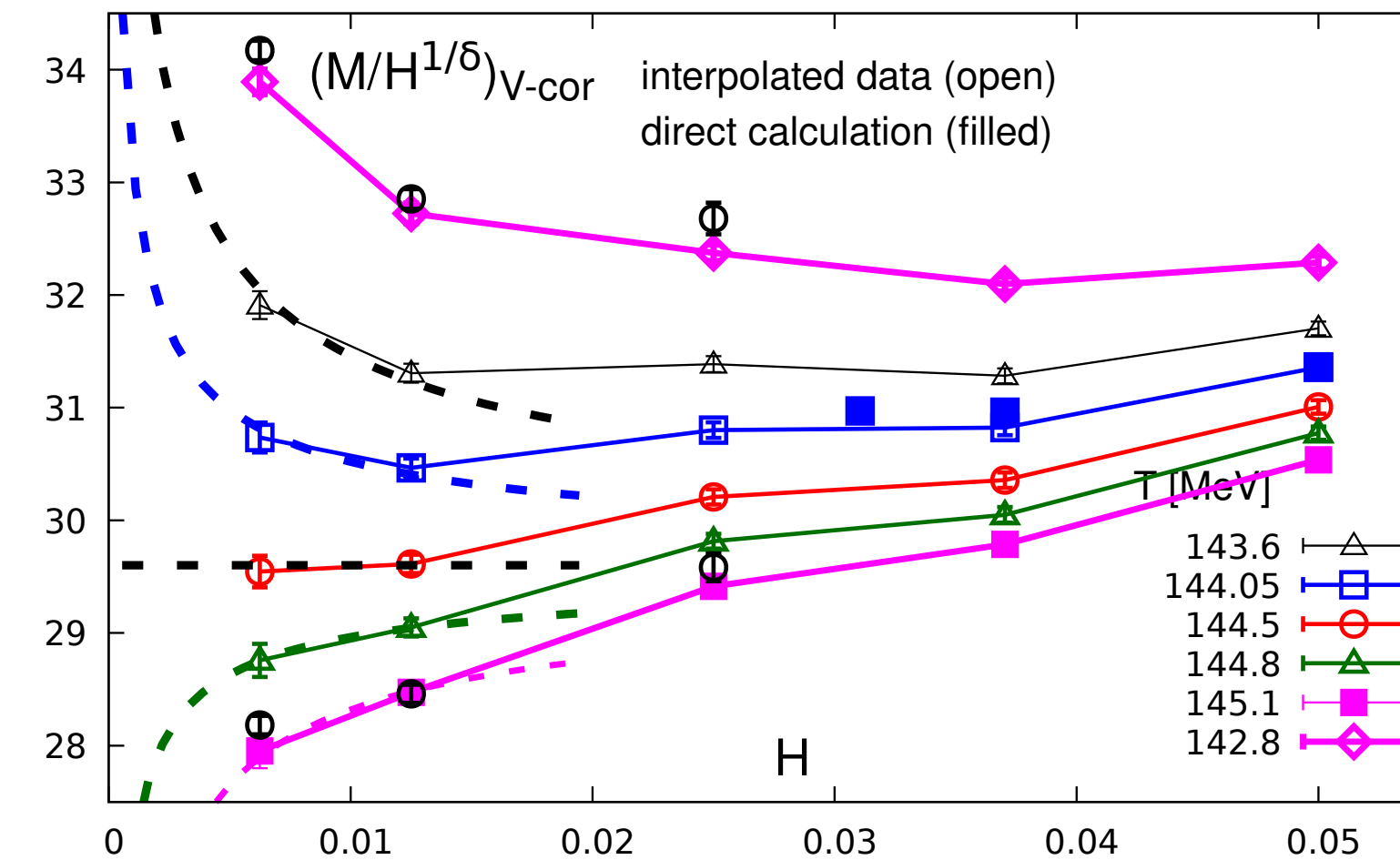
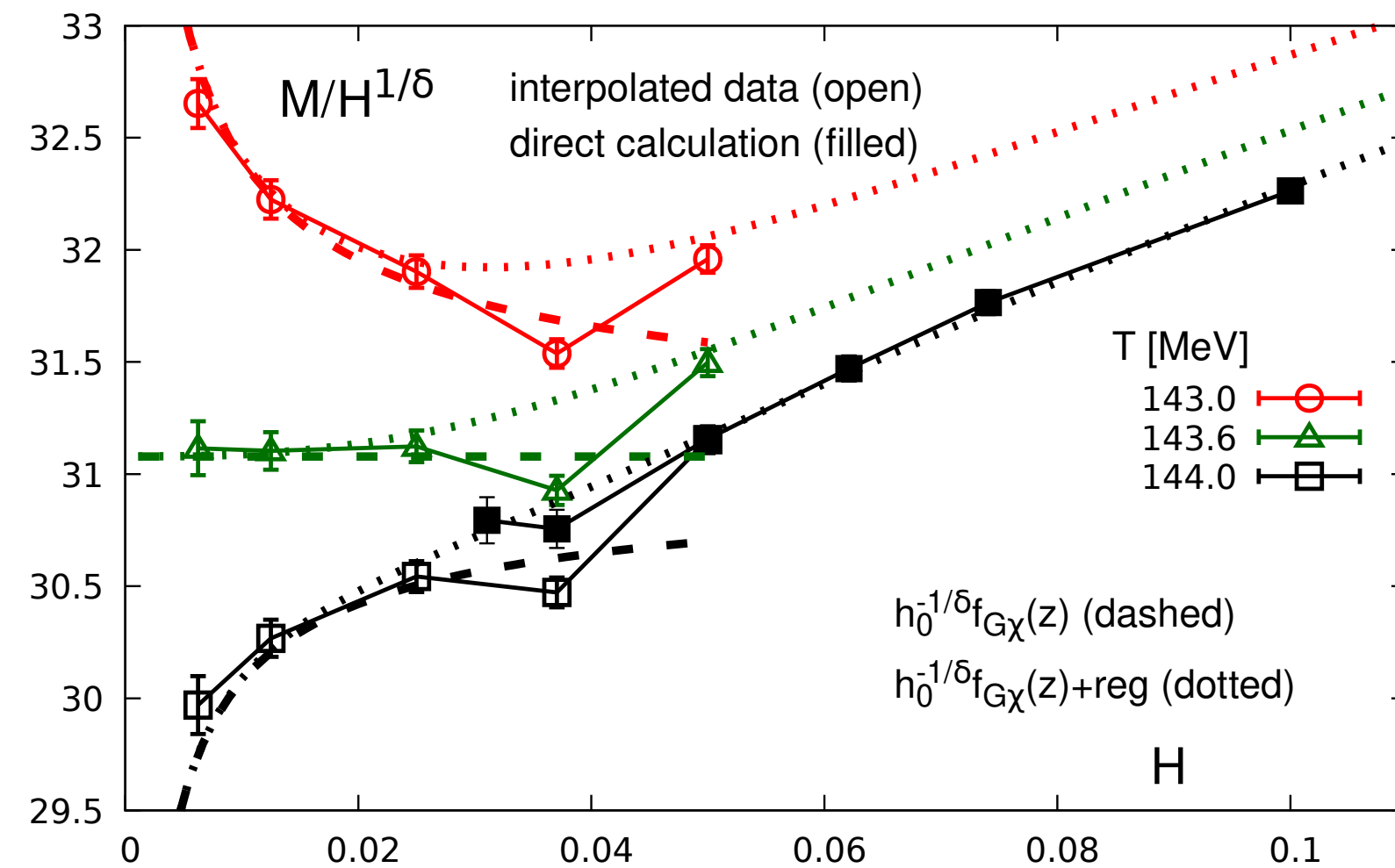
THANK YOU
for your
PATIENCE and
ATTENTION



Credit: Organisers

Backup slides

Lattice QCD results



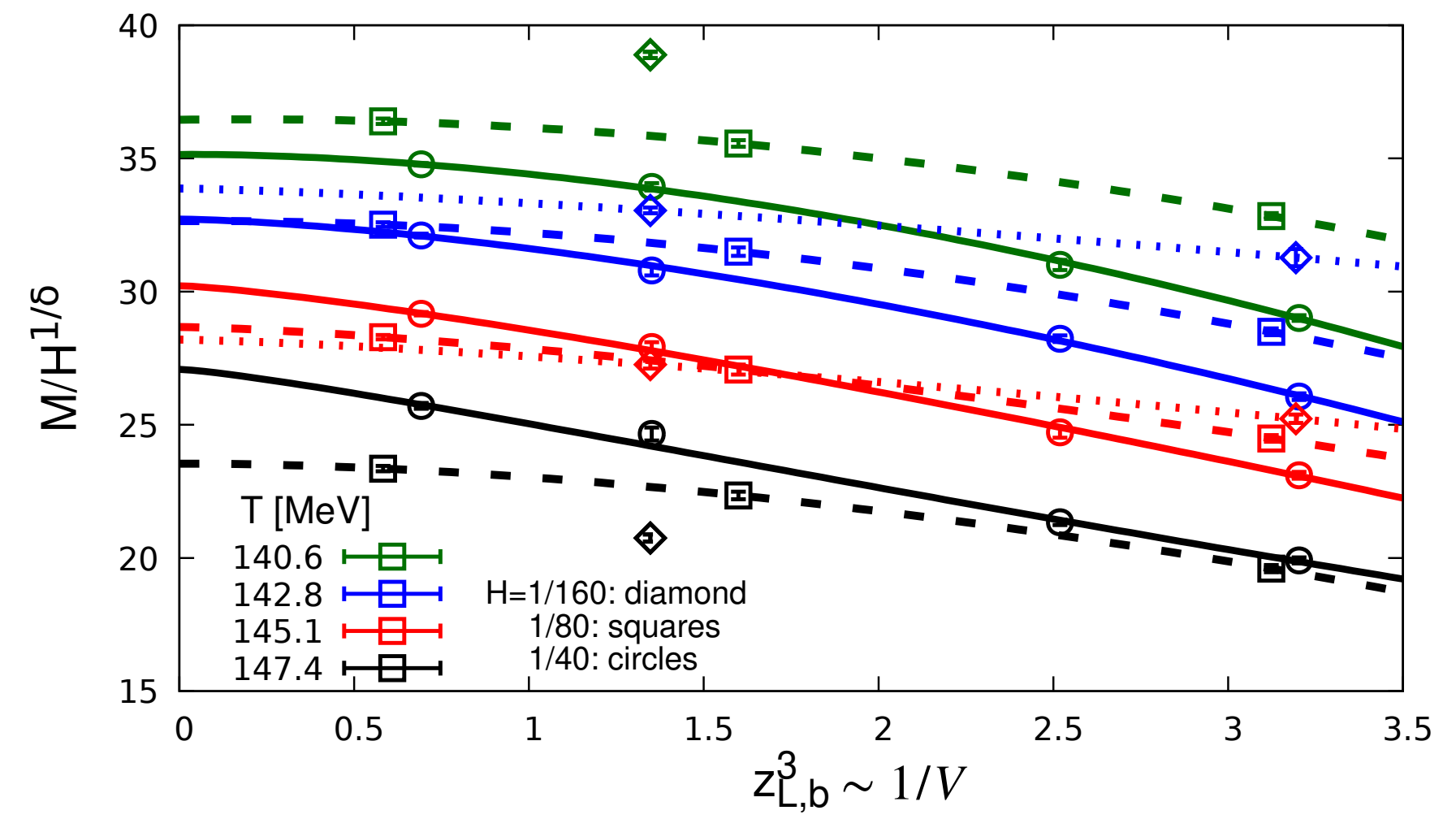
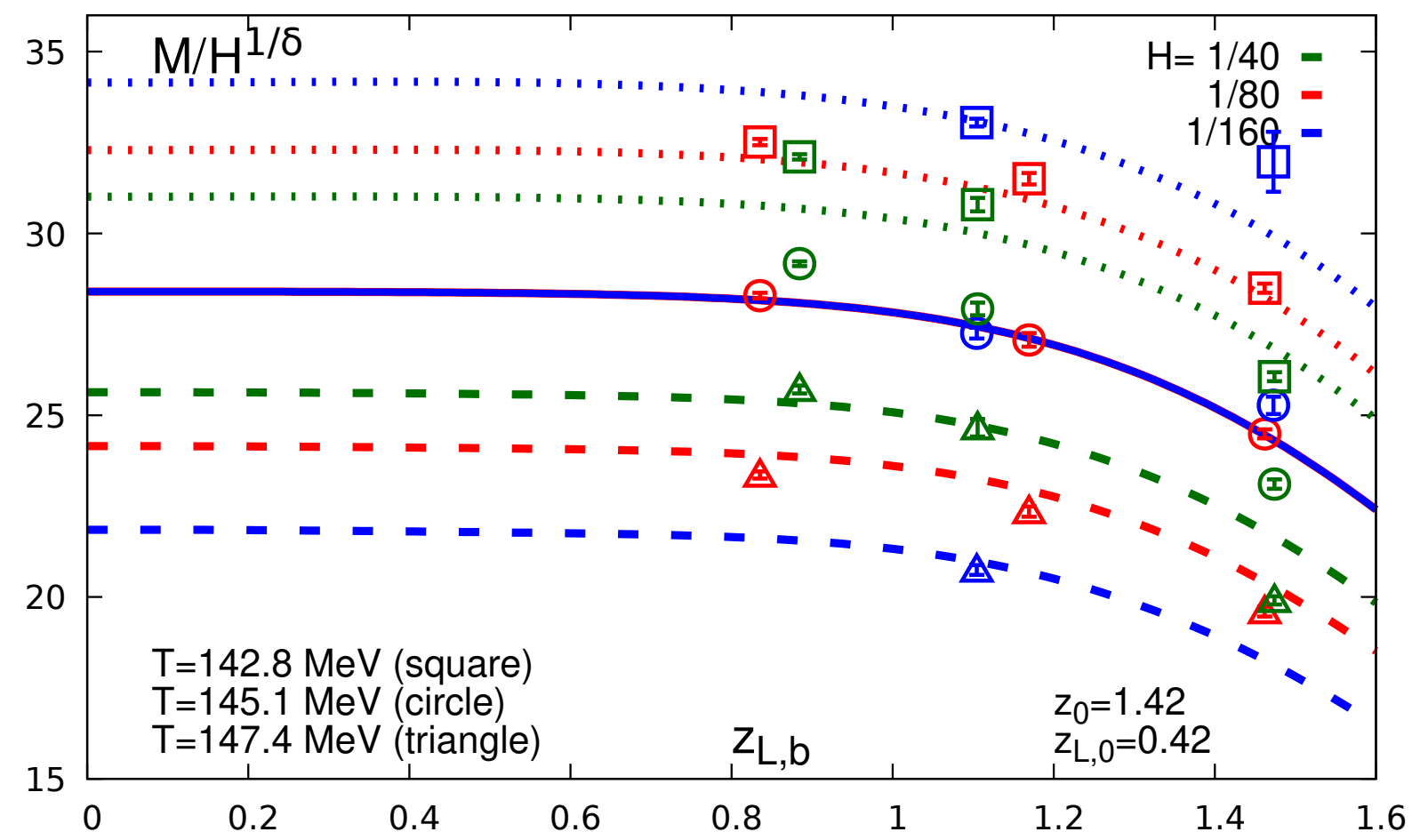
LQCD results for
142.8 , 144.05 , 145.1
MeV

- Good agreement with universal (singular) behaviour apparent for sufficiently low $H < 1/40 = 0.025$ (expected in chiral limit)
- Also, visible deviations from this for “heavier” $m_\ell \equiv$ higher H
- Regular contributions still to be quantified

- Similar story of agreement and deviations, offering universal manifestations in small H .
- More control on FVE's on lowest $H = 1/160$ data point , with two lattice data values computed

Causing new estimate of
 $T_c = 144.5 > 143.6$ MeV

More Work on higher N_σ in Progress ...



$$f(x) = a + b x^{4/3} + c x^{5/3} + \mathcal{O}\left(z_{L,b}^6\right)$$

$$x = z_{L,b}^3$$