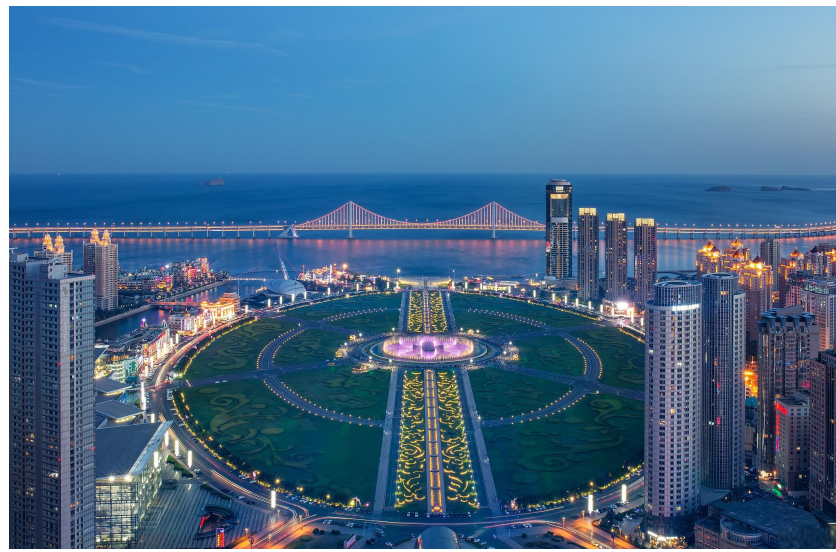




Disentangling the nature of the momentum fluctuations and radial flow at the RHIC-STAR experiment

Chunjian Zhang

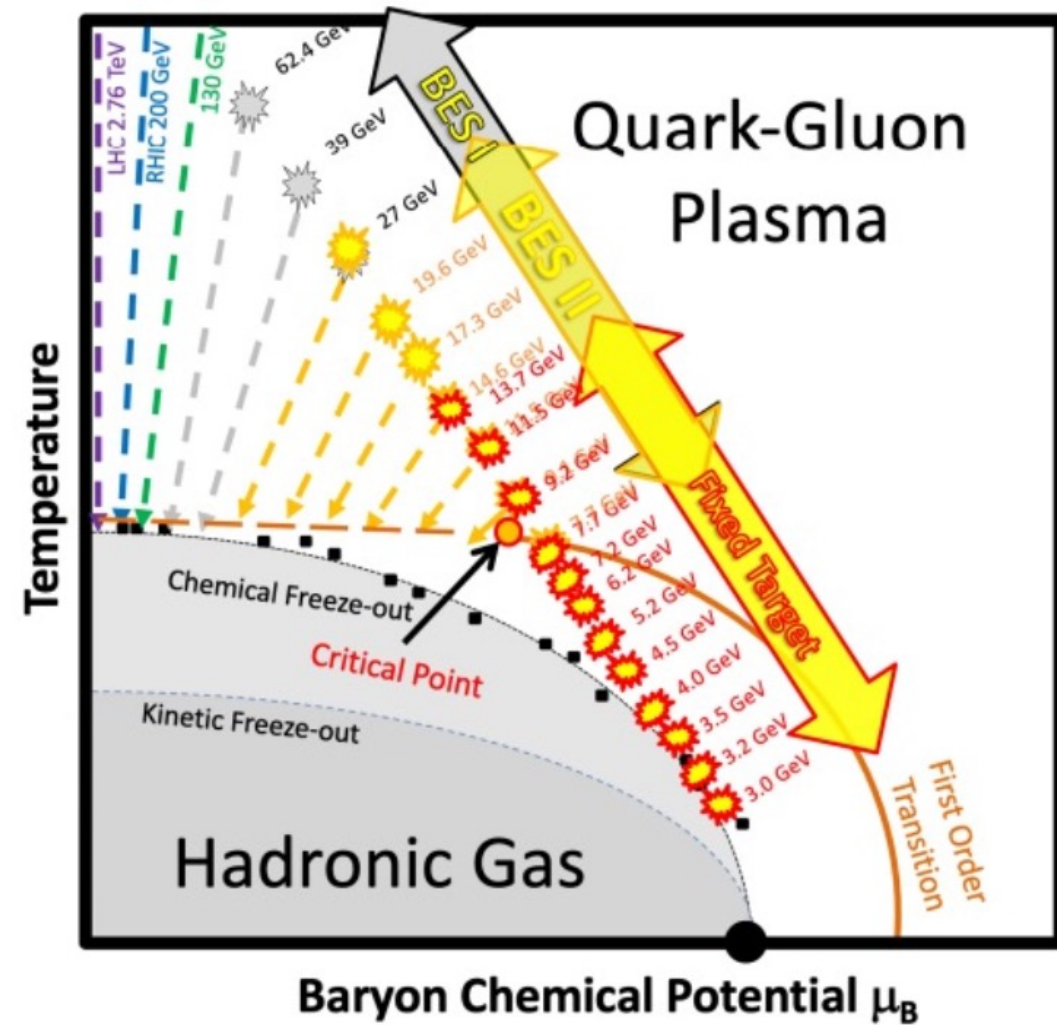
Fudan University

October 8-11, 2025, Dalian

The QCD Phase Diagram: From Theory to Experimental Signatures



QCD Phase Diagram and RHIC Beam Energy Scan



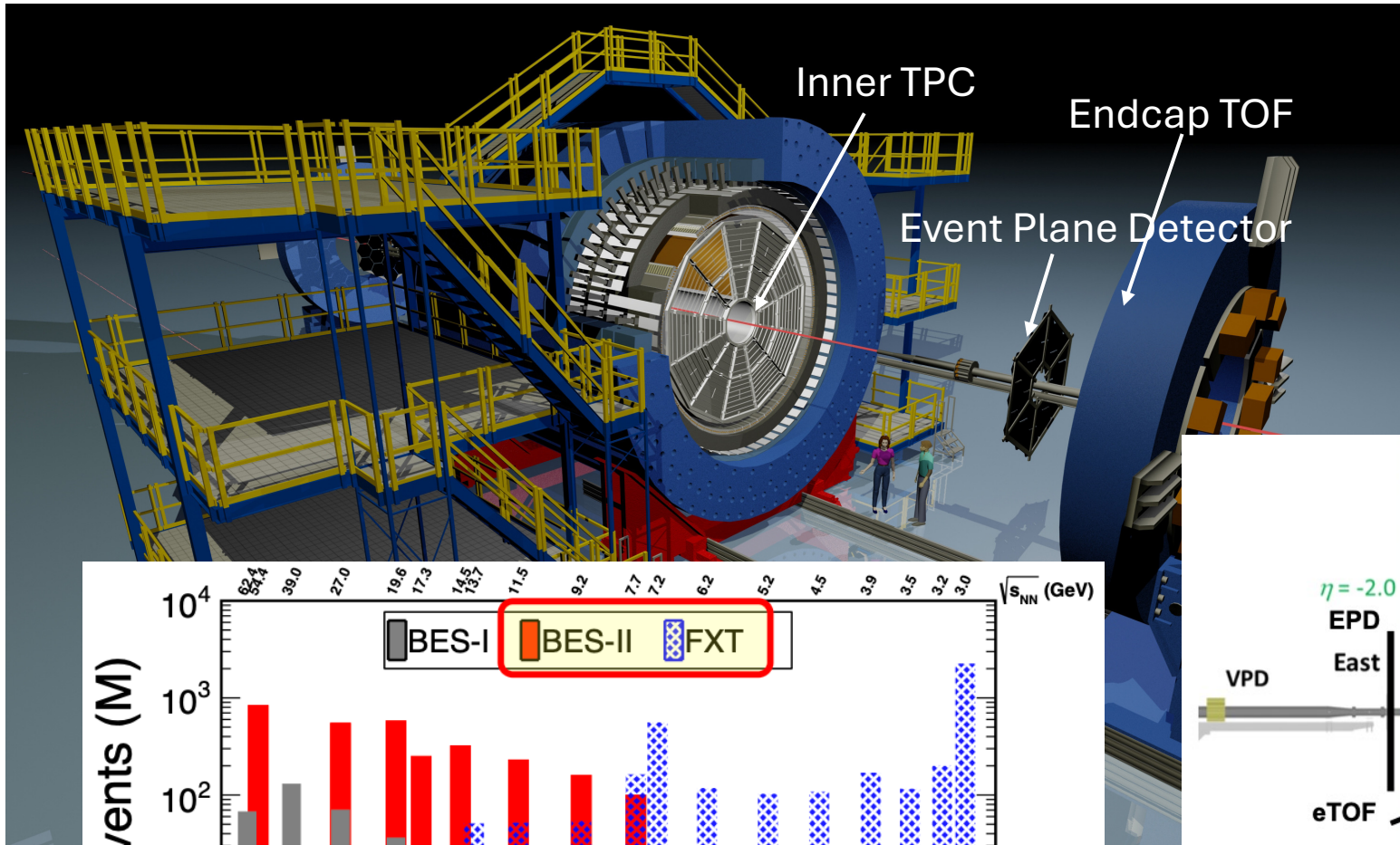
Determine QCD phase diagram via high energy heavy-ion collisions: seeking to map onset of deconfinement, and the predicted QCD critical point

Crossover at μ_B close to 0 and predicted 1st-order phase transition at high μ_B

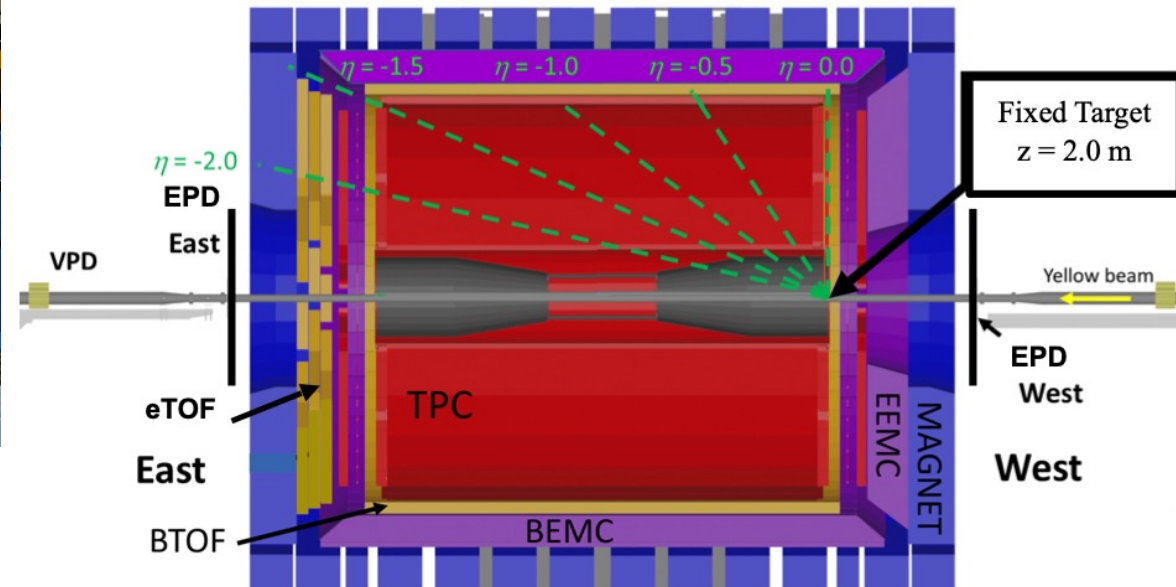
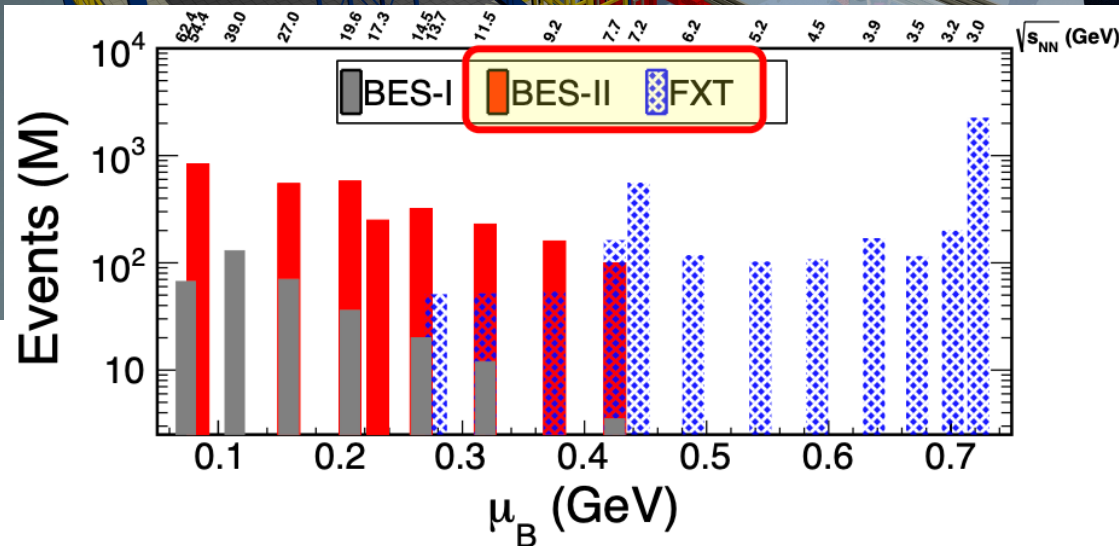
Beam Energy Scan program at RHIC:

- Vary initial T and μ_B
- Explore different trajectories across the phase boundary

The BES-II Upgrades and fixed-target (FXT) setup at STAR



Enhanced Acceptance (p_T and y)
Enhanced PID mid and forward
Enhanced Event Plane Resolution
Enhanced Centrality Definition
Enhanced $\sqrt{s}$ range



Mid-rapidity for 3.0 GeV is $y = 1.049$

Large datasets with iTPC upgrade $\sim 10 \times$ BES-I

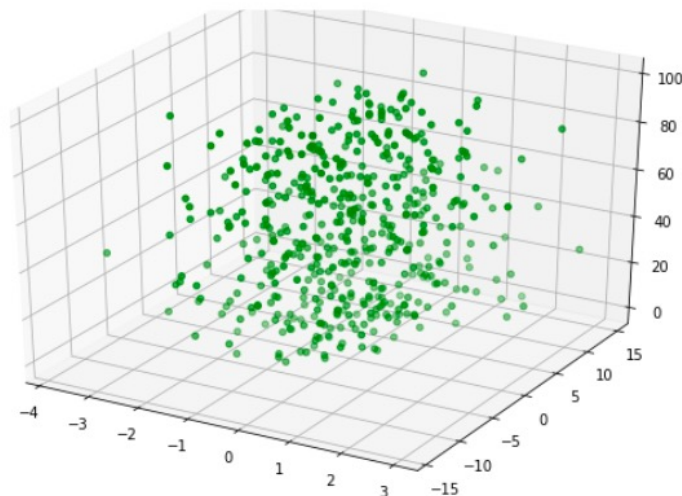
Measurement: mean transverse momentum fluctuations

J. Jia, S. Huang, C. Zhang, PRC 105 , 014906 (2022)

S. Bhatta, C. Zhang, J. Jia, PRC106, L031901 (2022)

C. Zhang, A. Behera, S. Bhatta, J.Jia, PLB 822, 136702 (2021)

$$f(p_{x1}, p_{y1}, p_{z1}, p_{x2}, p_{y2}, p_{z2} \dots)$$



Event-by-event temperature fluctuations

Mean:

$$\langle \langle p_t \rangle \rangle = \frac{\sum_{k=1}^{N_{\text{events}}} \langle p_t \rangle_k}{N_{\text{events}}} \quad \text{where} \quad \langle p_t \rangle_k = \frac{\sum_{i=1}^{N_k} p_{t,i}}{N_k}$$

Variance:

$$\langle (\delta p_T)^2 \rangle = \frac{\sum_i \sum_{j \neq i} \omega_i \omega_j (p_{T,i} - \langle [p_T] \rangle) (p_{T,j} - \langle [p_T] \rangle)}{\sum_i \sum_{j \neq i} \omega_i \omega_j},$$

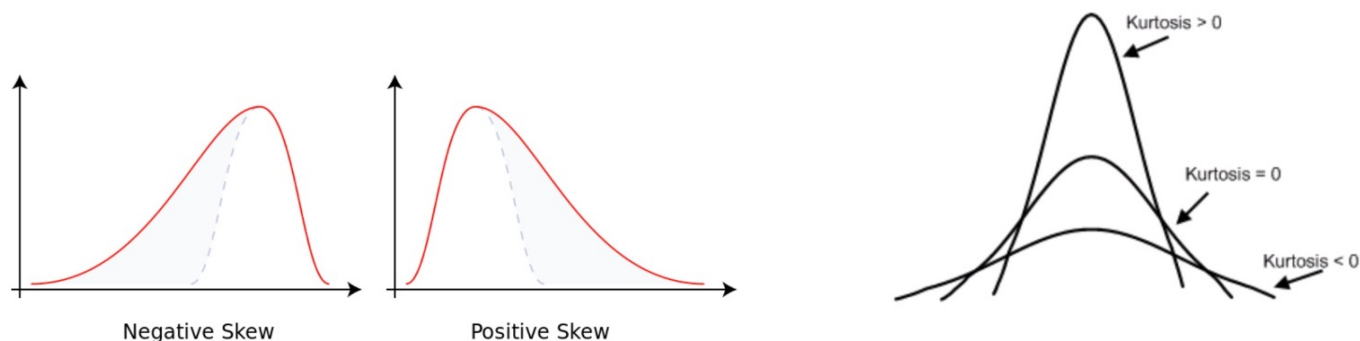
Skewness:

$$\langle (\delta p_T)^3 \rangle = \frac{\sum_{i,j \neq i, k \neq j \neq i} \omega_i \omega_j \omega_k (p_{T,i} - \langle [p_T] \rangle) (p_{T,j} - \langle [p_T] \rangle) (p_{T,k} - \langle [p_T] \rangle)}{\sum_{i,j \neq i, k \neq j \neq i} \omega_i \omega_j \omega_k},$$

Kurtosis:

$$\langle (\delta p_T)^4 \rangle = \frac{\sum_{i,j \neq i, k \neq j \neq i, l \neq k \neq j \neq i} \omega_i \omega_j \omega_k \omega_l (p_{T,i} - \langle [p_T] \rangle) (p_{T,j} - \langle [p_T] \rangle) (p_{T,k} - \langle [p_T] \rangle) (p_{T,l} - \langle [p_T] \rangle)}{\sum_{i,j \neq i, k \neq j \neq i, l \neq k \neq j \neq i} \omega_i \omega_j \omega_k \omega_l}$$

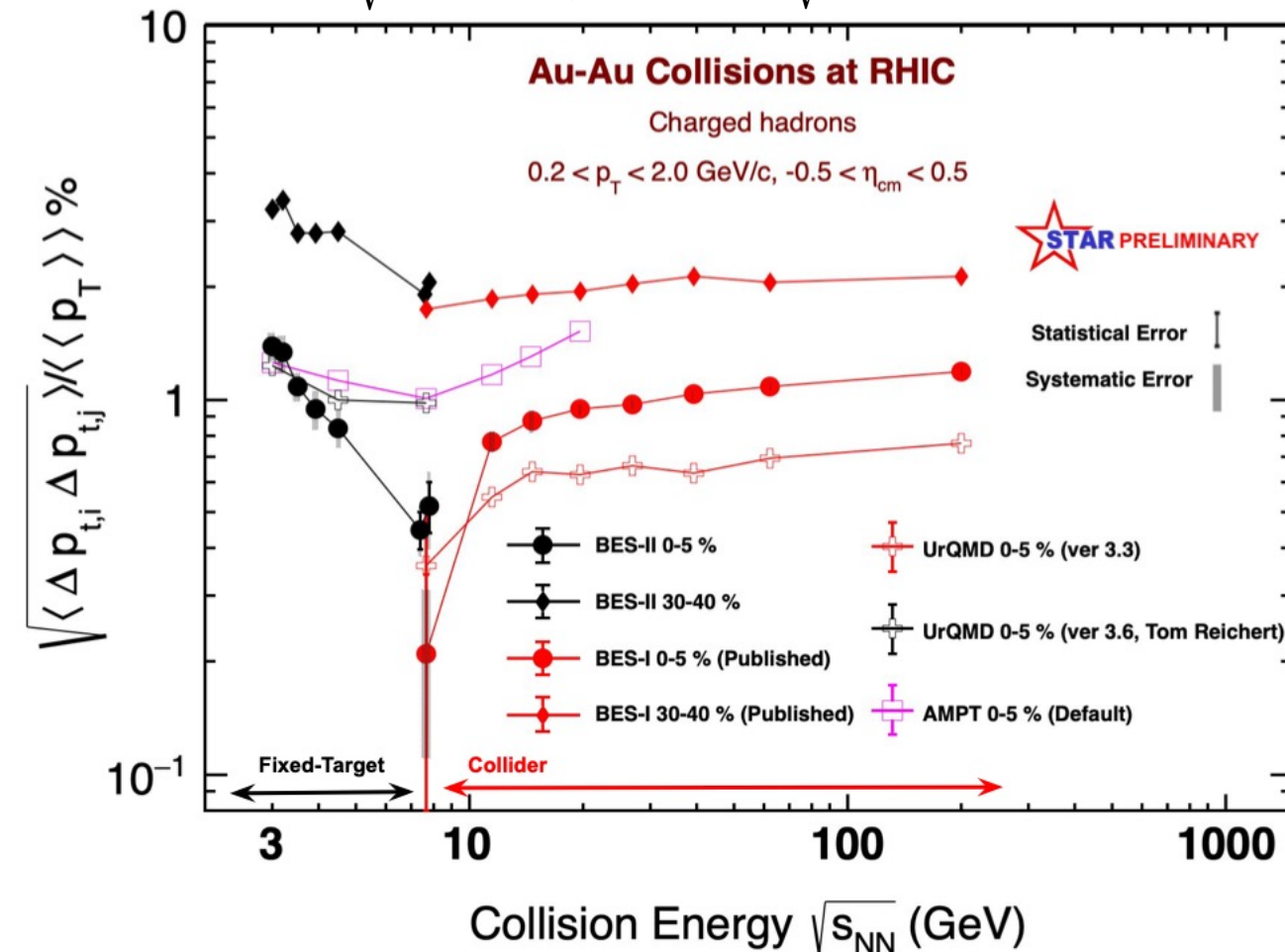
Related to the entropy of the initial source presumably deposited early in the collision before thermalization and interactions with QGP



Correlator $\sqrt{\Delta p_{T,i} \Delta p_{T,j}} / \langle \langle p_T \rangle \rangle$ vs collision energy in FXT mode

Rutik Manikandhan (UH) Quark Matter 2025

$$\sqrt{\langle \Delta p_{t,i} \Delta p_{t,j} \rangle} / \langle \langle p_T \rangle \rangle = \sqrt{\langle (\delta p_T)^2 \rangle} / \langle \langle p_T \rangle \rangle$$



❑ The correlator is scaled by the mean to remove trivial energy dependence

❑ Efficiency independent

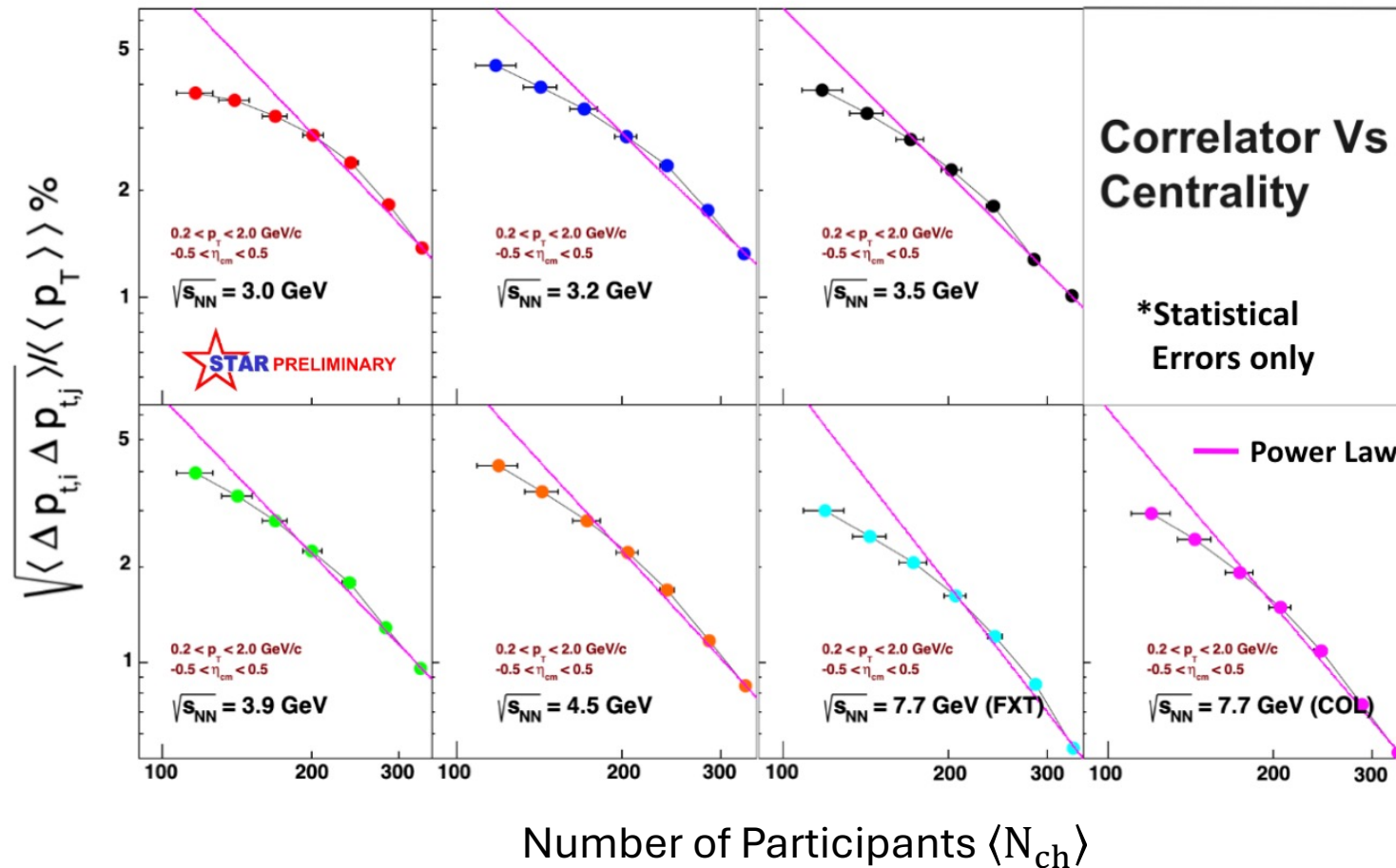
❑ Significant non-monotonic dependence in most central collisions.

❑ Transport codes provide a qualitative description of measurements [4],

Our measurements show non-monotonicity in the region predicted by various theoretical estimates for the location of the critical point.

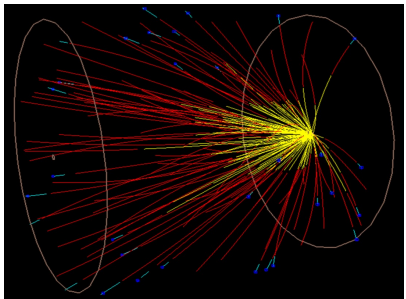
Correlator $\sqrt{\Delta p_{T,i} \Delta p_{T,j}} / \langle \langle p_T \rangle \rangle$ vs centrality in FXT mode

Rutik Manikandhan (UH) Quark Matter 2025



correlations larger in the peripheral collisions and decrease with centrality

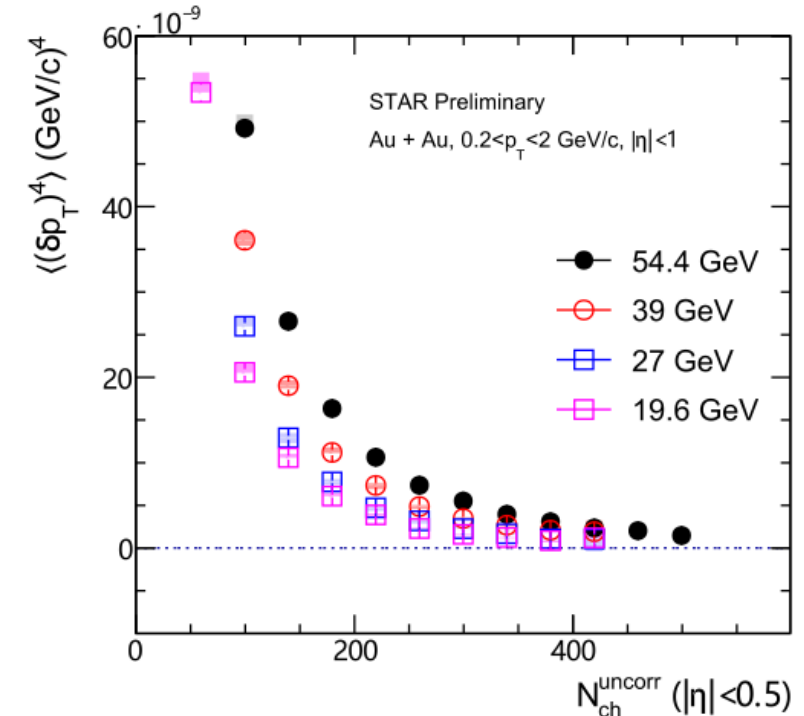
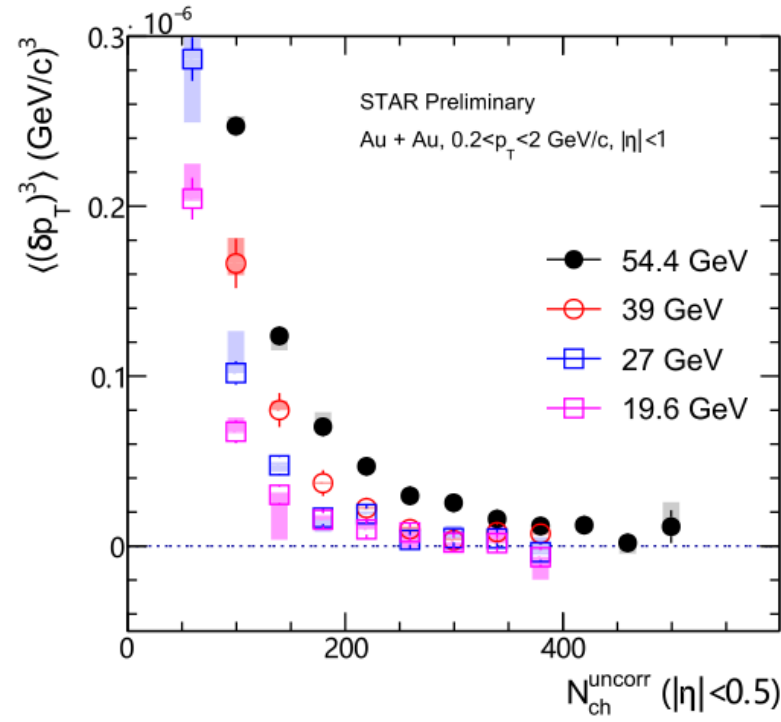
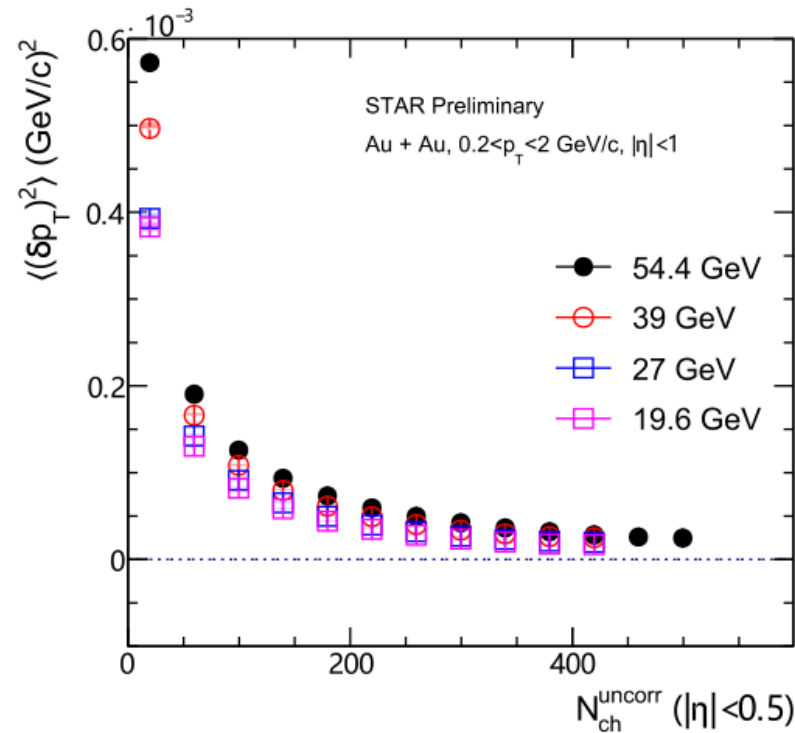
Measurements deviate from power law behavior, implying the centrality dependence does not follow a simple superposition scenario



Independent of the collision mode: good agreement between measurements made at 7.7 GeV in FXT and collider modes.

Correlator $\langle(\delta p_T)^n\rangle$ vs centrality in collider mode

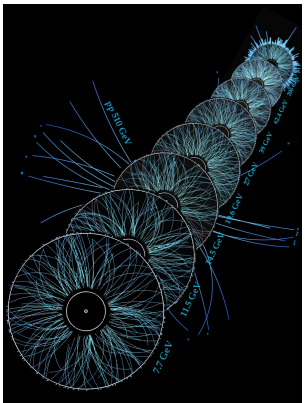
Yining Gao (Fudan) Quark Matter 2025



Efficiency-uncorrected N_{ch}

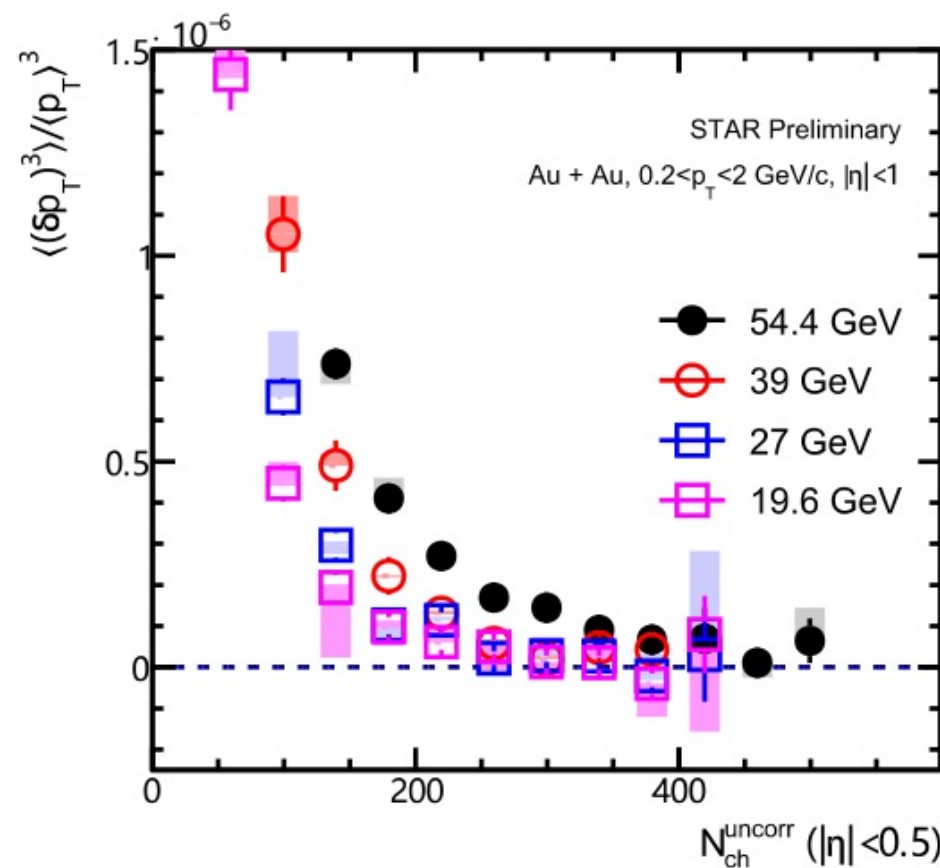
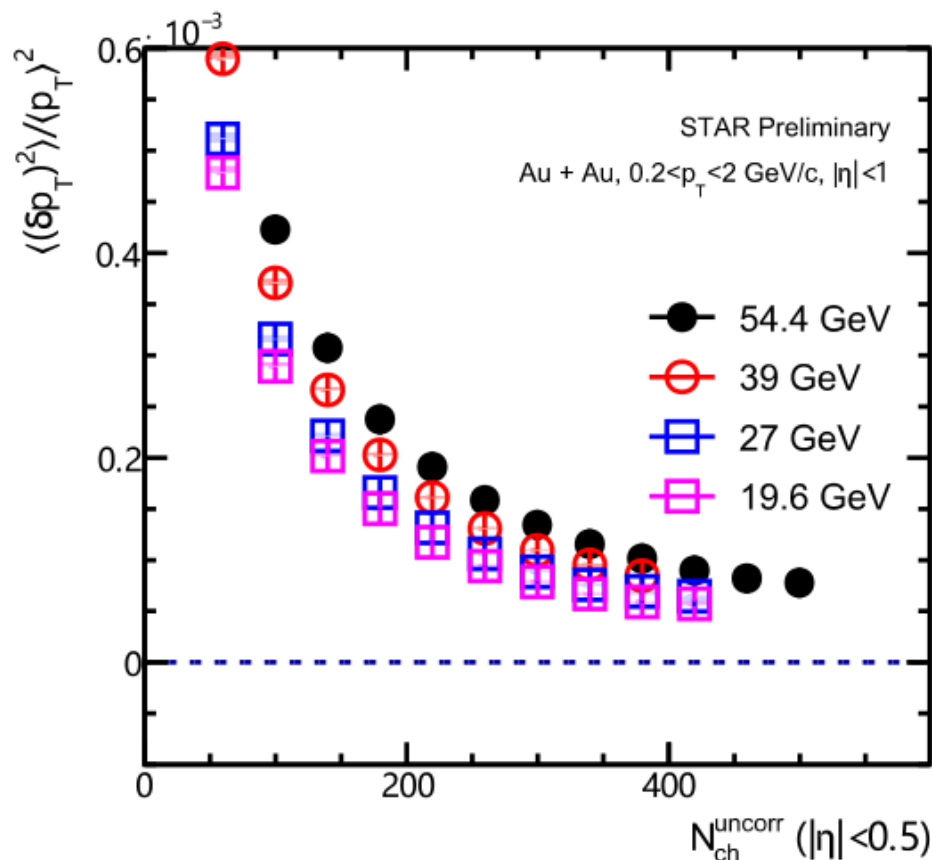
❑ Values of correlators of variance, skewness and kurtosis increase towards peripheral collisions.

❑ Slight energy dependence is observed.



Correlator $\langle(\delta p_T)^n\rangle$ vs centrality in collider mode

Yining Gao (Fudan) Quark Matter 2025



□ With increasing charge-multiplicity, the normalized correlators of variance and skewness decrease in magnitudes.

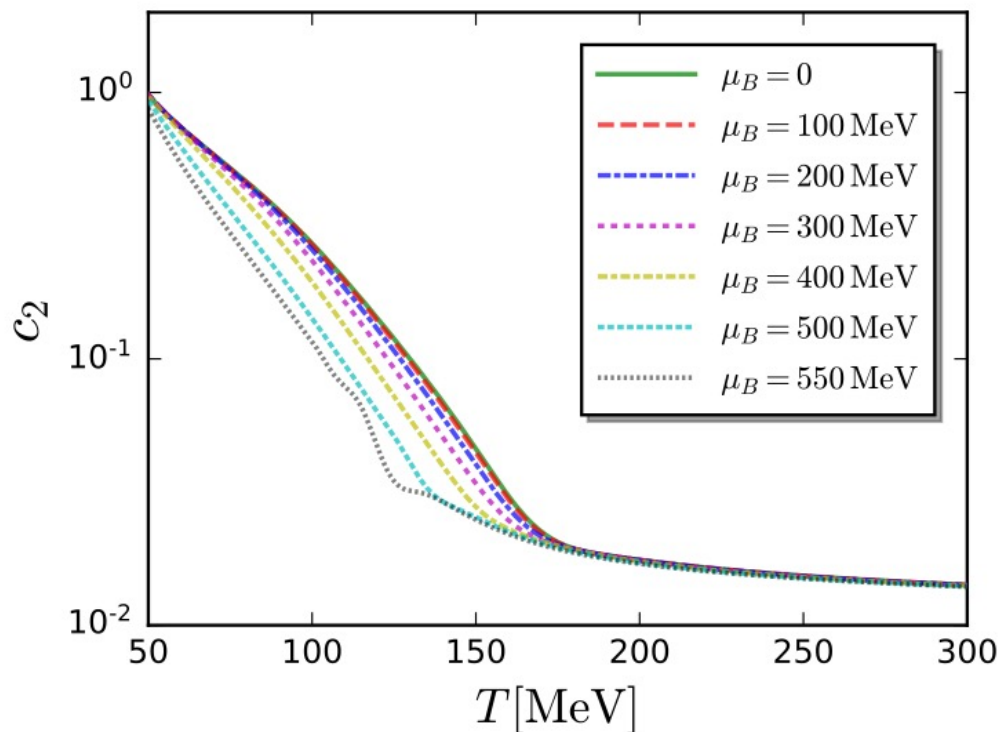
□ Energy dependence that increase towards peripheral collisions are observed.

Our measurement are useful to further understand initial conditions and EbE fluctuations across different collision energies.

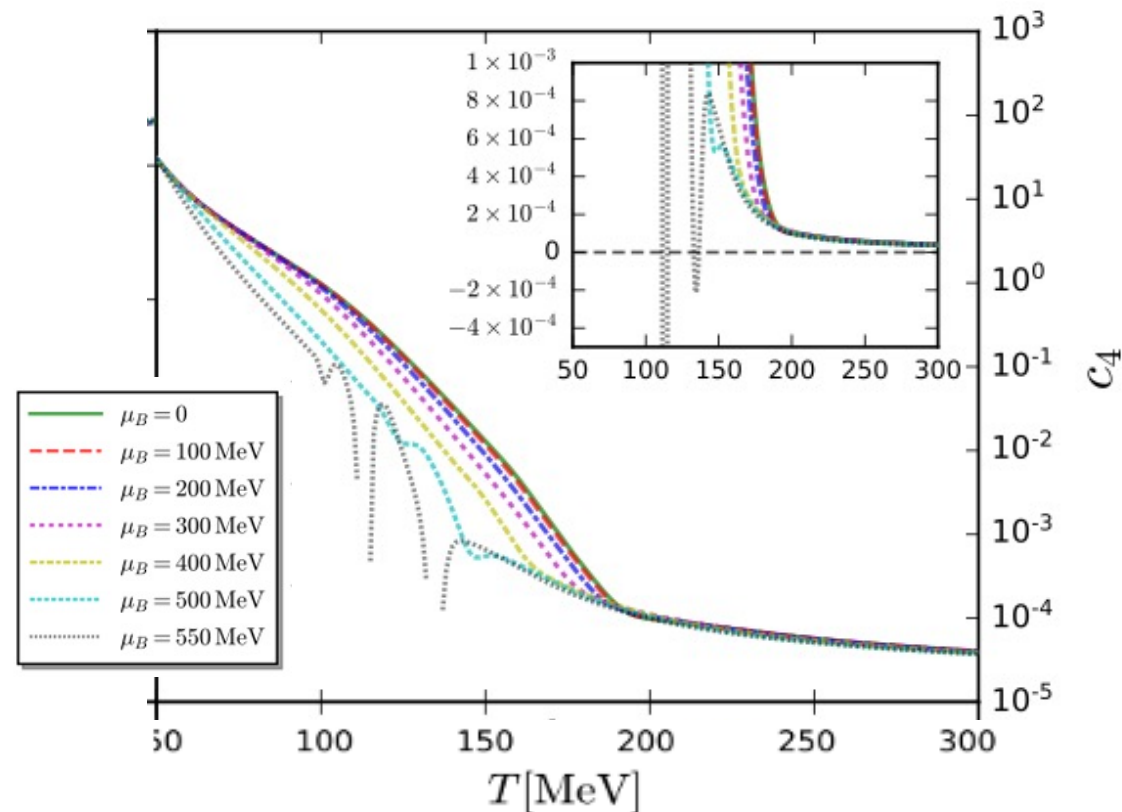
High-order fluctuations of temperature in fRG approach

2+1 flavor LEFT within the fRG approach

$$c_2 = \frac{1}{\chi_2}, \quad c_3 = -\frac{\chi_3}{\chi_2^3}, \quad c_4 = 3\frac{\chi_3^2}{\chi_2^5} - \frac{\chi_4}{\chi_2^4}$$



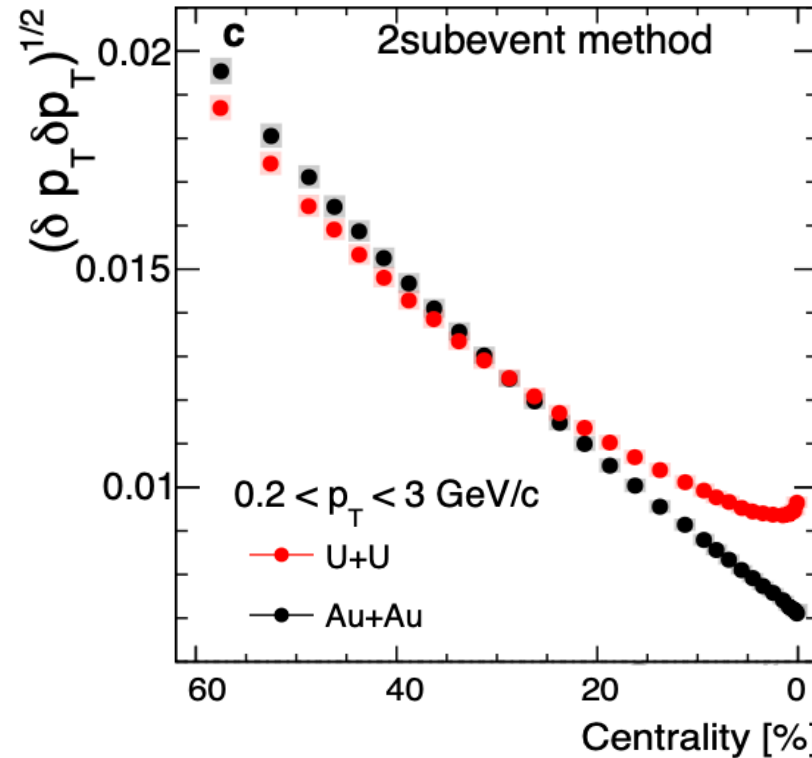
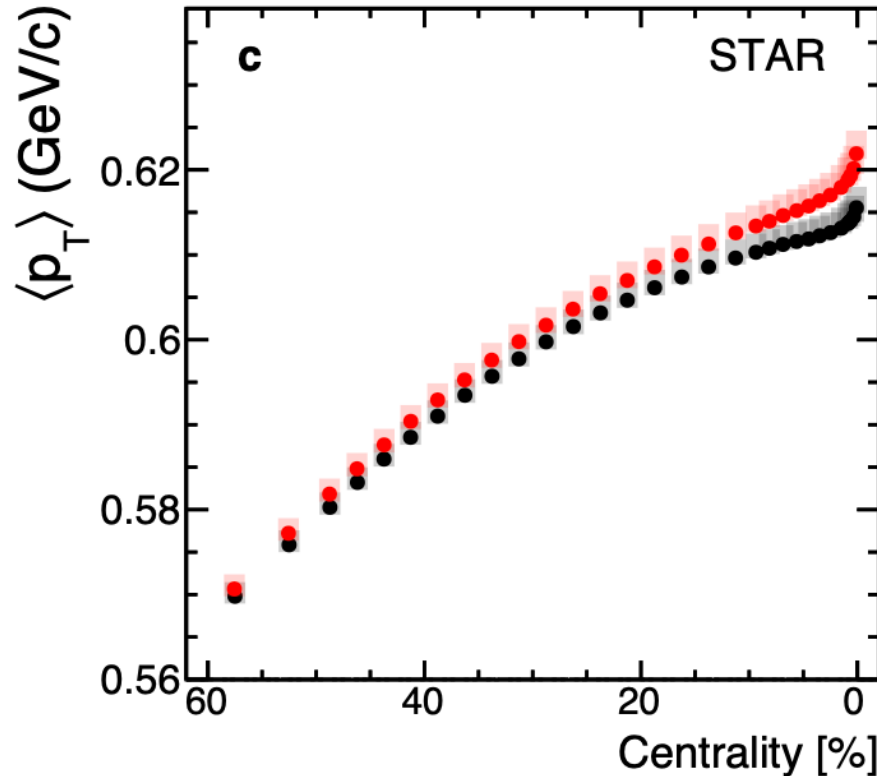
- Indicating a significant suppression of temperature fluctuations in QGP phase compared to those in HRG phase.



- The suppression is more remarkable for high-order temperature fluctuations

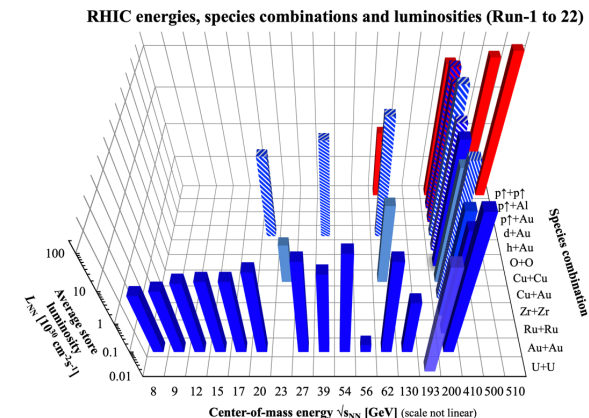
$[p_T]$ fluctuations as another novel tool to image nuclear structure

STAR, Reports on Progress in Physics: Original Research (accepted)



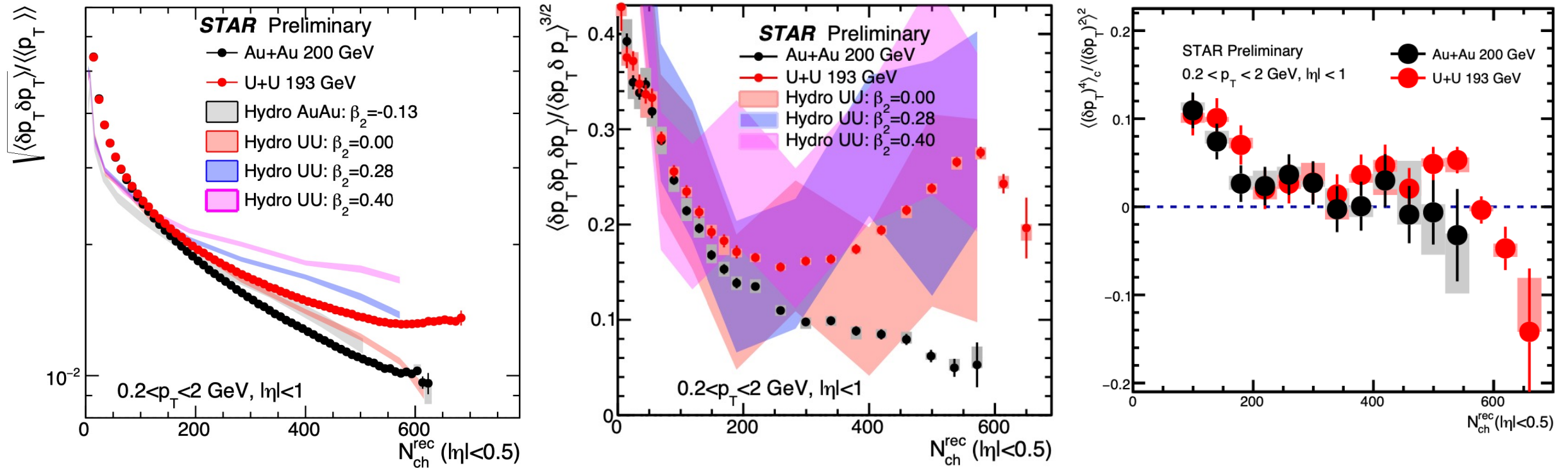
large enhancement in mean and variance $\rightarrow$ size fluctuations enhanced

$[p_T]$ fluctuations also serve as a good observable to image nuclear deformation.



$[p_T]$ fluctuations as another novel tool to image nuclear structure

C. Zhang, IJMPA 32, 2341001 (2023)



Au+Au: variance and skewness follow independent source scaling $1/N_s^{n-1}$ within power-law decrease

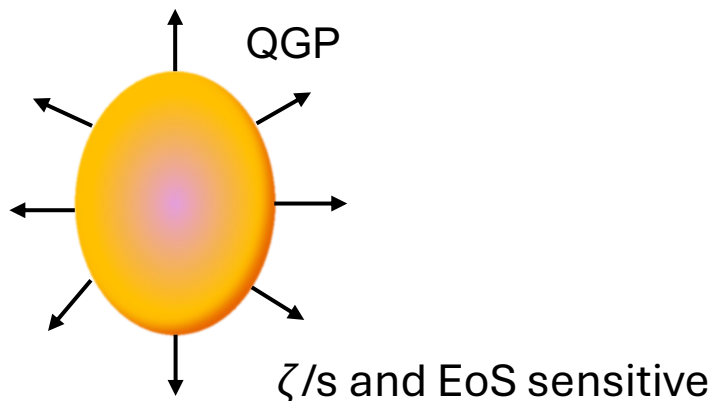
U+U: large enhancement in normalized variance and skewness and sign-change in normalized kurtosis
→ size fluctuations enhanced

The nuclear deformation role is further confirmed by hydro calculations.

$[p_T]$ fluctuations also serve as a good observable to image nuclear deformation.

Radial flow $v_0(p_T)$ measurement in 200 GeV

Radial expansion: the $n=0$ modulation



Spectra slope quantified by $[p_T]$ in each event:
Fluctuation in $[p_T]$ $\rightarrow$ Fluctuation in spectra.

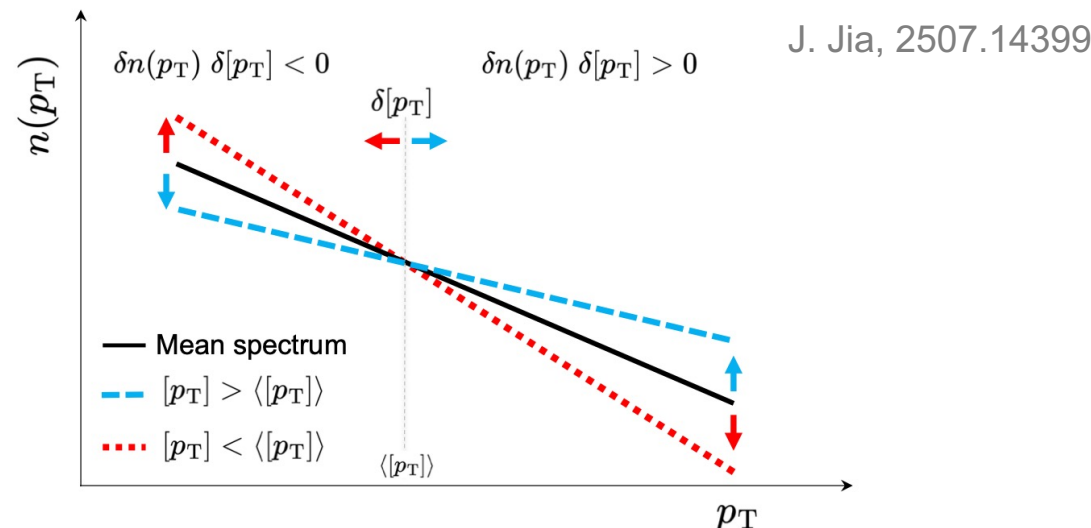
Quantify the correlation induced by radial flow using
covariance between $n(p_T)$ and $[p_T]$: $\langle \delta n(p_T) \delta [p_T] \rangle$

B. Schenke, C. Shen, D. Teaney, PRC 102, 034905 (2020)

T. Parida, R. Samanta, J.Y. Ollitrault, PLB 857, 138985 (2024)

S. A. Jahan, H. Roch, C. Shen. arXiv: 2507.11394

L. Du, 2508.07184



J. Jia, 2507.14399

Calculated within a reference range, p_T -independent.

$$v_0(p_T) = \frac{\langle \delta n(p_T) \delta [p_T] \rangle}{\langle n(p_T) \rangle \langle [p_T] \rangle} \times \frac{1}{v_{0,int}}$$

Pearson Correlations between local possibility of particle yield and reference mean p_T from spectra

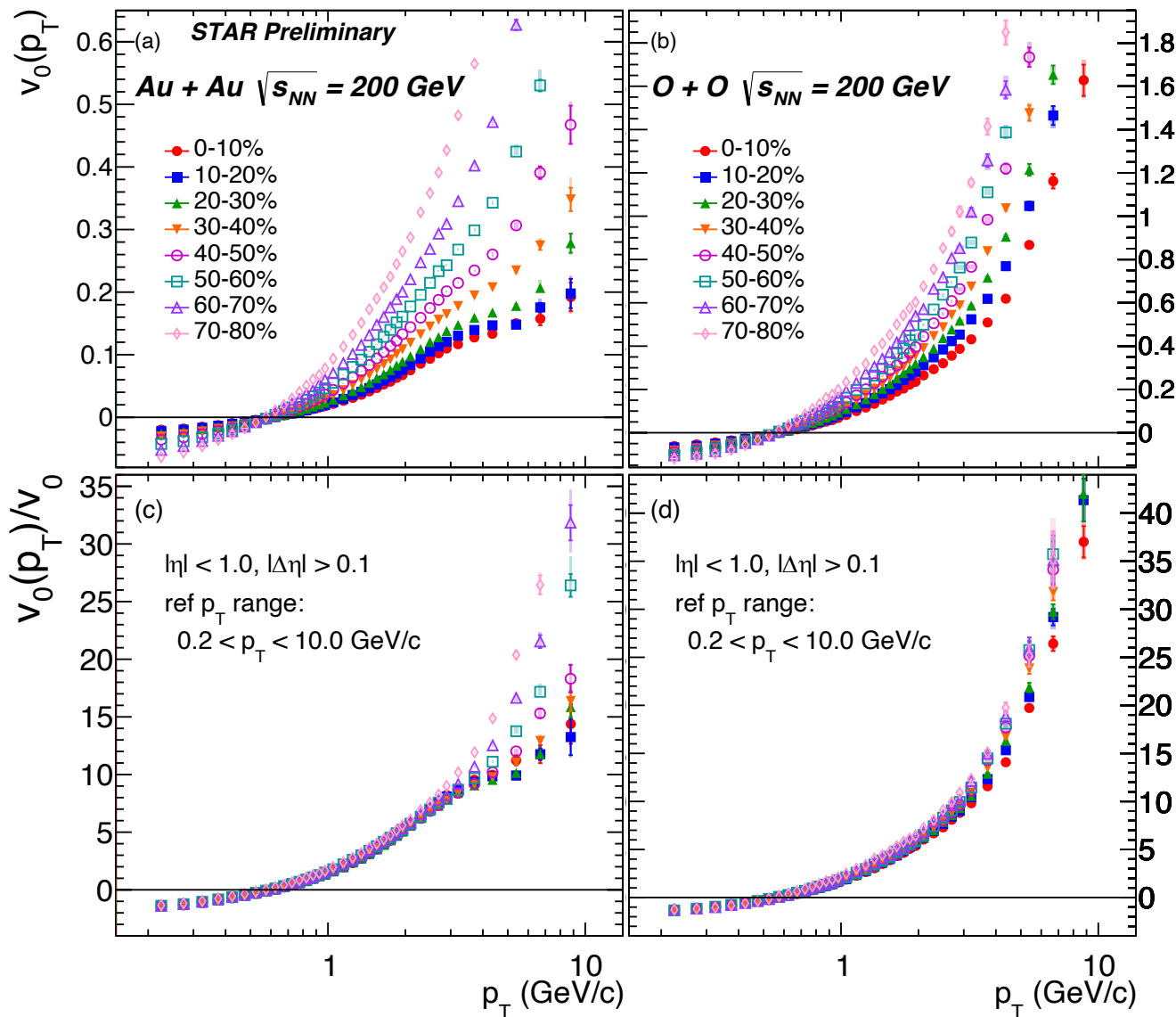
Remove influence from global fluctuations

$$v_{0,int} \equiv \frac{\sqrt{\langle (\delta [p_T])^2 \rangle}}{\langle [p_T] \rangle}$$

This allows us to verify the signatures of radial flow, like for anisotropic flow ($n \geq 1$)

Au+ Au vs. O+O

Zaining Wang (Fudan & Stony Brook), Initial Stages 2025



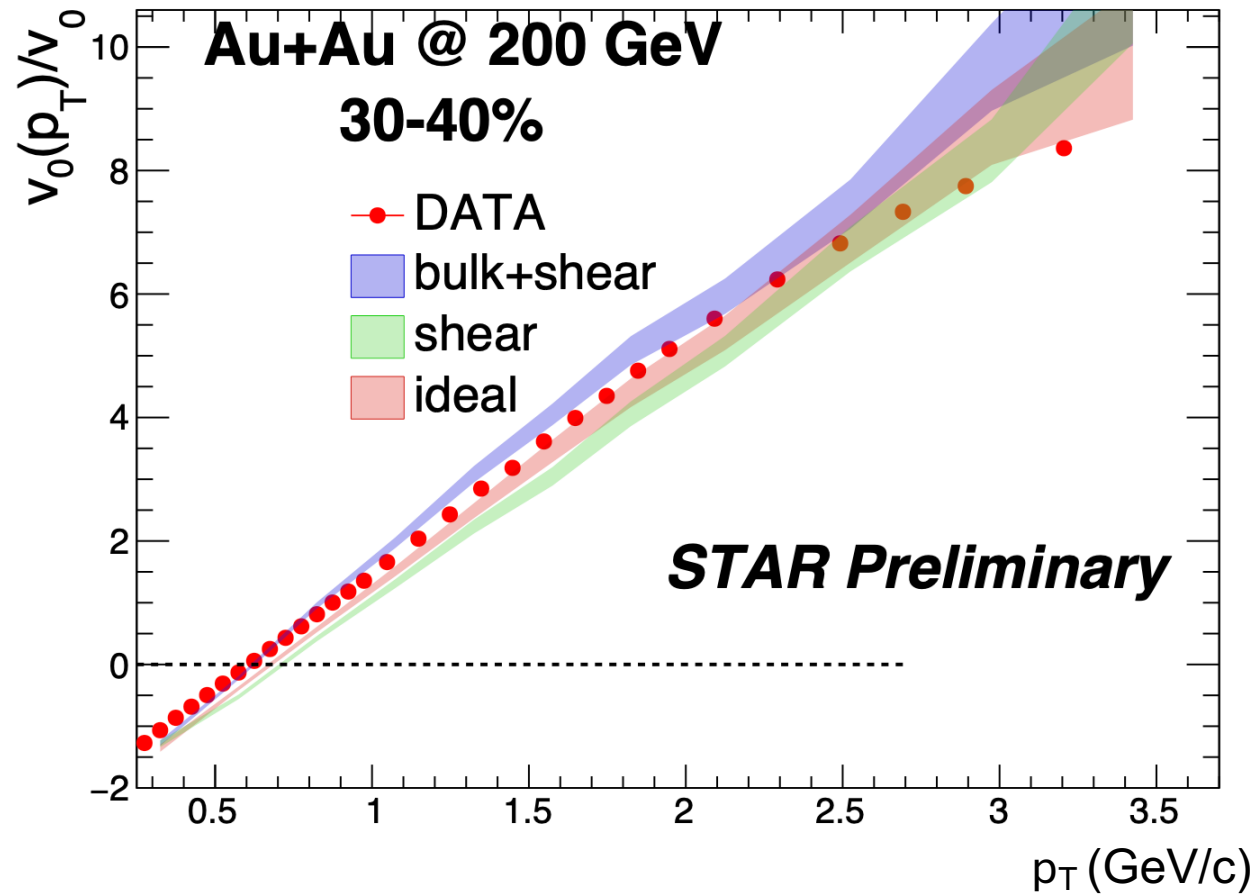
First measurement at RHIC energy

Increase with p_T

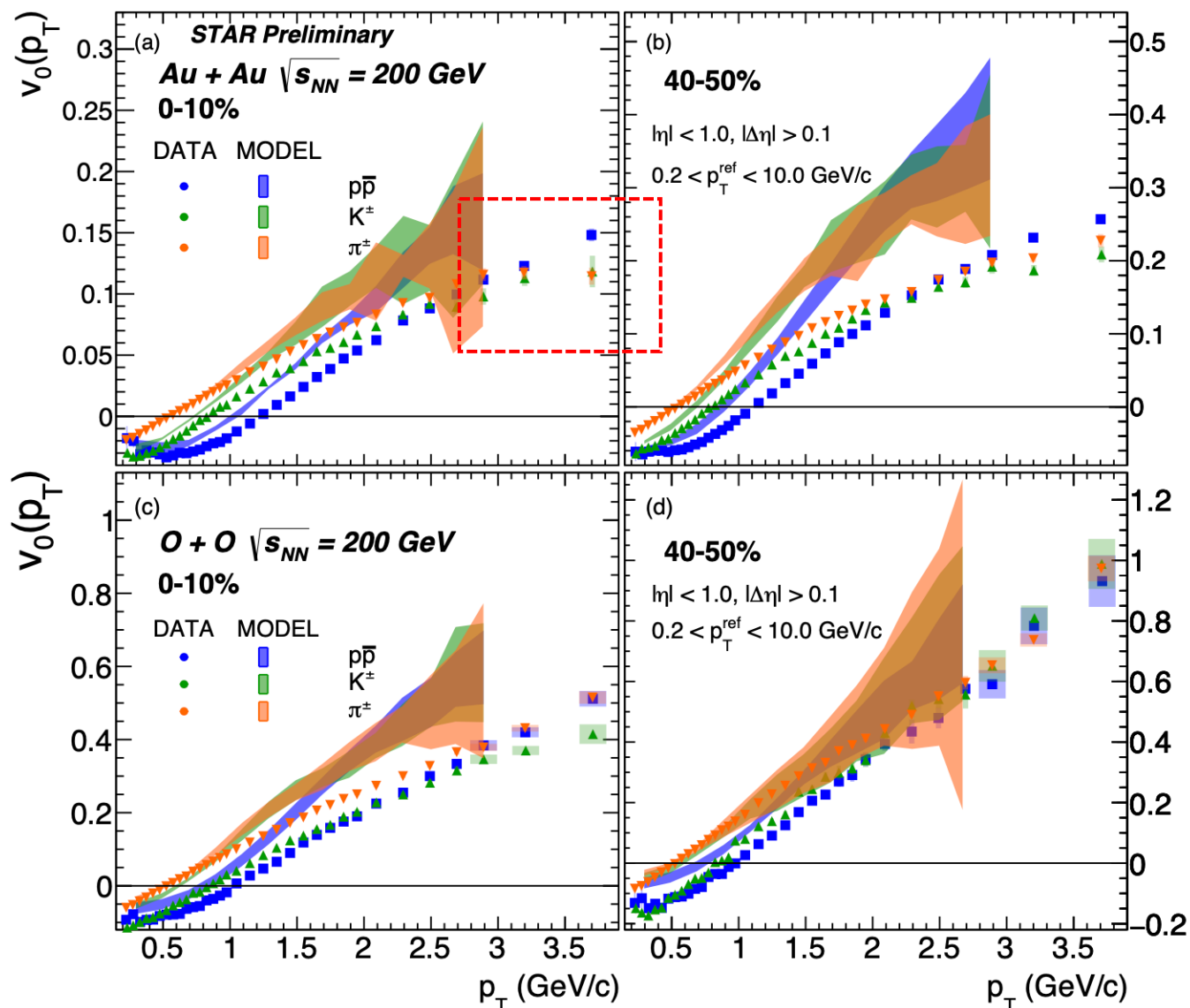
Exhibiting a strong centrality dependence.
The zero-crossing point coincides with $\langle [p_T] \rangle$,
where fluctuations vanish

After normalization, all centralities collapse onto a
single trend, indicating a hydrodynamic collective
response to global fluctuations

$$v_0(p_T) = \frac{\langle \delta n(p_T) \delta [p_T]_{p_T^{\text{ref}}} \rangle}{\langle \delta n(p_T) \rangle (\langle \delta [p_T] \rangle v_0)_{p_T^{\text{ref}}}}$$



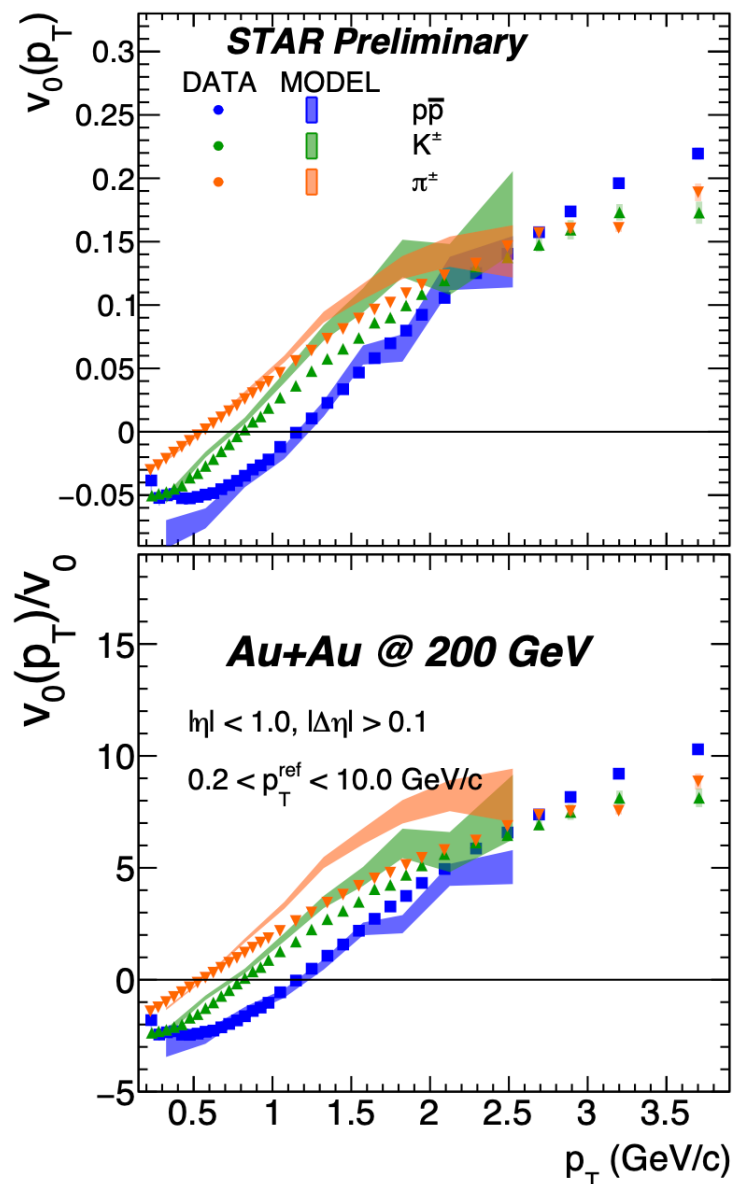
The observable shows sensitivities to bulk viscosity



$$v_0(p_T) = \frac{\langle \delta n(p_T) \delta [p_T]_{p_T^{\text{ref}}} \rangle}{\langle \delta n(p_T) \rangle (\langle \delta [p_T] \rangle v_0)_{p_T^{\text{ref}}}}$$

The stronger boost experienced by heavier particles gives rise to a clear mass ordering

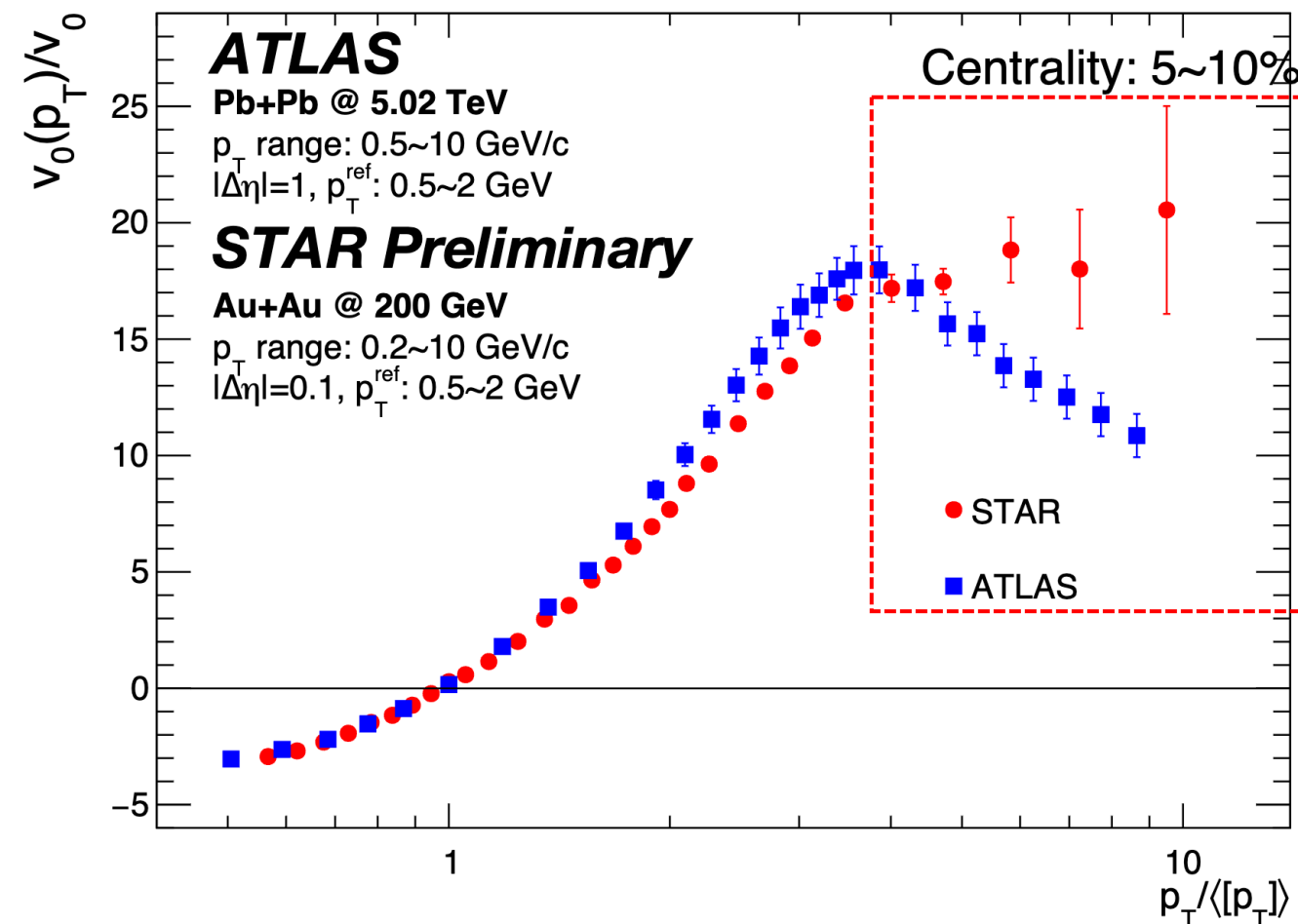
3D-Glauber+MUSIC+UrQMD model also show clear mass ordering, but quantitative differences compared to data



$$v_0(p_T) = \frac{\langle \delta n(p_T) \delta [p_T]_{p_T^{\text{ref}}} \rangle}{\langle \delta n(p_T) \rangle (\langle \delta [p_T] \rangle v_0)_{p_T^{\text{ref}}}}$$

TRENTTo+free-streaming+MUSIC+iSS+SMASH seems more comparable with data.

More model calculations are needed to tune for fitting this new observables at RHIC energies.



Future system- and energy-scans are interesting and useful for understanding radial flow

Conclusions and Outlooks

1. **Observation of non-monotonic dependence** of mean p_T fluctuations at high baryon density, **negative kurtosis** in highly-deformed U^{238} .
2. **Negative kurtosis** in highly deformed U^{238} indicates the nuclear structure effect.
3. New measurements to help narrow down the nature of radial flow and temperature fluctuation, including the **QGP collectivity behavior, initial conditions, and nuclear structure imaging**.

