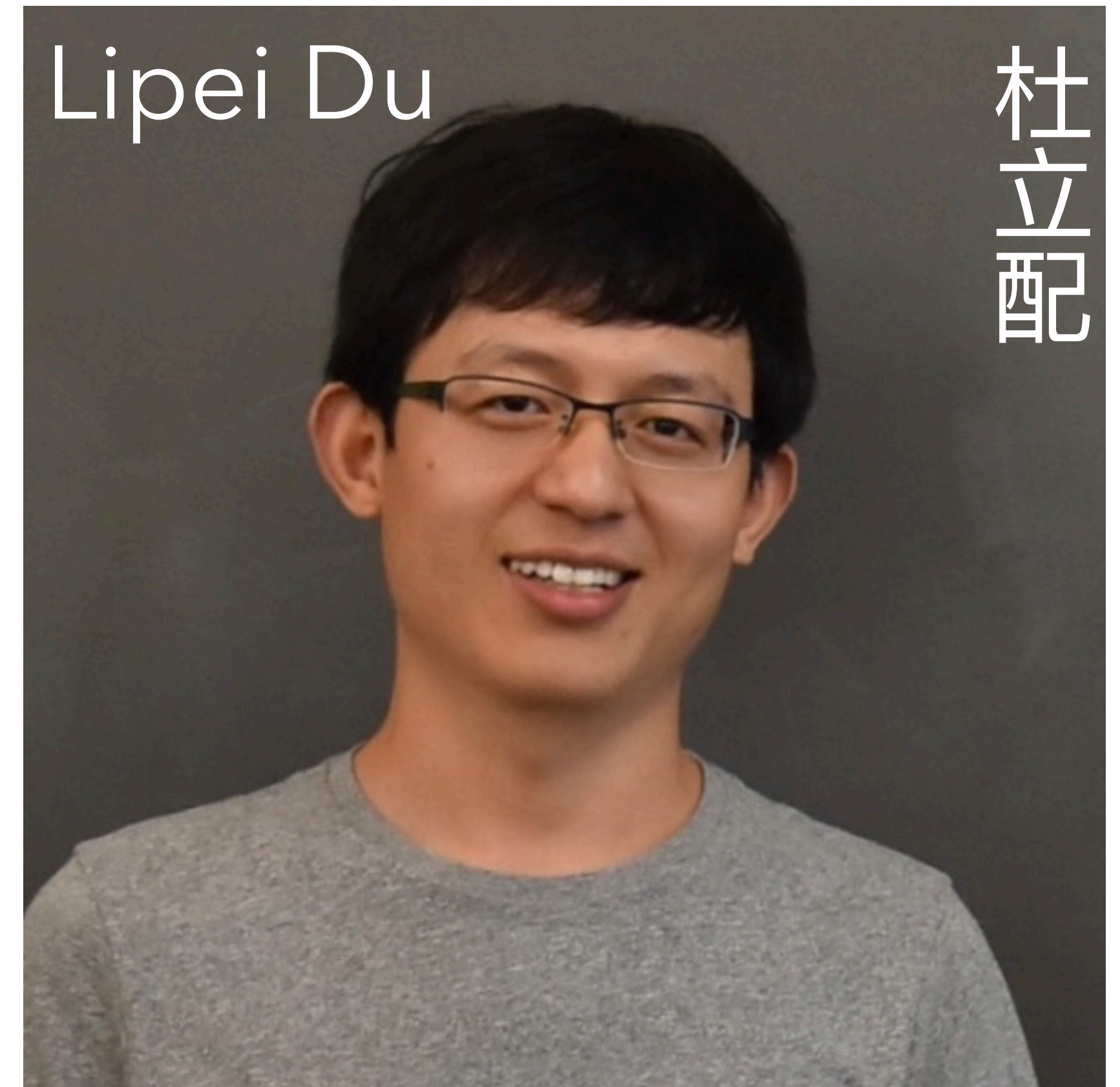




Jianing Li

李嘉宁



Lipei Du

杜立配

Correlations in an inhomogeneous medium

Shuzhe Shi (施舒哲)

Tsinghua University

w/Jianing Li & Lipei Du, PhysRevC.109.034906 + in preparation

$$\kappa_2 \equiv \langle (\delta N)^2 \rangle = \frac{T}{V} \xi^2,$$

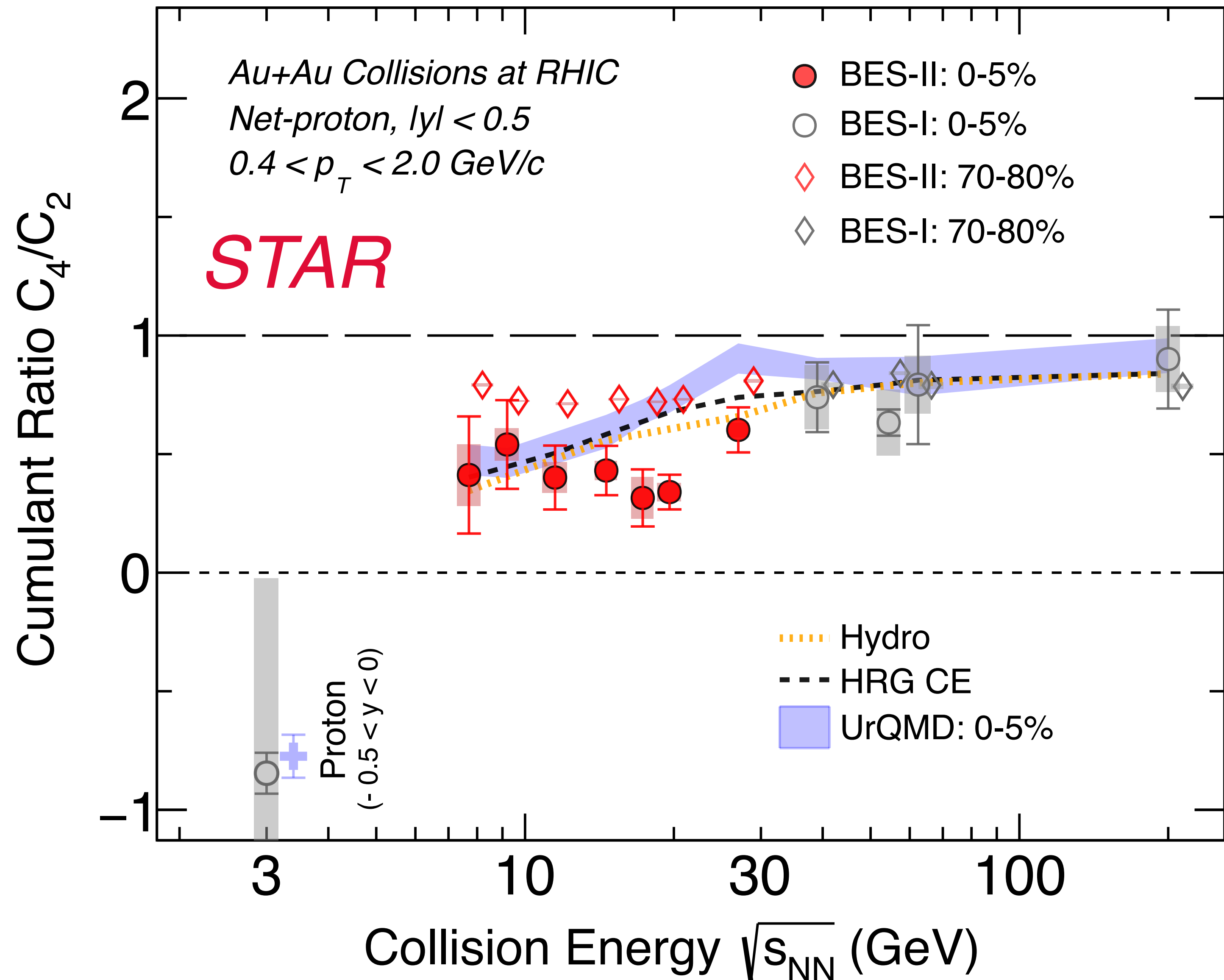
$$\kappa_4 \equiv \langle (\delta N)^4 \rangle - 3(\kappa_2)^2 = -6\tilde{\lambda}_4 \frac{T}{V} \xi^7.$$

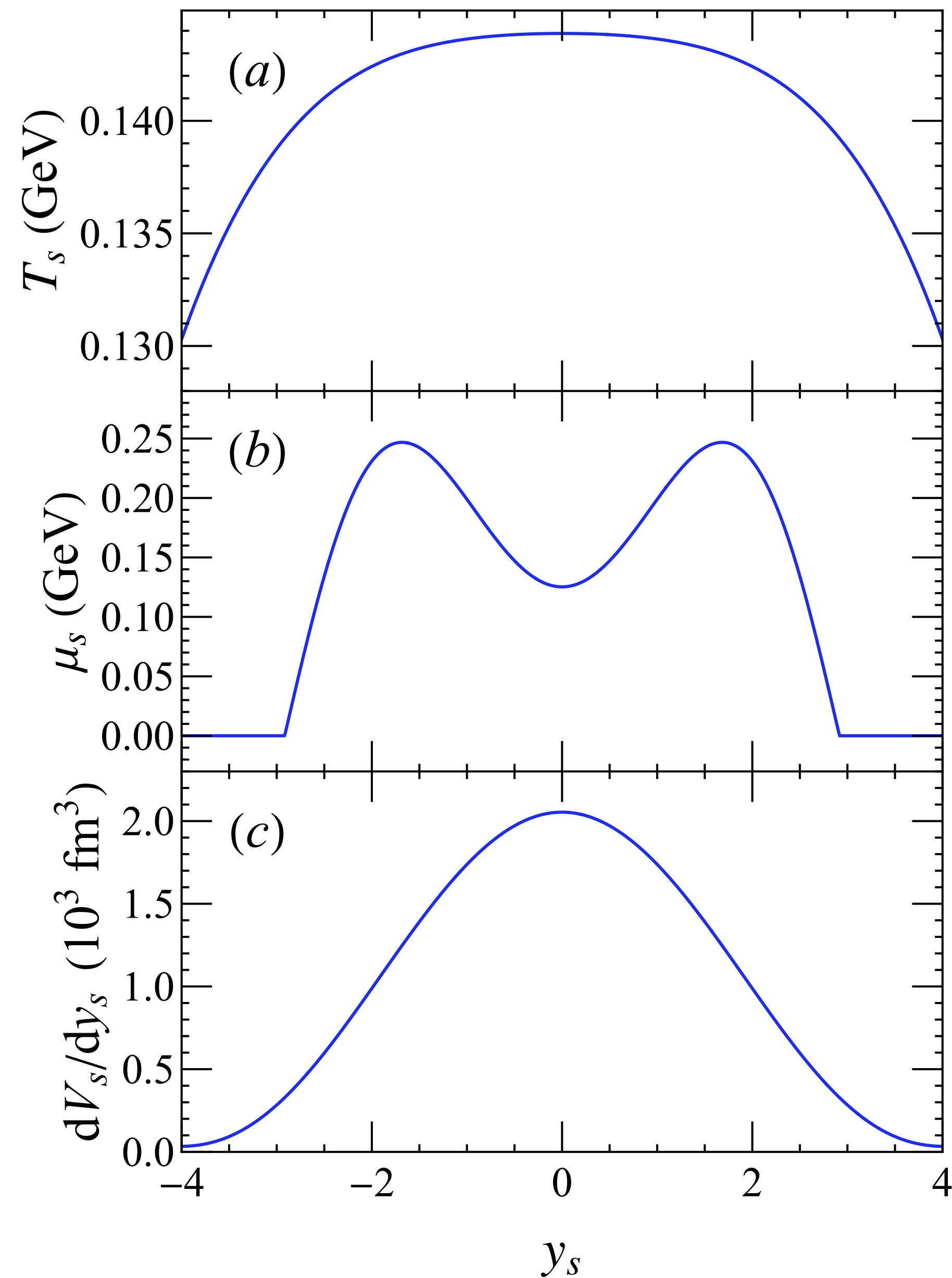
M. Stephanov,

Phys.Rev.Lett. 102 (2009) 032301

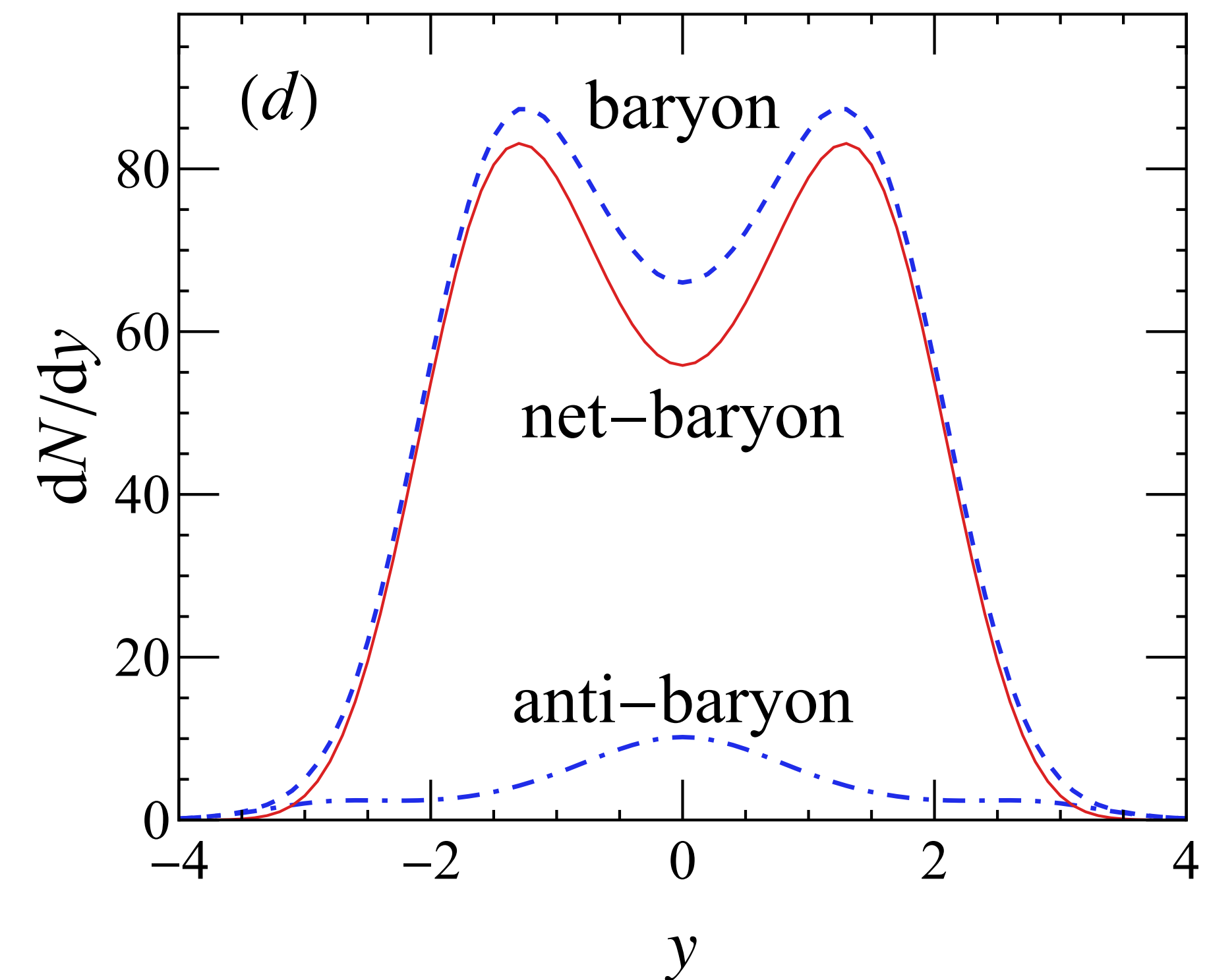
Presumptions:

1. *homogeneous*
2. *open thermal system*





*inhomogeneous rapidity-dependent hydro
needed to understand experiment*



L. Du, H. Gao, S. Jeon, C. Gale, PhysRevC.109.014907

J. Li, L. Du, SS, PhysRevC.109.034906



$$P\left(B_D\right)=\sum_{N_D, \bar{N}_D, N_{\bar{D}}, \bar{N}_{\bar{D}}=0}^{\infty}\left(\delta_{B_D, N_D-\bar{N}_D} P_A\left(N_D\right) P_A\left(\bar{N}_D\right)\right) \times\left(\delta_{B_{\bar{D}}, N_{\bar{D}}-\bar{N}_{\bar{D}}} P_{\bar{D}}\left(N_{\bar{D}}\right) P_{\bar{D}}\left(\bar{N}_{\bar{D}}\right)\right) \times \delta_{B, B_D+B_{\bar{D}}}.$$

conservation effect

D : detected,

$\bar{D}$: not detected;

N : baryon,

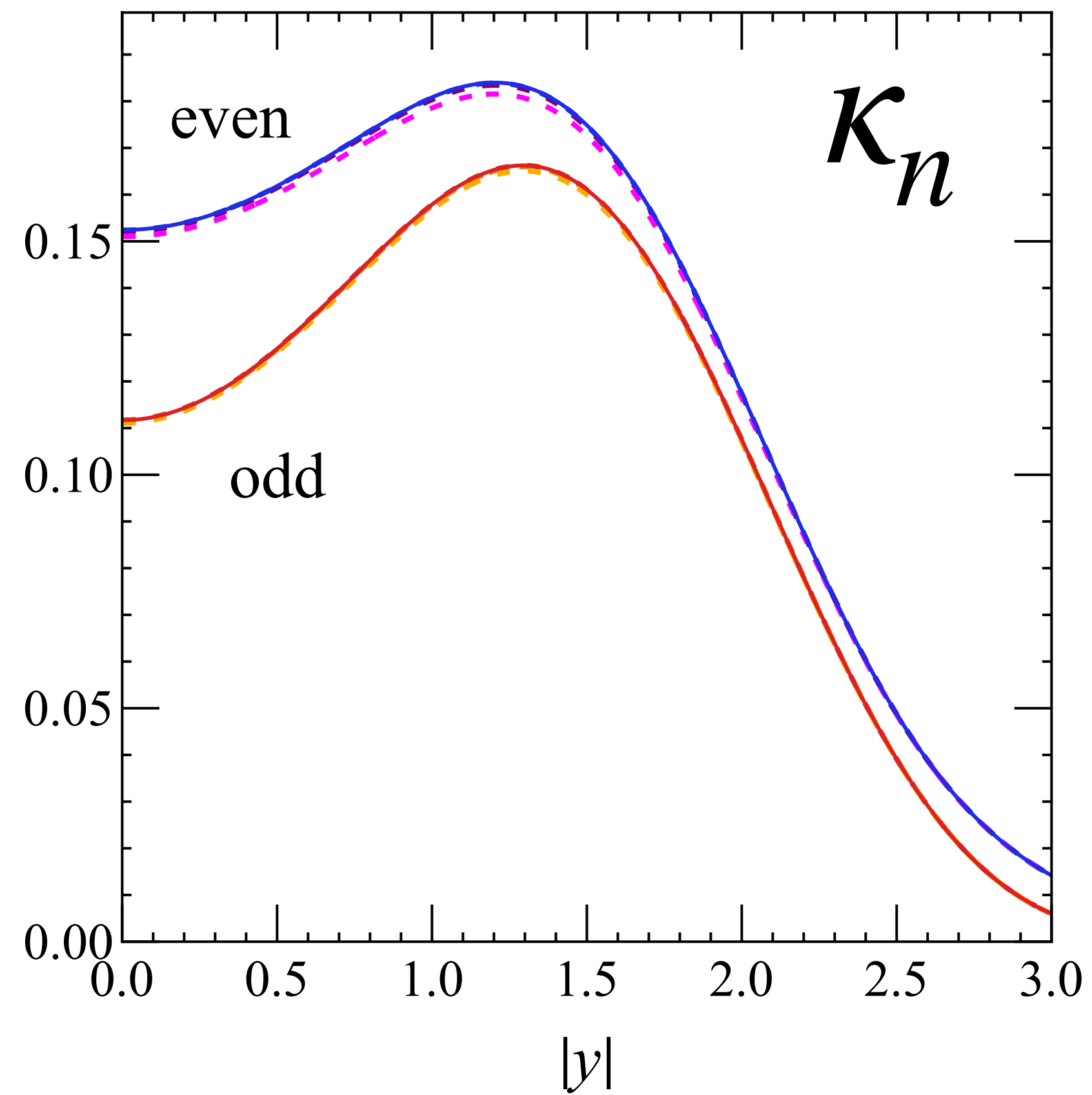
$\bar{N}$: anti – baryon;

J. Li, L. Du, SS, PhysRevC.109.034906

see also: Vovchenko, Koch, PhysRevC.103.044903

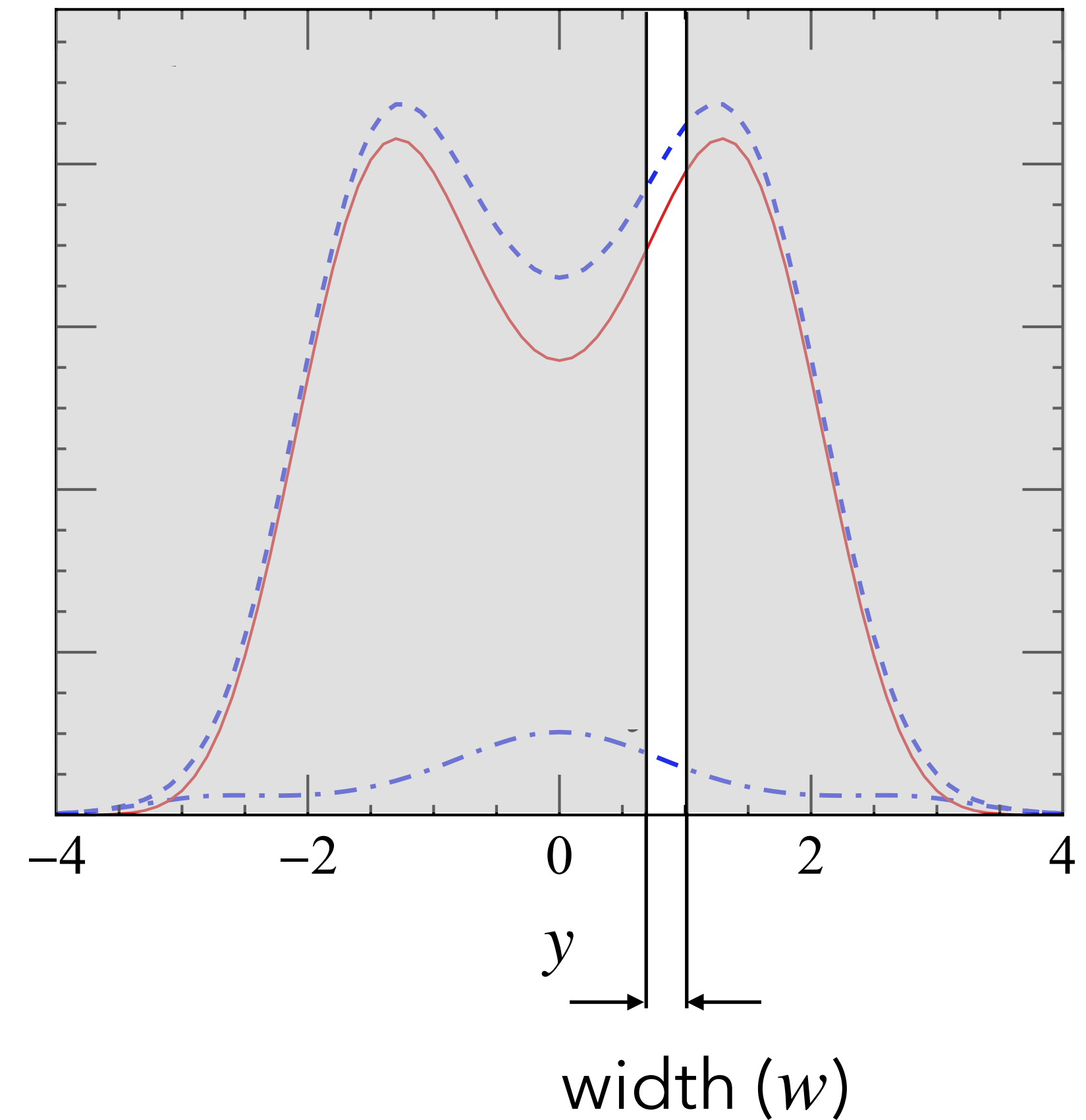
Braun-Munzinger, Friman, Karsch, Redlich, Skokov, PhysRevC.84.064911

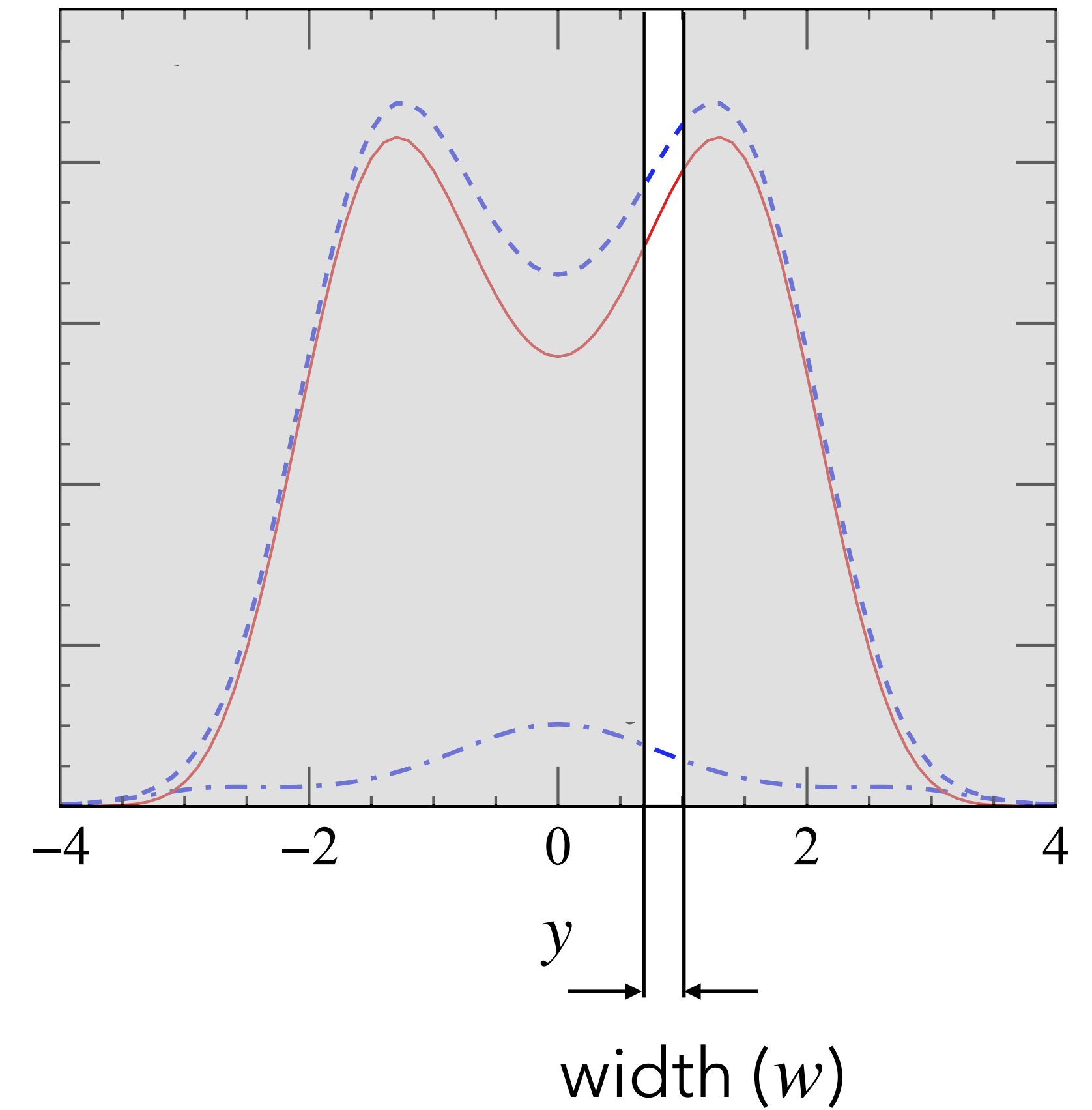
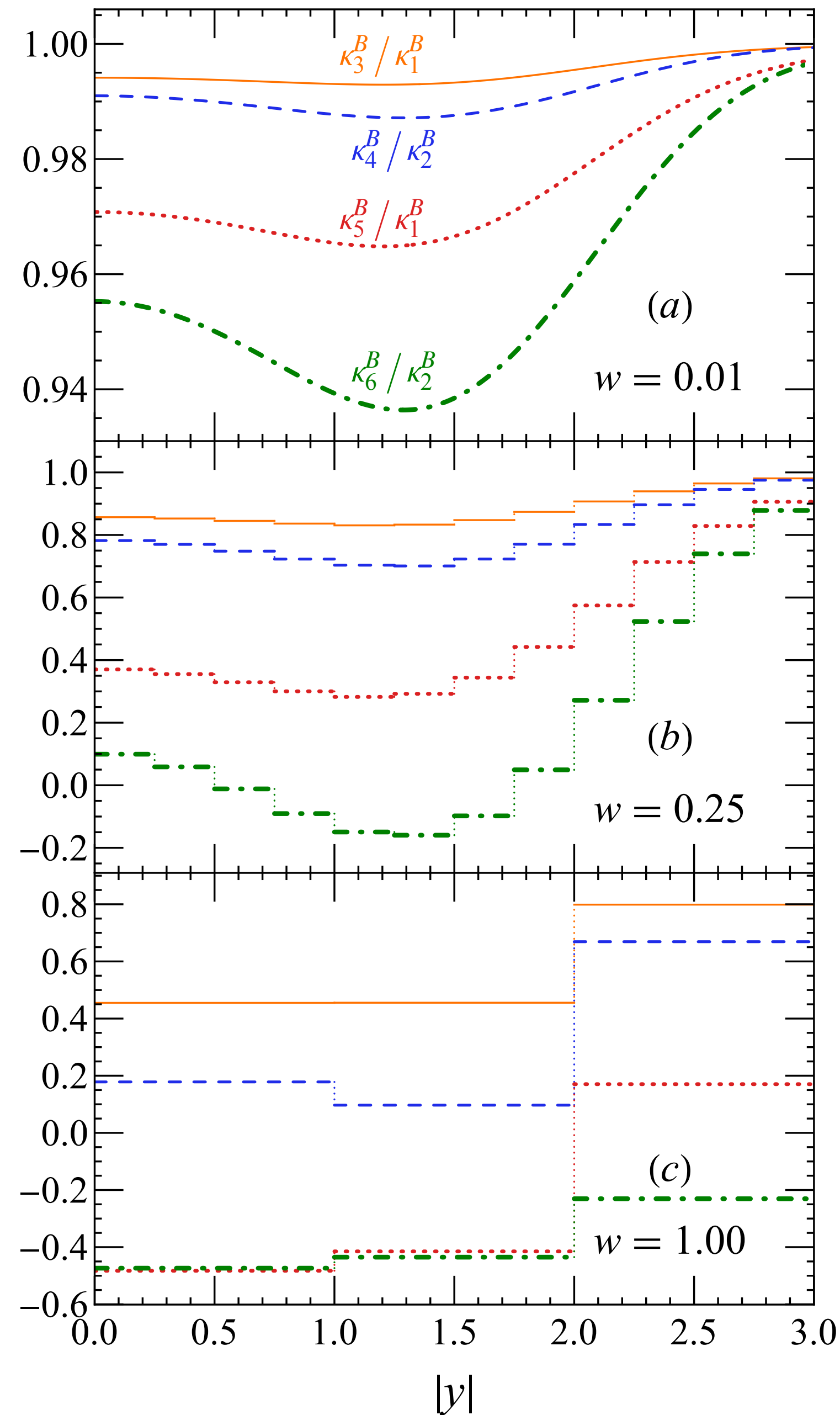
...



solid: open thermal system

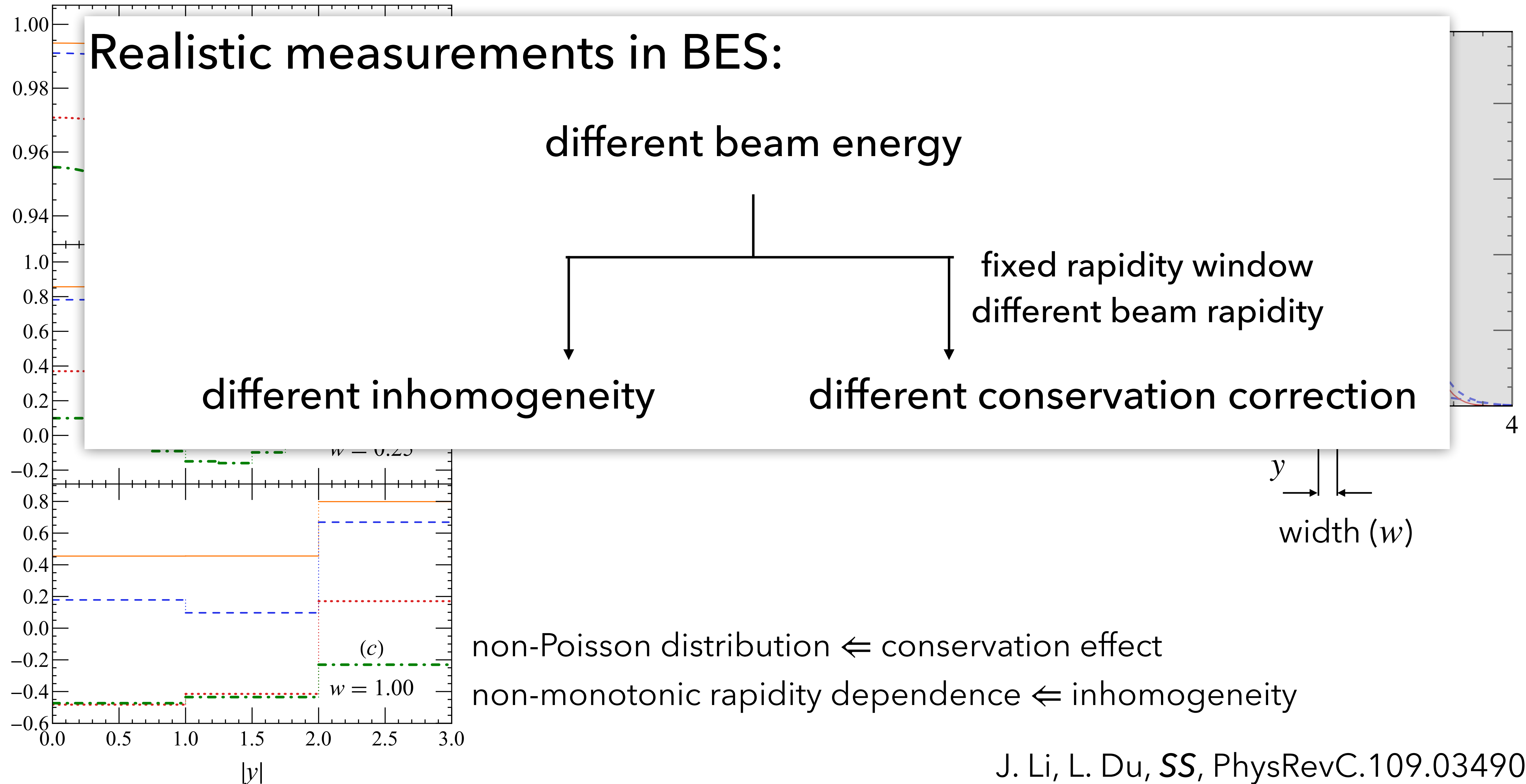
dashed: closed thermal system with small width ($w = 0.002$)

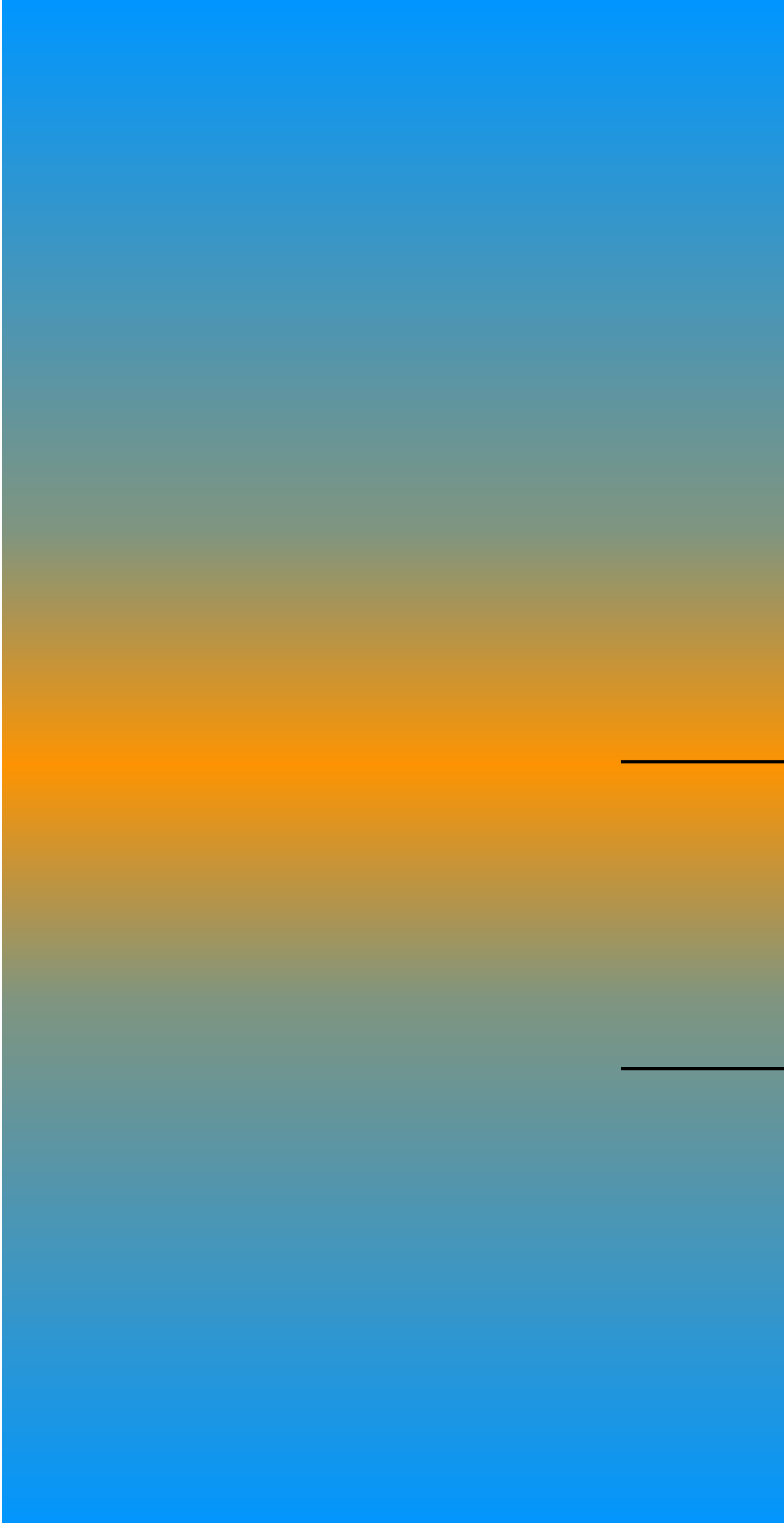




non-Poisson distribution $\Leftarrow$ conservation effect

non-monotonic rapidity dependence $\Leftarrow$ inhomogeneity





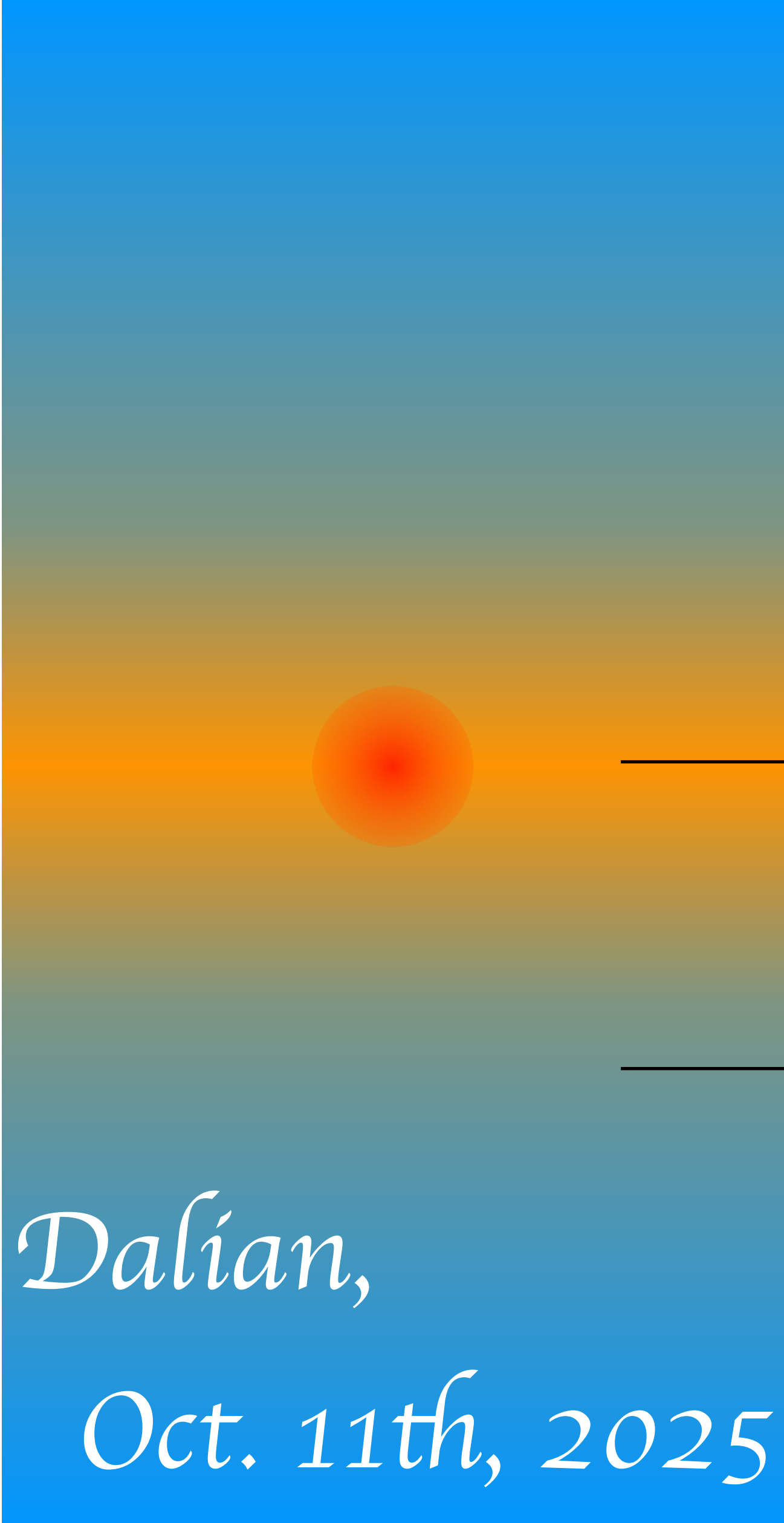
far from CP

near CP

$$\ell_{\text{inhomo}} \lesssim \xi_{\text{therm}}$$

far from CP

near CP


$$\ell_{\text{inhomo}} \lesssim \xi_{\text{therm}}$$

Dalian,

Oct. 11th, 2025

far from CP

Review: Criticality in a *homogeneous* system

Hadrons coupled with a scalar critical field

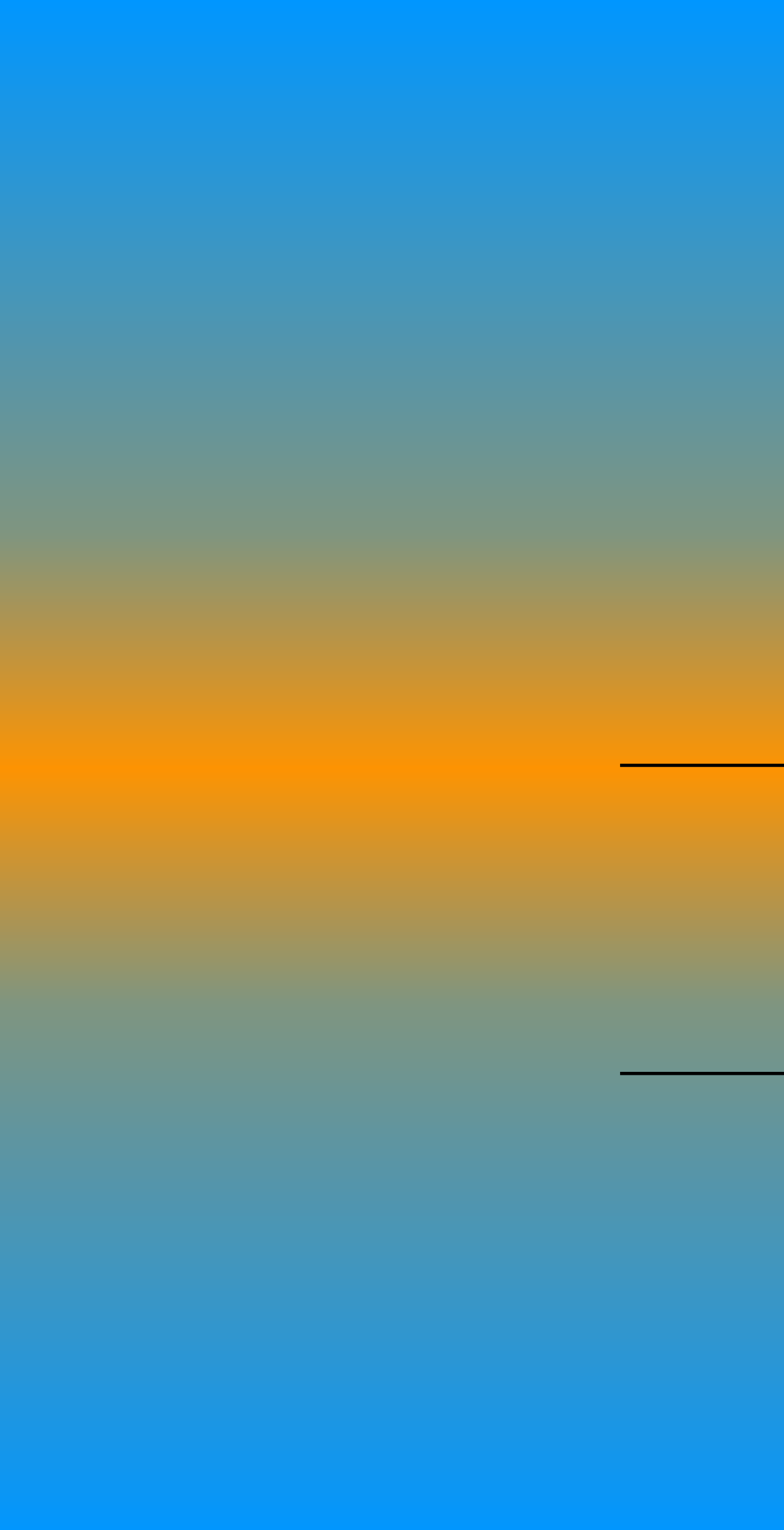
Free energy density:

$$\Omega[\sigma(\mathbf{x})] = -\frac{1}{2}\sigma(\mathbf{x}) \nabla^2 \sigma(\mathbf{x}) + \frac{1}{2} m^2 \sigma^2(\mathbf{x}) .$$

$m = 1/\xi(T, \mu)$

near CP

$\ell_{\text{inhomo}} \lesssim \xi_{\text{therm}}$



far from CP

near CP

$$\ell_{\text{inhomo}} \lesssim \xi_{\text{therm}}$$

Review: Criticality in a *homogeneous* system

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Density matrix:

$$\mathcal{W}[\sigma(\mathbf{x})] = \frac{e^{-\int \beta \Omega[\sigma(\mathbf{x})] d^3x}}{\mathcal{Z}} ,$$

Two-point correlation:

$$\langle \sigma(\mathbf{x}_1) \sigma(\mathbf{x}_2) \rangle = \int \mathcal{D}\sigma \mathcal{W}[\sigma(\mathbf{x})] \sigma(\mathbf{x}_1) \sigma(\mathbf{x}_2) .$$

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$$m = 1/\xi(T, \mu)$$

$$\int \Omega[\sigma(\mathbf{x})] d^3x = \int (p^2 + m^2) \tilde{\sigma}^*(\mathbf{p}) \tilde{\sigma}(\mathbf{p}) \frac{d^3p}{(2\pi)^3} ,$$
$$\sigma(\mathbf{x}) = \int \tilde{\sigma}(\mathbf{p}) e^{i\mathbf{p} \cdot \mathbf{x}} \frac{d^3p}{(2\pi)^3}$$

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$$\langle \sigma(\mathbf{x}_1) \sigma(\mathbf{x}_2) \rangle = \frac{T e^{-m |\mathbf{x}_1 - \mathbf{x}_2|}}{4\pi |\mathbf{x}_1 - \mathbf{x}_2|}.$$

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Criticality in an *inhomogeneous* system

Hadrons coupled with a scalar critical field

Free energy density:

$$\Omega[\sigma(\mathbf{x})] = -\frac{1}{2}\sigma(\mathbf{x})\nabla^2\sigma(\mathbf{x}) + \frac{1}{2}m^2(\mathbf{x})\sigma^2(\mathbf{x}).$$

$$m = 1/\xi(T, \mu)$$

$$\langle\sigma(\mathbf{x}_1)\sigma(\mathbf{x}_2)\rangle = \frac{T e^{-m|\mathbf{x}_1-\mathbf{x}_2|}}{4\pi|\mathbf{x}_1-\mathbf{x}_2|}.$$

$$\int \Omega[\sigma(\mathbf{x})]d^3x = \int (p^2 + m^2)\tilde{\sigma}^*(\mathbf{p})\tilde{\sigma}(\mathbf{p})\frac{d^3p}{(2\pi)^3},$$

$$\sigma(\mathbf{x}) = \int \tilde{\sigma}(\mathbf{p})e^{i\mathbf{p}\cdot\mathbf{x}}\frac{d^3p}{(2\pi)^3}$$

Density matrix:

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Two-point correlation:

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Criticality in an *inhomogeneous* system

Hadrons coupled with a scalar critical field

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Two-point correlation:

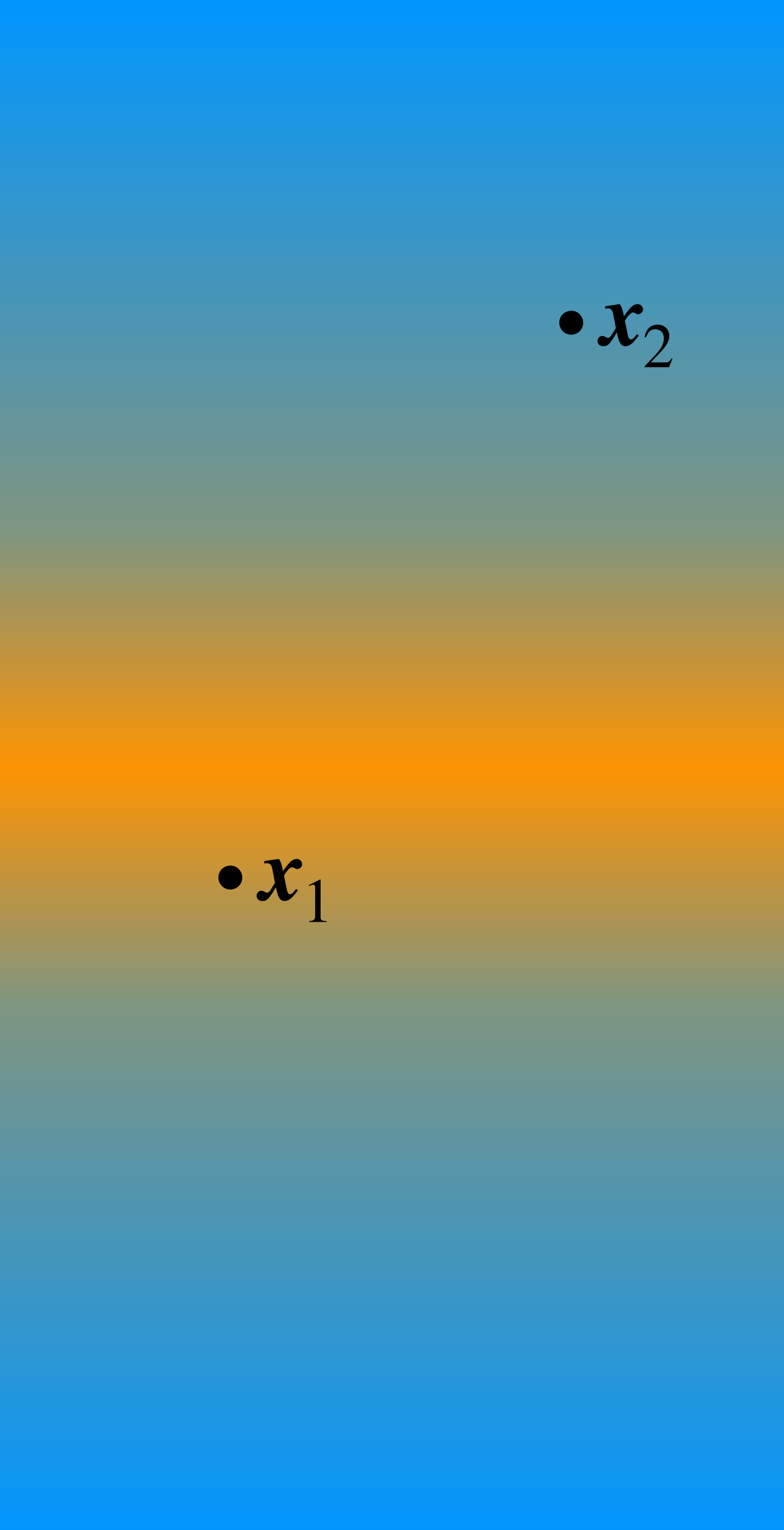
$$\langle \sigma(\mathbf{x}_1)\sigma(\mathbf{x}_2) \rangle = \int \mathcal{D}\sigma \mathcal{W}[\sigma(\mathbf{x})] \sigma(\mathbf{x}_1)\sigma(\mathbf{x}_2).$$

$$\left(-\nabla^2 - \frac{\nabla\beta}{\beta} \cdot \nabla + m^2(\mathbf{x}) - \frac{\nabla^2\beta(\mathbf{x})}{2\beta(\mathbf{x})} \right) \psi_n(\mathbf{x}) = \lambda_n \psi_n(\mathbf{x})$$

$$\langle \sigma(\mathbf{x}_1)\sigma(\mathbf{x}_2) \rangle = \sum_n \frac{\psi_n(\mathbf{x}_1)\psi_n(\mathbf{x}_2)}{\lambda_n}.$$

$$\int \beta(\mathbf{x})\Omega[\sigma(\mathbf{x})]d^3x = \sum_n \frac{1}{2}\lambda_n s_n^2,$$

$$\sigma(\mathbf{x}) = \sum_n s_n \psi_n(\mathbf{x})$$



• x_2

• x_1

$$m^2(\mathbf{x}) = m_0^2 + m_0^2 \omega^2 z^2, \quad \beta(\mathbf{x}) = \beta_0.$$

Eigenfunctions can be solved analytically:

$$\psi_{n,+}(\mathbf{x} | \mathbf{p}_\perp) = \cos(\mathbf{p}_\perp \cdot \mathbf{x}_\perp) \phi_n(z),$$

$$\psi_{n,-}(\mathbf{x} | \mathbf{p}_\perp) = \sin(\mathbf{p}_\perp \cdot \mathbf{x}_\perp) \phi_n(z),$$

$$\lambda_n(\mathbf{p}_\perp) = \frac{(2n+1)m_0\omega + m_0^2 + |\mathbf{p}_\perp|^2}{T},$$

• $\mathbf{x}_2$

• $\mathbf{x}_1$

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$$\begin{aligned} & \langle \sigma(\mathbf{x}_1) \sigma(\mathbf{x}_2) \rangle \\ &= \frac{T}{4\pi |\mathbf{x}_1 - \mathbf{x}_2|} e^{-m_0 \left(1 + \omega^2 \frac{z_1^2 + z_1 z_2 + z_2^2}{6} \right) |\mathbf{x}_1 - \mathbf{x}_2|} + \mathcal{O}(\omega^4). \end{aligned}$$

inhomogeneity correction

• $\mathbf{x}_2$ • $\mathbf{x}_1$

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inhomogeneity correction

reproduces the homogenous result if $\omega \rightarrow 0$

conservation effect → ◎ modifies the baseline w/o criticality

inhomogeneous corrections → ◎ corrects the critical correlation function

- can be calculated (semi)analytically

disentangling finite-size effect from the critical correlations?