

# Background Issues at DAFNE

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# Outline

 $DA\Phi NE$  background and beam lifetime are dominated by Touschek scattering, as typically happens in low energy rings

- · DADNE
- Experimental knobs used to control backgrounds in the detectors
- Simulation Studies
- Comparison between simulation and measurements
- Studies with the crab waist scheme
- First studies for the LER of the SUPERB Factory





## DAFNE peak luminosity for KLOE (red) DEAR (blue) and FINUDA (green)



	KLOE May 04 Nov 05	FINUDA Nov 06 Jun 07
$L_{peak}$ [cm <sup>-2</sup> s <sup>-1</sup> ]	<b>1.53·10</b> <sup>32</sup>	<b>1.6·10</b> <sup>32</sup>
$L^{MAX}_{\int day} [pb^{-1}]$	9.8	9.4
$L^{MAX}_{fmonth} [pb^{-1}]$	209	226
I-MAX [A]	2.2	1.5
I <sup>+MAX</sup> [A]	1.35	1.1
n <sub>bunches</sub>	111	110
$L_{\text{flogged}}$ [fb <sup>-1</sup> ]	2	0.966
$\beta_x^*$ [m]	1.5	2.0
$\beta_v^*$ [m]	0.018	0.019
$\varepsilon_{\rm x}$ [10 <sup>-6</sup> m rad]	0.37	0.37
к (%)	1.5	1.5



**DEAR** (DA $\Phi$ NE Exotic Atoms Research) and **SIDDHARTA** their goal is to study the creation of exotic atoms, where the charged K mesons produced in the  $\Phi$ resonance decays are captured in the electronic shells.

**FINUDA** (Fisica Nucleare a DA $\Phi$ NE) was an experiment dedicated to the study of hypernuclei The detector is embedded in a SC solenoid with the same

field integral as KLOE (2.4 Tm)

# Detectors Sensitivity to Backgrounds

 KLOE suffered from 'high' energy particles (E> 10 MeV) seen in overlap with physics (accidentals)

> also important higher energy products with E>150MeV (endcap trigger threshold)

 $4\pi$  acceptance- difficult shielding

• DEAR suffered from low energy photons (O(100) keV)-

no trigger, but small gas target detector could be shielded by lead all around

• SIDDHARTA is a gas target detector with trigger, shielded by lead

At the beginning of data taking, all these experiments suffered from large background.

## DAFNE machine background and lifetime dominated by Touschek effect:

elastic Coulomb scattering of pairs of particles within a bunch

The two emerging particles have the same momentum deviation  $\epsilon$ : one gains and the other looses it.

They start to oscillate with a betatron amplitude:

$$\mathbf{A} = \varepsilon \sqrt{\mathbf{H}\beta(\mathbf{s})} \qquad \text{where} \qquad \mathbf{H} = \gamma_{\mathbf{x}} \mathbf{D}_{\mathbf{x}}^{2} + 2\mathbf{a}_{\mathbf{x}} \mathbf{D}_{\mathbf{x}} \mathbf{D}_{\mathbf{x}}^{'} + \beta_{\mathbf{x}} \mathbf{D}_{\mathbf{x}}^{'2}$$

Off-momentum particles can exceed the momentum acceptance given by the RF bucket, or may hit the physical aperture when displaced by dispersion.

A betatron oscillation is excited if the momentum change happens in a dispersive region.

# SCALING of Touschek loss rate dN/dt and lifetime $\tau$ with beam parameters

The Touschek part. loss  $\dot{N}$  rate is approximately

$$\propto \frac{N^2}{\gamma^3 \epsilon^2 V}$$

N particles/bunch

V bunch volume

 $\boldsymbol{\epsilon}$  momentum acceptance

1 dN

 $\tau$  N dt

Touschek effect is determined by momentum acceptance and bunch density integrated over the lattice structure.

$$\begin{array}{ll} \mbox{Lifetime} & \tau \propto \frac{\sigma_x \, \sigma_y \, \sigma_z}{I} \\ & \tau \propto I^{-2/3} \end{array}$$

where  $\sigma_z \propto I^{1/3}$ 

 $dN/dt \propto I/\tau \propto I^{5/3}$ 

$$\frac{\mathrm{dN}}{\mathrm{dt}} \propto 1 / \sqrt{\kappa} \qquad \kappa = \varepsilon_{\mathrm{y}} / \varepsilon_{\mathrm{x}}$$

# Background Handling

- o Tracking studies/measurements and comparisons have been performed to reduce these backgrounds rates
- o Additional collimators inserted,

shape of existing and new ones optimized

- o Low- $\beta$  quads in KLOE IR: from triplet to doublet
- o Shielding implemented between pipe and KLOE low- $\beta$  quads, lead around gas target detectors
- o Optics Adjustments:
  - orbit optimization,
  - Sextupoles Optimization
  - Octupoles Optimization
  - ·Improved linear and non-linear knowledge of the machine
  - •Increased Dynamic aperture

with better  $\beta \textbf{s}$  on Sexts and Wigglers M. Boscolo- March 10th 2008

# **Backgrounds and Luminosity** versus years of **KLO**E data taking



	(10 <sup>31</sup> cm <sup>-2</sup> s <sup>-1</sup> )	(kHz)	(kHz 10 <sup>31</sup> cm <sup>2</sup> s <sup>1</sup> )	probability
2001	2.5	130	50	8%
2002	6	50	8.3	5%
2004	8	50	6.25	2.5%
2005	10	40	4	2.5%

# Program Flow Touschek simulation

**Optics** check

(nonlinearities included)

Beam parameters calculation (betatron tunes, emittance, synchrotron integrals, natural energy spread, bunch dimensions, optical functions and Twiss parameters all along the ring)

Calculation of Touschek energy spectra all along the ring averaging Tousc. probability density function over 3 magnetic elements

### Tracking of Touschek particles:

Start with transverse gaussian distribution and proper energy spectra every 3 elements: track over many turns or until they are lost

 •Estimation of IR and total Touschek particle losses (rates and longitudinal position)
•Estimation of Touschek lifetime

Touschek energy spectra related mostly to beam parameters (i.e. bunch volume, ε, σ<sub>p</sub>, bunch current...) With a given energy spectrum P(E) (see next slide) one can: 1. extract according to P(E) or 2.Use a uniform extraction and use P(E) as a weight





We use 2. to cope with tails of both distributions

# Calculation of energy spectra

Starting formula: Integrated Touschek probability

$$\frac{1}{\tau} = \frac{\sqrt{\pi}r_e^2 cN}{\gamma^3 (4\pi)^{3/2} V \sigma'_x \varepsilon^2} C(u_{\min})$$

$$\frac{1}{\tau} = \int_{\varepsilon}^{\infty} P_{Tou}(E) dE$$

$$\varepsilon = \frac{\Delta E}{E} \qquad u_{\min} = \left(\frac{\varepsilon}{\gamma \sigma_x}\right)$$
$$\sigma'_x = \sqrt{\frac{\varepsilon_x}{\beta} + \sigma_p^2 \left(D'_x + D_x \frac{\alpha_x}{\beta_x}\right)^2}$$

 $\sqrt{2}$ 

V = bunch volume=  $\sigma_x \cdot \sigma_y \cdot \sigma_l$ 

C(umin) accounts for Moller x-section (polarization is included) and momentum distribution

For a chosen machine section the Touschek probability is evaluated in small steps (9/element) to account for the beam parameters evolution for 100  $\epsilon$  values.

Use an interpolation between the calculated  $\varepsilon$  values according to the Touschek scaling law:  $A_1 \cdot \varepsilon^{-A_2}$ 

Effects of non linearities on Touschek particle losses

Tracking includes non linearities: sextupoles and octupoles relevant to account for the correct dynamical aperture



Expected beam losses at the KLOE IR for one bunch of 10 mA: contributions change when nonlinearities are taken into account







 $\beta_x$ ,  $\beta_y$ , H and dispersion functions for the DAFNE crab waist configuration

Touschek particles that get lost are scattered in high dispersive regions

Distribution of Touschek scattering position for losses all over the ring (upper plot) or only at the IR (lower plot).

## Optimal position for collimators along the rings









In DAFNE IP very close to last high dispersive region, about 10 m

when all collimators are inserted only Touschek particles scattered in the closest arc before the IP are lost at the IR

collimators needed close to the IP

careful study of collimator shape to avoid background generation

### phase advance between collimators and IR

in order to stop particles that get lost at IR, phase advance between collimator and QF should be ~  $\pi/2$ 



# Typical behaviour of the measured background rate as a function of collimators openings

The background rate measured by KLOE for an electron beam as a function of the KLOE collimator setting

Ncal (Hz/mA\*\*2)



### (dashed blue line):

the external jaw reduces the background by about a factor of 2.3 at a distance of ~25 mm from the beam axis, while the closed orbit was measured at -2.7 mm from the center of the chamber. A steep increase of the background rate is observed when collimators are closed to less than about 23 mm ( $9\sigma_x$ ) from the closed orbit position. Apparently, from this moment on, collimators are producing more background particles than they are stopping.

#### (full red line)

It is shown the behaviour of the **modified** collimator. The external jaw reduces the background by a factor 2.9 at a distance of  $\approx$ 20 mm from the chamber axis. No background reduction is found when moving in the internal jaw, however, the previously observed strong increase is no longer present (dashed line), indicating an improved stopping efficiency of the scraper blocks.

### **Modified Collimator**



The total scraper thickness of now 5.5 cm (about 16 r.l.) is reducing the punch through probability of 500 MeV electrons to below 10<sup>-6</sup>. The inner surface of the modified collimator blocks has been divided into two flat parts.

A first 1 cm long section has a slope of 100 mrad towards the beam, in order to **increase the impinging angle** into the block for most particles. This is followed by a second section

of 4.5 cm length with a slope of 10 mrad in the opposite direction to avoid foreward scattering of electrons back into the beam pipe.

## **Collimators Modeling**

Perfectly absorbing collimators

•No width

collimators assumed perfectly absorbing and infinitely thin

actual behaviour is reproduced but

Edge effect is missing

It has been found that most of the particles are scattered by the collimator edge, instead of being absorbed, thereby producing additional background to the experiments.

real collimator shape included in simulation and edge effect has been simulated

### Simulation results with real collimator

# Electron interactions at collimator edge

Electron interaction: Multiple scattering, Bremsstrahlung, de/dx simulated by a toy MC



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### IR collimator more effective with other ones in

In addition to a direct background reduction the collimators added in dispersive regions helped in making the collimator upstream the experiment more effective. In fact, they stop particles that would be just deviated by the IR collimator and eventually lost at the experiment.



Scan of the normalized background rate in the KLOE forward calorimeter versus position of the external jaw of the positron beam KLOE collimator with all other ones out (red dots) and in (blue squares).

A factor 1.6 is gained due to the fact that it can be inserted closer to the center of the pipe.

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### Collimators on/off measurements

at the end of the KLOE data taking -in optimized conditions in terms of L/bkg ratio and of vacuum conditions- (March 06)

scrapers	I[A]	σ <sub>x</sub>	σγ	τ <b>[S]</b>	ECM2 [KHz]
All OFF	474	1.30	0.15	3561	252.3
All ON	461	1.27	0.13	3125	12.5

**e**-

105 bunches, I<sub>bunch=</sub>5mA

$$\frac{\text{scr.OFF}}{\text{scr.ON}} = 20$$

<del>)</del> +	scrapers	I[A]	σ <sub>x</sub>	σγ	τ <b>[S]</b>	ECM4 [KHz]
	All OFF	524	1.12	0.21	2378	946.4
	All ON	524	1.08	0.20	1959	19.4

$$\frac{\text{scr.OFF}}{\text{scr.ON}} = 49$$

### 20/03/2006 measurement e+



Collimator behavior differently with respect to past now they are very efficient: In 2002 collimators were set not so inside the pipe,

Efficiency was not so high

In 2006

- Most of them were at maximum inner positions, reducing continously hot rates at calorimeter
- never reduce lifetime
- never enhance backgrounds
- Collimator need to be VERY clean in order to be efficient- very slow process;
- ⇒ Emittance decreased by a factor 2
- Copper tapering for preventing discontinuity in the pipe for impedance reasons has been removed from borders of collimators

**KLOE IR** 



Comparison between measured and calculated effectiveness of collimators

The calculated rate is evaluated by tracking Touschek scattered particles from their loss point in the pipe into the KLOE detector. The endcap acceptance has been taken into account by means of full detector simulation including the geometrical details of the IR.

### There is a qualitative agreement between measurement and simulation.



Scan of the background rate in the KLOE forward calorimeter versus position of the internal jaw of a collimator:

### measured normalized rate in red dots,

calculated normalized rate in black squares and normalized lifetime in blue triangles.

The collimator opening is measured from the beam pipe edge.

### Comparison between expected and measured background rates at the KLOE calorimeter versus sextupoles strengths

The MC reproduces actual behaviour of background vs sextupoles strengths



Simulated loss rates have been normalized to measurement (27/09/01): nb=47; Ibunch=5mA (Itot=100mA) and normalization is made for R=0.1 (R= $\sigma$ y/ $\sigma$ x)

# Masks added between pipe and low- $\beta$ quads



### IR layout + background event x-s view

Low  $\beta x$  in Dear



### Touschek particles trajectories at DEAR IR

### calculated background With last December 2001 optics

## calculated background with April 2002 optics



Studying the particles distribution at DEAR IR improved background rates both by the indication of lowering the  $\beta$ x-function at the closest quad to IP2 and by shielding at the calculated position



## Touschek particles trajectories at KLOE IR

triplet FDF low- $\beta$  quads

doublet low- $\beta$  quads



# Touschek Backgrounds view from the experiments- an example



# Backgrounds versus Luminosity in two typical periods of KLOE data taking

x 10<sup>2</sup>



## 2005 KLOE run: background rate/current vs day



# Touschek Backgrounds for the Crab waist scheme at DAFNE

Energy deviat.	0.003 -0.02	
տ <sub>p</sub> /p	4 e-4	
$\epsilon_x$ (m rad)	0.2.10-6	
coupling	0.005	
N <sub>p</sub>	2·10 <sup>10</sup>	
I <sub>bunch</sub> (mA)	10	

### BEAM DISTRIBUTION AT IP





	DAΦNE KLOE	DA <b>ΦNE</b> Upgrade
I <sub>bunch</sub> (mA)	13	13
N <sub>bunch</sub>	110	110
β <sub>y</sub> * (cm)	1.7	0.65
β <sub>x</sub> * (cm)	170	20
σ <sub>y</sub> * (μm)	7	2.6
σ <sub>x</sub> * (mm)	0.7	0.2
σ <sub>z</sub> (mm)	25	20
θ <sub>cross</sub> /2 (mrad)	12.5	25
<b>Φ</b> <sub>Piwinski</sub>	0.45	2.5
L (cm <sup>-2</sup> s <sup>-1</sup> ) x10 <sup>32</sup>	1.6	>5

Collision scheme with large Piwinski angle and Crab Waist

•Large Piwinski angle  $\Phi$  by:

smaller  $\sigma_x$ larger  $\theta$  $\varphi \approx \frac{\sigma_z}{\sigma_x} \frac{\theta}{2}$  $\zeta_y \propto \frac{N\sqrt{\beta_y}}{\sigma_z \theta}$   $\zeta_x \propto \frac{N}{(\sigma_z \theta)^2}$   $L \propto \frac{N\zeta_y}{\beta_y}$ 

•colliding beams overlap  $\Sigma$  becomes:











### Commissioning with Crab Waist scheme started last December



Peak Currents

I-≈ 1 A, I<sup>+</sup> ≈ 0.6 A



As expected, vertical beam size shrinks



Benefits of crabbed waist scheme -measurements:

- synchrobetatron resonances suppression with fine tune scans
- shrinking of vertical beam size switching on crab sextupoles
- luminosity increase by a factor 2

In progress

- •Coupling minimization by careful rotation IR QUADS done,
- orbit tuning, nonlinearities optimization
- •Minimization of initially very high background rates in calorimeter and Kaon monitor: lead shieldings added, with collimators in

•Short Lifetime due to initially bad vacuum, now investigating possible scrapering somewhere along the rings

## no collimators inserted in simulation

# with collimators inserted in simulation



## no collimators inserted in simulation

## with collimators inserted in simulation



about a factor 5 with collimators inserted at present machine set, in agreement with measurements

# collimators on/off

### measurements 26/2/08



$$\frac{\gamma^{-}(\text{scr.OFF})}{\gamma^{-}(\text{scr.ON})} = 7.35$$
$$\frac{K^{-}(\text{scr.OFF})}{K^{-}(\text{scr.ON})} = 7.4$$

## simulations

$$\frac{\dot{N}(\text{scr.OFF})}{\dot{N}(\text{scr.ON})} = 5.5$$

10 mA, 1bunch	lifetime
NO SCRAPERS	3220 s
WITH SCRAPERS	2410s

## Investigation of losses downstream the IP



particles lost downstream the IP, at the QFO, are Touschek scattered in the closest arc before the IP and cannot be stopped by IR collimator

these particles, as well as upstream ones can produce high showers in the close-by calorimeter, lead shielding has been put to protect calorimeter and SIDDHARTA detector

# Siddharta Set-up

SIDDHARTA

bhabha monitor quads gamma monitor







## Investigations on short beam Lifetime

Assuming a smaller physical aperture in dipoles we check Touschek lifetime to estimate if it can be that we are scrapering the beam at dipoles



Touschek Background at the LER of the SUPERB Factory

	CDR		
Luminosity x 10 <sup>36</sup>			1
			_
Circumference (m)		2250	_
Revolution frequency	(MHz)	0,13	-
Eff. long. polarization	(%)	0	_
RF frequency (MHz)		476	_
Harmonic number		3570	
Momentum spread		8,4E-04	
Momentum compaction	on	1,8E-04	
Rf Voltage (MV)		6	
Energy loss/turn (MeV	)	1,9	
Number of bunches		1733	
Particles per bunch x*	<b>10</b> <sup>10</sup>	6.16	
Beam current (A)		2,28	
Beta y* (mm)		0,30	
Beta x* (mm)		20	
Emit y (pmr)		4	
Emit x (nmr)		1,6	
Sigma y* (microns)		0,035	
Sigma x* (microns)		5,657	
Bunch length (mm)		6	
Full Crossing angle (m	nrad)	34	
Wigglers (#)		4	
Damping time (trans/l/	ong)(ms)	32/16	
Luminosity lifetime (m	nin)	10,4	
Touschek lifetime (mi	n)	5,5	
Effective beam lifetim	e (min)	3,6	
Injection rate pps (100	1%)	4,9E+11	
Tune shifts (x/y) (from	formula)	0.004/0.17	
RF Power (MW)			17

#### SBF Parameters June,15-2007

### last parameter set (Pantaleo)

DADAMETED	LED
	LER
Porticle type	
Factor (Co)	ет
Energy (Gev)	4
Luminosity x 10 <sup>55</sup>	1,0
	1700
Circumference (m)	1780
Revolution frequency (IVIHZ)	0,169
Eff. long. polarization (%)	U 470
RF frequency (IVIHZ)	4/6
Harmonic number	2824
Momentum spread	7,9E-04
Momentum compaction	2,6E-04
Rf Voltage (MV)	7
Energy loss/turn (MeV)	1,09
Number of bunches	1342
Particles per bunch x10 <sup>10</sup>	5.52
Beam current (A)	2,00
Beta y* (mm)	0,22
Beta x* (mm)	35
Emit y (pmr)	7
Emit x (nmr)	2,8
Sigma y* (microns)	0,039
Sigma x* (microns)	9,90
Bunch length (mm)	5
Full Crossing angle (mrad)	48
Wigglers (#) 10 meters each	0
Damping time (trans/long)(ms)	44/22
Luminosity lifetime (min)	7,18
Touschek lifetime (min)	13,8
Effective beam lifetime (min)	4,7
Injection rate pps (100%)	2,6E+11 2
Tune shift y (from formula)	0,15
Tune shift x (from formula)	0,0043
RF Power (MW)	12

Touschek particles lost at IR Rates are given for 1 bunch with  $I_{bunch}$  =1.49 mA



 $\Delta E/E = 0.1\% - 4\%$ rf accept. =2.9 % machine turns = 5 K=0.25%  $\epsilon_x$ =2.8 nm ;  $\sigma_z$ =5 mm

IR Losses = 1.7 MHz Tou. lifetime ≈ 24 min



#### work in progress to optimize collimators

for best longitudinal position along the ring and best trade-off between IR losses and lifetime

## Energy acceptance



### no collimators

# Conclusions

In DAFNE a lot of effort has been put in the backgrounds minimization being extremely high at the beginning of each run. DAFNE background is essentially Touschek

- Collimators: position and shape crucial
- Shieldings: very useful for small experiments
- Optics: IR design critical, small  $\beta_x$  required (synergy with L), nonlinearities and dynamic aperture optimization, as well as fine tuning of orbit at the IR

# Back-up slides DAFNE

### Bunch length measurement



# Half IR1 Magnetic Layout





### June 30<sup>th</sup> 2001: 24hours KLOE data taking





### October 6<sup>th</sup> 2001:

### 24hours KLOE data taking

