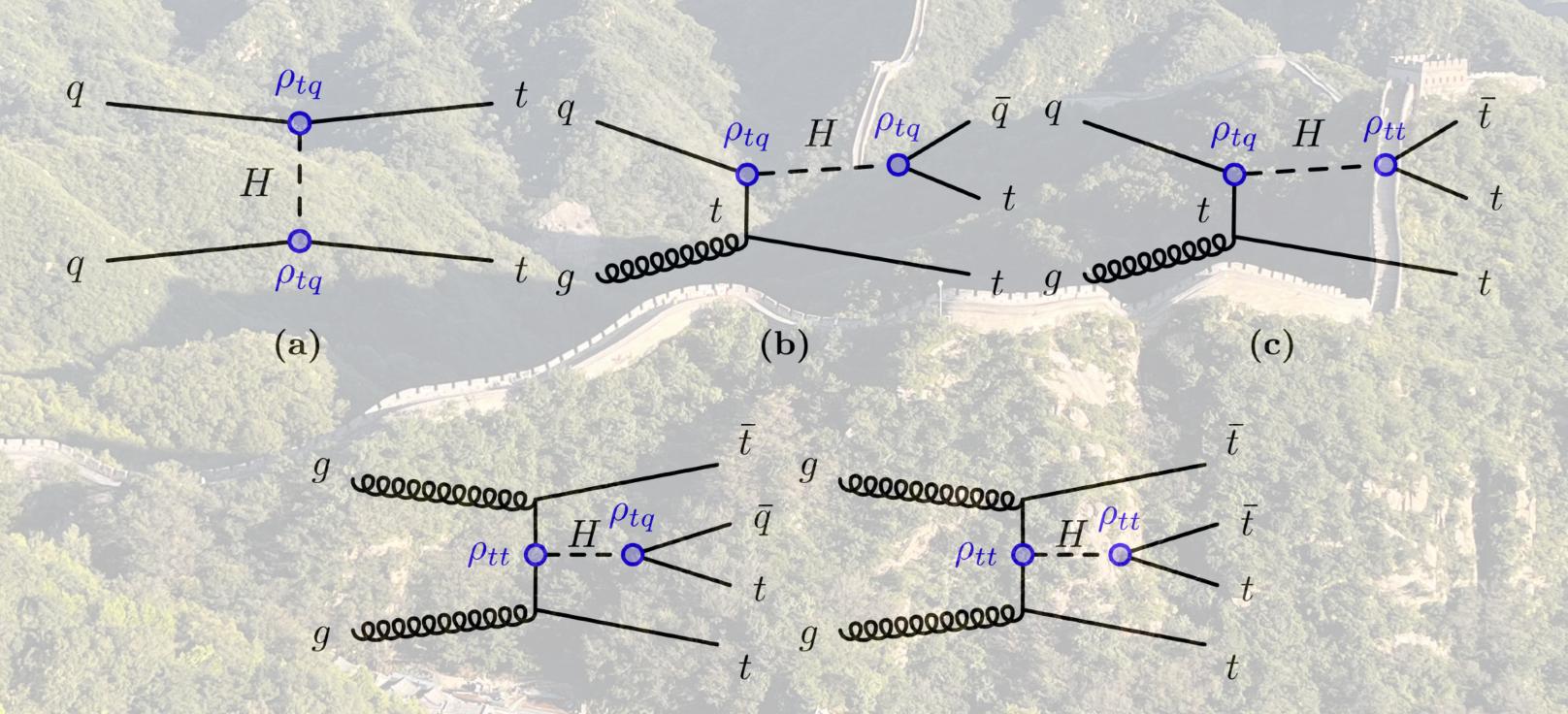




Interference effects from flavour physics in the General 2HDM and ATLAS multi-top probes



17th International Conference on Heavy Quarks and Leptons (HQL 2025), Beijing

Cristian Sierra, 2025年09月16日

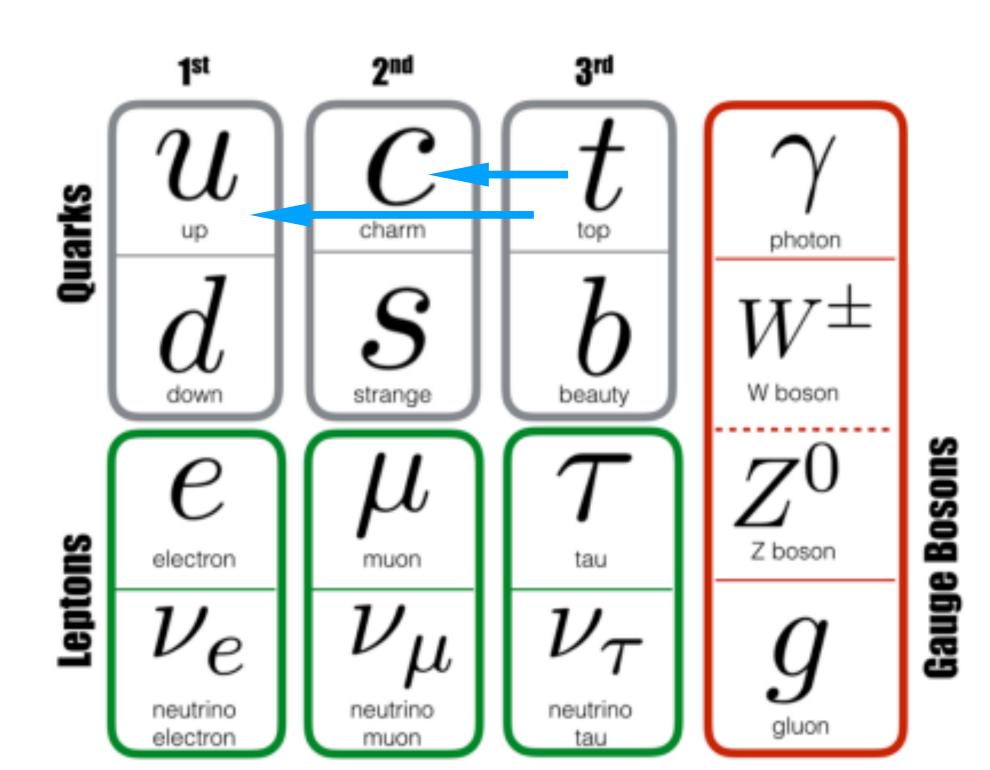


General 2HDM



Adding a second Higgs doublet is one of the simplest extensions of the SM

$$\Phi_i = \begin{pmatrix} \phi_i^+ \\ \frac{1}{\sqrt{2}}(v_i + \xi_i + i\eta_i) \end{pmatrix}, \quad i = 1, 2.$$

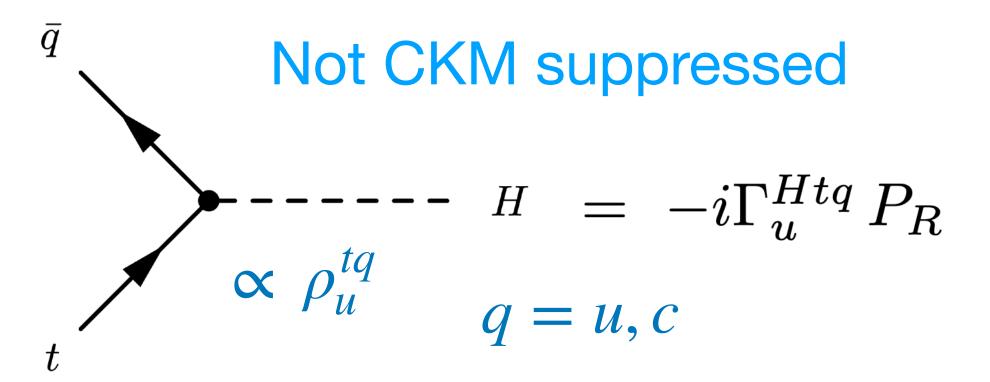


In the mass basis, the Yukawa Lagrangian has the following flavour violating couplings

$$-\mathcal{L}_{Yukawa} = \bar{u}_b \left(V_{bc} \rho_d^{ca} P_R - V_{ca} \rho_u^{cb*} P_L \right) d_a H^+ + \bar{\nu}_b \rho_\ell^{ba} P_R l_a H^+ + \text{h.c.}$$

$$+ \sum_{f=u,d,\ell} \sum_{\phi=h,H,A} \bar{f}_b \Gamma_f^{\phi ba} P_R f_a \phi + \text{h.c.},$$

$$\Gamma_f^{Hba} \equiv \frac{\bar{M}_f^{ba}}{v} c_{\beta\alpha} - \frac{1}{\sqrt{2}} \rho_f^{ba} s_{\beta\alpha}$$





General 2HDM



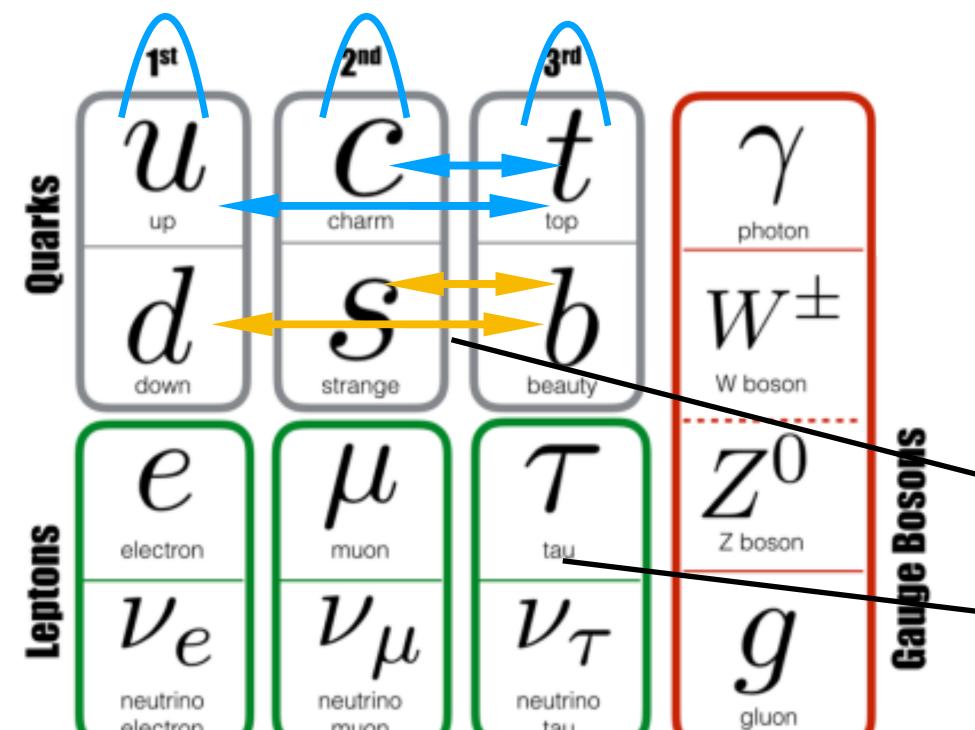
Adding a second Higgs doublet is one of the simplest extensions of the SM

$$\Phi_i = \begin{pmatrix} \phi_i^+ \\ \frac{1}{\sqrt{2}} (\upsilon_i + \xi_i + i\eta_i) \end{pmatrix}, \quad i = 1, 2.$$

$$\Phi_{i} = \begin{pmatrix} \phi_{i}^{+} \\ \frac{1}{\sqrt{2}}(v_{i} + \xi_{i} + i\eta_{i}) \end{pmatrix}, \quad i = 1, 2.$$

$$-\mathcal{L}_{Yukawa} = \bar{u}_{b} \left(V_{bc} \rho_{d}^{ca} P_{R} - V_{ca} \rho_{u}^{cb*} P_{L}\right) d_{a} H^{+} + \bar{\nu}_{b} \rho_{\ell}^{ba} P_{R} l_{a} H^{+} + \text{h.c.}$$

$$+ \sum_{f=u,d,\ell} \sum_{\phi=h,H,A} \bar{f}_{b} \Gamma_{f}^{\phi ba} P_{R} f_{a} \phi + \text{h.c.},$$



$$ho_u = \left(egin{array}{ccc}
ho_u^{uu} &
ho_u^{uc} &
ho_u^{ut} \
ho_u^{cu} &
ho_u^{cc} &
ho_u^{ct} \
ho_u^{tu} &
ho_u^{tc} &
ho_u^{tt} \end{array}
ight)$$

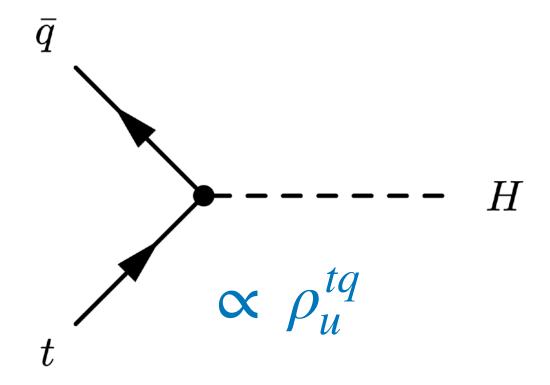
Very suppressed at tree level, only non-zero can be ρ_d^{bb} .

• Only non-zero are $ho_{\mathscr{E}}^{ au au}$ and possibly $ho_{\mathscr{E}}^{\mu au, au\mu}$.

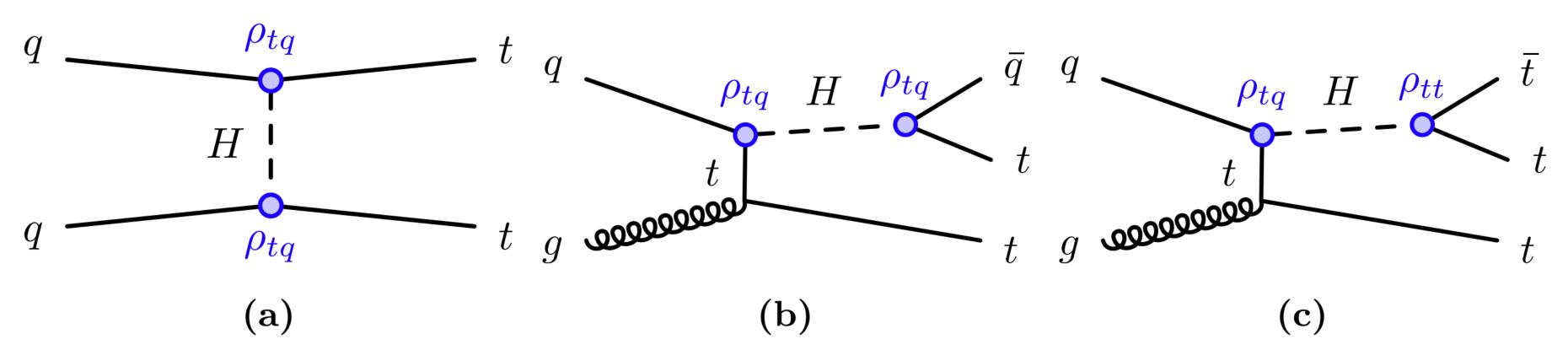


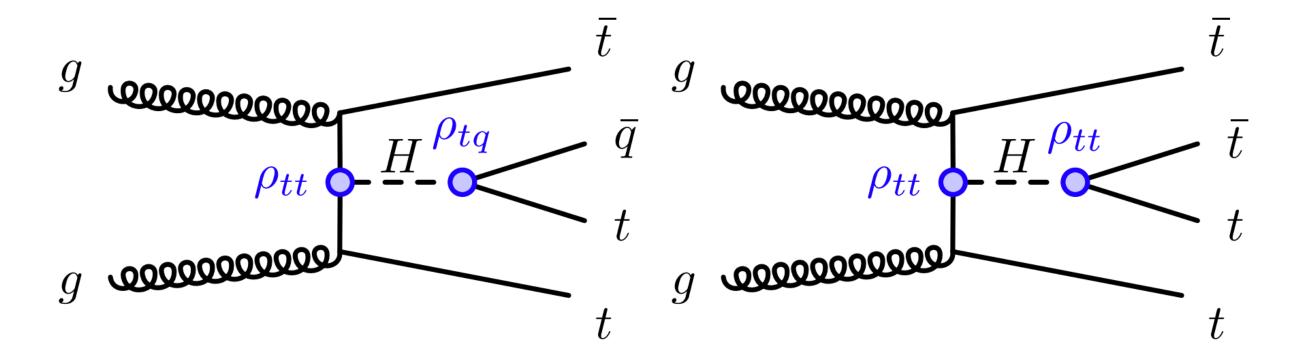
Multi-lepton + jets@ATLAS





ATLAS did a search for heavy Higgs bosons with FV couplings with the top quark https://arxiv.org/pdf/2307.14759

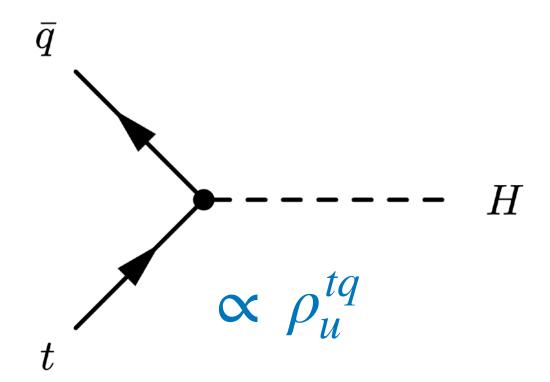






Multi-lepton + jets@ATLAS

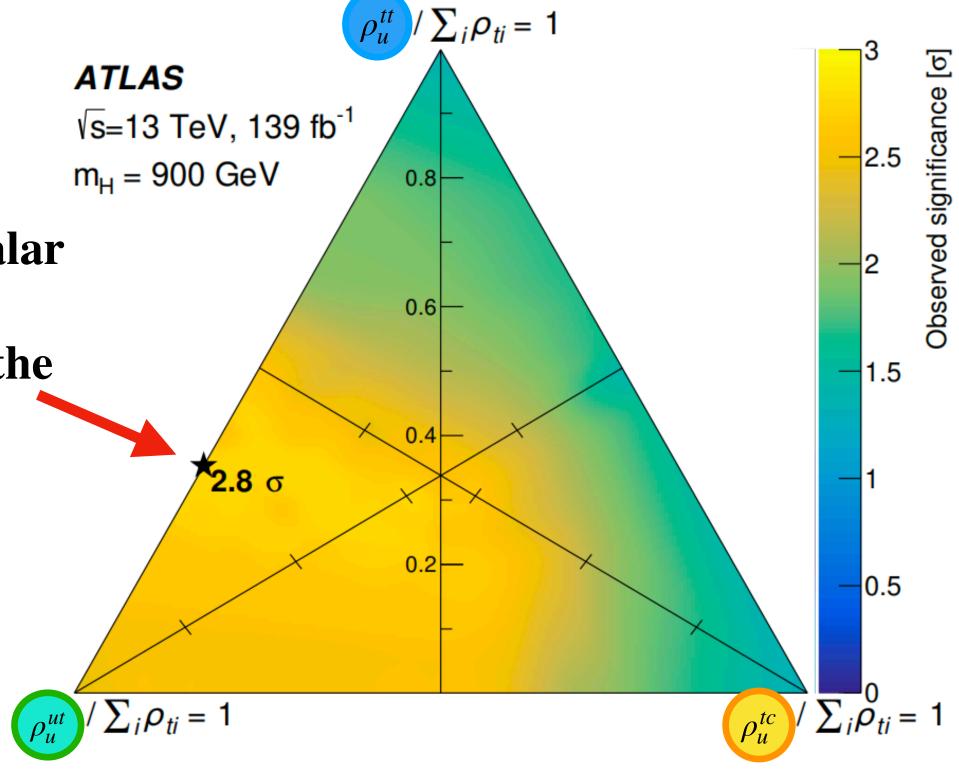




ATLAS did a search for heavy Higgs bosons with FV couplings with the top quark https://arxiv.org/pdf/2307.14759

 ρ_{tt} =0.6, ρ_{tc} =0.0, and ρ_{tu} =1.1 **ATLAS** \sqrt{s} =13 TeV, 139 fb⁻¹ 2.5 $m_H = 900 \text{ GeV}$ 0.8 Fit with no flavour constraints!

Observed significance for a heavy scalar with a mass of 900 GeV as a function of the three couplings normalised to the sum of the couplings



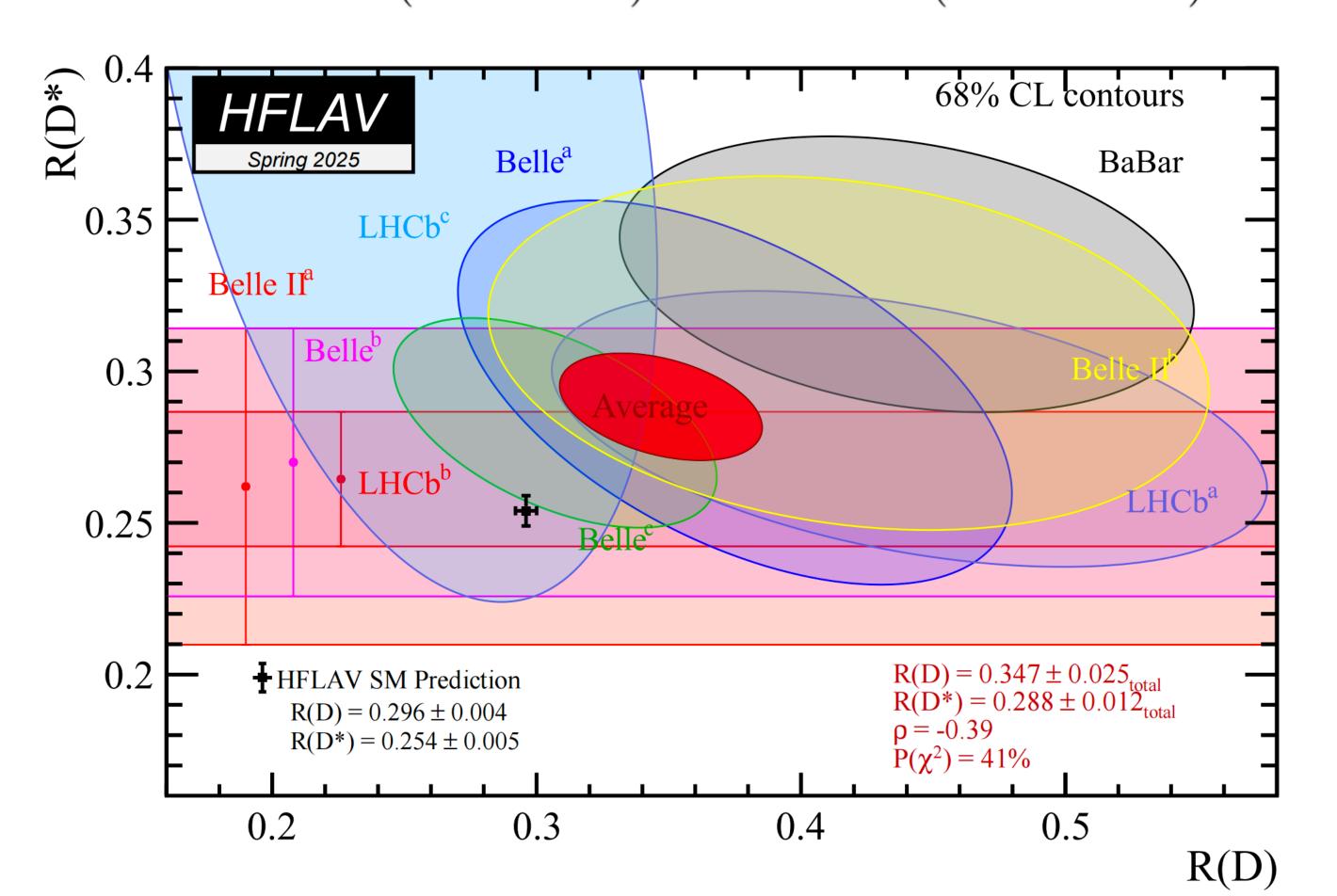
Motivation for this study:

- ρ_u^{ut} seems to big, might
 - not respect flavor constraints.
- Explaining the charged current anomalies.





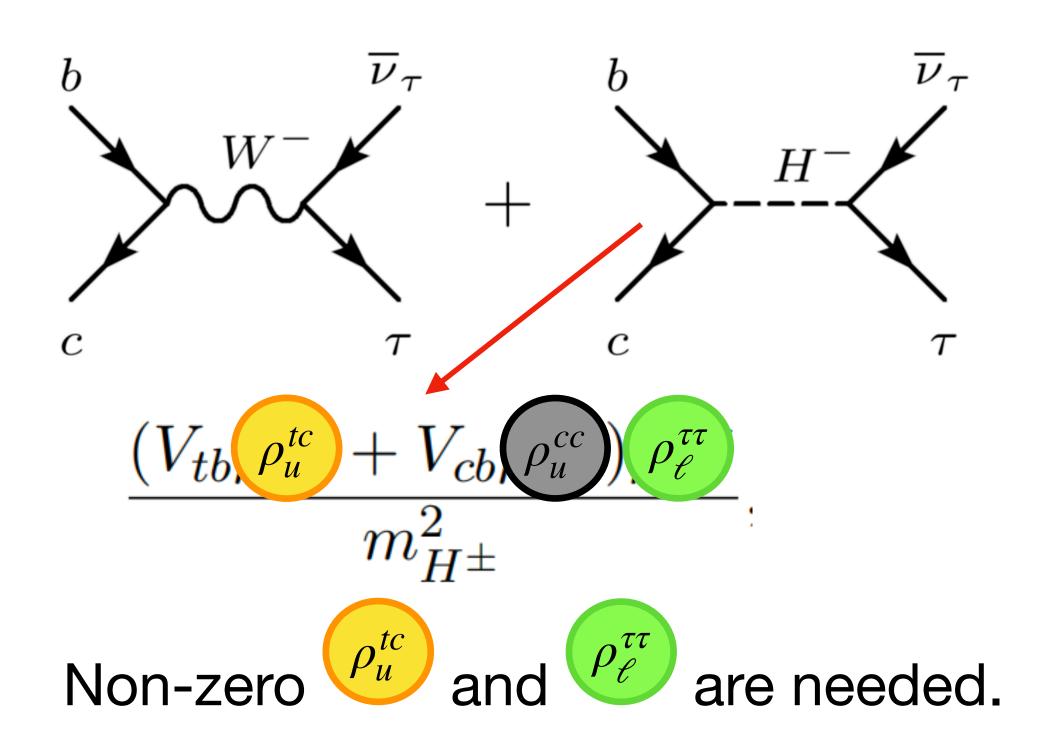
$$R_D = \frac{\Gamma(\overline{B} \to D\tau\overline{\nu})}{\Gamma(\overline{B} \to Dl\overline{\nu})} \quad R_{D^*} = \frac{\Gamma(\overline{B} \to D^*\tau\overline{\nu})}{\Gamma(\overline{B} \to D^*l\overline{\nu})}$$



$$R_{D^{(*)}}^{\mathrm{exp}} > R_{D^{(*)}}^{\mathrm{SM}}$$
 at 3

at 3.8σ

Possible interference with NP?



Taken from Moriond 2025





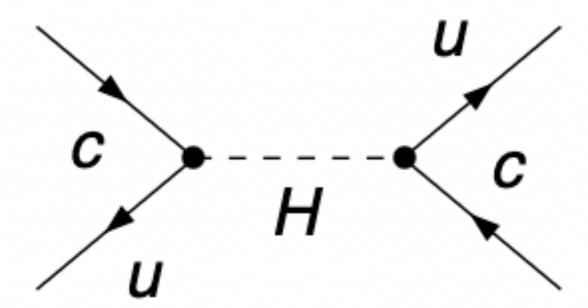
$$ho_u = \left(egin{array}{ccc}
ho_u^{uu} &
ho_u^{uc} &
ho_u^{ut} \
ho_u^{cu} &
ho_u^{cc} &
ho_u^{ct} \
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ho_u^{tc} &
ho_u^{tt} \end{array}
ight)$$





$$ho_u = \left(egin{array}{ccc}
ho_u^{uu} &
ho_u^{uc} &
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ho_u^{cc} &
ho_u^{ct} \
ho_u^{tu} &
ho_u^{tc} &
ho_u^{tt} \end{array}
ight)$$

 $-D^0 - \bar{D}^0$ mixing at tree level.

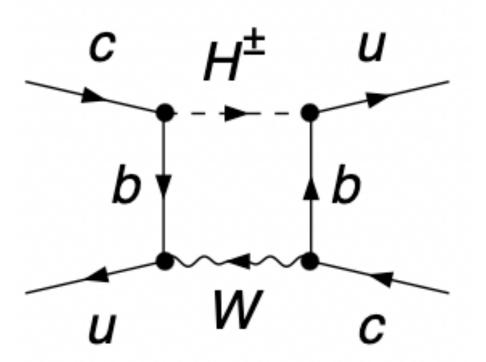


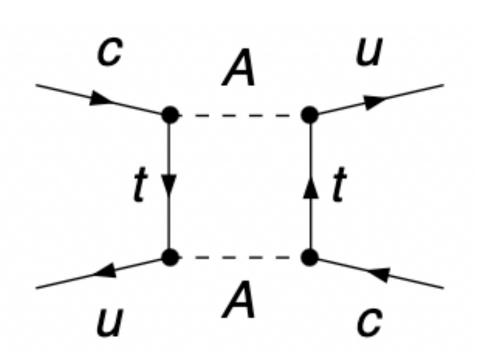


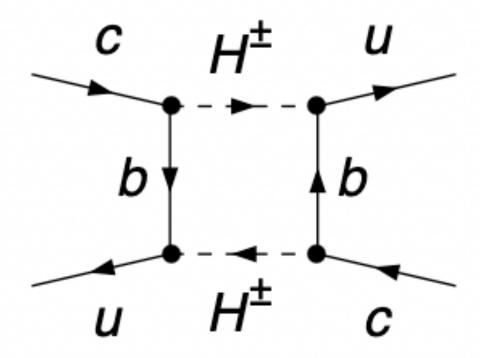


$$\rho_u = \begin{pmatrix} \rho_u^{uu} & \rho_u^{uc} & \rho_u^{ut} \\ \rho_u^{cu} & \rho_u^{cc} & \rho_u^{ct} \\ \rho_u^{tu} & \rho_u^{tc} & \rho_u^{tt} \end{pmatrix} \quad \begin{array}{l} -D^0 - \bar{D}^0 \text{ mixing at tree level.} \\ -D^0 - \bar{D}^0 \text{ mixing at one-loop.} \end{array}$$

 $-D^0 - \bar{D}^0$ mixing at tree level.





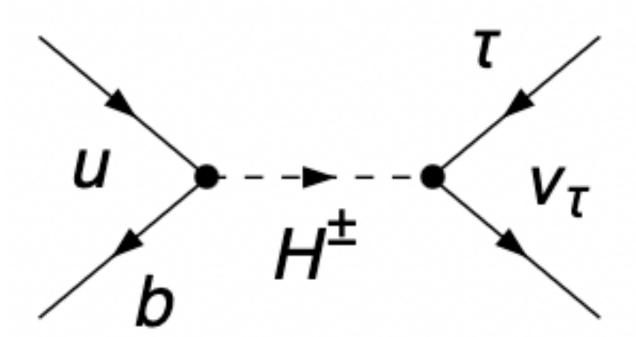






$$\rho_u = \begin{pmatrix} \rho_u^{uu} & \rho_u^{uc} & \rho_u^{ut} \\ \rho_u^{cu} & \rho_u^{cc} & \rho_u^{ct} \\ \rho_u^{tu} & \rho_u^{tc} & \rho_u^{tt} \end{pmatrix} \begin{array}{l} -D^0 - D^0 \text{ mixing at tree level.} \\ -D^0 - \bar{D}^0 \text{ mixing at one-loop.} \\ -B_u \to \tau \nu \text{ at tree level.} \end{array}$$

 $-D^0 - \bar{D}^0$ mixing at tree level.





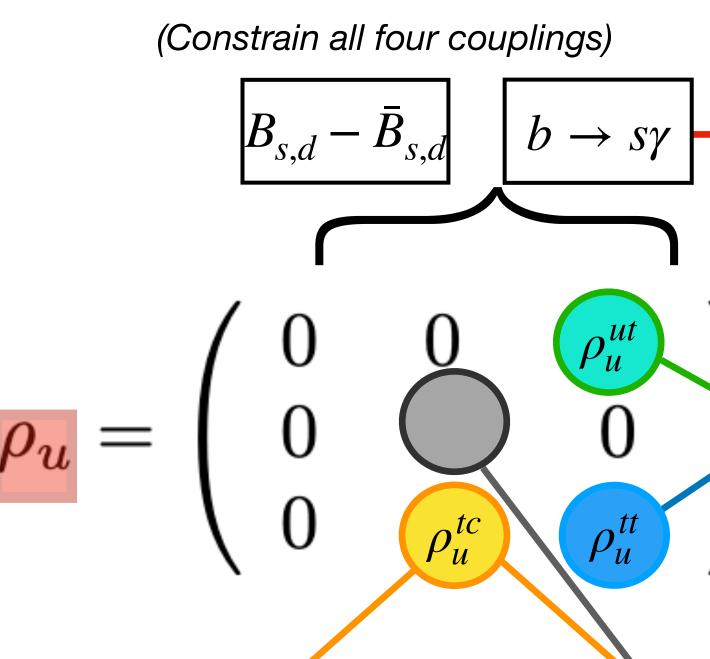
 $R(D^{(*)})$

ATLAS

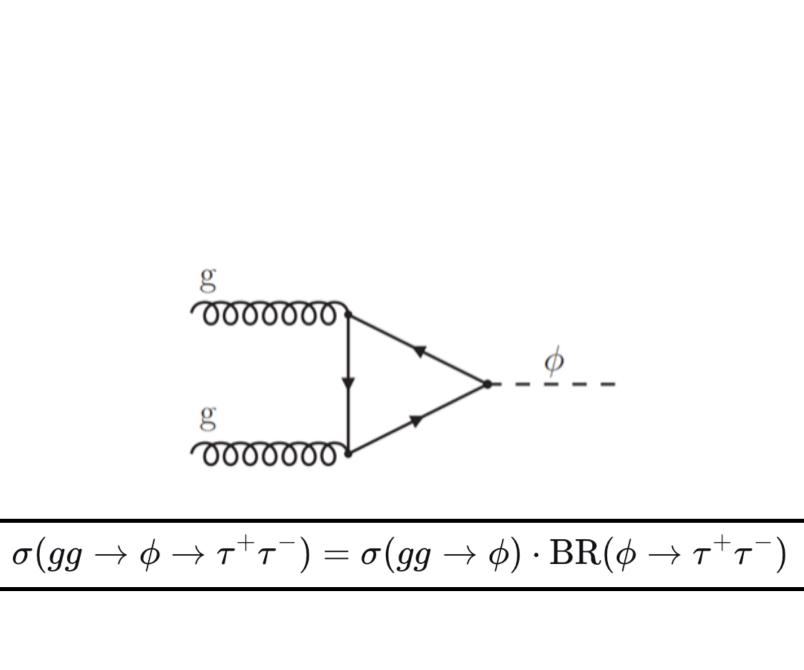
- Nested sampling scan using Bayesian inference.
- Resultant parameter
 space from flavour
 constraints.

$$m_A = m_{H^{\pm}}$$

$$m_A - m_H = 50 \,\text{GeV}$$



 $t \rightarrow hc$



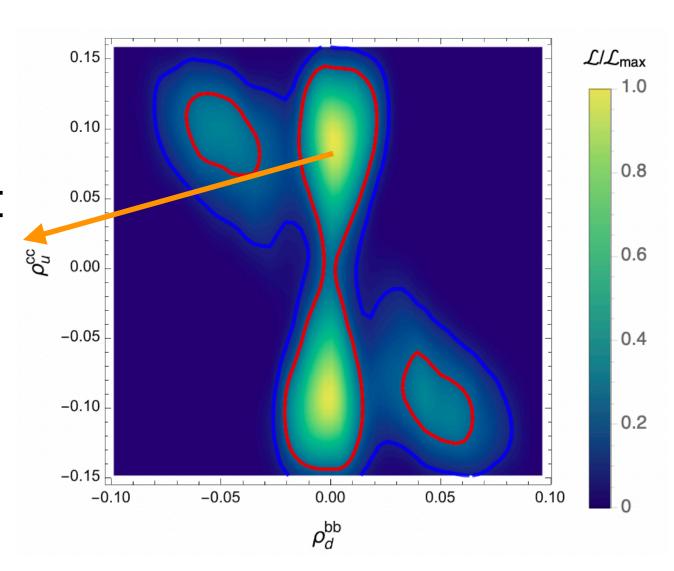
- Need to include multi-tau decays from CMS.
 - ullet eEDM constrain the imaginary part of $ho_\ell^{ au au}$.



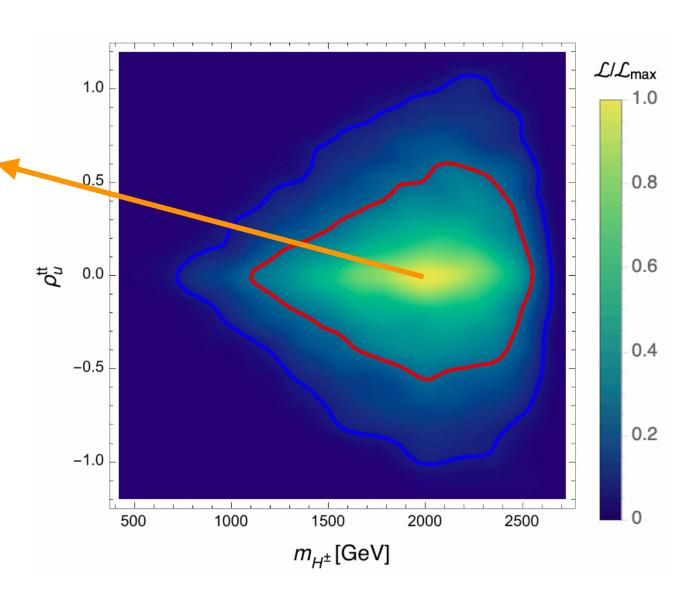
Results

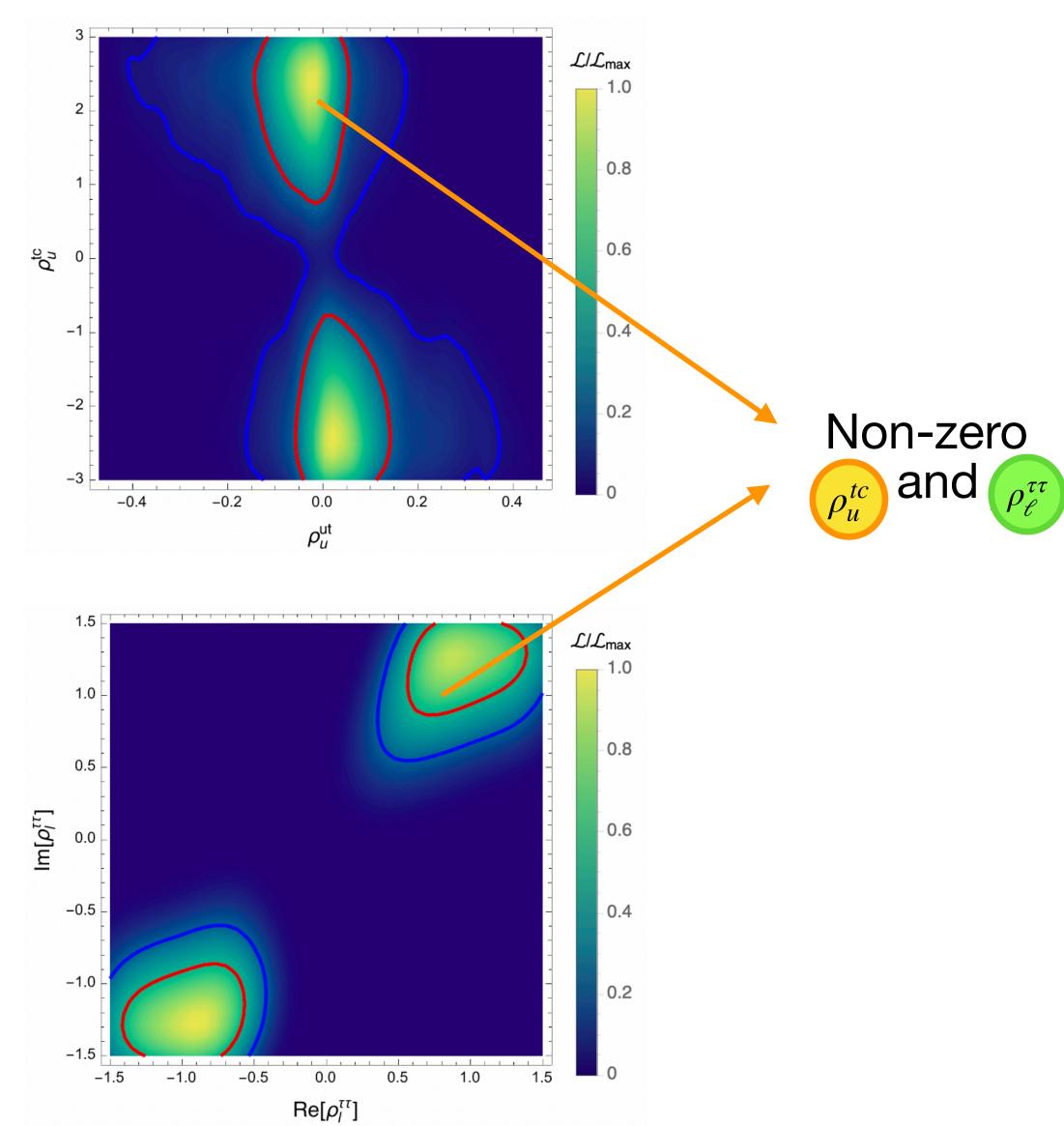


Non-zero best fit point for ρ_u^{cc} and small ρ_d^{bb}



Small ρ_u^{tt} and large masses



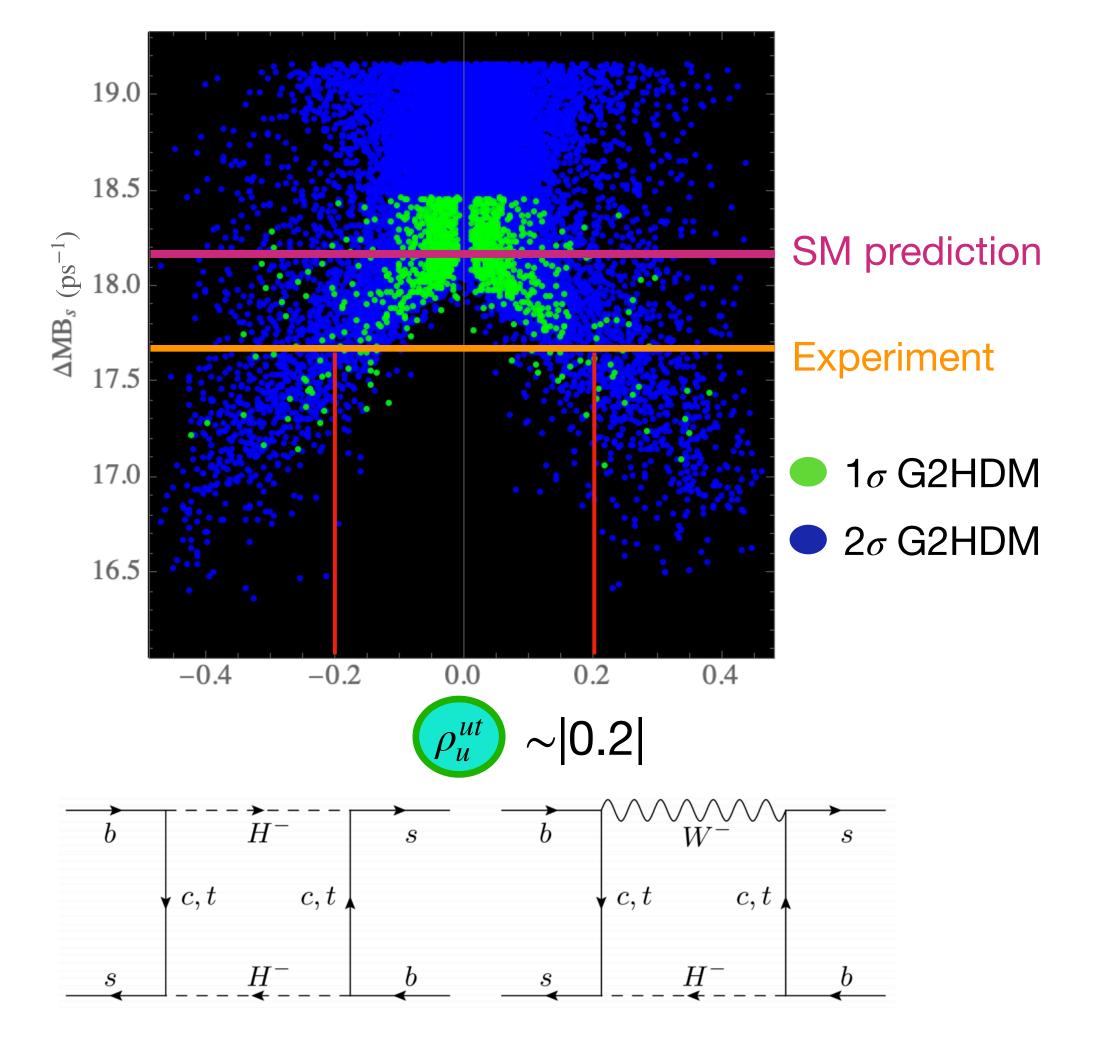




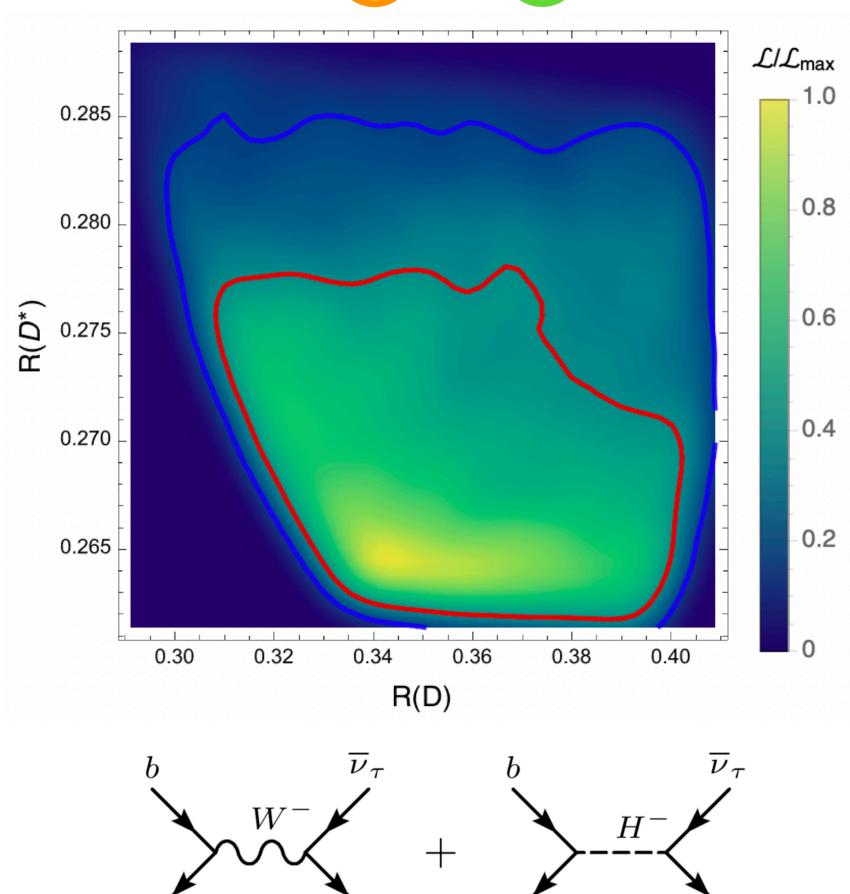
Results



SM has large uncertainties! ΔM_{B_s} = (18.2±0.7)ps⁻¹ Exp.measurement is very precise! (17.766±0.006)ps⁻¹



Fits simultaneously $R(D^*)$ ratios at 1σ with non-zero ρ_{μ}^{tc} and $\rho_{\mu}^{\tau\tau}$



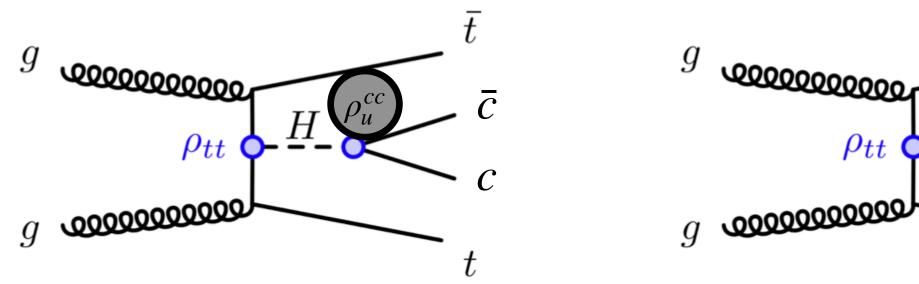


Summary



- General 2HDM offers FV couplings with rich phenomenology.
- ATLAS used the model to study multi-top production for the first time.
- Computed constraints for quark couplings from tree and loop level observables.
- ullet Interferences in meson oscillations along different up-like quark couplings improve the fit to $\Delta\,M_{B_s}$.
- The resultant parameter space can fit the $R(D^*)$ ratios at 1σ .
- Those constraints can be used for new analysis at ATLAS for multi-top decays including new

modes.



谢谢!





