

# Study of $\chi_c$ production in pPb collisions at $\sqrt{s_{NN}}$ = 8.16 TeV energy with the CMS experiment

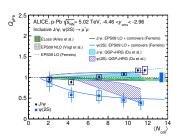
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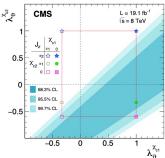
### Introduction

Proton-lead (pPb) collisions at the LHC provide an opportunity to study nuclear modification effects on quarkonia.

Excited S-wave state  $\psi(2s)$  shows different suppression than the ground state  $J/\psi$  by the ALICE experiment. A trend of increasing relative suppression of  $\psi(2s)$  to  $J/\psi$  is observed as multiplicity or related variables increases [1].



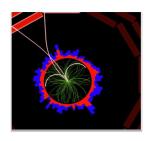
P-wave state charmonia  $\chi_c$  were studied in proton-proton collisions with 7 and 8TeV energies in the CMS experiment [2,3]. The results of polar anisotropy coefficients  $\lambda_{\vartheta}^{\chi_{c1}}$ ,  $\lambda_{\vartheta}^{\chi_{c2}}$  indicates that both  $\chi_{c1}$  and  $\chi_{c2}$  are strongly polarized [3].



The motivation of the analysis [4] is to study how  $\chi_c$  are affected in pPb compared to pp collisions based on data collected by CMS at the LHC with an integrated luminosity of  $175nb^{-1}$ . It is also the step toward the  $\chi_c$  measurements in PbPb collisions.

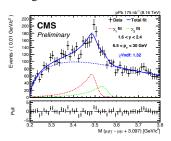
# Signal yield extraction

The  $\chi_c$  candidates are reconstructed through the radiative  $\chi_c \rightarrow J/\psi \gamma$  decays. The  $J/\psi$  is reconstructed through its decay to a muon pair, while the photon is reconstructed through its conversion to an  $e^-e^+$  pair detected in the silicon tracker.



The  $\chi_c$  mass is calculated using the world average of experimentally values of the  $I/\psi$  mass instead of the invariant mass of the di-muon:

 $m_{\chi_c} = m_{\mu\mu\gamma} - m_{\mu\mu} + 3.097 GeV,$ which improves  $\chi_c$  mass resolution by removing the di-muon resolution.

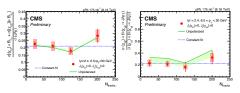


The yields of  $\chi_{c1}$   $\chi_{c2}$  are extracted from fitting the  $m_{\chi_c}$  spectrum. The nonprompt  $\chi_c$  contamination is evaluated and taken into account as a source of uncertainty.

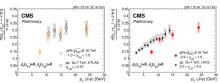
		Uncertainty [%]	
Source of uncertainty		$\chi_c$ -to-J/ $\psi$ ratio	$\chi_{c2}$ -to- $\chi_{c1}$ ratio
$\chi_c$ fit shape		0.03-4	1-4
$J/\psi$ fit shape	signal	0.3-2	_
	background	0.05-0.3	_
Conversion selection		12	23
Conversion selection - tag and probe		5	5
PYTHIA settings		6-12	0.1-0.5
Nonprompt contamination		5	8

## Results and conclusions

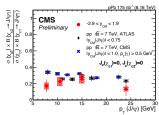
Results of the measurements are interpreted as the  $\chi_c$ -to- $J/\psi$  and  $\chi_{c2}$ -to- $\chi_{c1}$  cross section ratios. Both  $\chi_{c2}$  and  $\chi_{c1}$  are assumed to be fully polarized with  $J_z(\chi_{c2}) = J_z(\chi_{c1}) = 0$ .



 $\square$   $\chi_c$ -to- $J/\psi$  (left panel) and  $\chi_{c2}$ -to- $\chi_{c1}$  (right panel) ratio as a function of number of tracks in the event.



 $\square$   $\chi_c$ -to- $J/\psi$  ratio as a function of  $p_T(J/\psi)$  comparing to pp in a similar kinematic range and collision energy from ATLAS(left panel) and LHCb (right panel) results.



 $\square$   $\chi_{c2}$ -to- $\chi_{c1}$  ratio as a function of  $p_T(J/\psi)$ . No  $p_T$  trend and results are found to be similar to pp measurements from ATLAS [5] and CMS [2].

As a summary, no additional modification of  $\chi_c$ compared to  $I/\psi$  is observed in pPb collisions. In contrast to  $\psi(2s)[6]$ , the lack of dependence of the  $\chi_c$ -to- $J/\psi$  ratio on event multiplicity suggests weaker modification effects for  $\chi_c$  in pPb collision.

#### Reference

- [1] The ALICE Collaboration, JHEP 06 (2016) 050
- [2] The CMS Collaboration, EPJC 72 (2012) 225 [3] The CMS Collaboration, PRL 124 (2020) 162002 [4] The CMS Collaboration, CMS-PAS HIN-22-003
- [5] The ATLAS Collaboration, JHEP 07 (2014) 154[6] The CMS Collaboration, PRL 135 (2025) 092301