



Generalized symmetries in HEP and CMP

Low Entanglement Excitations in Invertible phases

Wenjie Ji

July 2025



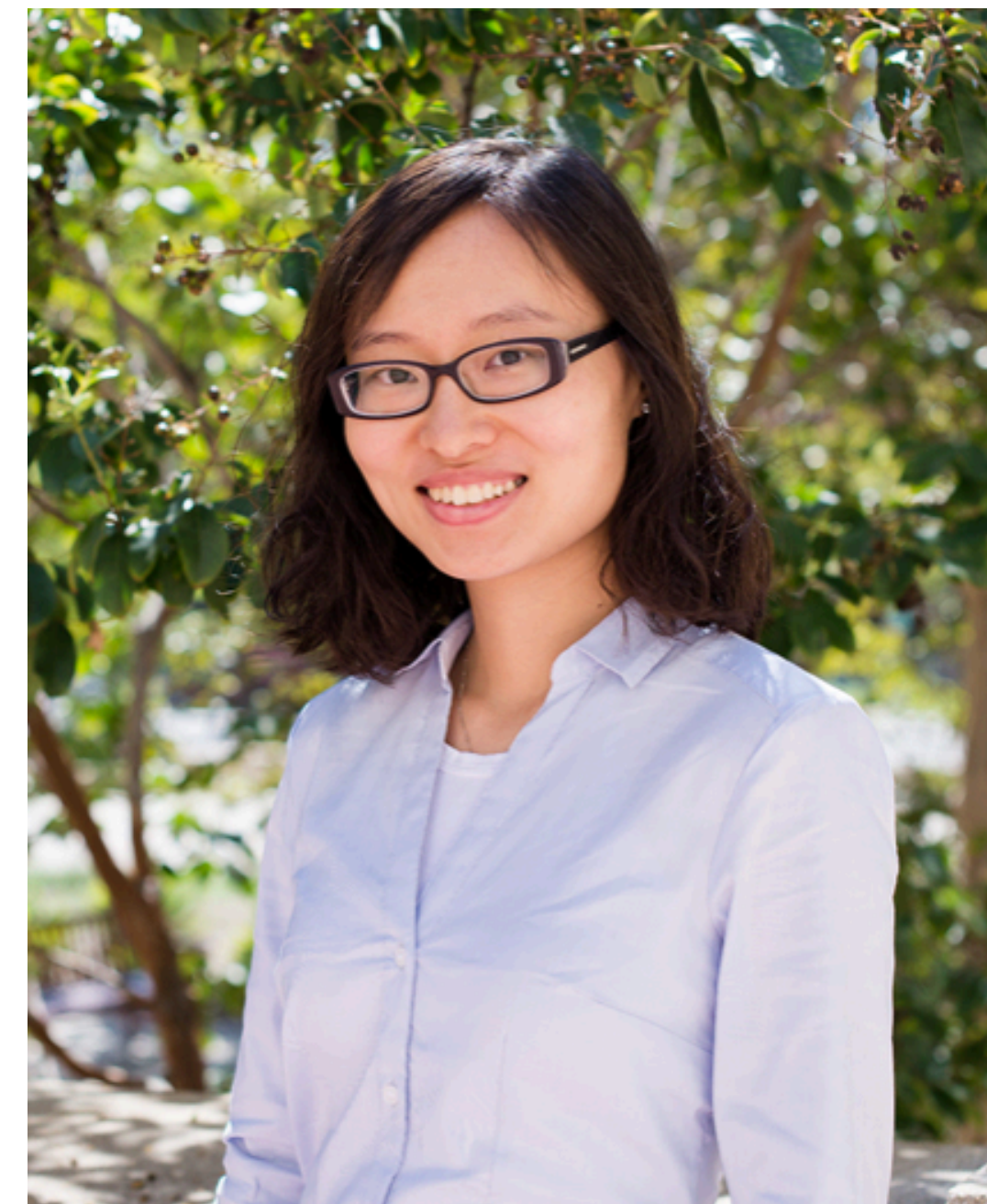
arXiv:2506.11288 with



David Stephen
University of Colorado, Boulder
→ Quantinuum

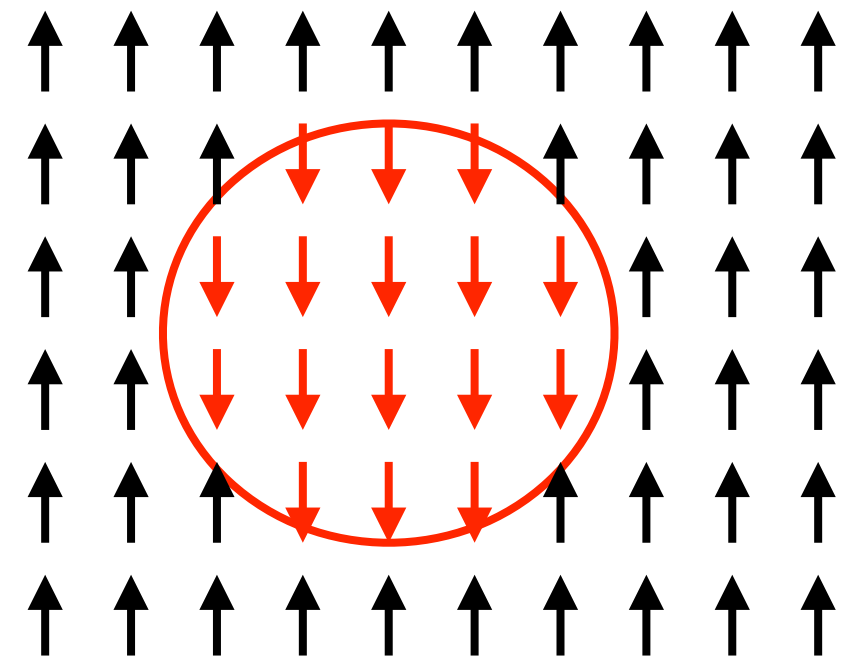


Michael Levin
University of Chicago

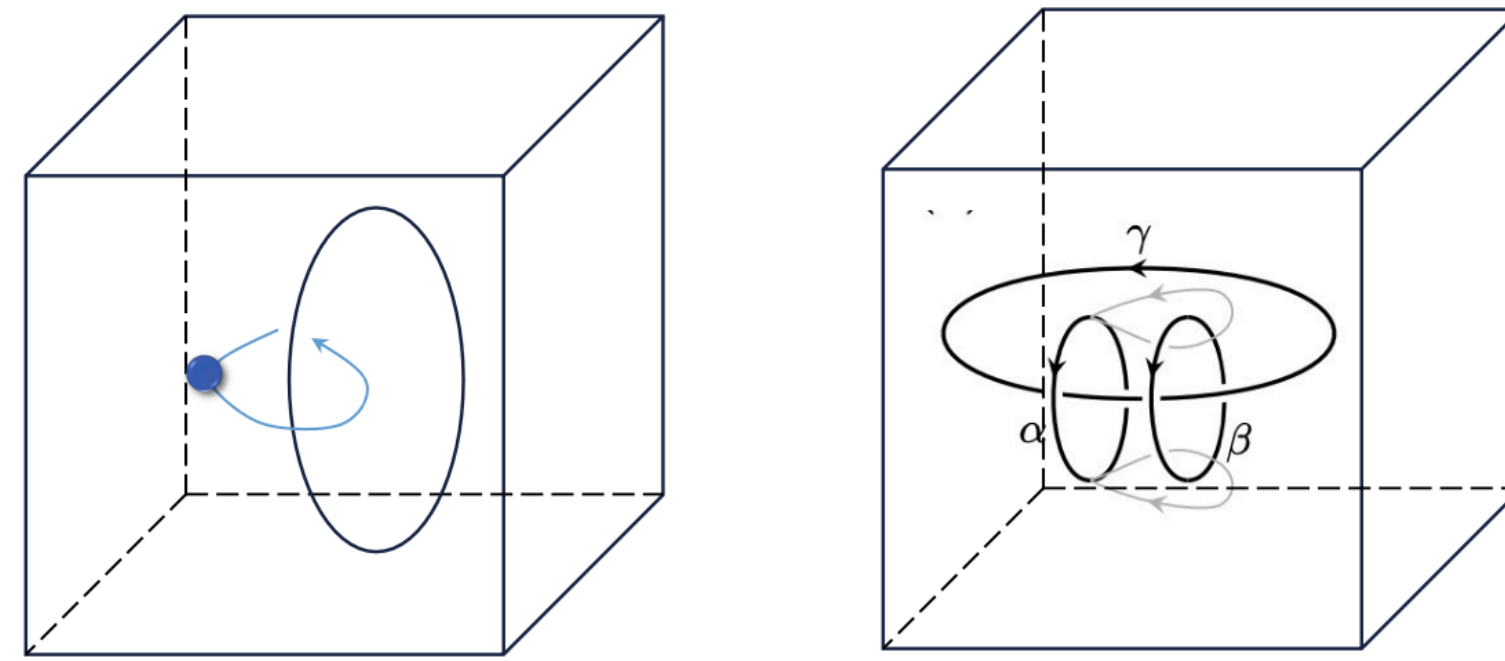


Xie Chen
Caltech

Higher dimensional excitations/defects

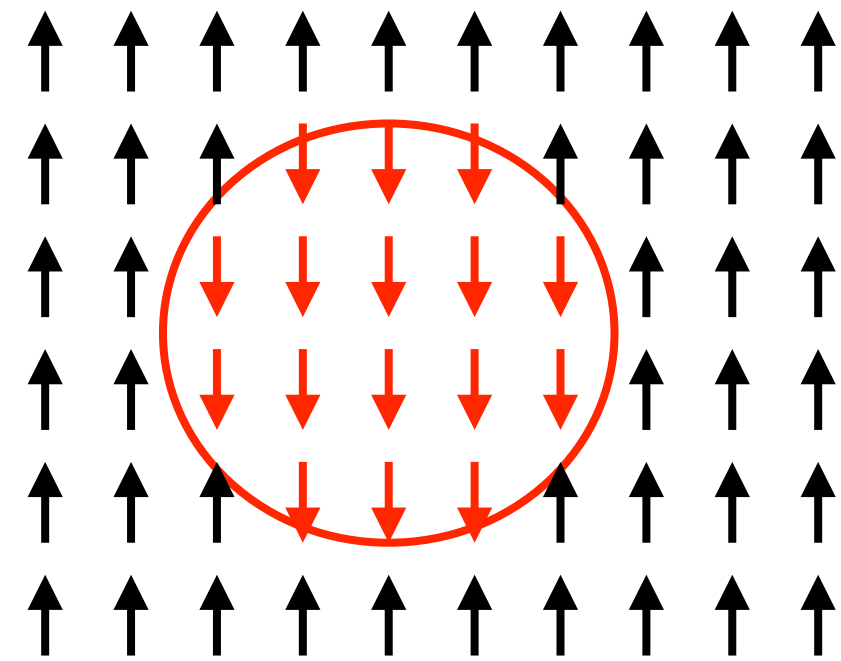


Domain wall
in Ising Ferromagnetic phase

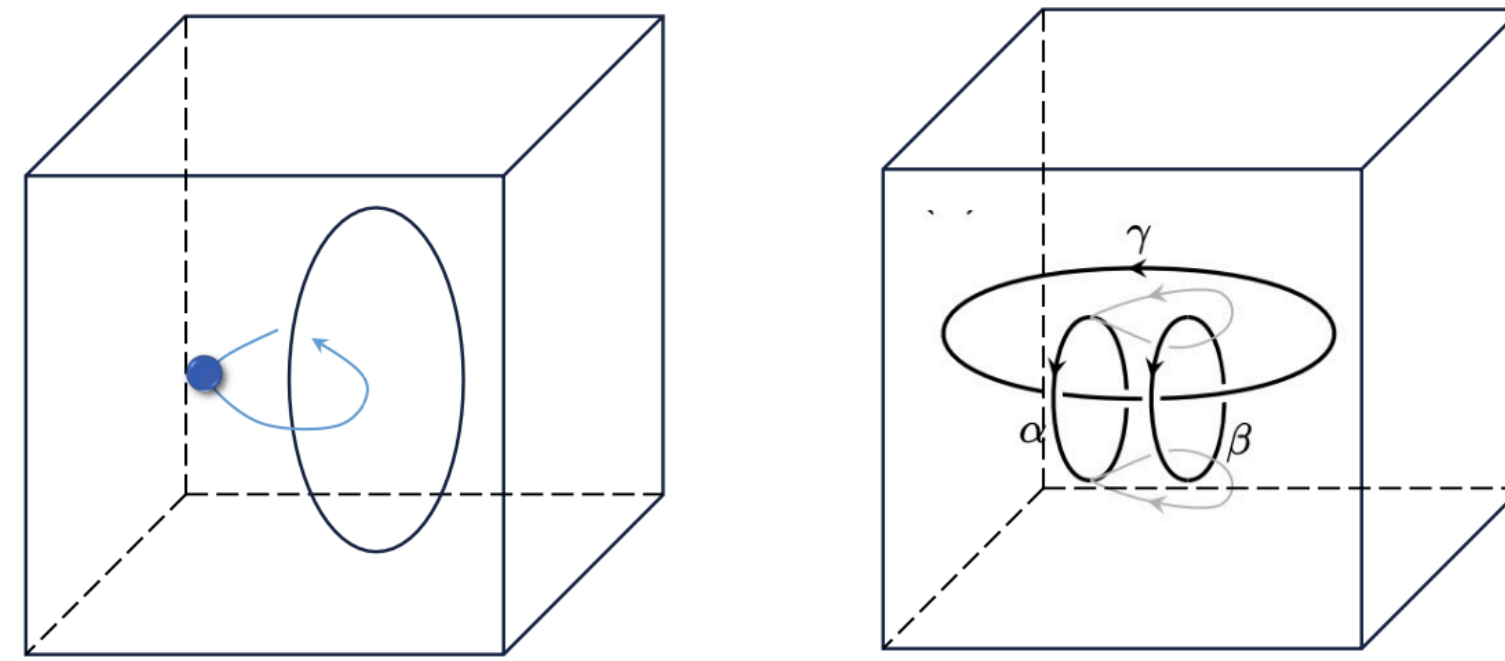


Flux loop in deconfined phase of
 G gauge theories
(G topological order)

Higher dimensional excitations/defects



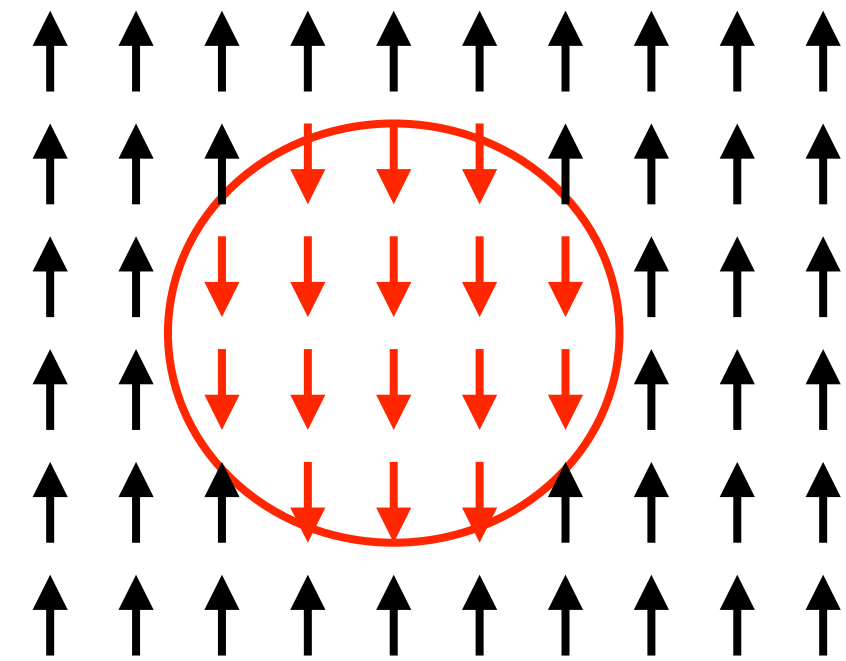
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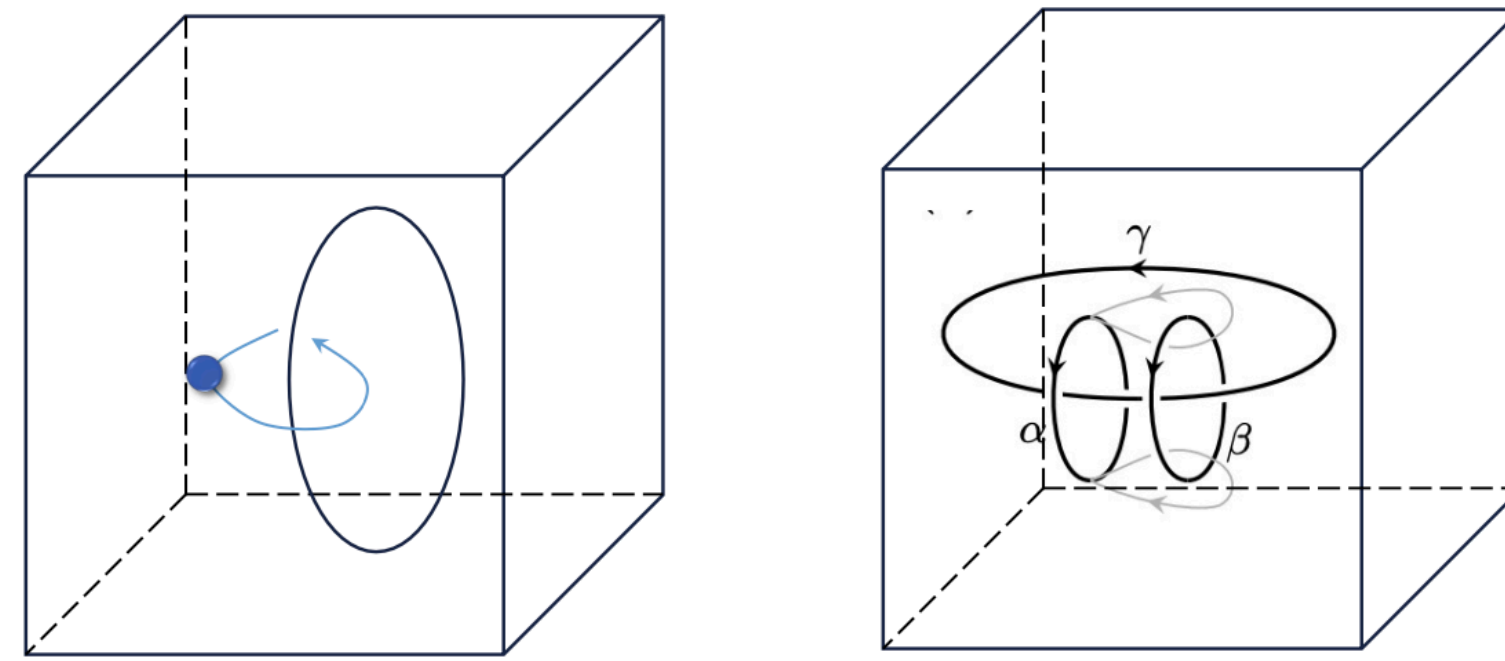
Flux loop in G gauge theories

- How to generate? 1. Modify the Hamiltonian, 2. A unitary operator on the ground state

Higher dimensional excitations/defects



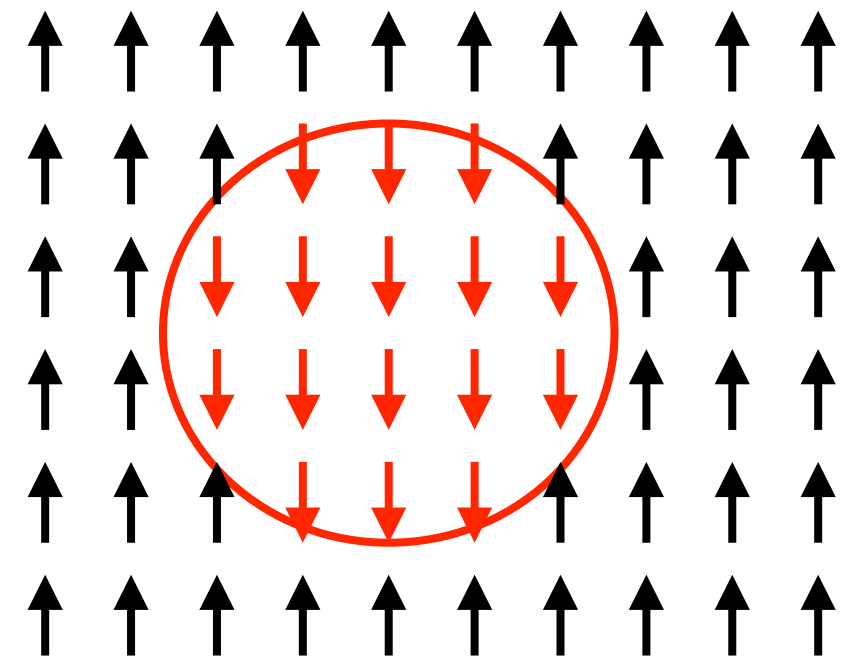
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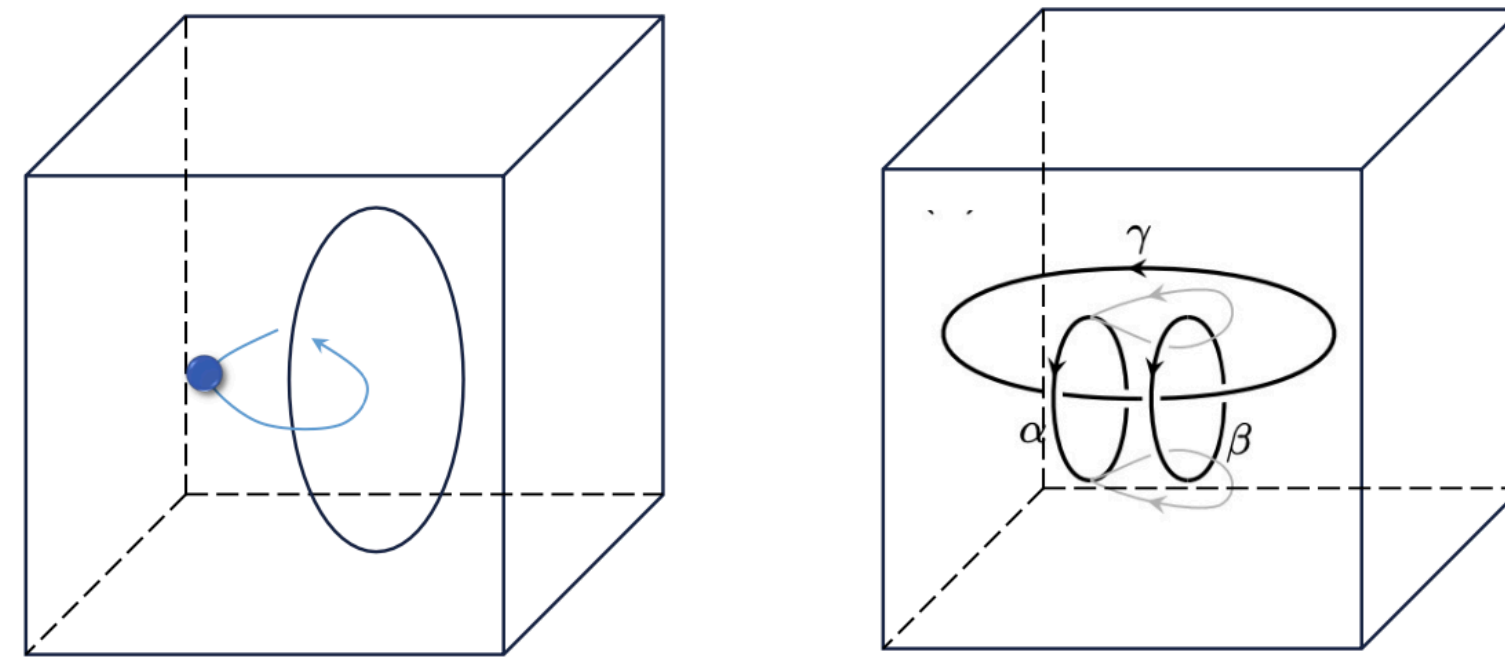
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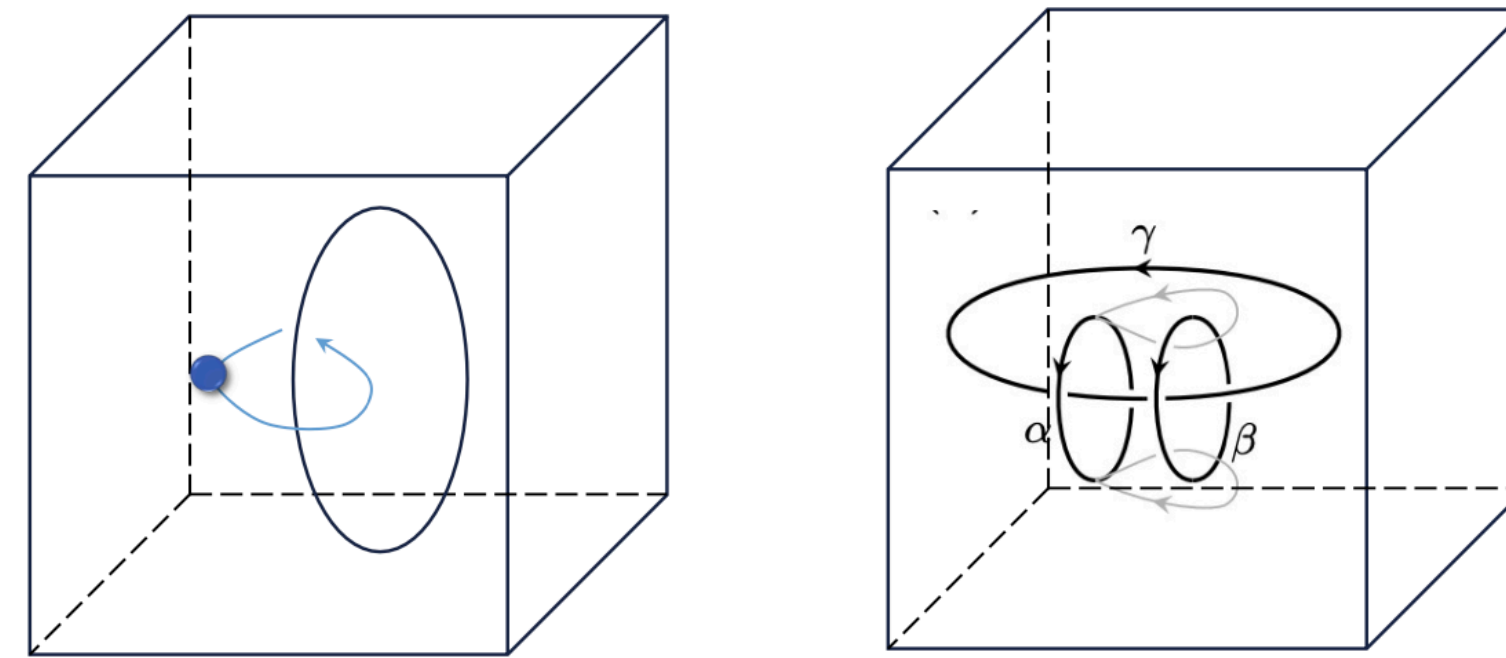
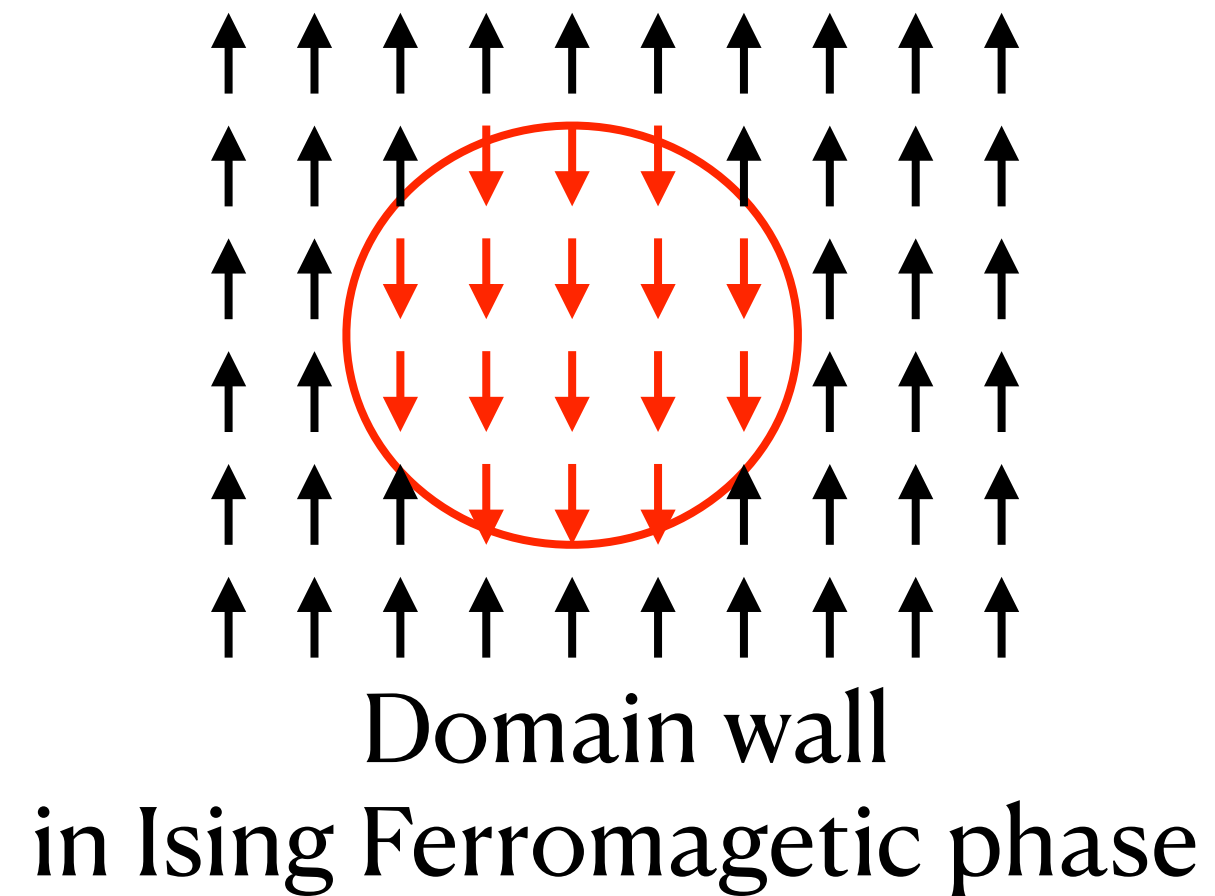
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Flux loop in G gauge theories

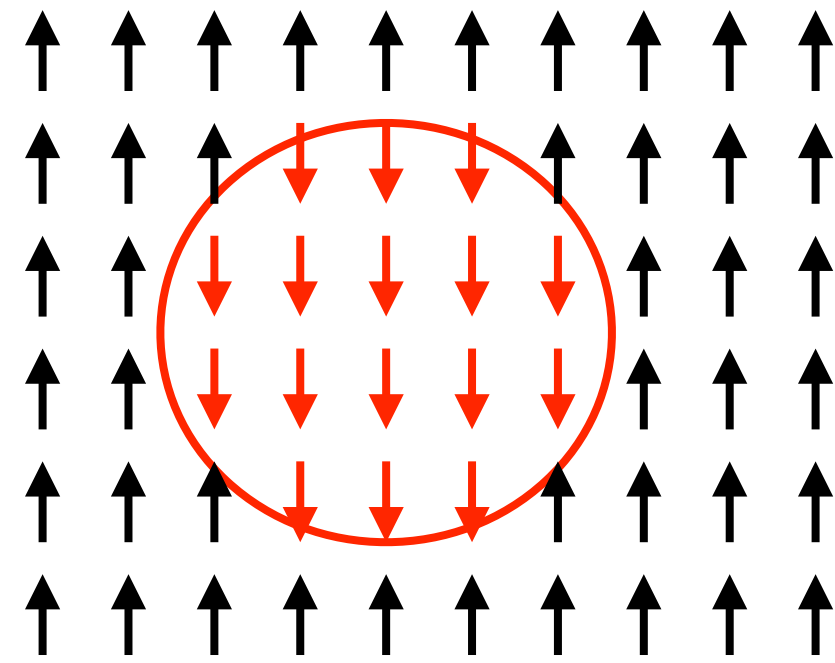
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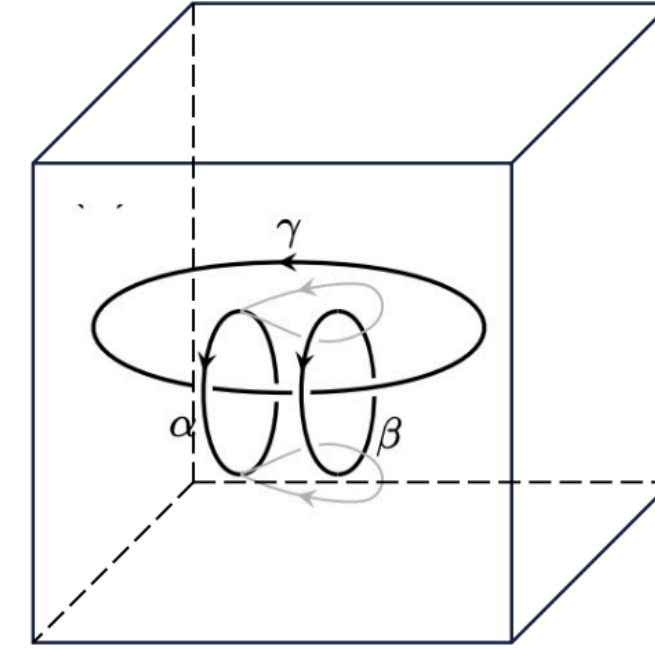
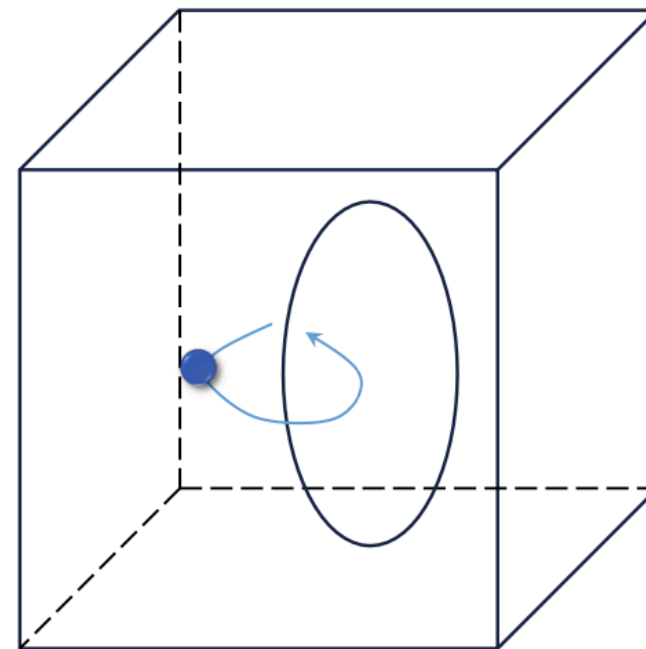


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- **Low entanglement** (area law since modified Hamiltonian has a gap)

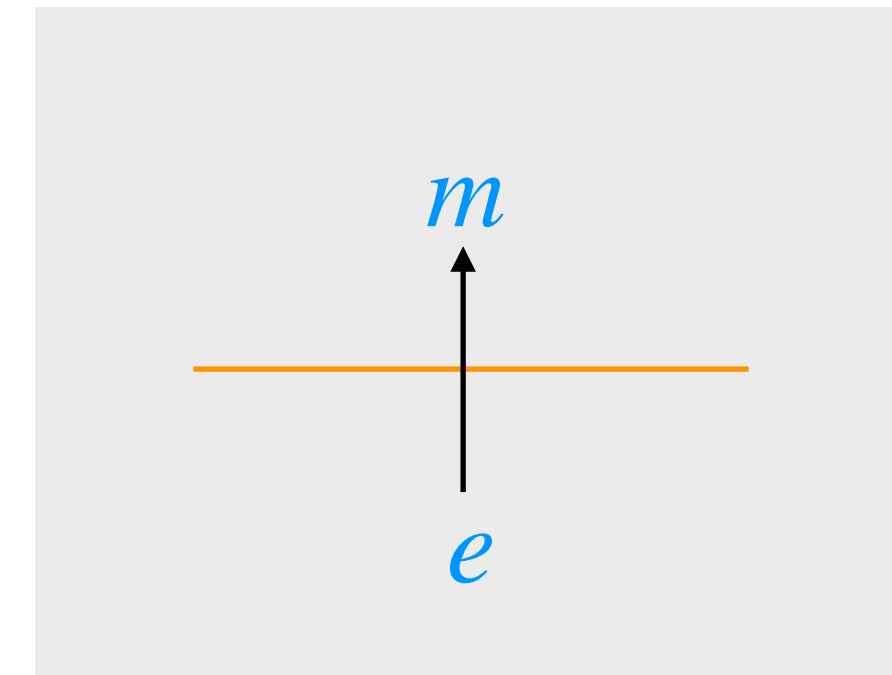
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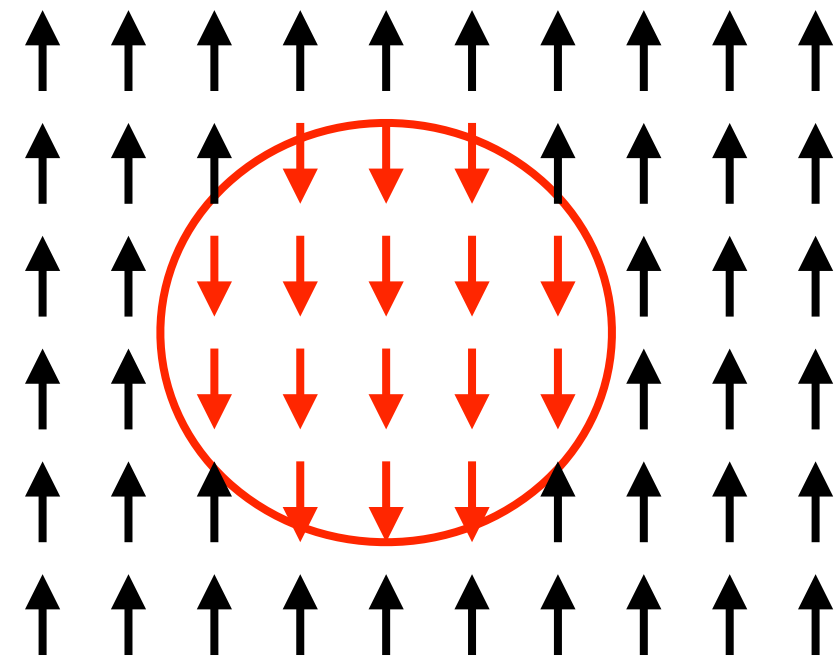


e-m exchange defect
in toric code

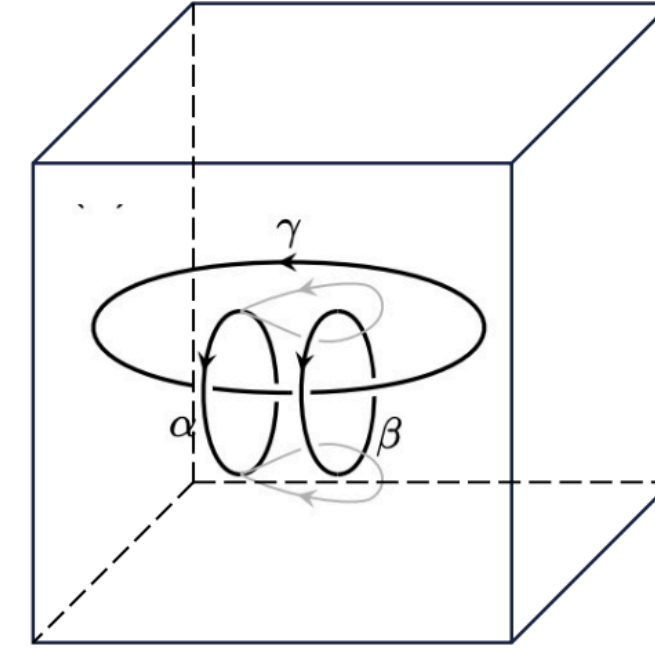
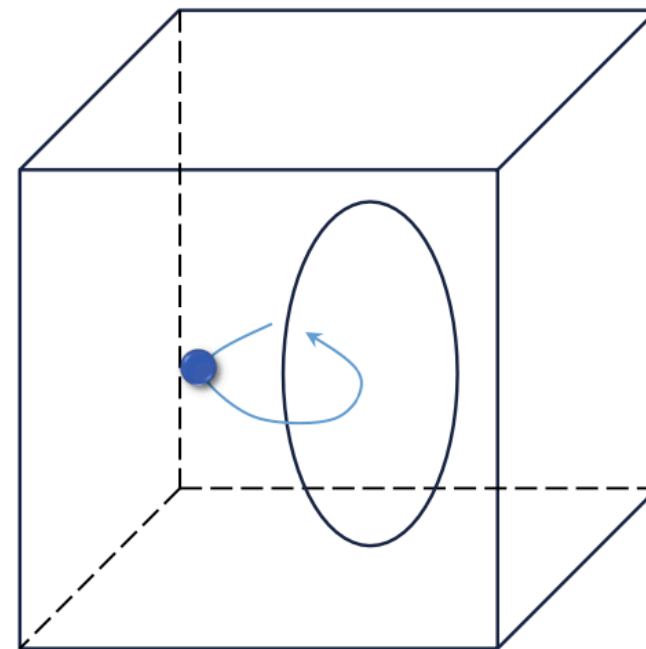
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Defects can share many of these properties

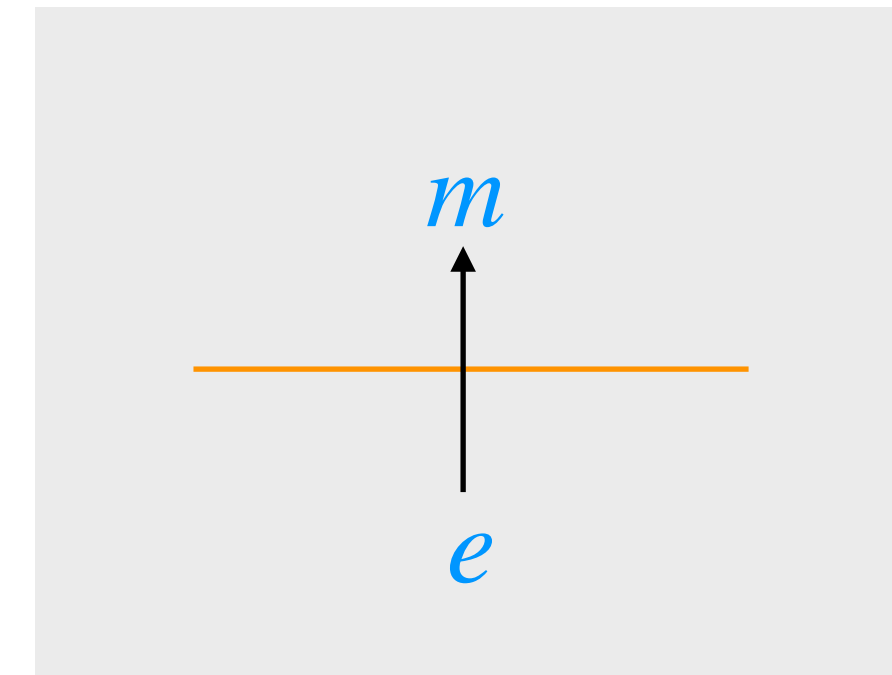
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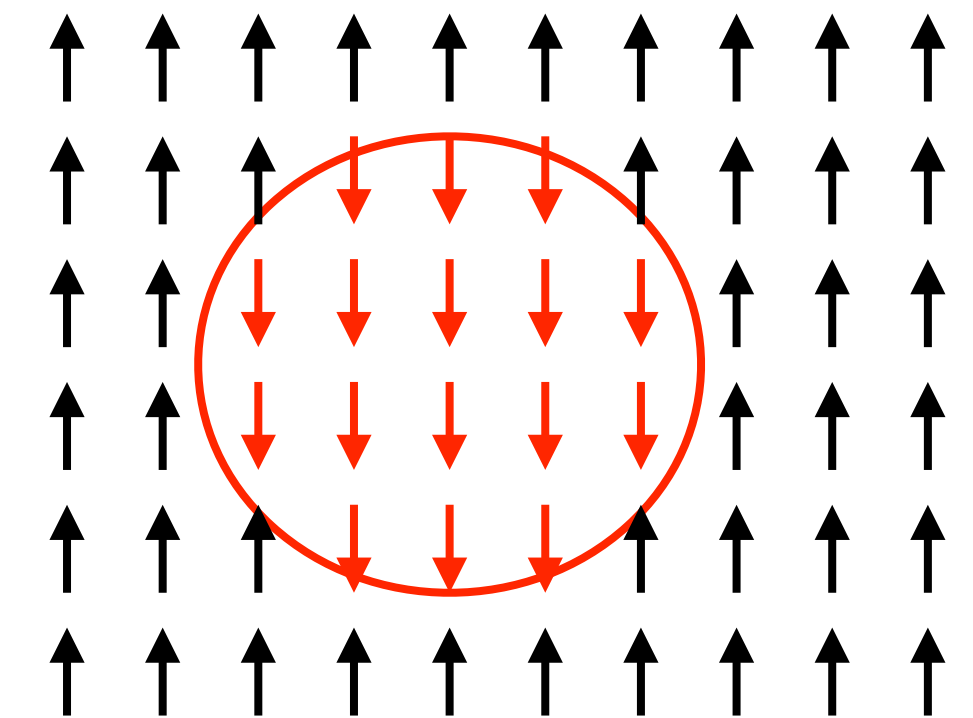
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A unified definition for excitations/defects ? Especially, those with low entanglement.

Low entanglement excitations

Use **quantum circuits** to create/define k -dimensional excitations

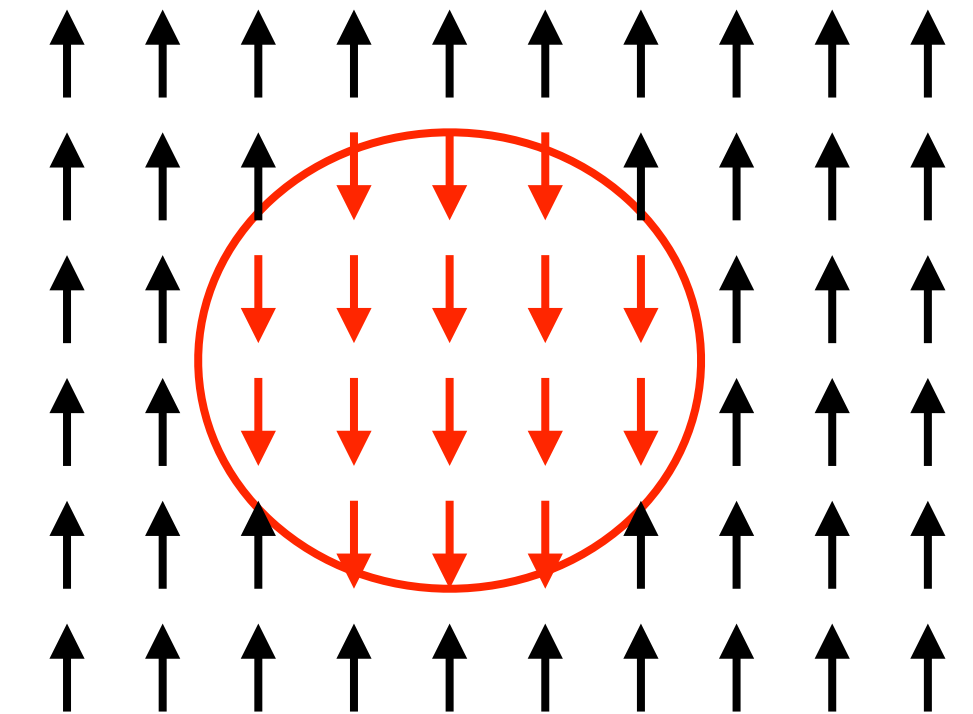


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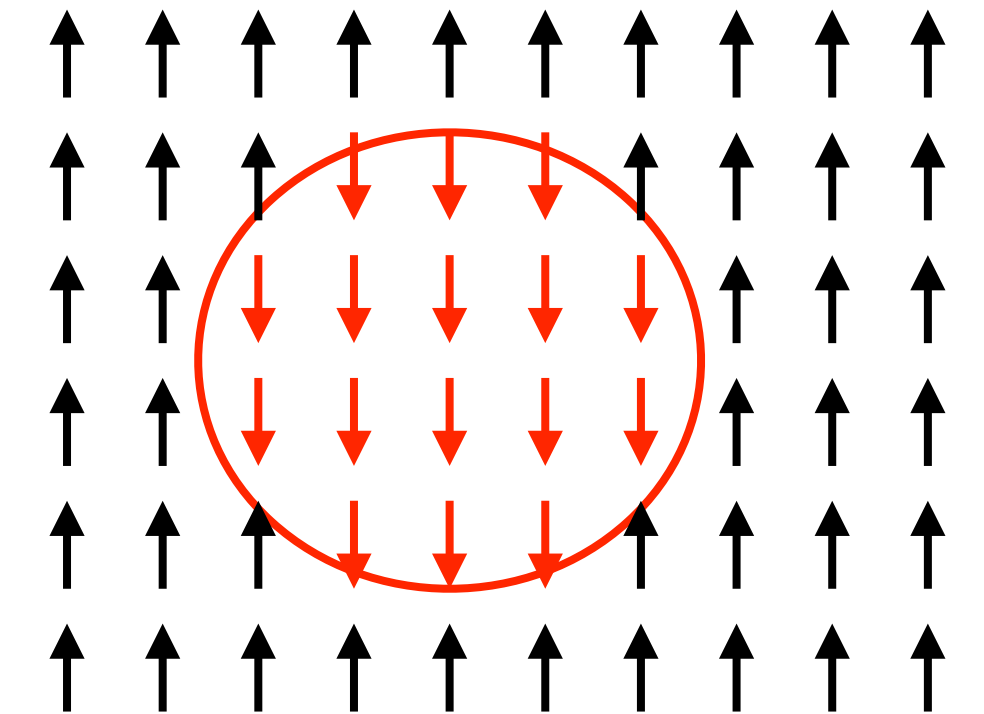


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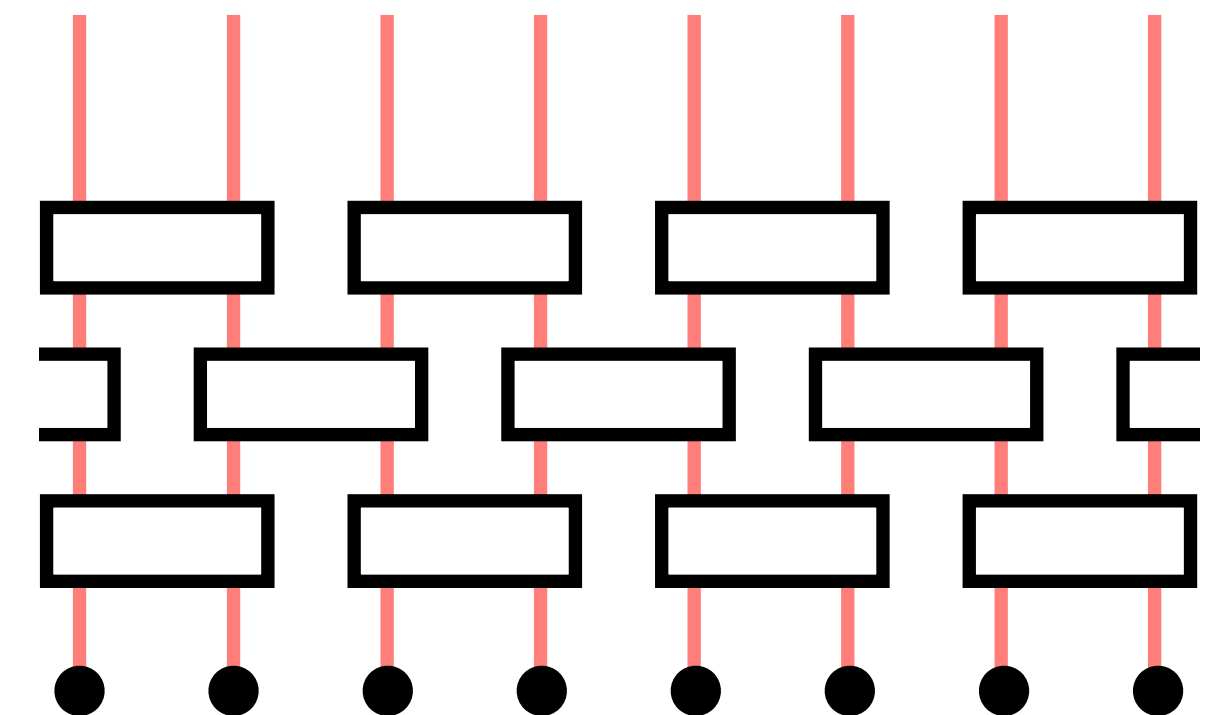
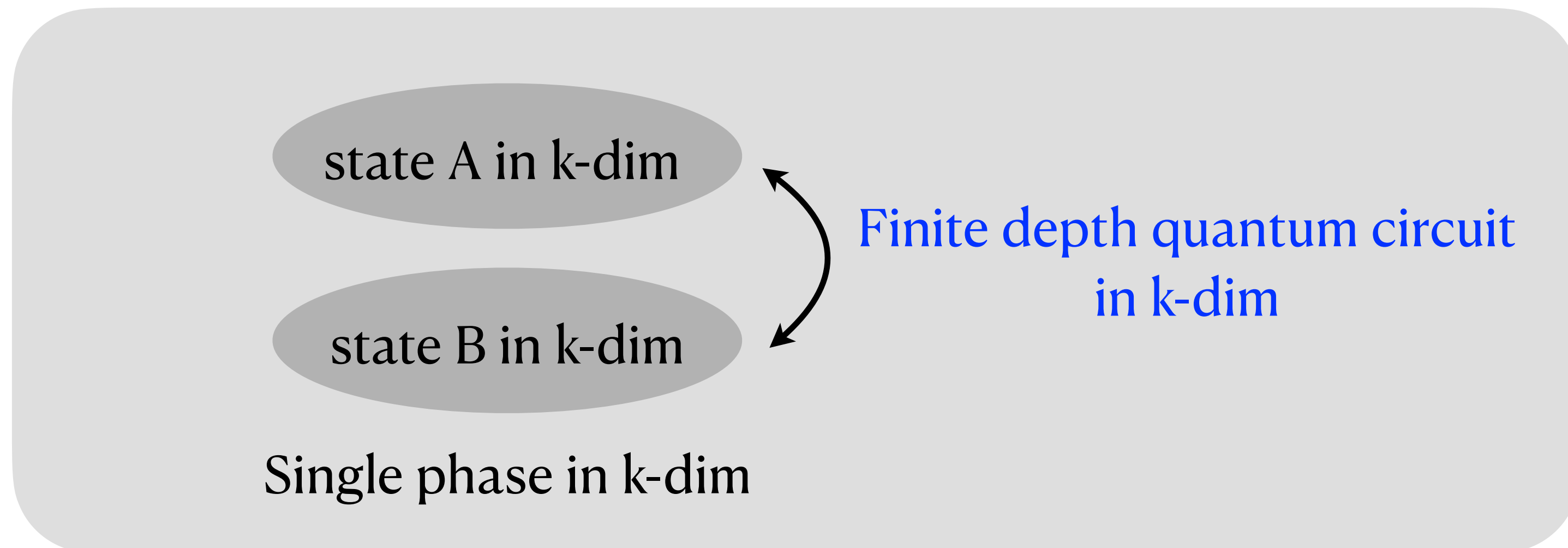
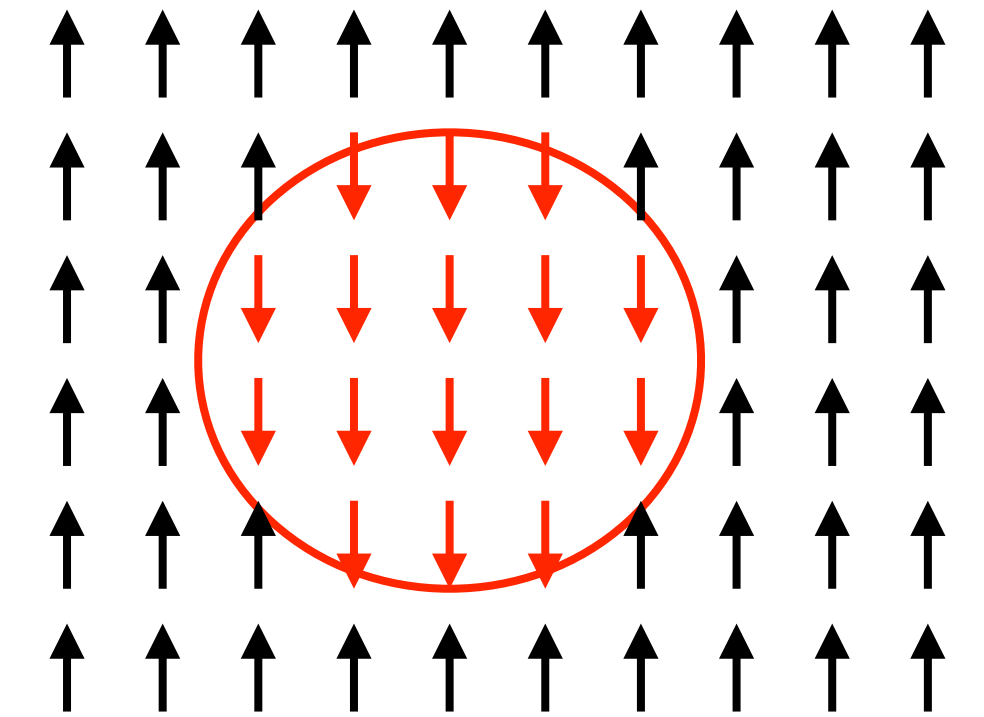


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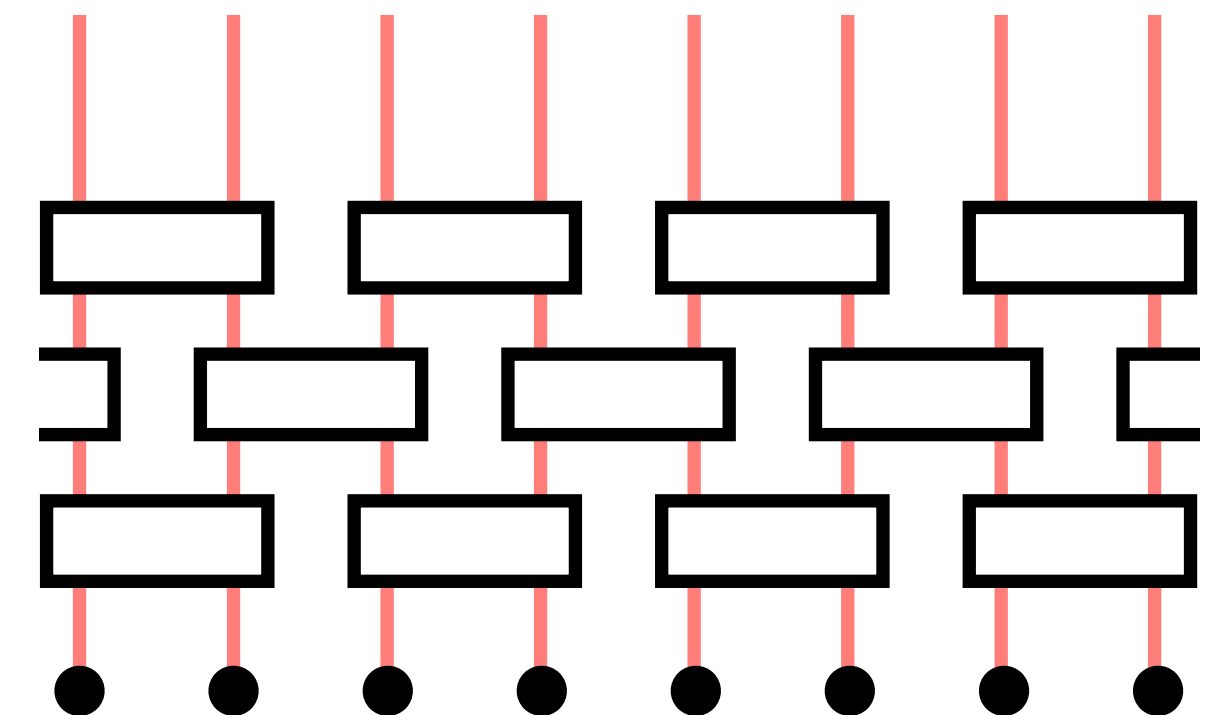
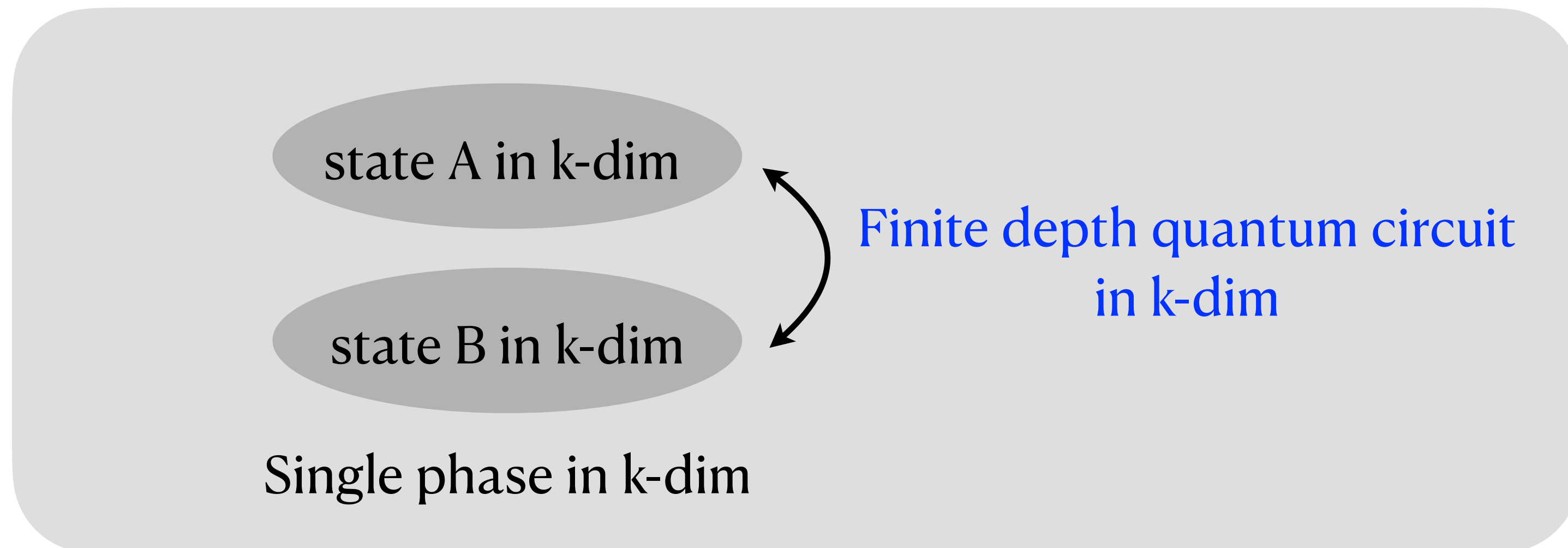
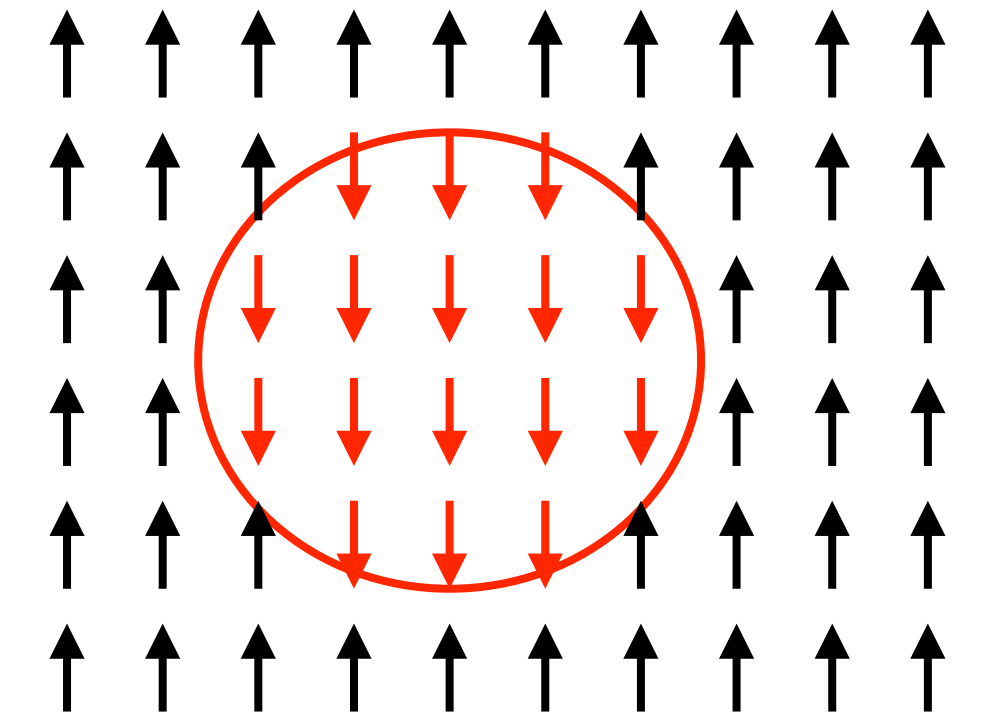
Chen-Gu-Wen

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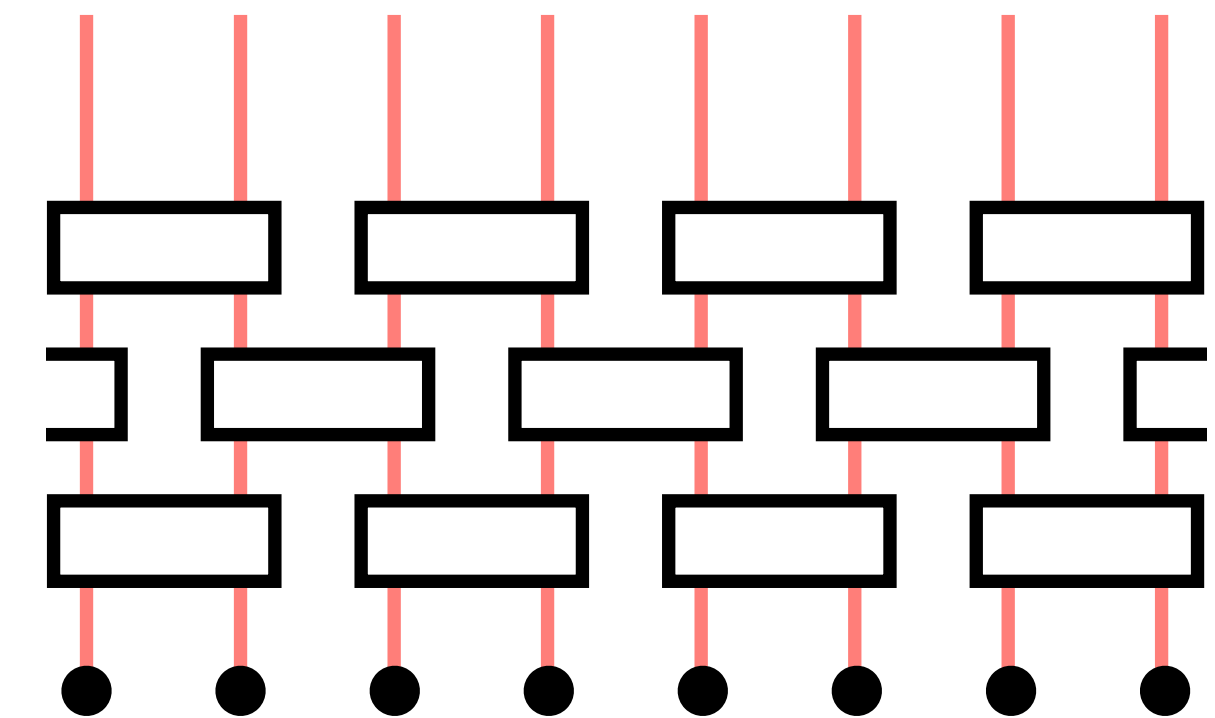
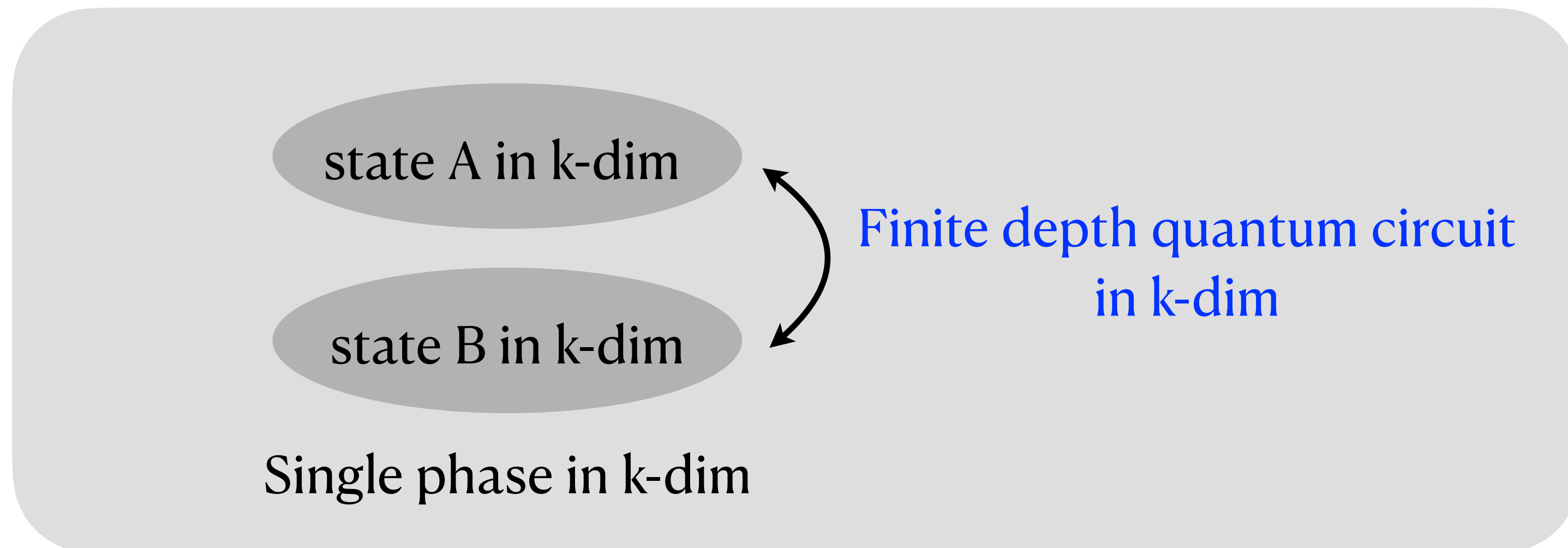
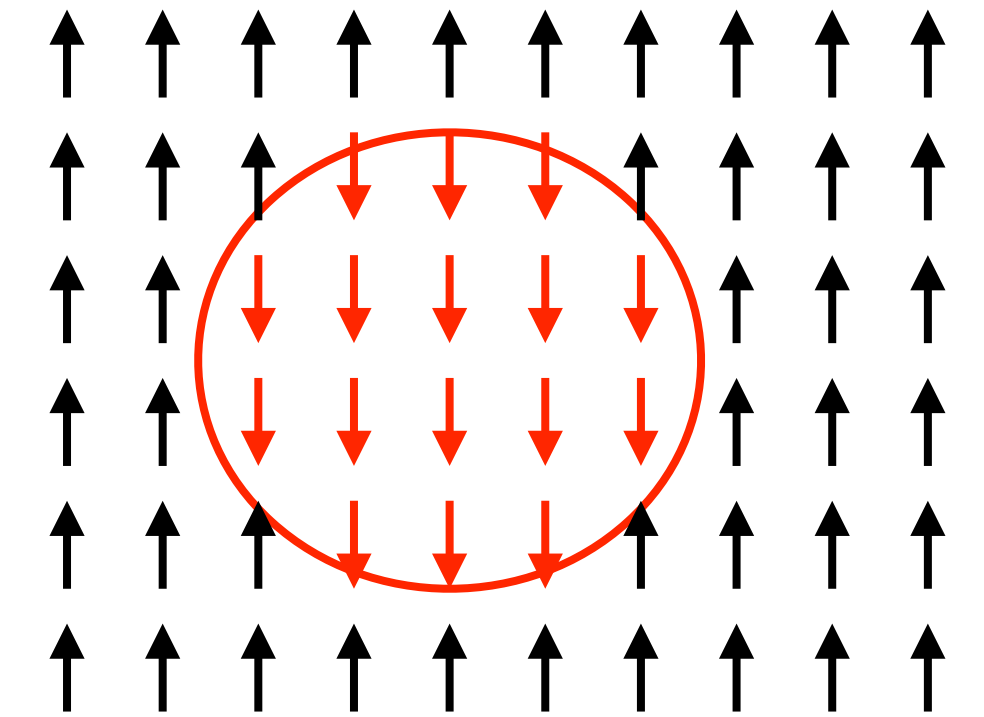
Chen-Gu-Wen

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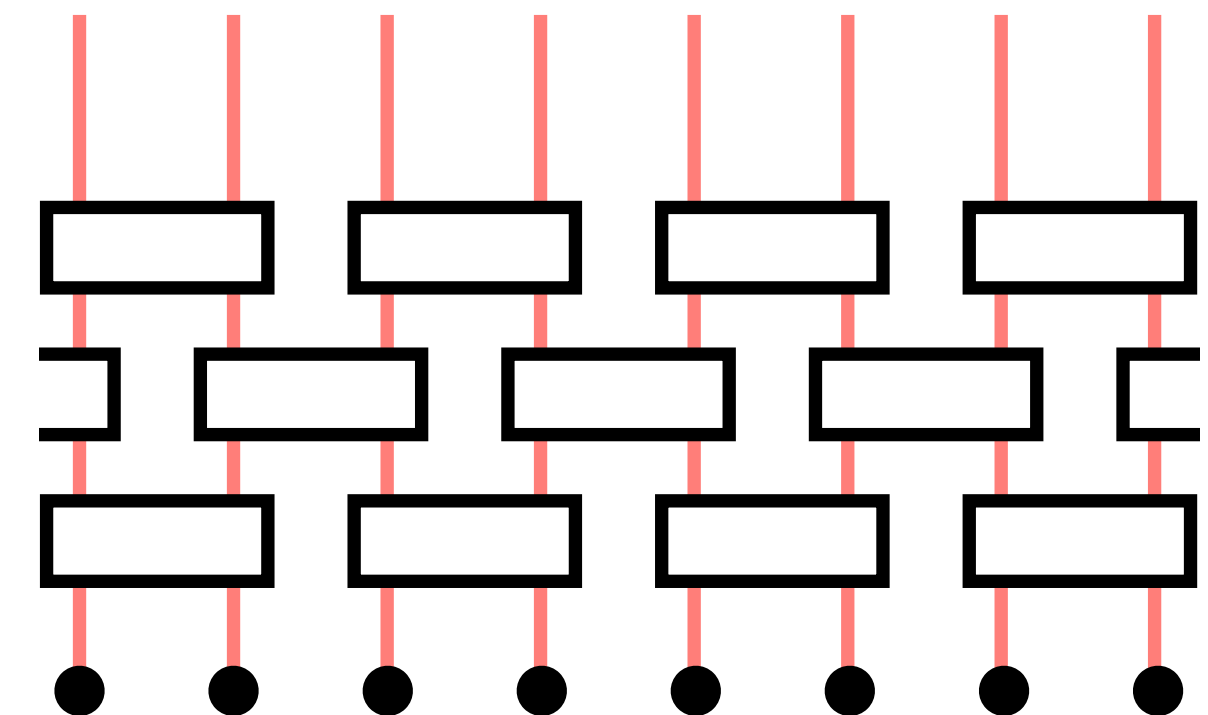
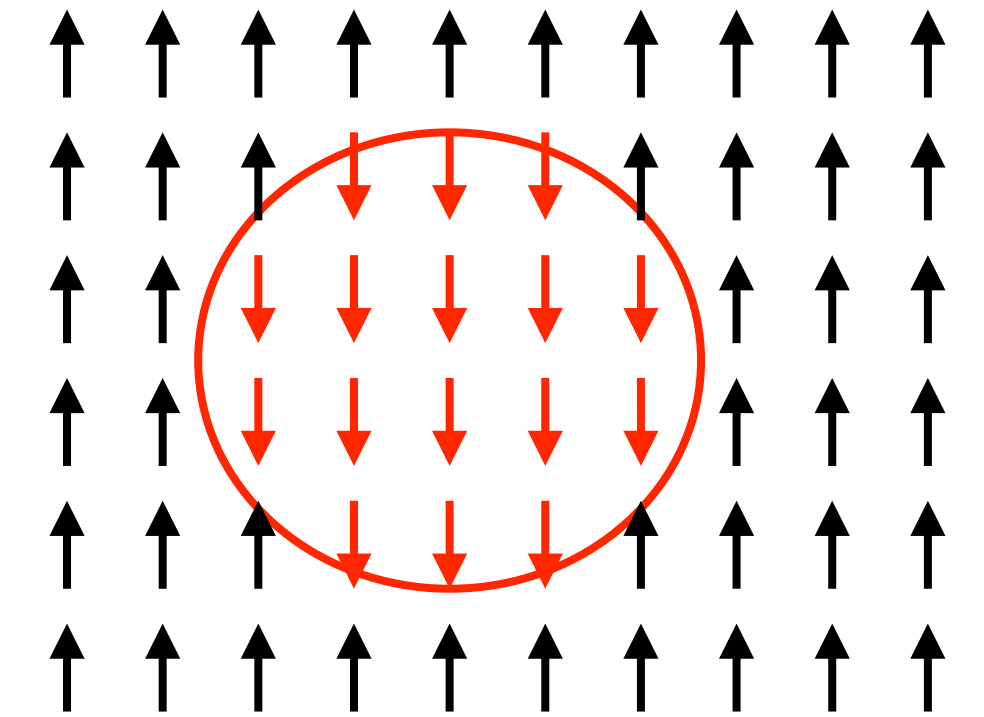
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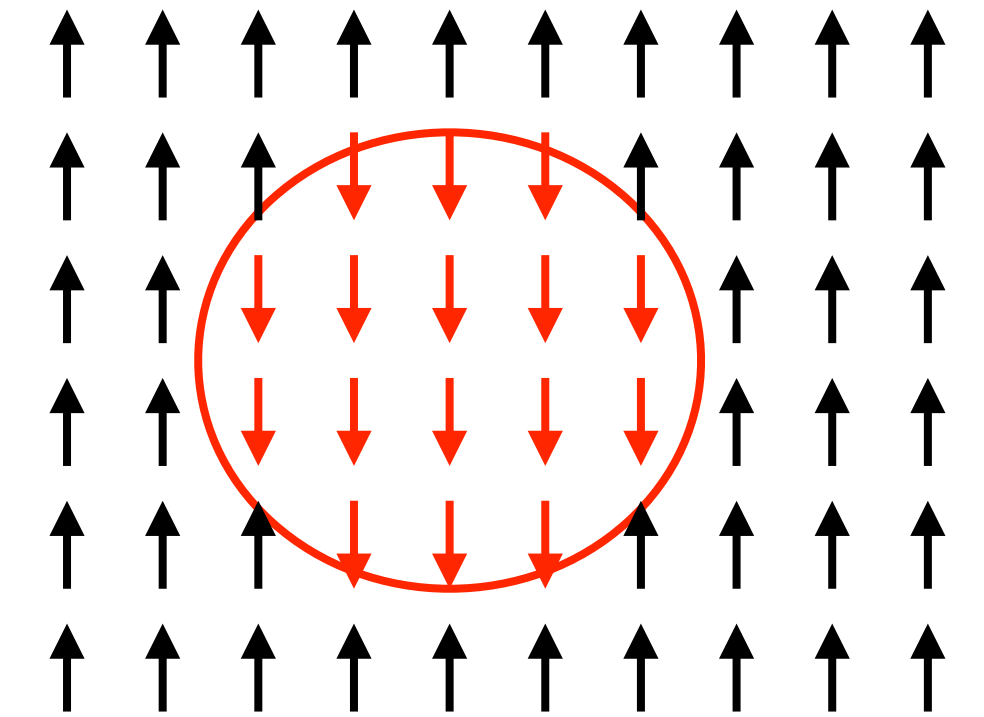


Low entanglement excitations

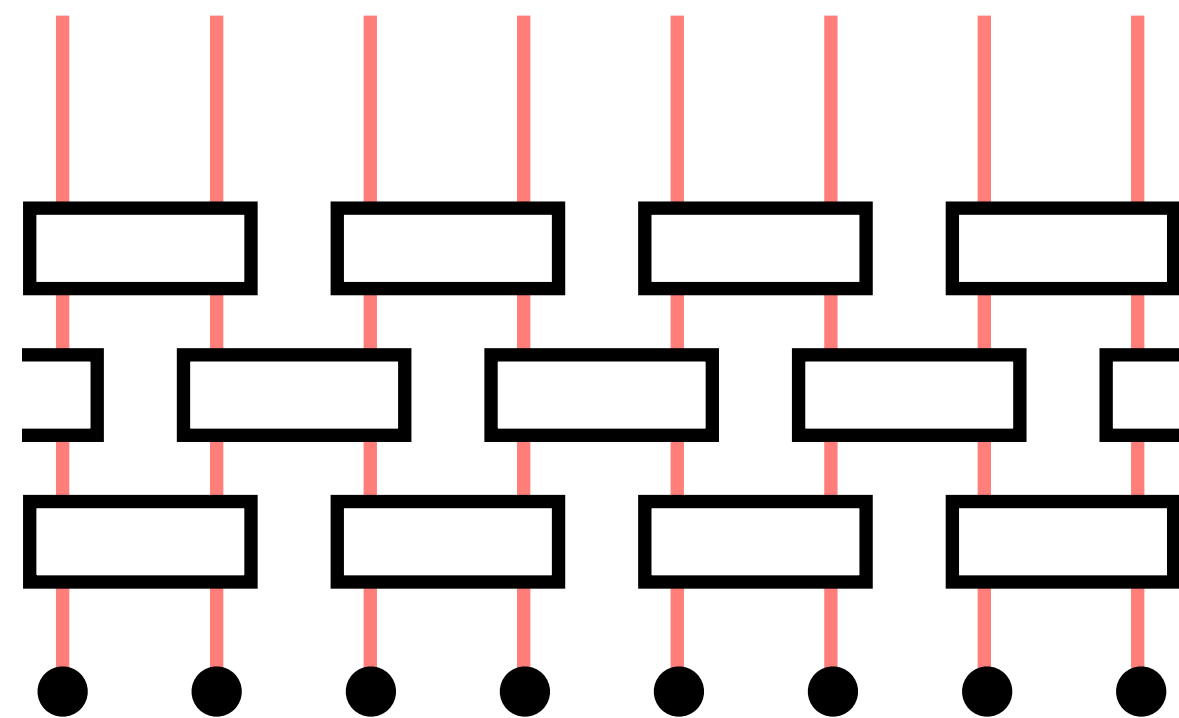
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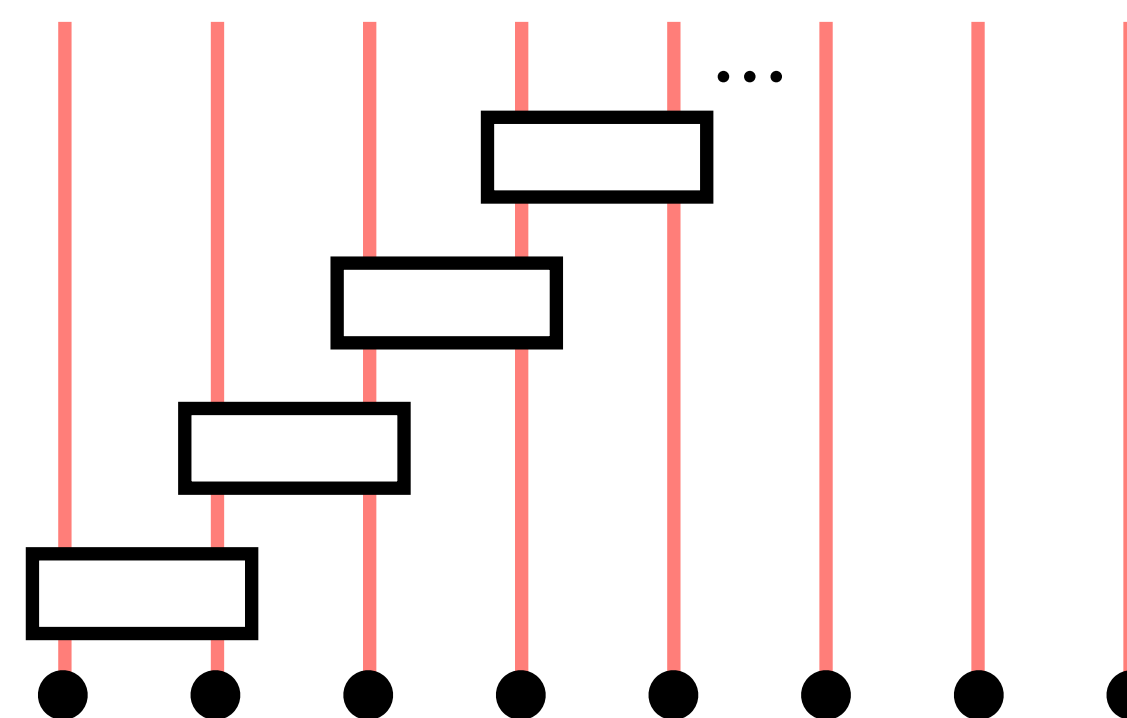
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finite depth quantum circuit



linear depth sequential circuit



Schon-Solano-Verstraete-Cirac-Wolf, 2005; Ho-Hsieh, 2019

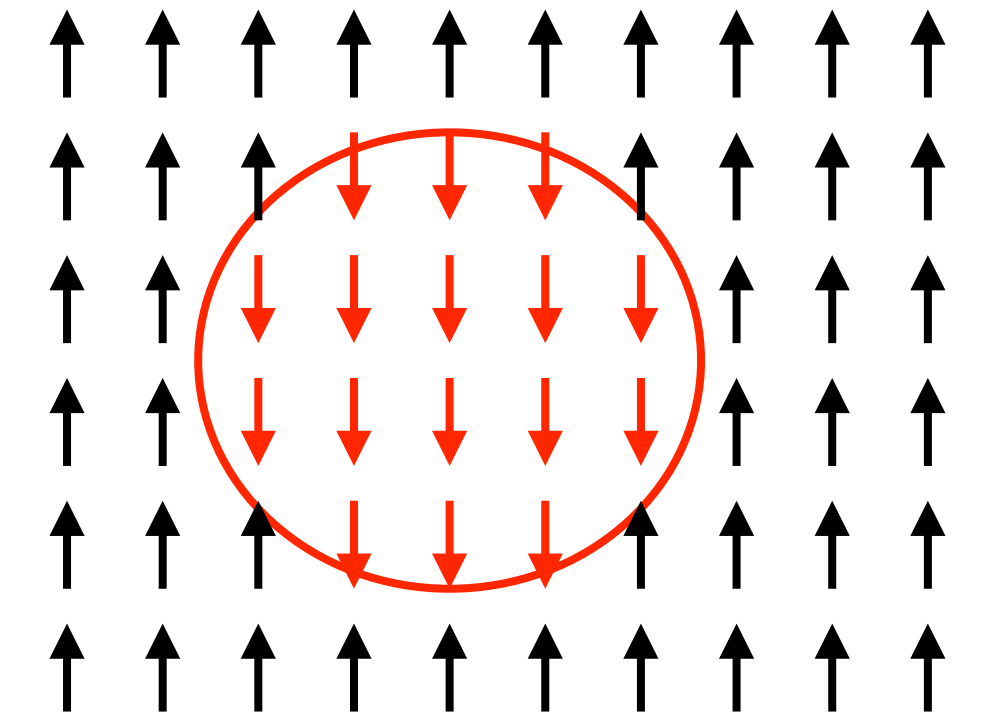
Chen-Dua-Hermele-Stephen-Tantivasadakarn-Vanhove-Zhao, 2023

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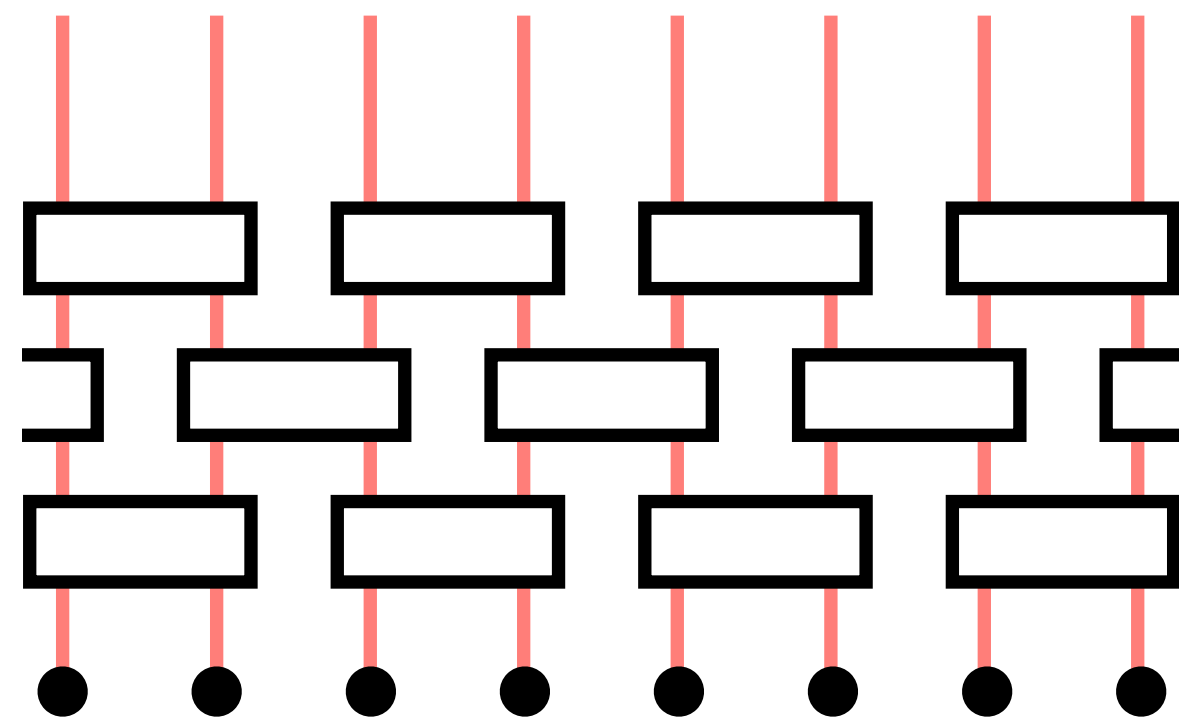
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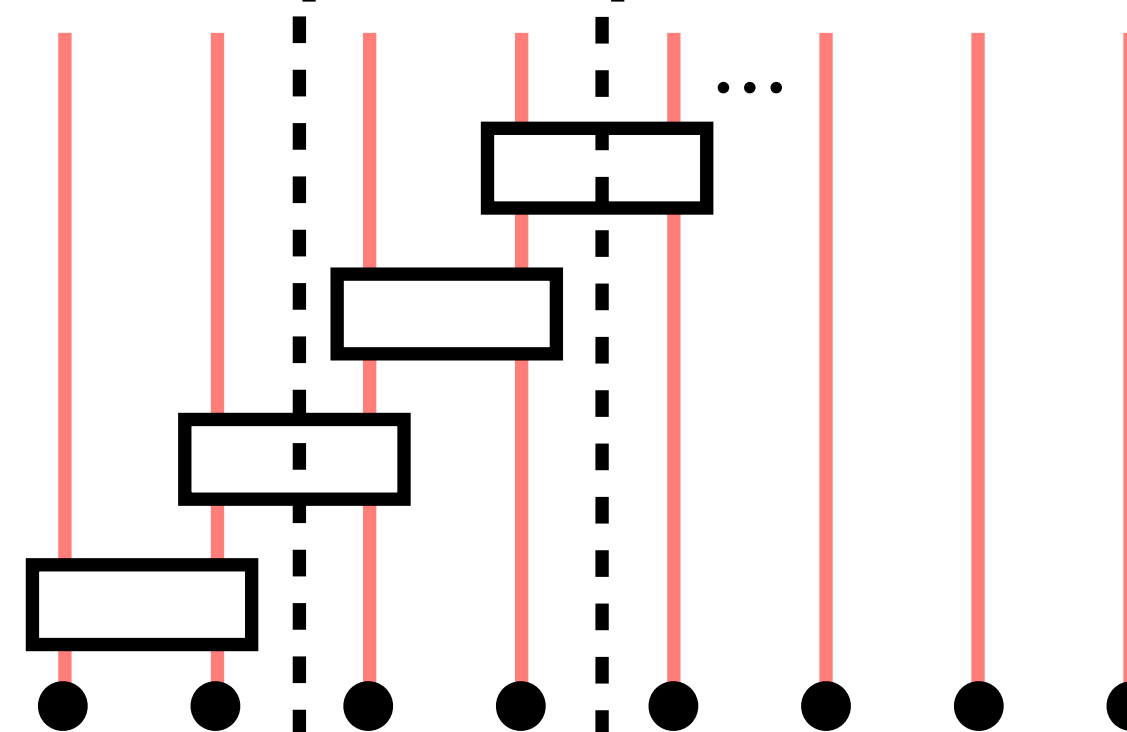
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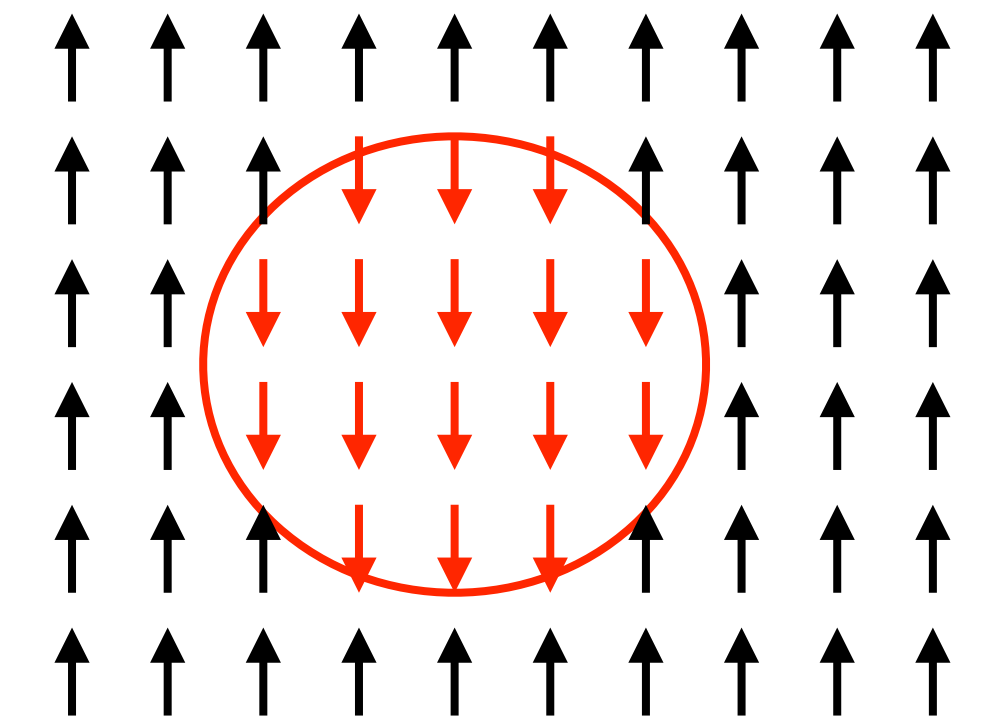
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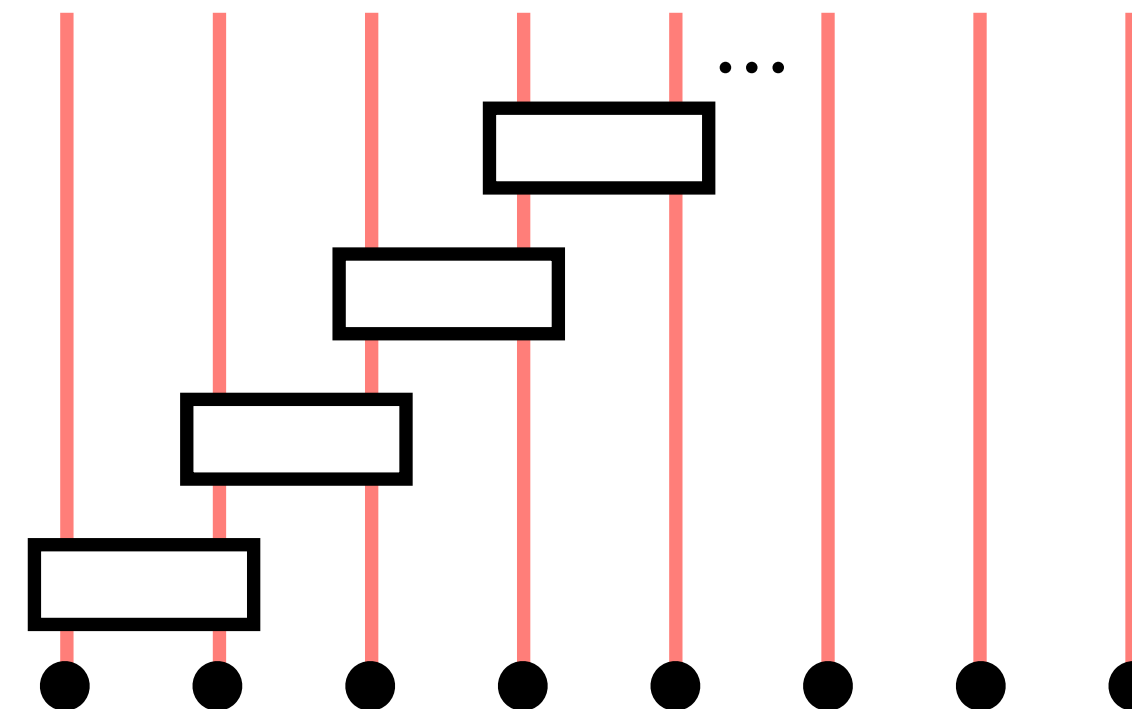
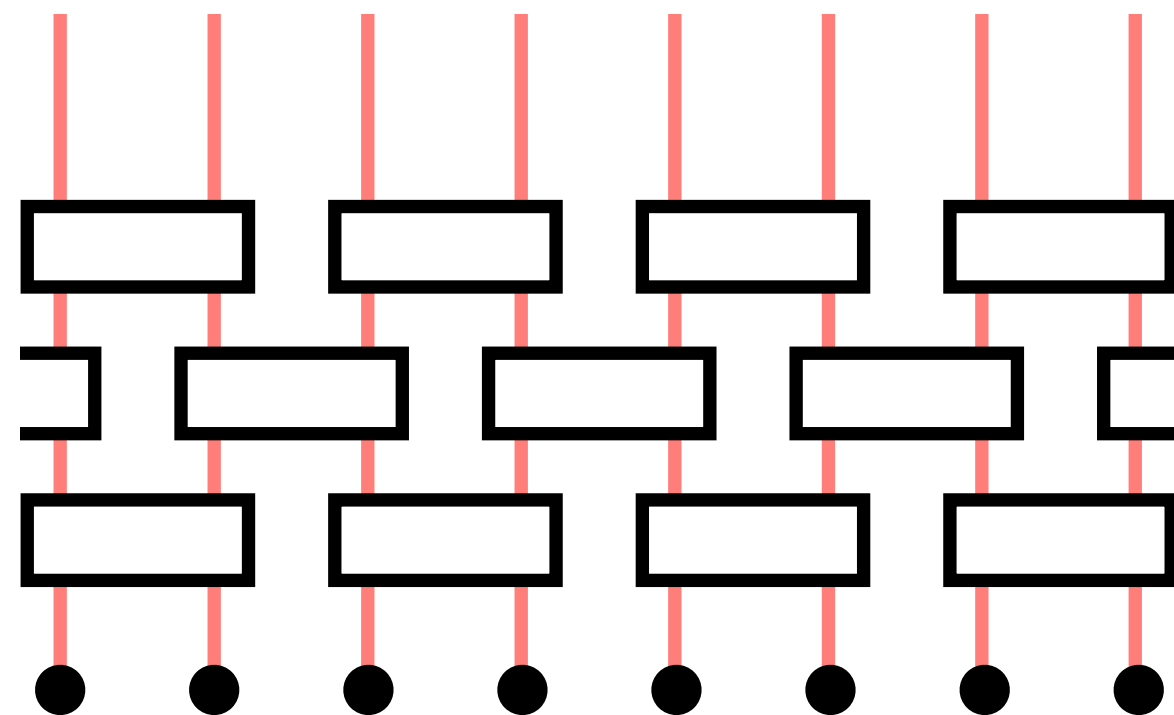
Use **quantum circuits** to create/define k -dimensional excitations

In the ground state of a gapped Hamiltonian in d dimensions, a k -dim. excitation

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- **Trivial types** are created from bulk ground state by a **k -dim** (symmetric) **FDQC**.
- **Non-trivial types** cannot be created by a k -dim (symmetric) FDQC;
can be created with a **$(k+1)$ -dim** (symmetric) **quantum circuit**,
or a **k -dim** (symmetric) **linear depth sequential circuit**.

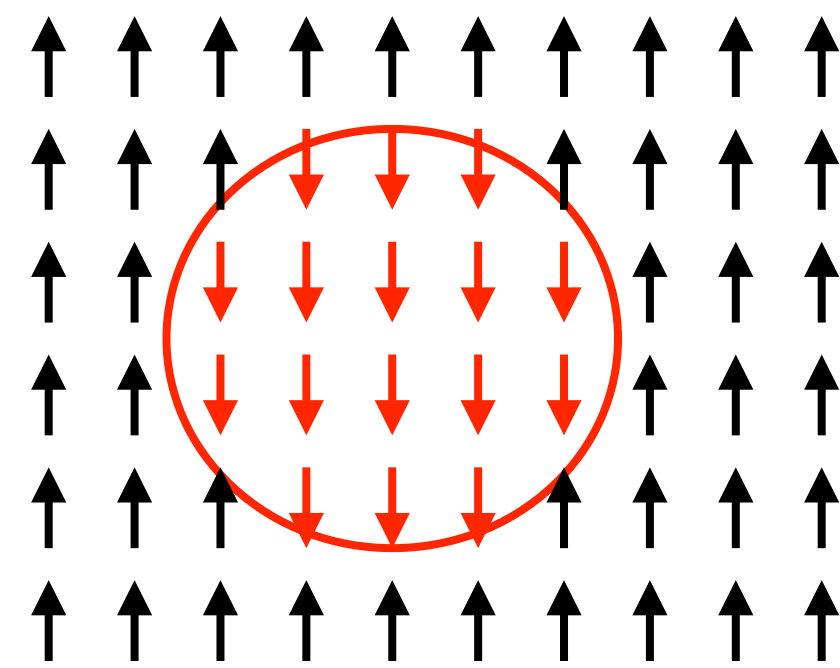


1d excitations
Hamiltonian modified
along a loop
2d FDLU

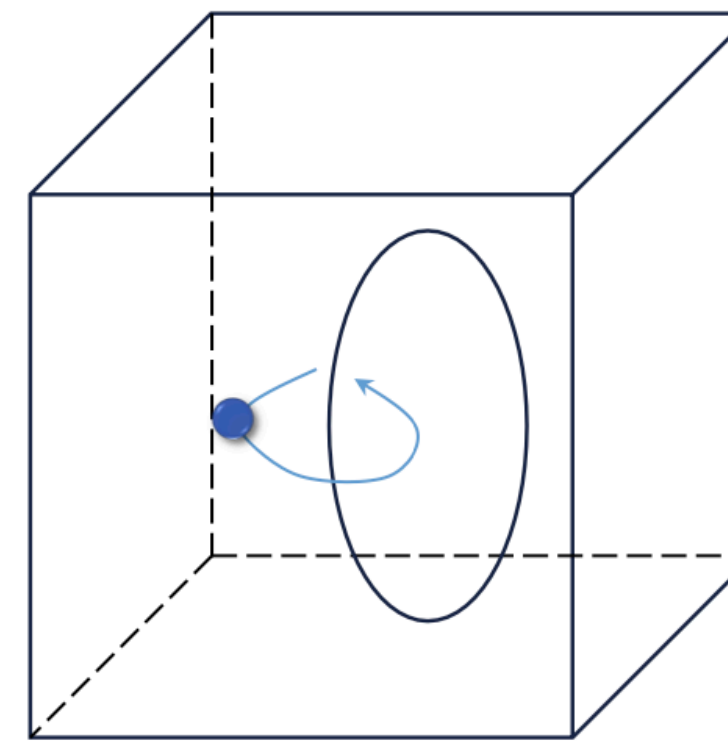


Low entanglement excitations

Line excitations created by 2d FDQC



1d excitations
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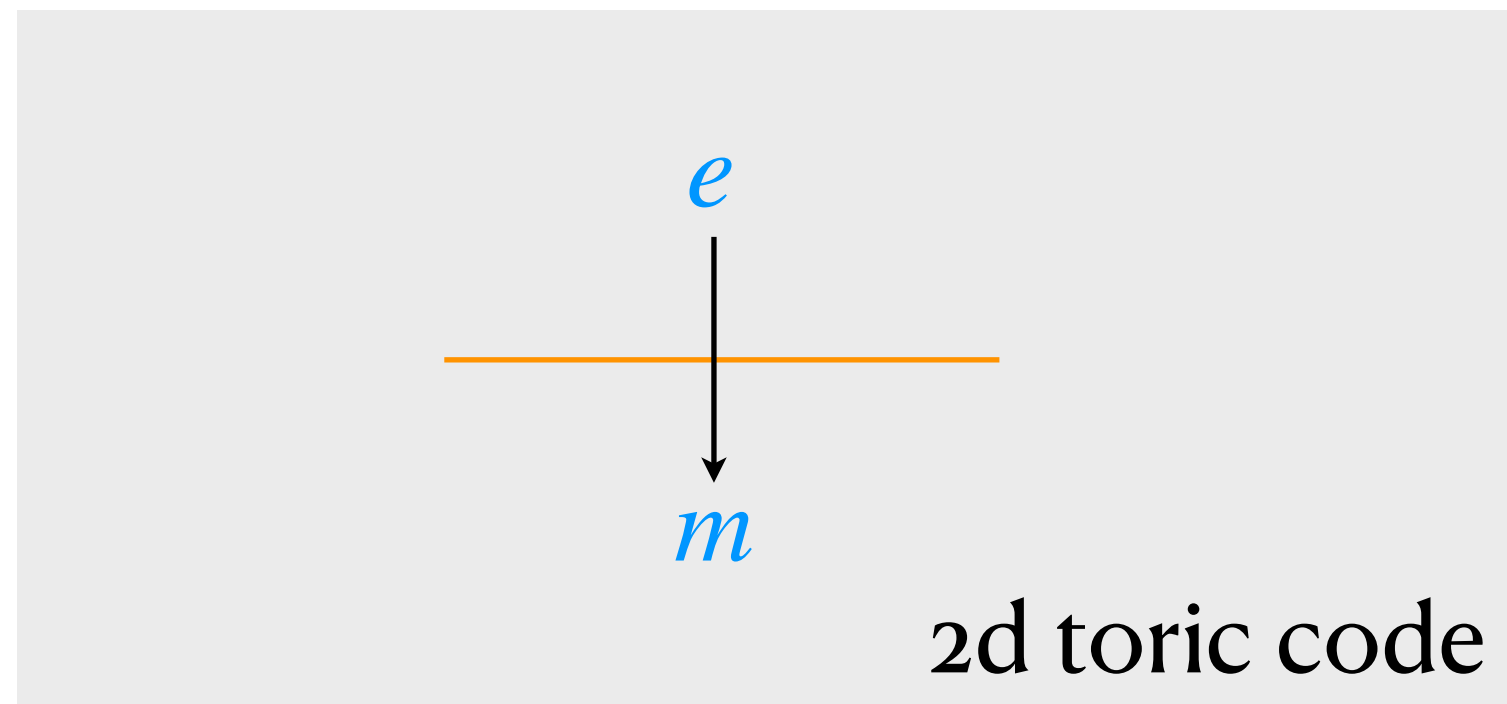


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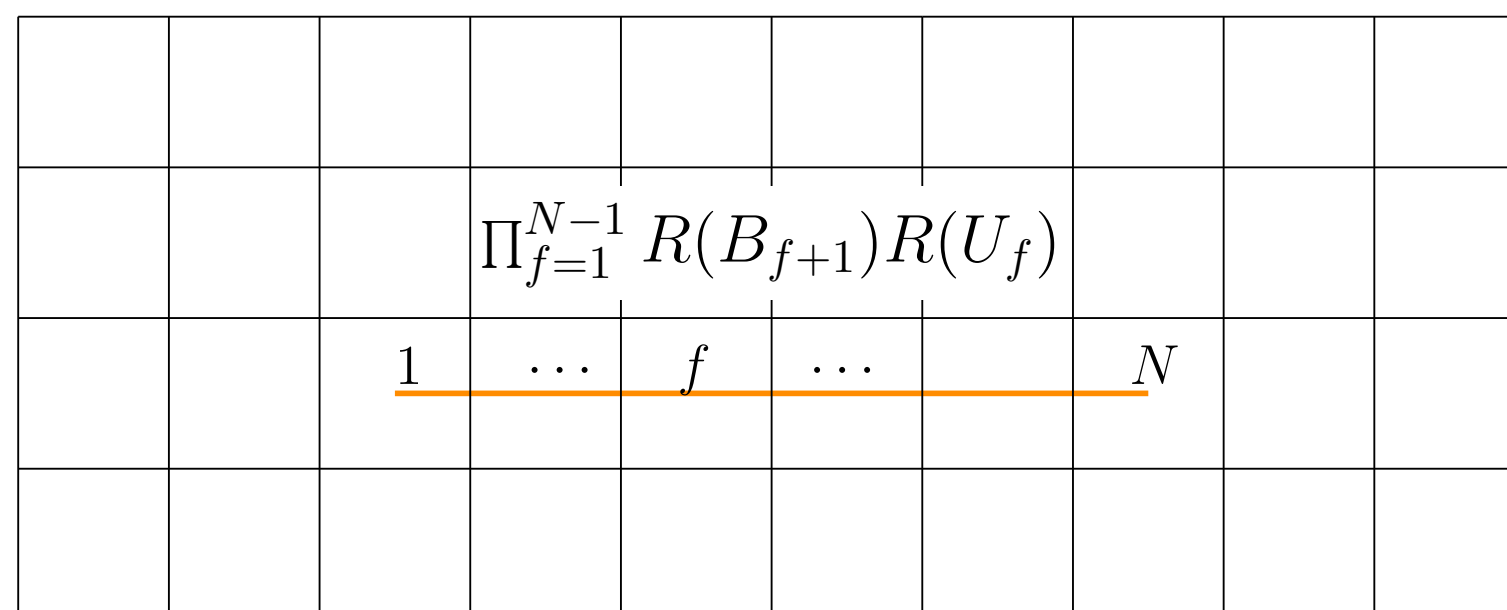
Low entanglement excitations

Line defect exchanging anyons

e-m exchange defect in toric code

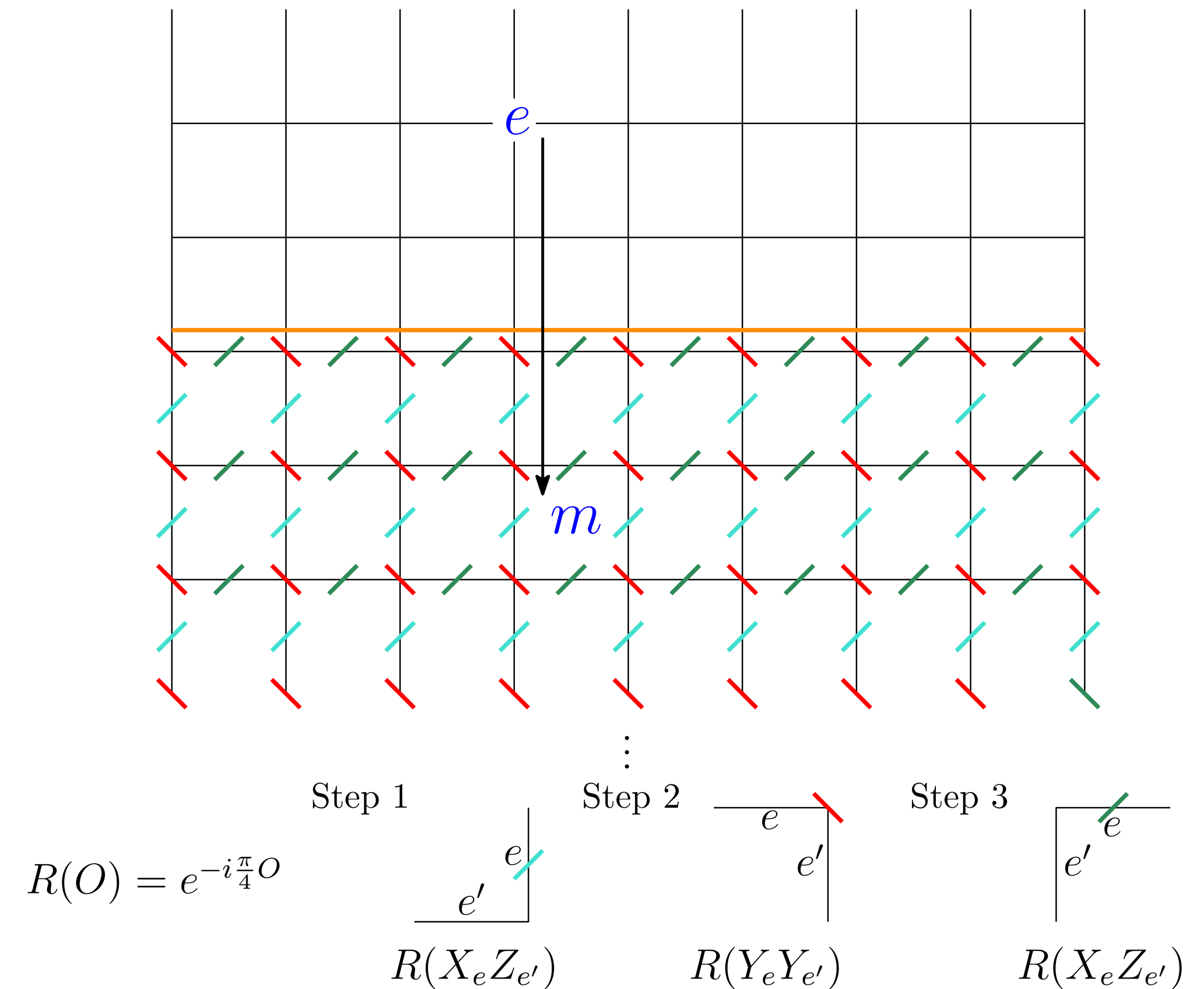


1D linear depth sequential circuit



$$U_f = \begin{array}{c|c} f & \\ \hline & Z \end{array} X \quad B_f = \begin{array}{c|c} Z & f \\ \hline Z & \end{array} Z$$

2D FDLU ("pump" unitary)



Low entanglement excitations

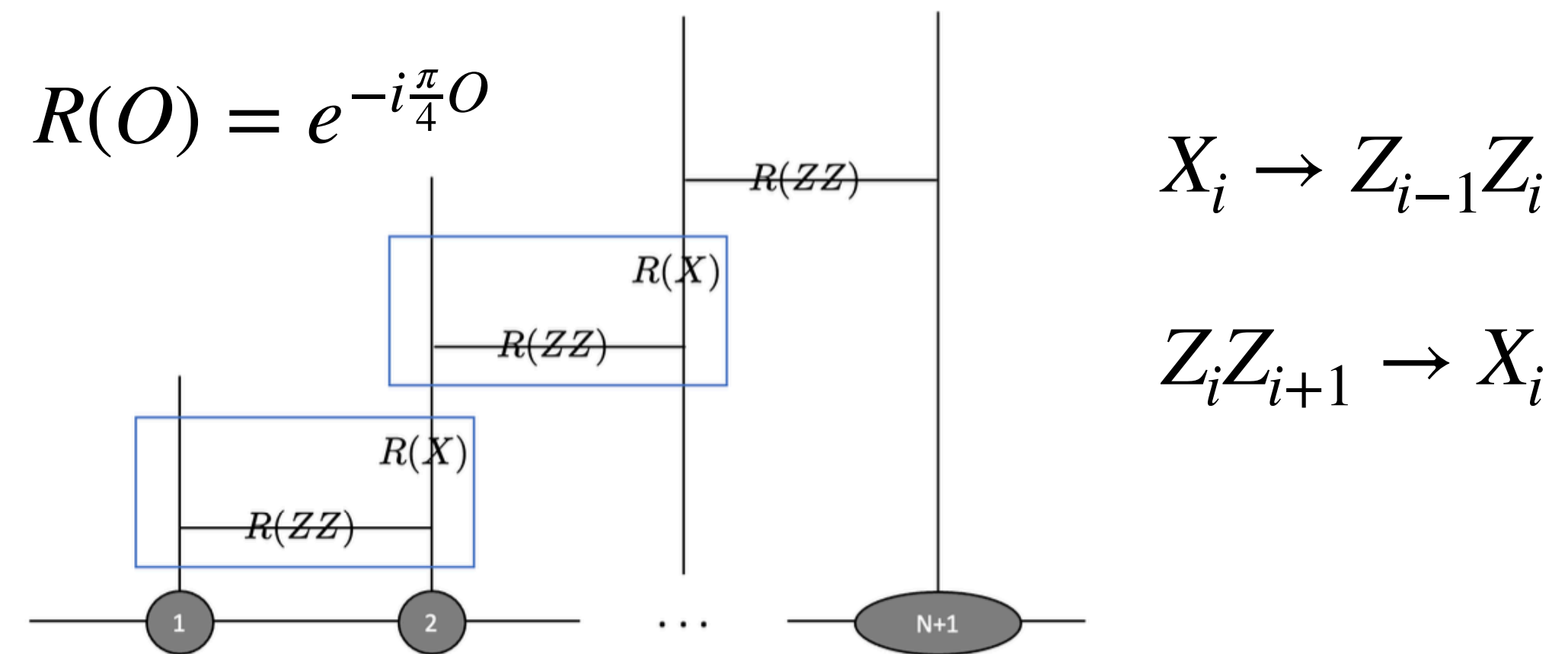
Line defect in a trivial product state

1d GHZ state (SSB)

$$|\cdots 000 \cdots\rangle + |\cdots 111 \cdots\rangle$$

2d symmetric state $|+\rangle^{\otimes N}$

1D sequential circuit



Schon-Solano-Verstraete-Cirac-Wolf, 2005; Ho-Hsieh, 2019

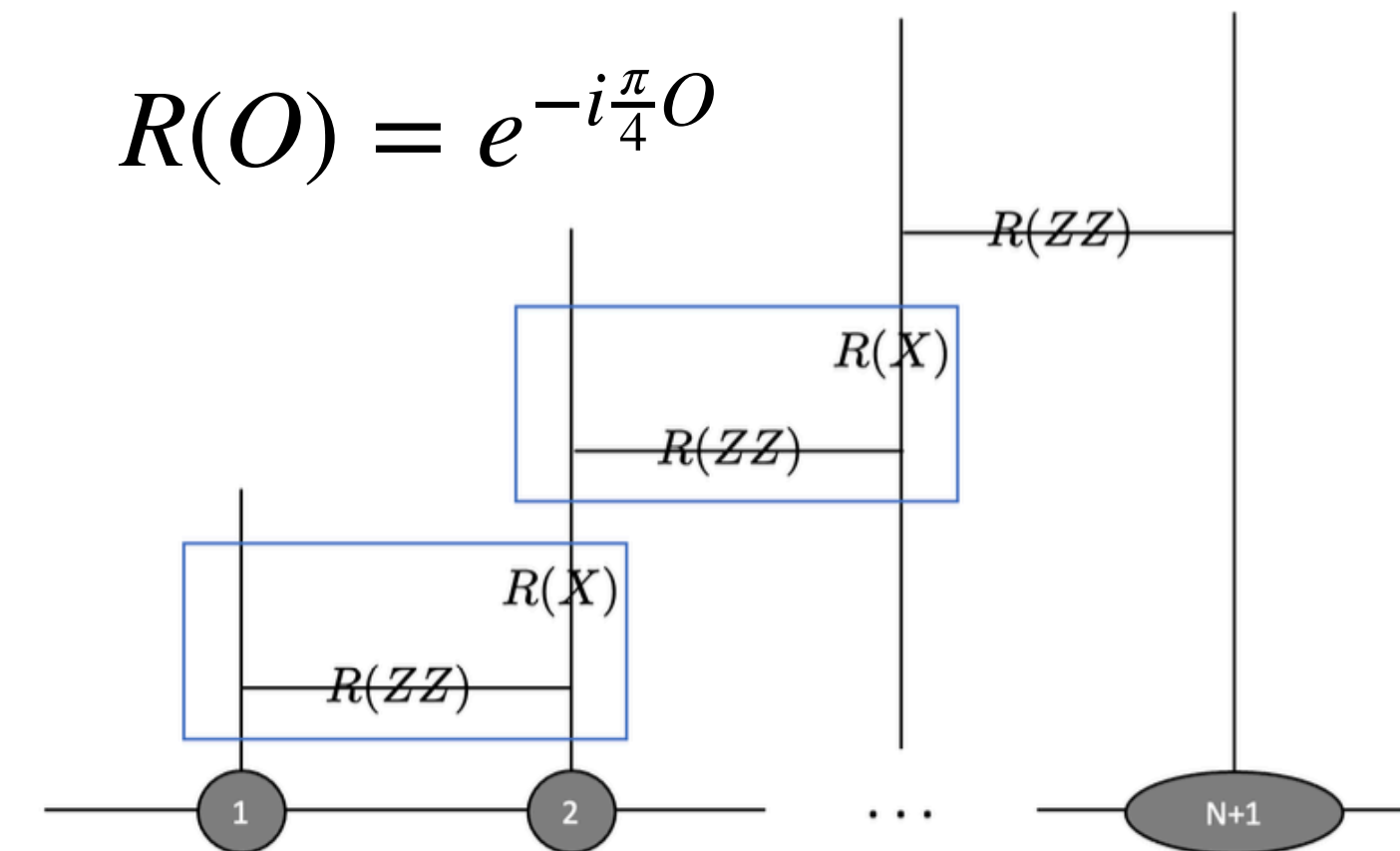
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Classifying Low entanglement excitations?

Low entanglement excitations

1d GHZ state (SSB)

2d symmetric state



All excitations in a 2d trivial product state

Ground state		2d \mathbb{Z}_2 paramagnet
Excitations	0d	\mathbb{Z}_2 charge (1d FDLU)
	1d	SSB state (1d SQC)

What about an entangled ground state?

LEE in an Invertible phase

Claim: Low entanglement excitations in an invertible phase and those in a product state have a one-to-one correspondence.

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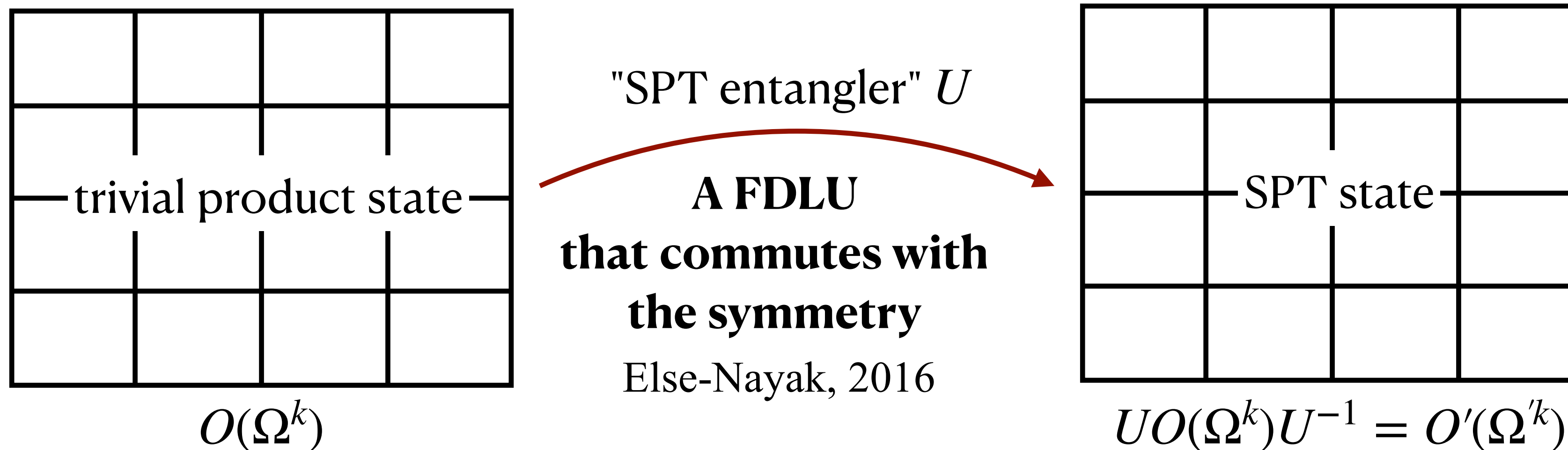
Ground state		2d \mathbb{Z}_2 paramagnet	2d \mathbb{Z}_2 SPT
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	1d	1d \mathbb{Z}_2 SSB 1d SQC	1d \mathbb{Z}_2 SSB 1d SQC

Bulk excitations in an Invertible phase

Claim: Low entanglement excitations in an invertible phase and those in a product state have a one-to-one correspondence.

For many examples, already easy to show.

Example SPT phases



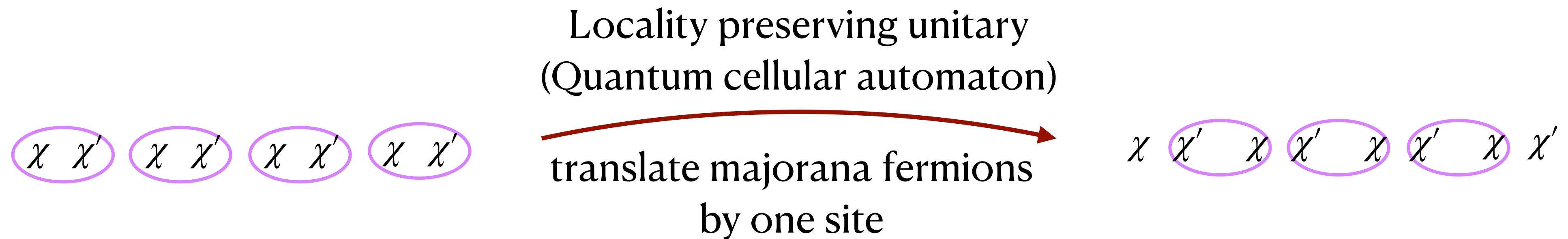
- Symmetric operators have one-to-one correspondence, including LEE

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Example 1d Kitaev chain (fermion parity symmetry)



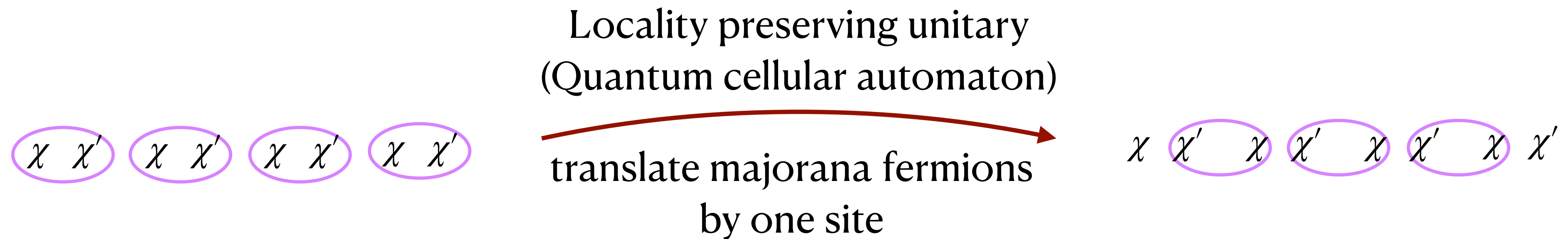
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- Symmetric operators have one-to-one correspondence, including LEE

SPT entangler approach does not apply to examples such as $p + ip$ superconductor

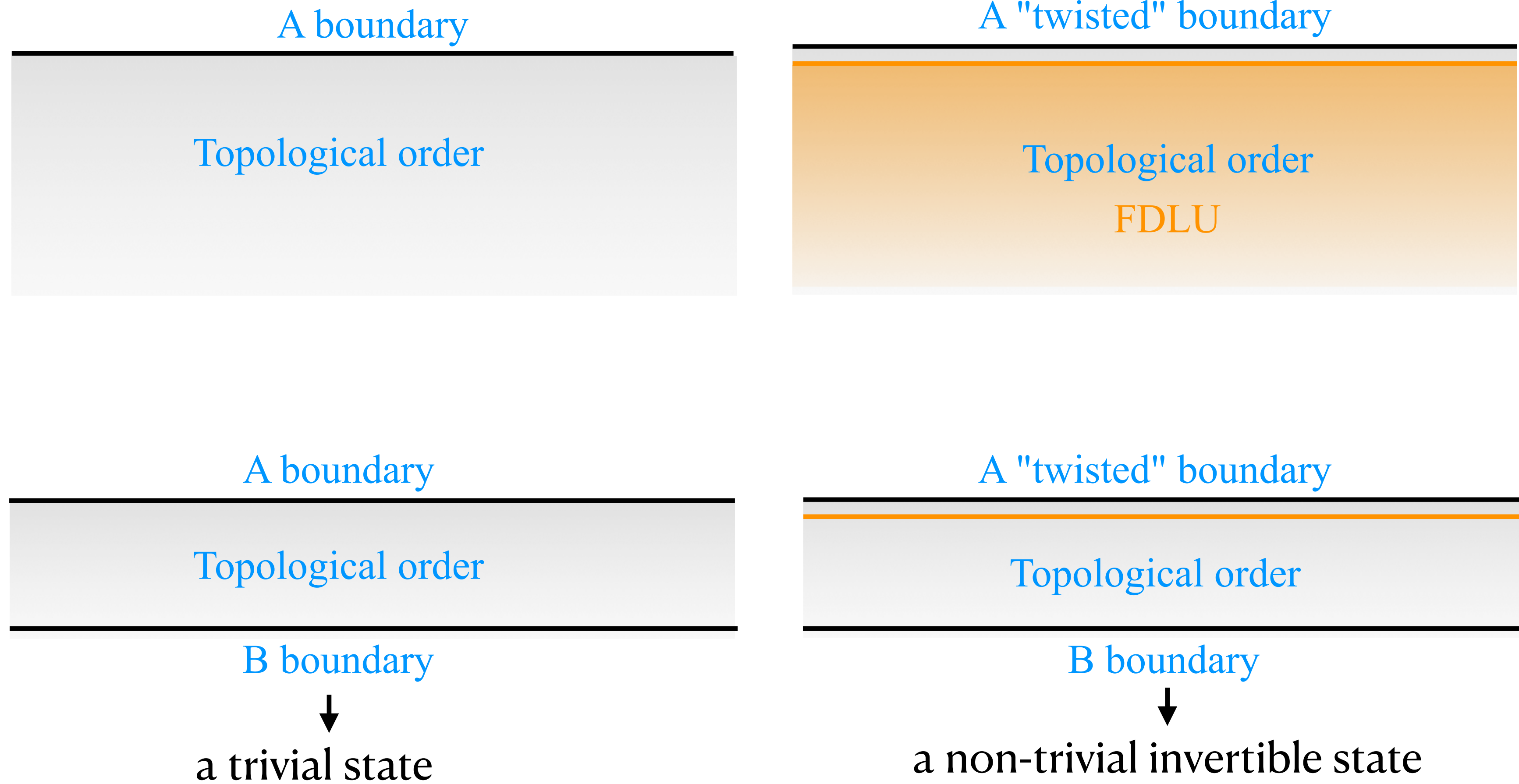
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Argument using "Symmetry TO/TFT"

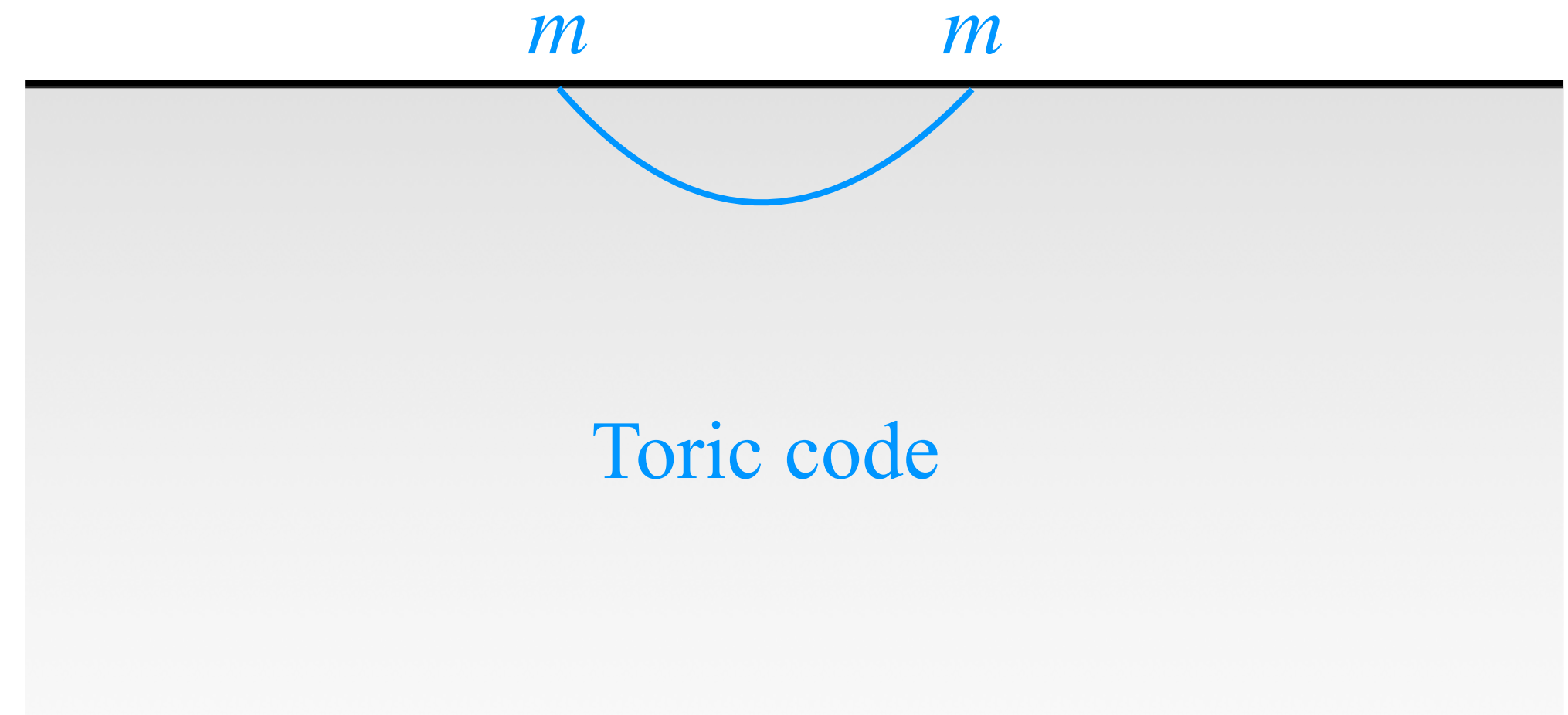
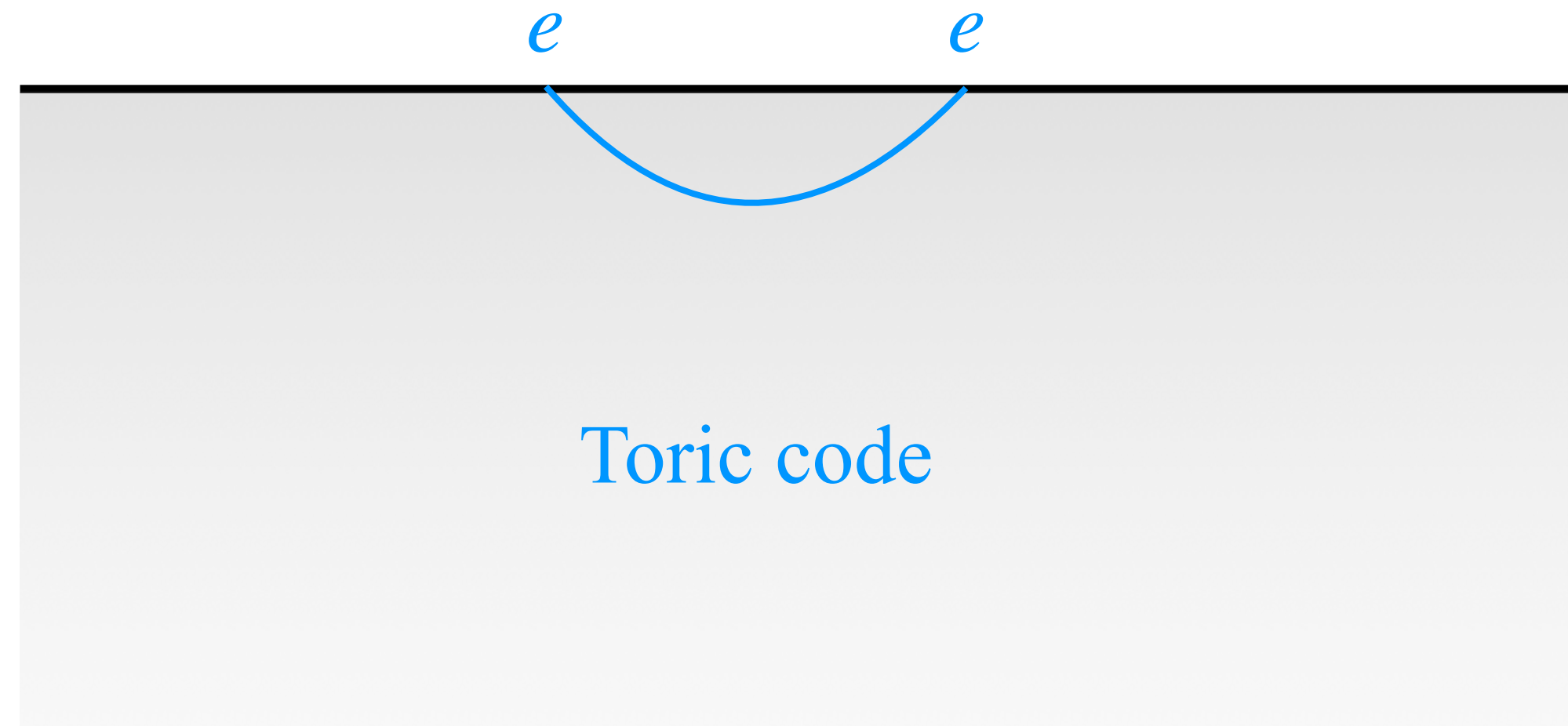
Argument using "Symmetry TO/TFT"

Step 1: Compare Boundaries of topological orders

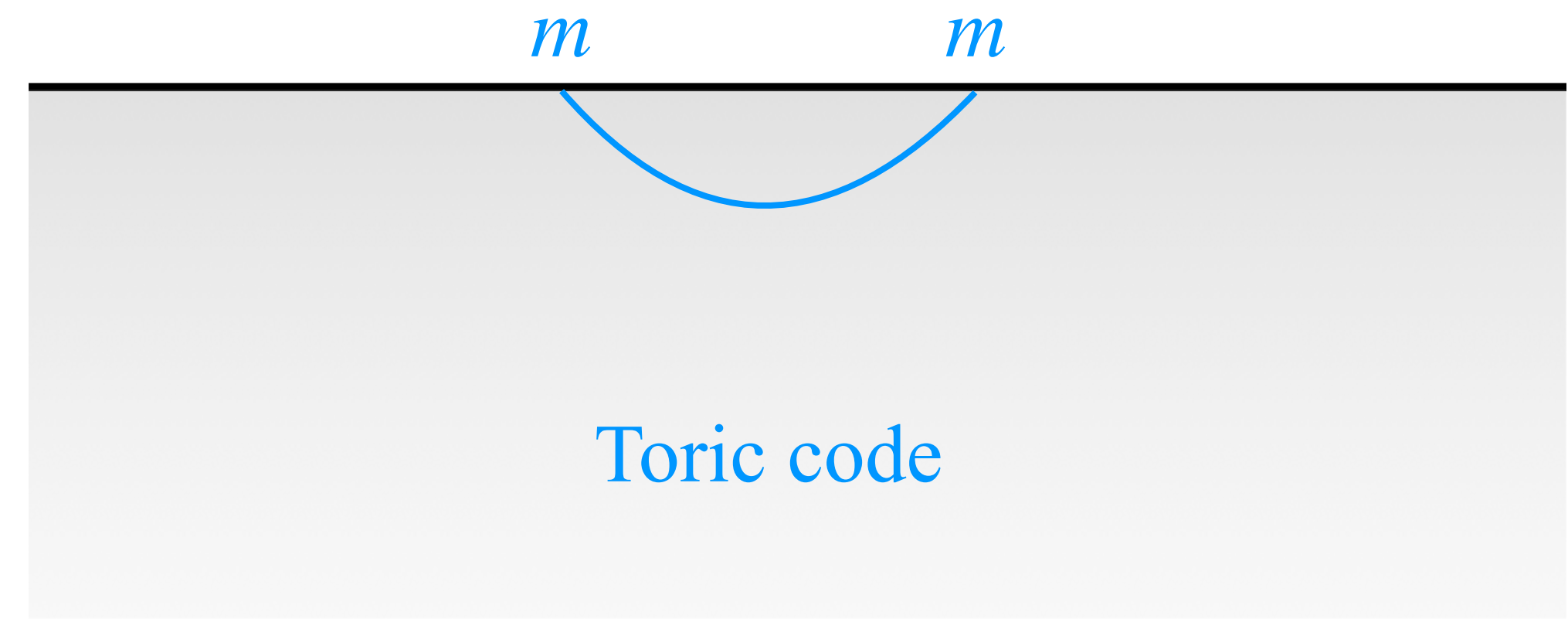
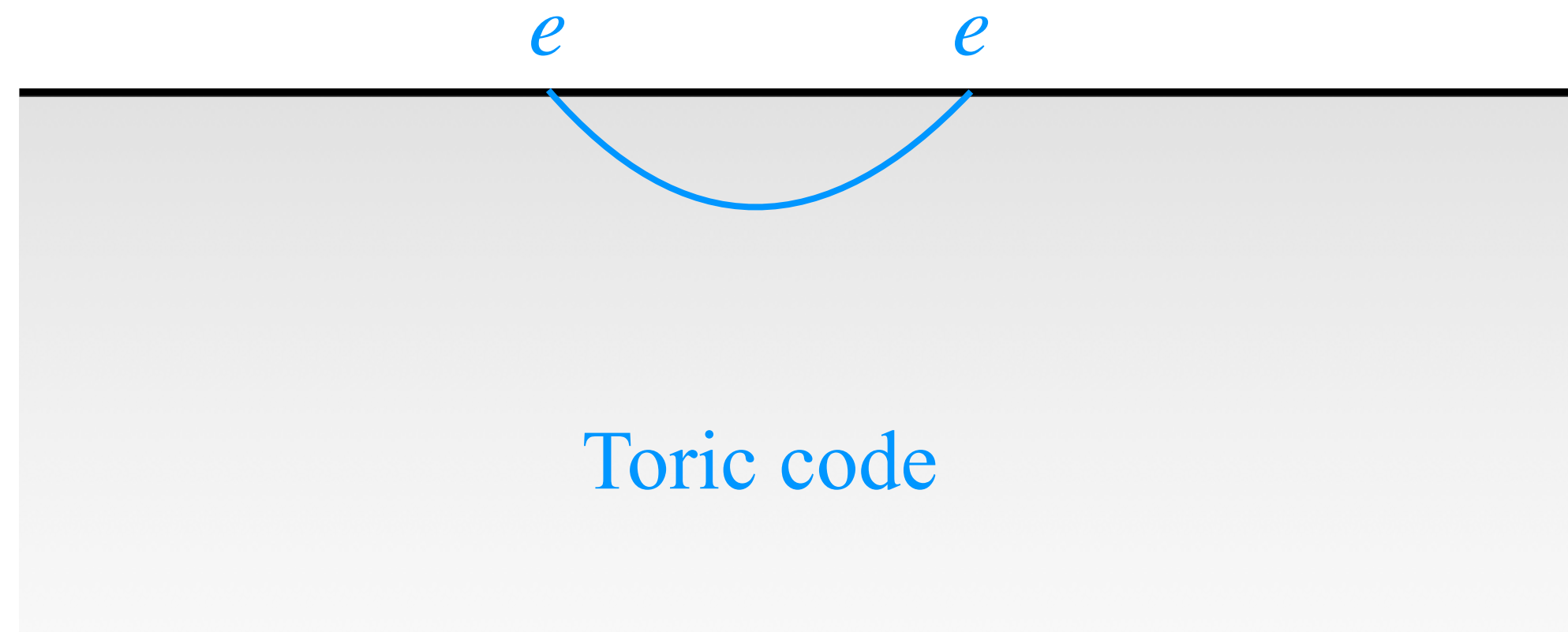


Step 3: Compare excitations

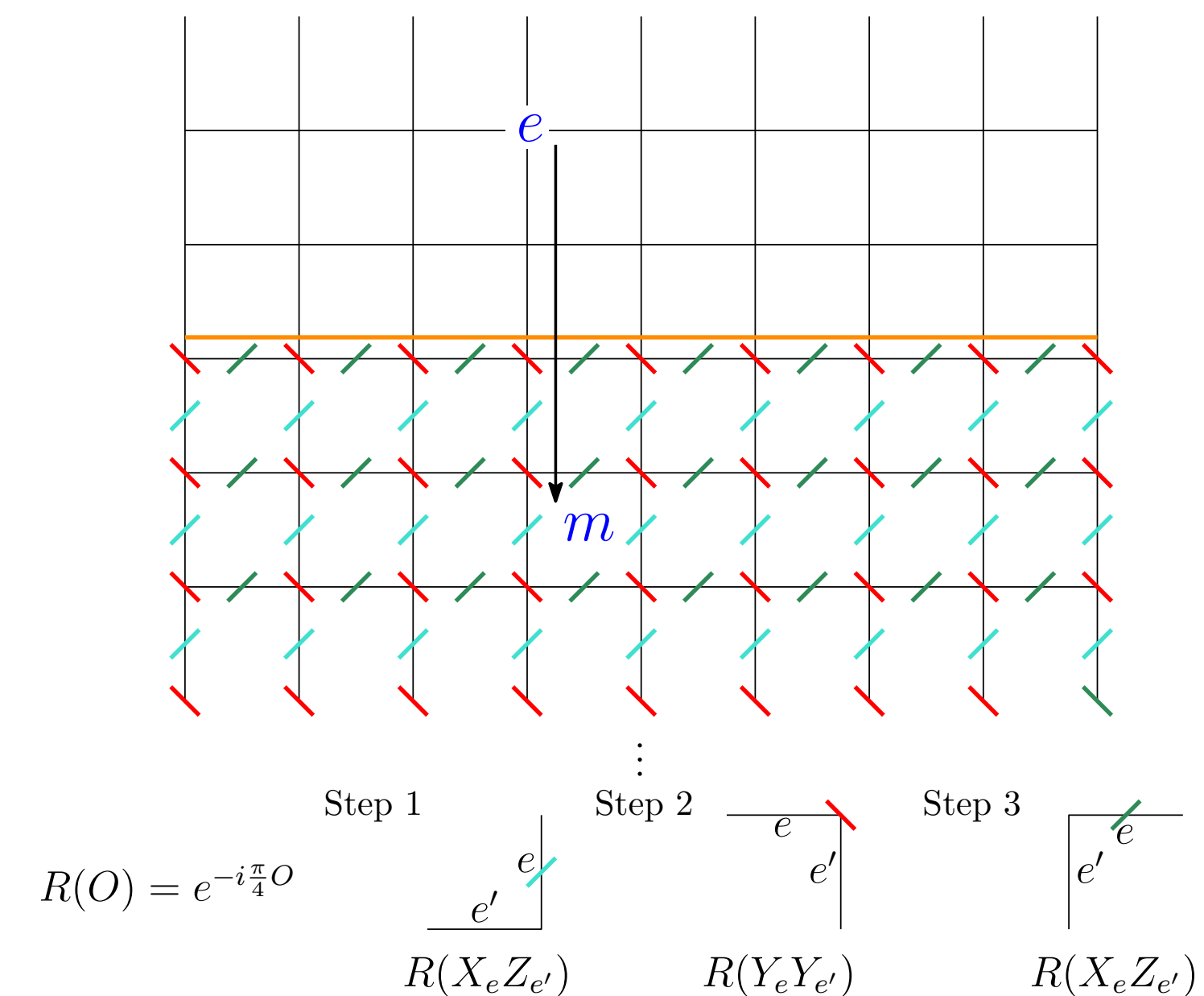
Topological order with a boundary



Topological order with a boundary



Emergent symmetry in toric code bulk:
anyon exchange symmetry $e \leftrightarrow m$

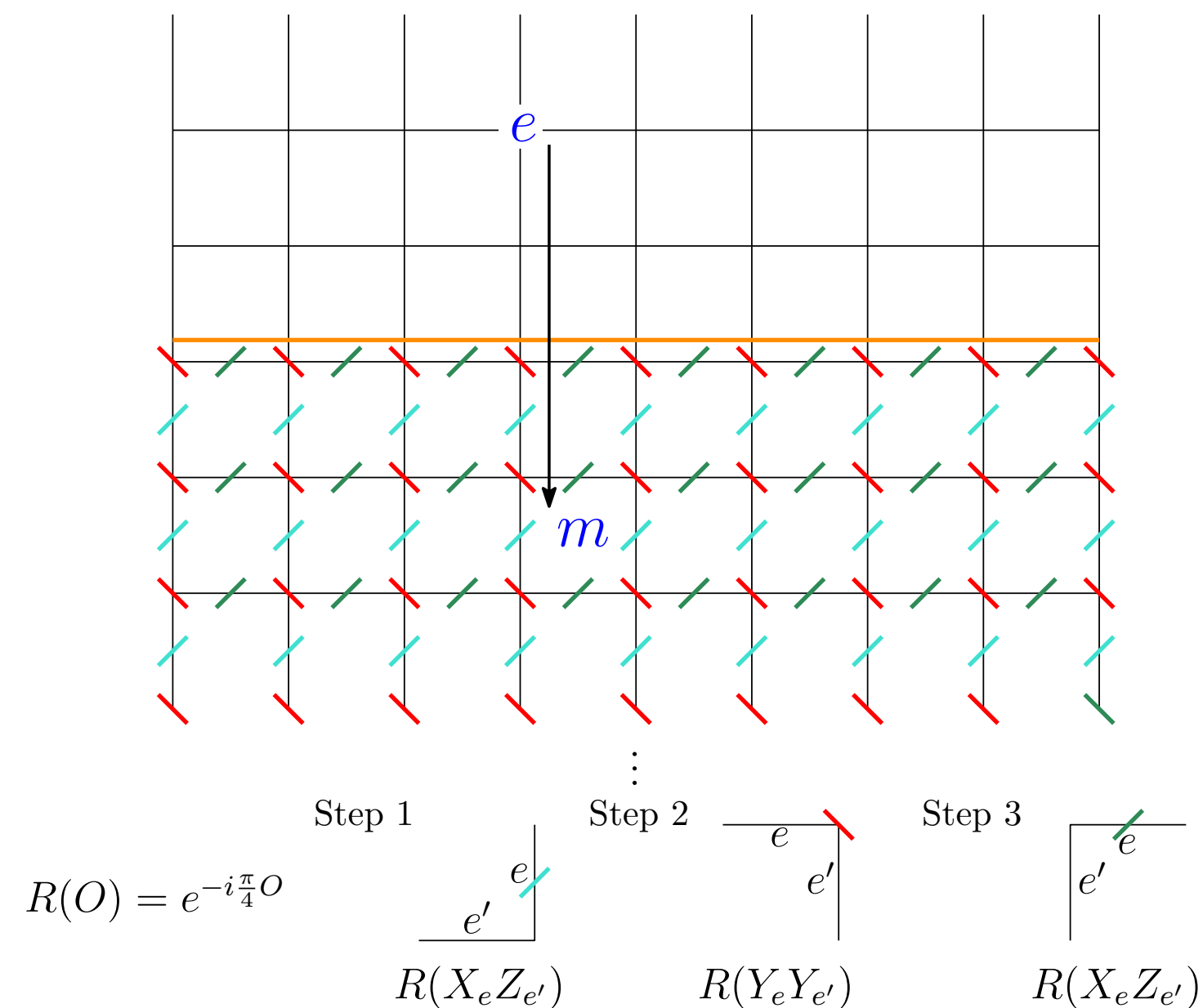
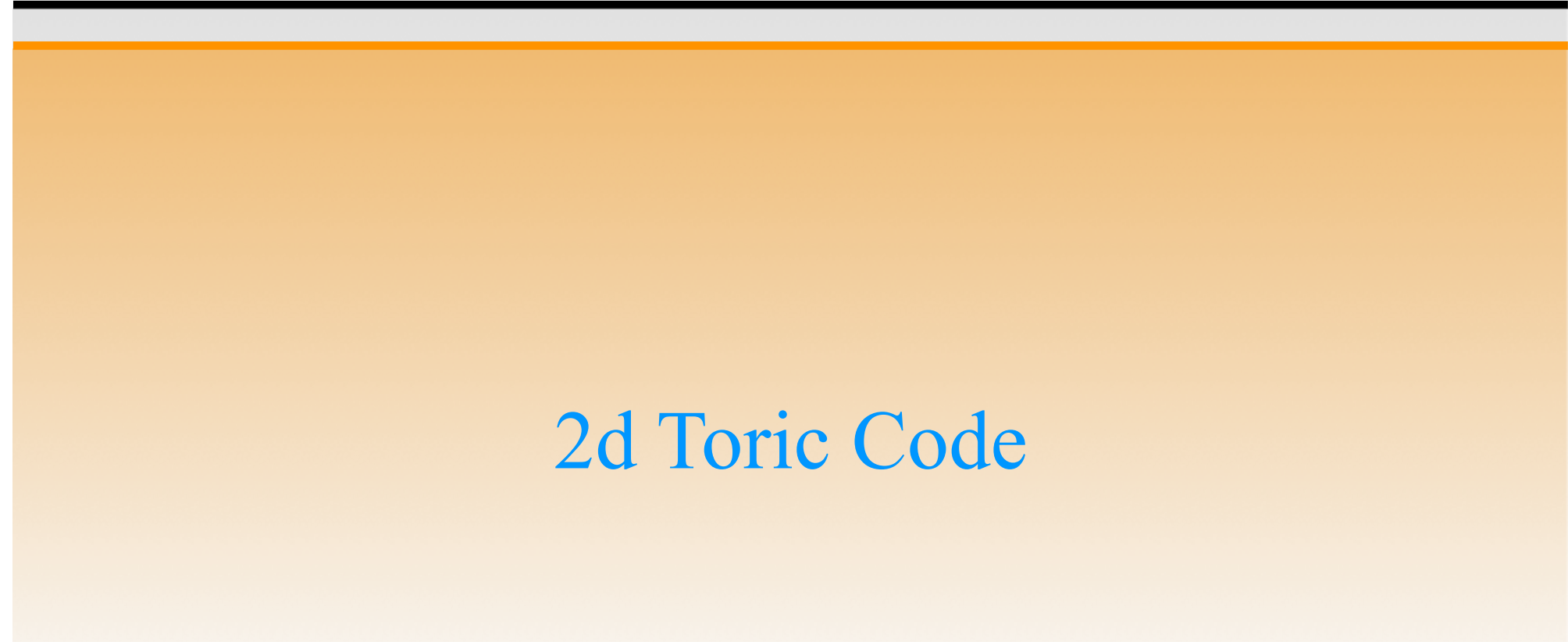


Topological order with a boundary

e condensed



m condensed



U: Sweeping the e-m exchange defect
2d FDQC

U generates an \mathbb{Z}_2 emergent symmetry
only modify the boundary terms non-trivially

Topological order with a boundary

No experiments near the boundary can pin down the boundary type.

e condensed

$$O(\Omega^k)$$

Toric Code

m condensed

$$UO(\Omega^k)U^{-1} = O'(\Omega'^k)$$

Toric Code

U: sweeping the e-m exchange defect
2d FDQC

Topological order with a boundary

No absolute distinction

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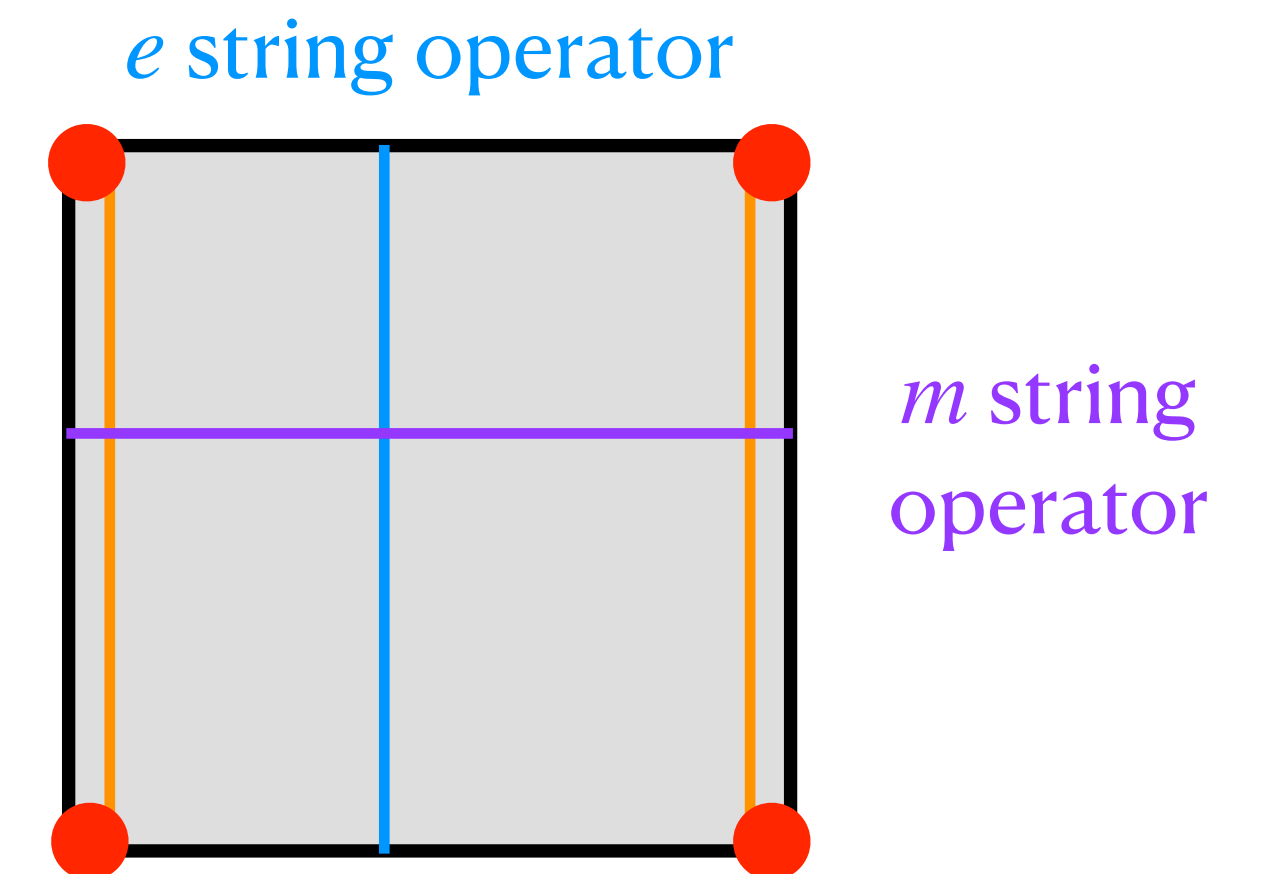
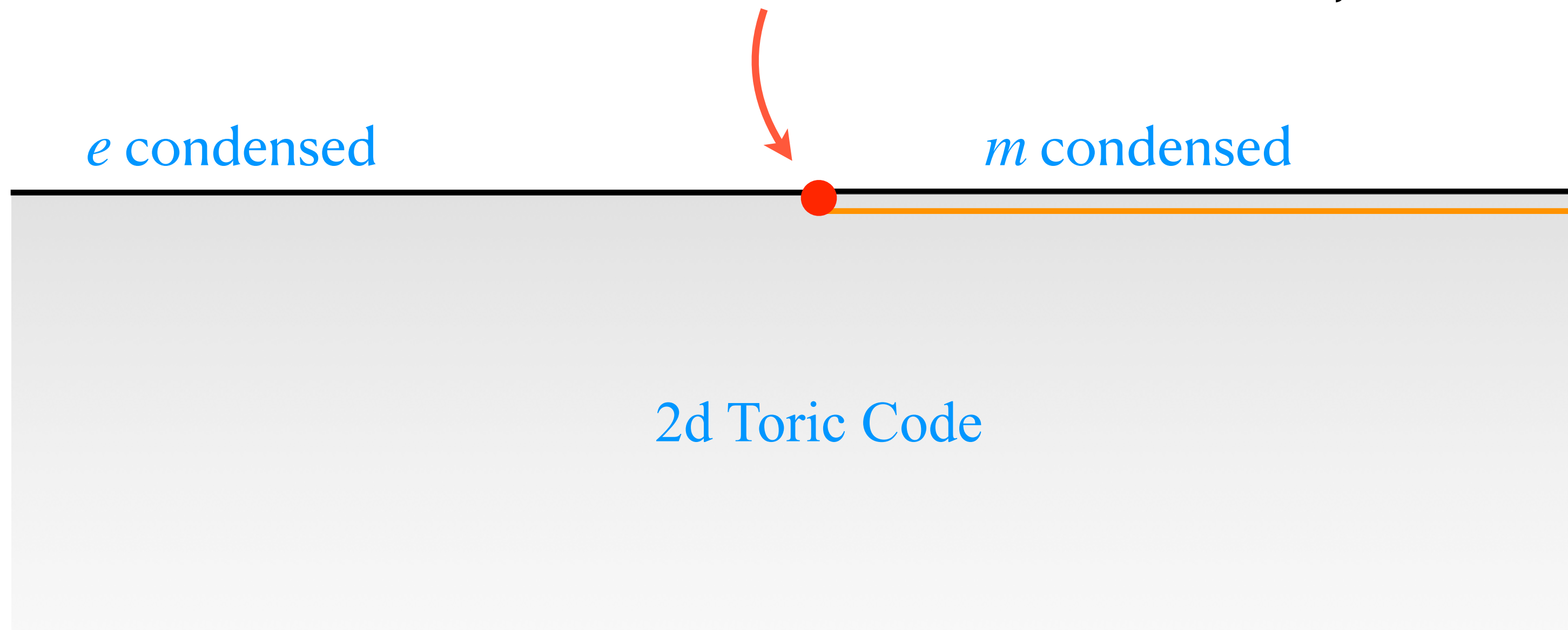
Toric Code

U: sweeping the e-m exchange defect
2d FDQC

Topological order with a boundary

Relative distinction

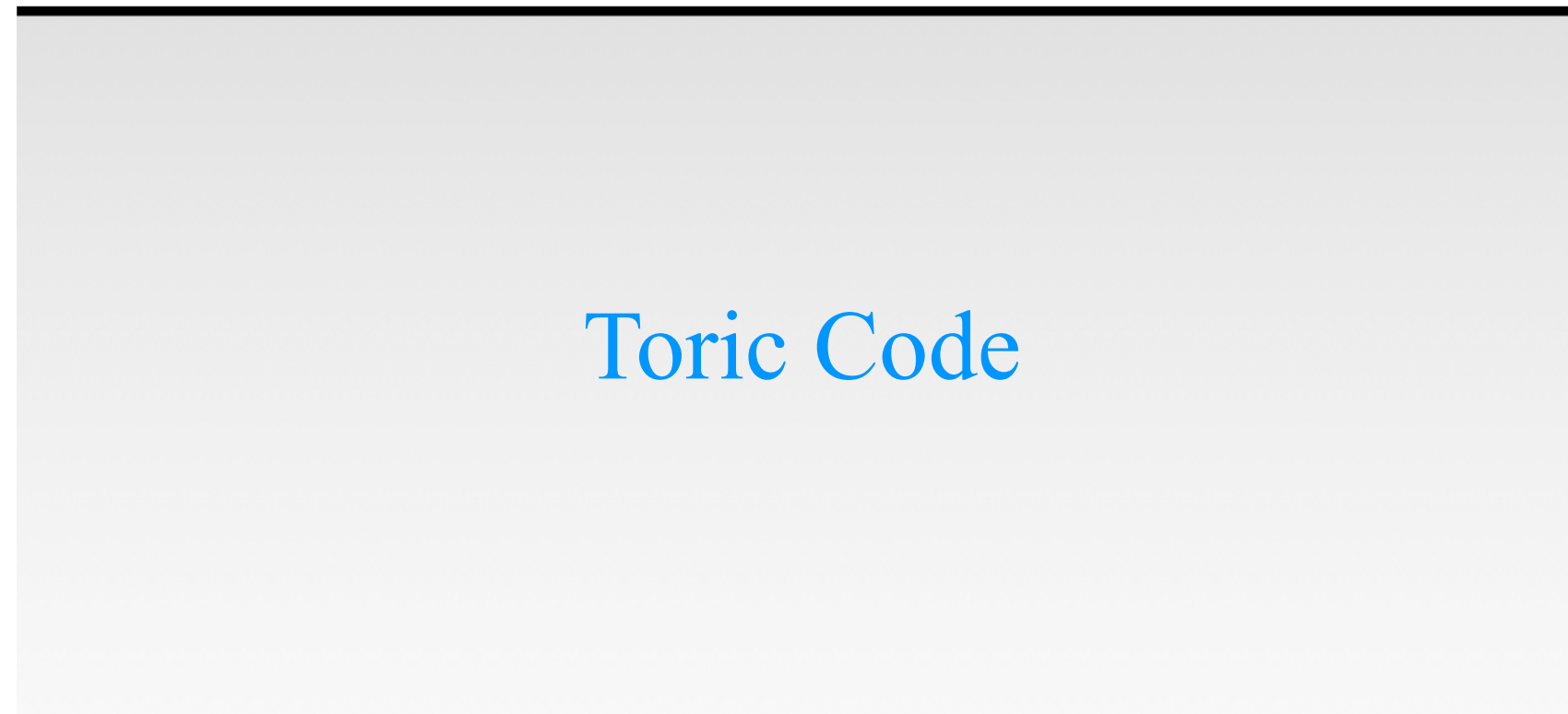
Non-trivial domain wall: no local unitary can remove it



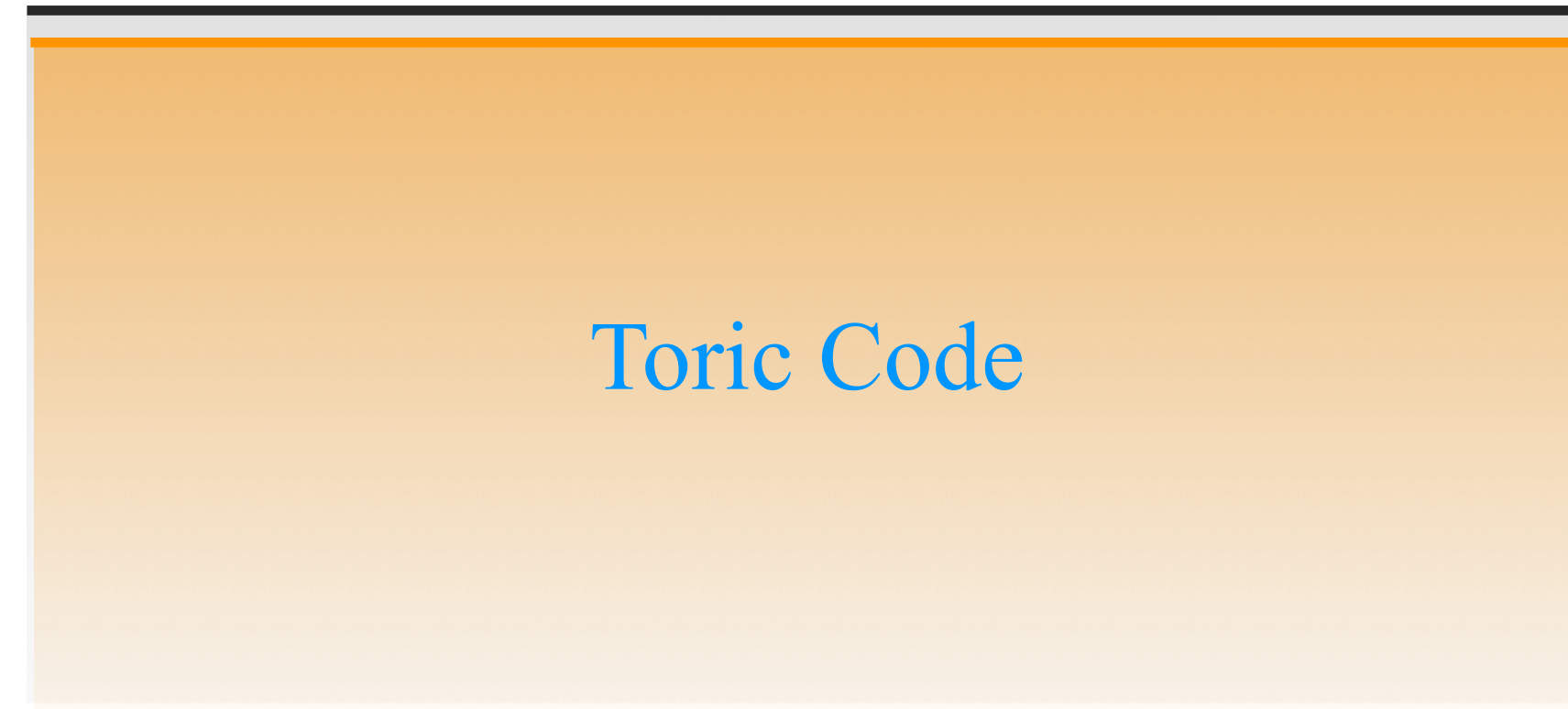
There is a logical qubit.

Topological order with a boundary

e condensed



m condensed



Sweep invertible line defect \longrightarrow emergent symmetry of the bulk \longrightarrow change boundary types
generated by **FDQC** only relatively distinguishable

Boundaries related by bulk FDQC are "twins".

Topological order with a boundary

Example II

m - flux condensed

3d Toric Code

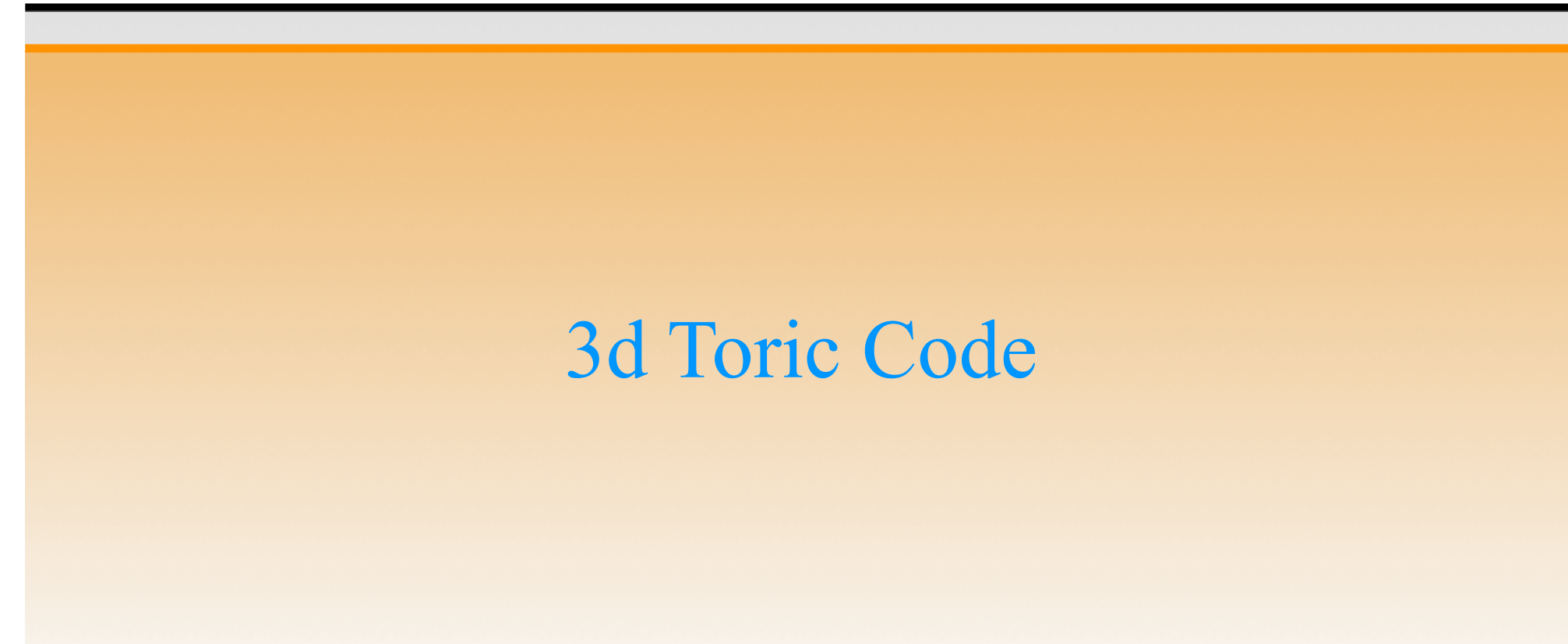
Topological order with a boundary

Example II

m -flux condensed



twisted m -flux condensed



Barkeshli-Chen-Hsin-Kobayashi, 2022;
WJ-Tantivasadakarn-Xu, 2023

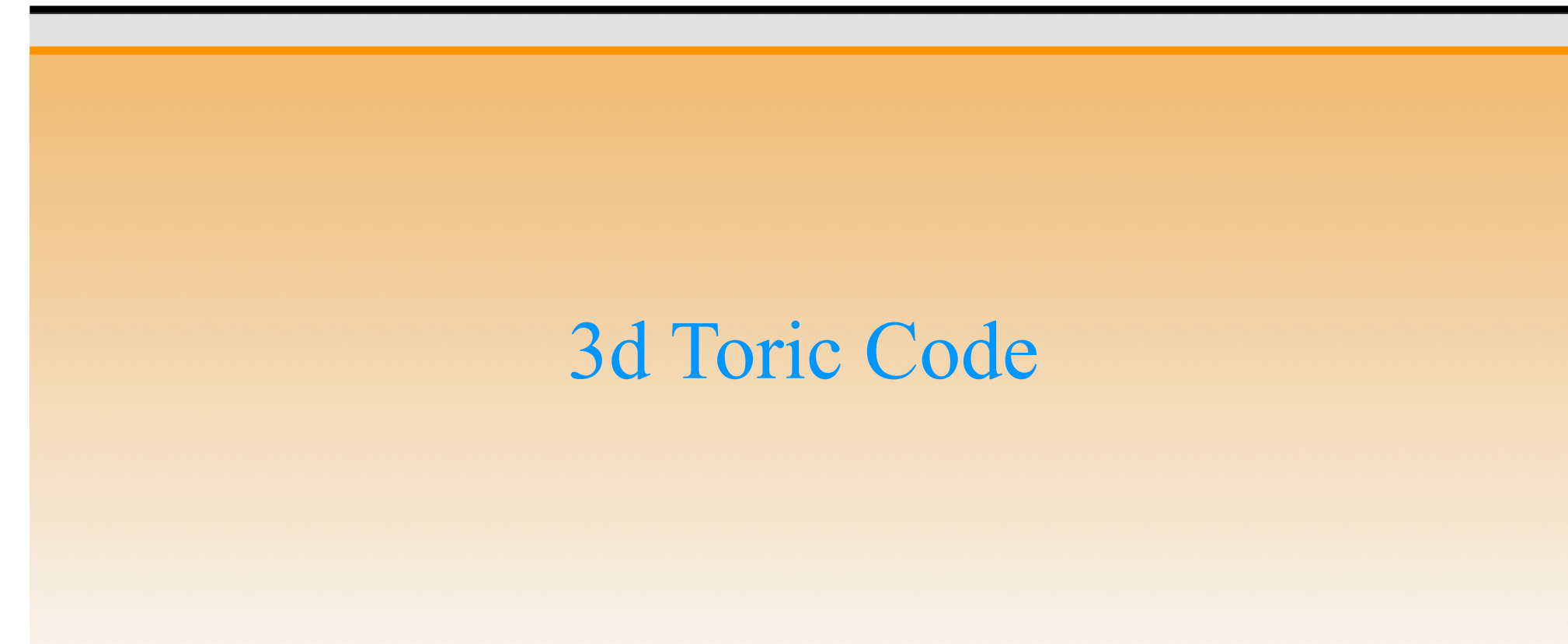
Topological order with a boundary

Example II

m-flux condensed



twisted *m*-flux condensed



Emergent \mathbb{Z}_2 symmetry: Sweep the "gauged \mathbb{Z}_2 SPT" defect (3d FDQC)

$$[U, H_{\text{tc}}] = 0 \text{ in ground state subspace}$$

Barkeshli-Chen-Hsin-Kobayashi, 2022;

WJ-Tantivasadakarn-Xu, 2023

Topological order with a boundary

No absolute distinction

m- flux condensed

$$O(\Omega^k)$$

3d Toric Code

twisted *m*- flux condensed

$$UO(\Omega^k)U^{-1} = O'(\Omega'^k)$$

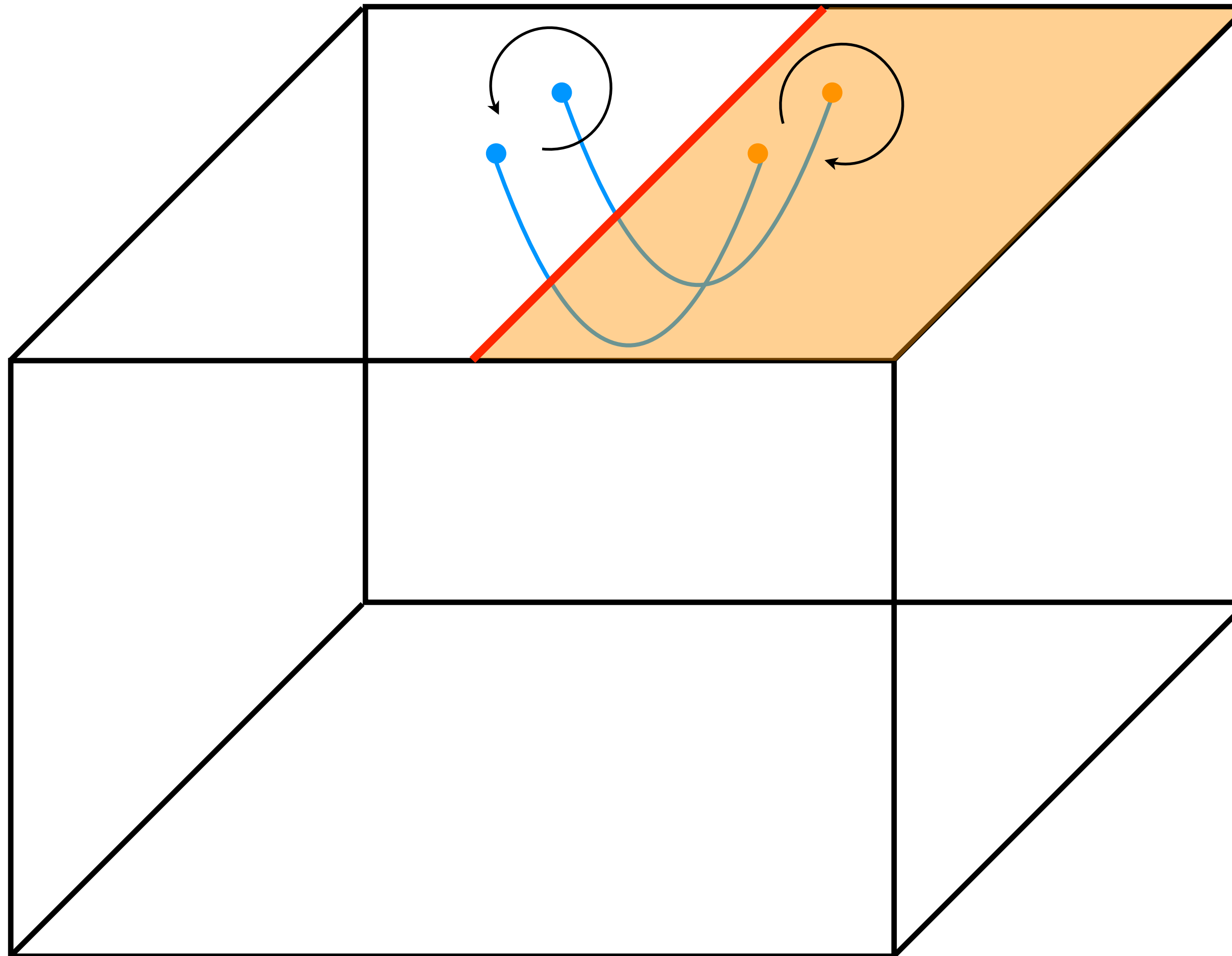
3d Toric Code

3d FDQC

Sweep the "gauged \mathbb{Z}_2 SPT" defect

Relative distinction

$$\theta = -1$$



Topological order with a boundary

Example III

m- flux condensed

fermionic Toric Code

(fermionic gauge charge, flux loop)

Topological order with a boundary

Example III

m - flux condensed

fermionic Toric Code

twisted m - flux condensed

fermionic Toric Code

Sweep the "p+ip" defect

Topological order with a boundary

Example III

m - flux condensed

fermionic Toric Code

twisted m - flux condensed

fermionic Toric Code

Sweep the "p+ip" defect

3d FDQLU preserving $(-1)^F$

Topological order with a boundary

Example III

No absolute distinction

m- flux condensed

fermionic Toric Code

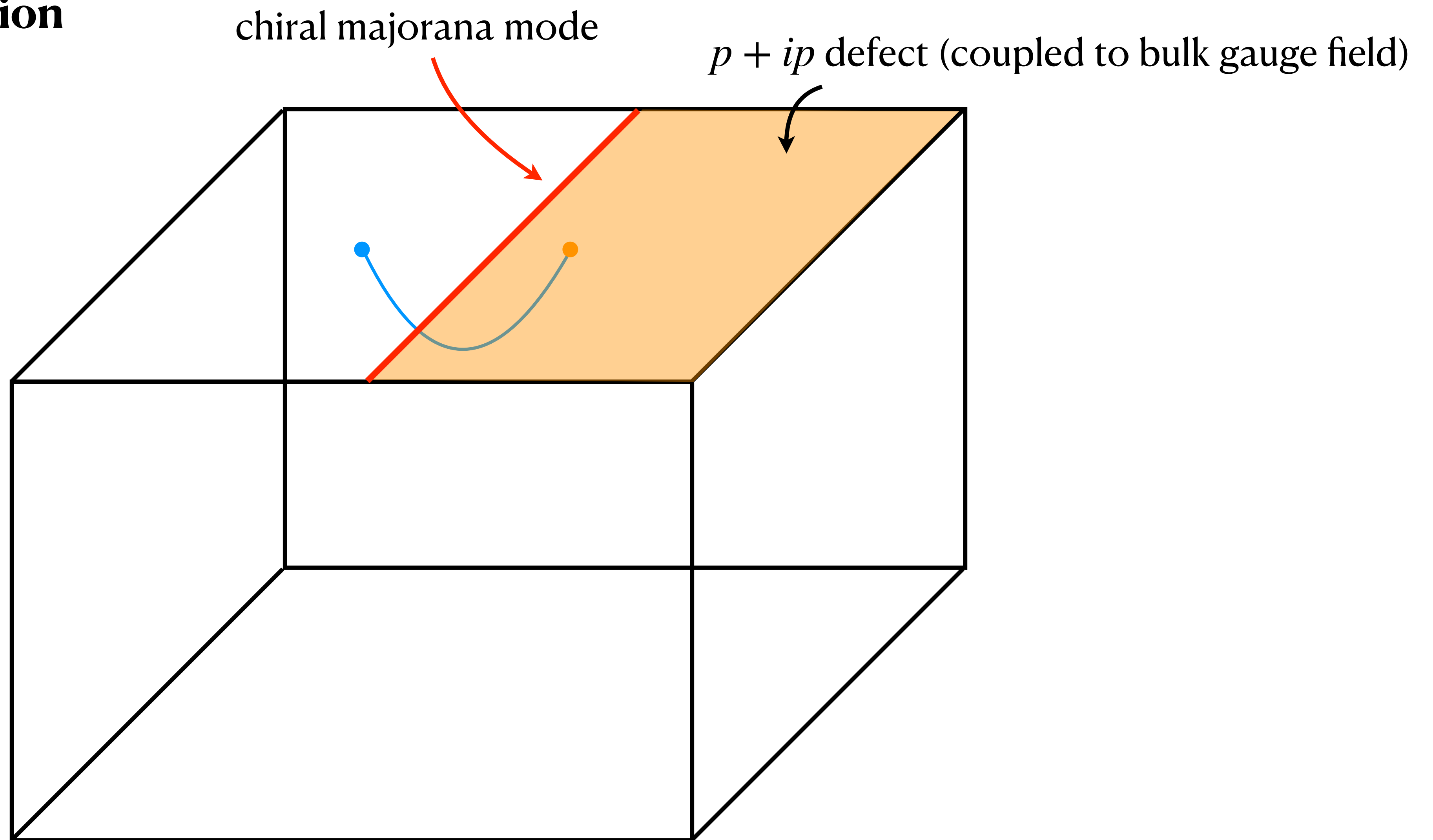
twisted *m*- flux condensed

fermionic Toric Code

Sweep the "p+ip" defect

3d FDQLU preserving $(-1)^F$

Relative distinction



A boundary

Topological order

Relatively
distinct
↔

A "twisted" boundary

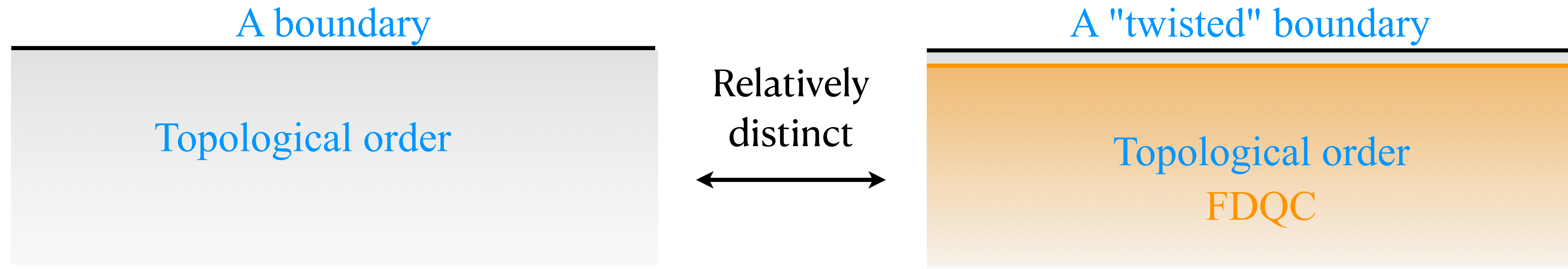
Topological order

FDQC

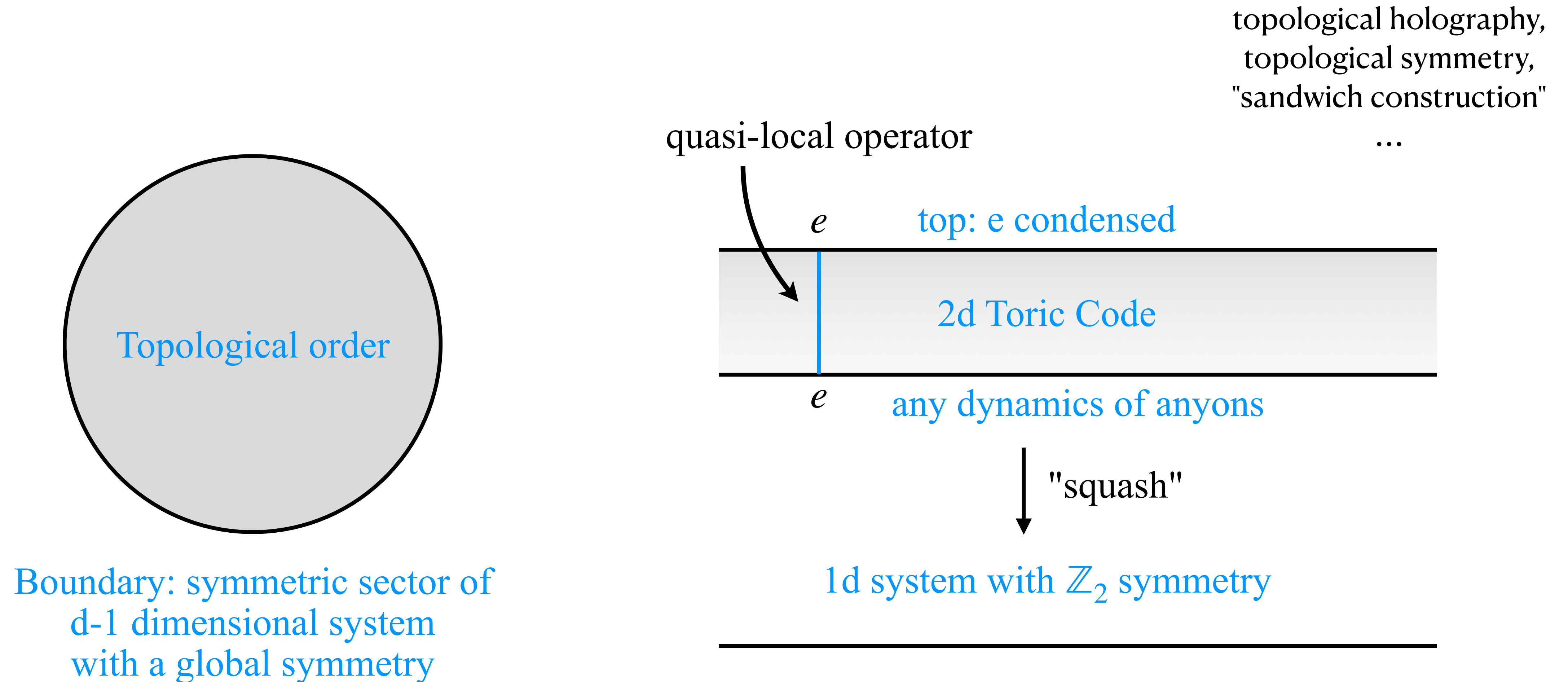
generate emergent symmetry
in the bulk

Argument using "Symmetry TO/TFT"

Step 1:

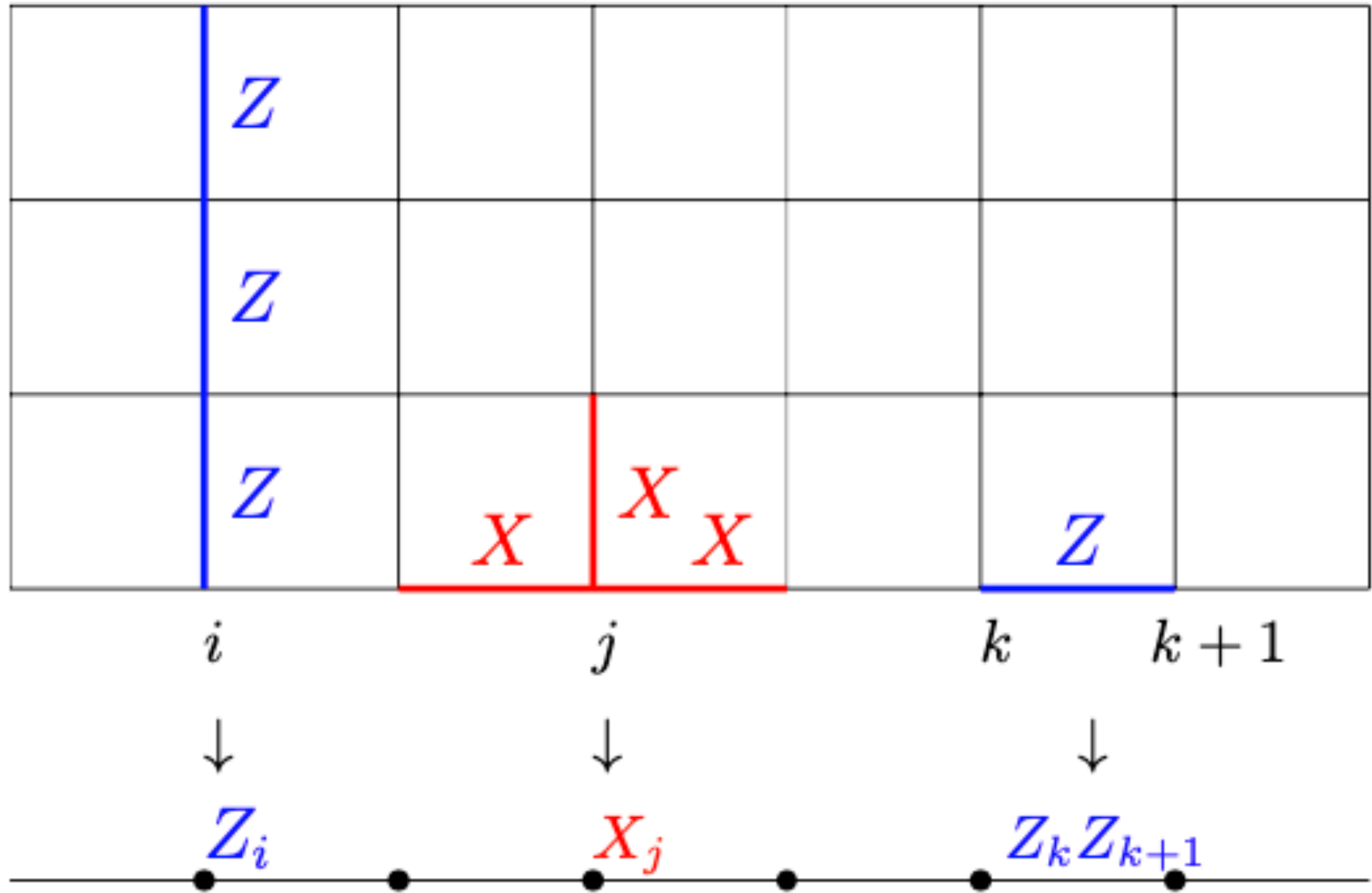
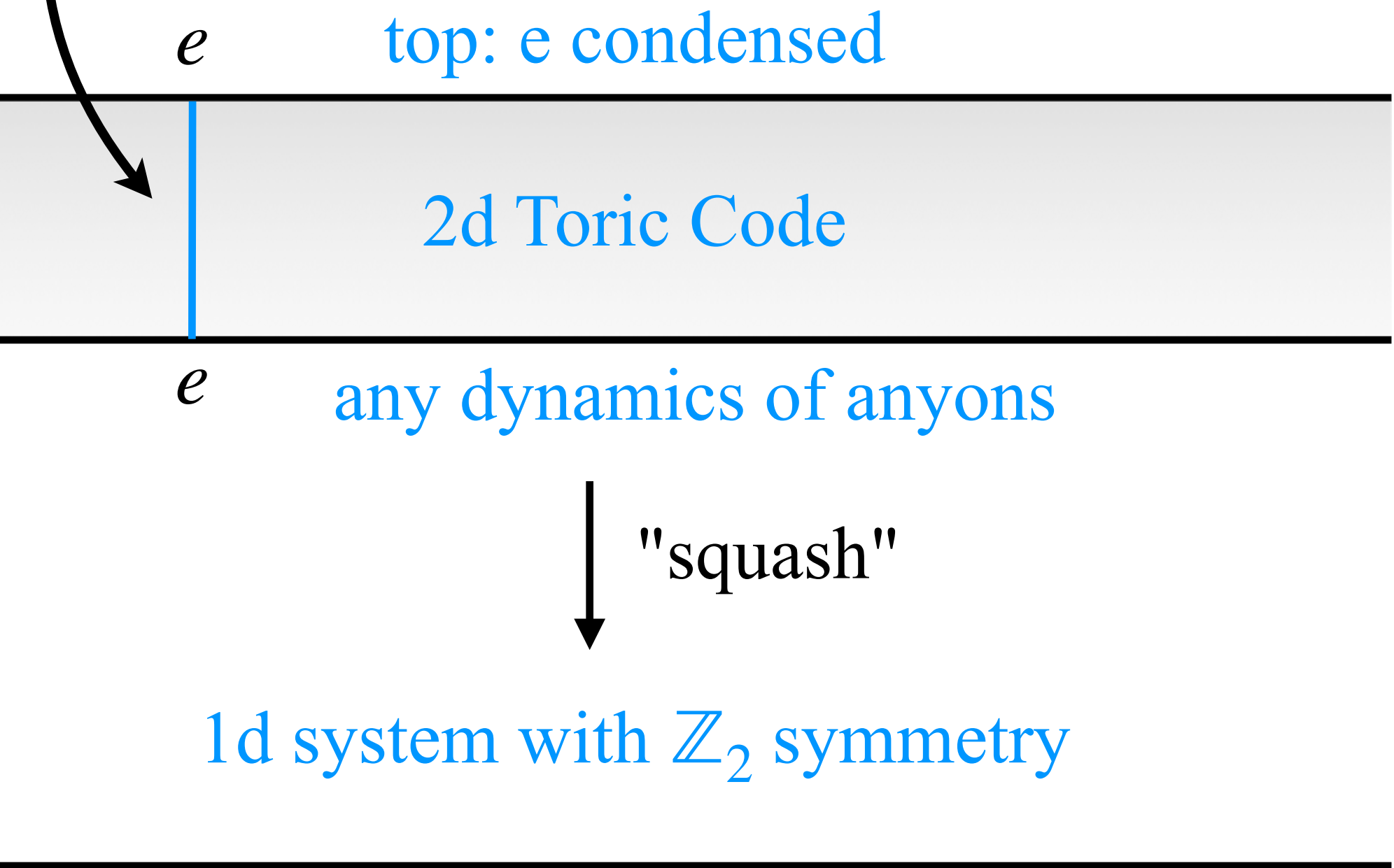


Invertible phases from symmetry TO/TFT construction



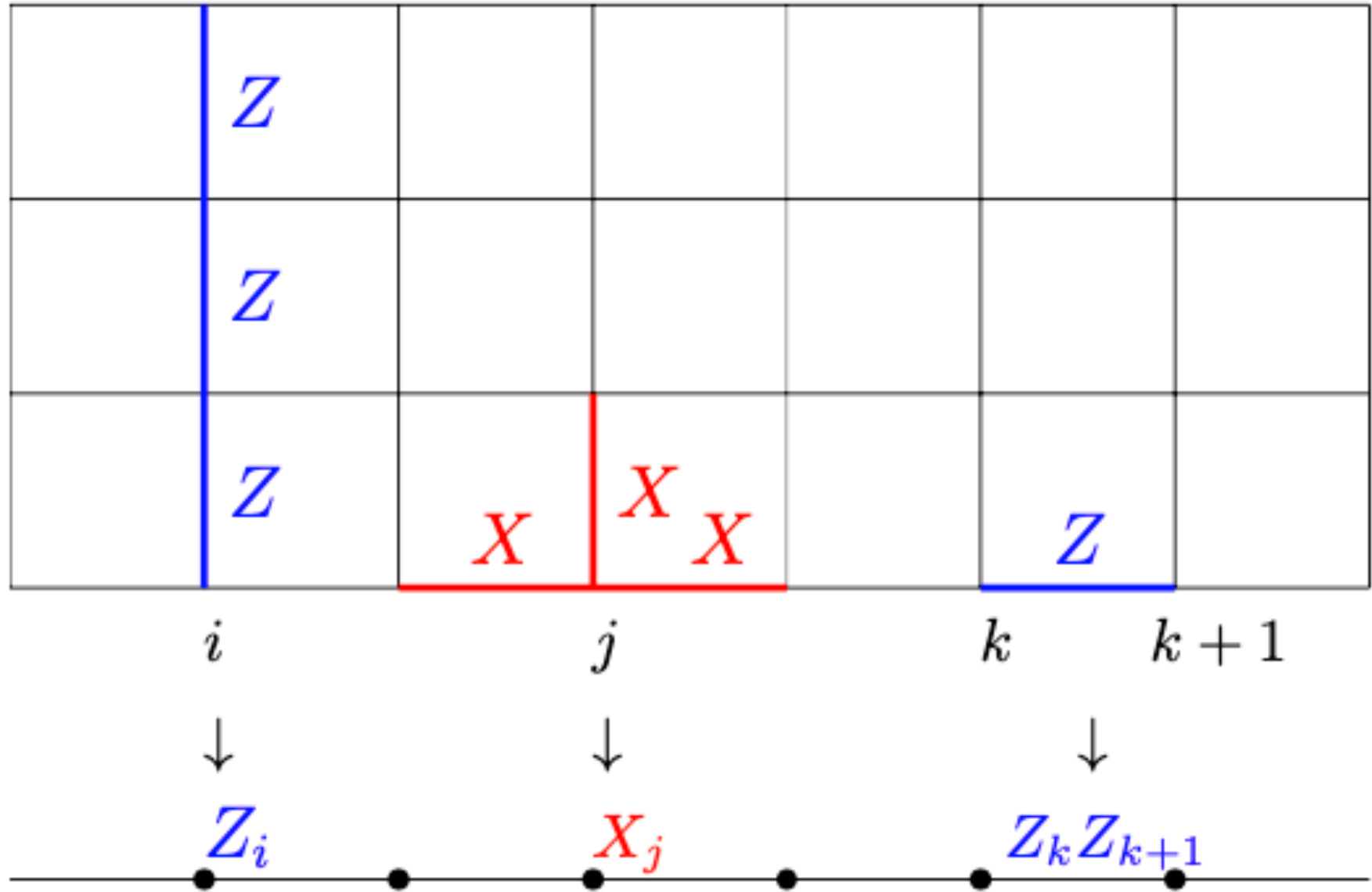
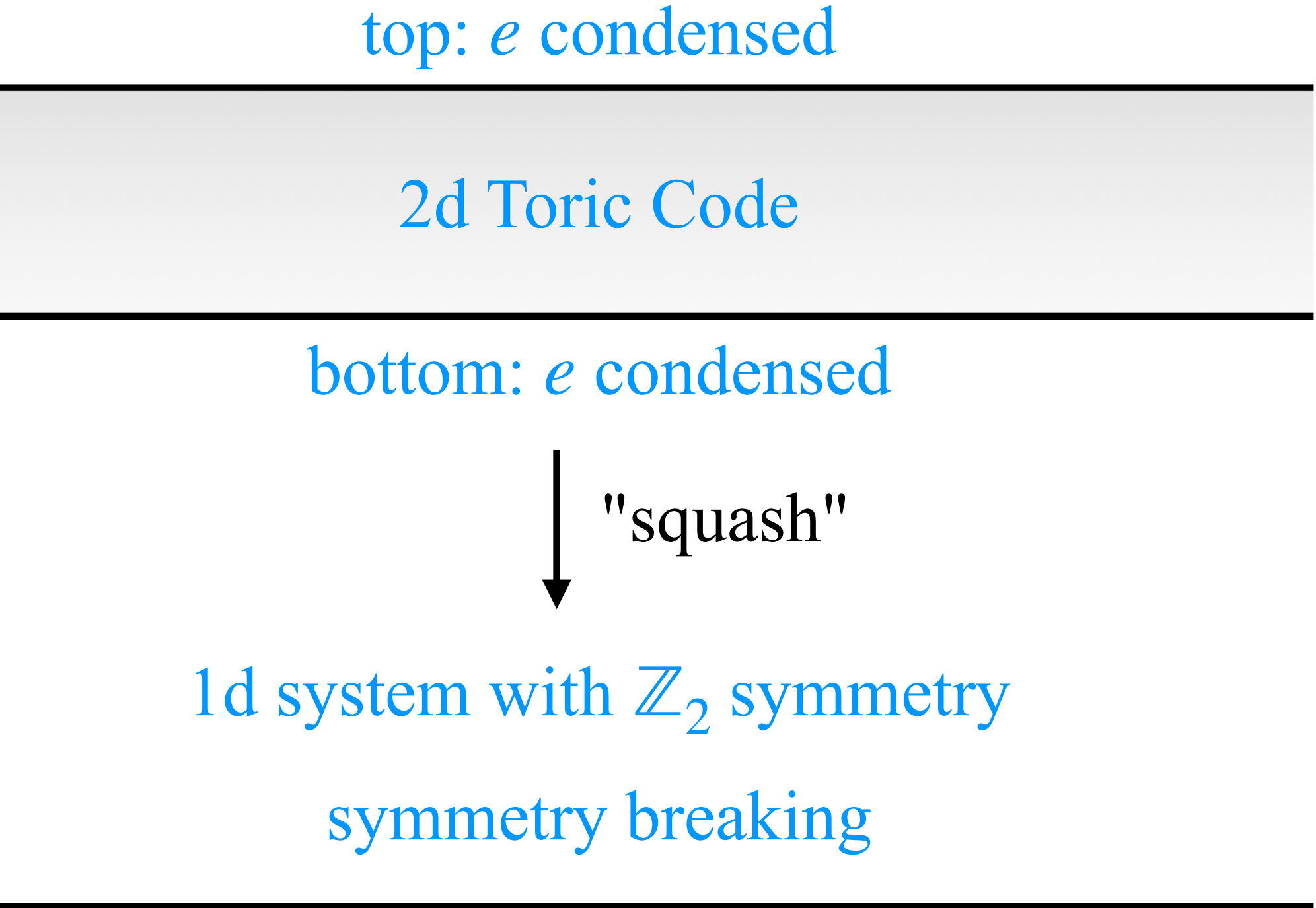
Invertible phases from symmetry TO/TFT construction

quasi-local operator



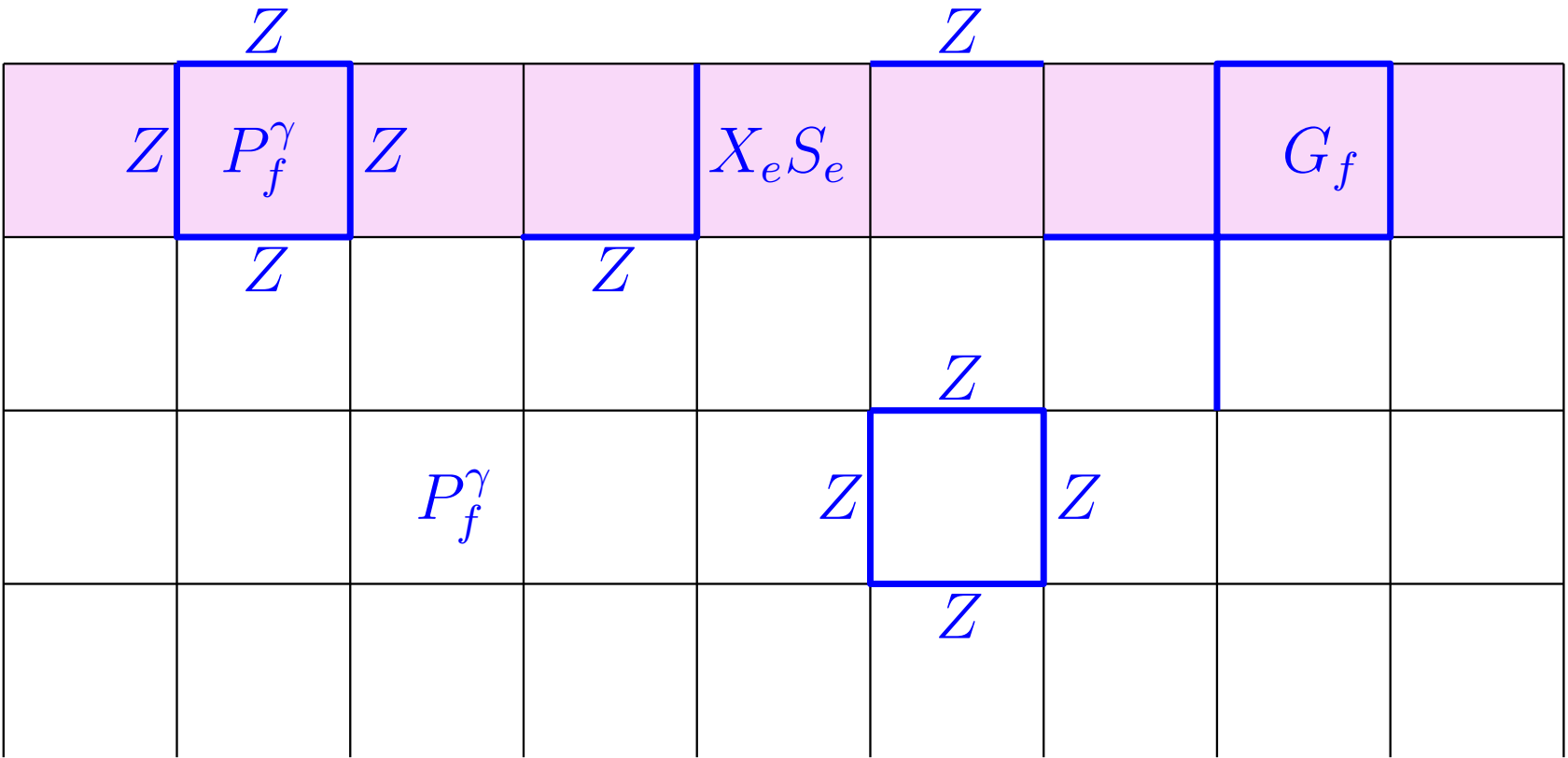
WJ-Wen, 2019; Lichtman, et al, 2020; Kong, et al, 2020; Kulp, 2020; Freed-Moore-Teleman, 2021; Apruzzi, et al, 2021 ...

Invertible phases from symmetry TO/TFT construction



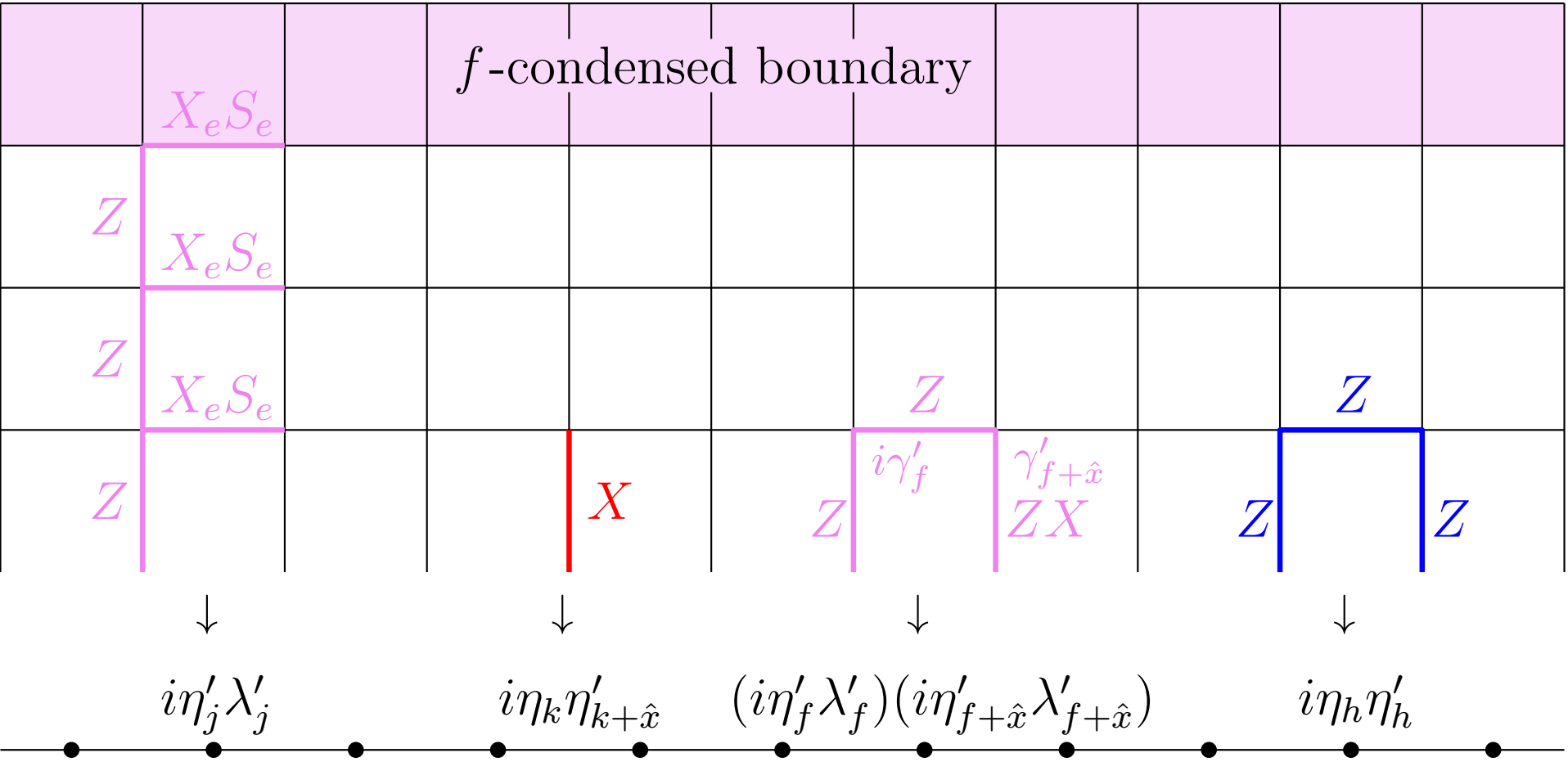
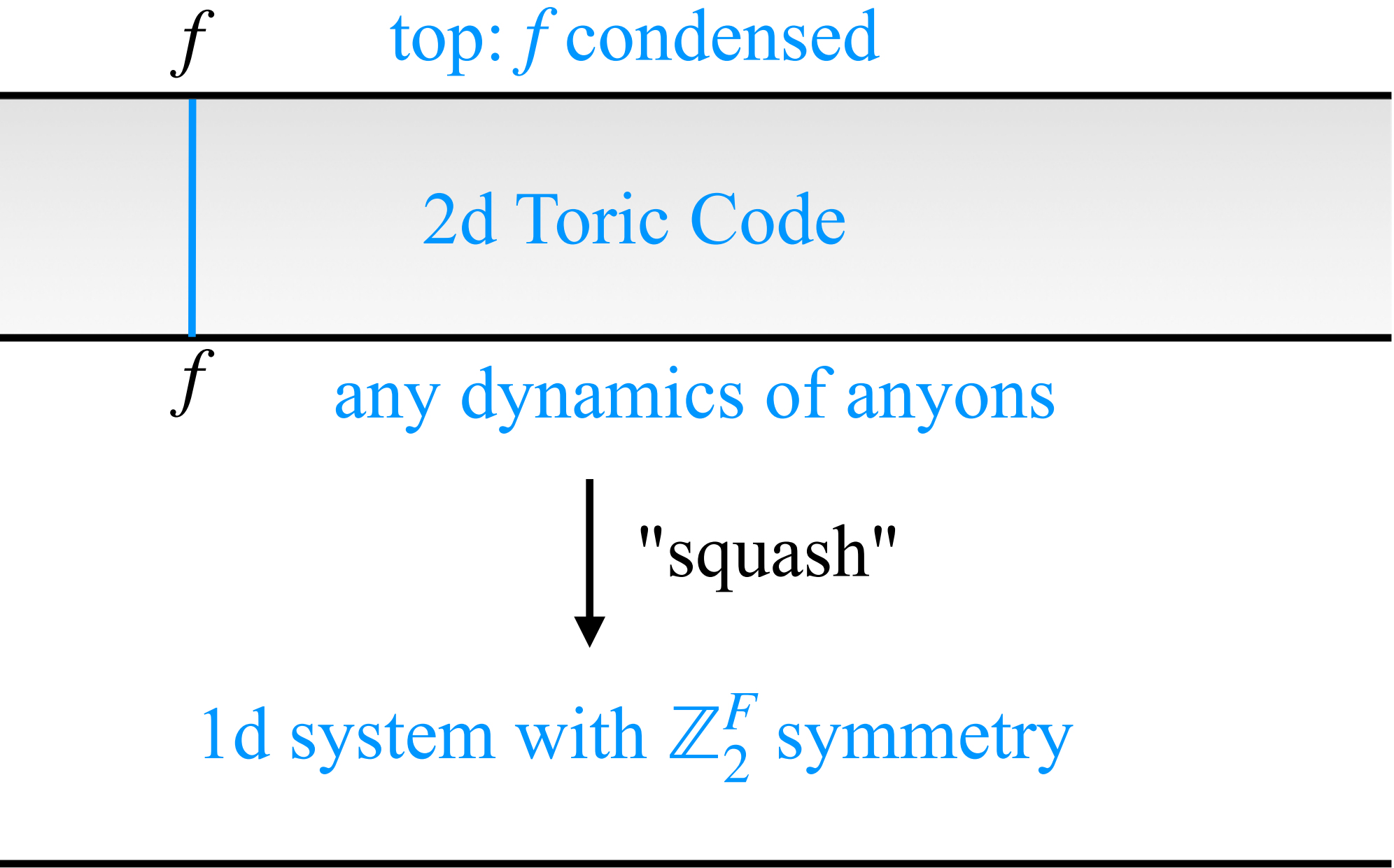
Invertible phases from symmetry TO/TFT construction

top: f condensed



$$P_f^\gamma = -i\gamma_f \gamma'_f \qquad S_e = i\gamma'_{e_L} \gamma_{e_R} \qquad G_f = \begin{array}{c} \begin{array}{cc} X & \begin{array}{c} Y \quad f \\ Y \end{array} \end{array} \\ X \end{array} \begin{array}{c} Z \\ Z \end{array}$$

Invertible phases from symmetry TO/TFT construction



Invertible phases from symmetry TFT construction

top: f condensed

2d Toric Code

e condensed

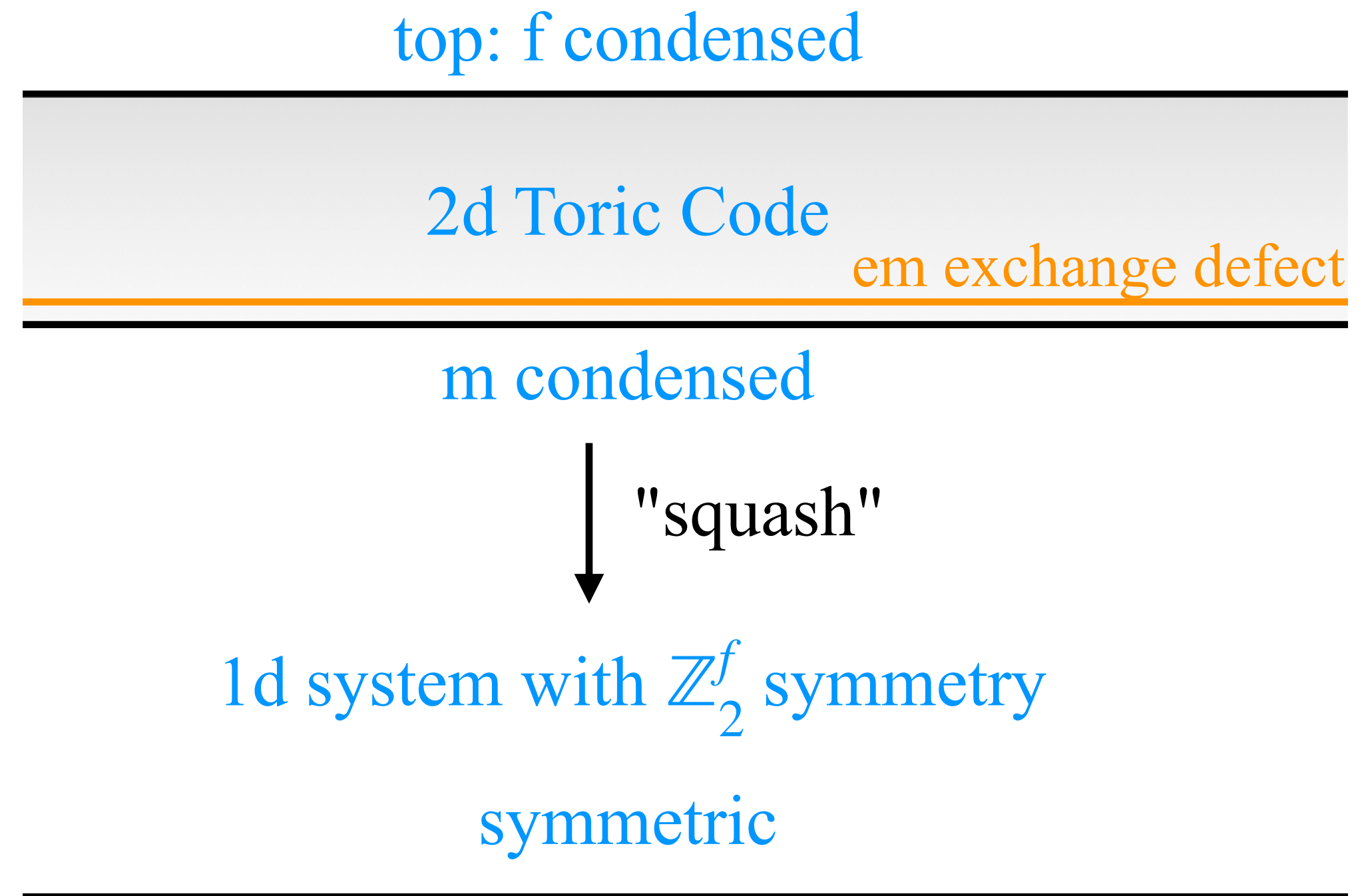
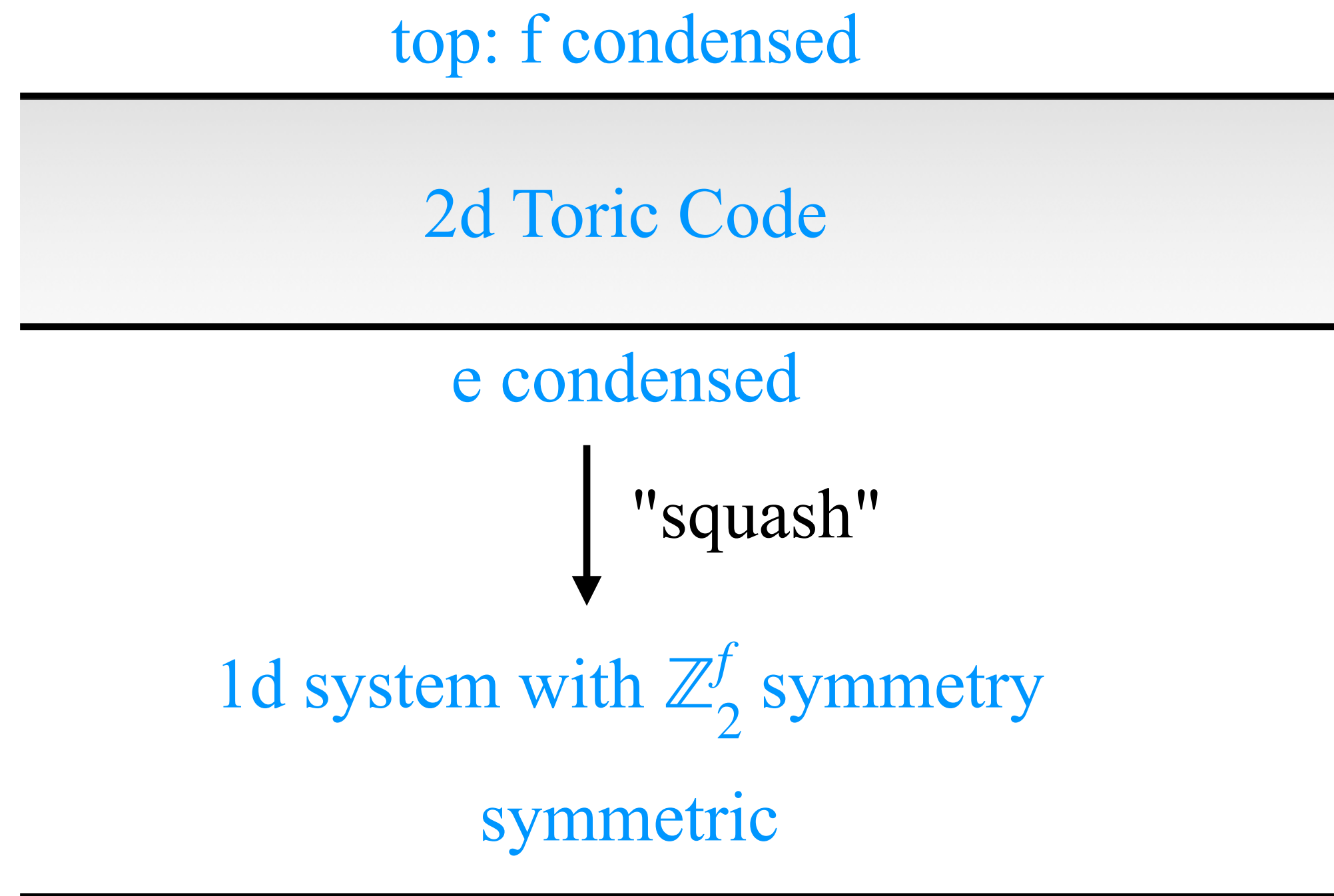


"squash"

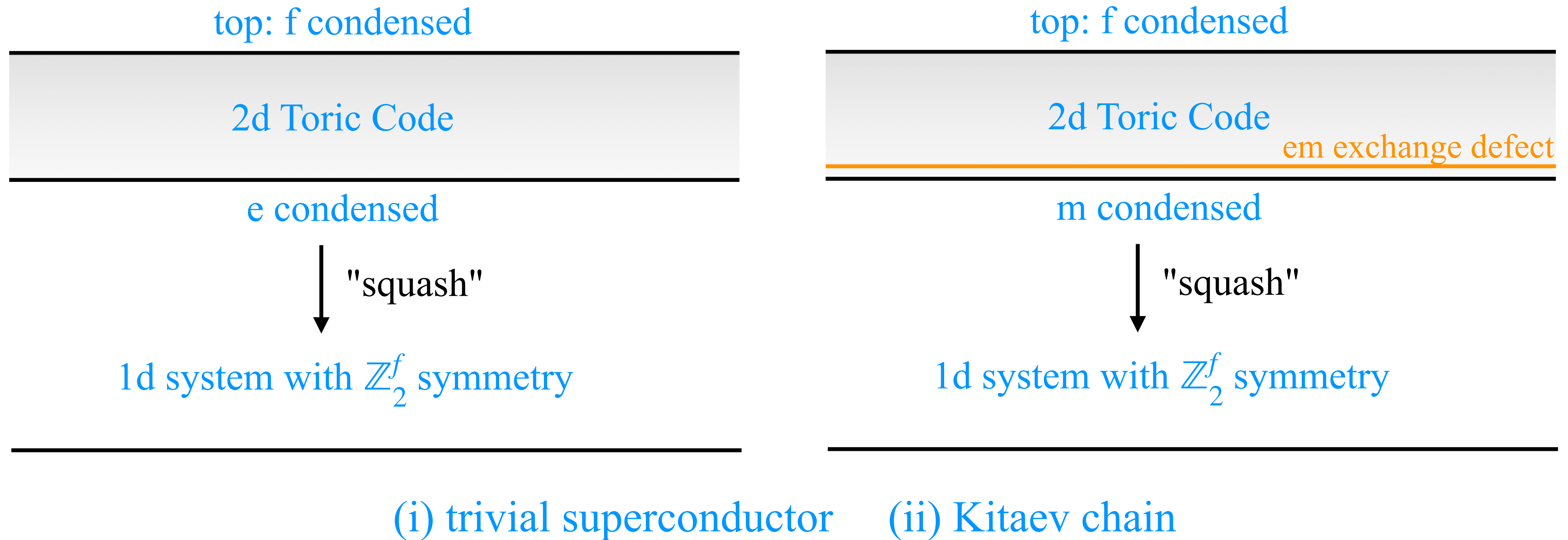
1d system with \mathbb{Z}_2^f symmetry

symmetric

Invertible phases from symmetry TFT construction



Invertible phases from symmetry TFT construction



Invertible phases from symmetry TFT construction

	Z			Z				
Z	P_f^γ	Z		$X_e S_e$			G_f	
	Z		Z					
					Z			
	Z	P_f^γ		Z		Z		
Z		Z			Z		P_f^γ	

(i) trivial superconductor

	Z			Z				
Z	P_f^γ	Z		$X_e S_e$			G_f	
	Z		Z					
					Z			
		P_f^γ		Z		Z		
X					Z		P_f^γ	

(ii) Kitaev chain

Invertible phases from symmetry TFT construction



(i) trivial superconductor (ii) Kitaev chain

- Symmetric operators have one-to-one correspondence
- **Low entanglement excitations of the two have one-to-one correspondence**
- Only using Hamiltonians, and symmetric operators, one cannot distinguish two phases.

Invertible phases from symmetry TFT construction

top: e condensed

3d Toric Codes

flux loop condensed



"squash"

2d system with \mathbb{Z}_2 symmetry

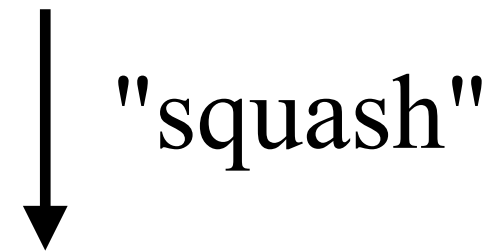
symmetric phase

Invertible phases from symmetry TFT construction

top: e condensed

3d Toric Codes

flux loop condensed



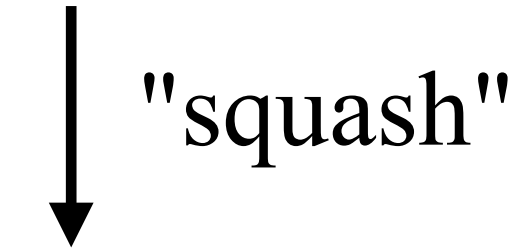
2d system with \mathbb{Z}_2 symmetry
symmetric phase

top: e condensed

3d Toric Codes

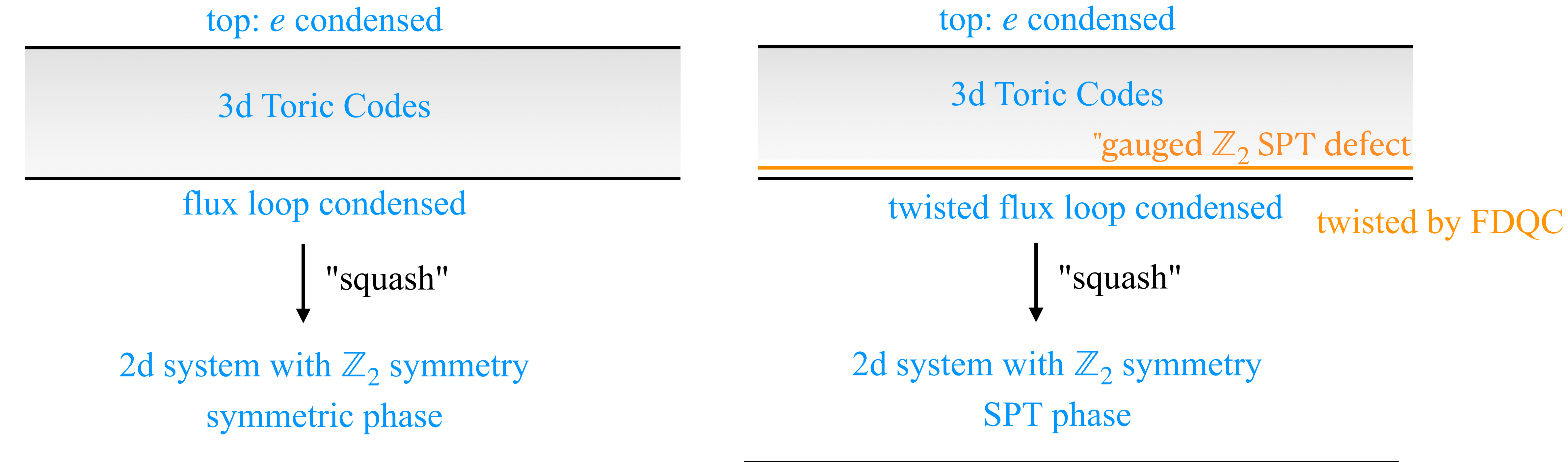
"gauged \mathbb{Z}_2 SPT defect"

twisted flux loop condensed



2d system with \mathbb{Z}_2 symmetry
SPT phase

Invertible phases from symmetry TFT construction



- Symmetric operators have one-to-one correspondence
- **Low entanglement excitations of the two have one-to-one correspondence**
- Only using Hamiltonians, and symmetric operators, one cannot distinguish two phases.

"Sandwiches" with fermionic Toric Code

top: fermion condensed

3d fermionic Toric Code

m - flux condensed

↓ "squash"

2d phase with fermion parity symmetry

trivial superconductor

top: fermion condensed

3d fermionic Toric Code

$p + ip$ defect

twisted m - flux condensed

↓ "squash"

2d phase with fermion parity symmetry

$p + ip$ superconductor

		2d trivial superconductor	$p + ip$ superconductor
non-trivial defect/excitation	0d	single fermion 1d QC	single fermion 1d QC
	1d	a Kitaev chain defect 1d SQC	fermion parity twist line 1d SQC

		2d \mathbb{Z}_2 paramagnet	2d \mathbb{Z}_2 SPT
non-trivial defect/excitation	0d	\mathbb{Z}_2 charge 1d FDLU	\mathbb{Z}_2 charge 1d FDLU
	1d	1d \mathbb{Z}_2 SSB 1d SQC	1d \mathbb{Z}_2 SSB 1d SQC

Conclusion & Further questions

Low entanglement excitations in an invertible phase and those in a product state have a one-to-one correspondence

SPTs, Kitaev chain, $p + ip$ superconductor

- Higher form SPTs
- Integer quantum Hall
- Non-invertible defects
- Implications for dynamics
- ...

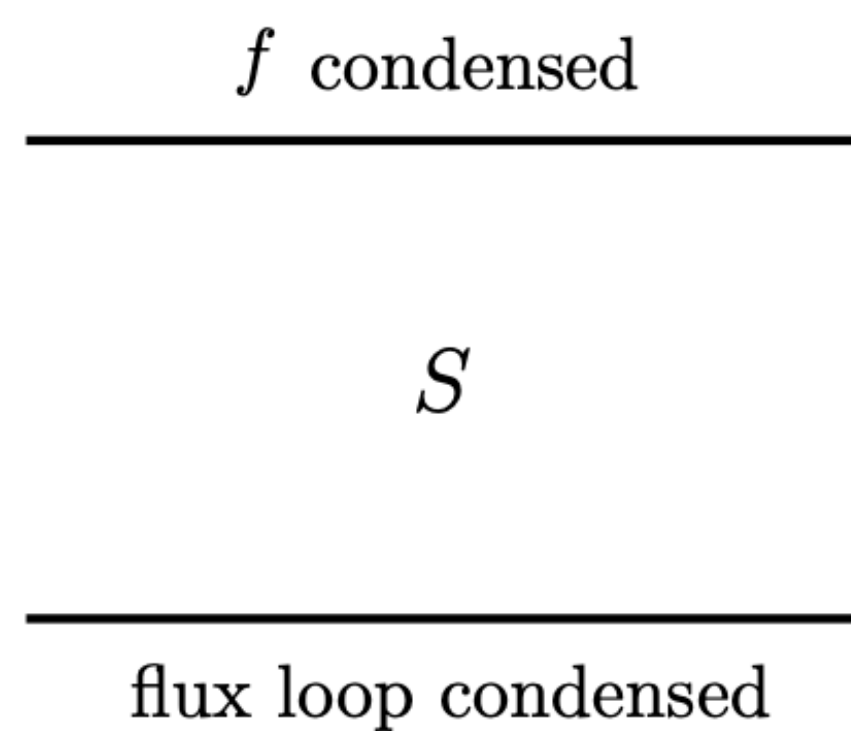
Conclusion & Further questions

Low entanglement excitations in an invertible phase and those in a product state have a one-to-one correspondence

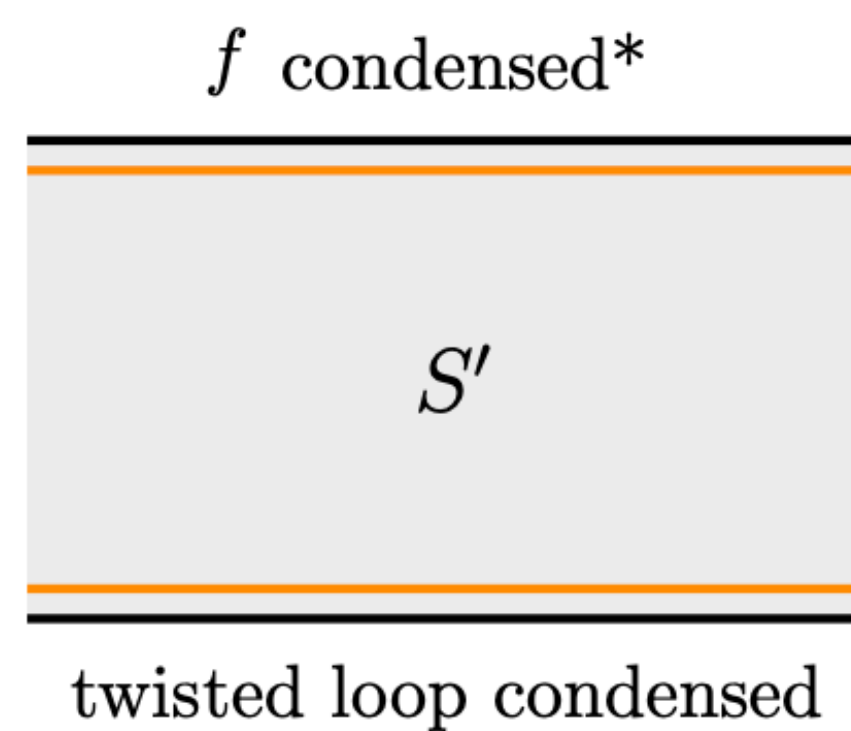
SPTs, Kitaev chain, $p + ip$ superconductor

- Higher form SPTs
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- Non-invertible defects
- Implications for dynamics
- ...

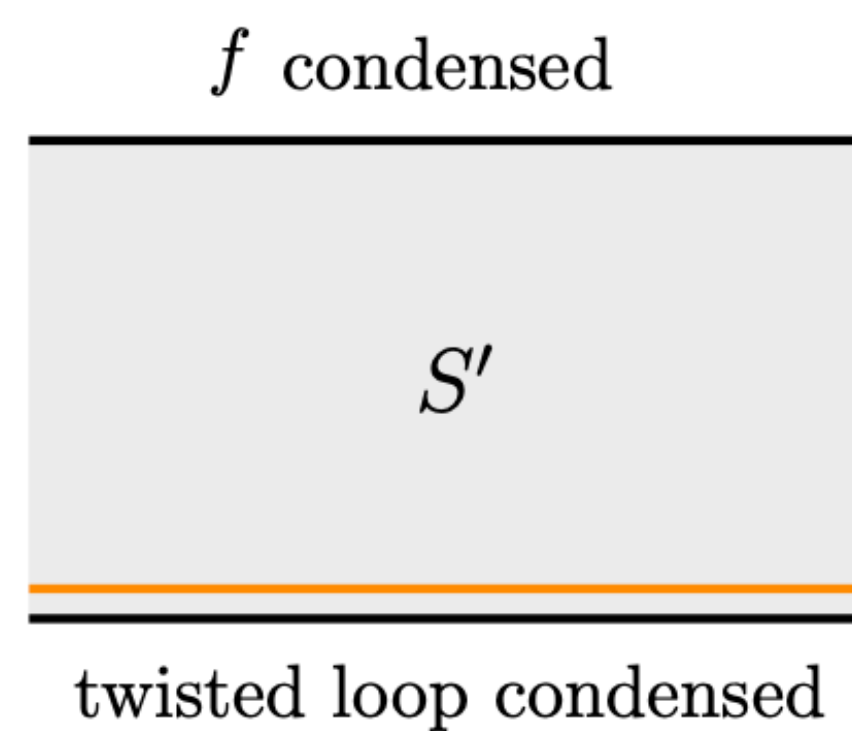
Thank you !



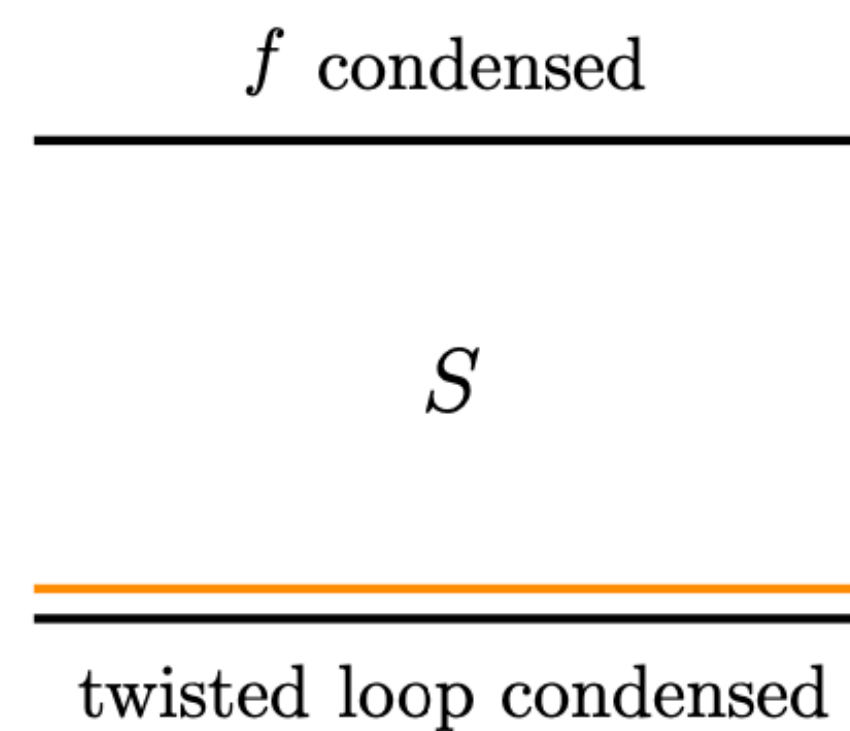
(a)



(b)



(c)



(d)